

# Two-meson form factors from dispersion relations

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**Future Challenges in Non-Leptonic B Decays:  
Theory and Experiment**

Bad Honnef, Germany, 10-12/02/2016

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# INTRODUCTION

## Three-Body I

08:30 – 09:20 Thomas Edward  
Latham

**Experimental results from  $B$  decays to  
charmless and open-charm 3-body  
final states**

09:20 – 10:10 Thomas Mannel

**Three-body non-leptonic  $B$  decays and  
QCD Factorization**

CPV

$\alpha, \gamma$

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## Two-pion systems

13:50 – 14:40 Pablo Roig Garcés

**Two-meson form factors from  
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14:40 – 15:30 Pere Masjuan

**Padé Theory: A toolkit for hadronic  
form factors**

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Dispersive methods in heavy-meson decays

16:45 – 17:30 Ignacio De Bediaga Hickman

CP violation and CPT invariance in charmless three-body  $B$  decays

17:30 – 18:15 Hai-Yang Cheng

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18:15 – 18:30 Patricia Magalhães

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Phenomenology of  $B \rightarrow K\pi\pi$  modes and prospects with LHCb and Belle II data

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$K\pi$

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  - Other approaches: Gounaris-Sakurai '68, Heyn-Lang '81, Truong'88, Kühn-Santamaría '90, Hannah '96-'97, Domínguez '01, Cillero-Pich '03, Bruch-Khodjamirian-Kühn '05, Czyz-Grzelinska-Kühn '10, Hanhart'12, Masjuan (**next talk**), ...

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**VFF**

$\tau^- \rightarrow \nu_\tau \pi^- \pi^0$   
( $s > 0$ , B/ $\Phi$  factories)

$e^+ e^- \rightarrow \pi^+ \pi^-$   
(Via Watson's FSI Th.)

$e^- \pi^+ \rightarrow e^- \pi^+$   
( $s < 0$ , NA7)

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$SU(2)_1$

Cirigliano-Ecker-Neufeld '01,'02  
Descotes-Genon—Moussallam '14



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→ The approximate  $\chi$  symmetry of low-E QCD determines the VFF and is a good starting point to reach GeV energies

# $\pi\pi$ VFF

→ Near threshold,  $SU(n_f)_L \times SU(n_f)_R \rightarrow SU(n_f)_V$ ,  $\chi$ PT, determines the low-energy expansion of the VFF:

$$F(s)^{\text{ChPT}} = 1 + \frac{2L_9(\mu)}{f_\pi^2} s - \frac{s}{96\pi^2 f_\pi^2} \times \left[ A(m_\pi^2/s, m_\pi^2/\mu^2) + \frac{1}{2} A(m_K^2/s, m_K^2/\mu^2) \right]$$

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→ To enlarge the domain of applicability, **DRs** and unitarization techniques (Oller-Oset-Palomar '01, De Trocóniz-Ynduráin '02) have been employed

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(Guerrero-Pich '97, Pich-Portolés '01, Gómez-Dumm—PR '13, Celis-Cirigliano-Passemar '14)

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- **Unitarity** in the elastic region implies

$$F_V^\pi(s) = P_n(s) \exp \left\{ \frac{s^n}{\pi} \int_{s_{\text{thr}}}^{\infty} ds' \frac{\delta_1^1(s')}{(s')^n (s' - s - i\epsilon)} \right\} \quad \log P_n(s) = \sum_{k=0}^{n-1} \alpha_k \frac{s^k}{k!} \quad (\text{Watson's FSI Th.})$$

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→ Guerrero-Pich'97 delays breaking

$$F(s) = \frac{M_\rho^2}{M_\rho^2 - s - iM_\rho \Gamma_\rho(s)} \exp\left\{ \frac{-s}{96\pi^2 f_\pi^2} \times \left[ \text{Re } A(m_\pi^2/s, m_\pi^2/M_\rho^2) + \frac{1}{2} \text{Re } A(m_K^2/s, m_K^2/M_\rho^2) \right] \right\}$$

$$\Gamma_\rho(s) = -\frac{M_\rho s}{96\pi^2 F_\pi^2} \text{Im} \left[ A_\pi(s) + \frac{1}{2} A_K(s) \right] = \frac{s M_\rho}{96\pi F_\pi^2} \left[ \theta(s - 4m_\pi^2) \sigma_\pi^3(s) + \frac{1}{2} \theta(s - 4m_K^2) \sigma_K^3(s) \right]$$

# $\pi\pi$ VFF

→  $F_V^\pi(s) = P_n(s) \exp\left\{ \frac{s^n}{\pi} \int_{s_{\text{thr}}}^\infty ds' \frac{\delta_1^1(s')}{(s')^n (s' - s - i\epsilon)} \right\} \quad \log P_n(s) = \sum_{k=0}^{n-1} \alpha_k \frac{s^k}{k!} \quad (\text{Omnès '58})$

→  $\tan \delta_1^1(s) = \frac{\text{Im } F_V^{\pi(0)}(s)}{\text{Re } F_V^{\pi(0)}(s)} \quad (\text{Boito-Escribano-Jamin '09})$

→ Guerrero-Pich'97 delays breaking

$$F(s) = \frac{M_\rho^2}{M_\rho^2 - s - iM_\rho \Gamma_\rho(s)} \exp\left\{ \frac{-s}{96\pi^2 f_\pi^2} \times \left[ \text{Re } A(m_\pi^2/s, m_\pi^2/M_\rho^2) + \frac{1}{2} \text{Re } A(m_K^2/s, m_K^2/M_\rho^2) \right] \right\}$$

→ Instead we **resum** the **loop functions** in the **effective propagator** (Gómez-Dumm—PR '13, Celis-Cirigliano-Passemar '14)

$$F_V^{\pi(0)}(s) = \frac{M_\rho^2}{M_\rho^2 \left[ 1 + \frac{s}{96\pi^2 F_\pi^2} (A_\pi(s) + \frac{1}{2} A_K(s)) \right] - s} = \frac{M_\rho^2}{M_\rho^2 \left[ 1 + \frac{s}{96\pi^2 F_\pi^2} \text{Re}(A_\pi(s) + \frac{1}{2} A_K(s)) \right] - s - iM_\rho \Gamma_\rho(s)}$$

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(See Guerrero'98 for the reproduction of NNLO  $\chi$ PT results by this resummation)



# $\pi\pi$ VFF

$$F_V^\pi(s) = P_n(s) \exp\left\{\frac{s^n}{\pi} \int_{s_{\text{thr}}}^\infty ds' \frac{\delta_1^1(s')}{(s')^n (s' - s - i\epsilon)}\right\} \quad \log P_n(s) = \sum_{k=0}^{n-1} \alpha_k \frac{s^k}{k!} \quad (\text{Omnès '58})$$

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And we use the DR with n=3

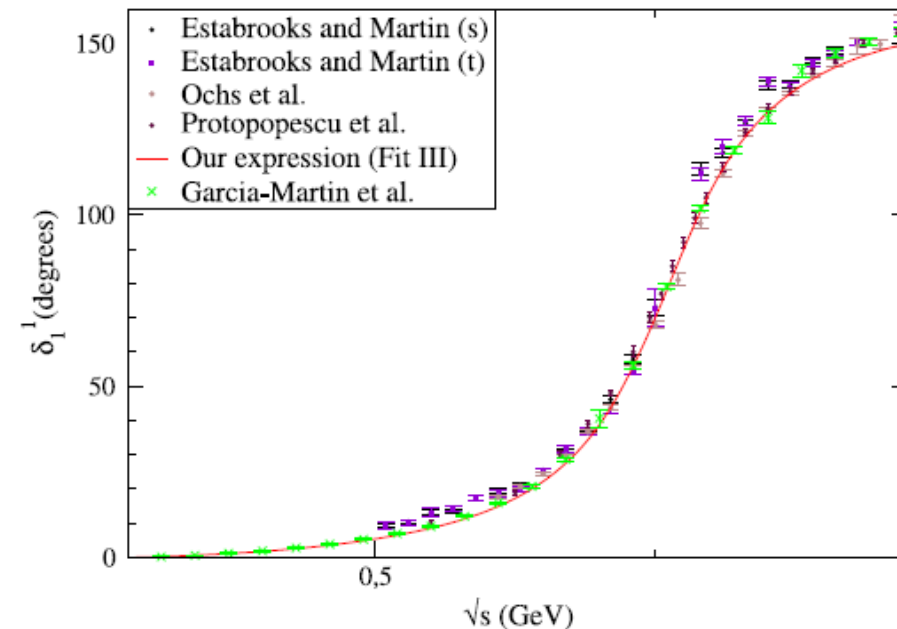
(See Guerrero'98 for the reproduction of NNLO  $\chi$ PT results by this resummation)

$$F_V^\pi(s) = \exp\left[\alpha_1 s + \frac{\alpha_2}{2} s^2 + \frac{s^3}{\pi} \int_{s_{\text{thr}}}^\infty ds' \frac{\delta_1^1(s')}{(s')^3 (s' - s - i\epsilon)}\right] \quad \delta_1^1(s \rightarrow \infty) = \pi$$

# $\pi\pi$ VFF

As a result of this setting, good description of phaseshift data (from  $\pi\pi$  scattering) is found [same for the modulus, which is shown later]

$F_V^\pi(s)$  is fitted and  $\delta_1^1(s)$  is prediction



**Fig. 1** Two-pion phase shift  $\delta_1^1$  as function of the  $\pi\pi$  invariant mass squared. Our theoretical expression (*red curve*) is shown to be in good agreement with experimental data (from Ochs et al. [65, 66], Estabrooks and Martin [67] in the  $s$  and  $t$  channels and Protopopescu et al. [68]) up to the opening of the two-kaon threshold,  $s_1 \simeq 1 \text{ GeV}^2$ . At very low energies, where no data are available, our prediction agrees with the results of García-Martín et al. [91] (Color figure online)

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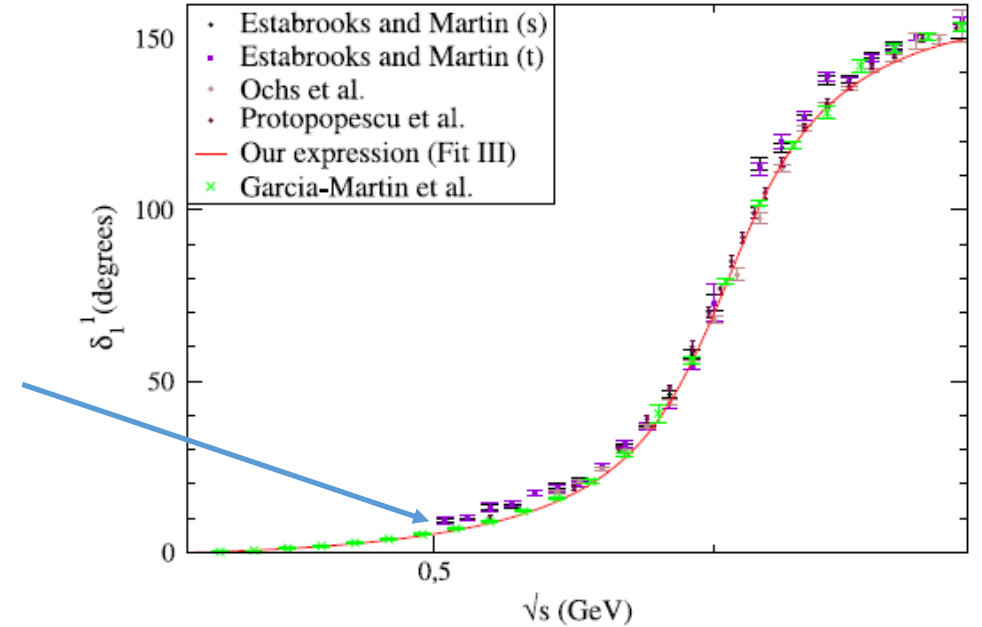
For low energies theory is ahead of experiment (also on the modulus)

$$\sqrt{s_{\pi/2}} = (775.0 \pm 0.2) \text{ MeV}$$

Our result

$$\sqrt{s_{\pi/2}} = (774 \pm 3) \text{ MeV}$$

Bern group



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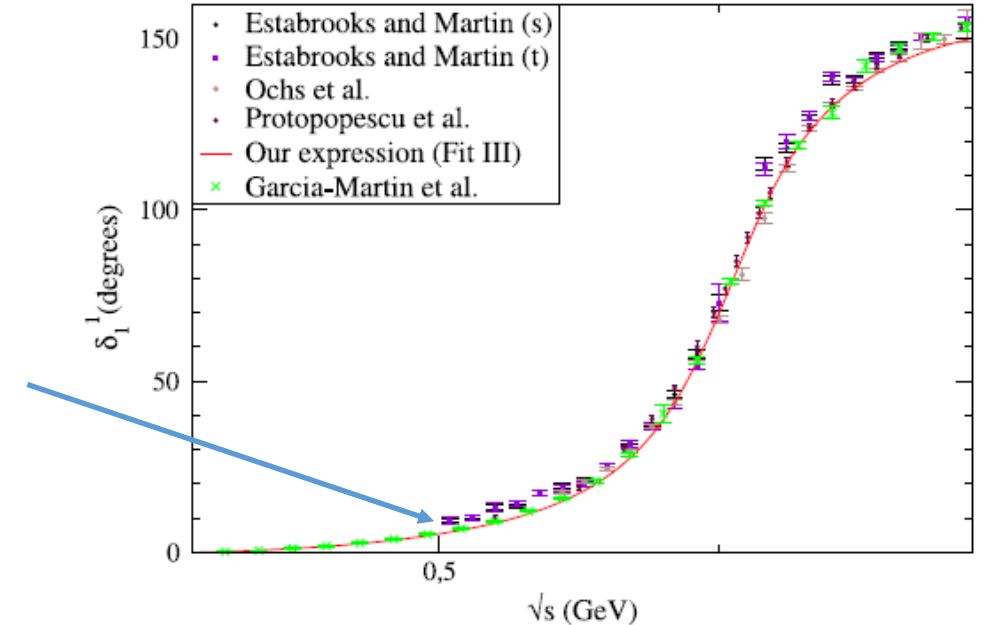
$\sqrt{s_{\pi/2}} = (774 \pm 3)$  MeV      Bern group

$$F_V^\pi(s) = 1 + \frac{1}{6} \langle r^2 \rangle_V^\pi s + c_V^\pi s^2 + d_V^\pi s^3 + \dots \quad \langle r^2 \rangle_V^\pi = 6\alpha_1, \quad c_V^\pi = \frac{1}{2}(\alpha_2 + \alpha_1^2)$$

Our values agree well with previous determinations

$$d_V^\pi = \frac{1}{6}(\alpha_3 + 3\alpha_1\alpha_2 + \alpha_1^3) = 9.84 \pm 0.05 \text{ GeV}^{-6}$$

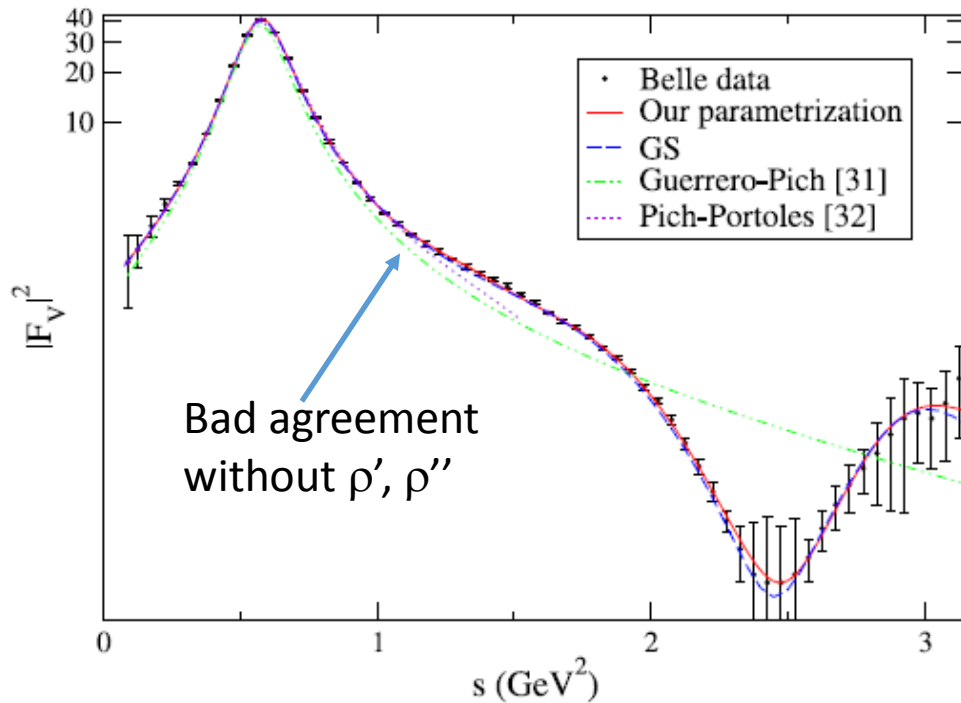
Previous determinations ( $\text{GeV}^{-6}$ ):  $9.70 \pm 0.40$ ,  $10.18 \pm 0.27$



**Fig. 1** Two-pion phase shift  $\delta_1^1$  as function of the  $\pi\pi$  invariant mass squared. Our theoretical expression (red curve) is shown to be in good agreement with experimental data (from Ochs et al. [65, 66], Estabrooks and Martin [67] in the  $s$  and  $t$  channels and Protopopescu et al. [68]) up to the opening of the two-kaon threshold,  $s_1 \simeq 1 \text{ GeV}^2$ . At very low energies, where no data are available, our prediction agrees with the results of García-Martín et al. [91] (Color figure online)

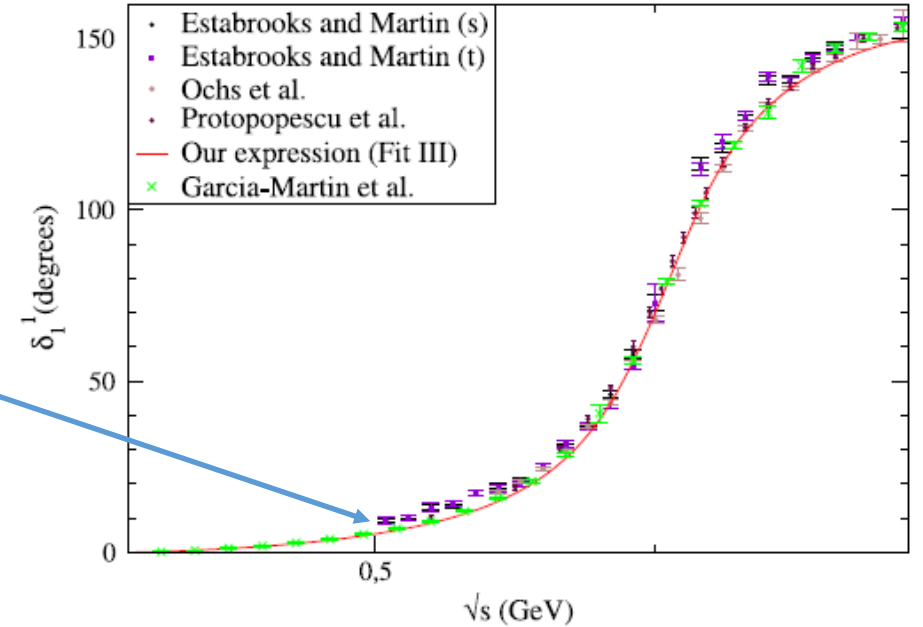
# $\pi\pi$ VFF

$F_V^\pi(s)$  is fitted and  $\delta_1^1(s)$  is prediction



**Fig. 2** Pion vector form factor  $F_V^\pi(s)$  compared to Belle data [3] (black dots). Solid and dashed lines correspond to our description and the GS parametrization, respectively. The dashed-dotted curve stands for the result from Ref. [42] (for  $M_\rho = 775$  MeV), while the dotted line corresponds to the dispersive representation in Ref. [43] (for  $\alpha_1 = 1.83$  GeV<sup>-2</sup>,  $\alpha_2 = 4.32$  GeV<sup>-4</sup> and  $M_\rho = 774.2$  MeV)

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# $\pi\pi$ VFF

$$F_V^{\pi(0)}(s) = \frac{M_\rho^2}{M_\rho^2[1 + \frac{s}{96\pi^2 F_\pi^2}(A_\pi(s) + \frac{1}{2}A_K(s))] - s} = \frac{M_\rho^2}{M_\rho^2[1 + \frac{s}{96\pi^2 F_\pi^2} \text{Re}(A_\pi(s) + \frac{1}{2}A_K(s))] - s - iM_\rho \Gamma_\rho(s)}$$

$$\Gamma_\rho(s) = -\frac{M_\rho s}{96\pi^2 F_\pi^2} \text{Im}\left[A_\pi(s) + \frac{1}{2}A_K(s)\right] = \frac{sM_\rho}{96\pi F_\pi^2} \left[\theta(s - 4m_\pi^2)\sigma_\pi^3(s) + \frac{1}{2}\theta(s - 4m_K^2)\sigma_K^3(s)\right]$$

(Gómez-Dumm—PR '13, Celis-Cirigliano-Passemar '14)

$$F_V^\pi(s) = \frac{M_\rho^2 + (\alpha' e^{i\phi'} + \alpha'' e^{i\phi''})s}{M_\rho^2[1 + \frac{s}{96\pi^2 F_\pi^2}(A_\pi(s) + \frac{1}{2}A_K(s))] - s} - \frac{\alpha' e^{i\phi'} s}{M_{\rho'}^2[1 + sC_{\rho'} A_\pi(s)] - s} - \frac{\alpha'' e^{i\phi''} s}{M_{\rho''}^2[1 + sC_{\rho''} A_\pi(s)] - s} \quad C_R = \frac{\Gamma_R}{\pi M_R^3 \sigma_\pi^3(M_R^2)}$$

$$M_\rho^{\text{pole}} = (760 \pm 2) \text{ MeV}, \quad \Gamma_\rho^{\text{pole}} = (147 \pm 6) \text{ MeV}$$

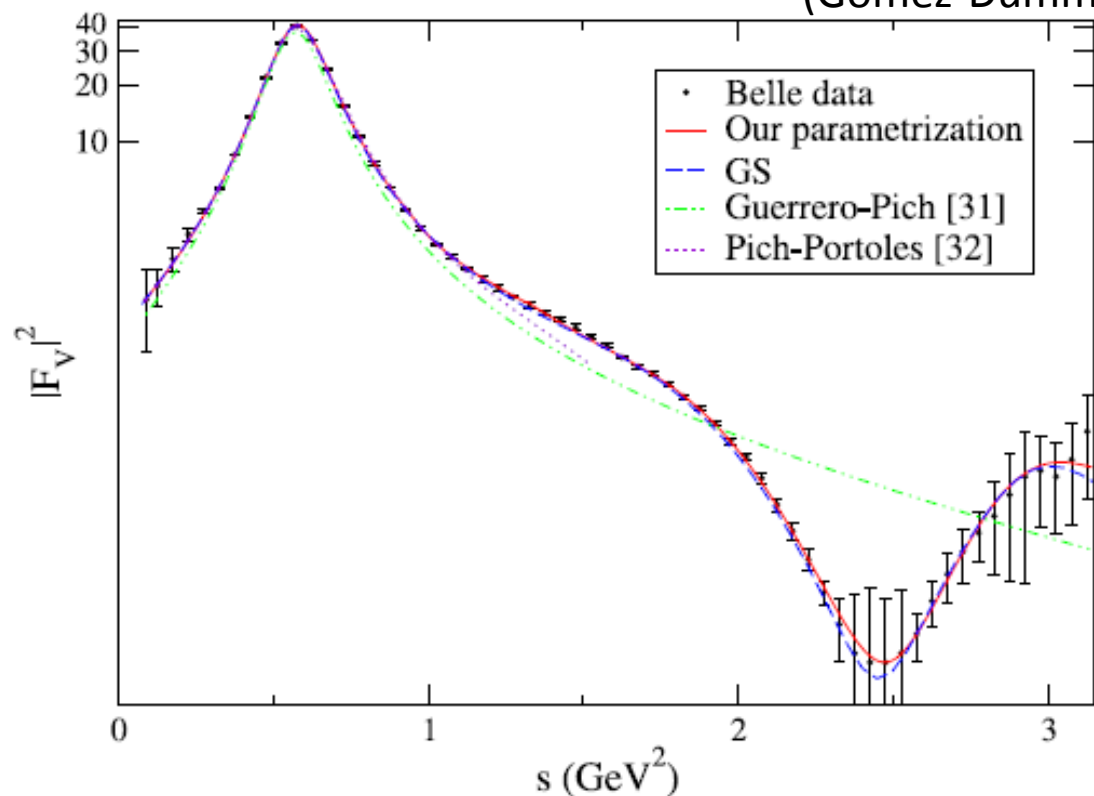
$$M_{\rho'}^{\text{pole}} = (1.44 \pm 0.08) \text{ GeV}, \quad M_{\rho''}^{\text{pole}} = (1.72 \pm 0.09) \text{ GeV}, \quad \alpha' = 0.08_{-0.01}^{+0.03}, \quad \phi' = 0.14_{-0.08}^{+0.10},$$

$$\Gamma_{\rho'}^{\text{pole}} = (0.32 \pm 0.08) \text{ GeV}, \quad \Gamma_{\rho''}^{\text{pole}} = (0.18 \pm 0.09) \text{ GeV}, \quad \alpha'' = 0.03 \pm 0.01, \quad \phi'' = 3.14_{-0.06}^{+0.50}.$$

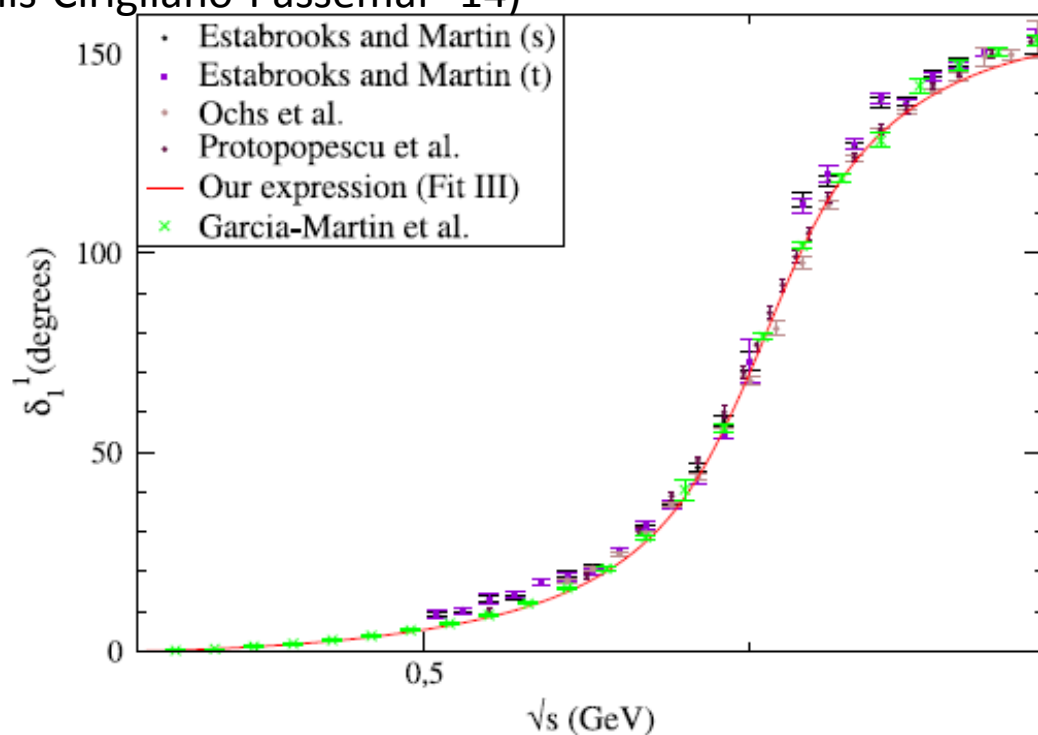
→ One can also use the DR up to some energy and match it to  $F_V^{\pi(0)}(s)$ . Results are almost undistinguishable [see later].

# $\pi\pi$ VFF

(Gómez-Dumm—PR '13, Celis-Cirigliano-Passemar '14)



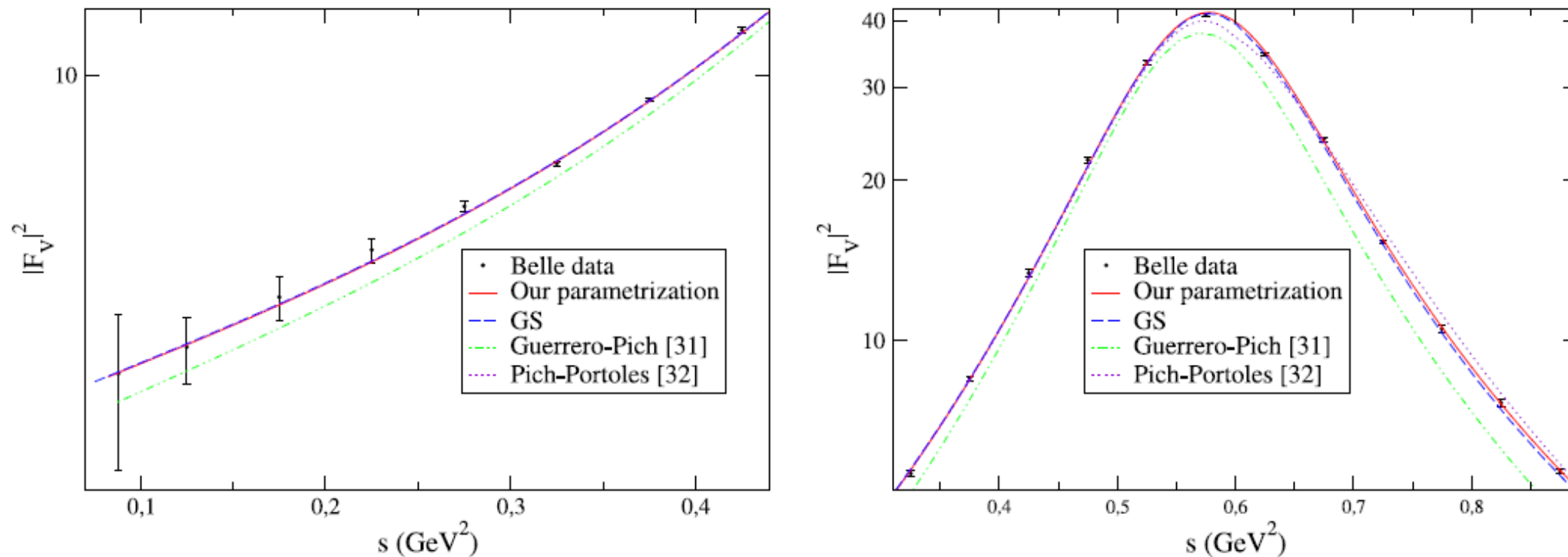
**Fig. 2** Pion vector form factor  $F_V^\pi(s)$  compared to Belle data [3] (black dots). Solid and dashed lines correspond to our description and the GS parametrization, respectively. The dashed-dotted curve stands for the result from Ref. [42] (for  $M_\rho = 775$  MeV), while the dotted line corresponds to the dispersive representation in Ref. [43] (for  $\alpha_1 = 1.83$  GeV $^{-2}$ ,  $\alpha_2 = 4.32$  GeV $^{-4}$  and  $M_\rho = 774.2$  MeV)



**Fig. 1** Two-pion phase shift  $\delta_1^1$  as function of the  $\pi\pi$  invariant mass squared. Our theoretical expression (red curve) is shown to be in good agreement with experimental data (from Ochs et al. [65, 66], Estabrooks and Martin [67] in the  $s$  and  $t$  channels and Protopopescu et al. [68]) up to the opening of the two-kaon threshold,  $s_1 \simeq 1$  GeV $^2$ . At very low energies, where no data are available, our prediction agrees with the results of García-Martín et al. [91] (Color figure online)

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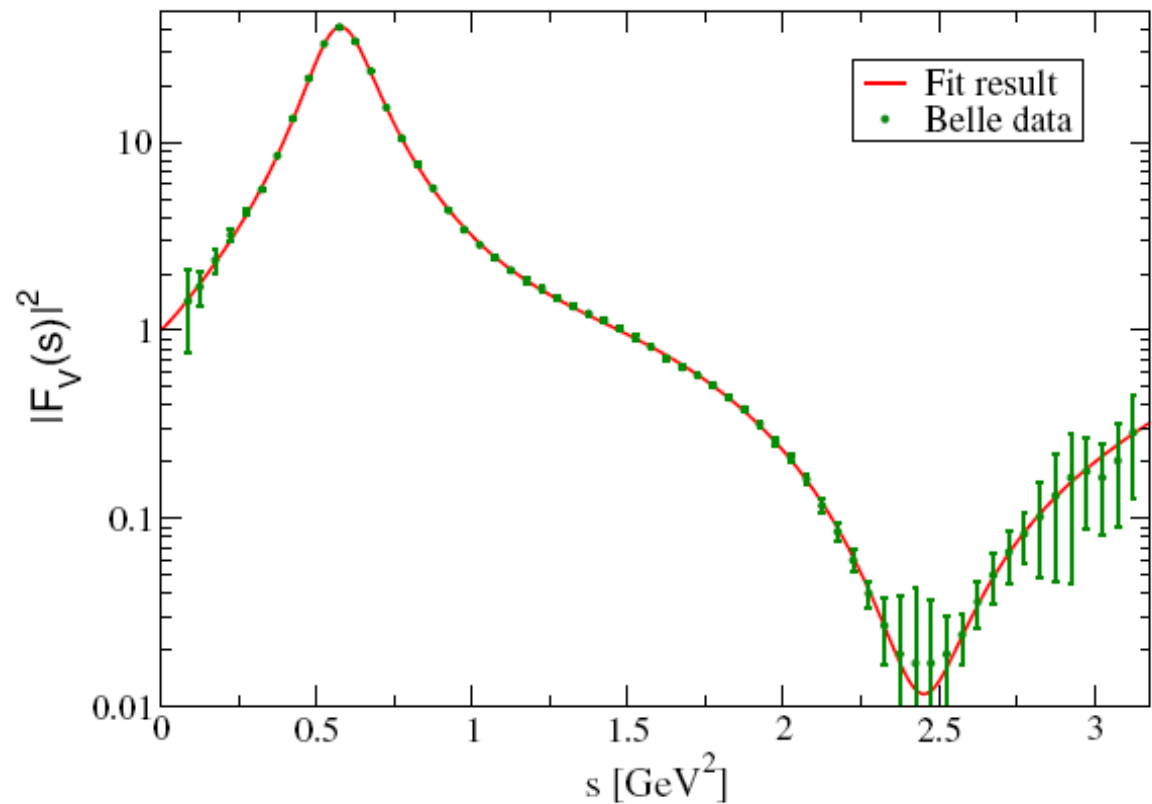
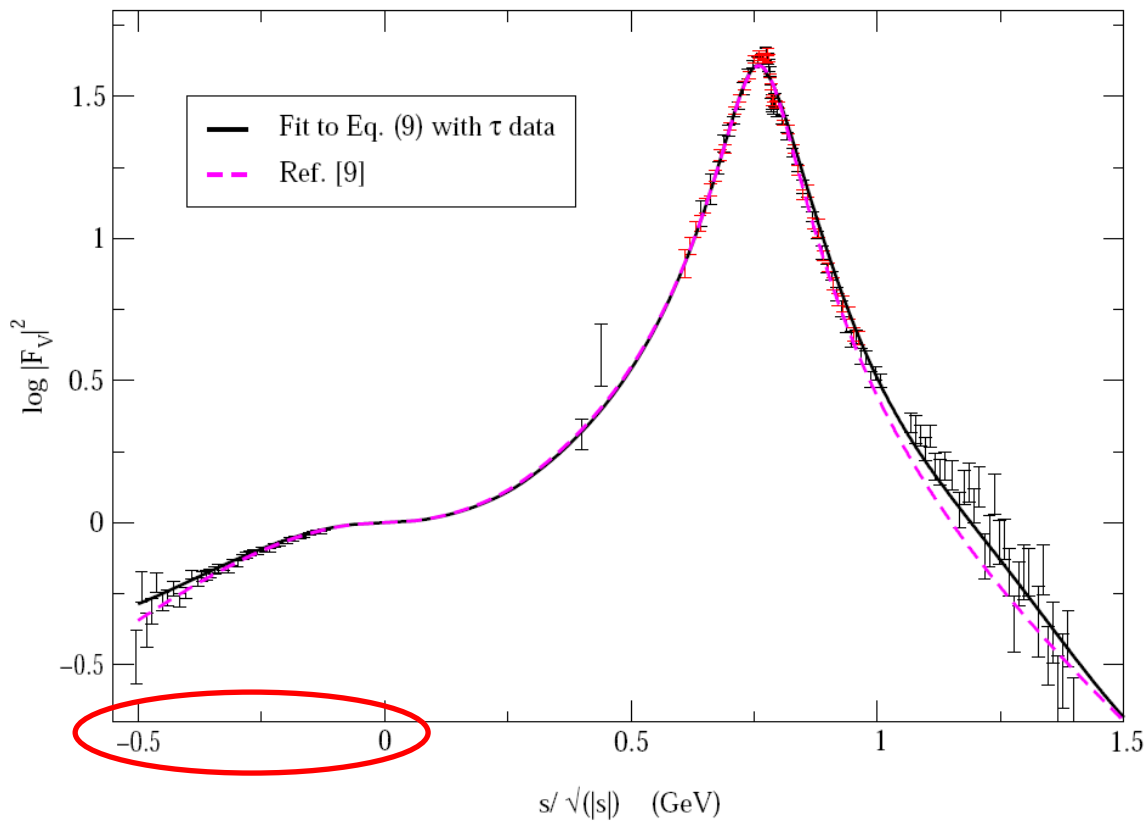
**Fig. 3** Two close-ups of Fig. 2 are displayed, corresponding to the low-energy region (*left panel*) and the peak region (*right panel*)

N.B. : GS has wrong phaseshift, bad high-E behaviour, excited resonances do not decouple at low-E ( $\chi$  limit is lost), ...



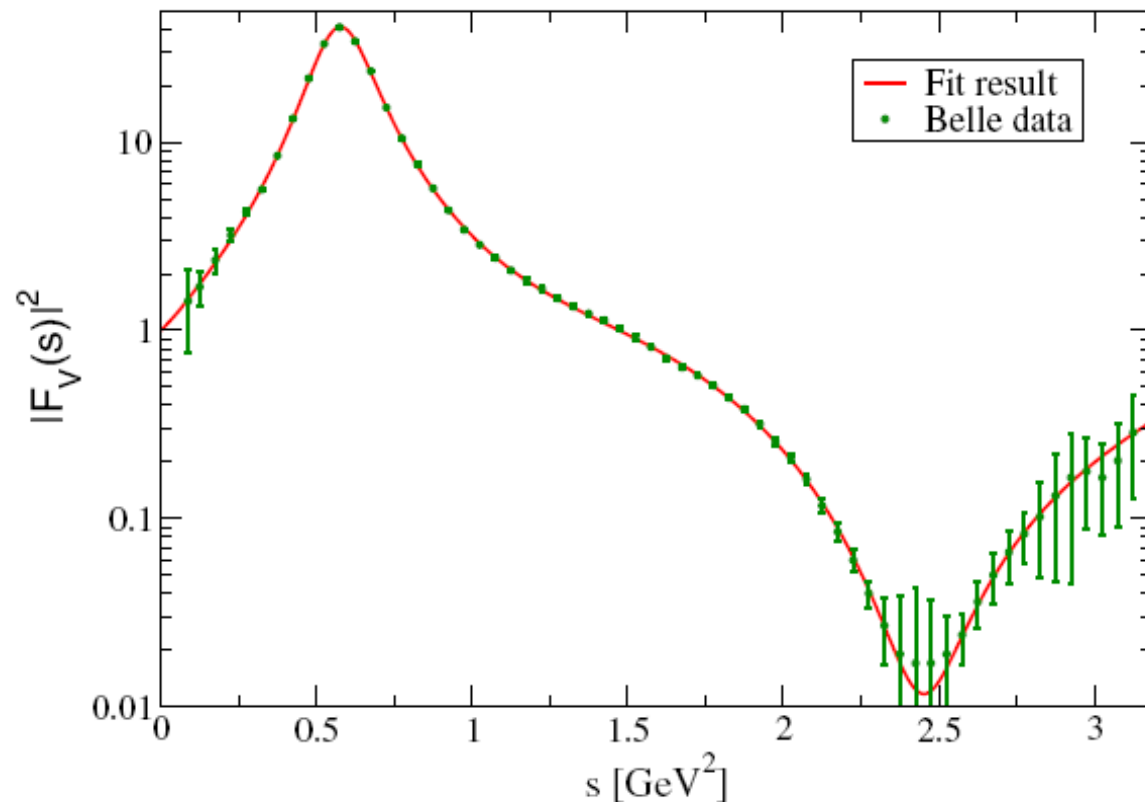
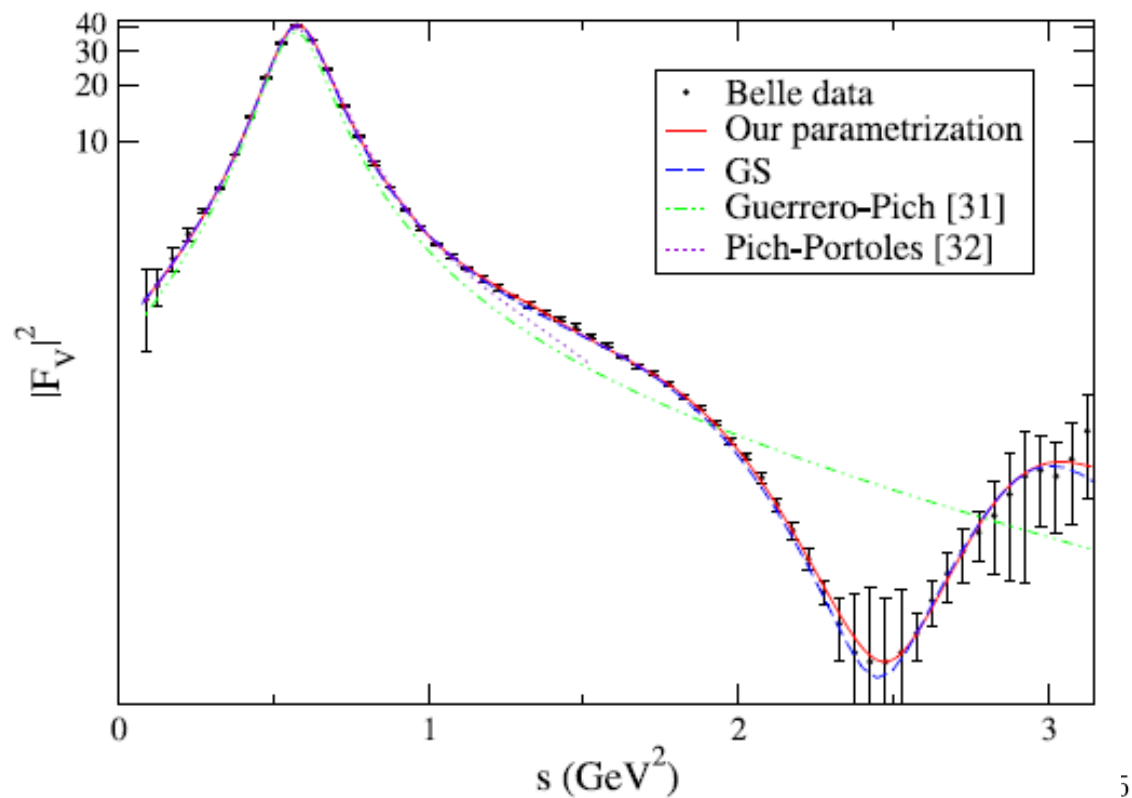
# $\pi\pi$ VFF

(Pich-Portolés '01, Gómez-Dumm—PR '13, Celis-Cirigliano-Passemar '14)



# $\pi\pi$ VFF

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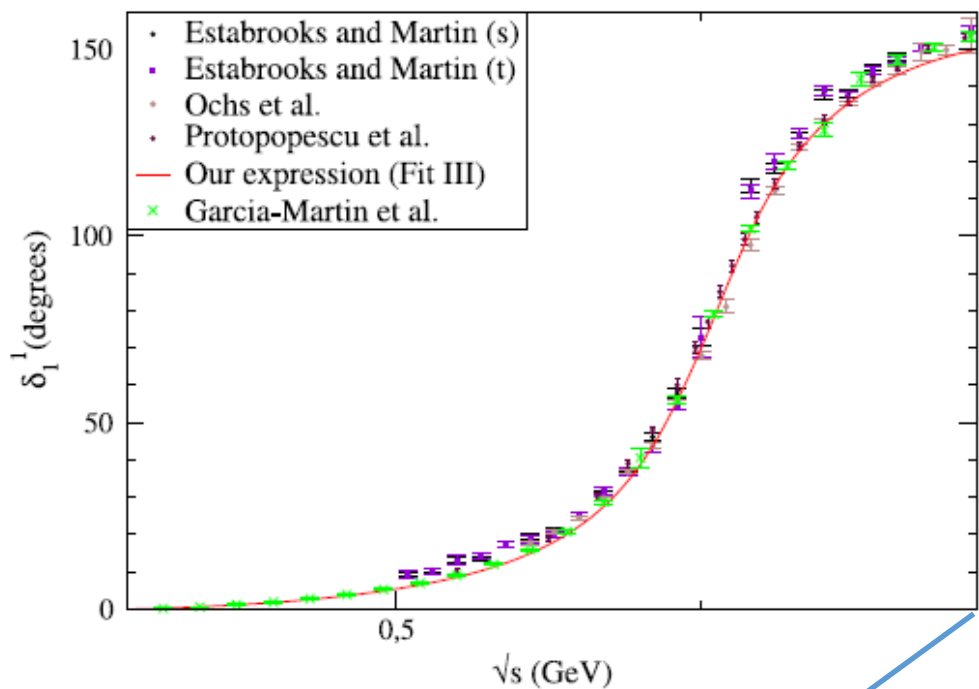
(Gómez-Dumm—PR '13, Celis-Cirigliano-Passemar '14)

→ According to data it might seem that little else remains to be done ( $\rho$ - $\omega$  mixing & other isospin corrections are 'easy').

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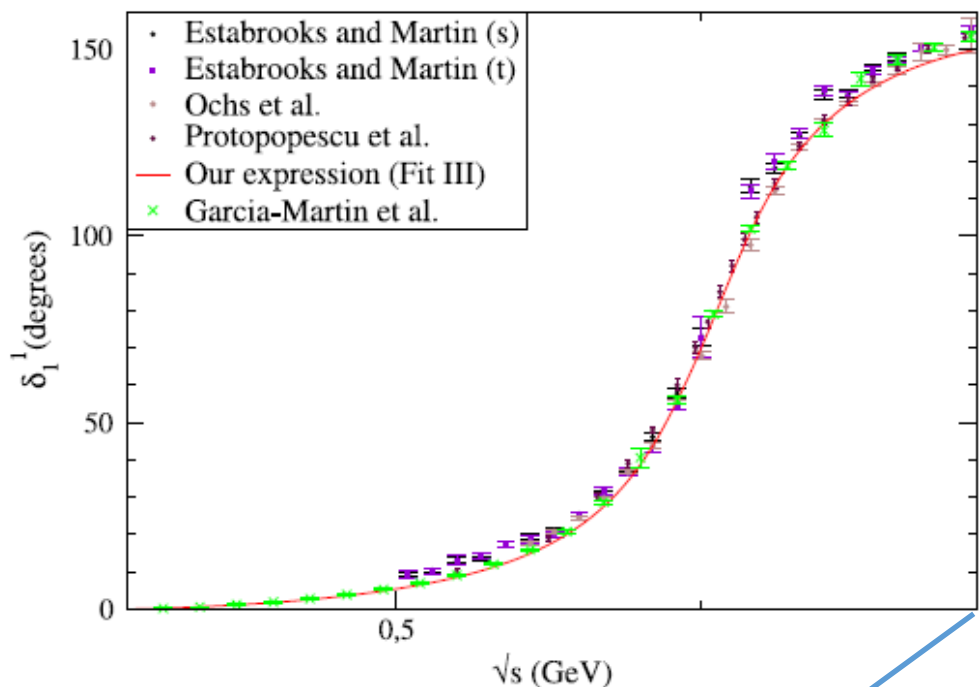


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According to data it might seem that little else remains to be done.



Inelasticities need to be included & a **coupled-channels** treatment is needed!

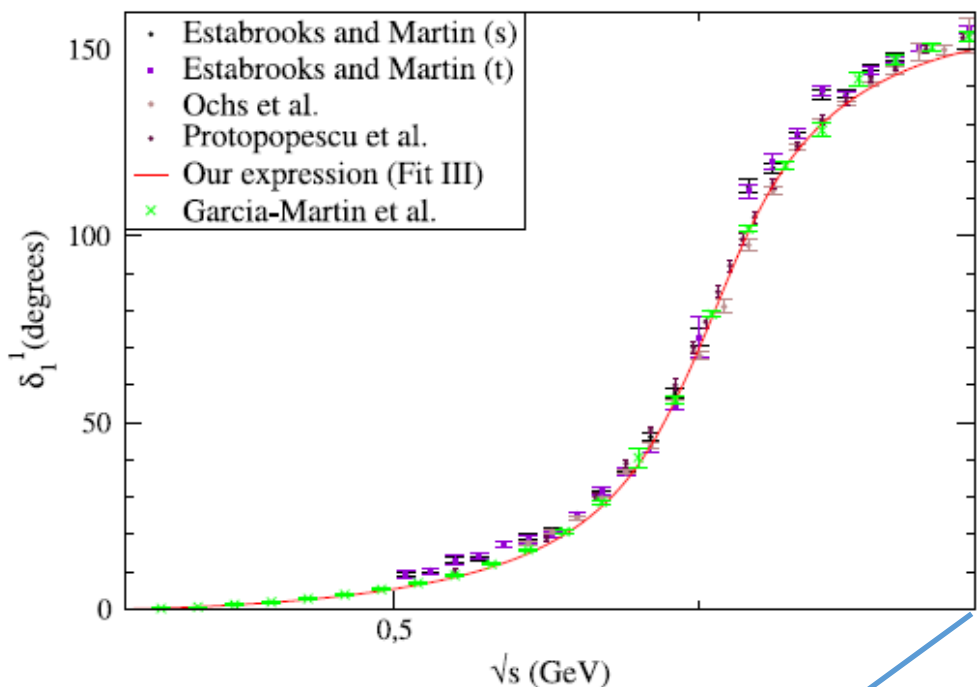
According to our experience in SFFs I believe that is feasible and, moreover, quite fast (see also Moussallam '08 and Kubis et. al.).

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As a by-product we would get the unitarized KK VFF

$$F_V^{K^+K^-}(s) = \frac{1}{2} \frac{M_\rho^2}{M_\rho^2 - s - iM_\rho\Gamma_\rho(s)} \exp \left[ 2 \operatorname{Re} \left( \tilde{H}_{\pi\pi}(s) \right) + \operatorname{Re} \left( \tilde{H}_{KK}(s) \right) \right] + \frac{1}{2} \left[ \sin^2 \theta_V \frac{M_\omega^2}{M_\omega^2 - s - iM_\omega\Gamma_\omega} + \cos^2 \theta_V \frac{M_\phi^2}{M_\phi^2 - s - iM_\phi\Gamma_\phi} \right] \times \exp \left[ 3 \operatorname{Re} \left( \tilde{H}_{KK}(s) \right) \right]$$

$$F_V^{K^0\bar{K}^0}(s) = -\frac{1}{2} \frac{M_\rho^2}{M_\rho^2 - s - iM_\rho\Gamma_\rho(s)} \exp \left[ 2 \operatorname{Re} \left( \tilde{H}_{\pi\pi}(s) \right) + \operatorname{Re} \left( \tilde{H}_{KK}(s) \right) \right] + \frac{1}{2} \left[ \sin^2 \theta_V \frac{M_\omega^2}{M_\omega^2 - s - iM_\omega\Gamma_\omega} + \cos^2 \theta_V \frac{M_\phi^2}{M_\phi^2 - s - iM_\phi\Gamma_\phi} \right] \times \exp \left[ 3 \operatorname{Re} \left( \tilde{H}_{KK}(s) \right) \right]$$

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(Arganda-Herrero-Portolés '08, Guerrero-Oller '99)

# Other two-meson VFFs: $K\pi$ , $(KK)$

Again, a lot of work on the subject: Jamin-Pich-Portolés '06, '08; Moussallam '08; Boito-Escribano-Jamin '09, '10; Escribano--González-Solís—Jamin--Roig '14;; Jamin-Oller-Pich '00, '00, '02,'06; Bernard-Kaiser-Meissner '91, (Buettiker)--Descotes-Genon—Moussallam ('04), '06

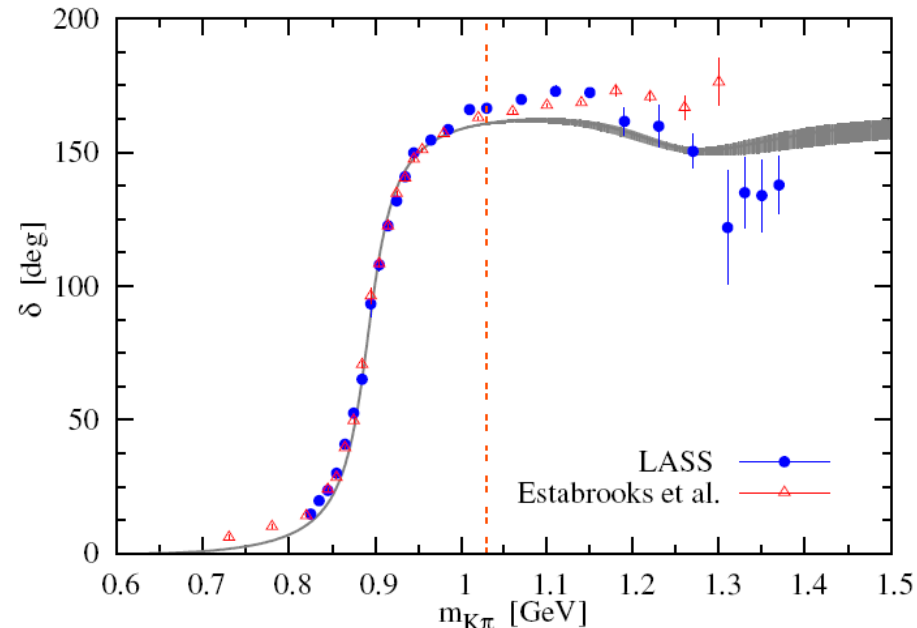
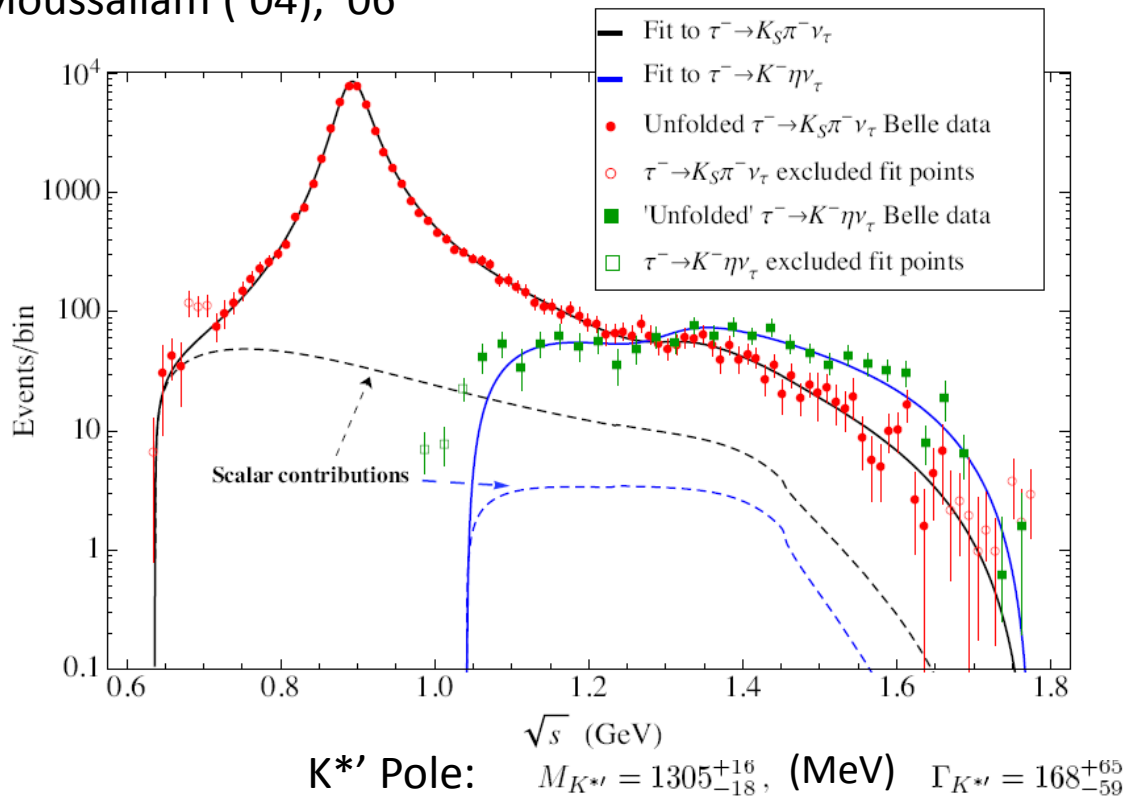
$$\tilde{f}_+^{K\pi}(s) = \exp \left[ \alpha_1 \frac{s}{M_{\pi^-}^2} + \frac{1}{2} \alpha_2 \frac{s^2}{M_{\pi^-}^4} + \frac{s^3}{\pi} \int_{s_{K\pi}}^{s_{\text{cut}}} ds' \frac{\delta_1^{K\pi}(s')}{(s')^3 (s' - s - i0)} \right] \quad \tan \delta_1^{K\pi}(s) = \frac{\text{Im} \tilde{f}_+^{K\pi}(s)}{\text{Re} \tilde{f}_+^{K\pi}(s)}$$

$$\tilde{f}_+^{K\pi}(s) = \frac{m_{K^*}^2 - \kappa_{K^*} \tilde{H}_{K\pi}(0) + \gamma s}{D(m_{K^*}, \gamma_{K^*})} - \frac{\gamma s}{D(m_{K^*}, \gamma_{K^*})} \quad D(m_n, \gamma_n) = m_n^2 - s - \kappa_n \tilde{H}_{K\pi}(s) \quad \kappa_n = \frac{192\pi}{\sigma_{K\pi}(m_n^2)^3} \frac{\gamma_n}{m_n}$$

$$\tilde{H}_{K\pi}(s) = \frac{2}{3} \tilde{H}_{K^0\pi^-}(s) + \frac{1}{3} \tilde{H}_{K^-\pi^0}(s)$$

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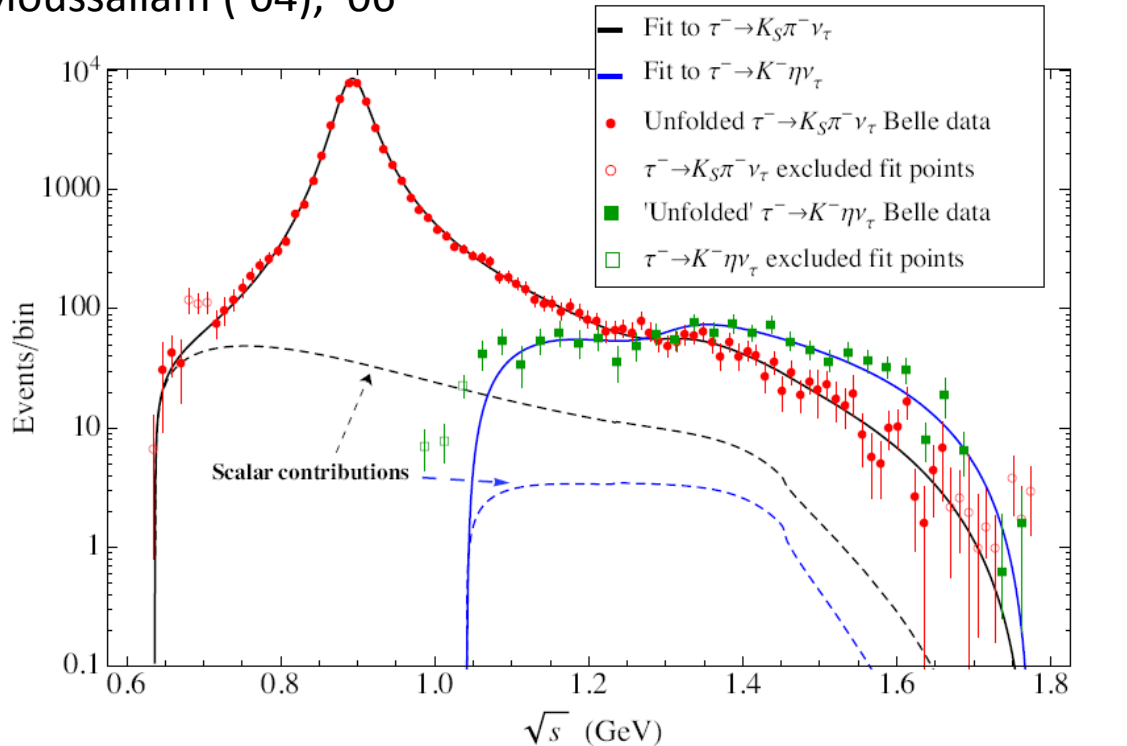


**Figure 2:** Phase of the form factor  $F_+(s)$  together with experimental results from LASS [46] and Estabrooks *et al.* [47]. The opening of the first inelastic channel,  $K^*\pi$ , is indicated by the dashed vertical line. The gray band represents the extrema from the fits of Tab. 3.



# Other two-meson VFFs: $K\pi$ , $(KK)$

Again, a lot of work on the subject: Jamin-Pich-Portolés '06, '08; Moussallam '08; Boito-Escribano-Jamin '09, '10; Escribano--González-Solís—Jamin--Roig '14;; Jamin-Oller-Pich '00, '00, '02,'06; Bernard-Kaiser-Meissner '91, (Buettiker)--Descotes-Genon—Moussallam ('04), '06



$K^{*}$  Pole:  $M_{K^{*}} = 1305_{-18}^{+16}$ , (MeV)  $\Gamma_{K^{*}} = 168_{-59}^{+65}$

→ A coupled-channel analysis was done for the SFF(s) but remains to be done for the VFF(s)

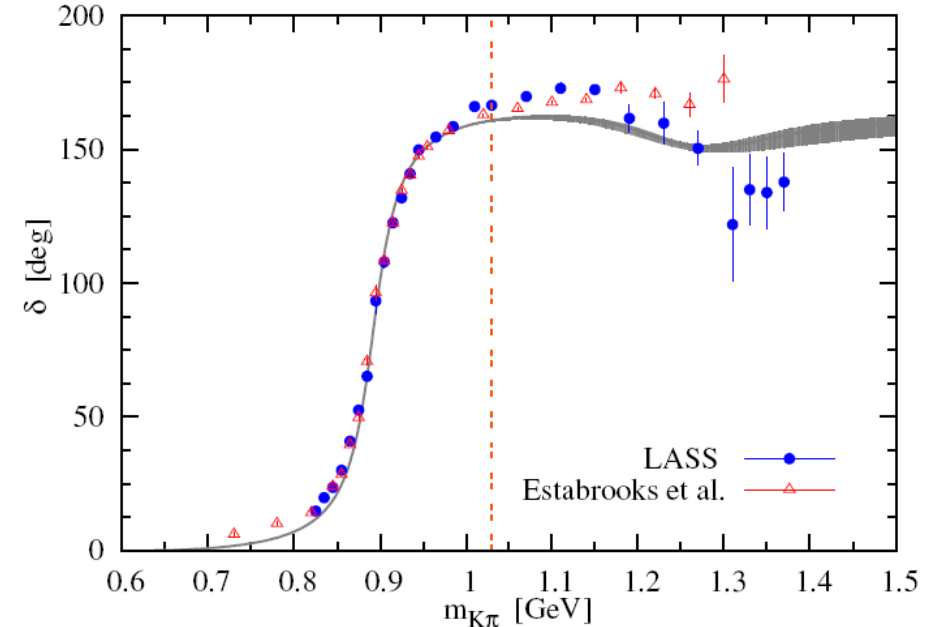
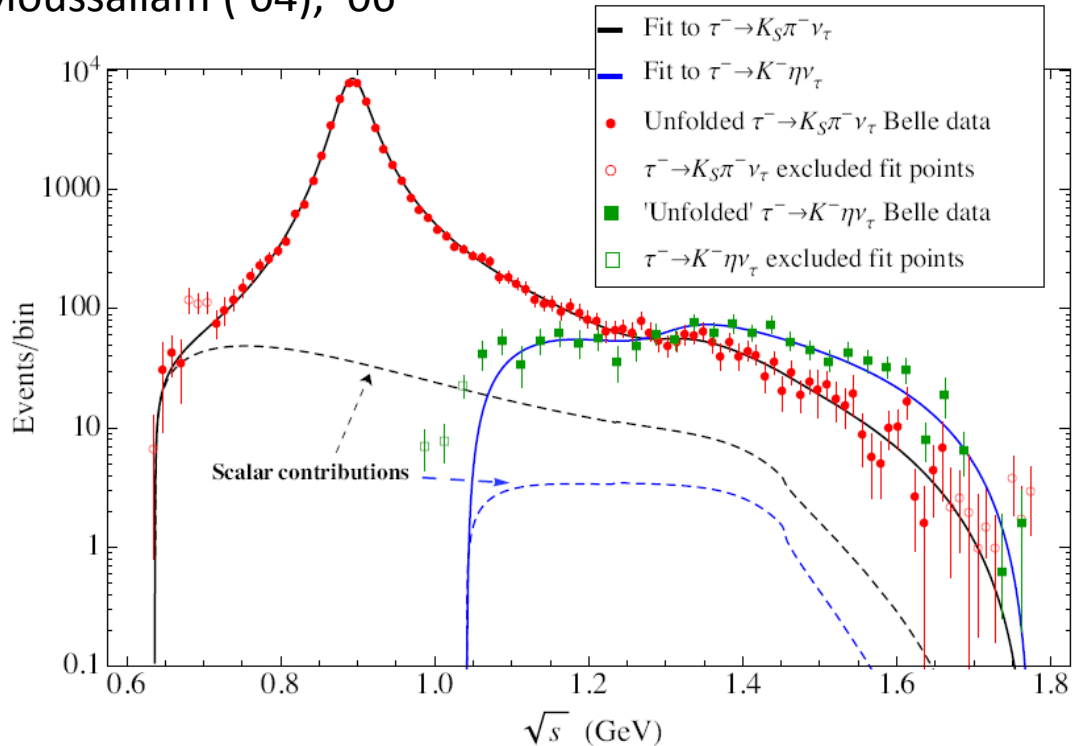


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→ We also have the  $K\eta$ ,  $K\eta'$ ,  $\pi\eta$ ,  $\pi\eta'$  (S)VFFs available, in case you need them

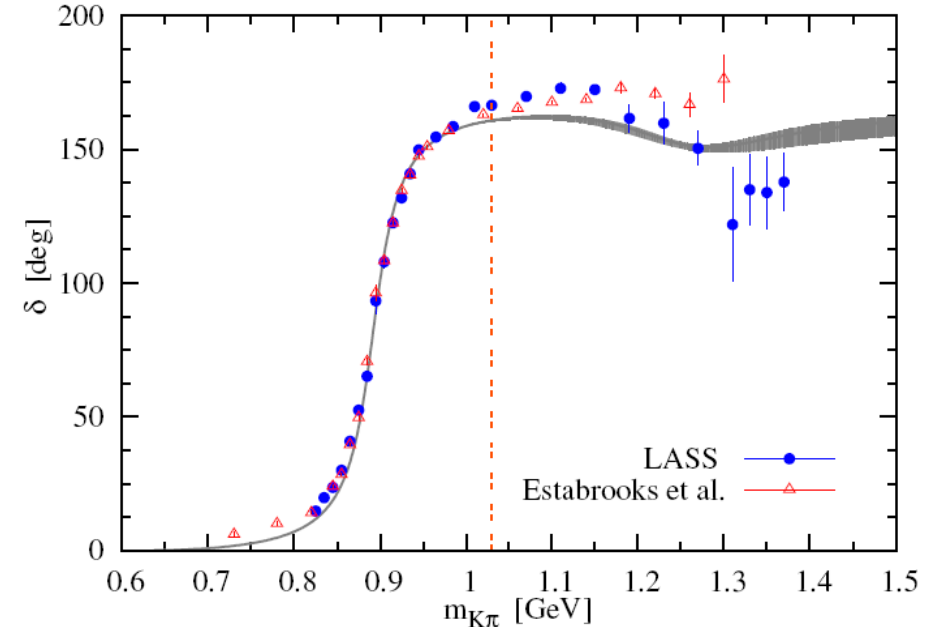
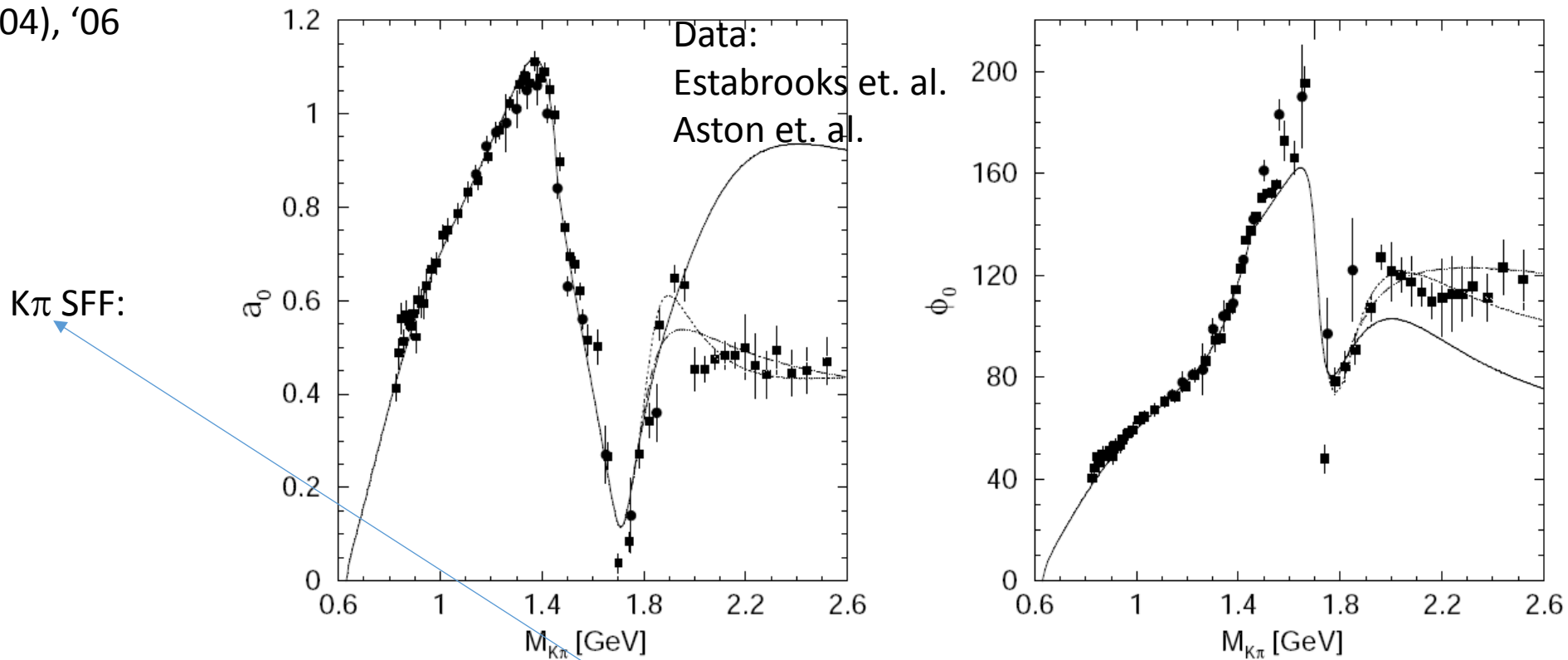


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- All corresponding SFFs obtained within c-c are at your disposal.





# LatinoAmerican Workshop on High Energy Physics: Particles and Strings

Havana, 15-21 July 2012

Me gusta 0 Tweet +

## Overview

The **Latinoamerican Workshop on High Energy Physics: Particles and Strings** aims to bring together scientist of Europe and Latin-american countries with the purpose of develop and strengthen scientific links and collaborations between the various physics communities. The program of the conference will include a set of review lectures on the following topics:

- Physics of the Standard Model and Beyond.
- Theories of fundamental interactions: QCD, Nuclear physics, Astrophysics, Dark matter.
- LHC: Experiments and first signatures.
- String phenomenology
- D-branes, holography, black hole physics.

The program will also includes informal talks given by young scientist from Latin-American countries and short talks by other participants in the meeting. Open discussions on scientific areas of common interests will be included.



### Special Information

- [This is the Workshop Programme](#)
- [Workshop Poster.](#)
- [Full Talk List and Materials.](#)
- [Picture Gallery 2012.](#)

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Though the figure of Ernest Hemingway is...  
1 2 3 4 5 6 7  
8 next › last »



# Latin-American Conference on High Energy Physics: Particles and Strings II

Havana, 18-22 July 2016

- Home
- Program
- Registration
- Participants
- The Venue
- Accommodation

## Overview

The **Latin-American Conference on High Energy Physics: Particles and Strings II**, aims to bring together scientist of Europe and Latin-american countries with the purpose of develop and strengthen scientific links and collaborations between the various physics communities. The program of the conference will include a set of review lectures on the following topics:

- Physics of the Standard Model and Beyond.
- Theories of fundamental interactions: QCD.
- LHC: Experiments and first signatures.
- String phenomenology, amplitudes.
- D-branes, holography, black hole physics.



The program will also include short talks, informal discussion sessions and a student gong show.

The conference is the second of a conference series starting with the ["Latin-American Workshop on High Energy Physics: Particle Strings"](#), Havana, July 21, 2012.

## Special Information

- Registration from December 01, 2015
- Registration deadline: May 31, 2016

# Scientific Program

The program of the conference will include review lectures, short talks, discussion sessions and a student gong show. We request speakers to keep all talks at the pedagogical level, focusing on the main conceptual insights and results in the field reviewed in accessible and self-consistent terms.

## Invited Speakers (confirmed):

- **Albert De Roeck** (CERN) " *Searches at the LHC: present and future* "
- **Michelangelo Mangano** (CERN) " *Future Circular Collider* "
- **Mark Wise** (Caltech) " *Effective Theories and Physics Beyond the Standard Model* "
- **Fernando Quevedo** (ICTP) "String Phenomenology"
- **Dieter Lust (ICTP)** (MPIP) "Classical and Quantum Black Hole Hair"
- **Eva Silverstein** (Stanford) " *Horizon physics in string theory* "
- **Burt Ovrut** (Pennsylvania) "The Minimal Supersymmetric B-L Model--from Unification to the LHC"
- **Massimo Bianchi** (Rome) "(Soft) String amplitudes"
- **Margarete Muehleitner** (KIT) " *Phenomenology of Beyond the Standard Model at the LHC* "
- **Thomas Mueller** (KIT) " *Collider physics and detector performance: present and future accelerators* "
- **Jonas Rademacker** (Bristol) " *Heavy flavour physics at the LHC* "
- **Gino Isidori** (Zurich) "Flavor Physics beyond the SM"
- **Jose Ocariz** (CERN/LPNHE) " *Higgs Physics at the LHC* "
- **Sameer Murthy** (King's college) "Quantum black holes in string theory"
- **Rafael Nepomechie** (Miami) "Integrability in AdS/CFT"

## Invited speakers (to be confirmed):

- **Michael Peskin** (SLAC) " *Standard Model and Symmetry Breaking* "
- **Lisa Randall** (Harvard) " *TBA* "
- **Joseph Lykken** (Fermilab) " *Neutrino Physics* "
- **Marcela Carena** Fermilab/Chicago) " *Phenomenology of Higgs boson models at the LHC* "
- **Javier Redondo** (Zaragoza) "Dark matter candidates: Axions"
- **Sean Hartnoll** (Stanford) " *TBA* "
- **Alberto Lerda** (Turin/Alessandria) " *TBA* "



We'll have it in Mexico  
in summer 2017!!





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THANK  
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