Precision Tests At High Energy

Francesco Riva - CERN

In collaboration with Liu, Pomarol, Rattazzi (appear next week)

Also Contino, Falkowski, Goertz, Grojean (YR4) and Biekotter, Knochel, Kraemer, Liu (2014)

Why EFT Parametrization?

Necessary for precision tests:

Motivation (SM test -> New Physics Search)

Organisation

(e.g. E/M expansion indicates hierarchy between departures from SM)

Self-Consistency Check perturbativity of physical expansion

perturbativity of physical expansion (no need to invoke unitarity)

What is the problem?

Validity: E/M << 1But experimental access to CE^2/M^2

Wilson coefficient

Can EFT validity be established model-independently?

No. Question on EFT validity depends on (broad) BSM hypotheses.

Example: Fermi theory
$$\frac{2}{v^2} \bar{\psi}_{\nu_\mu} \gamma^\mu \psi_\mu \bar{\psi}_{\nu_e} \gamma^\mu \psi_e$$
 is it valid ip to v=246 GeV?

No, only to
$$E=m_W=rac{g}{2}vpprox81~{
m GeV}$$

- * Weak couplings reduce the validity range of the EFT (as naively expected)
- * Strong couplings extend it (for g=41 Fermi theory ok to 3 TeV!)
 - → LHC relies on this

How can we know about couplings without resorting to simplified models?

* Scale

EFT is actually an expansion in

- * Coupling
- * Numerical coefficients

 \rightarrow 0(1) naturally

→ for symmetries <<1

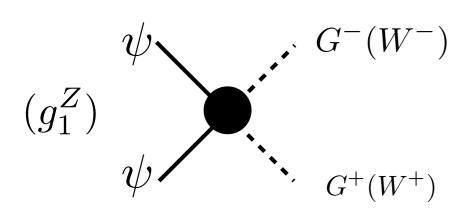
- * Perivatives and fields -> Powers of 1/M (matches units of energy)
- * Every field \rightarrow One power of coupling (matches units of \hbar)

$$\mathcal{L}_{\text{eff}} = \frac{M^4}{g_*^2} \mathcal{L}\left(\frac{D_{\mu}}{M}, \frac{g_H H}{M}, \frac{g_{\Psi} \Psi_{L,R}}{M^{3/2}}, \frac{g_V F_{\mu\nu}}{M^2}\right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \cdots$$

$$rac{g_*^2}{M^2} H^\dagger D_\mu H ar{\Psi} \gamma^\mu \Psi \ (g_1^Z)$$

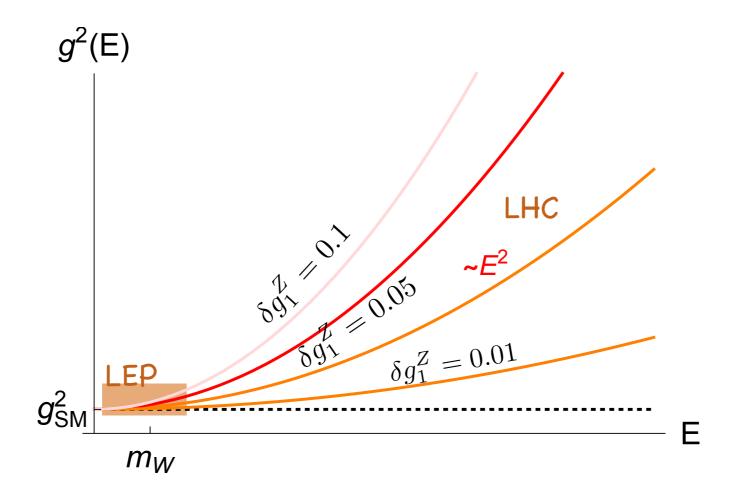
$$\frac{g_*}{M^2} \epsilon_{abc} W^{a \nu}_{\mu} W^{b}_{\nu \rho} W^{c \rho \mu}$$

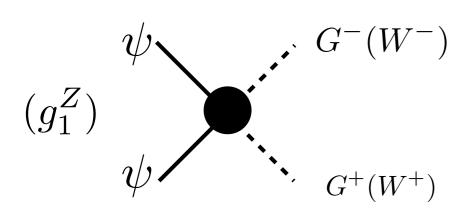
$$(\lambda_{\gamma})$$



Anomalous Coupling Version

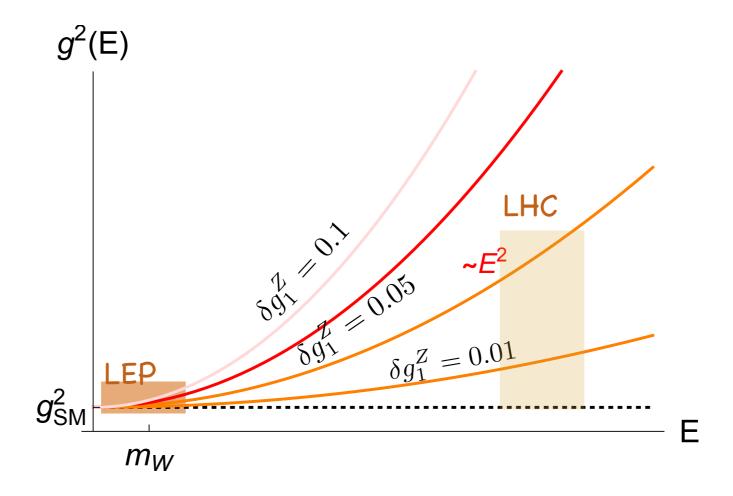
$$A \simeq g^2 \left(1 + \frac{\delta g_1^Z}{g} \frac{E^2}{m_W^2} \right) \equiv g^2(E)$$

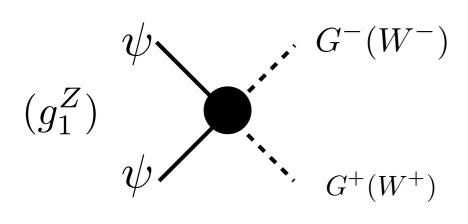




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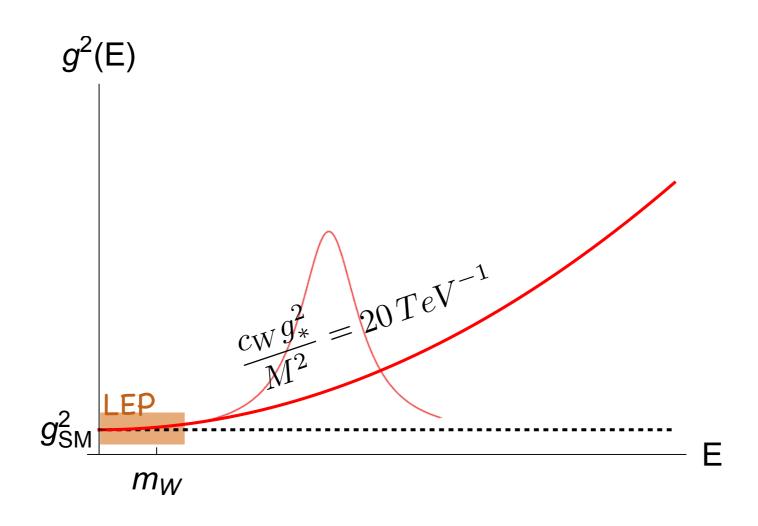
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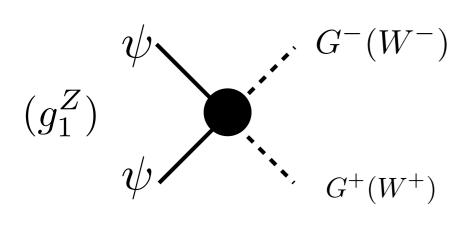




EFT Version

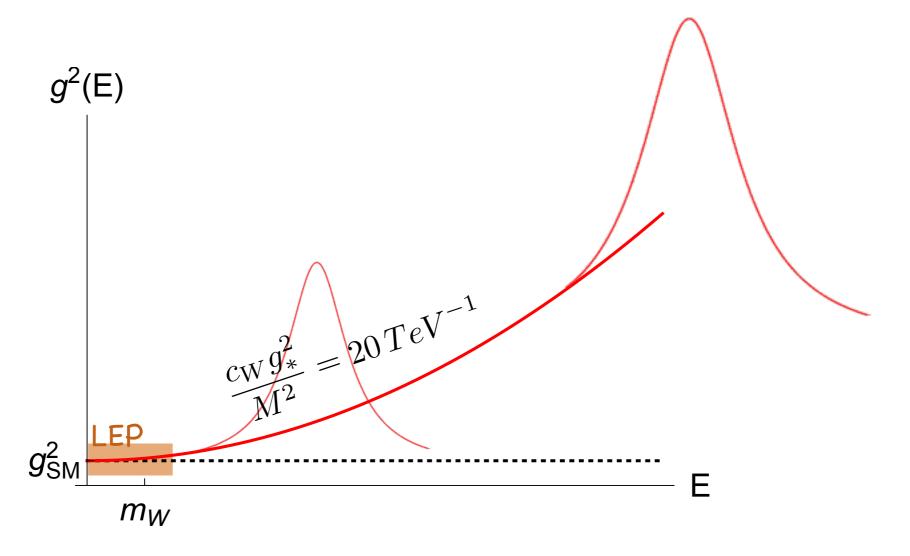
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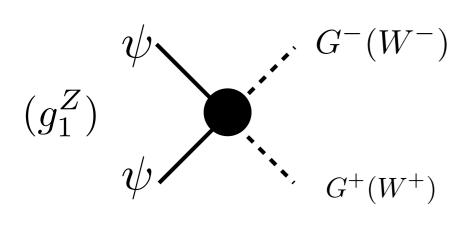




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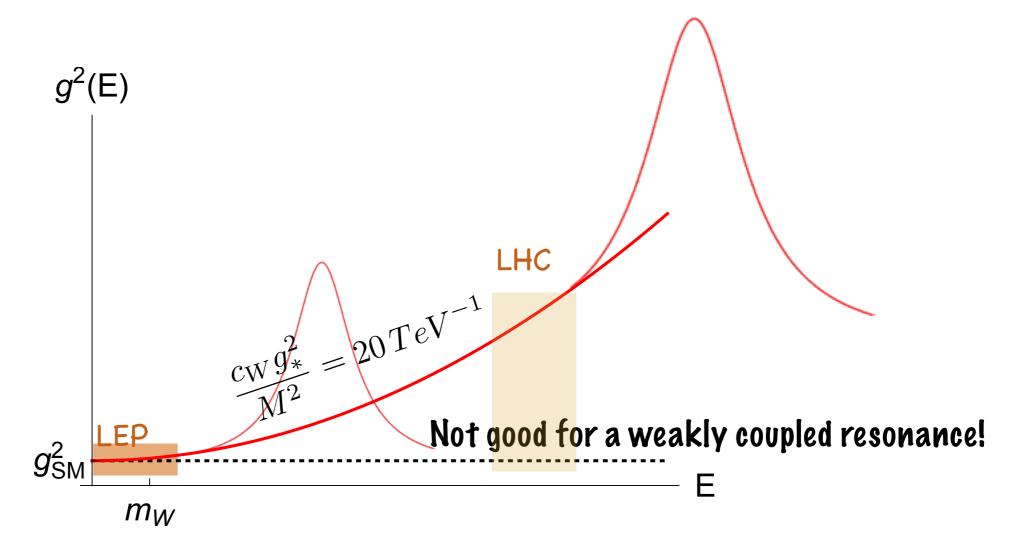
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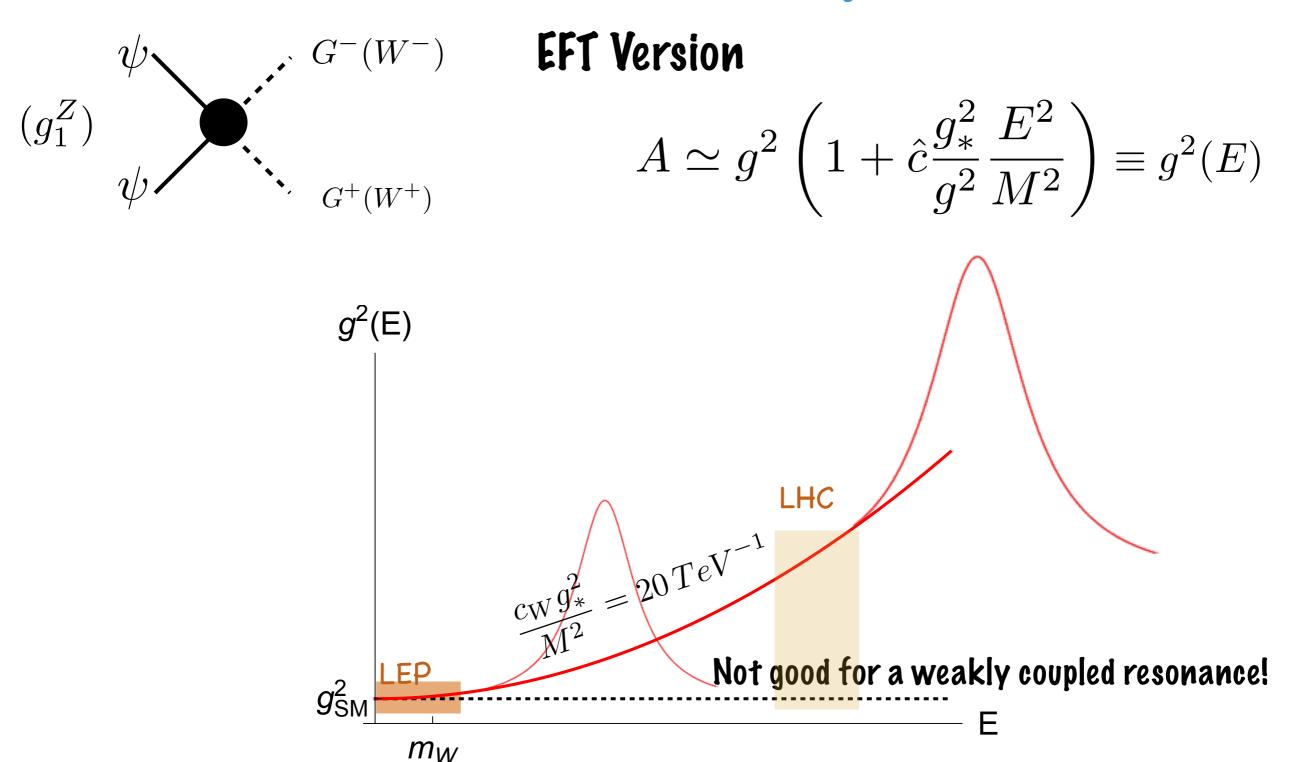




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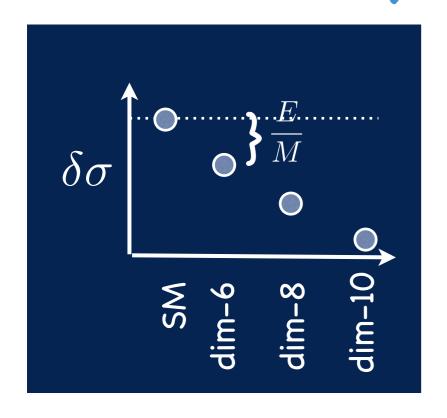
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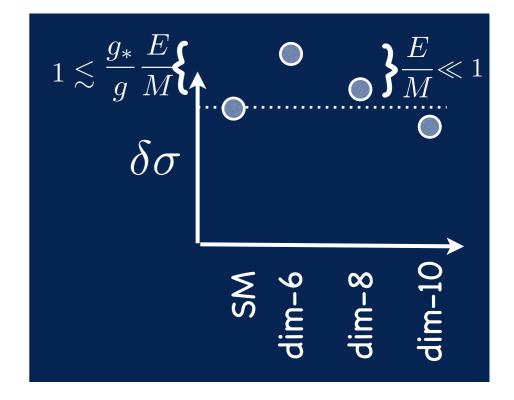


Although similar constraint as LEP on $cg*^2/M^2$, the LHC one consistent only for large g*>g

But if BSM>SM, then the EFT brakes down?

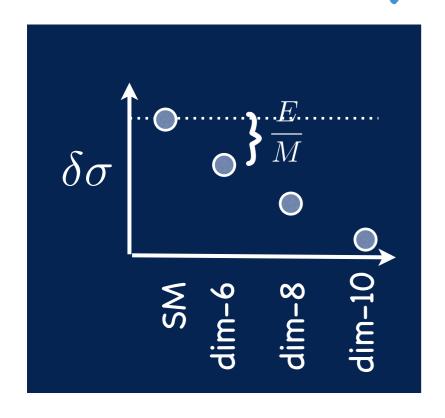


Not necessarily.

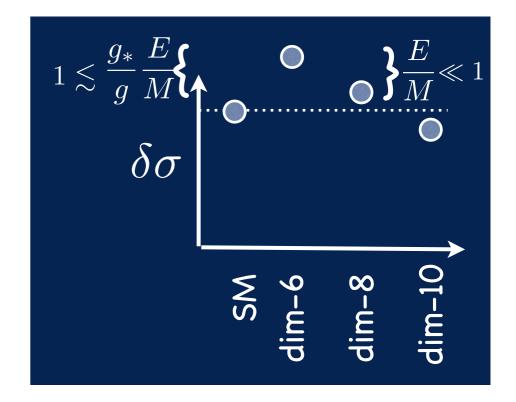


Strong Coupling g*/g can undo energy expansion E/M

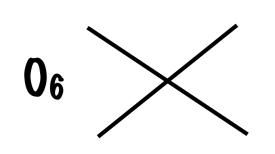
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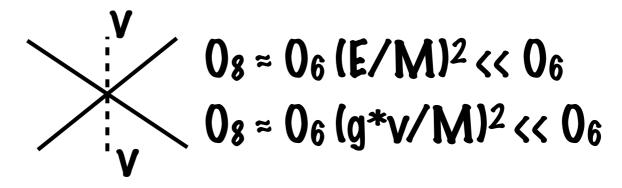


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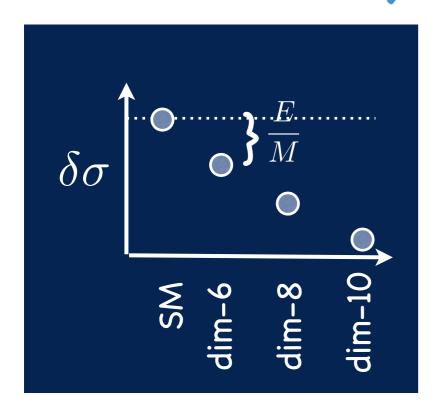


Strong Coupling g*/g can undo energy expansion E/M Dim-6 vs. Dim-8 contributing to same vertex:

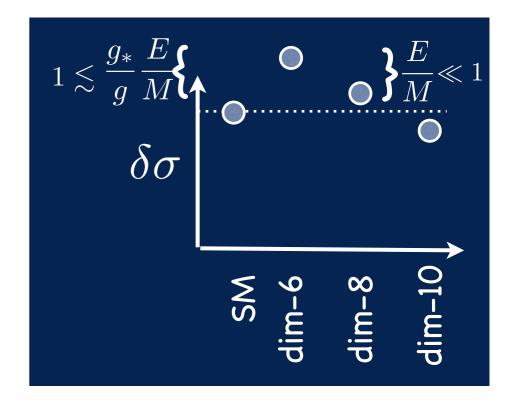




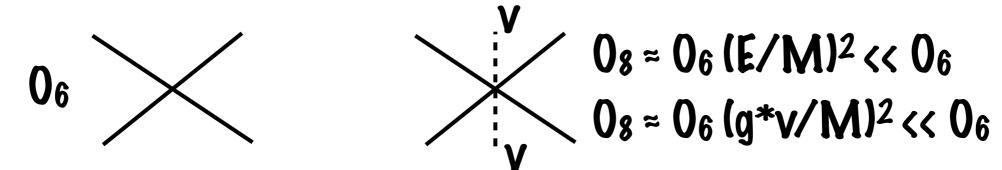
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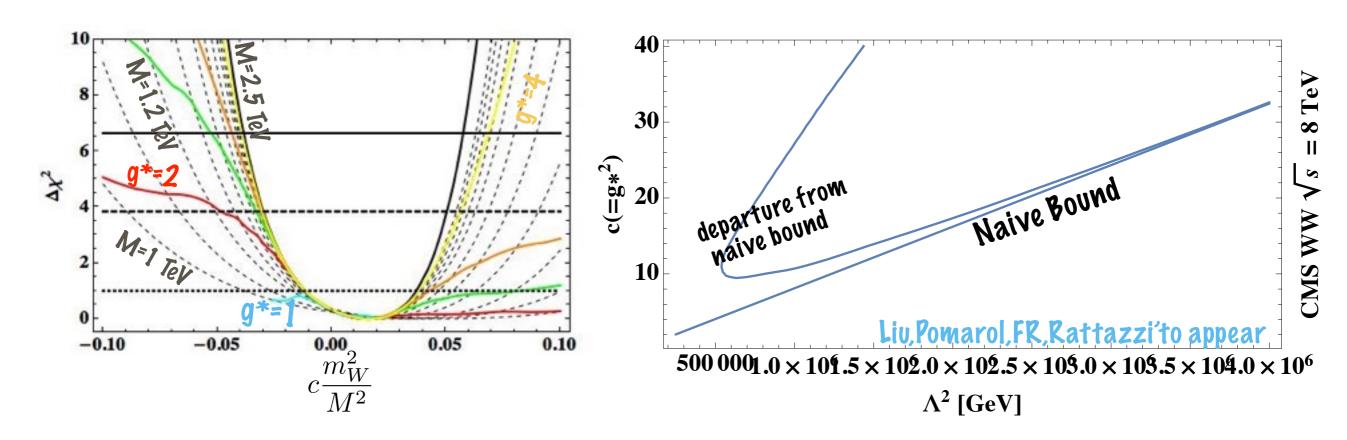
- → In crossection dim- 6^2 have the same powers of $(1/M^4)$ as dim-8xSM ...but more powers of $g^*!$
- \rightarrow It's always important to keep this dim-6² term: it's either small and negligible or large and requires strong coupling assumption

Recipe for incorporating this into an experimental analysis?

- * Unlike LEP, energy unknown...
- * Repeat analysis with extra cut on $\sqrt{s} < M$, for different values of M (if CoM energy unknown: similar techniques as DM monojet searches)

Racco, Wulzer, Zwirner' 2015; Brugisser, Mahbubani, FR, Urbano' to appear

* Bounds on c/M^2 can be shown for different g^* using $c=g^2$ and in (g^*,M) plane



-> Contains all information in terms of transparent physical parameters

Target:

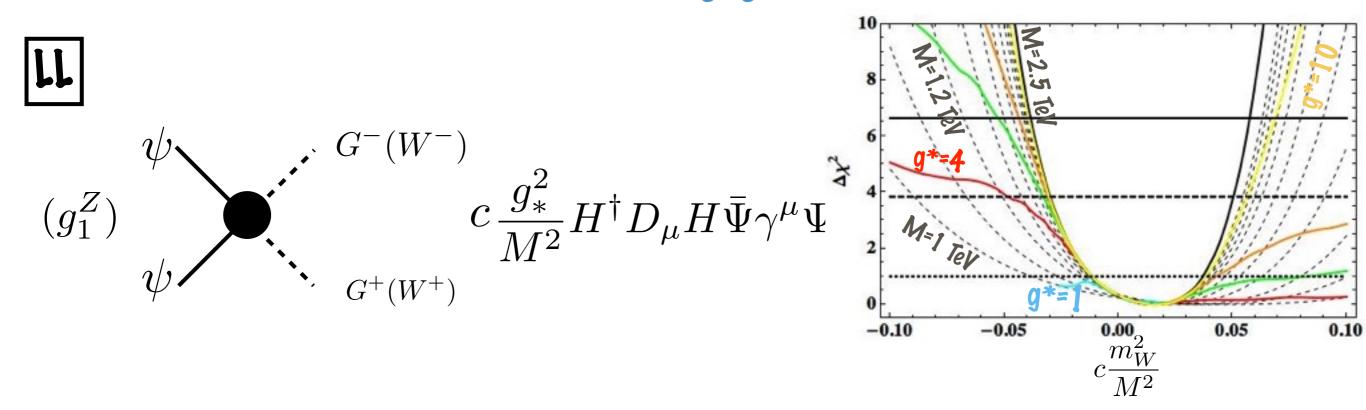
Many interesting BSM theories are weakly coupled:

$$A \simeq g^2 \left(1 + \hat{c} \frac{g_*^2}{g^2} \frac{E^2}{M^2} \right)$$
0(1)

they require sensitivity to SM effects also at high energy!

This is the ultimate target for LHC experiments

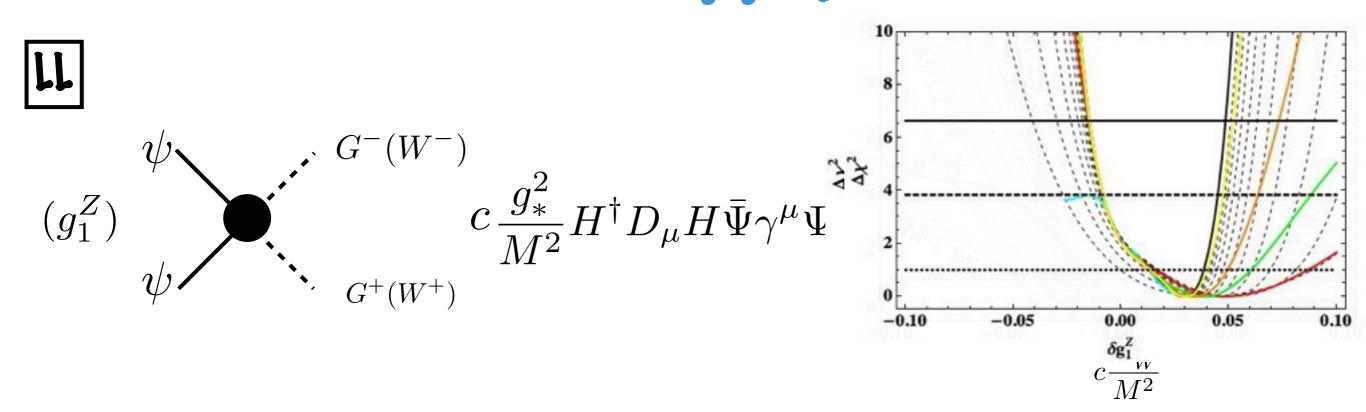
How does this apply to aGC?



-Strong Coupling necessary (otherwise no consistent bounds)

Improvable?

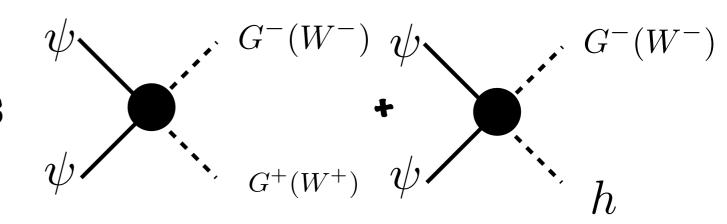
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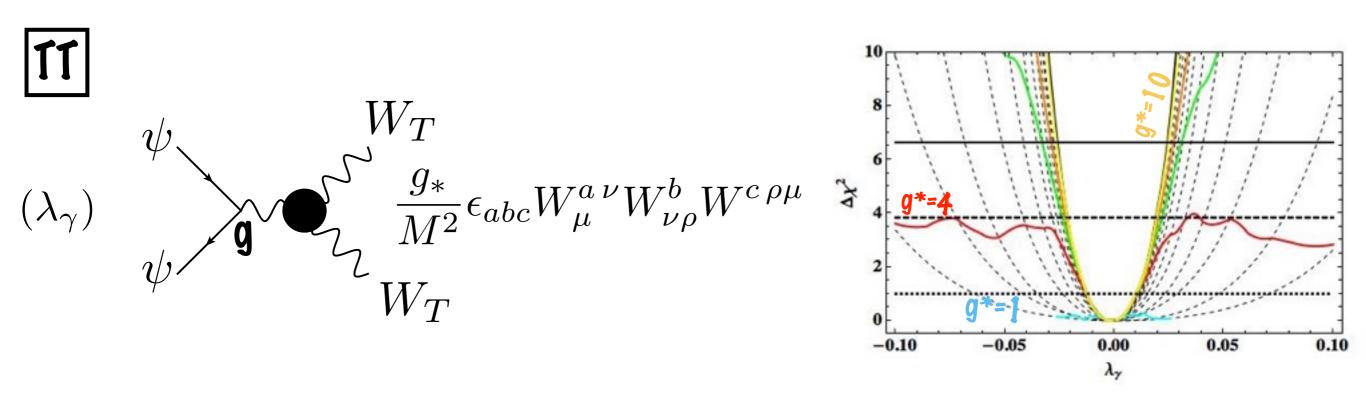
Improvable?

- Combination with Higgs physics



- Enhance signal through polarization tagging (In SM TT=10xLL)

How does this apply to aGC?



-Strong Coupling necessary (otherwise no consistent bounds)

Liu, Pomarol, FR, Rattazzi'to appear

Important: only one power of g*!

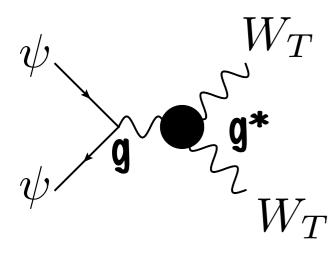
$$A \simeq g^2 \left(1 + \hat{c} \frac{g_*}{g} \frac{E^2}{M^2} \right)$$

Summary so far

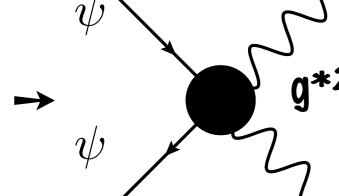
- Sensitivity to effects larger than SM at high-E
 → consistent with EFT only for strong coupling
- * Recipe: Analysis of c/M^2 with different cutoffs Using $c=g^{*2}$ or $c=g^*$, the results are presentable in (g^*,M) -plane and contain all info
- * This can be done consistently (quadratic terms in crossection can be kept) and no dimension-8 operators must be kept

...however:

New Signals for aGC



g* always involves a weak coupling!
$$A \simeq g^2 \left(1 + \hat{c} \frac{g_*}{g} \frac{E^2}{M^2}\right)$$



g*2 in some theories dimension-8 can have g*2 $W_T = i \frac{g_*^2}{M^4} \bar{\psi}_{L,R} \gamma^\mu D^\nu \psi_{L,R} W^a_{\mu\rho} W^{a\,\rho}_\nu$

$$W_T \qquad i \frac{g_*^2}{M^4} \bar{\psi}_{L,R} \gamma^\mu D^\nu \psi_{L,R} W^a_{\mu\rho} W^a_\nu$$

$$A \simeq g^2 \left(1 + \hat{c} \frac{g_*^2}{g^2} \frac{E^4}{M^4} \right)$$

- -> there are theories where this dominates
- → New effects with neutral vectors only (Z,photons) in final state
- -> Modifications of helicity +- amplitude in TGCs
- → EFT expansion under control

Conclusions

- Sensitivity to dimension-6 effects larger than SM at high-E
 → consistent with EFT for strong coupling
- * Recipe: Analysis of c/M^2 with different cutoffs Using $c=g^{*2}$ or $c=g^*$, the results are presentable in (g^*,M) -plane and contain all info
- * Target: higher sensitivity to effects comparable to the SM
- * In the same class of theories, there are regions where dimension-8 can be consistently studied and dominate, despite EFT valid.