

TMDs from MC evolution

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- Calculation of TMDs
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Summary

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- ▶ $xf(x, t)$,
- ▶ LO in $P(z)$,
- ▶ 1-loop- α_s .

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- ▶ **DGLAP evolution**,
- ▶ $xf(x, t)$,
- ▶ LO in $P(z)$,
- ▶ 1-loop- α_s .

We show:

- ▶ **integrated PDFs** and compare them with semi analytical results from **QCDNum17**,
- ▶ first results for **TMDs**.

Introduction to the method

Sudakov formalism

Evolution equation for parton density

$$t \frac{\partial f(x, t)}{\partial t} = \frac{\alpha_s}{2\pi} \int \frac{dz}{z} P(z) f\left(\frac{x}{z}, t\right) - \frac{\alpha_s}{2\pi} f(x, t) \int dz P(z). \quad (1)$$

Introducing *Sudakov form factor*

$$\Delta_s(t, t_0) \equiv \Delta_s(t) = \exp\left(- \int_x^{z_{max}} dz \int_{t_0}^t \frac{\alpha_s}{2\pi} \frac{dt'}{t'} P(z)\right) \quad (2)$$

we can rewrite (1)

$$t \frac{\partial f(x, t)}{\partial t} = \frac{\alpha_s}{2\pi} \int \frac{dz}{z} P(z) f\left(\frac{x}{z}, t\right) + f(x, t) \frac{t}{\Delta_s(t)} \frac{\partial \Delta_s(t)}{\partial t}. \quad (3)$$

Sudakov formalism

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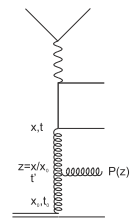
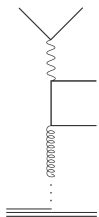
After integration

$$f(x,t) = f(x, t_0) \Delta_s(t) + \frac{\alpha_s}{2\pi} \int \frac{dt'}{t'} \frac{\Delta_s(t)}{\Delta_s(t')} \int \frac{dz}{z} P(z) f\left(\frac{x}{z}, t'\right). \quad (4)$$

Sudakov: probability of evolving from t_0 to t without any resolvable branching.

iterative solution:

$$f(x,t) = \lim_{n \rightarrow \infty} f_n(x,t) = \lim_{n \rightarrow \infty} \sum_n \frac{1}{n!} \log^n\left(\frac{t}{t_0}\right) A^n \otimes \Delta_s(t) f\left(\frac{x}{z}, t_0\right), \quad (5)$$

where $A = \frac{\alpha_s}{2\pi} \int \frac{dz}{z} P(z)$.

MC solution of the evolution equation

MC solution:

First branching: evolve from t_0 to t' obtained from $\Delta_s(t')$:

$$R_1 = \Delta_s(t'), \quad (6)$$

where R_1 is a random number in the interval $(0, 1)$.

If $t' > t$ evolution is stopped without any branching. If $t' < t$ branching is generated according to $P(z)$

$$\int_{z_{min}}^z dz' P(z') = R_2 \int_{z_{min}}^{z_{max}} dz' P(z') \quad (7)$$

and the evolution continues.

Second branching: evolve from t' to t'' generated according to $\Delta_s(t'', t')$. If $t'' > t$ evolution is stopped only with one branching. If $t'' < t$ branching is generated according to $P(z)$ and the evolution continues...etc.

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Observation:

$$\frac{\partial}{\partial t'} \frac{\Delta_s(t)}{\Delta_s(t')} = \frac{\alpha_s}{2\pi} \frac{\Delta_s(t)}{\Delta_s(t')} \frac{1}{t'} \int_x^{z_{max}} dz P(z) \quad (8)$$

rewrite (4)

$$f(x, t) = f_0(x, t) \Delta_s(t) + \int_x^1 \frac{dz'}{z'} \int_{t_0}^t d\Delta_s(t, t') P(z') f_0\left(\frac{x}{z'}, t'\right) \left(\int_x^{z_{max}} dz P(z)\right)^{-1} \quad (9)$$

Momentum weighted parton densities & momentum sum rule

To include all flavours in the evolution & use the Sudakov formalism we need to switch from $f(x, t)$ to $xf(x, t)$ & use momentum sum rule:

$$\sum_a \int_0^1 z P_{ab}(\alpha_s, z) dz = 0. \quad (10)$$

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example for gluon:

P - regularized splitting function (with *plus prescription*), \hat{P} - unregularized splitting function (without *plus prescription*)

$$t \frac{\partial xg(x, t)}{\partial t} = \frac{\alpha_s}{2\pi} \int_x^1 \frac{dz}{z} \left[z P_{gq}(z) \frac{x}{z} q\left(\frac{x}{z}, t\right) + z P_{gg}(z) \frac{x}{z} g\left(\frac{x}{z}, t\right) \right] \quad (11)$$

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Momentum weighted parton densities & momentum sum rule

To include all flavours in the evolution & use the Sudakov formalism we need to switch from $f(x, t)$ to $xf(x, t)$ & use momentum sum rule:

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If we define now Sudakov form factor as:

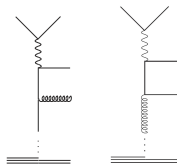
$$\Delta_s(t, t_0)_i \equiv \Delta_s(t)_i = \exp \left(- \int_x^{zmax} dz \int_{t_0}^t \frac{\alpha_s}{2\pi} \frac{dt'}{t'} z \sum_j \hat{P}(z)_{ji} \right) \quad (12)$$

where i, j are q or g ,

we obtain equation of the same form as eq. (3), just multiplied from both sides by x and with $zP(z)$ in $\Delta_s(t)_i \rightarrow$ it can be solved in analogical way.

Evolution in the code

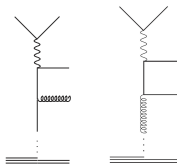
We consider ep collisions in which we can measure different pdfs:



Forward evolution: final parton is not specified when the evolution begins.

Evolution in the code

We consider ep collisions in which we can measure different pdfs:



Forward evolution: final parton is not specified when the evolution begins.

Four different situations:

- ▶ gluon at the beginning and at the end,
- ▶ quark (valence or sea) at the beginning and gluon at the end
- ▶ quark (valence or sea) at the beginning and quark (valence or sea) at the end and
- ▶ gluon at the beginning and sea quark at the end.

Valence quark at the end can come only from valence quark at the beginning.

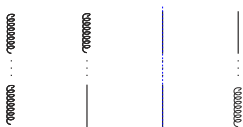
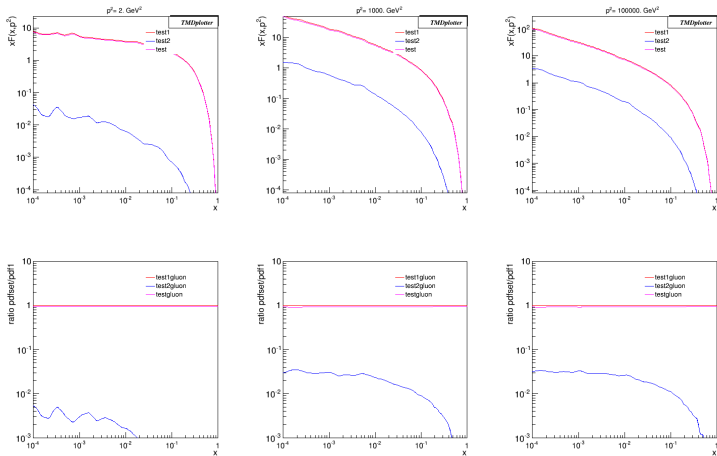


Figure : K_{gg} , K_{gq} , K_{qq} , K_{qg}

Contribution from quark and gluon evolution

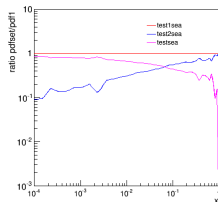
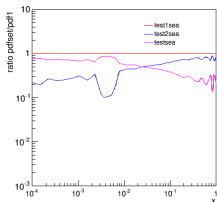
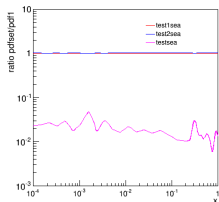
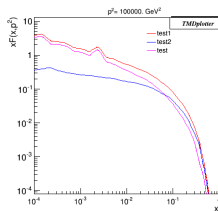
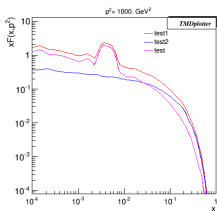
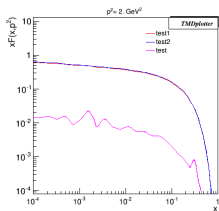
GLUON density: test1: merged quark and gluon evolution chain test2: quark evolution chain test: gluon evolution chain



Contribution from gluon kernel dominates, contribution from quark kernel (10^{-2}) times smaller.

Contribution from quark and gluon evolution

SEA density: test1: merged quark and gluon evolution chain test2: quark evolution chain test: gluon evolution chain

At small x main contribution to sea quark density from gluon kernel, for large x quark kernel dominates.

Integrated PDFs from MC solution

Ordering dependence- z_{\max} origin

Some of the splitting functions are divergent for $z \rightarrow 1$.

To avoid divergences:

$$\begin{aligned} \frac{\partial x f(x, t)}{\partial t} &= \frac{\alpha_s}{2\pi} \frac{1}{t} \int_x^1 dz P(z) \frac{x}{z} f\left(\frac{x}{z}, t\right) - \frac{\alpha_s}{2\pi} x f(x, t) \frac{1}{t} \int_x^1 dz P(z) \approx \\ &\approx \frac{\alpha_s}{2\pi} \frac{1}{t} \int_x^{z_{\max}} dz P(z) \frac{x}{z} f\left(\frac{x}{z}, t\right) - \frac{\alpha_s}{2\pi} x f(x, t) \frac{1}{t} \int_x^{z_{\max}} dz P(z). \end{aligned} \quad (13)$$

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it can be shown that terms $\int_{z_{max}}^1$ skipped in the integral in eq. (13) are of order $\mathcal{O}(1 - z_{max})$ multiplied by $x f(x, t)$ or $x \frac{df(x, t)}{dt}$

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Different choices of z_{max} :

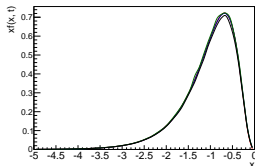
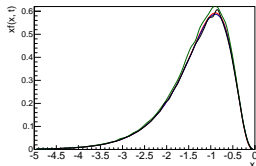
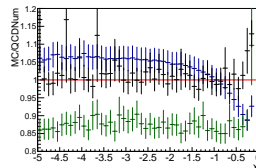
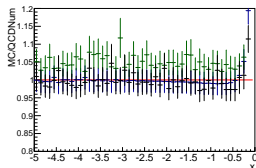
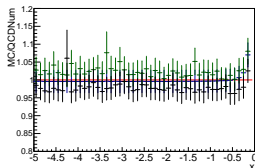
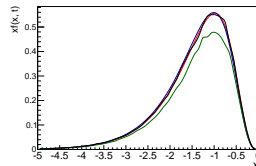
- ▶ z_{max} - fixed
- ▶ z_{max} - can change dynamically with the scale, for example:
 - angular ordering: $z_{max} = \left(\frac{Q_0}{Q}\right)^2$

up-val quarks

2GeV^2 : QCDNum, angular ordering, $z_{\text{max}} = 0.99$ $z_{\text{max}} = 0.97$

1000GeV^2 : QCDNum, angular ordering, $z_{\text{max}} = 0.999$ $z_{\text{max}} = 0.997$

100000GeV^2 : QCDNum, angular ordering, $z_{\text{max}} = 0.999$ $z_{\text{max}} = 0.9999$

up-val at scale 2 GeV^2 up-val at scale 1000 GeV^2 up-val at scale 100000 GeV^2 

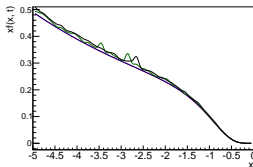
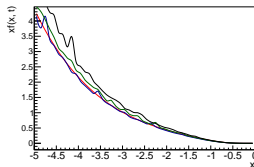
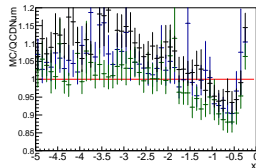
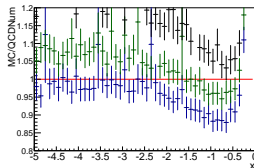
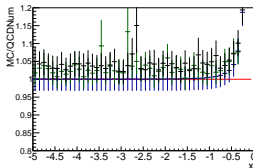
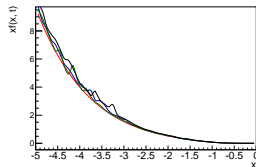
MC results close to the QCDNum results. Effect on z_{max} observed. z_{max} value giving the results the closest to QCDNum depends on scale and shape of the distribution (terms of order $\mathcal{O}(1 - z_{\text{max}})$ multiplied by $xf(x, t)$ or $x \frac{df(x, t)}{dt}$ skipped).

sea quarks

2GeV^2 : QCDNum, angular ordering, $z_{max} = 0.99$ $z_{max} = 0.995$

1000GeV^2 : QCDNum, $z_{max} = 0.97$, $z_{max} = 0.98$ $z_{max} = 0.99$

100000GeV^2 : QCDNum, $z_{max} = 0.996$, $z_{max} = 0.995$ $z_{max} = 0.997$

sea at scale 2 GeV²sea at scale 1000 GeV²sea at scale 100000 GeV²

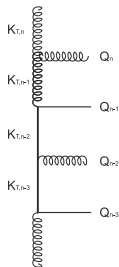
MC results close to the QCDNum results.

Effect on z_{max} observed. z_{max} value giving the results the closest to QCDNum depends on scale and shape of the distribution.

Results for TMDs

k_T dependence

MC solution: for every branching Q is generated and Q_x and Q_y are calculated.



k_T contains the whole history of the evolution:

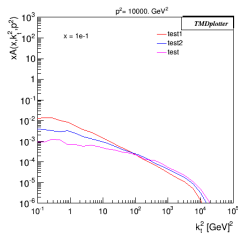
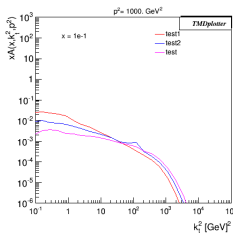
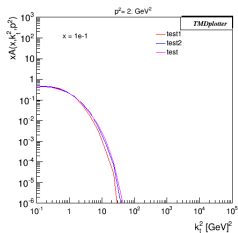
$$\vec{k}_{T,n} = \vec{k}_{T,n-1} + \vec{Q}_{T,n-1}.$$

The information about k_T is available for every branching.

up-val

$$z_{max} = 0.99, z_{max} = 0.999, z_{max} = 0.9999$$

$$x = 0.1$$

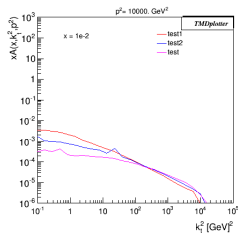
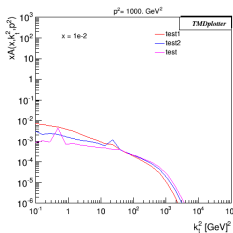
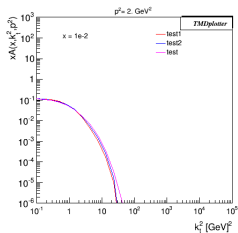


- ▶ Initial scale: intrinsic k_T distribution.
- ▶ Different choices of z_{max} lead to different uTMDs, especially different large k_T tails.

up-val

$$z_{max} = 0.99, z_{max} = 0.999, z_{max} = 0.9999$$

$$x = 0.01$$

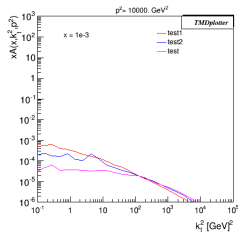
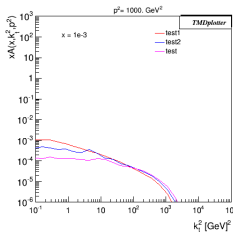
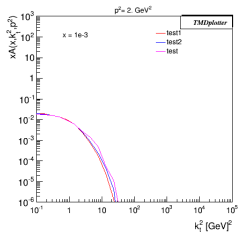


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up-val

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$$x = 0.001$$



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Fitting method

Evolution kernels

Two different evolution kernels are defined:

- ▶ initial quark (valence or sea) → quark grid,
- ▶ initial gluon → gluon grid.

Kernels for evolution initiated by gluons and quarks are calculated separately only once per run of the code and combined at the end → fitting procedure is fast.



Figure : $K_{gg}, K_{gq}, K_{qq}, K_{qg}$: gluon grid, quark grid

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Figure : $K_{gg}, K_{gq}, K_{qq}, K_{qg}$: gluon grid, quark grid

To get the final pdf: evolution kernel is folded with starting distribution

$$xf(x, t) = x \int dx_0 \int dz f_0(x_0) K(z, t) \delta(zx_0 - x) \quad (14)$$

For example for gluon:

$$xf(x, t)_g = x \int dx_0 \int dz (f_{0g}(x_0) K_{gg} + f_{0q}(x_0) K_{gq}) \delta(zx_0 - x), \quad (15)$$

and for up-valence quark:

$$xf(x, t)_{q, up} = x \int dx_0 \int dz (f_{0, up}(x_0) K_{qq}) \delta(zx_0 - x). \quad (16)$$

Summary

Summary

Solution of DGLAP evolution using MC method was shown.

Advantages:

- ▶ reproduce analytical solution (results consistent with QCDNum17),
- ▶ applicable for TMDs evolution,
- ▶ direct usage in PS matched calculation.

Prospects:

- ▶ implementation into xFitter: first full TMD distributions,
- ▶ include NLO in $P(z)$,
- ▶ development of full TMD MC - CASCADE.

Thank you!

Back up

evolution in the code

Different kind of splittings can happen during the evolution process:



Commonly used flavour decomposition:

▶ see

$$u + \bar{u} + d + \bar{d} + s + \bar{s} = 2(\bar{u} + \bar{d} + \bar{s}) + u_v + d_v$$



▶ valence

$$\frac{\bar{u} - u}{\bar{d} - d}$$



▶ gluon



Flavour decomposition used in the method:

▶ see

$$2(\bar{u} + \bar{d} + \bar{s})$$



▶ valence

$$\frac{\bar{u} - u}{\bar{d} - d}$$



▶ gluon



Ordering dependence- z_{max} origin

Some of the splitting functions are divergent for $z \rightarrow 1$.

To avoid divergences:

$$\begin{aligned} \frac{\partial xf(x, t)}{\partial t} &= \frac{\alpha_s}{2\pi} \frac{1}{t} \int_x^1 dz P(z) \frac{x}{z} f\left(\frac{x}{z}, t\right) - \frac{\alpha_s}{2\pi} xf(x, t) \frac{1}{t} \int_x^1 dz P(z) \approx \\ &\approx \frac{\alpha_s}{2\pi} \frac{1}{t} \int_x^{z_{max}} dz P(z) \frac{x}{z} f\left(\frac{x}{z}, t\right) - \frac{\alpha_s}{2\pi} xf(x, t) \frac{1}{t} \int_x^{z_{max}} dz P(z). \end{aligned} \quad (17)$$

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Using the form of the splitting functions:

$$P_{ab}(\alpha_s, z) = A_{ab}(\alpha_s)\delta(1-z) + K_{ab}(\alpha_s)\frac{1}{(1-z)_+} + R_{ab}(\alpha_s) \quad (18)$$

and the expansion of $xf(\frac{x}{z}, t)$

$$xf\left(\frac{x}{z}, t\right) = xf(x, t) + x^2 f'(x, t)(1-z) + \mathcal{O}(1-z_{max}) \quad (19)$$

it can be shown that terms $\int_x^{z_{max}}$ skipped in the integral in eq. (17) are of order $\mathcal{O}(1-z_{max})$ multiplied by $xf(x, t)$ or $x \frac{df(x', t)}{dt}$

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example:

$$\begin{aligned} \frac{\alpha_s}{2\pi} \frac{1}{t} \int_{z_{max}}^1 dz K_{ab}(\alpha_s) \frac{1}{1-z} \frac{x}{z} f\left(\frac{x}{z}, t\right) - \frac{\alpha_s}{2\pi} xf(x, t) \frac{1}{t} \int_{z_{max}}^1 dz K_{ab}(\alpha_s) \frac{1}{1-z} = \\ = \frac{\alpha_s}{2\pi} \frac{1}{t} K_{ab}(\alpha_s) \left[x^2 f'(x, t)(1-z_{max}) + \mathcal{O}(1-z_{max}) \right] \end{aligned} \quad (20)$$

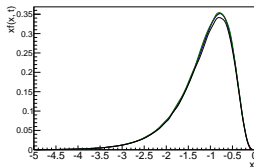
down-val quarks

2GeV^2 : QCDNum, angular ordering, $z_{max} = 0.99$ $z_{max} = 0.97$

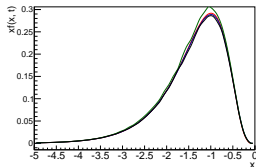
1000GeV^2 : QCDNum, angular ordering, $z_{max} = 0.999$ $z_{max} = 0.997$

100000GeV^2 : QCDNum, angular ordering, $z_{max} = 0.999$ $z_{max} = 0.9999$

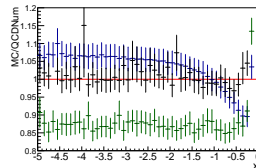
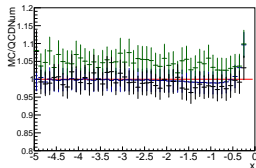
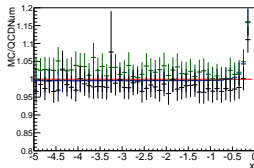
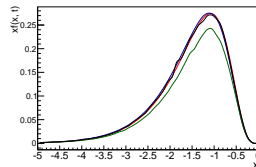
down-val at scale 2GeV^2



down-val at scale 1000GeV^2



down-val at scale 100000GeV^2



MC results close to the QCDNum results.

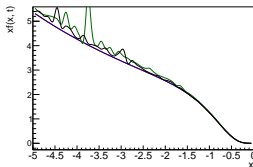
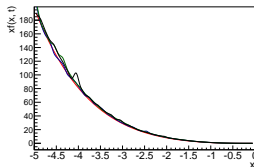
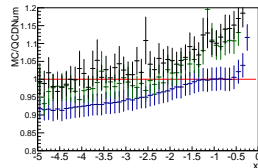
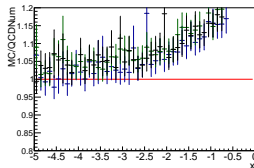
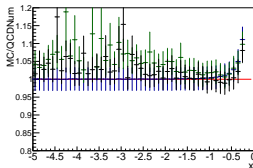
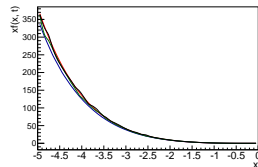
Effect on z_{max} observed. z_{max} value giving the results the closest to QCDNum depends on scale and shape of the distribution.

gluon

2GeV²:QCDNum, angular ordering, $z_{max} = 0.9$ $z_{max} = 0.87$

1000GeV²:QCDNum, $z_{max} = 0.99$, $z_{max} = 0.995$ $z_{max} = 0.993$

100000GeV²:QCDNum, angular ordering, $z_{max} = 0.995$ $z_{max} = 0.997$

gluon at scale 2 GeV²gluon at scale 1000 GeV²gluon at scale 100000 GeV²

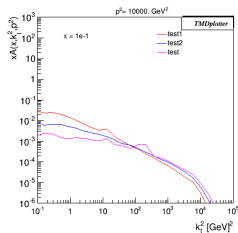
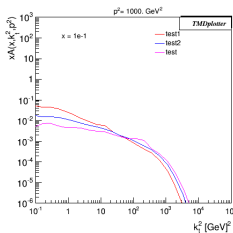
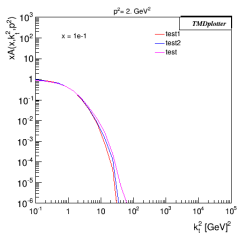
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down-val

$$z_{max} = 0.99, z_{max} = 0.999, z_{max} = 0.9999$$

$$x = 0.1$$

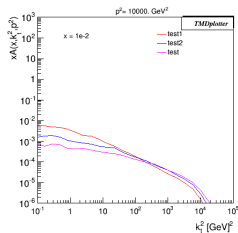
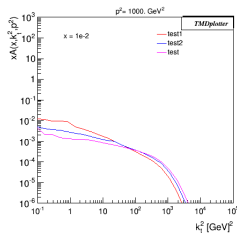
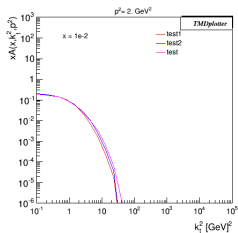


- ▶ Initial scale: intrinsic k_t distribution.
- ▶ Different choices of z_{max} lead to different uTMDs, especially different large k_T tails.

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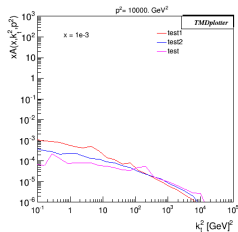
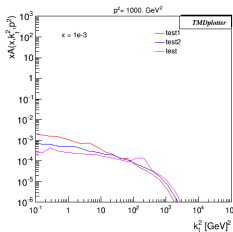
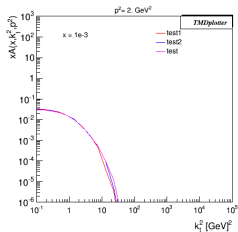


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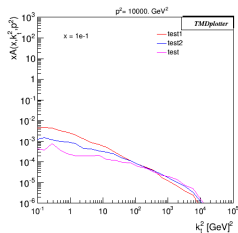
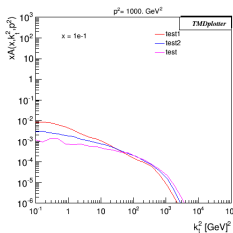
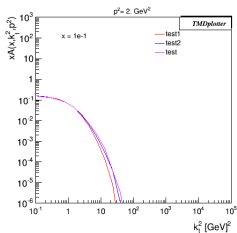


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sea

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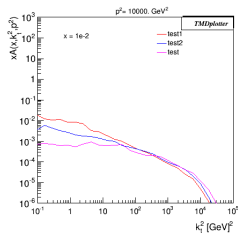
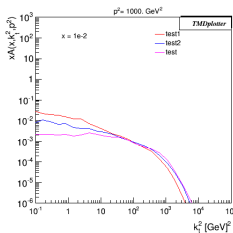
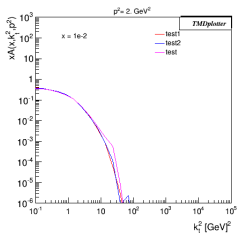


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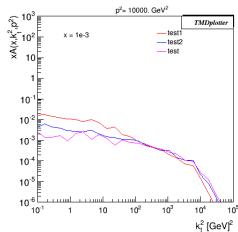
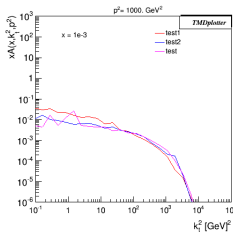
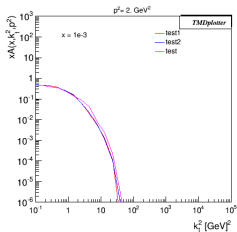


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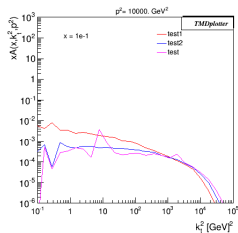
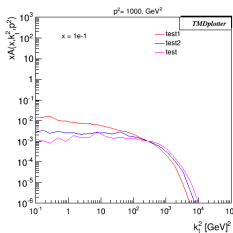
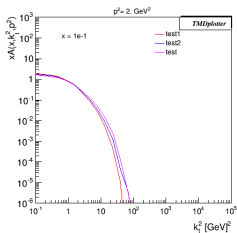


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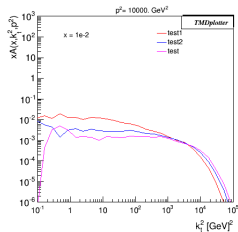
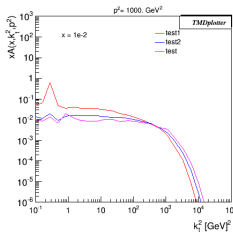
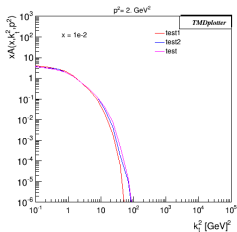


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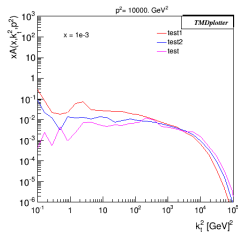
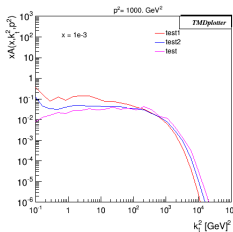
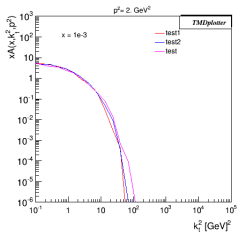


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