# Mass-mass plots for s-channel mediators.

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based on **arXiv:1503.05916** and **arXiv:1510.02110** in collaboration with Mikael Chala, Matthew McCullough, Germano Nardini, Kai Schmidt-Hoberg, Thomas Schwetz and Stefan Vogl





#### **Mass-mass plots**



For concreteness, I will focus on a spin-1 mediator and Dirac fermion DM:

$$\mathcal{L} = -\sum_{f=q,l,\nu} Z^{\prime\mu} \,\bar{f} \left[ g_f^V \gamma_\mu + g_f^A \gamma_\mu \gamma^5 \right] f - Z^{\prime\mu} \,\bar{\psi} \left[ g_{\rm DM}^V \gamma_\mu + g_{\rm DM}^A \gamma_\mu \gamma^5 \right] \psi$$



In the EFT limit, cross sections just depend on the effective suppression scale

$$M^* = m_{\rm med} / (g_{\rm DM} g_{\rm q})^{1/2}$$

and on the DM mass  $m_{\rm DM}$ .

However, for perturbative couplings such large mediator masses are very difficult to reach with monojet searches.

Moreover, it is typically very difficult to reproduce the observed DM relic abundance from thermal freeze-out for very heavy mediators.



# The off-shell region



## The on-shell region



In the on-shell region things become more complicated. However, in many cases we can use the narrow-width approximation (NWA):

$$\sigma(q\bar{q} \to Z' + Y \to xy + Y)$$
  
=  $\sigma(q\bar{q} \to Z' + Y) \cdot BR(Z' \to xy)$ 

In this approximation the monojet cross section is proportional to the invisible branching ratio of the mediator.

 $\Gamma_R/M_R = 0.4$ 

 $\Gamma_R/M_R = 0.2$ 

200



#### **Advanced rescaling attempts**

One could try to use different rescaling rules in the on-shell region and in the off-shell region.

 $\sigma \propto \begin{cases} g_q^2 g_{\rm DM}^2 / \Gamma_{\rm OS} \text{ if } M > 2m_{\rm DM} \\ g_q^2 g_{\rm DM}^2 & \text{if } M < 2m_{\rm DM} \end{cases}$ 





#### How to present results without assumed rescaling

Option 1: Fix both couplings and quote an upper bound on the signal strength.

Option 2: Fix one coupling (or the ratio of the couplings) and quote an upper bound on the other coupling (or the product of the couplings).





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- + The mediator width is fixed for each value of  $m_{\rm DM}$  and  $m_{\rm med}$ , so only two parameters need to be scanned for the signal generation.
- + It is immediately clear how to include other constraints (dijets, direct detection...) in these plots.
- Requires a (somewhat arbitrary) choice of couplings. Translation to different couplings is non-trivial.

Option 2: Fix one coupling (or the ratio of the couplings) and quote an upper bound on the other coupling (or the product of the couplings).

- + Results can be understood intuitively.
- The width (and therefore kinematic distributions) depend on the coupling, so signal generation becomes more complicated.
- Need to make sure to avoid nonperturbative couplings and large mediator width.
- Difficult to show other constraints in the same parameter plot.



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What are good choices for  $g_{\rm DM}$  and  $g_{\rm q}$ ?

DM Forum recommendation (for spin-1 mediators):  $g_{\rm DM} = 1$ ,  $g_{\rm q} = 0.25$ 

- Reasonable mediator width:  $\Gamma_{med} / m_{med} \sim 0.2$
- $BR(Z' \rightarrow qq) \sim BR(Z' \rightarrow invisible) \sim 50\%$
- No overwhelming constraints (but potential sensitivity) from dijet searches



#### What other constraints should be shown?





# **Relic density**

 $\overline{q}$ 



- It is a useful guideline to indicate the parameters where interactions between DM and quarks alone are sufficient to explain the observed DM abundance.
- Should be thought of more as a model prediction (for one specific model) rather than a model-independent bound.





# Unitarity

 $\overline{q}$ 



- > Moreover, for axial couplings the longitudinal component of the mediator couples to fermions with a coupling strength that is given by  $2 g_f^A m_f / m_{Z'}$ .
- > Perturbative unitarity thus bounds the DM mass:  $m_f \lesssim$

FK et al., arXiv:1510.02110



DES

## **Direct detection**

 $\overline{q}$ 



- To show direct detection bounds, two important assumptions are necessary:
  - Relative sign between up-quark coupling and down-quark coupling.
  - Rescaling of direct detection bounds for DM sub-components.



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# **Dijets**







See talk by Matthew McCullough (tomorrow)



#### **Combined constraints**



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#### Conclusions

- Mass-mass plots are a useful way to present results from mono-X searches at the LHC and provide an interpretation in terms of simplified models.
- Since there is no simple scaling rule relating different couplings, it will be necessary to fix one (or both) coupling(s) in order to produce such a plot.
- Fixing both couplings has the advantage that one can compare the results to other bounds (relic density, direct detection, ...).
- Such a comparison always requires additional assumptions, so one should either specify the assumptions clearly or produce several bounds for different assumptions.
- In the end it may be necessary to produce plots for different coupling combinations in order to capture the full phenomenology of the simplified model (and the complementarity of different search stategies).

