



Summary of DM searches in CMS

LHC DM Working Group 10/12/2015

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Dark Matter Madness

- Strong evidence of physics beyond the Standard Model
- We know very little about the nature of DM (gravitational interaction but what else?)
- If DM talks to SM particles (eg WIMPs) we can detect it via...





Increasing interest in dark matter (DM) searches at the LHC

- Complementary strategies to direct detection and spectral data
- High hopes to find evidence of DM production for Run-2!

Run-1 a huge success for CMS and the LHC!



Mono-Mania (MET+X)

DM particles produced in p-p collisions escape the detector without interacting.

Search for an abundance of events with an imbalance of energy in the transverse plane

Look for additional particles recoiling against DM particles.

Provides a rich assortment of $X + E_T^{miss}$ searches at the LHC...



*NB// also many SUSY searches cover DM but not covered here

Monojet

High-momentum jet, from initial state radiation (ISR), recoils against DM

- ♦ 1 central jet with high p_T > 150 GeV, |η|<2 with Δφ(j,E_T^{miss}) >2
- ♦ 2nd jet allowed provided if $\Delta \phi(j_1, j_2) < 2$
- ♦ Veto leptons photons with $p_T > 10$ GeV and taus with $p_T > 15$ GeV

Signal extraction from fit using multiple bins in $\mathsf{E}_{\mathsf{T}}^{\mathsf{miss}}$

Dominant backgrounds from Z(vv)+jets and W(lv)+jets

 \rightarrow Estimate the contributions with data using Z \rightarrow II +jet, γ +jet and W \rightarrow Iv +jet events





Mono-V (leptonic)

Associated production of DM with a W or Z (V) boson Smaller production cross-section but also lower backgrounds than jets final state.



W(Iv) + E_T^{miss}

- ♦ 1 isolated, high- p_T lepton
- ♦ Backgrounds from W/Z production decays with lepton out of acceptance.
- ♦ Discrimination from transverse mass variable



$Z(II) + E_T^{miss}$

- ♦ Look for 2 opposite-charge, same-flavor leptons
- ♦ Invariant mass consistent with Z boson $|m_{II}-m_{Z}| < 10 \text{ GeV}$
- ♦ Look for excess in m_T spectrum



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Mono-V (leptonic)

Interpret in terms of lower limit on new physics scale (Λ) as a function of **DM mass**





- ♦ Increased sensitivity to models with enhanced couplings to vector bosons
- Need to consider regions of Validity for EFT approach (eg via truncation)

Mono-V (Hadronic)

Low $p_T W$ or Z bosons decay to well separated jets (fully reconstructed)

Look at mass and pT of the dijet system to distinguish backgrounds.

Exploit jet/di-jet properties:

- ♦ Quark-gluon likelihood discriminator
- \diamond Jet-pull (color flow between jets) [1]
- ♦ "Mass-drop" [2]

Combine into multivariate discriminator



Use photon CR with photon as 'fake' E_t^{miss} to dramatically reduce statistical uncertainty on **Z(vv)+jets** backgrounds





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[2] PRL.114.041802

[1] arXiv:1001.5027

Mono-V (Hadronic)

Boosted (high p_T) vector bosons decaying to jets will form a single "fat"-jet

Jet substructure techniques to identify V-bosons:

- $\diamond~$ Look for high-p_T fat jet with $\boldsymbol{m_J}$ close to $\boldsymbol{m_W}$ or $\boldsymbol{m_Z}$
- \diamond N-subjettiness (τ_N) (likelihood for N-daughter hypotheses of lead jets to tag as having originated from a W or Z boson.
- Use (di)-lepton and photon CR to determine W(I) and Z(vv) bkgs







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Mono-photon

Typically smaller cross-sections but high- E_{τ} photon easy to trigger

 \rightarrow Lower E_T^{miss} thresholds accessible

- \diamond One high energy, isolated photon, E_T>145
- $\Leftrightarrow E_T^{miss} > 140 \text{ GeV}$
- ↔ Reject events with overlapping photon and E_T^{miss} (Δφ(γ, E_T^{miss})<2





- Largest backgrounds from Wγ/Zγ estimated from lepton CR
- γ+jet background from MC corrected using data-MC scale-factors



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Heavy Flavor

Scalar interactions with DM favor heavy-flavor quarks

- \diamond Top-quark and bottom-quark coupling enhanced ($\sim m_q$)
- Searches for top-quark pairs recoiling against the DM particles

\rightarrow tt+E_T^{miss}

- \diamond One lepton + at least 1 b-tagged jet, $E_T^{miss} > 320$
- $\diamond~m_{\tau}\! >\! 160$ GeV to remove W and tt backgrounds
- \diamond Normalize tt background from data CR







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Heavy Flavor

Several BSM models predict single top + E_T^{miss}

- → Consider resonant and non-resonant production modes
- Hadronic final states used for mono-top analysis, veto events with isolated leptons
- Cut on invariant mass of 3 jet (top) system < 250 GeV</p>





Di-jet/bb + E_Tmiss

Additional jets produced in association with DM particles

- \diamond Additional information from second jet allows for good discrimination against SM backgrounds (W/Z/tt) 18.8 fb⁻¹ (8 TeV)
- ♦ Fit "Razor-variable"

$$\begin{split} M_R &\equiv \sqrt{(|\vec{p}_{J_1}| + |\vec{p}_{J_2}|)^2 - (p_z^{J_1} + p_z^{J_2})^2} \\ M_T^R &\equiv \sqrt{\frac{E_T^{miss}(p_T^{J_1} + p_T^{J_2}) - \vec{E}_T^{miss} \cdot (\vec{p}_T^{J_1} + \vec{p}_T^{J_2})}{2}} \\ & \text{In 4x bins of } M_R \end{split}$$



- bb and b tagged bins target scenario where DM preferentially couples to b-quarks [1]
- \Rightarrow Backgrounds estimated from $\mu\mu(b)+jj$ or $\mu\mu(b)+tt$ control regions in data



[1] Phys.Lett.B697:412-428,2011

Moving on from EFT(LHC DM forum)

EFT interpretation useful as a benchmark for sensitivity and comparison to DD, However, validity break-down where LHC can reach mediator mass-scale



$H \rightarrow$ Invisibles (concrete scalar model)

If DM is massive then $DM \leftarrow \rightarrow SM$ mediated via the **Higgs boson** (eg Higgs portal Models*)





Gluon-fusion

Dominant production but low acceptance after kinematic selection: >0 jets, large E_T^{miss} as in monojet search



- Look for pair of charged leptons / b-quarks consistent with Z decay
- \diamond V(had)+E_T^{miss} also targets this mode

Smaller x-section (~10x less than ggF) but Large S/B ratio (~70%!)

- Search for two jets with large η-separation and large invariant mass
- \diamond Robust counting analyses in E_T^{miss} / m_{ii} tails



$H \rightarrow Invisibles$



$H \rightarrow Invisibles$

Combine several searches and interpret as limits on decay



Combination of monojet + mono-V (hadronic) searches interpreted under simplified model



Scalar and pseudo-scalar case * assume Yukawa-couplings to SM $g_{SM} = g_q m_q / v_c$

Fermion only scenario assumes no couplings to vector bosons (results in greatly reduced sensitivity for spin-0 mediators)



^{*}mono-V topology signal generated with JHUgen

Simplified Model interpretation (CMS-PAS-EXO-12-055)

Comparison to direct-detection in terms of DM-nucleon scattering cross-section vs DM mass planes (comparison on the direct-detections own turf)

 \rightarrow Re-interpretation of **exclusion contours** in m_{DM}-m_{MED} plane [1,2]



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Spin-dependent case

$$\sigma_{SD}^0 = \frac{3g_{DM}^2 g_{SM}^2 (\Delta_u + \Delta_d + \Delta_s) \mu_{n\chi}^2(m_{DM})}{\pi m_{med}^4}$$



 $\Delta_u = -0.42, \, \Delta_d = 0.85 \,\, {
m and} \,\, \Delta_s = -0.08 \,$ (Jhep 07 (2012) 009)

Spin-0 scenario interpretations assume only couplings to heavy quarks

→ limits the sensitivity in direct detection but avoids additional assumptions on light-quark couplings

DM-nucleon scattering velocity suppressed in pseudoscalar scenario

→ Compare sensitivity in terms of DM-annihilation cross-section (CMS limit assumes only bb mode is relevant)



Plans for Run 2 (LHC DM forum)

arXiv.org > hep-ex > arXiv:1507.00966

High Energy Physics – Experiment

Dark Matter Benchmark Models for Early LHC Run-2 Searches: Report of the ATLAS/CMS Dark Matter Forum

Daniel Abercrombie, Nural Akchurin, Ece Akilli, Juan Alcaraz Maestre, Brandon Allen, Barbara Alvarez Gonzalez, Jeremy Andrea, Alexandre Arbey, Georges Azuelos, Patrizia Azzi, Mihailo Backović, Yang Bai, Swagato Banerjee, James Beacham, Alexander Belyaev, Antonio Boveia, Amelia Jean Brennan, Oliver Buchmueller, Matthew R. Buckley, Giorgio Busoni, Michael Buttignol, Giacomo Cacciapaglia, Regina Caputo, Linda Carpenter, Nuno Filipe Castro, Guillelmo Gomez Ceballos, Yangyang Cheng, John Paul Chou, Arely Cortes Gonzalez, Chris Cowden, Francesco D'Eramo, Annapaola De Cosa, Michele De Gruttola, Albert De Roeck, Andrea De Simone, Aldo Deandrea, Zeynep Demiragli, Anthony DiFranzo, Caterina Doglioni, Tristan du Pree, Robin Erbacher, Johannes Erdmann, Cora Fischer, Henning Flaecher, Patrick J. Fox, et al. (94 additional authors not shown)

(Submitted on 3 Jul 2015)

This document is the final report of the ATLAS-CMS Dark Matter Forum, a forum organized by the ATLAS and CMS collaborations with the participation of experts on theories of Dark Matter, to select a minimal basis set of dark matter simplified models that should support the design of the early LHC Run-2 searches. A prioritized, compact set of benchmark models is proposed, accompanied by studies of the parameter space of these models and a repository of generator implementations. This report also addresses how to apply the Effective Field Theory formalism for collider searches and present the results of such interpretations.

Study of production cross-sections and kinematics of mediator models with different final states

ATLAS + CMS + Theorists collaborate to produce agreed set of minimal benchmark simplified models to be considered in early Run-2 DM searches

Search or

Plans for Run 2

LHC DM Forum recommendation of common grid of model points as benchmark for all DM searches (all final states) to cover variation in kinematics



Need a finer grained sampling the parameter space with additional points to produce smooth contours.

 \rightarrow The use of shape-based analyses, multiple observable requires a careful consideration of the appropriate distributions to interpolate

Example procedure

High statistics generator level samples (eg MG+PS) used to fill in the gaps.

 \rightarrow + Fast sim of the detector?

Use key analysis variables at generator level to reweight* full RECO samples at intermediate points. $\rightarrow EG E_T^{miss} vs \Delta \varphi(Z, E_T^{miss})$ for mono-Z

Repeat for each mediator and for different coupling scenarios

ightarrow Different effects on kinematics for different analyses





Summary

Wide variety of searches for DM at CMS in Run-1

- \diamond Mono-mania (X+E_T^{miss}) searches are simple yet powerful tools to look for DM production
- \diamond Interpretation of the data under EFT models \rightarrow useful for comparing sensitivities between searches and direct detection
- \Rightarrow First use of simplified models to interpret DM searches in Run-1

[GeV] 4000 90% CL Observe 90% CL Expected 3000 $\Lambda < m/2$ 2500 2000 CMS Preliminary µ + E^{mis} L dt = 20 fb⁻¹ 1500 - 300 GeV 👌 – 200 G 1000 500 10⁶ Events 10 WIMP mass m, [GeV] 10 10-4 10⁻³ 10-4 10-5 500 1000 1500 2000 2500 M_T (GeV) 19.7 fb⁻¹ (8 TeV) 3) (S) 10 **CMS** Preliminarv Annihilation Cross Section (cm g__=g_=1 10⁻²⁶ Pseudoscalar 10⁻²⁷ onoie -taggec ermion Only 10⁻²⁹ miLAT Exclusion ermil AT Best Fit (m) miLAT Best Fit (gg) ermil AT Best Fit (bb) 10 10 10 m_ (GeV)

Run-2 is in progress!

- \diamond 2015 pp run complete (Results coming next week!)
- \diamond Benefit from increased cross-sections and harker kinematics at 13TeV
- \diamond More boost @ 13 TeV will require more use of substructure techniques
- \diamond DM is an exciting topic for Run-2 at the LHC, exciting times ahead!

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Thank you!





Let's hope we will find it in here!

BACKUP SLIDES

Effective operators describing DM-SM interaction (Effective Field Theory, EFT)

- Scalar/Vector (spin-independent, SI)
- Pseudo-scalar/Axial-vector (spin-dependent, SD)

Only 2 parameters

- Interaction scale $M^* = M_{med}/\sqrt{(g_X g_q)}$
- DM mass M_X

| Name | Initial state | Type | Operator |
|------|---------------|--------------|---|
| C1 | qq | scalar | $\frac{m_{\bar{q}}}{M_*^2}\chi^{\dagger}\chi\bar{q}q$ |
| C5 | gg | scalar | $\frac{1}{4M_*^2}\chi^\dagger\chi\alpha_{\rm s}(G^a_{\mu\nu})^2$ |
| D1 | qq | scalar | $\frac{m_q}{M_\star^3} \bar{\chi} \chi \bar{q} q$ |
| D5 | qq | vector | $\frac{1}{M_*^2} \bar{\chi} \gamma^\mu \chi \bar{q} \gamma_\mu q$ |
| D8 | qq | axial-vector | $\frac{1}{M_{\star}^2}\bar{\chi}\gamma^{\mu}\gamma^5\chi\bar{q}\gamma_{\mu}\gamma^5q$ |
| D9 | qq | tensor | $\frac{1}{M_{\star}^2} \bar{\chi} \sigma^{\mu\nu} \chi \bar{q} \sigma_{\mu\nu} q$ |
| D11 | gg | scalar | $\frac{1}{4M_\star^3}\bar\chi\chi\alpha_{\rm s}(G^a_{\mu\nu})^2$ |

Validity issue for high transferred momentum, Q_{tr}: Q_{tr}>M_{med}

 Truncation techniques: account for the fraction of events not satisfying this assumption

A. de Cosa (LHCP 2015)





CMS Monojet+mono-V (EXO-12-055)

Simultaneous multi binned likelihood fit to 3 Control Regions (photon, muon, dimuon) (example, monojet category)





CMS Razor

Categorise events based on MR variable





Table 2: Definition of the event categories based on the M_R value, the muon multiplicity, and the output of the CSV b-tagging algorithm. For all the samples, $R^2 > 0.5$ is required.

| 00 | 0 0 | 1 ' 1 |
|------------------------------|--------------------------------|---|
| Sample | b-tagging selection | M _R selection |
| | no CSV loose jet | $200 < M_R \le 300 \text{ GeV} (VL)$ $300 < M_R \le 400 \text{ GeV} (L)$ |
| 0μ , 1μ , and 2μ | | $400 < M_R \le 400 \text{ GeV}$ (L) $400 < M_R \le 600 \text{ GeV}$ (H) |
| | | $M_R > 600 \text{ GeV} (VH)$ |
| 0 <i>µ</i> bb | \geq 2 CSV tight jets | |
| 0µb | =1 CSV tight jets | |
| 1 <i>µ</i> b | > 1 CSV tight into | $M_R > 200 \; GeV$ |
| 2 <i>µ</i> b | $\geq 1 \cos \alpha$ ugin jets | |
| $Z(\mu\mu)b$ | \geq 1 CSV loose jets | |

Events with >2 reconstructed jets are included by forming 2 "megajets" (sum jet 4momenta) to calculate Razor variables B-tagging brings sensitivity in scalar/pseudoscalar models



V/A Comparison to Direct Detectors



$$\begin{split} \sigma^{SD}_{N-\chi} &= 0.33 \frac{\mu^2}{\pi \Lambda_{LL}^4}, \\ \sigma^{SI}_{N-\chi} &= 9 \frac{\mu^2}{\pi \Lambda^4}, \end{split}$$

$$\mu = rac{\mathrm{M}_{\chi}\mathrm{M}_{p}}{\mathrm{M}_{\chi} + \mathrm{M}_{p}}$$
 ,

Comparison of mono-V analyses



Improved sensitivity in models interference in couplings to up/ down quarks



not particularly favoured model since potentially violates gauge invariance*

*arxiv.org/abs/1503.07874

Cross-section scaling 13 TeV/ 8 TeV(monojet)



Invisible Standard Model

Compare the predictions with the data in the CRs \rightarrow Maximize likelihood to obtain corrected prediction for backgrounds in signal regions



Simultaneous likelihood used in several places instead of simple `count and scale approach'

Interpretations of searches

Even a successful search at the LHC won't tell us...

- \diamond What is DM (fermion (dirac/majorana?)
- \diamond It's even the same as what's seen in galaxies (~stable for detector != stable)?
- $\diamond\,$ Is any of these the right one



Invisible Standard Model

Common enemies shared by many of the searches for DM

Neutrinos escape detection to mimic a DM signal!

Most common are Z decays to neutrinos and leptonic W decays where the lepton falls outside of the acceptance or isn't reconstructed.



Relatively large cross-sections for these processes mean backgrounds are sizable compared to real signals.

Invisible Standard Model

Z→vv W→lv Signal Region ν Lost µ/e Extrapolate yields/shapes in well understood control regions CR) in data to constrain backgrounds in signal region (SR) **Data Control Regions** \rightarrow Transfer Factors (TF) taken from simulation/theory used for extrapolation: eg BR($Z \rightarrow vv$)/BR($Z \rightarrow \mu\mu$) µ/e μ/e μ/e In CRs, calculate "Fake" E_t^{miss} (recoil) $\rightarrow |\vec{E}_T^{miss} + \sum \vec{p}_T^{ll/\gamma}|$

Monojet

High-momentum jet, from initial state radiation (ISR), recoils against DM

- \diamond 1 central jet with high p_T > 110 GeV
- ♦ 2nd jet allowed provided if $\Delta \phi(j_1, j_2) < 2.5$
- \diamond Veto leptons, photons

Dominant backgrounds from Z(vv)+jets and W(lv)+jets

 \rightarrow Estimate the contributions with data using $Z \rightarrow II$ and $W \rightarrow Iv$ events

Define signal regions with increasing E_T^{miss} thresholds \rightarrow Model independent cross-section limits EPJC75 (2015) 2





Comparisons to Direct Detectors

Translate EFT limits into upper limits on spin-dependent and spin-independent χ -N cross-sections \rightarrow Complementary to DD searches



$$\begin{split} \mathcal{L}_{\text{scalar}} \supset &-\frac{1}{2} m_{\text{MED}}^2 S^2 - g_{\text{DM}} S \, \bar{\chi} \chi - \sum_{q=b,t} g_{SM}^q S \, \bar{q} q - m_{\text{DM}} \bar{\chi} \chi \,, \\ \mathcal{L}_{\text{pseudo-scalar}} \supset &-\frac{1}{2} m_{\text{MED}}^2 P^2 - i g_{\text{DM}} P \, \bar{\chi} \gamma^5 \chi - \sum_{q=b,t} i g_{SM}^q P \, \bar{q} \gamma^5 q - m_{\text{DM}} \bar{\chi} \chi \,, \\ \mathcal{L}_{\text{vector}} \supset &\frac{1}{2} m_{\text{MED}}^2 Z_{\mu}' Z'^{\mu} - g_{\text{DM}} Z_{\mu}' \bar{\chi} \gamma^{\mu} \chi - \sum_{q} g_{SM}^q Z_{\mu}' \bar{q} \gamma^{\mu} q - m_{\text{DM}} \bar{\chi} \chi \,, \\ \mathcal{L}_{\text{axial}} \supset &\frac{1}{2} m_{\text{MED}}^2 Z_{\mu}'' Z'^{\mu} - g_{\text{DM}} Z_{\mu}'' \bar{\chi} \gamma^{\mu} \gamma^5 \chi - \sum_{q} g_{SM}^q Z_{\mu}'' \bar{q} \gamma^{\mu} \gamma^5 q - m_{\text{DM}} \bar{\chi} \chi \,. \end{split}$$