

# Precision predictions for Higgs production at the LHC

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in collaboration with C. Anastasiou, F. Dulat, E. Furlan,  
T. Gehrmann, F. Herzog, A. Lazopoulos, B. Mistlberger

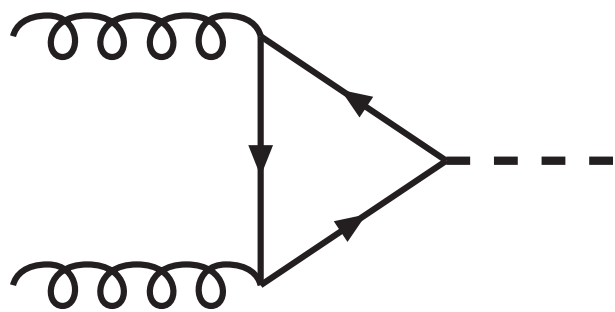
Zurich Phenomenology Workshop, 06/01/2016



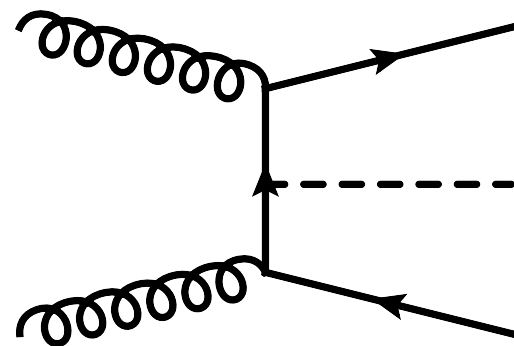
# Higgs production at the LHC



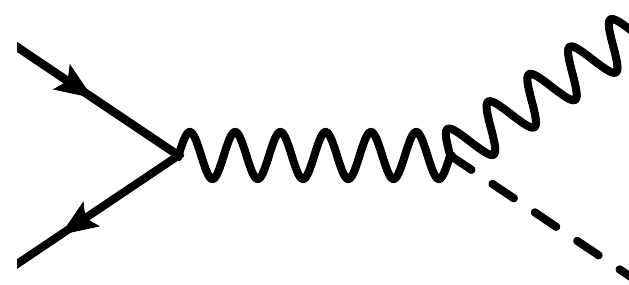
- Establishing whether the BEH mechanism and its boson is SM-like will be of utmost importance for the run of the LHC.
- Higgs-boson production modes at the LHC:



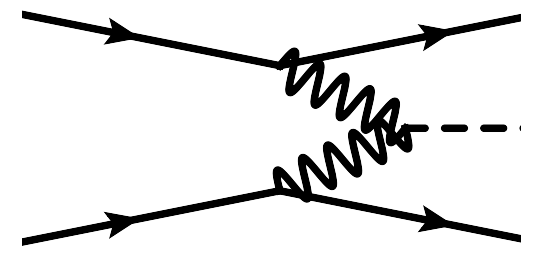
Gluon fusion



TTH



Higgs strahlung



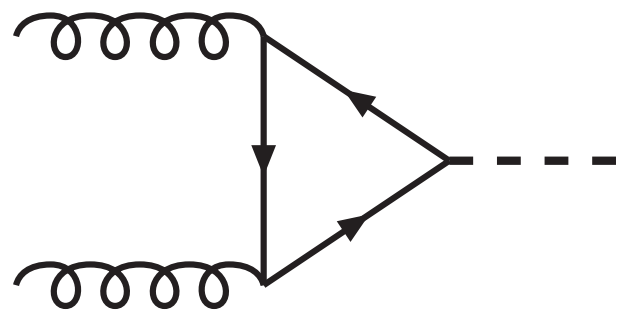
VBF



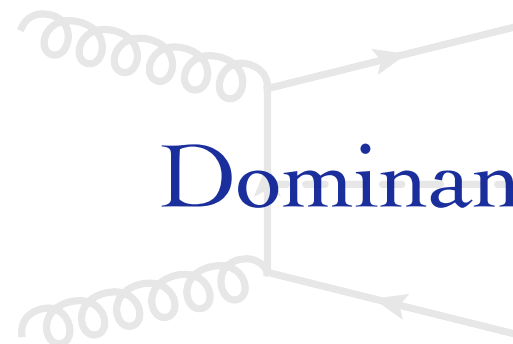
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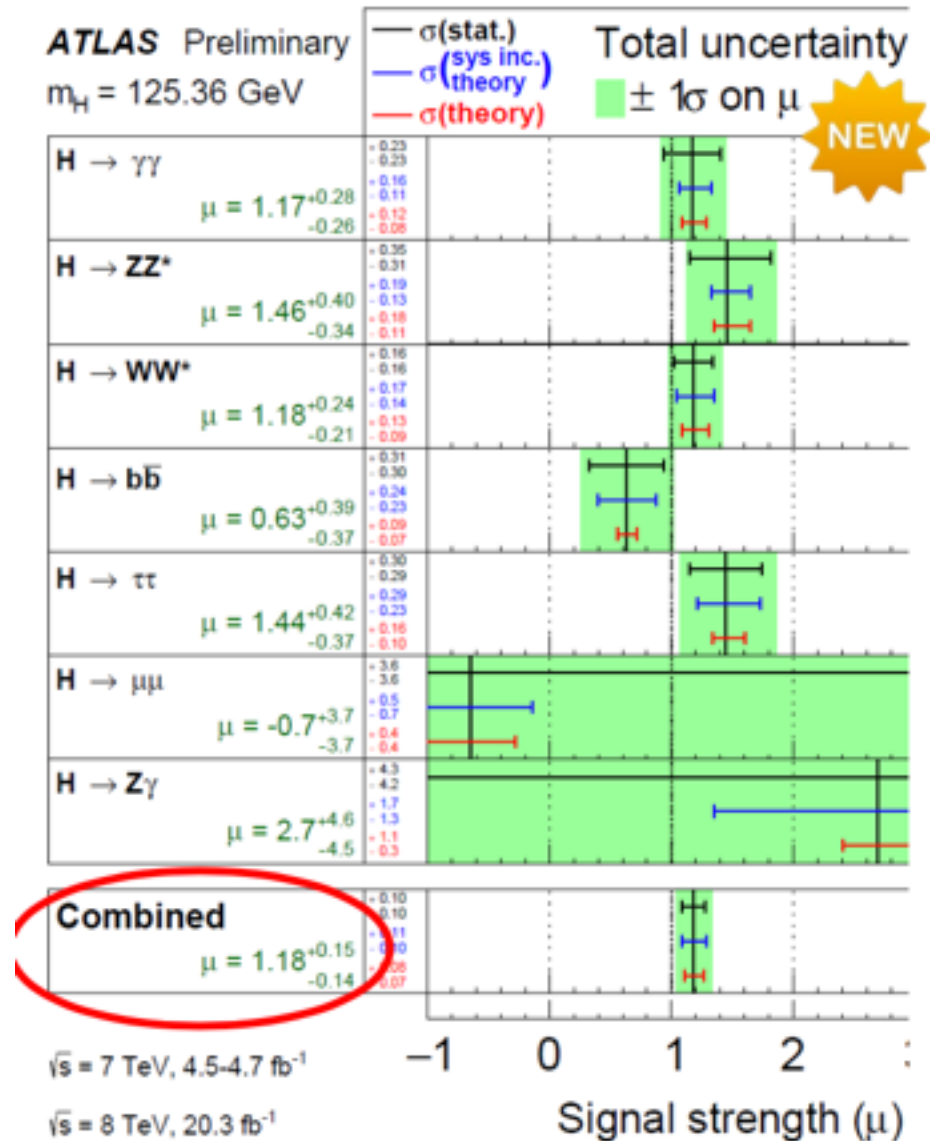
VBF

Dominant production mode at the LHC

- We want to know the gluon-fusion cross section precisely!

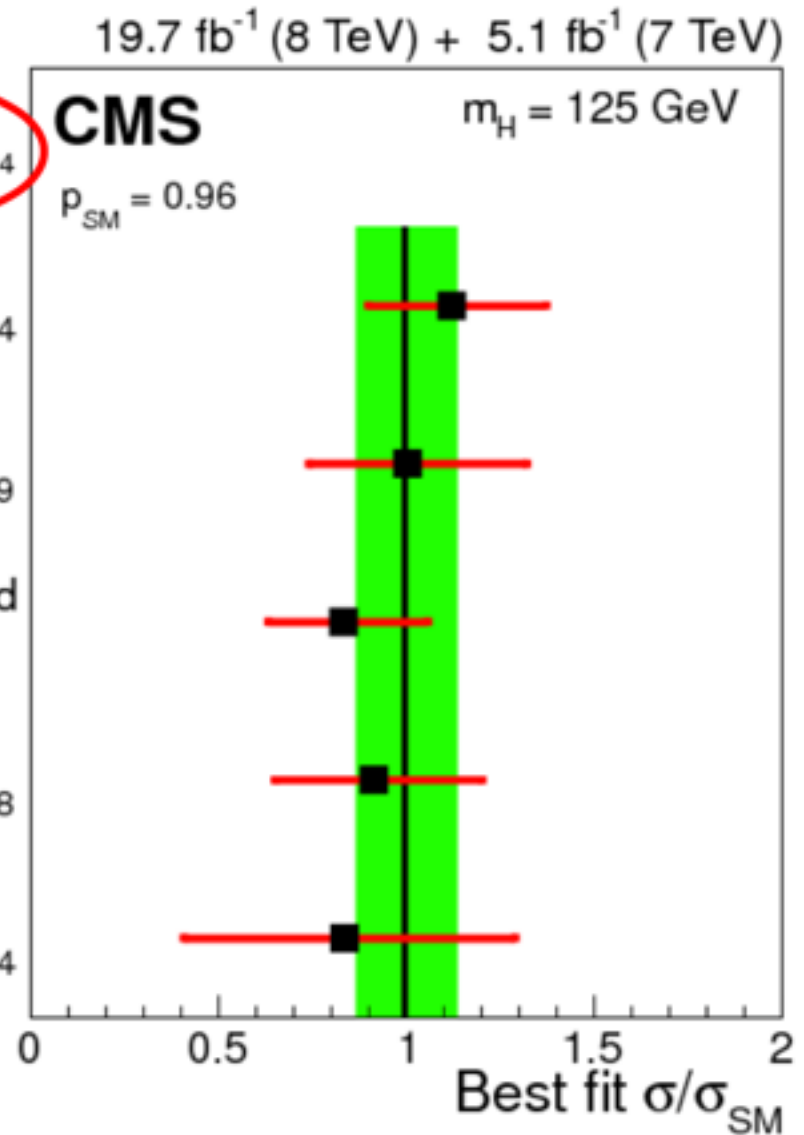


# Higgs production at the LHC



Combined  
 $\mu = 1.00 \pm 0.14$

- $H \rightarrow \gamma\gamma$  tagged  $\mu = 1.12 \pm 0.24$
- $H \rightarrow ZZ$  tagged  $\mu = 1.00 \pm 0.29$
- $H \rightarrow WW$  tagged  $\mu = 0.83 \pm 0.21$
- $H \rightarrow \tau\tau$  tagged  $\mu = 0.91 \pm 0.28$
- $H \rightarrow b\bar{b}$  tagged  $\mu = 0.84 \pm 0.44$



$$\mu_{CMS} = 1.00 \pm 0.14$$

$$\mu_{ATLAS} = 1.18^{+0.15}_{-0.14}$$

$$\text{stat.} = +0.10_{-0.10}$$

$$\text{sys. (inc. theo.)} = +0.11_{-0.10}$$

$$\text{theory} = +0.08_{-0.07}$$

[M. Dührssen @ Moriond EW 2015]



# Higgs production at the LHC



- Aim: Combine recent computation of N<sup>3</sup>LO cross section in the large  $m_t$  limit with other known effects.
  - ➔ finite quark-mass effects.
  - ➔ state-of-the-art PDFs.
  - ➔ electroweak corrections.
  - ➔ resummation of threshold logarithms.
- Assess residual uncertainty on the cross section at the LHC:
  - ➔ Scale, PDF,  $\alpha_S$ .
  - ➔ Truncation of threshold series.
  - ➔ Missing higher orders (N<sup>4</sup>LO and beyond).
  - ➔ Missing N<sup>3</sup>LO PDFs.
  - ➔ Scheme for quark masses + parametric uncertainties.



# Outline



- The N<sup>3</sup>LO cross section in the large  $m_t$  limit:
  - ➔ Convergence of the threshold expansion.
  - ➔ Scale variation.
  - ➔ Effects beyond N<sup>4</sup>LO.
- Other corrections:
  - ➔ Quark mass effects.
  - ➔ Electroweak corrections.
- PDF uncertainties.
- Final prediction for the cross section

The N<sup>3</sup>LO cross section  
in the large  $m_t$  limit

Scale variation  
&  
higher orders in QCD



# The large $m_t$ limit



- In the limit  $m_t \rightarrow \infty$ , the Higgs boson couples directly to gluons:

$$\mathcal{L} = \mathcal{L}_{QCD,5} - \frac{1}{4v} C_1 H G_{\mu\nu}^a G_a^{\mu\nu}$$

- In this limit, the cross section is known

➔ at NLO.

[Dawson; Djouadi, Spira, Zerwas]

➔ at NNLO.

[Anastasiou, Melnikov; Harlander, Kilgore;  
Ravindran, Smith, van Neerven]

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➔ at N3LO.

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- The N3LO cross section is only known as an expansion around threshold:

$$\sigma = \tau \sum_{ij} \int_{\tau}^1 \frac{dz}{z} \mathcal{L}_{ij}(\tau/z) \frac{\hat{\sigma}_{ij}(z)}{z}$$

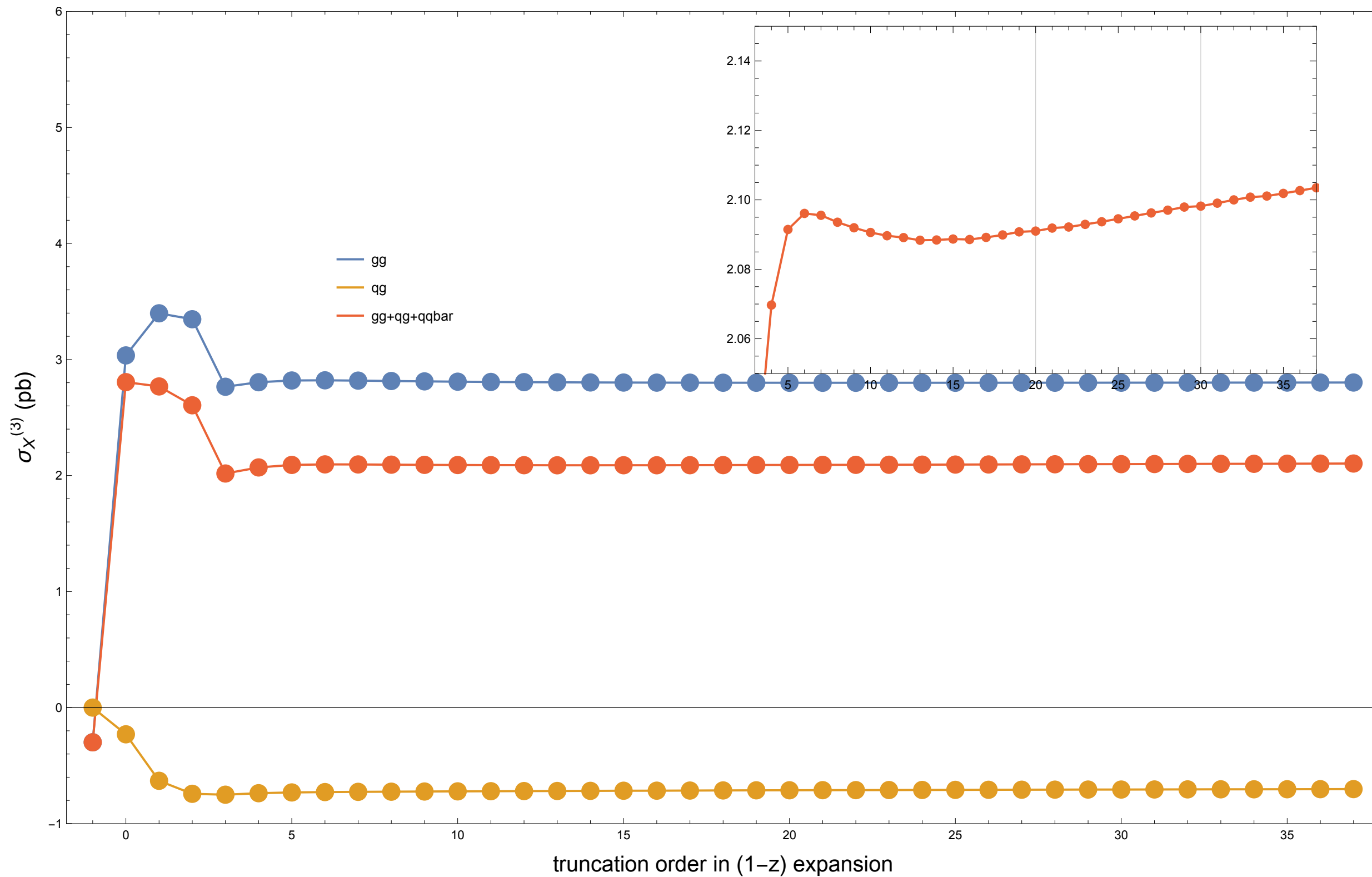
$$z = \frac{m_H^2}{\hat{s}}$$

$$\hat{\sigma}(z) = \sigma_{-1} + \sigma_0 + (1-z)\sigma_1 + \mathcal{O}(1-z)^2$$

$$\tau = \frac{m_H^2}{S} \simeq 10^{-4}$$



# Truncation uncertainty





# Truncation uncertainty



- We estimate the uncertainty from the truncation of the series as:

Conservative factor  $5 \times \frac{\sigma_{EFT}^{(3)}(30) - \sigma_{EFT}^{(3)}(20)}{\sigma_{EFT}^{N^3LO}} = 0.25\%$



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- This is consistent with known exact results for logarithms:

$$\sigma_{EFT}^{(3)} = f_0(z) + f_1(z) \log(1 - z) + f_2(z) \log^2(1 - z) + f_3(z) \log^3(1 - z) \\ + f_4(z) \log^4(1 - z) + f_5(z) \log^5(1 - z)$$



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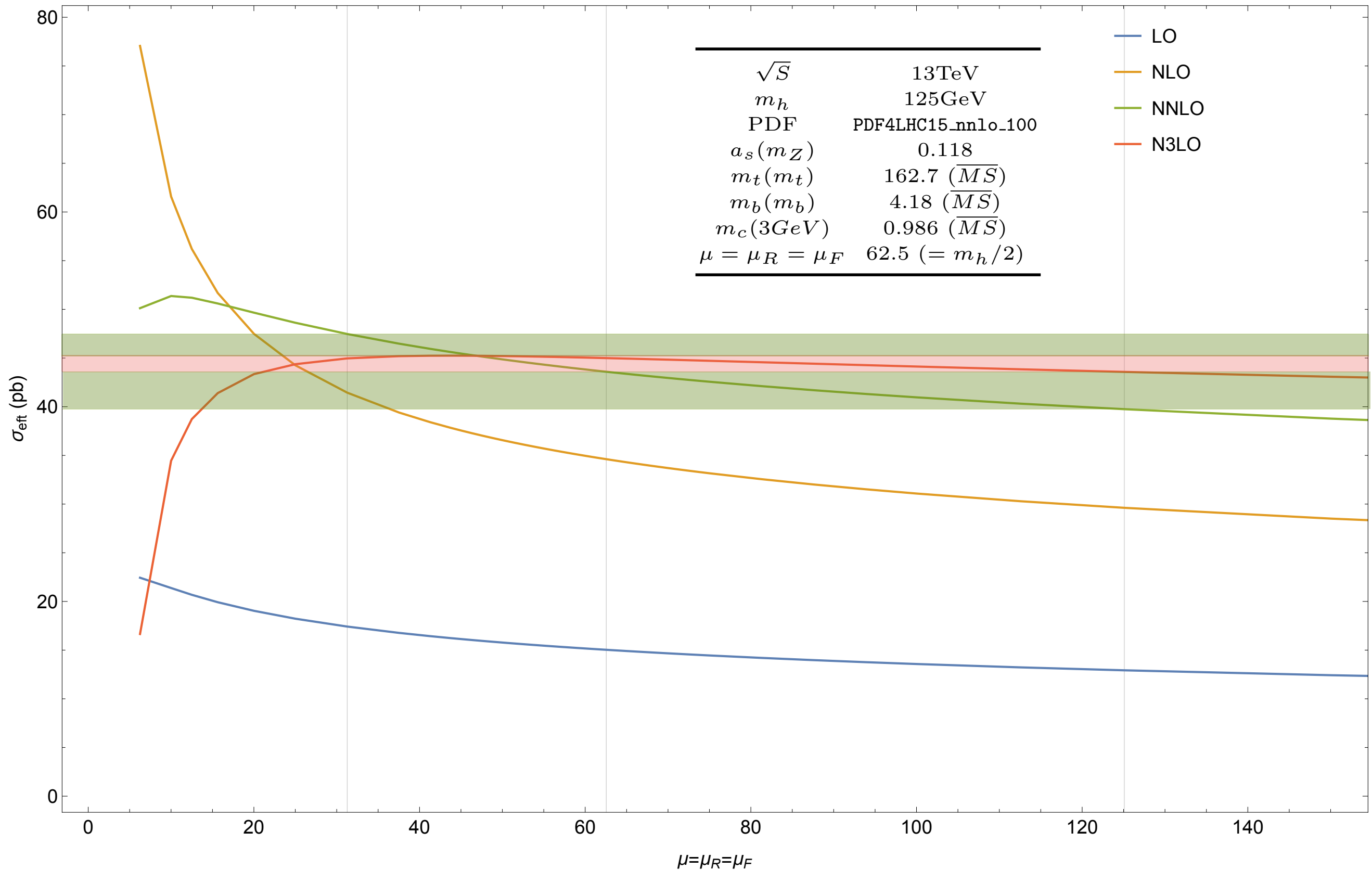
$$\begin{aligned} \sigma_{EFT}^{(3)} = & f_0(z) + f_1(z) \log(1-z) + f_2(z) \log^2(1-z) + f_3(z) \log^3(1-z) \\ & + f_4(z) \log^4(1-z) + f_5(z) \log^5(1-z), \quad \text{Known exactly} \end{aligned}$$

$$\sigma_{EFT}^{(3)} \Big|_{\text{expansion}} - \sigma_{EFT}^{(3)} \Big|_{\text{full logs}} = 0.006 \text{ pb}$$

- The threshold expansion gives a reliable result for the N3LO cross section!

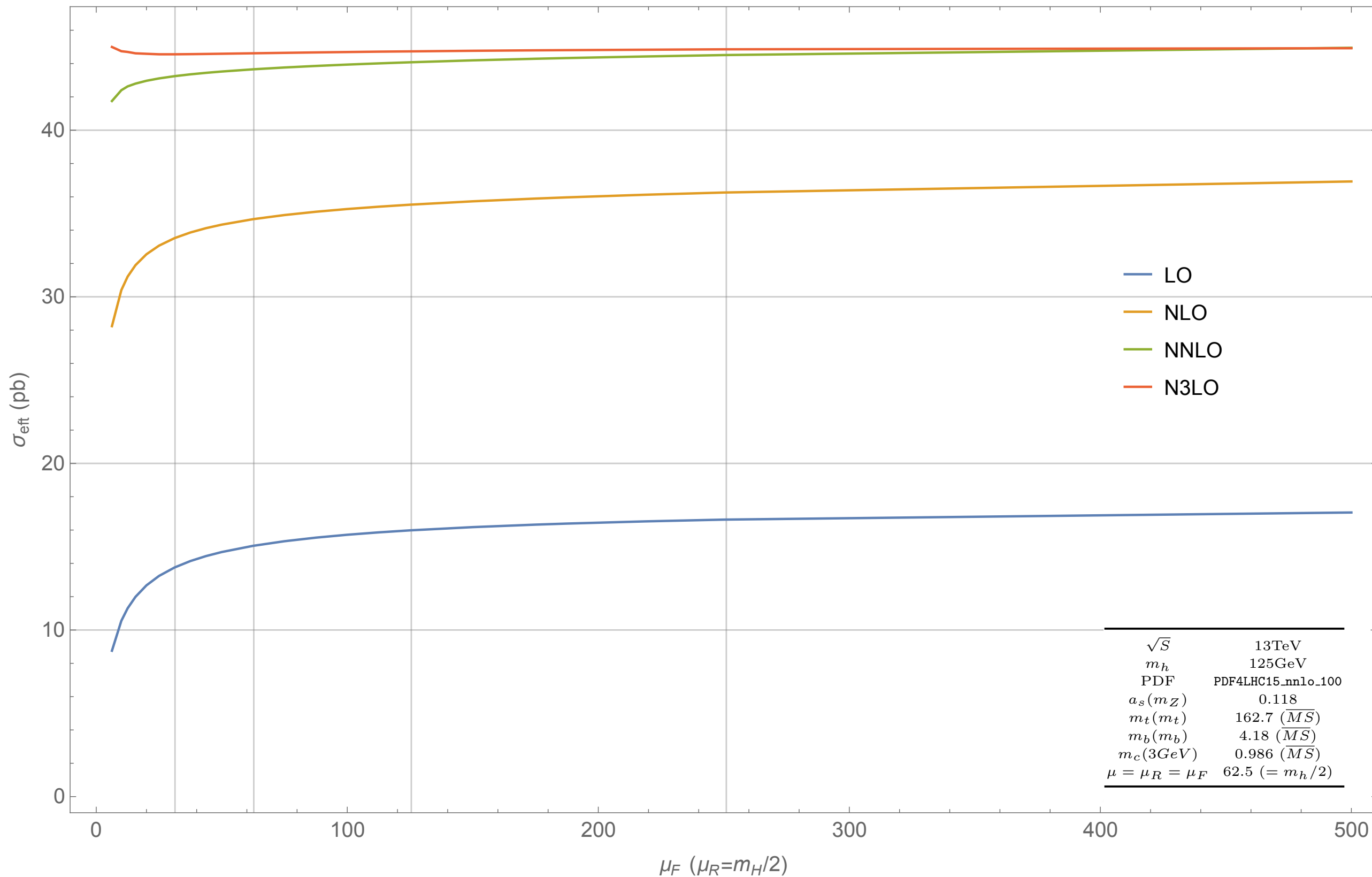


# Scale variation



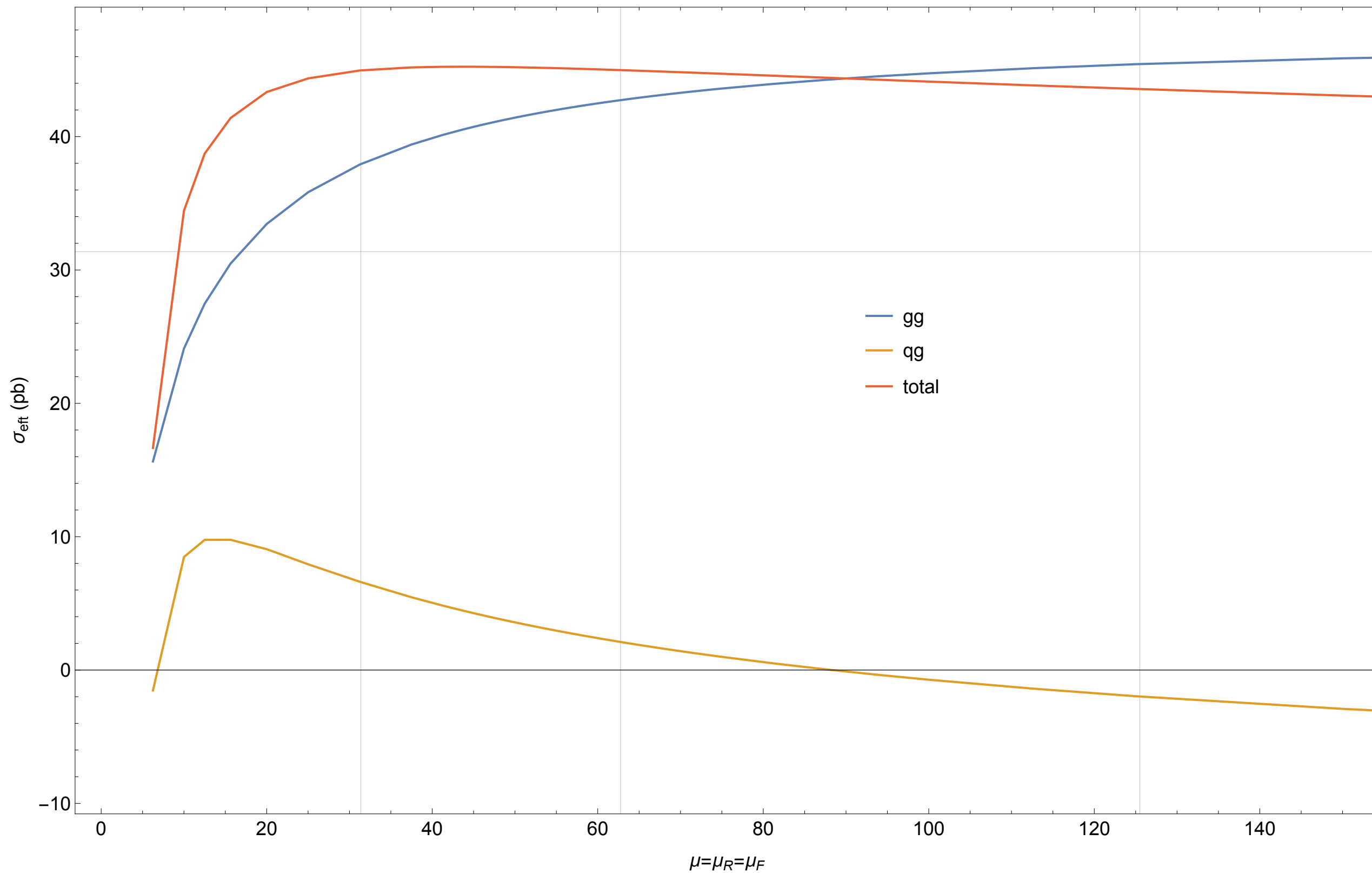


# Scale variation





# Scale variation







# Scale variation



- Scale variation at N<sup>3</sup>LO almost entirely due to renormalisation scale.
- Scale uncertainty  $\mu \in [m_H/4, m_H]$  per order:

$$\Delta_{EFT,k}^{\text{scale}} = \pm \frac{\sigma_{EFT,k}^{\text{max}} - \sigma_{EFT,k}^{\text{min}}}{\sigma_{EFT,k}^{\text{max}} + \sigma_{EFT,k}^{\text{min}}} 100\%$$

		$\Delta_{EFT,k}^{\text{scale}}$
LO	( $k = 0$ )	$\pm 22.0\%$
NLO	( $k = 1$ )	$\pm 19.2\%$
NNLO	( $k = 2$ )	$\pm 9.5\%$
N <sup>3</sup> LO	( $k = 3$ )	$\pm 1.9\%$

- Important question: Is scale variation a reliable estimator of missing higher-order corrections?  
→ We know that it is not at low orders!



# Missing higher orders



- We estimate the effect of missing higher orders in different ways.
  - ➔ Factorisation of the Wilson coefficient.
  - ➔ Threshold resummation in Mellin space (using different prescriptions).
  - ➔ Threshold resummation in SCET (including  $\pi^2$  resummation).



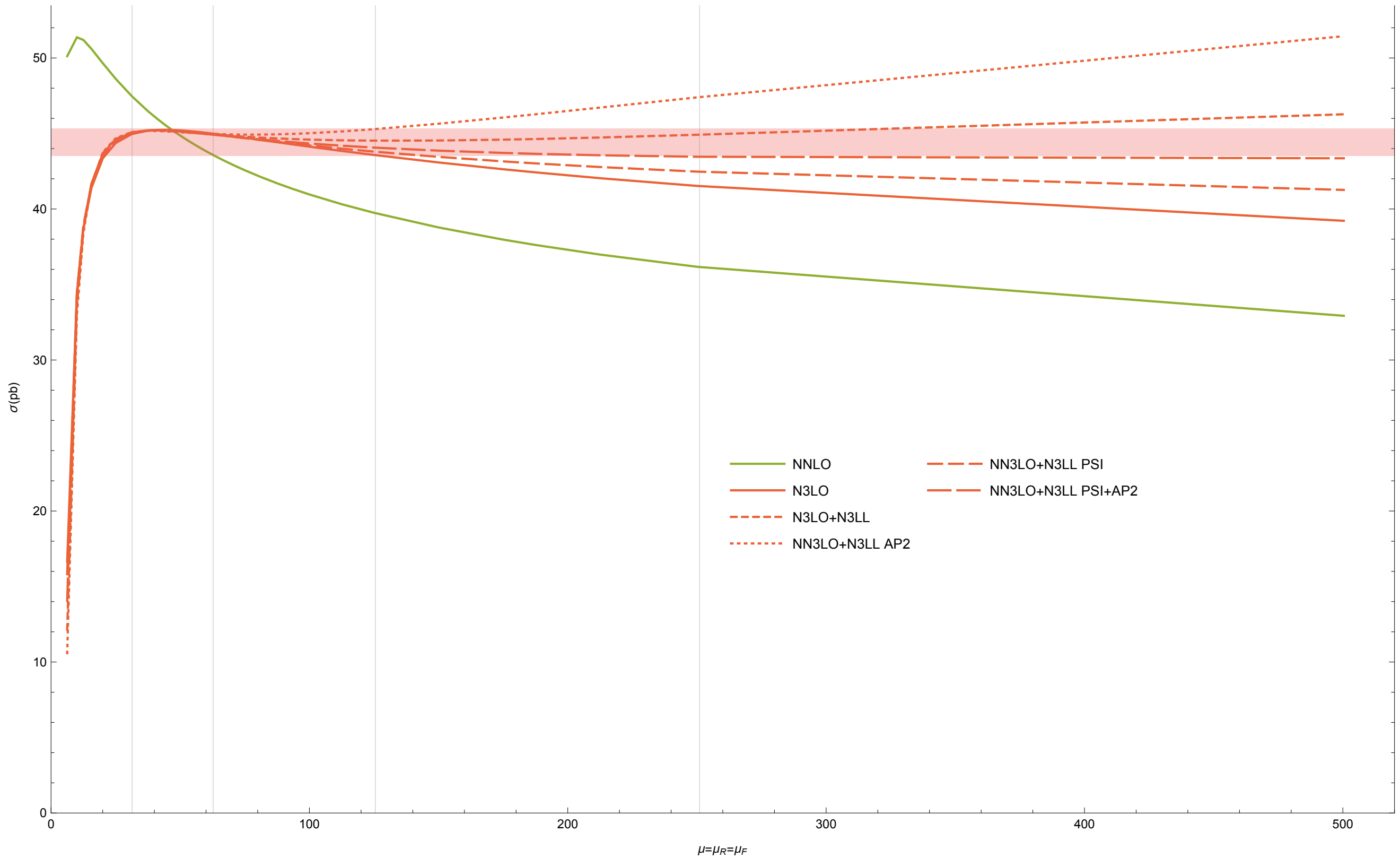
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- In all cases, we find that the higher-order effects are captured by the scale variation with  $\mu \in [m_H/4, m_H]$ .
  - ➔ Scale variation is a reliable estimator of missing QCD corrections for scales  $\mu \in [m_H/4, m_H]$ .
  - ➔ Outside this interval different prescriptions may differ wildly.



# Threshold resummation



Other corrections

Quark masses  
&  
Electroweak corrections



# Quark-mass effects



- At LO and NLO, we know the exact result including all quark mass effects.

➔ EFT works well if rescaled by the LO ratio.

$$R_{LO} \equiv \frac{\sigma_{ex;t}^{LO}}{\sigma_{EFT}^{LO}}$$

$\sigma_{EFT}^{LO}$	15.05	$\sigma_{EFT}^{NLO}$	34.66
$R_{LO} \sigma_{EFT}^{LO}$	16.00	$R_{LO} \sigma_{EFT}^{NLO}$	36.84
$\sigma_{ex;t}^{LO}$	16.00	$\sigma_{ex;t}^{NLO}$	36.60
$\sigma_{ex;t+b}^{LO}$	14.94	$\sigma_{ex;t+b}^{NLO}$	34.96
$\sigma_{ex;t+b+c}^{LO}$	14.83	$\sigma_{ex;t+b+c}^{NLO}$	34.77

➔ Scale uncertainty in rescaled EFT is 1.6% (vs. 1.9% in EFT).



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- At NNLO, we do not know any quark mass effects exactly.

→ Can include top-mass corrections as  $1/m_t$  expansion.

[Harlander, Mantler, Marzani, Ozeren]

→ We do not know t-b interference at NNLO.

$$\Delta_{rEFT}^{NNLO} \frac{\Delta_{t+b}^{NLO} - \Delta_t^{NLO}}{\Delta_t^{NLO}} \simeq \pm 0.38 \text{pb}$$



# Electroweak corrections



- Exact NLO EW corrections are known. [Actis, Passarino, Sturm, Uccirati]
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- ➔ Modified Wilson coefficient.

$$C_{QCD} \rightarrow C_{QCD} + \lambda_{EW} (1 + C_{1w} a_s + C_{2w} a_s^2 + \dots)$$

NLO EW      EW-QCD in EFT approach



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NLO EW

EW-QCD in EFT approach

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NLO EW      EW-QCD in EFT approach

- ➔ Numerical impact is similar to 'complete factorisation' for EW corrections,  $\sim 5.1\%$ .
- We estimate the uncertainty of this approach by varying the mixed QCD-EW Wilson coefficient,

$$\lambda_{EW} (1 + C_{1w} a_s + \dots) \rightarrow \lambda_{EW} (1 + y_\lambda \cdot C_{1w} a_s + \dots)$$

- ➔ Cross section varies by  $+0.4\%$  /  $-0.2\%$  if  $y_\lambda$  is varied by a factor 3.

PDF + aS uncertainties



# PDF + $\alpha_s$ uncertainty



- We follow the PDF4LHC recommendation:
  - ➔ PDF and  $\alpha_s$  error are added in quadrature

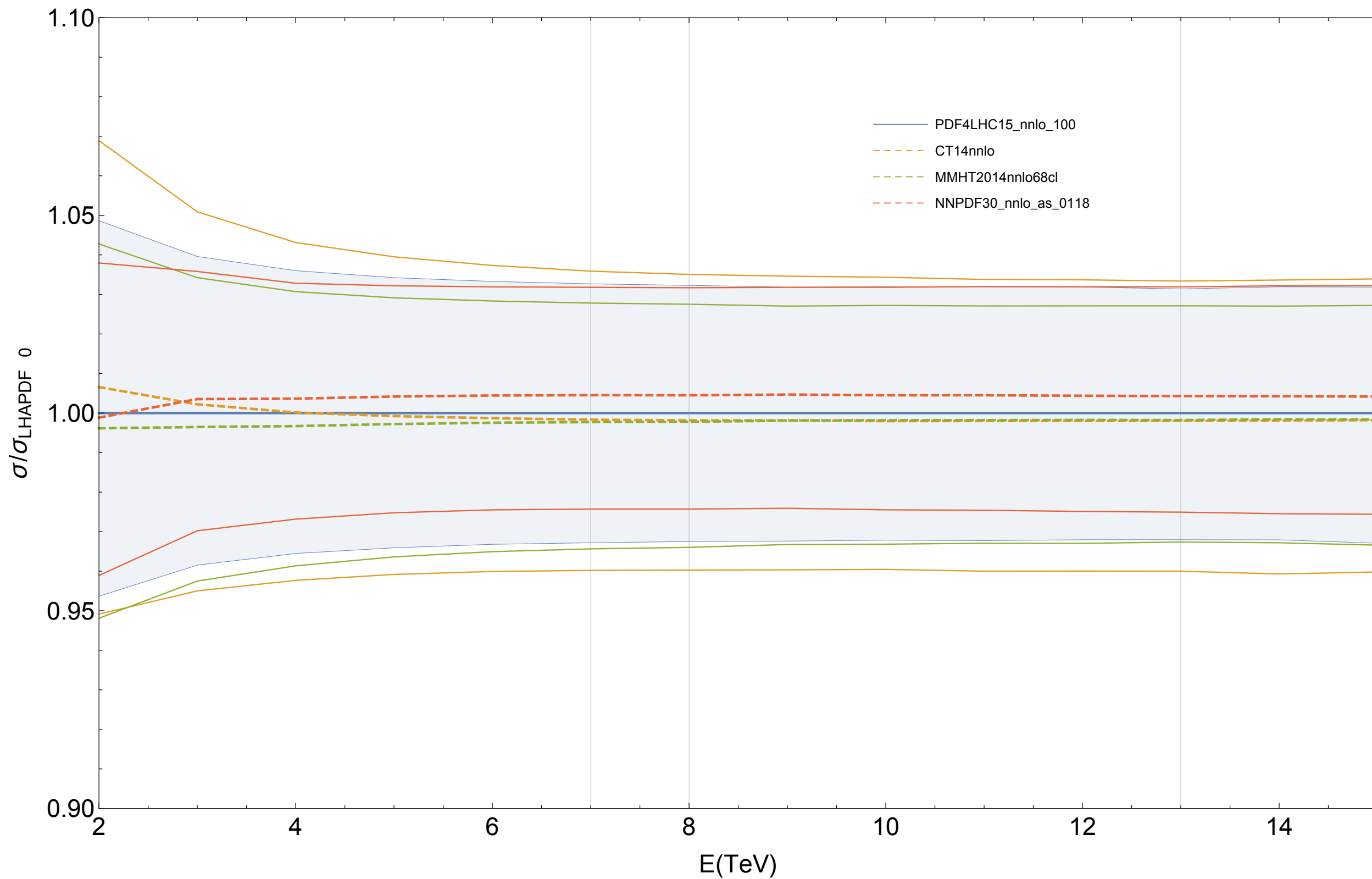
$$\delta_{\text{PDF}+\alpha_s} = \sqrt{\delta_{\text{PDF}}^2 + \delta_{\alpha_s}^2}$$

$$\delta_{\alpha_s} = \frac{\sigma(\alpha_s = 0.1195) - \sigma(\alpha_s = 0.1165)}{2}$$

- ➔  $\delta_{\text{PDF}}$  obtained by using Hessian method.
- We also studied other PDF sets (ABM, HERAPDF).
  - ➔ At LHC energies, HERAPDF in good agreement with PDF4LHC, while ABM 7-9% lower.



# PDF + $\alpha_S$ uncertainty

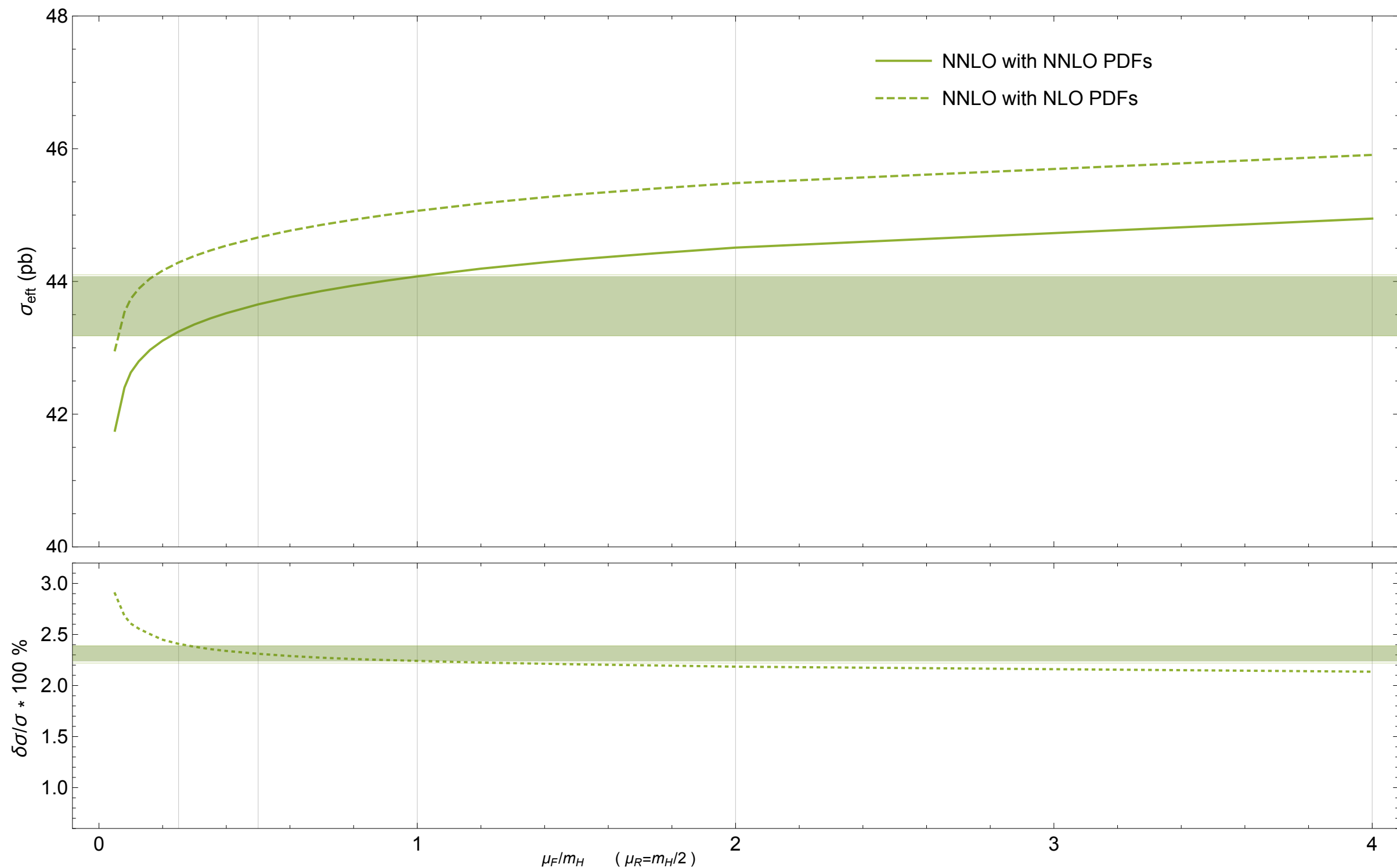




# Missing N3LO PDFs



- All our predictions were made using NNLO PDFs.





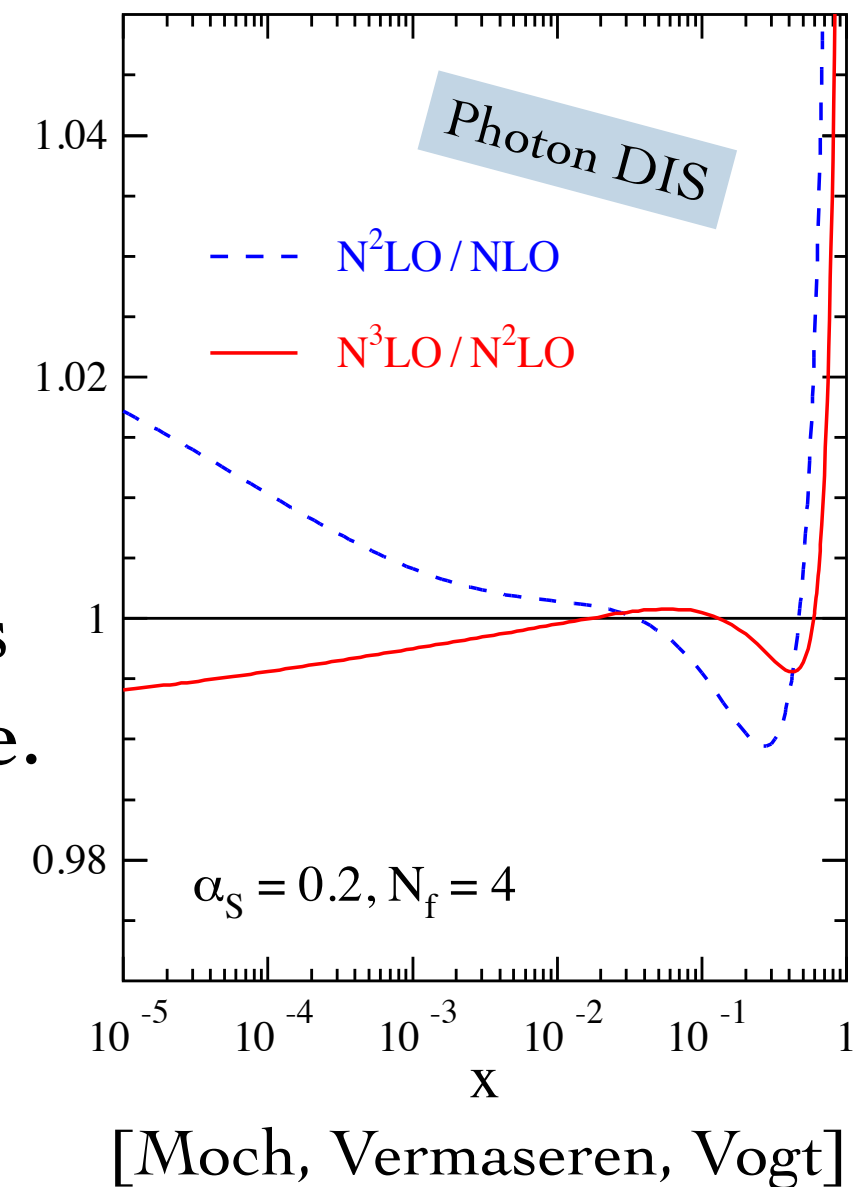
# Missing N3LO PDFs



- Using NLO PDFs at NNLO results in a 2-2.5% error at NNLO.
- From this, we estimate the uncertainty of using NNLO PDFs at N3LO

$$\pm \frac{1}{2} \frac{\sigma_{\text{NNLO PDF}}^{\text{NNLO}} - \sigma_{\text{NLO PDF}}^{\text{NNLO}}}{\sigma_{\text{NLO PDF}}^{\text{NNLO}}} \sigma_{\text{NNLO PDF}}^{\text{N}^3\text{LO}}$$
$$\simeq \pm 0.55 \text{ pb} \simeq \pm 1.15\%$$

- The factor 1/2 takes into account that this estimate is most likely overly conservative.  
→ cf. convergence pattern of DIS.







# Summary



$\sigma$ [pb]	$\delta_{\text{PDF}}$	$\delta_{\alpha_s}$	$\delta_{\text{scale}}$	$\delta_{\text{trunc}}$	$\delta_{\text{PDF-TH}}$	$\delta_{\text{EW}}$	$\delta_{tb}$	$\delta_{1/m_t}$
48.48	$\pm 0.90$ pb	$\pm 1.26$ pb	$^{+0.09}_{-1.11}$ pb	$\pm 0.12$	$\pm 0.56$	$\pm 0.48$	$\pm 0.34$	$\pm 0.48$
	$\pm 1.86\%$	$\pm 2.60\%$	$^{+0.2}_{-2.3}\%$	$\pm 0.25\%$	$\pm 1.15\%$	$\pm 1.00\%$	$\pm 0.70\%$	$\pm 1.00\%$

- Scale choice  $\mu_F = \mu_R \in [m_H/4, m_H]$
- $\delta_{\text{PDF}}$  and  $\delta_{\alpha_s}$  are computed using the PDF4LHC recommendation.
- We have also considered parametric uncertainties on quark masses, and change of renormalisation scheme.
  - ➔ Negligible.
- We do not include threshold resummation effects.
  - ➔ Captured in N3LO scale variation.
- Combination of errors:
  - ➔ PDF and  $\alpha_s$  in quadrature.
  - ➔ the rest is added linearly.



# Summary



$$\sigma = 48.48 \pm 1.55 \begin{matrix} +2.07 \\ -3.09 \end{matrix} \text{ pb} = 48.48 \text{ pb} \pm 3.19\% \begin{matrix} +4.27\% \\ -6.37\% \end{matrix}$$

- Most precise prediction of the Higgs cross section to date!
- Perturbative stability of the cross section under control.
  - ➔ Scale variation gives a reliable estimate of higher-order QCD corrections.
- Places where we can improve:
  - ➔ top-bottom interference at NNLO in QCD.
  - ➔ N3LO PDFs.
  - ➔ Exact mixed QCD-EW corrections.
  - ➔ NNLO corrections including exact top-mass dependence.