

ENTROPY/INFORMATION FLUX IN HAWKING RADIATION

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INTRODUCTION

- It is well known that there exists an **entropy/information “puzzle”** associated with the black hole evaporation process
 - Assuming **unitarity**, we can consider that the information comes out with the radiation (**Page calculations** → firewalls)
- We think that a key point for understanding the problem is to start looking at the entropy/information budget in a **pure statistical mechanics process** where all the physics is under control
- The next step is to apply these considerations to the Hawking evaporation process in **black hole physics**
- For black hole systems, we calculate the **classical and quantum entropies** to show that the Hawking evaporation is relatively benign, no worse than the previous case

ON BURNING A LUMP OF COAL

- It is a unitary process → strict conservation of the **von Neumann entropy**
- Standard statistical mechanics applied to a furnace with a small hole → **blackbody radiation**
 - Planck spectrum: implies some coarse graining
 - Every photon transfers thermodynamic entropy to the radiation field (Clausius entropy)

$$S = \frac{E}{T} = \frac{\hbar \omega}{T}$$

- Now consider the effect of **coarse graining** the (von Neumann) entropy

$$S_{\text{coarse grained}} = S_{\text{before coarse graining}} + S_{\text{correlations}}$$



$$S_{\text{before coarse graining}} = S_{\text{coarse grained}} + I_{\text{correlations}}$$

Information as negative entropy

- On our single-photon definition of **entropy**, it is often convenient to measure the entropy in “natural units of information” (nat)

$$\hat{S} = \frac{S}{k_B} \quad \longrightarrow \quad \text{In an equivalent number of bits: } \hat{S}_2 = \frac{S}{k_B \ln(2)}$$

- The **average energy** per photon in blackbody radiation

$$\langle E \rangle = \hbar \langle \omega \rangle = \hbar \frac{\int \omega f(\omega) d\omega}{\int f(\omega) d\omega} = \frac{\pi^4}{30 \zeta(3)} k_B T$$

- Consequently, the **average entropy**: $\langle \hat{S} \rangle = \frac{\langle E \rangle}{k_B T \ln(2)} \approx 3.897$ bits/photon

- It is easy to calculate also the **standard deviation**

$$\sigma_{\hat{S}_2} = \frac{\sqrt{\langle E^2 \rangle - \langle E \rangle^2}}{k_B T} \approx 2.521 \quad \text{bits/photon}$$

- So, the **main result** we obtained is that the entropy per photon in blackbody radiation is

$$3.897 \pm 2.521 \text{ bits/photon}$$

- Relevant for a photon for which the only thing you know is that it was emitted as part of some blackbody spectrum
- Since we know the underlying physics is unitary —▶ this entropy is compensated by an equal **“hidden information”**
- Hidden in the correlations we choose to ignore in our coarse graining procedure
- The blackbody photons arising from burning a lump of coal carry both thermodynamic entropy and **“hidden information”**

HAWKING EVAPORATION OF A BLACK HOLE

- Using the previous results as a **starting point** to understand what happens in a general relativistic black hole system
 - We have to deal with the existence of horizons —▶ **apparent/trapping horizons**, so the information could escape
 - **Complete** evaporation process
 - We assume then that the evaporation has to be **unitary**
- ➡ Does the information comes out throughout the Hawking emission?
How is the process?
- Page in 1993 developed a similar idea that led to a paradoxical behaviour of the black holes
- We consider a different **model**
 - Compare and contrast the behaviour of the classical thermodynamic entropy with the quantum entanglement entropy

THERMODYNAMIC ENTROPY IN THE HAWKING FLUX

- Loss of **Bekenstein entropy** of a Schwarzschild black hole

$$\frac{dS}{dN} = \frac{dS/dt}{dN/dt} = \frac{d(4\pi k_B GM^2/\hbar c)/dt}{dN/dt} = \frac{k_B \pi^4}{30 \zeta(3)}$$

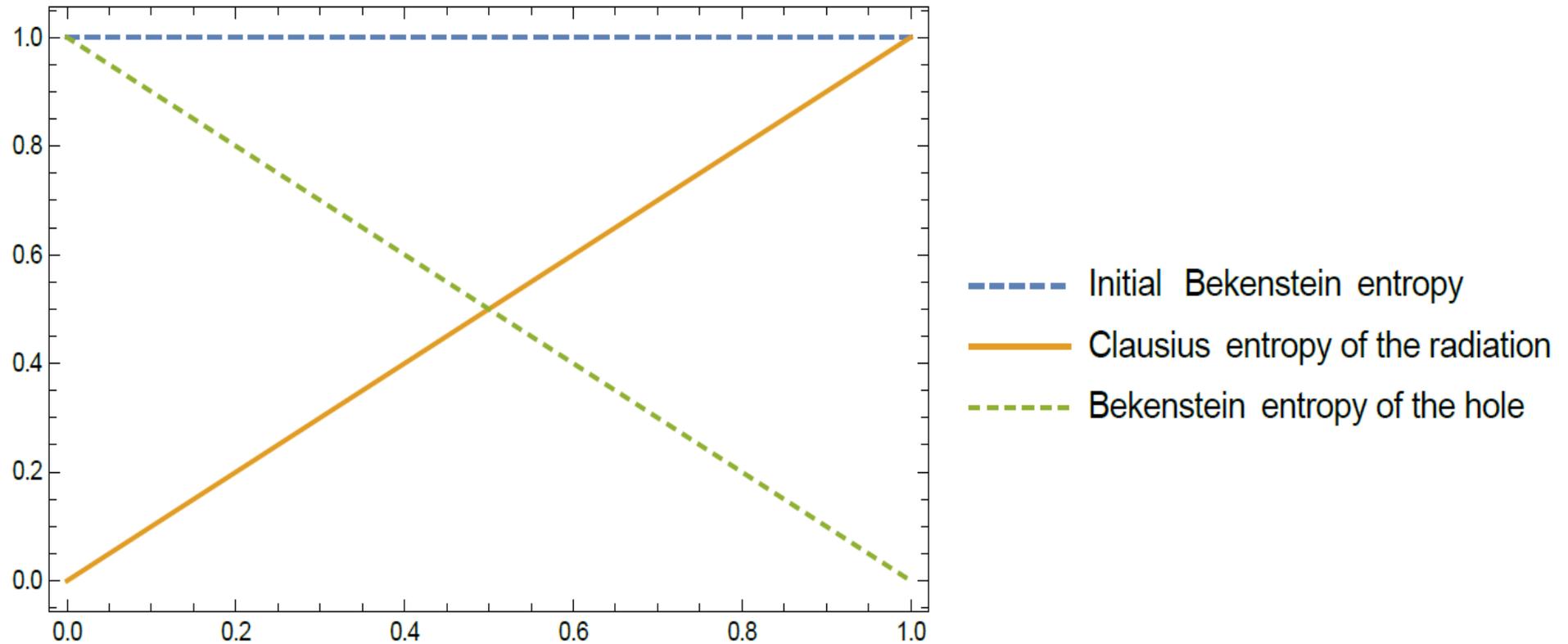
- Thermodynamic entropy gain (**Clausius entropy** gain) of the external radiation field per emitted quanta → entropy gain of the Hawking flux

$$\frac{dS}{dN} = \frac{dE/dT_H}{dN} = \frac{\hbar \langle \omega \rangle dN}{T_H dN} = \frac{k_B \pi^4}{30 \zeta(3)}$$

- The Hawking radiation is essentially (adiabatically) **transferring** Bekenstein entropy from the hole into Clausius entropy of the radiation field

- If we measure the entropy in bits: $\frac{d\hat{S}_2}{dN} \approx 3.897$ bits/quanta

➤ Throughout the **evaporation process** we have $S_{Bekenstein}(t) + S_{Clausius}(t) = S_{Bekenstein,0}$



➤ So, **semiclassically** everything holds together very well

ENTANGLEMENT ENTROPY IN THE HAWKING FLUX

➤ Average subsystem entropies (Page, 1993)

➤ Consider a Hilbert space that factorizes $H_{AB} = H_A \otimes H_B$

Pure state $\rho_{AB} = |\psi\rangle\langle\psi| \longrightarrow$ Subsystem density matrices $\rho_A = \text{tr}_B(|\psi\rangle\langle\psi|)$

Subsystem von Neumann entanglement entropy $\hat{S}_A = -\text{tr}(\rho_A \ln \rho_A)$

➤ Uniform average over all pure states, taking:

$$n_1 = \dim(H_A), \quad n_2 = \dim(H_B) \quad \text{and} \quad m = \min[n_1, n_2]$$

The central result $\hat{S}_{n_1, n_2} = \langle \hat{S}_A \rangle = \langle \hat{S}_B \rangle \leq \ln m$

➤ Average subsystem entropy is very close to its maximum possible value

Strict bound (combined with our results): $\hat{S}_{n_1, n_2} = \langle \hat{S}_A \rangle = \langle \hat{S}_B \rangle \in (\ln m - \frac{1}{2}, \ln m)$

Bipartite entanglement: black hole + Hawking radiation (Page, 1993)

➤ “Closed box” argument

➤ **Initially** there is no yet any Hawking radiation $\left\{ \begin{array}{l} H_R \text{ trivial} \\ H_H \text{ enormous} \end{array} \right. \longrightarrow (\hat{S}_{n_H, n_R})_0 = 0$

➤ After the black hole has **completely evaporated**: H_H is trivial $\longrightarrow (\hat{S}_{n_H, n_R})_\infty = 0$

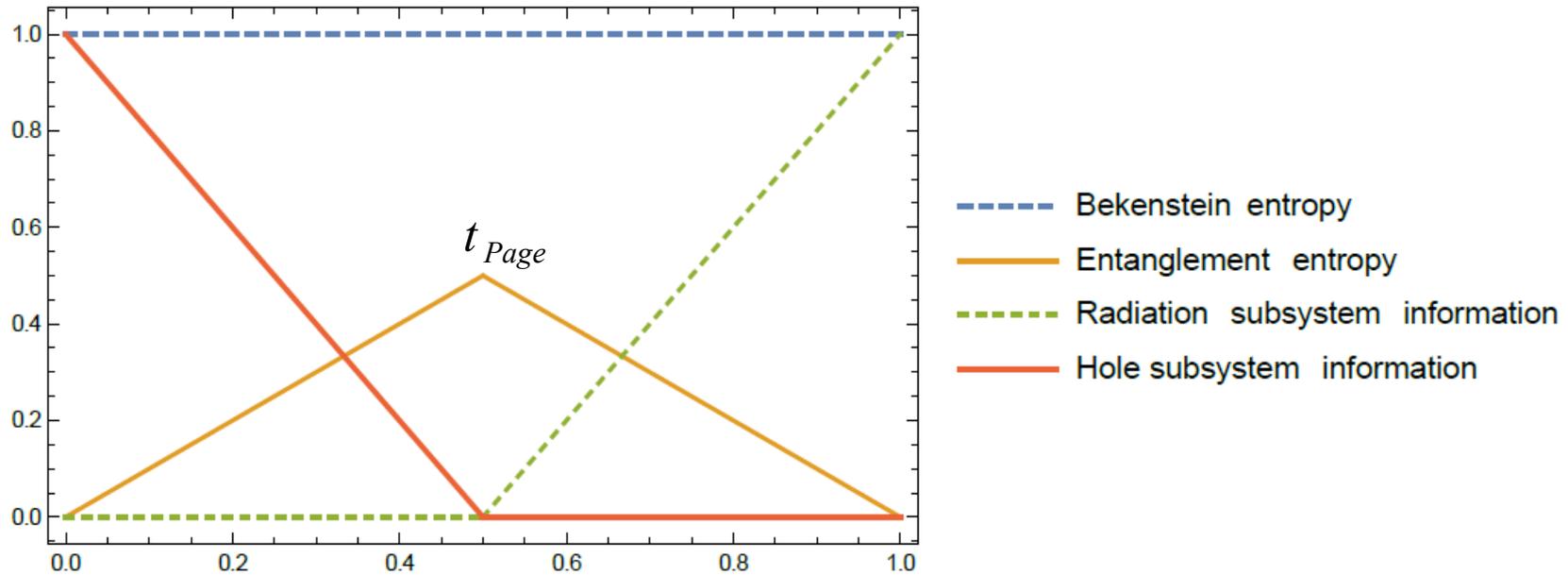
➤ At **intermediate times** both dimensionalities are nontrivial $\longrightarrow (\hat{S}_{n_H, n_R})_t \neq 0$

- Since the evolution is assumed **unitary**: total Hilbert space is constant

$$n_H(t)n_R(t) = n_{H_0} = n_{R_\infty} \longrightarrow (\hat{S}_{n_H, n_R})_t = \ln \left\{ n_H(t), \frac{n_{H_0}}{n_H(t)} \right\}$$

Maximized when $n_H(t) \approx \sqrt{n_{H_0}} \longrightarrow \hat{S}_{n_H, n_R}(t = t_{Page}) \approx \frac{1}{2} \ln n_{H_0}$

➤ Page curve:



- It is the shape of this curve that underlies much of the modern discussion surrounding the “information puzzle”
- Subsystem entropy is initially zero —► tension with Bekenstein entropy
- If we entangle the black hole with the environment, then the total state is not pure

➤ Pages defines a novel **asymmetric** version of the **subsystem information**

$$\tilde{I}_{n_1, n_2} = \ln n_1 - \hat{S}_{n_1, n_2}$$

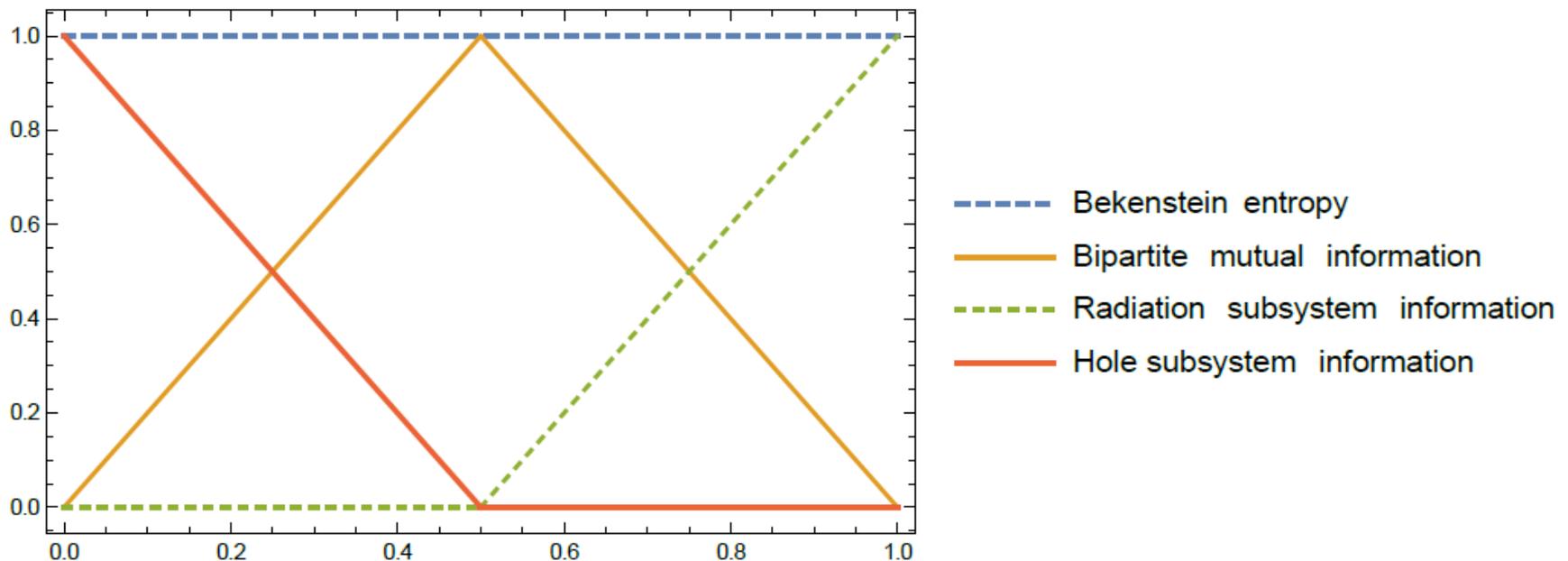
Mutual information

➤ In general: $I_{A:B} = S_A + S_B - S_{AB}$

➤ **Bipartite** entanglement: $I_{H:R} = 2S_H = 2S_R$

➤ Applying the “average subsystem” argument and combined with the Page's asymmetric subsystem information, we found:

$$\langle \tilde{I}_{H,R} \rangle + \langle \tilde{I}_{R,H} \rangle + \langle \hat{I}_{H:R} \rangle = \ln(n_H n_R) \approx \hat{S}_{\text{Bekenstein},0}$$



➤ What happen with these results?

- The late radiation is maximally entangled with the early radiation and with the black hole subsystems



Monogamy of entanglement



Firewalls

➤ We though that the point is the considered “closed box” system

- May be it is more appropriate to consider a tripartite system, including the interaction with the environment

Tripartite entanglement: bh + Hawking radiation + rest of the universe

➤ The Hilbert space is now $H_{HRE} = H_H \otimes H_R \otimes H_E$

➤ Take the entire universe be in a pure state $S_{HRE}(t) = 0$

And now the subsystem entropies $S_H(t) = S_{ER}(t)$, $S_R(t) = S_{HE}(t)$, $S_E(t) = S_{HR}(t)$

➤ Initially $S_{H_0} = S_{E_0}$, $S_{R_0} = 0 = S_{HE_0}$

➤ Once the black hole has completely evaporated $S_{H_\infty} = 0 = S_{ER_\infty}$, $S_{R_\infty} = S_{E_\infty}$

➤ The evolution is assumed unitary, with the unitary time evolution operator factorized as

$$U_{HRE} = U_{HR}(t) \otimes U_E(t)$$

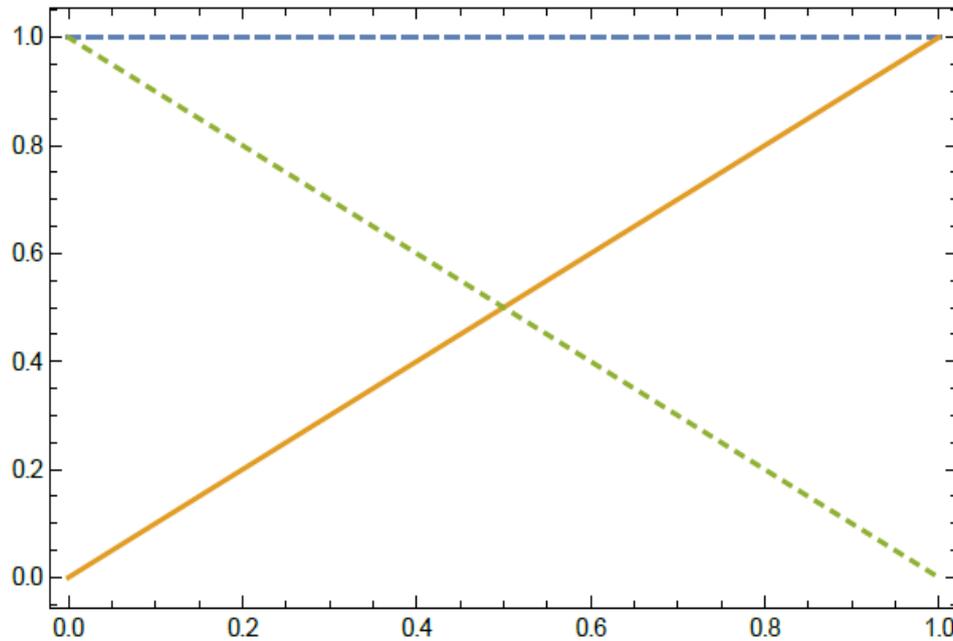
Therefore
$$\begin{cases} n_{E_0} = n_{E_\infty} \equiv n_E \\ n_H(t) n_R(t) = n_{H_0} = n_{R_\infty} \end{cases}$$

➤ We make an additional assumption: That the Bekenstein entropy can be identified with the average entanglement entropy

➤ Then, the entropies:

- $\langle \hat{S}_H(t) \rangle \approx \ln \min[n_H(t), n_R(t)n_E] \approx \ln n_H(t)$
- $\langle \hat{S}_R(t) \rangle \approx \ln \min[n_R(t), n_H(t)n_E] \approx \ln n_R(t)$

$$\left. \begin{array}{l} \langle \hat{S}_H(t) \rangle + \langle \hat{S}_R(t) \rangle \approx \ln [n_H(t)n_R(t)] = \ln n_{H_0} \\ \hat{S}_{Bekenstein}(t) + \langle \hat{S}_{Hawking\ radiation}(t) \rangle \approx \hat{S}_{Bekenstein,0} \end{array} \right\}$$



--- Initial Bekenstein entropy
 — Entanglement entropy of the radiation
 - - - Entanglement entropy of the hole

➤ The “rest of the universe environment”: the extent to which the subsystem is entangled

$$\langle \hat{S}_E(t) \rangle \approx \ln \min[n_E(t), n_H(t)n_R] \approx \ln n_{H_0} \approx \hat{S}_{Bekenstein,0}$$

Mutual information

➤ For the **tripartite system**: $I_{H:R} = S_H + S_R - S_{HR}$ 

➤ **Averaging** over the pure states in the total system we obtain

$$\langle \hat{I}_{H:R} \rangle \leq \frac{n_{H_0}}{2n_E} \leq \frac{1}{2}$$

➡ So the average mutual information never exceeds $\frac{1}{2}$ nat throughout the entire evaporation process

Infinite-dimensional environment

➤ For the **tripartite** system: environment is used to initially entangle the black hole with the rest of the universe

- There is no loss of generality in taking the limit $n_E \rightarrow \infty$

➤ In this limit we therefore have the **equality**

$$\lim_{n_E \rightarrow \infty} \left(\langle S_H \rangle + \langle S_R \rangle \right) = \lim_{n_E \rightarrow \infty} \langle S_E \rangle \quad \text{➡} \quad \lim_{n_E \rightarrow \infty} \langle I_{H:R} \rangle = 0$$

DISCUSSION

- There is no “information puzzle” in [burning a lump of coal](#)
 - Entropy per photon in blackbody radiation is compensated (since the process is unitary) by an equal “hidden information” in the correlations
- We calculated the [classical](#) thermodynamics entropy and the Bekenstein entropy and they compensate perfectly
- We calculate the [quantum \(entanglement\) entropy](#) by means of a tripartite system, which results that is completely agree with the classical expected results, to within 1 nat
- When we restrict attention to a particular subsystem we perceive an amount of [entanglement entropy](#) (a loss of information), there exists a complementary amount of entropy that is codified in the correlations between the subsystems
- Assuming the unitarity of the evolution of the (black hole) + (Hawking radiation) subsystem, we showed that there are [no weird](#) physical effects in the evaporation of a black hole

Thank you for your attention!

References:

- [1] Ana Alonso-Serrano, Matt Visser,
On burning a lump of coal,
Phys. Lett. B 757 (2016) 383.
- [2] Ana Alonso-Serrano, Matt Visser,
Entropy/information flux in Hawking radiation,
arXiv:1512.01890.
- [3] Ana Alonso-Serrano, Matt Visser,
Entropy budget in black hole evaporation process,
Work in progress.