Testing the homogeneity and isotropy of the Universe with recent observations

Elisabetta Majerotto

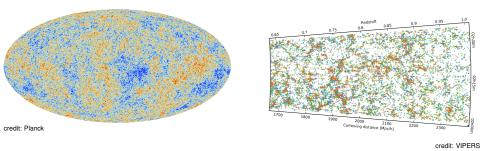
In collaboration with Domenico Sapone and Savvas Nesseris

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Invisibles '15 Workshop: "Invisibles meets Visibles" Madrid, 23rd of June 2015

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Motivation



The CMB and large scale structure show that the Universe is nearly homogeneous and isotropic on large scales.

Motivation

The metric which describes a perfectly homogeneous and isotropic Universe is the **Friedmann-Lemaitre-Robertson-Walker metric**

$$g_{\mu\nu}dx^{\mu}dx^{\nu} = -dt^2 + a^2(t)\left(\frac{dr^2}{1 - kr^2} + d\Omega^2\right)$$

Inserting the FLRW metric into the Einstein equations

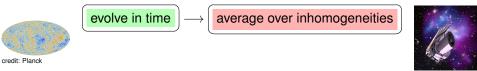
$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R + g_{\mu\nu}\Lambda = \frac{8\pi G}{3T_{\mu\nu}}$$

one obtains the Friedmann equations:

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{k}{a^2} - \frac{\Lambda}{3} = \frac{8\pi G}{3}\rho$$
$$3\frac{\ddot{a}}{a} + 4\pi G\rho - \Lambda = 0$$

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credit: ESA

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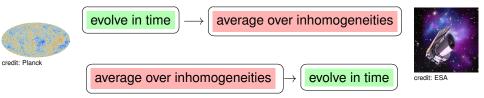
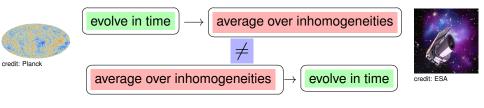


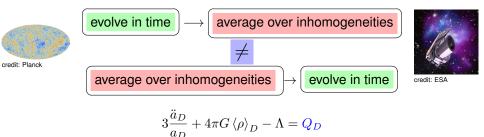
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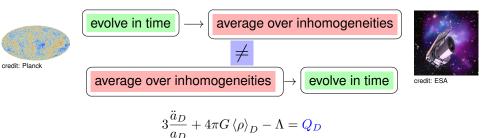


$$3\frac{\ddot{a}_D}{a_D} + 4\pi G \left< \rho \right>_D - \Lambda = Q_D$$

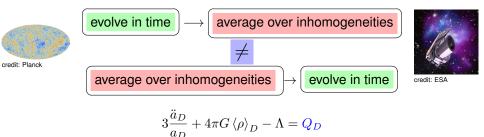
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- Interesting! It could explain the present apparent acceleration of the Universe without finely tuned Λ, extra scalar fields or modifications of gravity
- Difficult: defining univocally an averaging procedure, finding a metric for the coarse-grained spacetime
- Up to now it seems that the effect cannot account for observed Λ, but very interesting anyway for precision cosmology

This talk: testing the FLRW metric

Consistency test for FLRW

of C. Clarkson, B. Bassett and T. Hui-Ching Lu (2008)

Comoving distance in a general FLRW model with curvature:

$$D(z) = \frac{c}{H_0 \sqrt{-\Omega_K}} \sin\left(\sqrt{-\Omega_K} \int_0^z dz' \frac{H_0}{H(z')}\right)$$

 $\Omega_K =$ curvature parameter today.

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Invert the previous equation \Rightarrow expression for Ω_K independent of the specific cosmology:

$$\Omega_K(z) = \frac{\left[H(z)D_{,z}(z)\right]^2 - 1}{\left[H_0D(z)\right]^2} = \begin{cases} const & \text{for FLRW} \\ \Omega_K(z) & \text{otherwise} \end{cases}$$

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By measuring independently H(z), D(z) and H_0 one can test whether Ω_K deviates from a constant.

Data sets

- 19 H(z) data from passively evolving galaxies, *Moresco et al. (2012)*
 - select most massive red elliptical galaxies, with no signature of star formation $\to t_{gal} \sim t_{Universe}$
 - compute their ages $t_{gal} \Rightarrow dt_{gal}/dz$
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compute SNIa distance modulus

$$\mu(z) = m(z) - M = 5\log_{10}(d_L(z)) + 25$$

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H₀ from HST and Wide Field Camera 3, *Riess et al. (2011)*:

$$H_0 = 73.8 \pm 2.4 \ \mathrm{km \, s^{-1} \, Mpc^{-1}}$$

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Discrete approach: binning $\rightarrow \Omega_K(z_k)$ Continuous approach: reconstructing $\rightarrow \Omega_K(z), z \in \{z_{\min}, z_{\max}\}$

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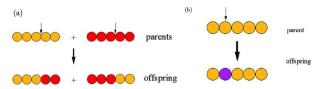
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Best model-independent reconstruction: Genetic Algorithms

Based on principles of evolution through natural selection:

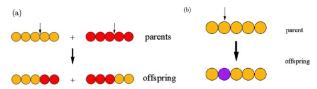


(a) crossover = combination of different individuals
(b) mutation = a random change in an individual

Probability of reproductive success \propto fitness of individual

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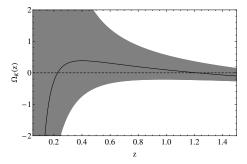
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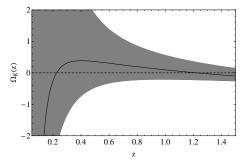
In our case:

- fitness of individual $\rightarrow \chi^2$ function
- initial population \rightarrow set of functions, the grammar, and set of operators
- in each generation, crossover and mutation are applied
- process repeated several thousand times until e.g. max number of generations or desired convergence reached.

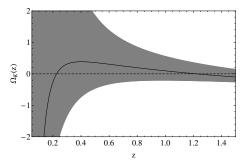
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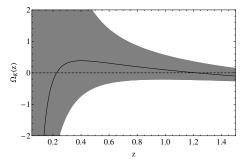




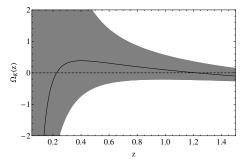
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- GA give the smallest errors with respect to other techniques: $\sigma_{\Omega_K} \sim 0.1$ due to smooth and analytical expression at all *z*
- reconstruction consistent with $\Omega_K = 0$

Error regions computation in GA

Analytical method devised in *Nesseris & Garcia-Bellido (2012) based on* path integral formalism.

To calculate the 1σ error δf_i around the best-fit $f_{bf}(x)$ at a point x_i :

$$CI(x_i,\delta f_i) \equiv \int_{f_{bf}(x_i)-\delta f_i}^{f_{bf}(x_i)+\delta f_i} df_i \frac{1}{(2\pi)^{1/2}\sigma_i} \exp\left(-\frac{1}{2}\left(\frac{y_i-f_i}{\sigma_i}\right)^2\right) = \operatorname{erf}\left(1/\sqrt{2}\right)$$

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This gives a discrete set of $\{x_i, \delta f_i\}$.

We want a smooth function, so we assume a shape for df:

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and minimise simultaneously

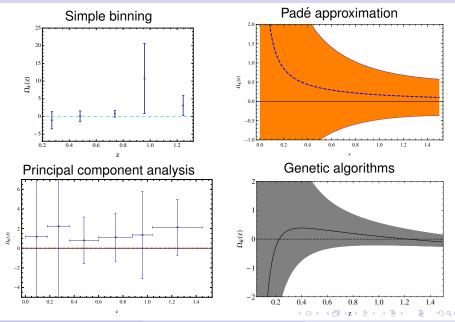
$$\chi^2_{CI}(\delta f_i) = \sum_{i=1}^N \left(CI(x_i, \delta f_i) - \operatorname{erf}\left(1/\sqrt{2}\right) \right)^2$$

and

$$\chi^2(f_{bf} + df_i) \equiv \sum_{i=1}^N \left(\frac{y_i - (f_{bf} + \delta f_i)}{\sigma_i}\right)^2$$

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Other model-independent methods



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Forecasts for future experiments

Euclid

Medium-size mission of the ESA Cosmic Vision programme, launch planned for 2020. The spectroscopic survey will measure ~ 50 million galaxies with slitless spectroscopy, over 15,000 square deg.

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Observed power spectrum:

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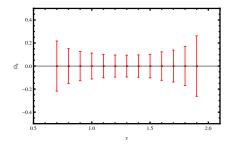
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- $0.65 < z < 2.05, \Delta z = 0.1$
- scale-independent bias approximation, OK for large scales, as in Orsi et al. (2011)
- number densities as in EM, L. Guzzo et al (2012), from a sophisticated simulation
- WMAP-7 ΛCDM fiducial model

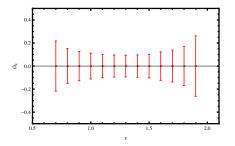
Forecasts for future experiments

Euclid spectroscopy alone



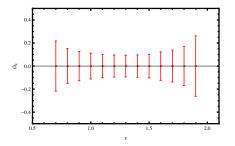
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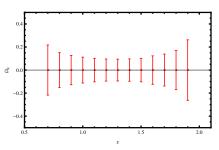


• errors are smallest when 1.1 < z < 1.5

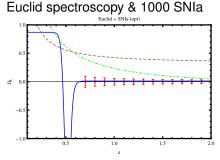
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- errors are smallest when 1.1 < z < 1.5
- at lower and higher *z* constraints become worse by up to a factor of 2.5

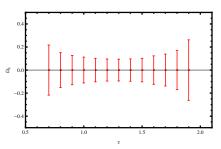


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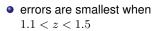


LTB, Tardis model, timescape.

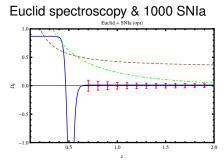
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Euclid spectroscopy alone

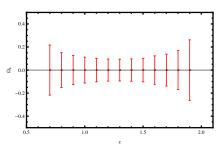


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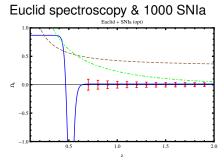
LTB, Tardis model, timescape.

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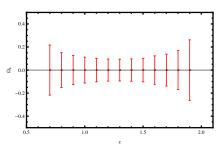
Euclid spectroscopy alone

- errors are smallest when 1.1 < z < 1.5
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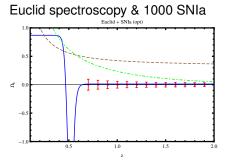
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LTB, Tardis model, timescape.

- $\bullet~$ errors improve by a factor of ~ 2
- best constrained area: z > 1.5
- LTB cannot be distinguished, Tardis model and timescape may (but need to include error on *H*(*z*), SNIa systematics etc...)

Image: Image:

Why is Ω_K badly constrained at small *z*?

Model-independent proof

Expand $\Omega_K(z)$ for small *z*:

 $\tilde{\Omega}_{K} = \Omega_{k} + \frac{1 - \frac{H_{1,0}^{2}}{H_{0}^{2}}}{z^{2}} + \frac{\frac{2H_{0}H'(0) - H_{1,0}H'_{1}(0)}{H_{0}^{2}} - \frac{H'_{1}(0)}{H_{1,0}}}{z} + \dots$

' = d/dz

 H_0 and $H_{1,0}$ are the values of the two Hubble parameters at z = 0.

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$$\Rightarrow$$
 divergence as z^2 unless $H_{1,0}=H_0$

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- Future data from Euclid improve considerably the errors: Euclid only by a factor \sim 10, Euclid + SNIa factor up to 40

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