

Global constraints on vector-like effective WIMP interactions

Pedro A N Machado

in collaboration with

M Blennow, P Coloma, E Fernandez-Martinez, B Zaldivar



WIMP searches: Where do we stand?

What is the global status of WIMP searches?

Different experiments are sensitive to different DM masses, channels, etc

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How can we be model independent?

We can't.



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Effective field theory

We will examine a part of the WIMP – SM EFT parameter space:

WIMP searches: Where do we stand?

Effective field theory

We will examine a part of the WIMP – SM EFT parameter space:

- 1) Vector couplings
- 2) Dirac fermion dark matter
- 3) DM is a standard model singlet
- 4) $SU(2) \times U(1)$ invariant operators
- 5) Dimension 6, DM-DM--- Ψ - Ψ operators

$$c_i (\bar{\chi} \gamma^\mu \chi) (\bar{f}_i \gamma_\mu f_i)$$

EFT proxies

“General model”

$c_{e_R}, c_{\mu_R}, c_{\tau_R}$

$c_{\ell e_L}, c_{\ell \mu_L}, c_{\ell \tau_L}$

$c_{u_R}, c_{c_R}, c_{t_R}$

$c_{d_R}, c_{s_R}, c_{b_R}$

$c_{Q1_L}, c_{Q2_L}, c_{Q3_L}$

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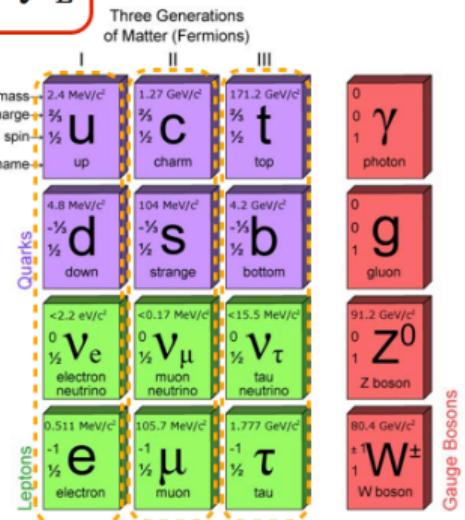
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Three Generations of Matter (Fermions)								
	I	II	III					
mass-	2.4 MeV/c ²	1.27 GeV/c ²	171.2 GeV/c ²					
charge	2/3	2/3	2/3					
spin	1/2	1/2	1/2					
name-	u	c	t					
Quarks	d	s	b					
Leptons	e	ν_μ	ν_τ					
Gauge Bosons	γ	g	Z ⁰					

Detailed description: The table illustrates the three generations of fermions in the Standard Model. It is organized into three columns (I, II, III) representing generations and three rows (top, middle, bottom) representing quarks and leptons. Each cell contains the name of the particle, its mass, charge, spin, and a small icon. The top row contains the up-type quarks (u, c, t) and the photon. The middle row contains the down-type quarks (d, s, b) and the gluon. The bottom row contains the electron neutrino (ν_e), muon neutrino (ν_μ), tau neutrino (ν_τ), and the Z boson. The last two rows are empty.

“Fa

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Ingredients

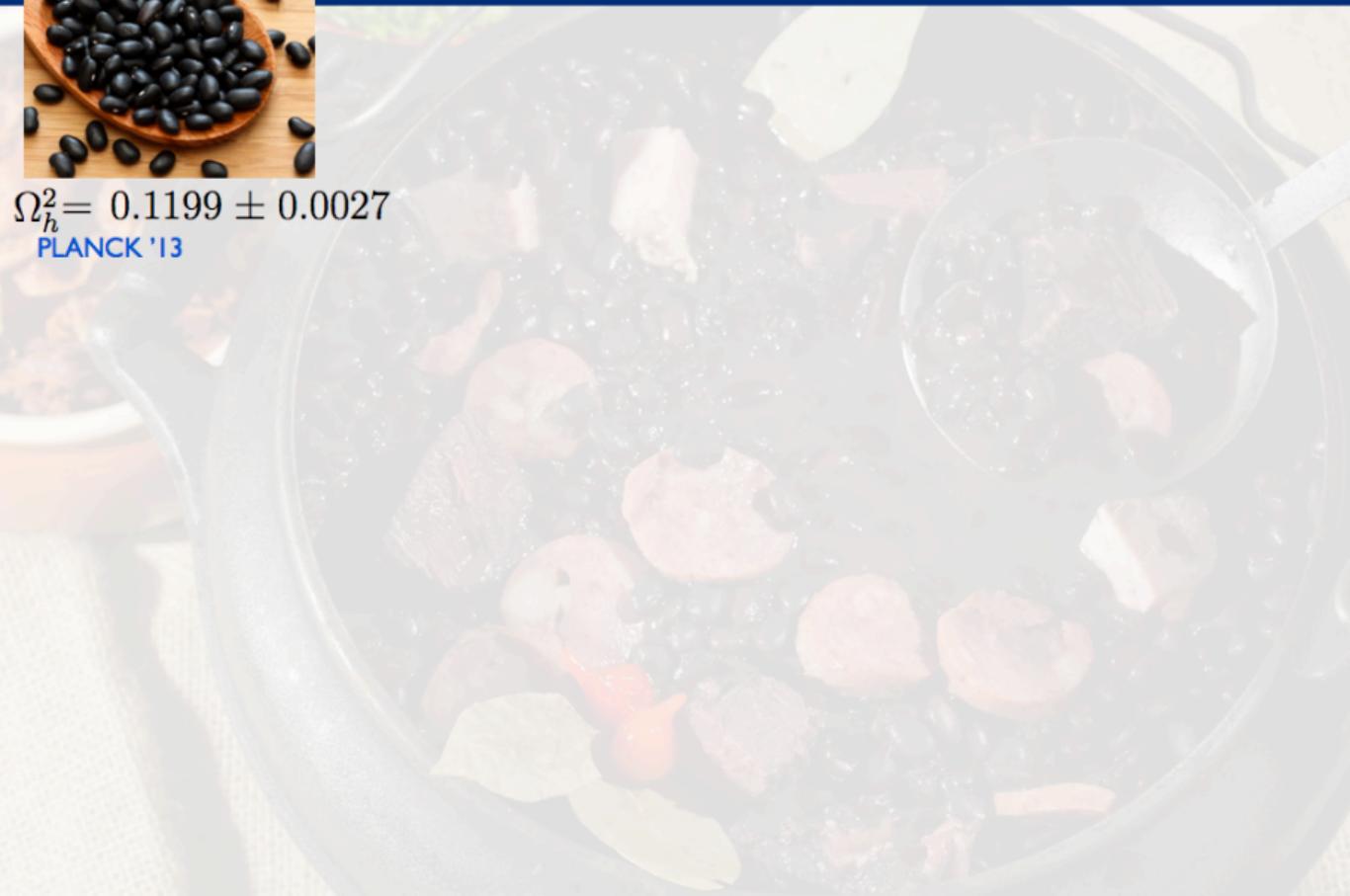




Constraints

$$\Omega_h^2 = 0.1199 \pm 0.0027$$

PLANCK '13



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CMB constraints
Slatyer Padmanabhan Finkbeiner '09
Lopez-Honorez et al '13

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Dwarf galaxies: Fermi-LAT
Fermi-LAT '15

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LHC direct searches
ATLAS '15



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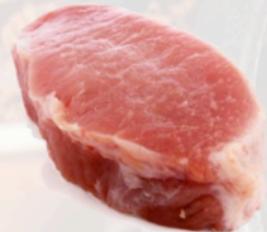


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AMS positron fraction
AMS '14

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Direct detection
LUX '14 EDELWEISS-II '11



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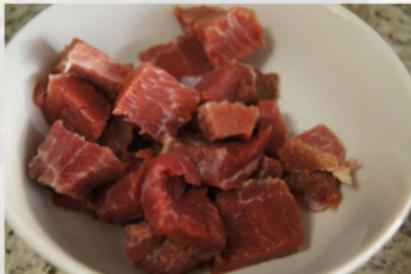
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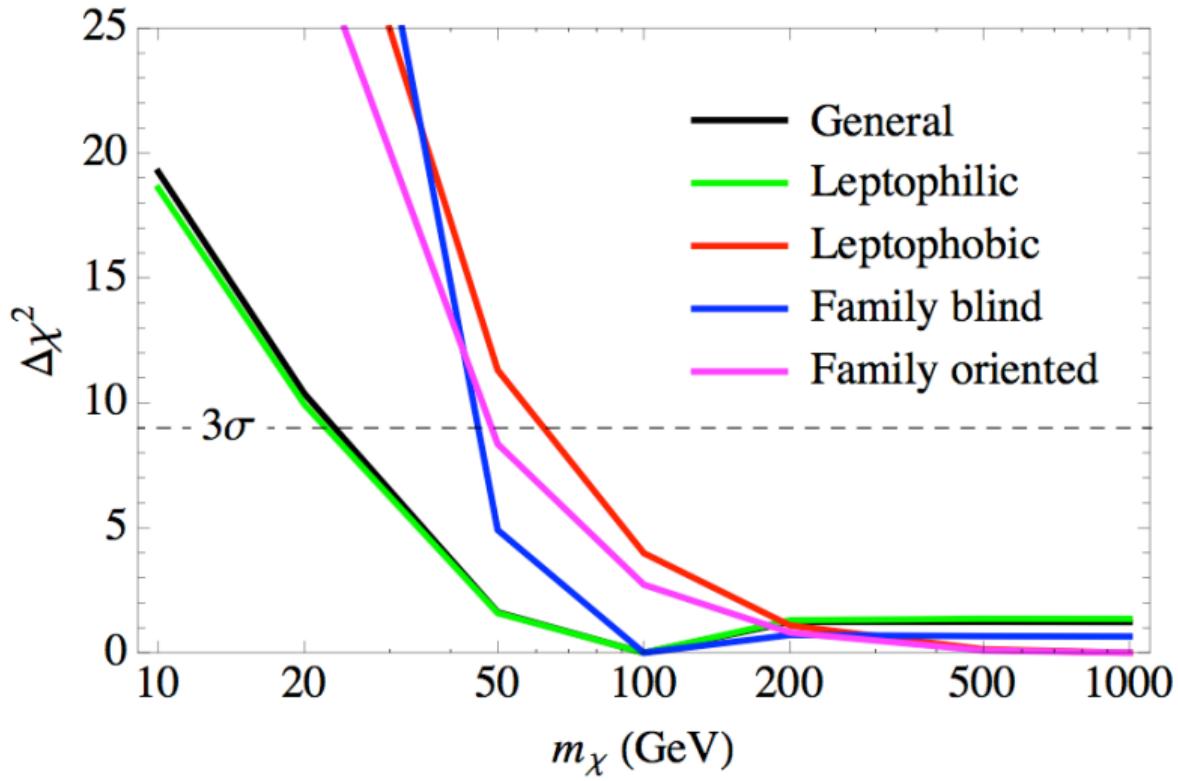


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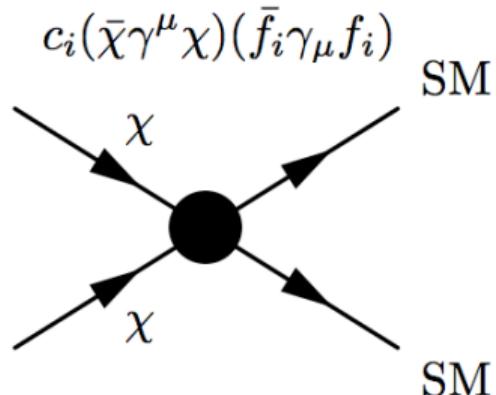


Micromegas + MultiNest + CalcHEP

Dependence with DM mass



Relic abundance



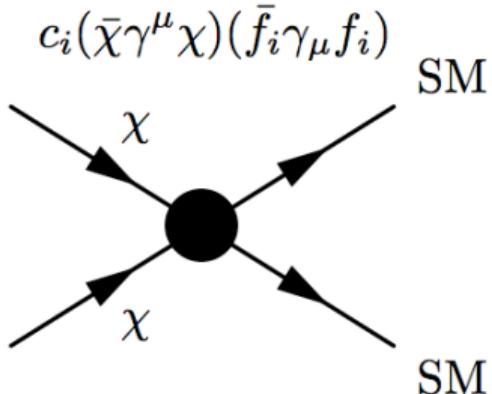
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PLANCK '13

We assume thermal relic

Larger c_i means more annihilation
which leads to less DM

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$$(\sigma v_{\text{rel}})_i \approx w_i \frac{(c_{i_L} + c_{i_R})^2}{24\pi} (3m_\chi^2 + m_\chi^2 v^2)$$

$$\Sigma_i \equiv \sqrt{\sum_i w_i c_i^2}$$

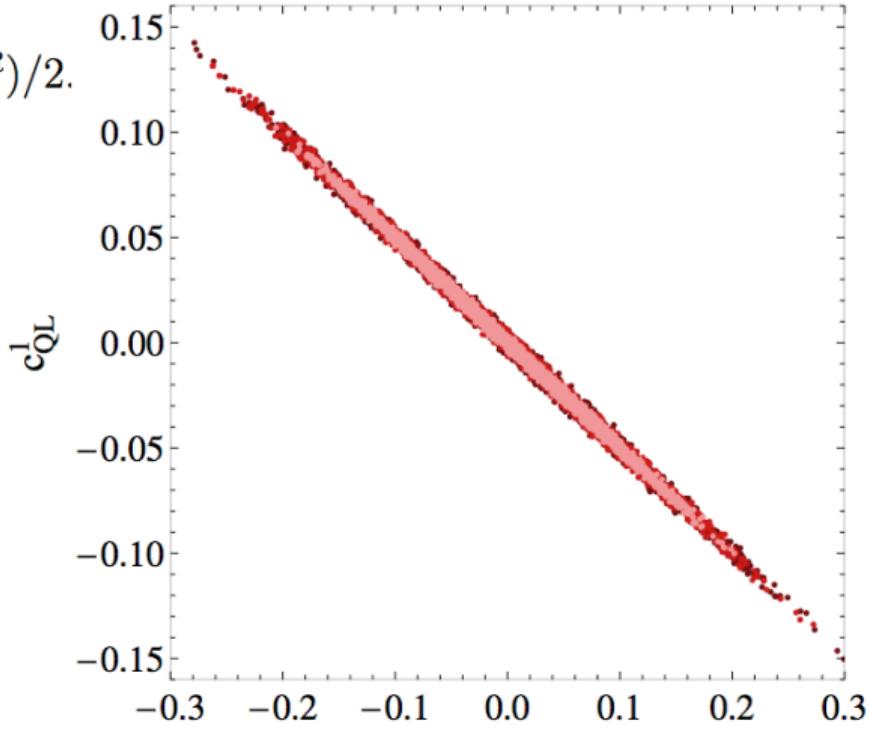
General model

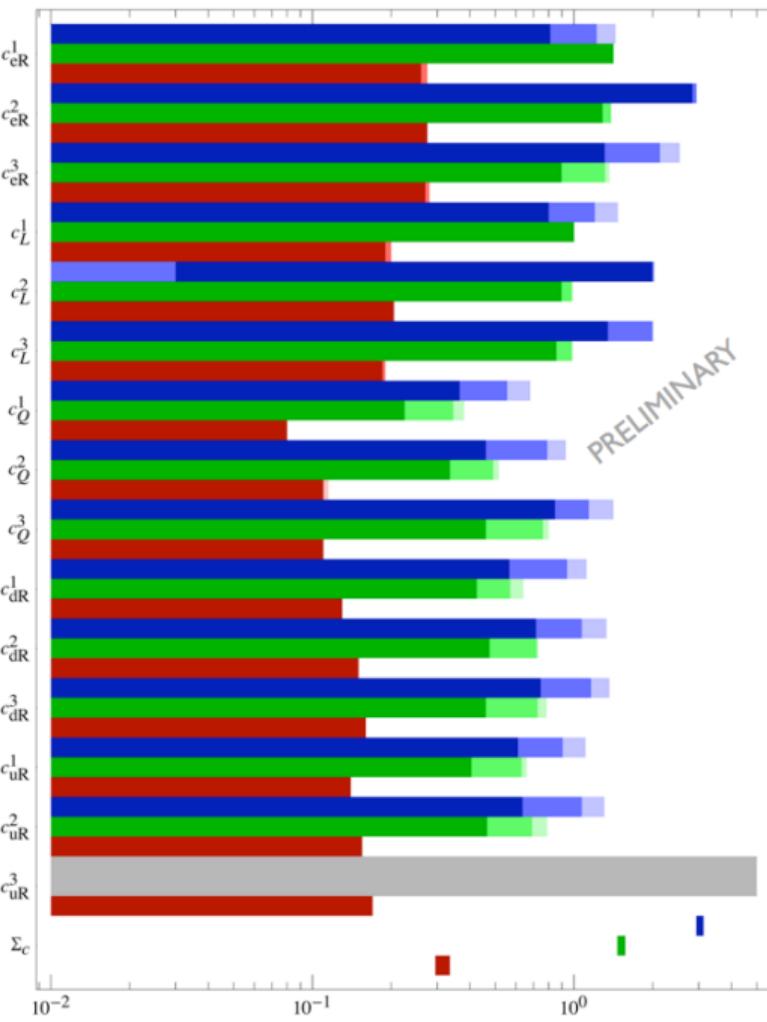
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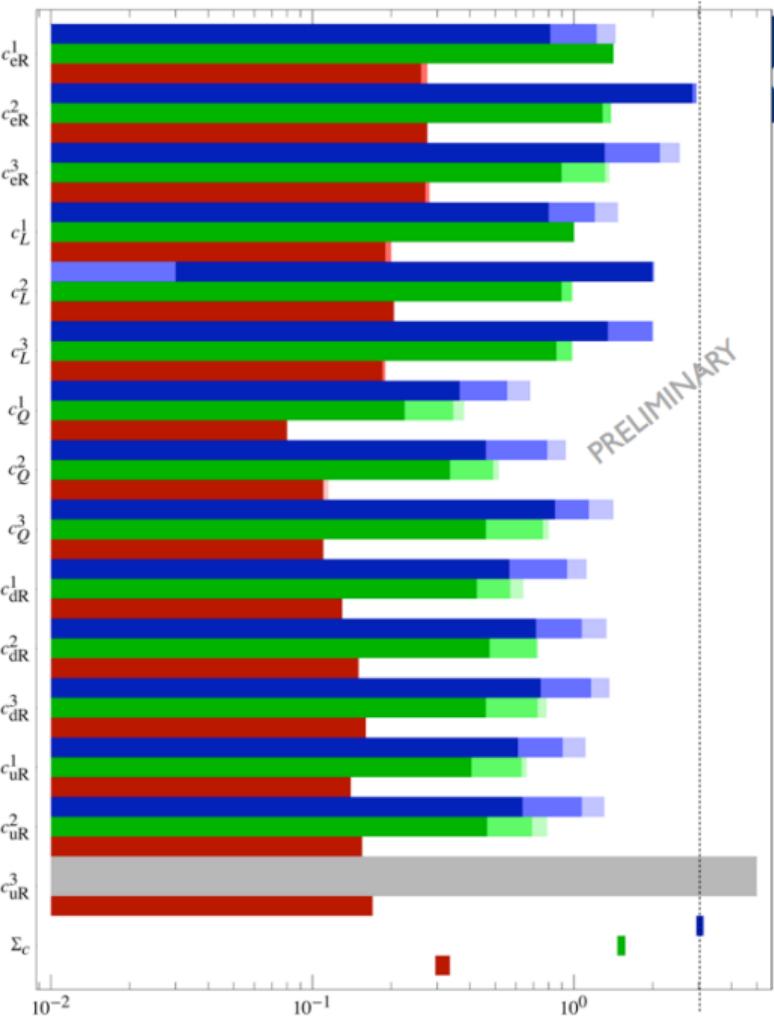
$$g_{Xe} \propto [(N_u + N_d)g_{(u,s)}^L + N_u g_u^R + N_d g_d^R]$$

$$N_u \approx N_d$$

$$g_{u,d}^L \approx -(g_u^R + g_d^R)/2.$$





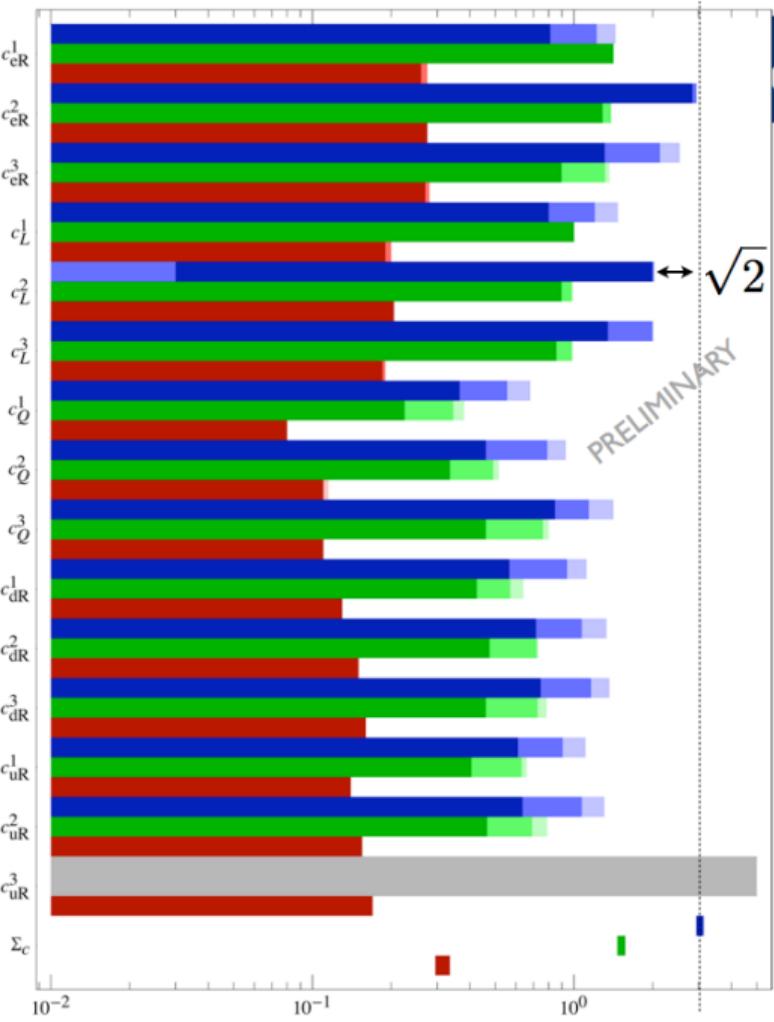


odel

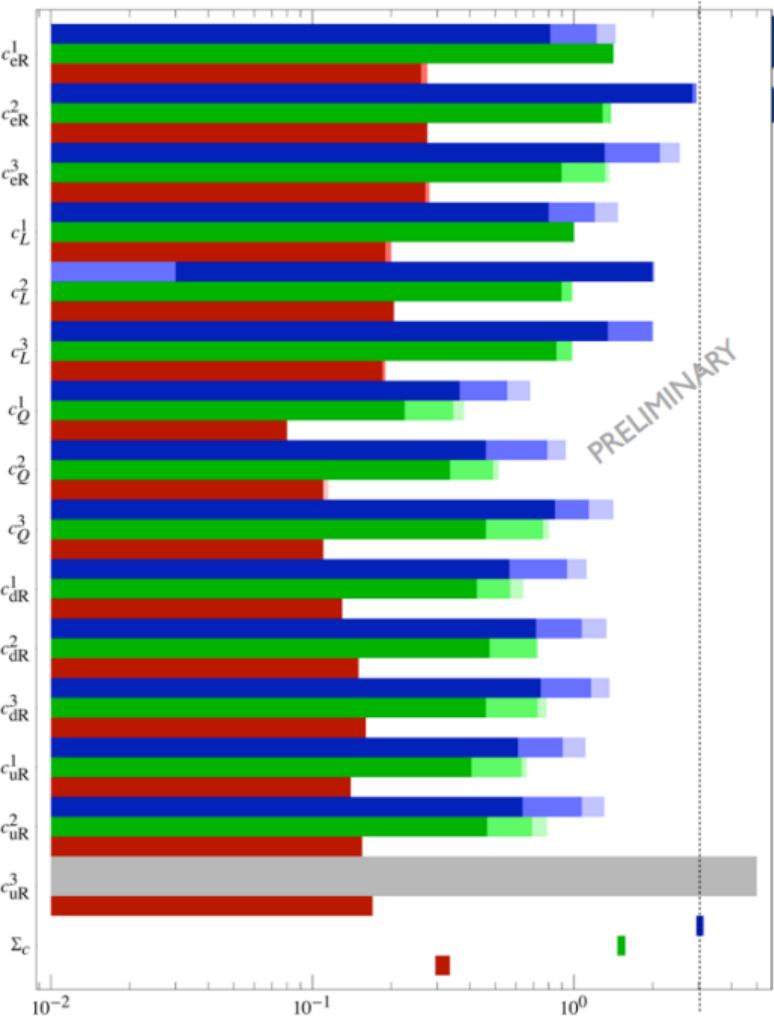
50 GeV

100 GeV

500 GeV



odel

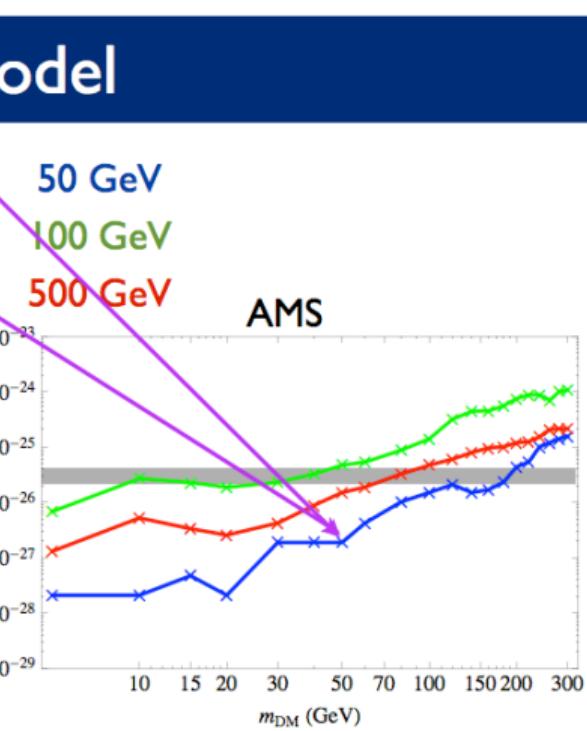
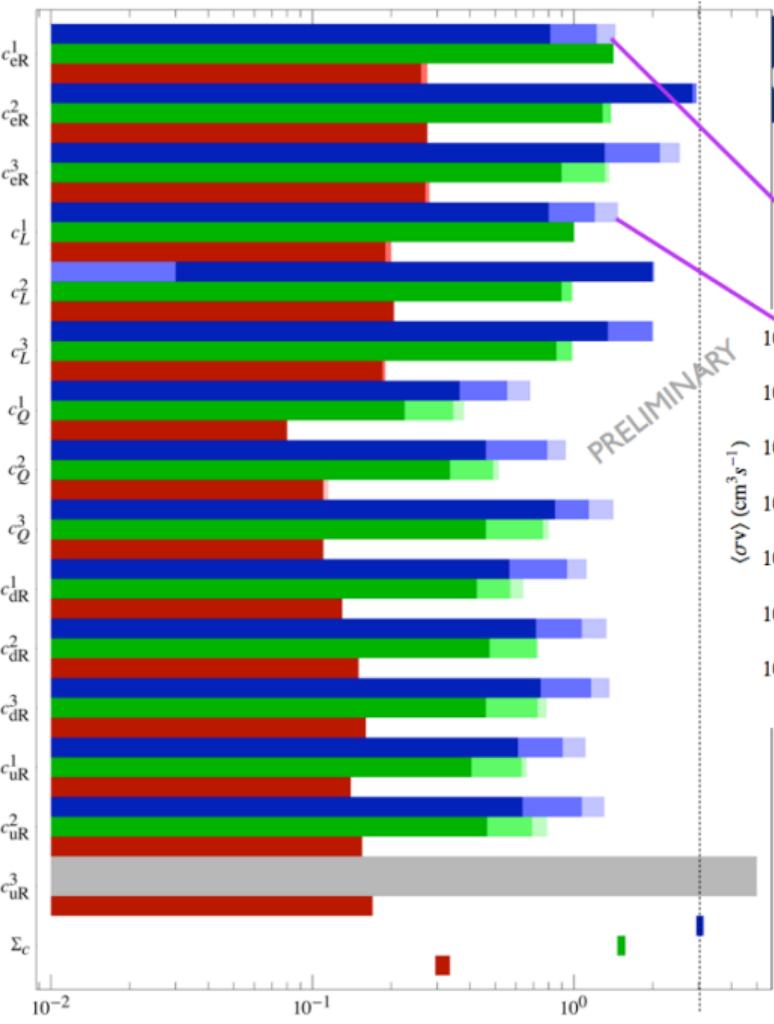


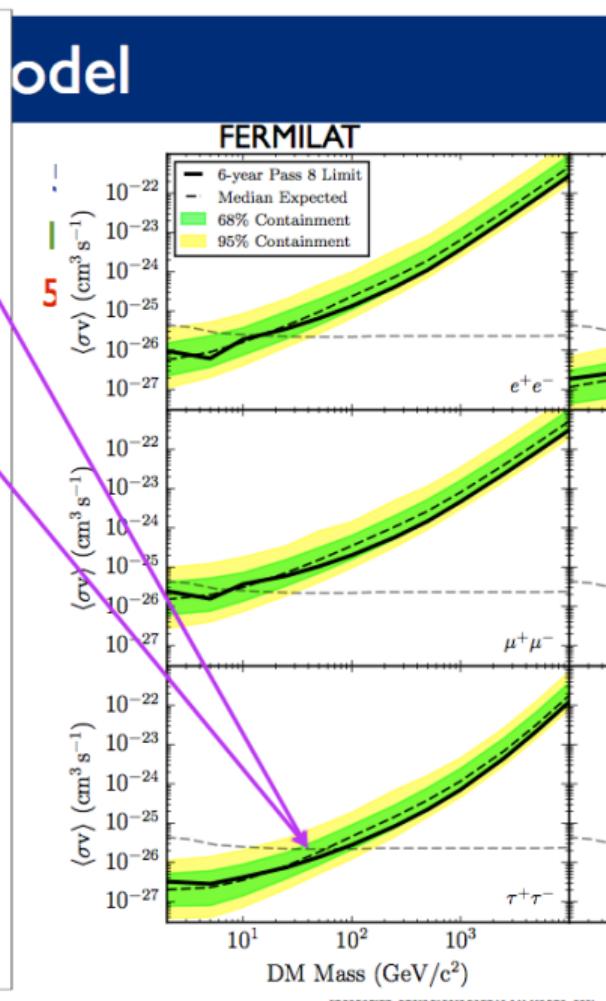
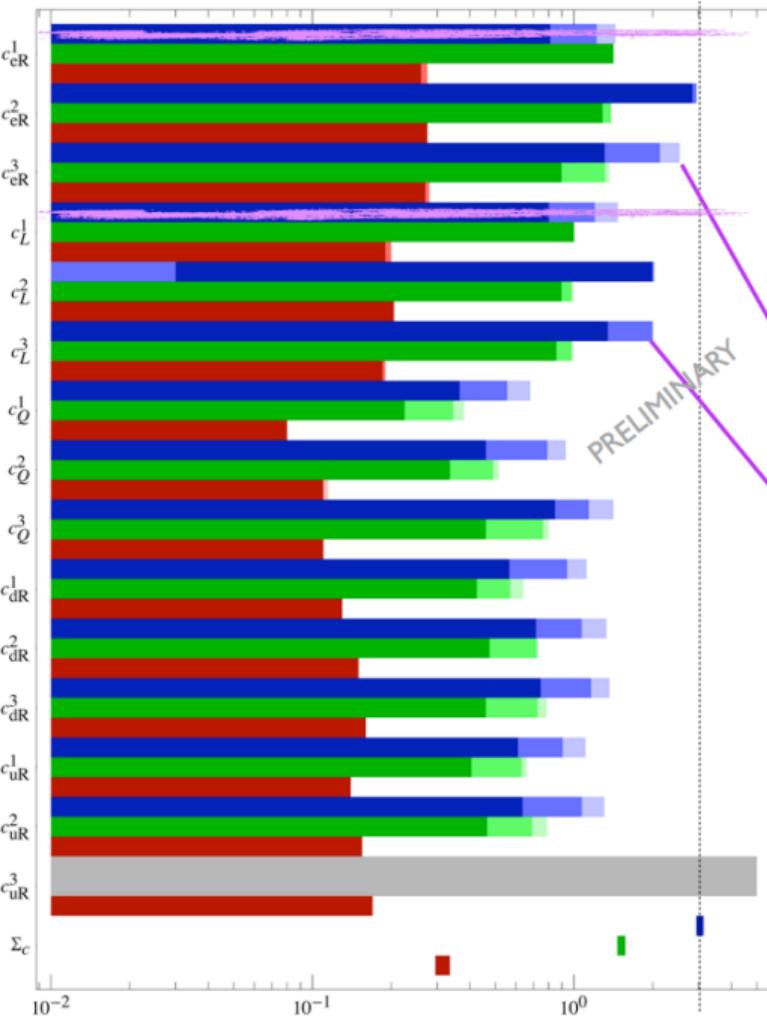
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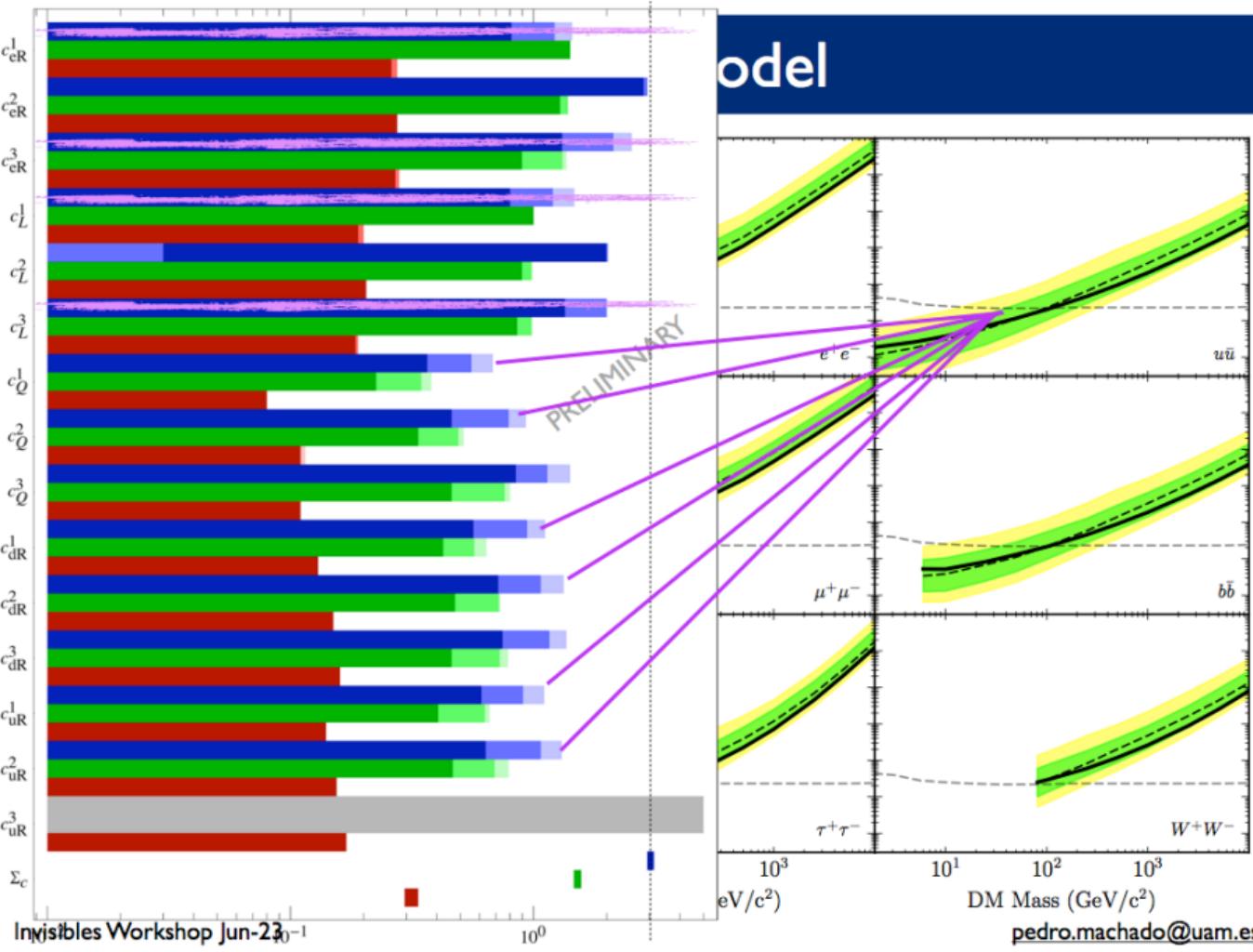
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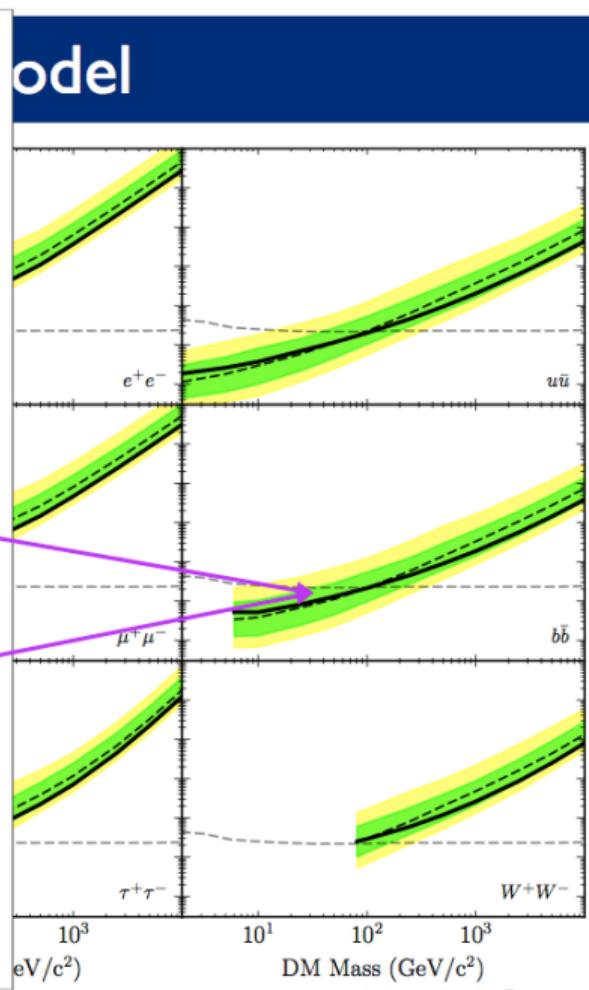
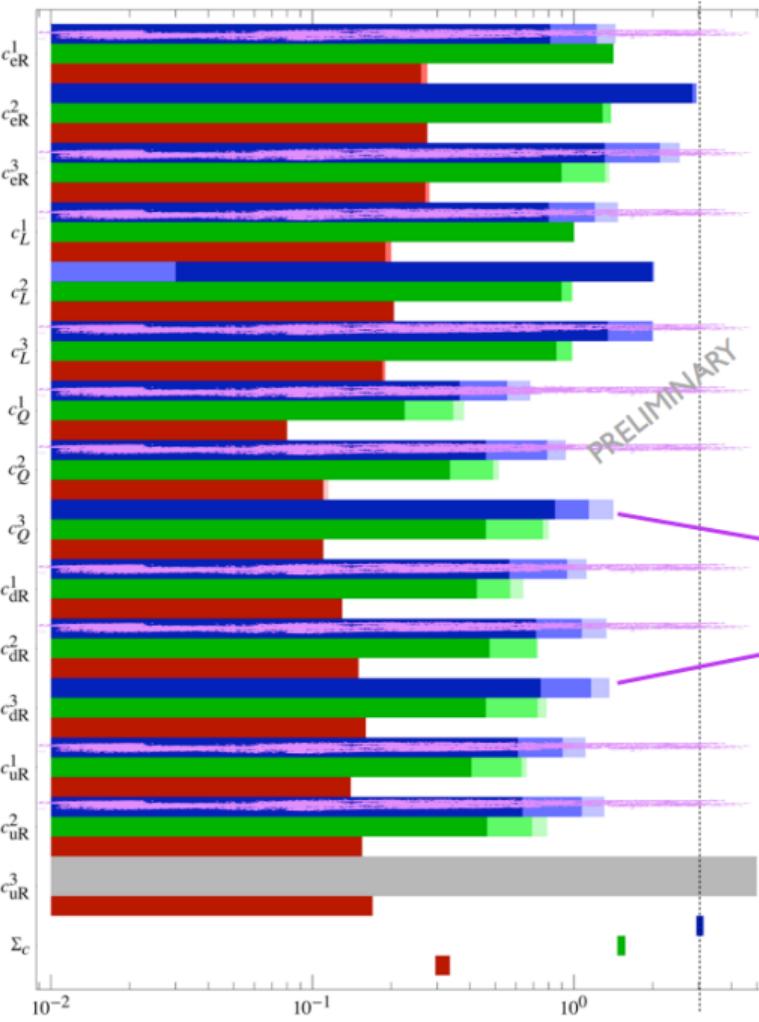
100 GeV

500 GeV

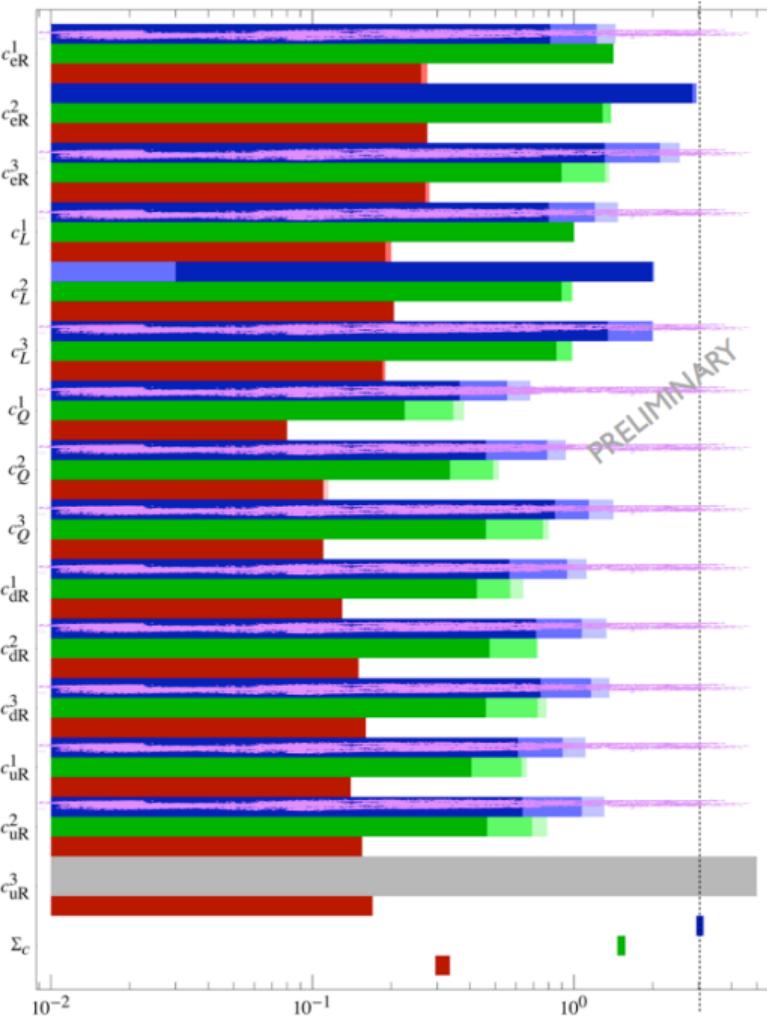








pedro.machado@uam.es



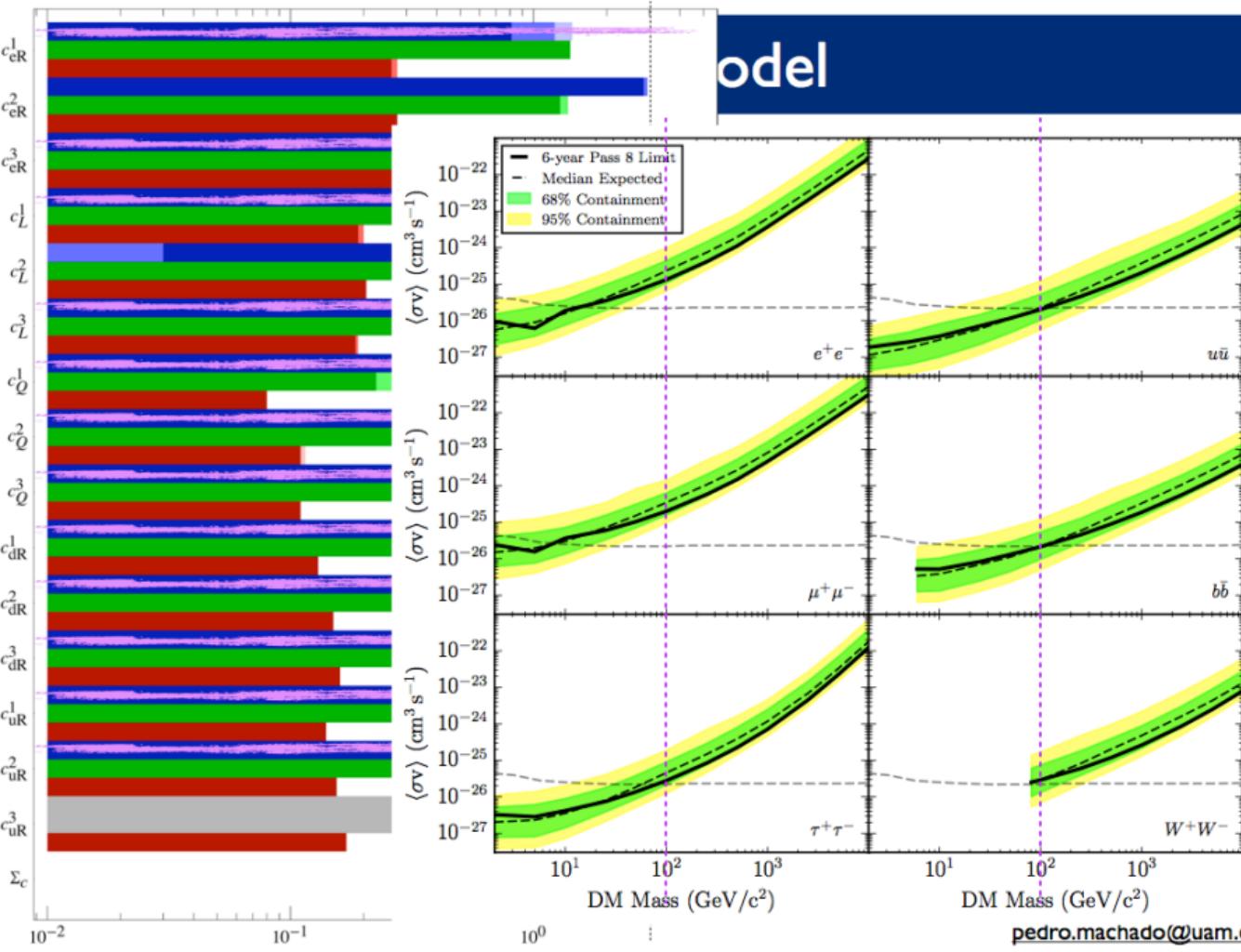
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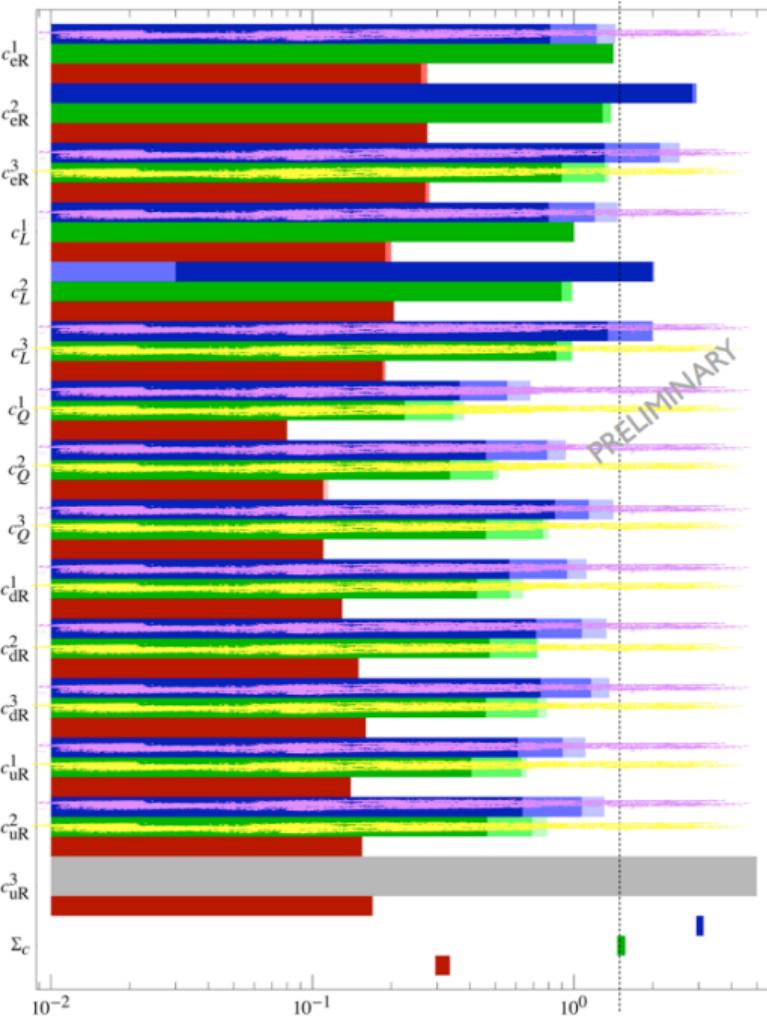
50 GeV

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500 GeV

In this scenario, for a 50 GeV DM, the right abundance is obtained by the coupling to muons and muon neutrinos!





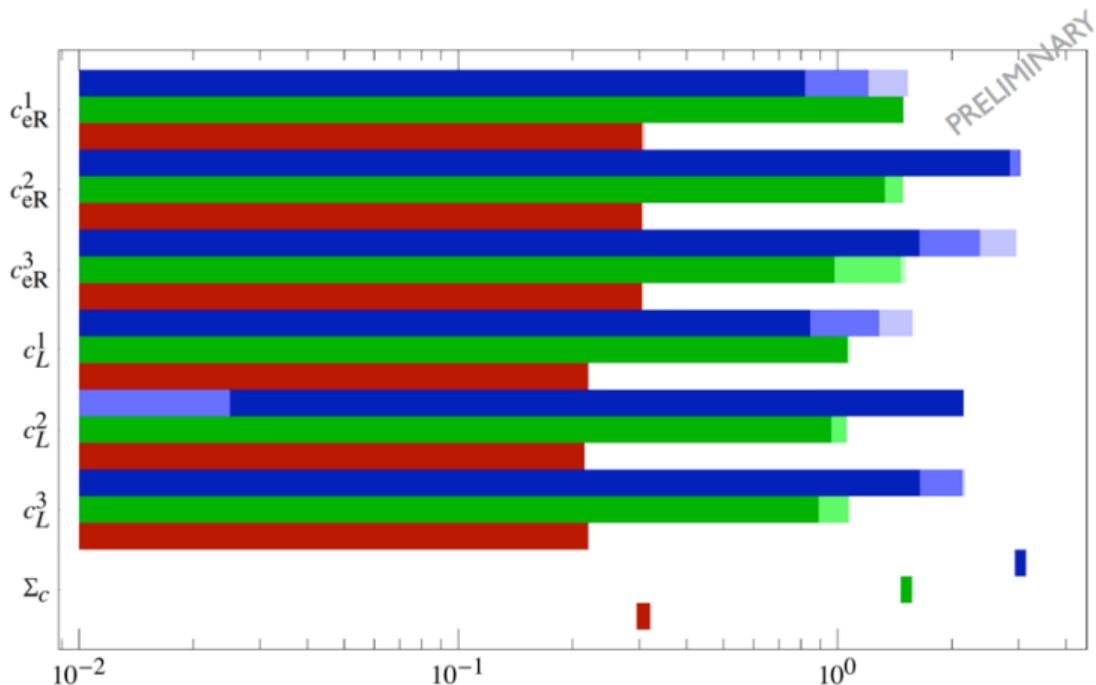
odel

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100 GeV

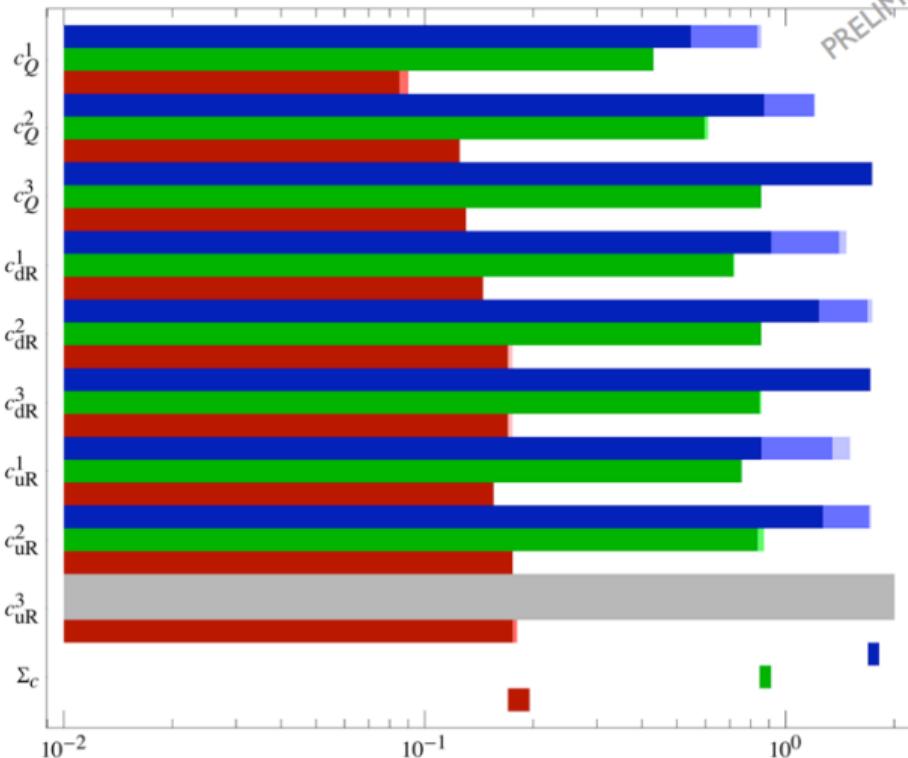
500 GeV

Leptophilic

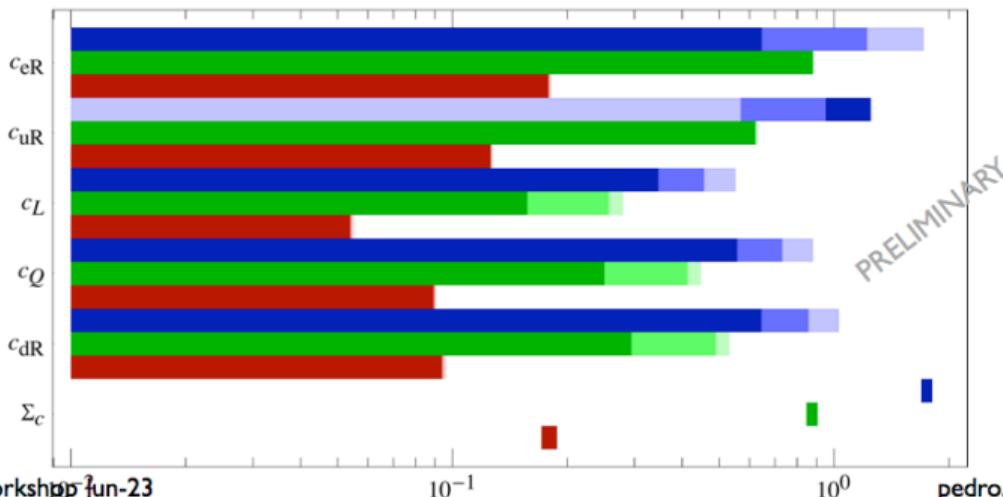
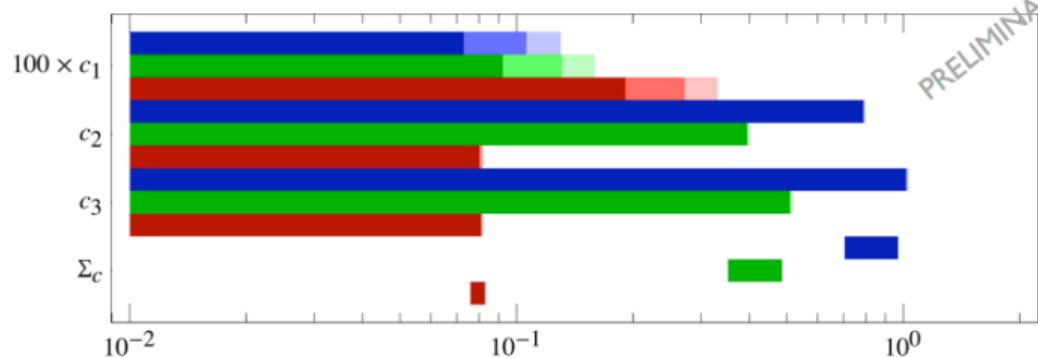


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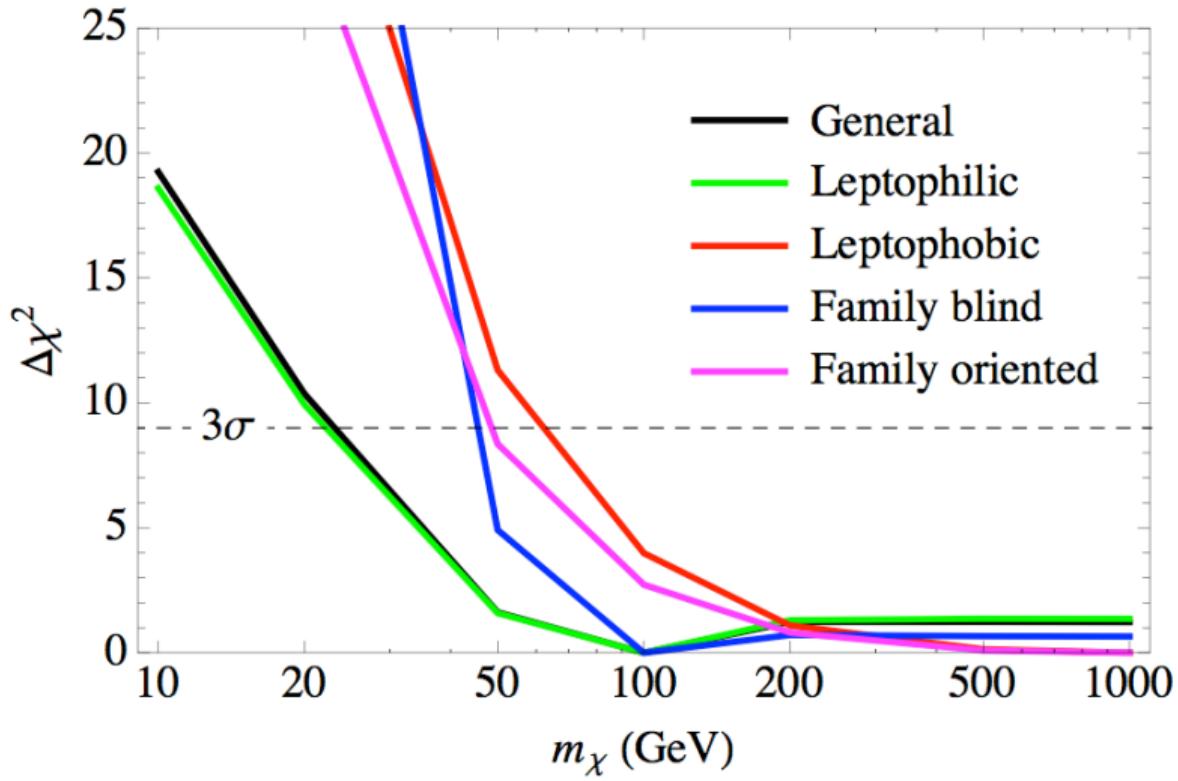
PRELIMINARY



Other models



Dependence with DM mass



Conclusions

We have analyzed the status of thermal WIMP dark matter
in the context of effective field theory

In some cases, specific couplings are required
to obtain the abundance

For the leptophilic and general models,
masses below 20 GeV are disfavored

For the other proxies,
masses below 50~60 GeV are disfavored

Backup

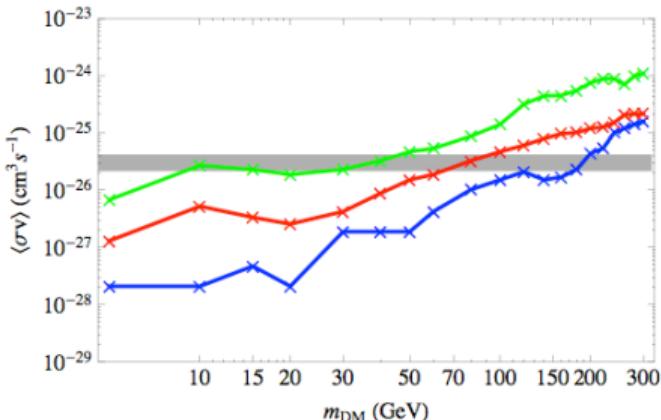
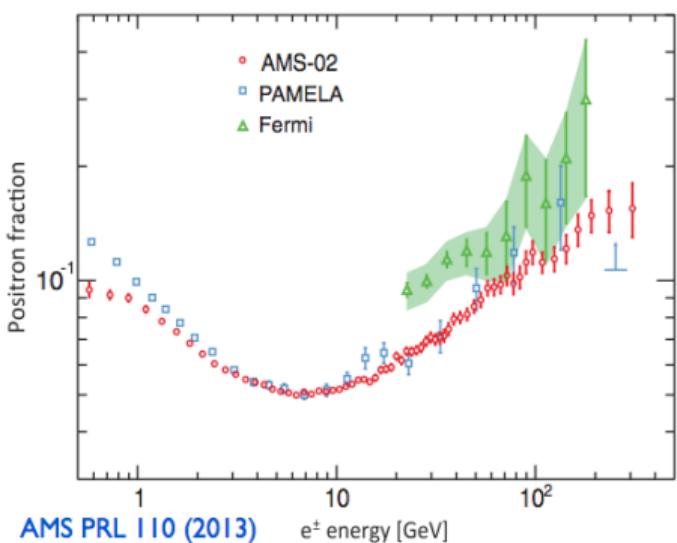
AMS

Backgrounds:

$$\Phi_{e^+} = C_{e^+} E^{-\gamma_{e^+}} + C_s E^{-\gamma_s} e^{-E/E_s};$$

$$\Phi_{e^-} = C_{e^-} E^{-\gamma_{e^-}} + C_s E^{-\gamma_s} e^{-E/E_s}$$

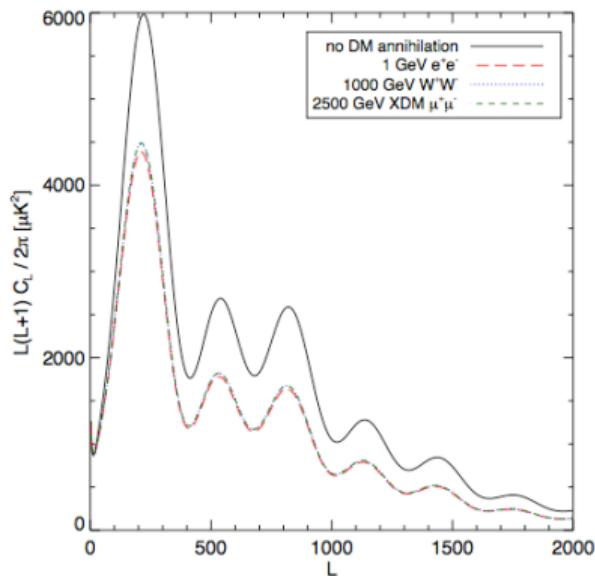
We marginalized the positron background, fixed the electron background, and fitted the DM signal.



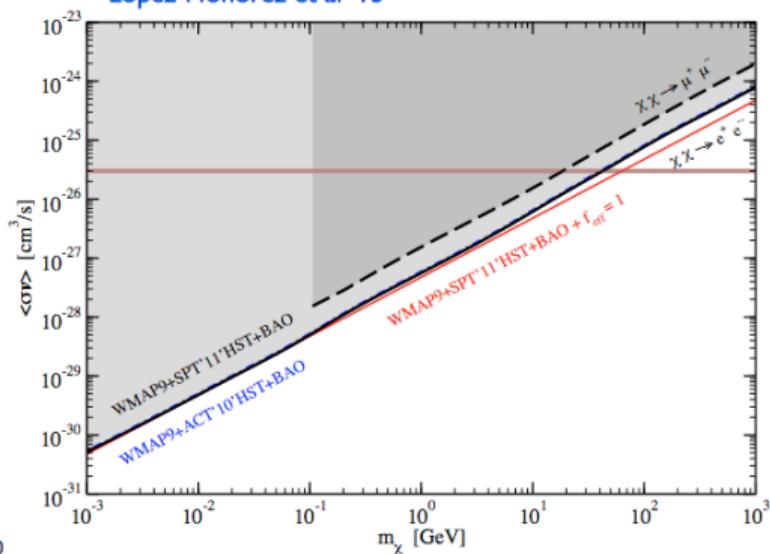
CMB

Dark matter annihilations can heat up and ionize the photon-baryon plasma, changing the CMB temperature and polarization angular power spectra

Slater Padmanabhan Finkbeiner '09

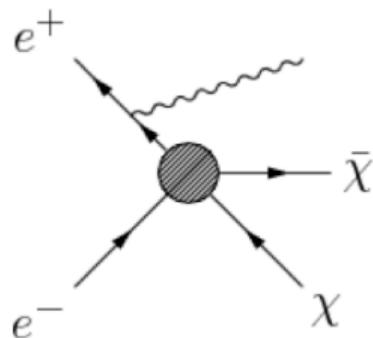
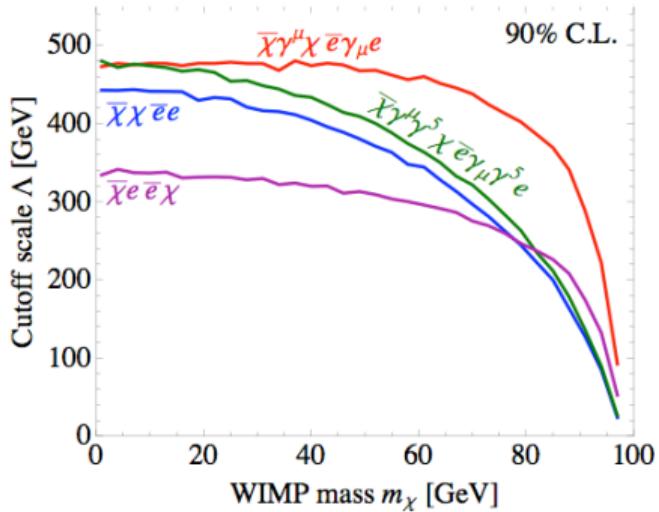


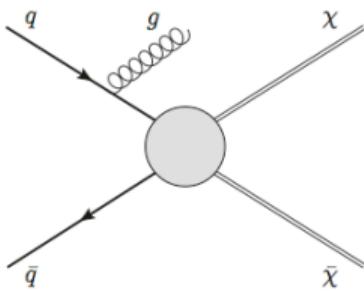
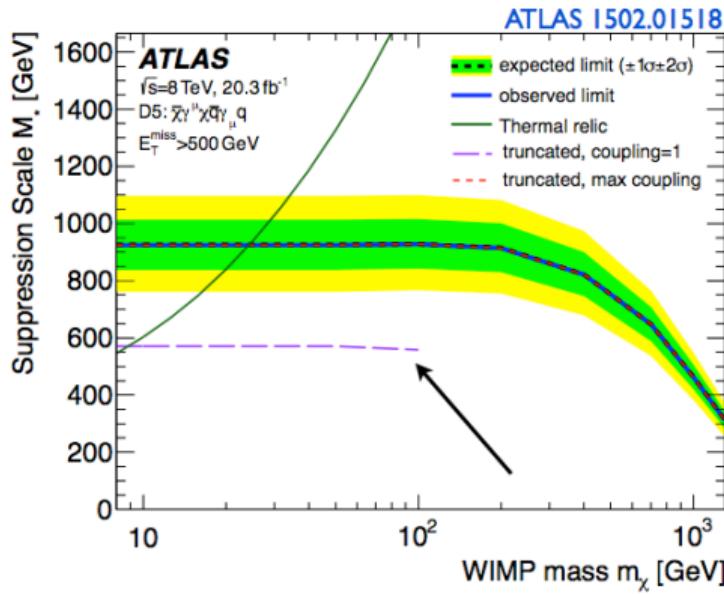
Lopez-Honorez et al '13



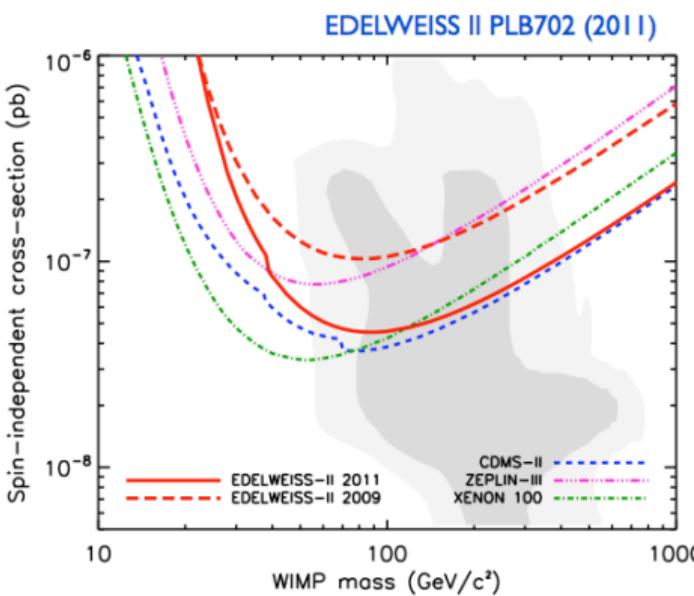
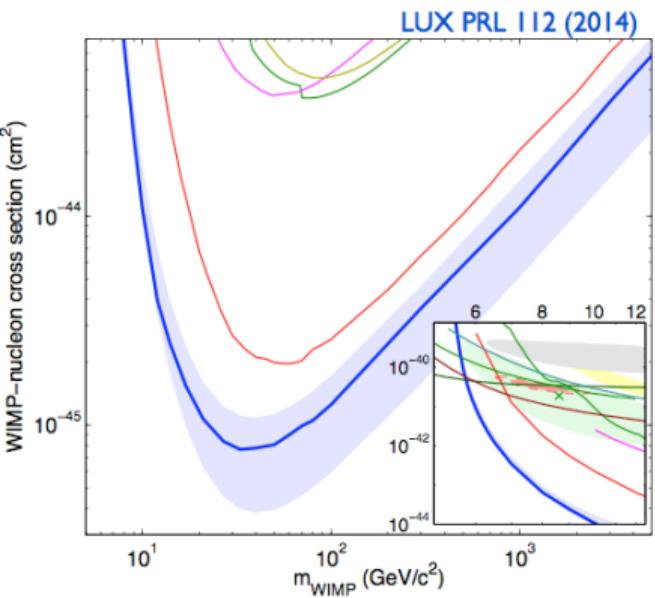
LEP

Fox Harnik Kopp Tsai PRD84 (2011)





Direct detection



FermiLAT

$$\phi_s(\Delta\Omega) = \underbrace{\frac{1}{4\pi} \frac{\langle\sigma v\rangle}{2m_{\text{DM}}^2} \int_{E_{\min}}^{E_{\max}} \frac{dN_\gamma}{dE_\gamma} dE_\gamma}_{\text{particle physics}} \times \underbrace{\int_{\Delta\Omega} \int_{\text{l.o.s.}} \rho_{\text{DM}}^2(\mathbf{r}) dl d\Omega'}_{\text{J-factor}}.$$

