



Carla Biggio
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Minimal muon anomalous magnetic moment

Based on JHEP 02 (2015) 099
in collaboration with M. Bordone

INVISIBLES15 Workshop
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The $(g-2)_\mu$ anomaly

$$a_\mu^{exp} = 116592080(63) \cdot 10^{-11}$$

Brookhaven 2006

$$a_\mu^{SM} = 116591790(65) \cdot 10^{-11}$$

Jegerlehner, Nyffeler 2009
($e+e^-$ annih.)

$$\Delta a_\mu = 290(90) \cdot 10^{-11}$$

3.1σ

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- **theoretical**: bad estimation of hadronic contribution
→ disfavoured Passera et al. 2008-09
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J-PARC
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← adopt this optimistic option :)

Explaining Δa_μ with a single new particle

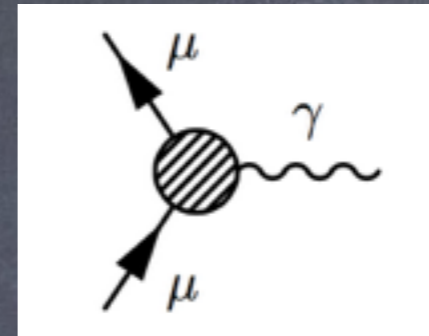
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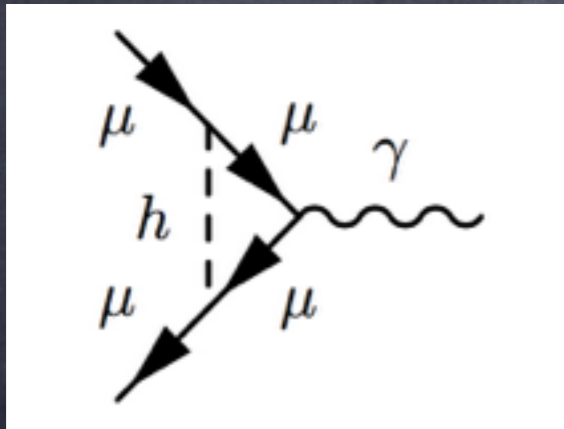
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In the SM (g-2) is a D=6 operator

$$\frac{1}{\Lambda^2} \bar{L} \sigma^{\mu\nu} e_R \phi F_{\mu\nu} + h.c.$$



At 1 loop, for example:

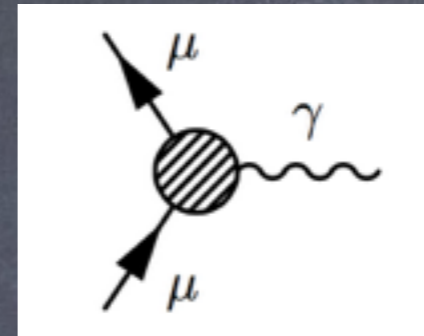


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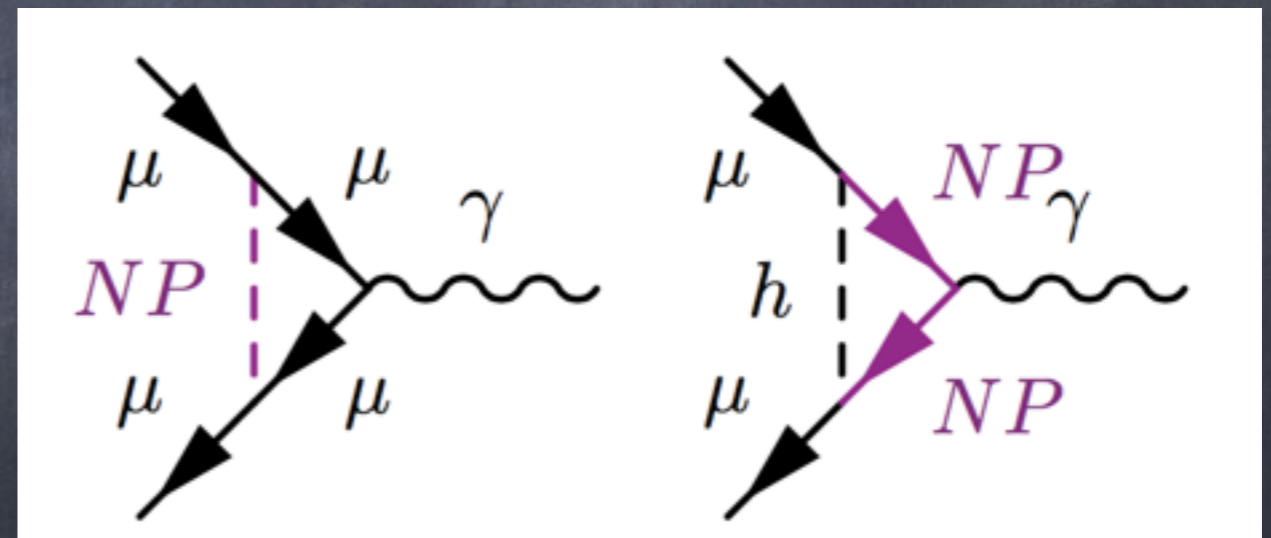
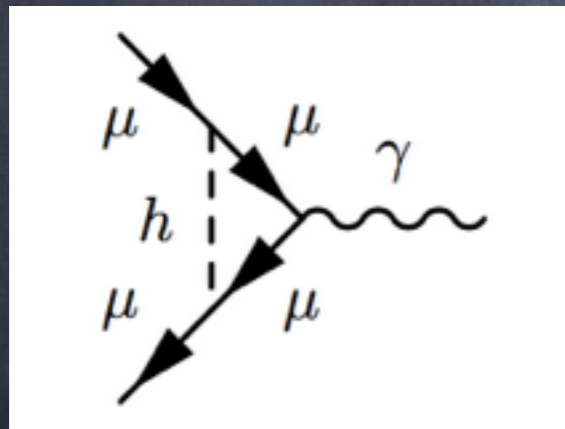
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- Requirements:
1. Lorentz invariance
 2. Gauge invariance
 3. Renormalizability

New fermions

All vector-like: masses indep. of EWSB, no anomalies

N_R	$(1, 1, 0)$	type I seesaw
Σ_R	$(1, 3, 1)$	type III seesaw
E_4	$(1, 1, -1)$	} 4 th generation.
L_4	$(1, 2, -1/2)$	
T	$(1, 3, -1)$	
D	$(1, 2, -3/2)$	

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but too small (EWPT)

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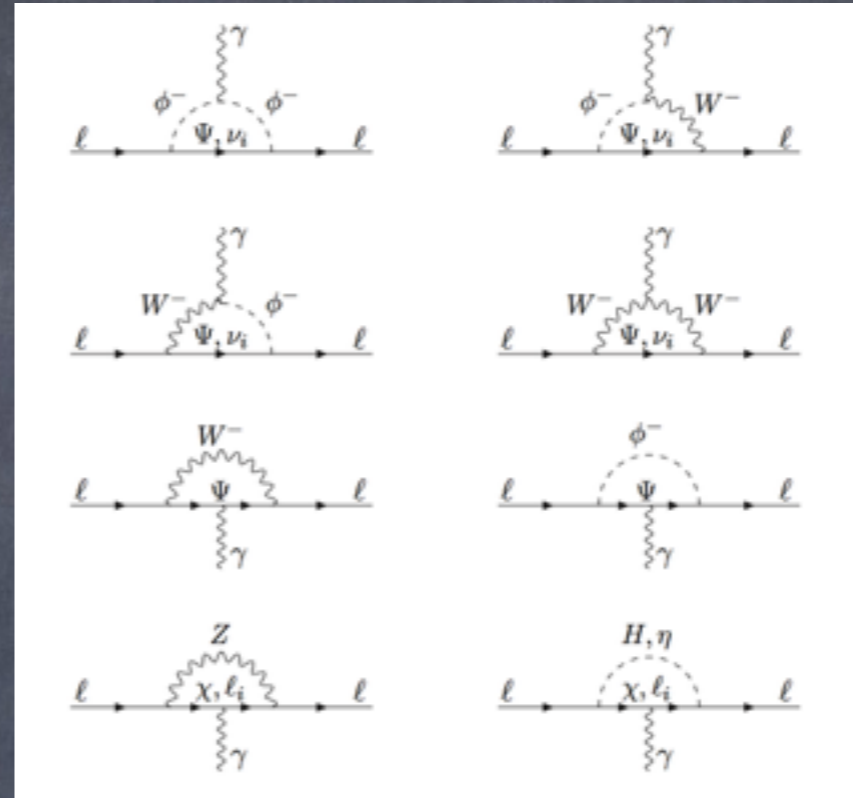
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The contribution of $D \sim (1,2,-3/2)$

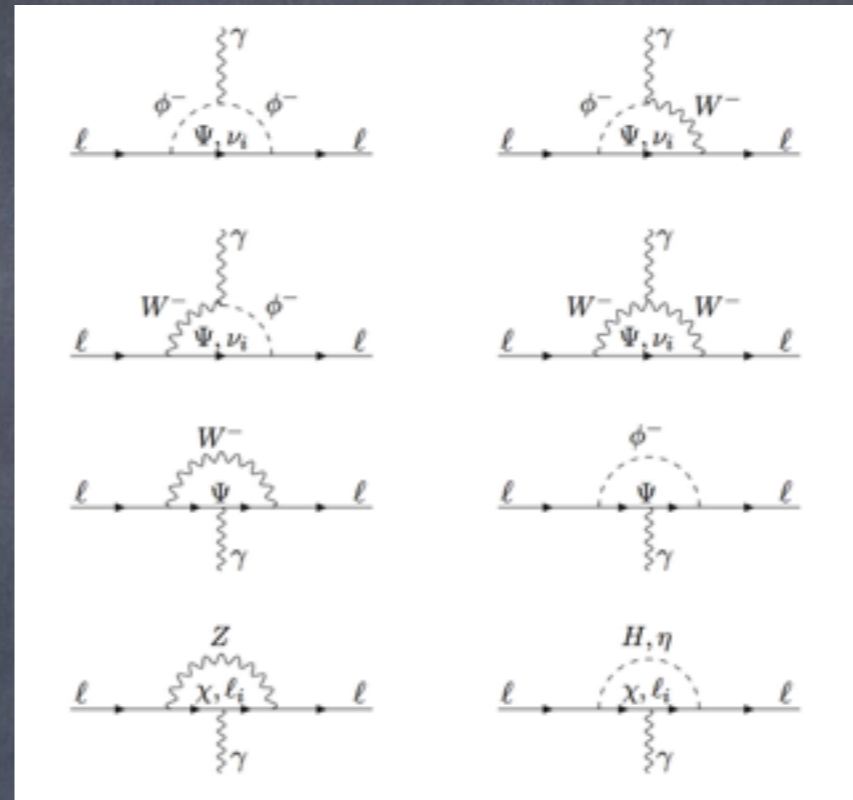
$$D = \begin{pmatrix} \chi \\ \Psi \end{pmatrix} \begin{matrix} \rightarrow Q=1 \\ \rightarrow Q=2 \end{matrix}$$



$$a_{\mu}^{\text{SM}+D} = \frac{m_{\mu}^2 G_F}{24\sqrt{2}\pi^2} \left\{ \overset{\text{SM}}{(3 - 4 \cos^2 \theta_W)^2 + 5} + \frac{v^2 |\lambda_{D\mu}|^2}{M_D^2} \left[-\frac{11}{4} - 4 \cos^2 \theta_W + \right. \right. \\ \left. \left. + F_{\text{NC}} \left(\frac{M_D^2}{M_Z^2} \right) + F_{\text{h}} \left(\frac{M_D^2}{M_H^2} \right) + F_{\text{CC}} \left(\frac{M_D^2}{M_W^2} \right) \right] \right\},$$

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CB, Bordone 14

The new contribution is negative
it cannot explain the discrepancy

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It's not possible to explain the discrepancy
adding to the SM a single fermion

New scalars

$$S_1 \quad (1, 1, 1)$$

$$S_2 \quad (1, 1, 2)$$

$$H_2 \quad (1, 2, 1/2)$$

II Higgs doublet

$$\Delta \quad (1, 3, 1)$$

type II seesaw

$$T_c^{1/3} \quad (3, 3, -1/3)$$

$$S_c^{1/3} \quad (3, 1, -1/3)$$

$$S_c^{4/3} \quad (3, 1, -4/3)$$

leptoquarks

$$D_c^{7/6} \quad (3, 2, 7/6)$$

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✗ Coaraza Perez et al. 95

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➤ enhancement with top (charm) mass

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In a particular 2HDM
if lighter than 200 GeV

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➤ enhancement with top (charm) mass

The \tilde{b}_R of the Higgsinoless MSSM

Riva, CB, Pomarol 12

It's a SUSY model without R-parity (with $U(1)_R$)
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In our case $l_\alpha \equiv \mu$, the b_R coupling is fixed to be Y_b

$$a_\mu^{LQ} \propto \frac{v^2 Y_b^2}{M_{LQ}^2}$$

we have a prediction for M_{LQ}

$$M_{\tilde{b}_R} \sim 500 \text{ GeV}$$

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ATLAS 13 Bounds from decay into $b\nu$:

Br=1	$M > 620 \text{ GeV}$
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In our model: $\tilde{b}_R \rightarrow$

$b_L \nu_L$	same BR
$t_L l_L$	
$(b_R \tilde{G})$	

This possibility is viable, to confirm/exclude it

look for final states with top and charged leptons!!!

Conclusions

- We have considered single particle extensions of the SM (scalar & fermion)
- A single new fermion cannot explain the $(g-2)_\mu$ anomaly
- Only 3 scalars - 2 leptoquarks and a second Higgs doublet- can do it
- The \tilde{b}_R of the Higgsinoless MSSM could solve the $(g-2)_\mu$ puzzle and we have a prediction for its mass: $M_{\tilde{b}_R} \sim 500 \text{ GeV}$
- Most of these solutions are going to be tested @ LHC13
- Wait for new LHC run and new $(g-2)_\mu$ experiment!

Muchas gracias!