

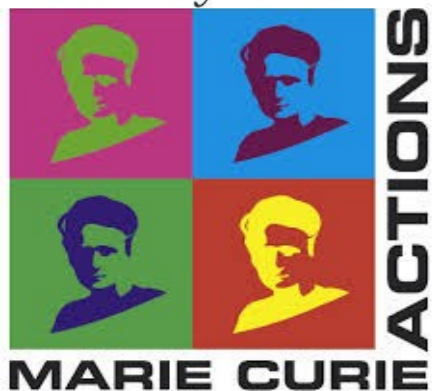


# More Higgses at LHC: The EW Road to Baryogenesis

Jose Miguel No, Sussex U.

Based on 1405.5537 (PRL), 1507.xxxxx with G. Dorsch, S. Huber, K. Mimasu.

US  
University of Sussex



InVisibles meets Visibles, 22<sup>nd</sup> June 2015



# The SM is NOT the End of the Story...

⇒ EW Hierarchy Problem...

Expect New  
Physics@TeV

⇒ Dark Matter

⇒ Neutrino Masses

⇒ **Matter-Antimatter Asymmetry** → **Baryogenesis**

$$\frac{n_B - n_{\bar{B}}}{n_\gamma} \sim 10^{-9}$$

*What is the Origin of the Baryon Asymmetry?*

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What is the Origin of the Baryon Asymmetry?

SAKHAROV CONDITIONS (for dynamical generation of baryon asymmetry)

B Violation ✓ **Sphalerons**

V. A. Kuzmin, V. A. Rubakov, M. Shaposhnikov, Phys. Lett. B155 (1985) 36

C/CP Violation **CKM**

Departure from Thermal Equilibrium **EW Phase Transition**





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C/CP Violation ✗ **not enough**

M. B. Gavela, P. Hernandez, J. Orloff, O. Pene, C. Quimbay, Nucl. Phys. **B430** (1994) 382

Departure from Thermal Equilibrium ✗ **not enough**

In the SM ( $m_h = 125$  GeV) EW Phase Transition Smooth CrossOver

K. Kajantie, M. Laine, K. Rummukainen, M. Shaposhnikov, Phys. Rev. Lett. **77** (1996) 2887



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C/CP Violation ?

Departure from Thermal Equilibrium ?

New CP Sources **EDMs**

EW Phase Transition **LHC**

**BSM**



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EW Phase Transition **LHC**

**BSM**

**LHC Signatures Revealing 1<sup>st</sup> Order EW Phase Transition**

# A 30s **CRASH** Course on the EW Phase Transition

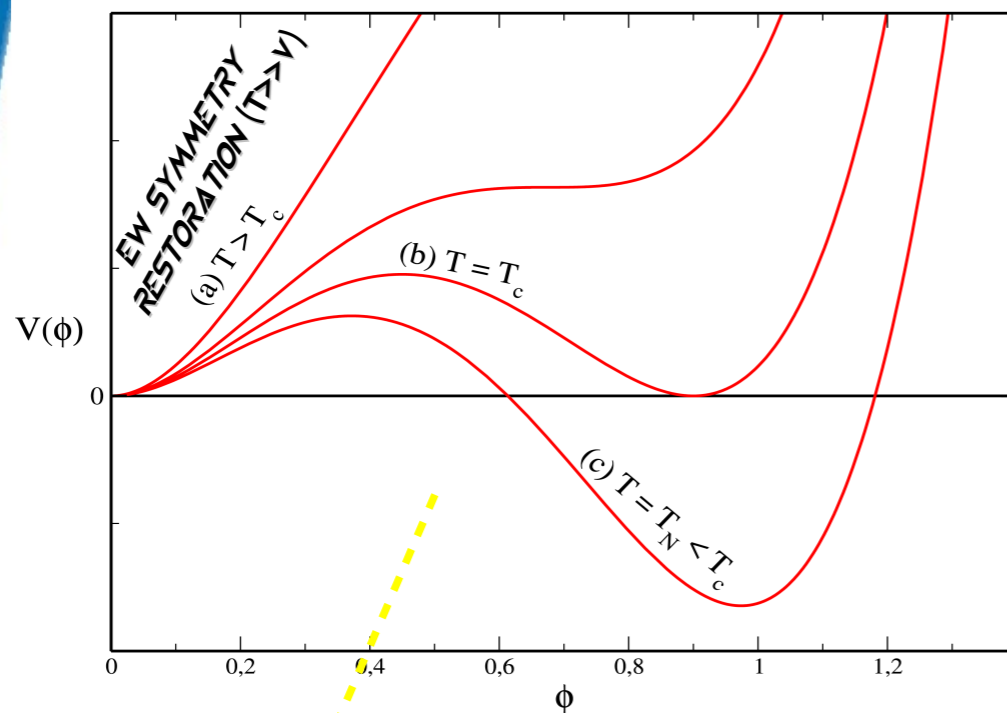
Universe Expands Adiabatically  $\Rightarrow$  Equilibrium Thermal Field Theory

Finite-T Effective Potential  $V(h, T)$  for the Higgs

$$V(h, T) \approx (a T^2 - \mu^2) h^2 - E(T) h^3 + \lambda(T) h^4$$

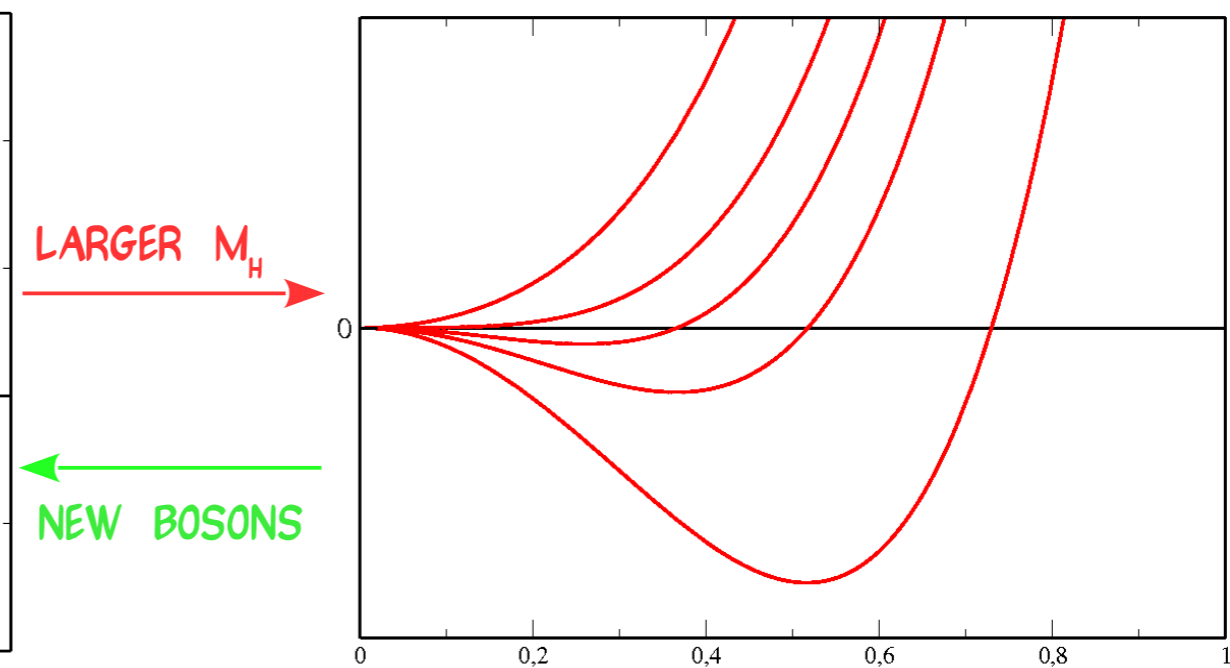
1<sup>st</sup> Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$  Discontinuous

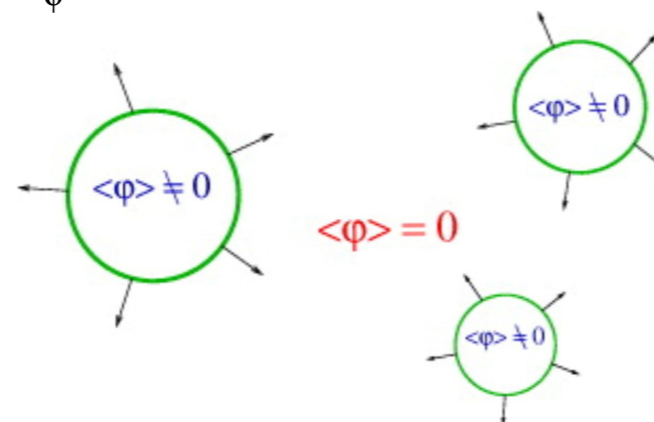


2<sup>nd</sup> Order:

$\langle \phi \rangle = 0 \rightarrow \langle \phi \rangle = \phi(T)$  Continuous



Bubble Nucleation & Growth



SM EW Phase Transition Smooth CrossOver for  $m_h > 80$  GeV

K. Kajantie, M. Laine, K. Rummukainen, M. Shaposhnikov, *Phys. Rev. Lett.* **77** (1996) 2887

# EW Scale Baryogenesis Needs:

→ *New Bosons (EW Scale)*

→ *Coupled to SM Higgs*

Strong 1<sup>st</sup> Order  
EW Phase Transition

Archetype Scenario: *Extended Higgs Sectors*

- SIMPLE EXTENSIONS OF THE SM
- PROVIDE MISSING INGREDIENTS FOR EW BARYOGENESIS

**More Higgses!**

**2HDM** (Add a Second Scalar Doublet to the SM)

Goal: *LHC signals of EW Phase Transition*



# 2HDM

$$\begin{aligned}
 V_s(\Phi_1, \Phi_2) = & -\mu_1^2 \Phi_1^\dagger \Phi_1 - \mu_2^2 \Phi_2^\dagger \Phi_2 - \frac{\mu^2}{2} (\Phi_1^\dagger \Phi_2 + h.c.) \\
 & + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) \\
 & + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + \frac{\lambda_5}{2} \left( (\Phi_1^\dagger \Phi_2)^2 + h.c. \right)
 \end{aligned}$$

$\mathbb{Z}_2$  Symmetric  
(softly broken)

(For simplicity, we do not consider CP Violation)

$$\begin{aligned}
 \Phi_1 = \begin{pmatrix} \varphi_1^+ \\ \frac{v_1 + h_1 + i\eta_1}{\sqrt{2}} \end{pmatrix} \quad \Phi_2 = \begin{pmatrix} \varphi_2^+ \\ \frac{v_2 + h_2 + i\eta_2}{\sqrt{2}} \end{pmatrix} \quad H^\pm = -\sin\beta \varphi_1^\pm + \cos\beta \varphi_2^\pm \\
 A_0 = -\sin\beta \eta_1 + \cos\beta \eta_2 \quad h = \cos\alpha h_1 + \sin\alpha h_2 \\
 H_0 = -\sin\alpha h_1 + \cos\alpha h_2
 \end{aligned}$$

If  $\alpha = \beta$   
h is SM Higgs

- New "Heavy" Scalars  $H_0$  (CP-Even),  $A_0$  (CP-Odd) and  $H^\pm$
- 6 New Parameters  $m_{H_0}$   $m_{A_0}$   $m_{H^\pm}$   $\mu$   $\alpha$   $\tan\beta$

We Focus on Type I 2HDM (all fermions coupled to same scalar doublet)

⇒ EW PHASE TRANSITION DOES NOT DEPEND ON THE TYPE

⇒ EXPERIMENTAL CONSTRAINTS DO DEPEND ON THE TYPE

# EW Phase Transition in 2HDM

→ Scan in  $m_{H_0}$   $m_{A_0}$   $m_{H^\pm}$   $\mu$   $\alpha$   $\tan\beta$

*G. Dorsch, S. Huber, J. M. No, JHEP **1310** (2013) 029*

⇒ Stability of the Effective Potential

⇒ Interface to 2HDMC & HiggsBounds

*D. Eriksson, J. Rathsman, O. Stal, Comput. Phys. Commun. **181** (2010) 189*

*P. Bechtle, O. Brein, S. Heinemeyer, G. Weiglein, K. Williams, Comput. Phys. Commun. **181** (2010) 138*

⇒ Flavour Constraints (mainly  $b \rightarrow s \gamma$ )

⇒ Global Fit to light Higgs Properties

*C. Chen, S. Dawson, M. Sher, Phys. Rev D **88** (2013) 015018*

*G. Belanger, D. Dumont, U. Ellwanger, J. F. Gunion, S. Kraml, Phys. Rev D **88** (2013) 075008*

*N. Craig, F. D'Eramo, P. Draper, S. Thomas, H. Zhang, 1504.04630*

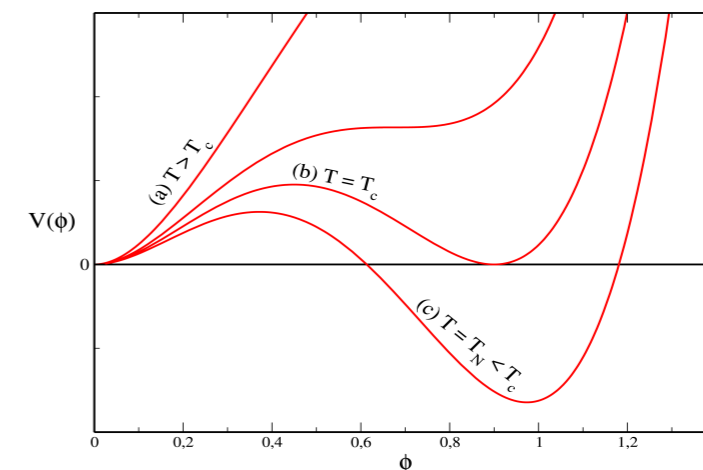
Unitarity, Perturbativity, EWPO,  
LEP/Tevatron/LHC Bounds

Constraints on  $\beta$ - $\alpha$  and  $\tan\beta$

Points satisfying all above constraints are “Physical”

→ Strength of the EW Phase Transition:

⇒ Daisy Resummed 1-loop  
Thermal Eff. Potential  $V_{\text{eff}}(\phi, T)$

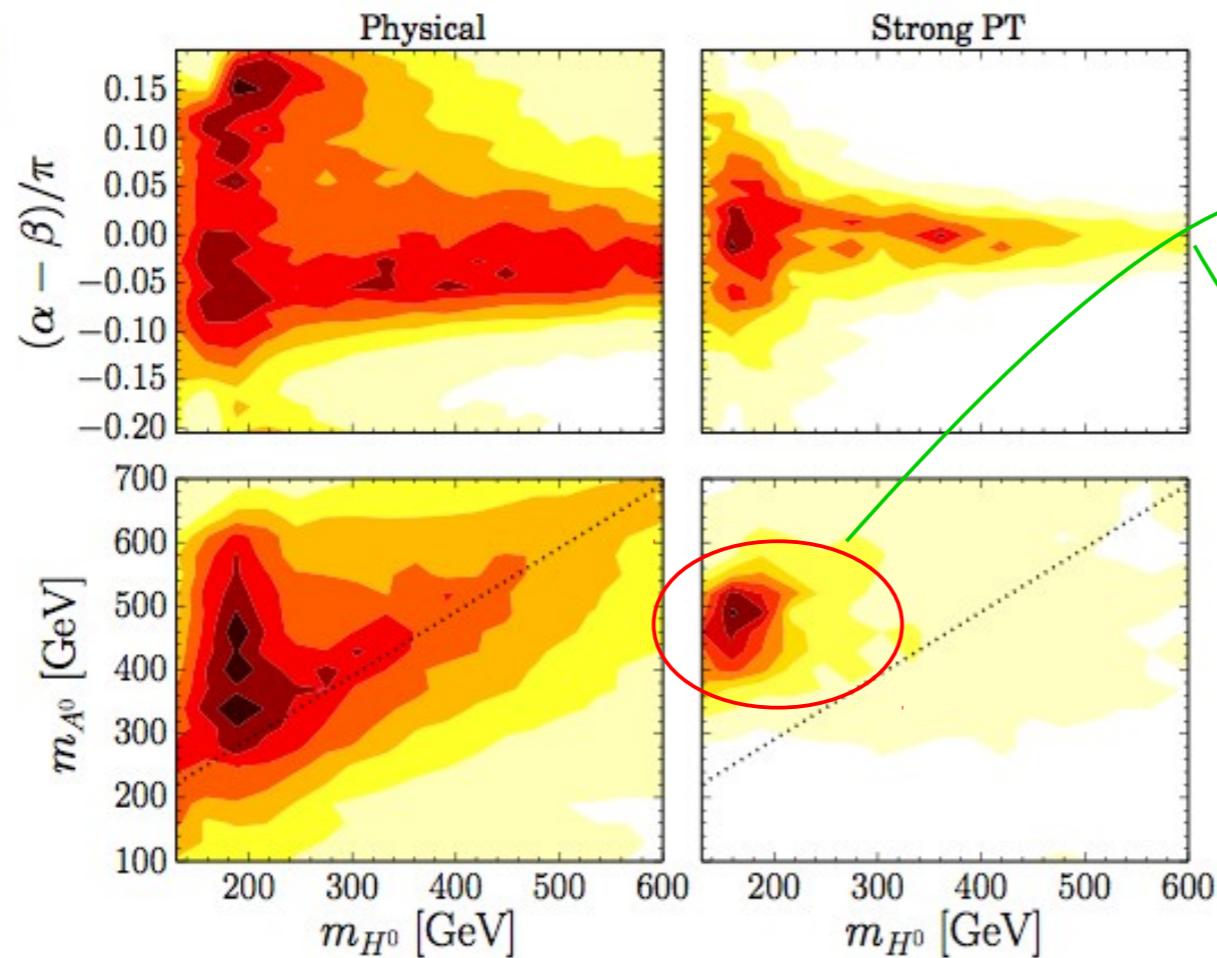




# EW Phase Transition in 2HDM

## Strong EW Phase Transition (vs Physically Allowed)

G. Dorsch, S. Huber, K. Mimasu, J. M. No, *Phys. Rev. Lett.* **113** (2014) 211802



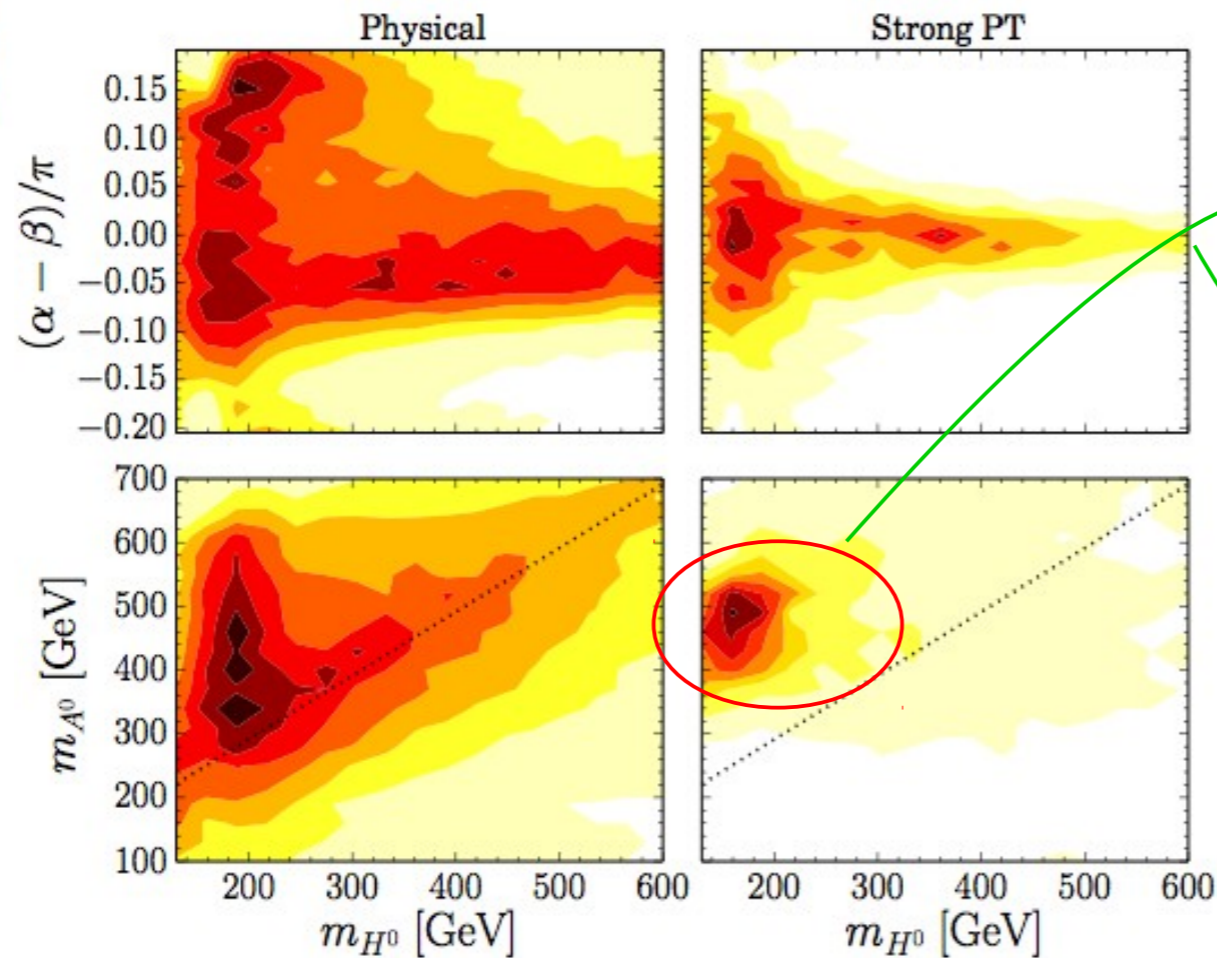
- $m_{A_0} - m_{H_0} \sim v$   
(&  $m_{A_0} > 300$  GeV)
- $m_{H_0} < 250$  GeV  
(Small  $\alpha - \beta$  and  $\tan\beta \gtrsim 1$ )



# EW Phase Transition in 2HDM

## Strong EW Phase Transition (vs Physically Allowed)

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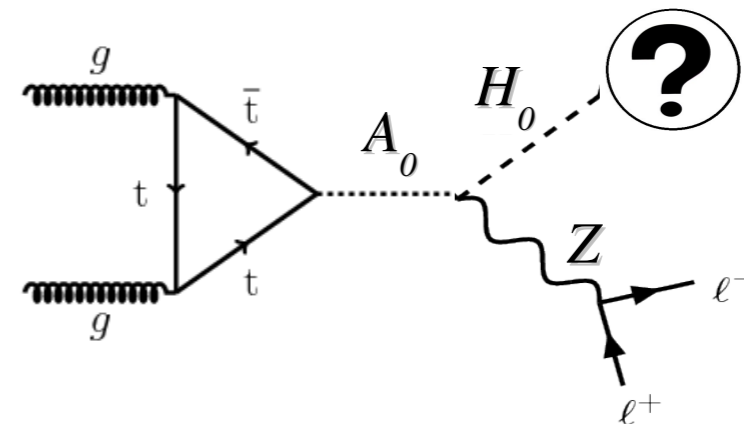
(Small  $\alpha - \beta$  and  $\tan\beta \gtrsim 1$ )

## Impact on 2HDM Searches at LHC!

*New Decay Channels*  $\phi_i \rightarrow V \phi_j$   
(not widely considered: Not Accessible in MSSM)

EW PHASE TRANSITION SIGNATURE

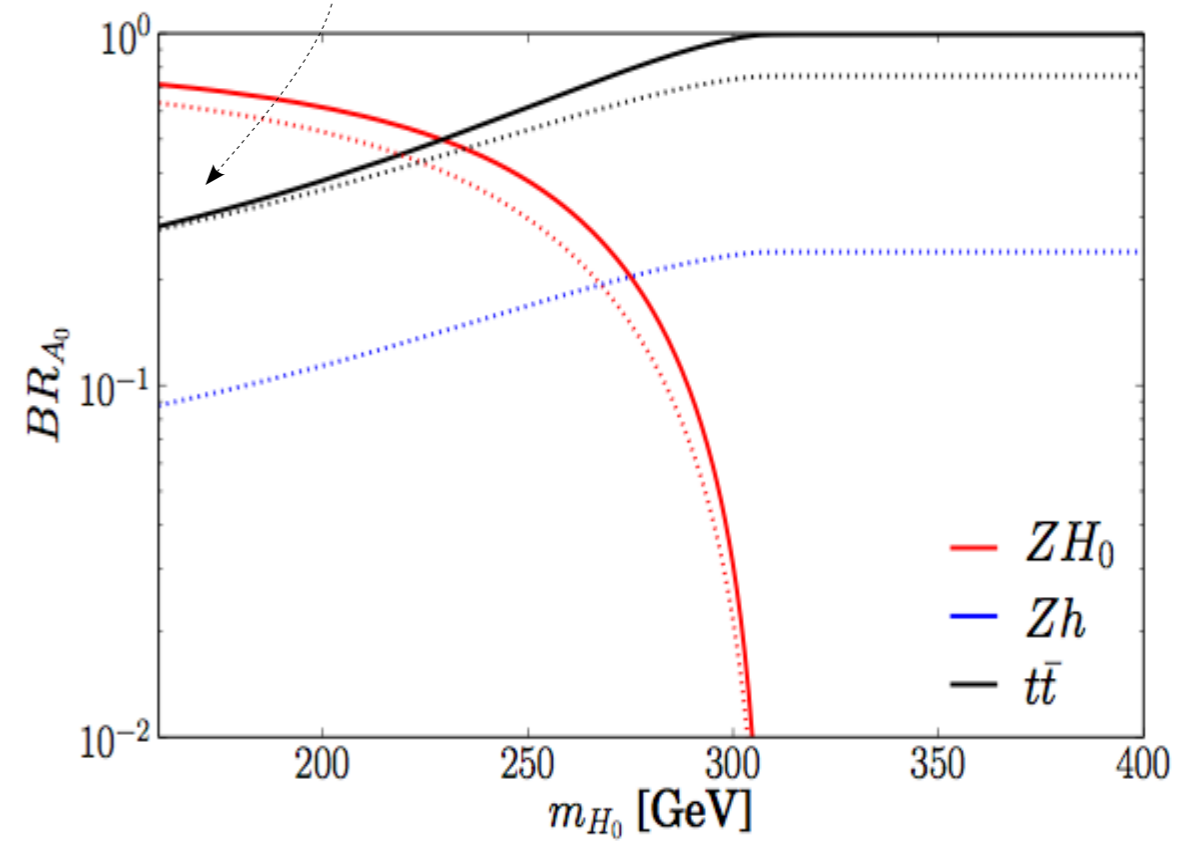
$$A_0 \rightarrow H_0 Z$$



$\bar{b}b$   
or  
 $WW, ZZ$

# EW Phase Transition in 2HDM@LHC

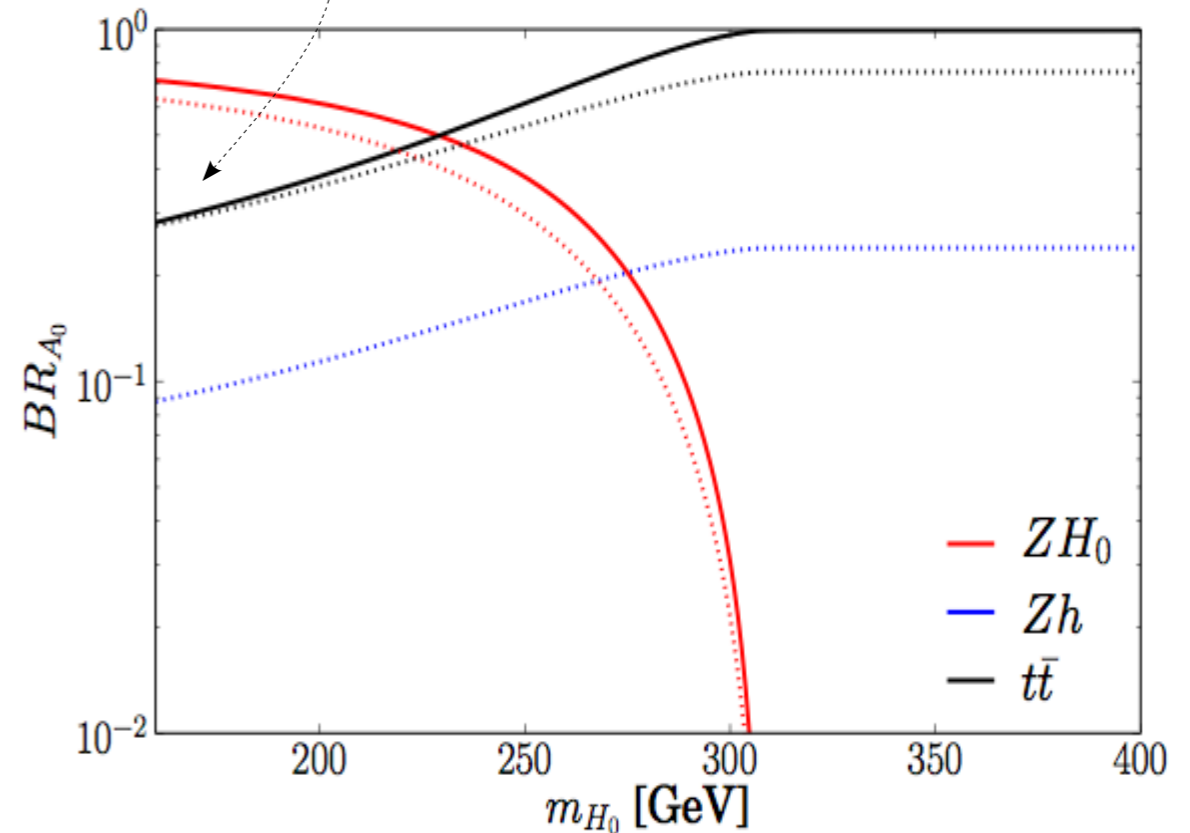
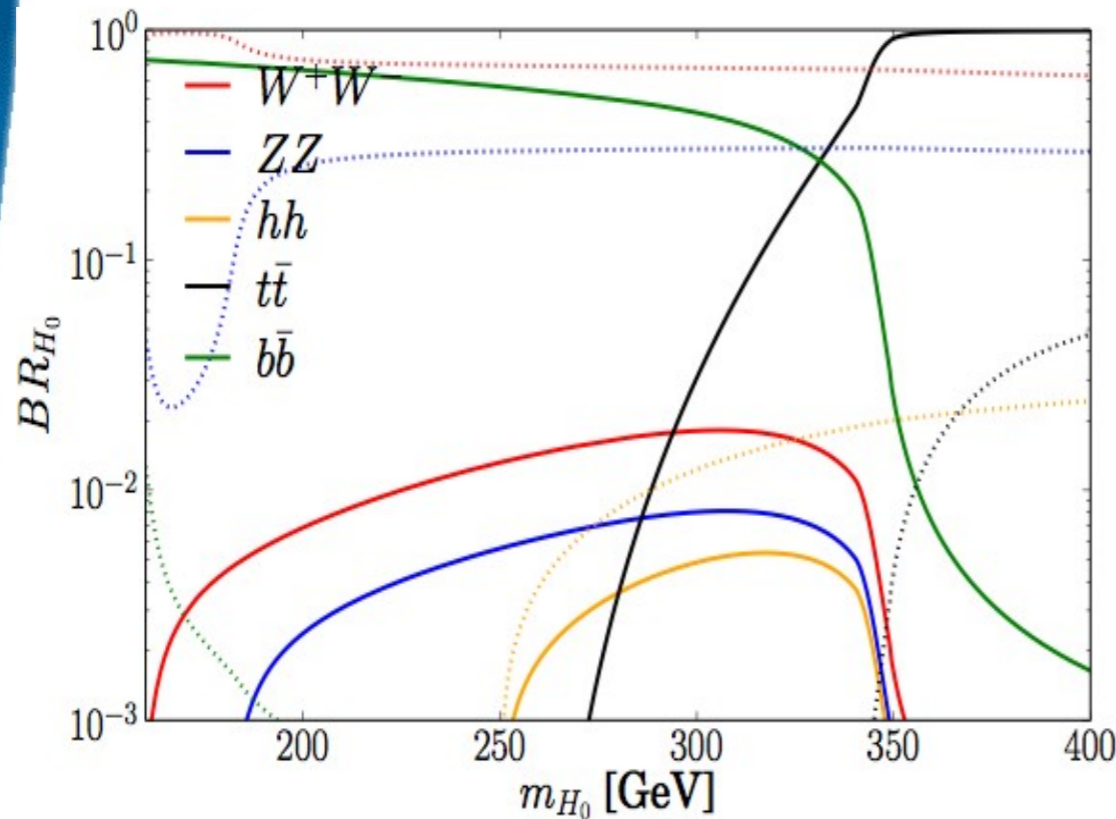
- Decay  $A_0 \rightarrow H_0 Z$  Dominant for  $m_{A_0} - m_{H_0} \sim v$





# EW Phase Transition in 2HDM@LHC

- Decay  $A_0 \rightarrow H_0 Z$  Dominant for  $m_{A_0} - m_{H_0} \sim v$



- Simple Benchmarks for a Strong EW Phase Transition:

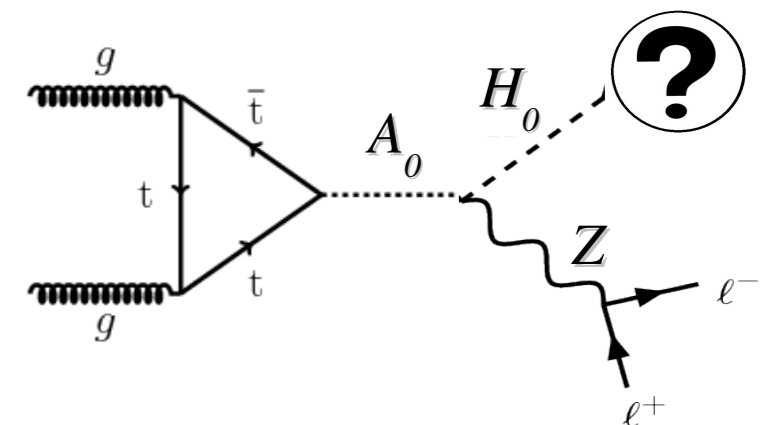
$$m_{A_0} = m_{H^\pm} = 400, m_{H_0} = 180, \mu = 100$$

$$\tan\beta = 2 \quad (\text{controls } gg \rightarrow A_0 \text{ production})$$

- Search Strategy Dictated by Dominant Decay Mode of  $H_0$

— A:  $\alpha - \beta = 0.001\pi$  (aligned)  $b\bar{b}$

..... B:  $\alpha - \beta = 0.1\pi$  (non-aligned)  $WW, ZZ$





# 2HDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

### ① *A few words on the Analysis...*

- ⇒ Type I 2HDM implemented in FeynRules (including gluon-fusion).
- ⇒ Signal & relevant backgrounds generated using MadGraph5\_aMC@NLO. Generated events passed on to Pythia for Parton Showering and Hadronization and subsequently to Delphes for detector simulation.
  - Use of NLO flat K-factors for signal (SusHi) and dominant backgrounds.
- ⇒ “Cut & Count” analysis on a small set of kinematical variables, to extract signal over background.
- ⇒ Determined required Integrated Luminosity at 14 TeV to achieve  $5\sigma$  statistical significance via a C.L.s hypothesis test.
  - *Only statistical uncertainties.*
  - *10% systematic uncertainty on background.*

# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

② *Benchmark A:*  $A_0 \rightarrow H_0 Z \rightarrow \bar{b}b \ell\ell$  ( $\alpha-\beta = 0.001\pi$ )

⇒ Irreducible backgrounds are  $Z\bar{b}b$ ,  $t\bar{t}$ ,  $ZZ$ ,  $hZ$

⇒ *Analysis at 14 TeV: Event Selection*

→ *Anti- $k_T$  Jets with distance parameter  $R = 0.6$*

→ *2  $b$ -tagged Jets with  $|\eta| < 2.5$*

→ *2 Isolated (within a cone of 0.3), Same-flavour leptons.  $|\eta| < 2.5$  (2.7) for electrons (muons)*

→  $P_T^{\ell_1} > 40 \text{ GeV}$ ,  $P_T^{\ell_2} > 20 \text{ GeV}$ .

	K-factor:	1.6	1.5	1.4	-	-
	Signal	$t\bar{t}$	$Z\bar{b}b$	$ZZ$	$Zh$	
Event selection	14.6	1578	424	7.3	2.7	
$80 < m_{\ell\ell} < 100 \text{ GeV}$	13.1	240	388	6.6	2.5	
$H_T^{\bar{b}b} > 150 \text{ GeV}$	8.2	57	83	0.8	0.74	
$H_T^{\ell\ell\bar{b}b} > 280 \text{ GeV}$	5.3	5.4	28.3	0.75	0.68	
$\Delta R_{bb} < 2.5$ , $\Delta R_{\ell\ell} < 1.6$	5.3	5.4	28.3	0.75	0.68	
$m_{bb}$ , $m_{\ell\ell\bar{b}b}$ signal region	3.2	1.37	3.2	< 0.01	< 0.02	



# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

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$\rightarrow P_T^{\ell_1} > 40 \text{ GeV}$ ,  $P_T^{\ell_2} > 20 \text{ GeV}$ .

14 TeV LHC,  $\mathcal{L} = 20 \text{ fb}^{-1}$

Invariant mass windows:

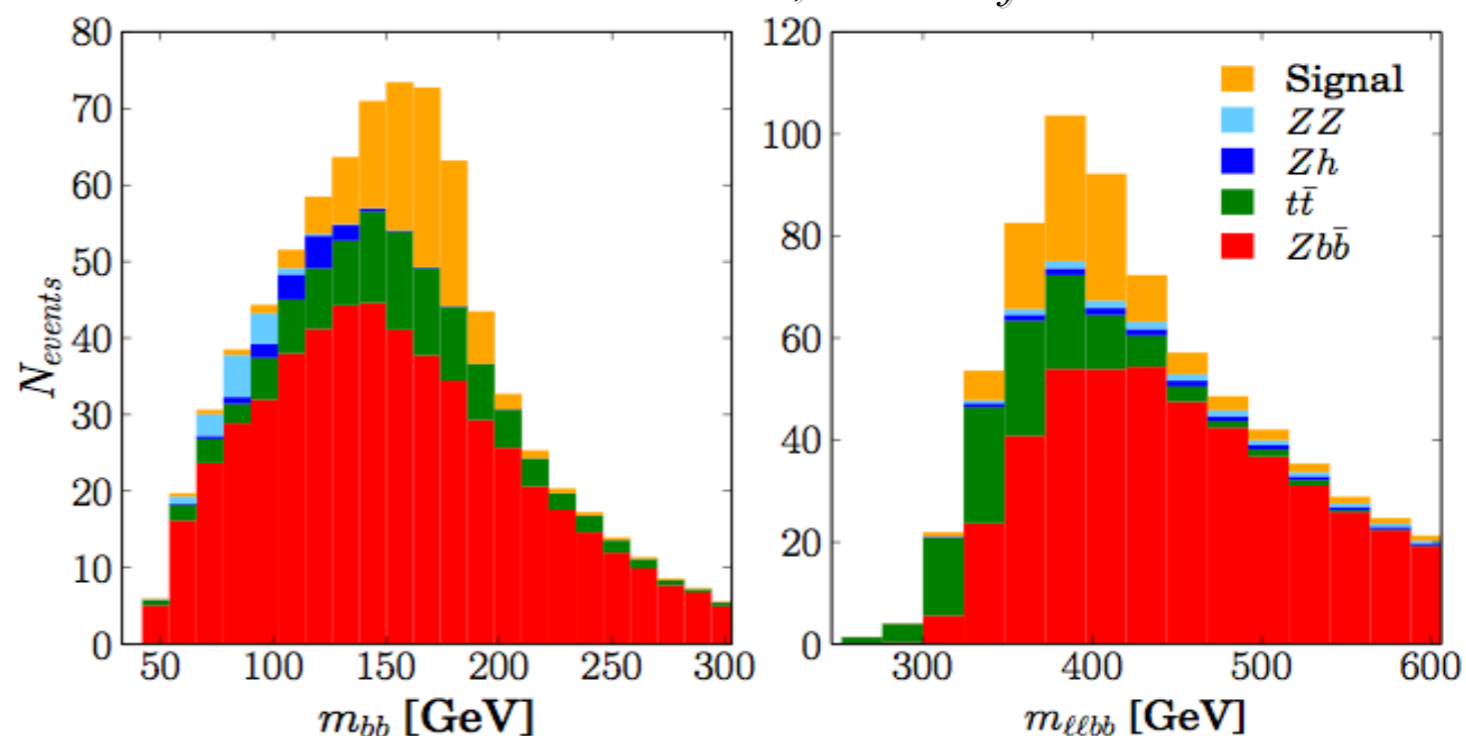
$$m_{\bar{b}b} \rightarrow (m_{H_0} - 20) \pm 30 \text{ GeV}$$

$$m_{\ell\ell\bar{b}b} \rightarrow (m_{A_0} - 20) \pm 40 \text{ GeV}$$

$5\sigma$  signal significance for:

$$\mathcal{L} \doteq 15 \text{ fb}^{-1} \quad (\text{statistics only})$$

$$\mathcal{L} = 40 \text{ fb}^{-1} \quad (10\% \text{ systematics})$$





# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

③ **Benchmark B:**  $A_0 \rightarrow H_0 Z \rightarrow W^+W^- \ell\ell \rightarrow 4\ell + 2\nu$  ( $\alpha-\beta = 0.1\pi$ )

⇒ Most sensitive  $A_0$  search channel away from alignment

⇒  $A_0 \rightarrow H_0 Z \rightarrow ZZ\ell\ell \rightarrow 4\ell + 2j$  also promising

*B. Coleppa, F. Kling, S. Su, JHEP 1409 (2014) 161*

⇒ Main backgrounds are  $ZZ$ ,  $Zt\bar{t}$ ,  $hZ$ ,  $ZWW$  subdominant

⇒ Analysis & Event Selection similar to previous case:

→ 4 Isolated (cone of 0.3) leptons, same-flavour pairs.  $|\eta| < 2.5$  (2.7) for electrons (muons)

→  $P_T^{\ell 1} > 40 \text{ GeV}$ ,  $P_T^{\ell 2,3,4} > 20 \text{ GeV}$ .

→ 1 pair of SF leptons must reconstruct  $m_Z$

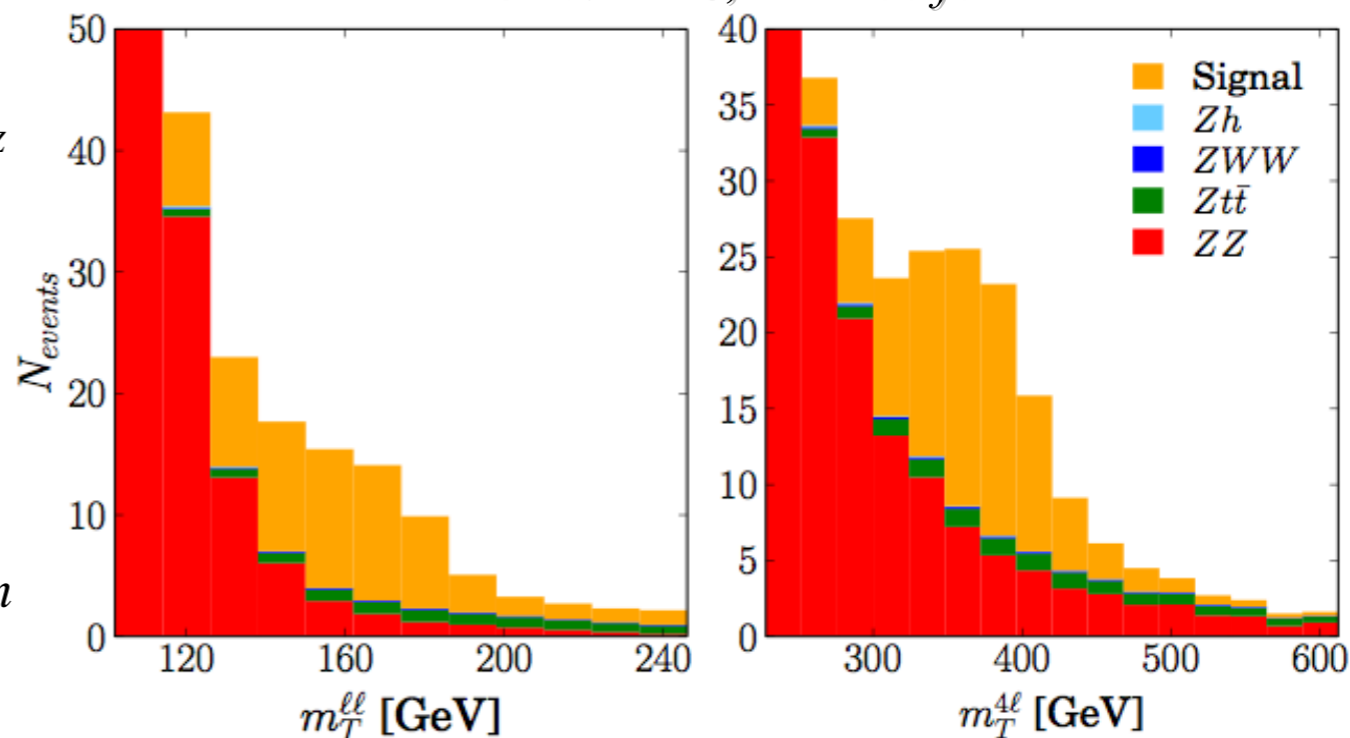
→ Transverse mass variables:

$$(m_T^{\ell\ell})^2 = (\sqrt{p_{T,\ell\ell}^2 + m_{\ell\ell}^2} + \cancel{p}_T)^2 - (\vec{p}_{T,\ell\ell} + \vec{\cancel{p}}_T)^2$$

$$m_T^{4\ell} = \sqrt{p_{T,\ell'\ell'}^2 + m_{\ell'\ell'}^2} + \sqrt{p_{T,\ell\ell}^2 + (m_T^{\ell\ell})^2}$$

$m_T^{4\ell} > 260 \text{ GeV}$  allows for Signal Extraction

14 TeV LHC,  $\mathcal{L} = 60 \text{ fb}^{-1}$



# ZHDM@LHC

## LHC DISCOVERY POTENTIAL OF BENCHMARK SCENARIOS

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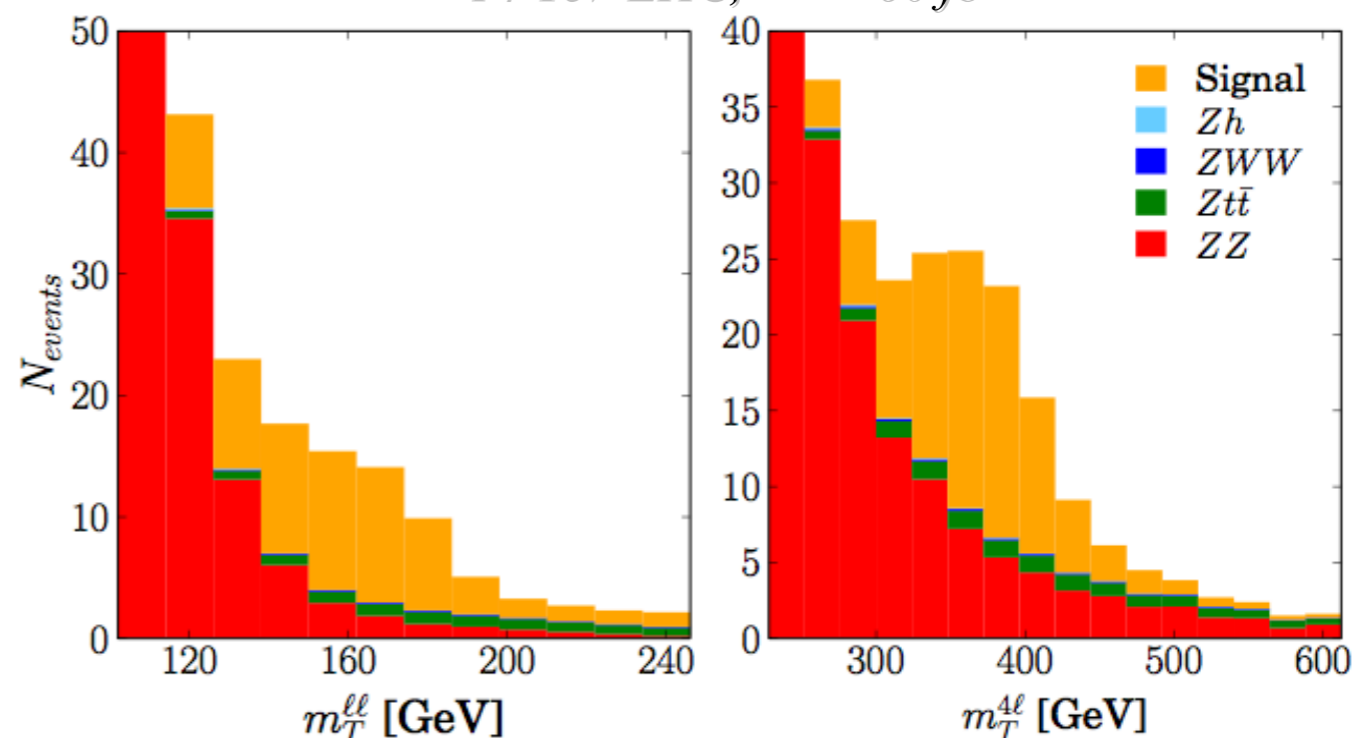
5 $\sigma$  signal significance for:

$\mathcal{L} = 60 \text{ fb}^{-1}$  (statistics only)

$\mathcal{L} = 200 \text{ fb}^{-1}$  (10% systematics)

(conservative...)

14 TeV LHC,  $\mathcal{L} = 60 \text{ fb}^{-1}$



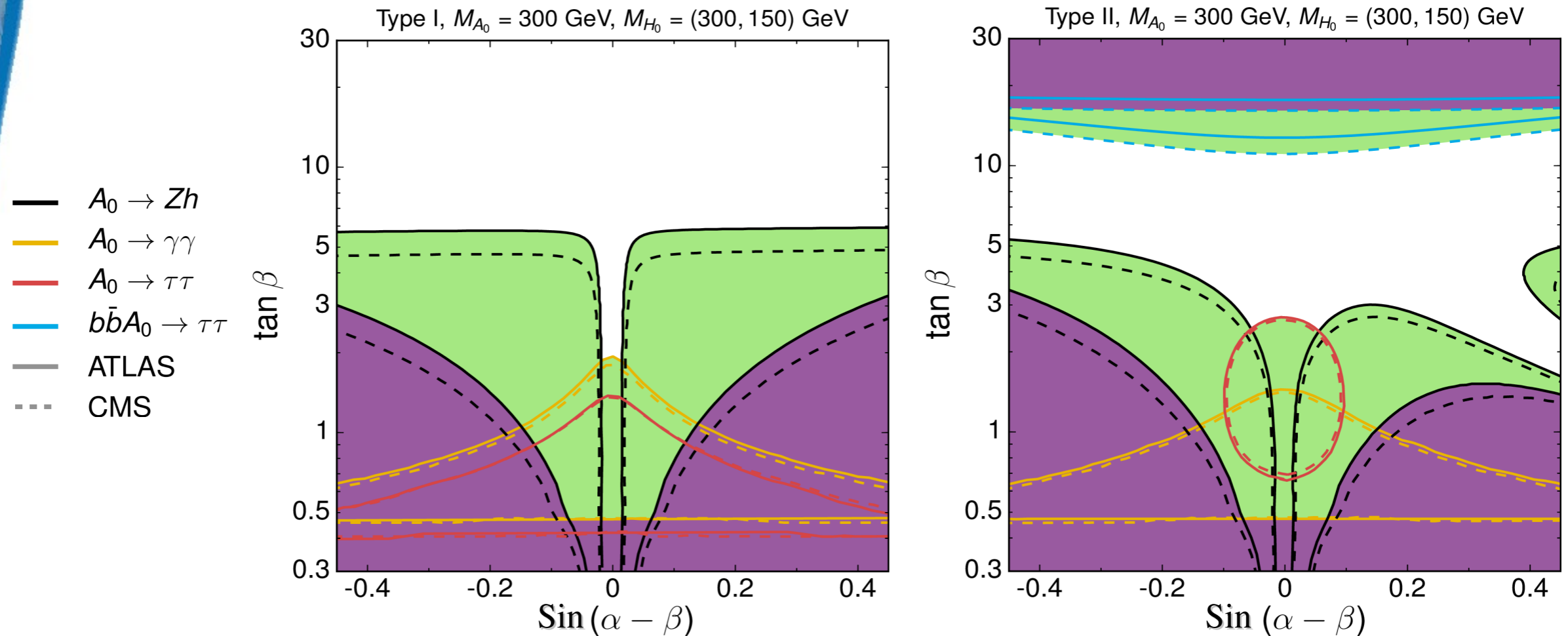


# Remarks on "Standard" 2HDM LHC Searches

$A_0 \rightarrow H_0 Z$  channel Open  $\rightarrow$  BR to other decay channels drastically reduced

LIMITS FROM STANDARD LHC SEARCHES SIGNIFICANTLY WEAKENED....

*G. Dorsch, S. Huber, K. Mimasu, J. M. No, 1507.xxxxx*



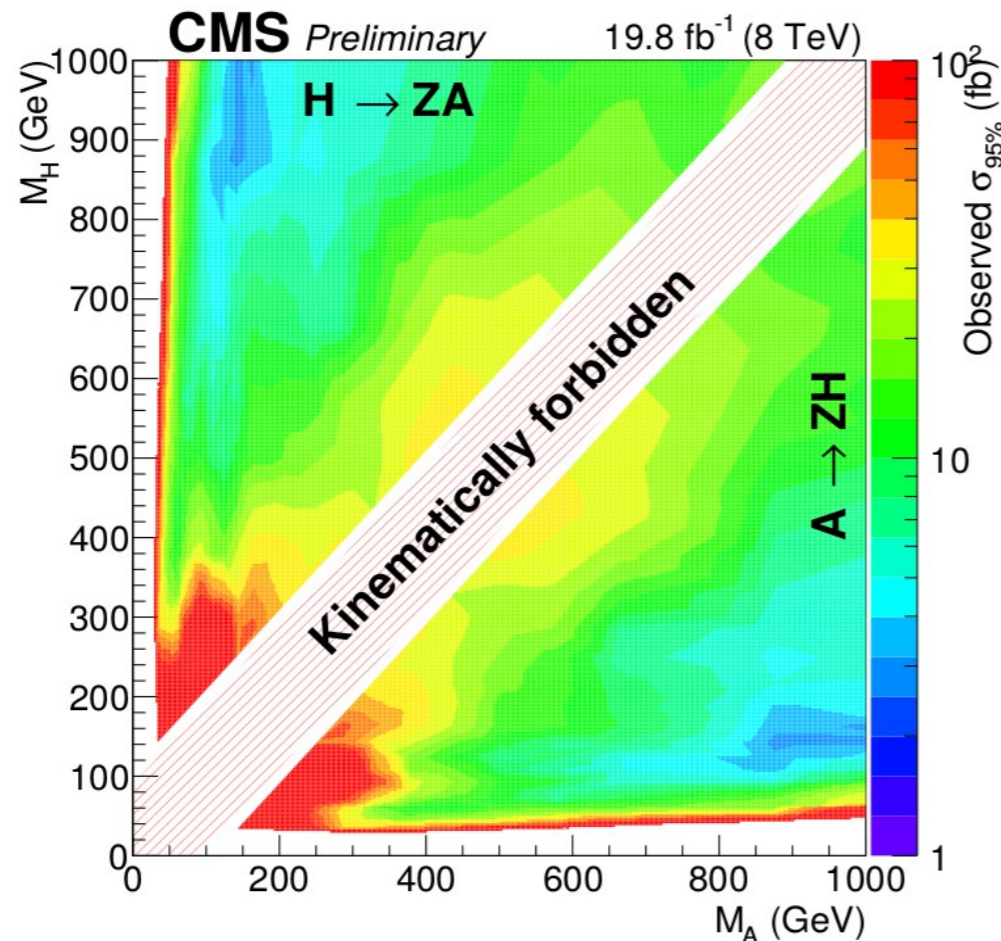


# New Searches “Fill in the Gap”

## CMS-PAS-HIG-15-001

Search for H/A decaying into Z and A/H, with  $Z \rightarrow \ell\ell$  and  $A/H \rightarrow bb$  or  $A/H \rightarrow \tau\tau$

The CMS Collaboration



One important motivation for 2HDMs is that these models provide a way to explain the asymmetry between matter and anti-matter observed in the Universe [4, 5]. Another important motivation is Supersymmetry [6], which is a theory that falls in the broad class of 2HDMs. Axion models [7], which would explain how the strong interaction does not violate the CP symmetry, would give rise to an effective low-energy theory with two Higgs doublets. Finally, it has also been recently noted [8] that certain realizations of 2HDMs can accommodate the muon  $g - 2$  anomaly [9] without violating the present theoretical and experimental constraints.

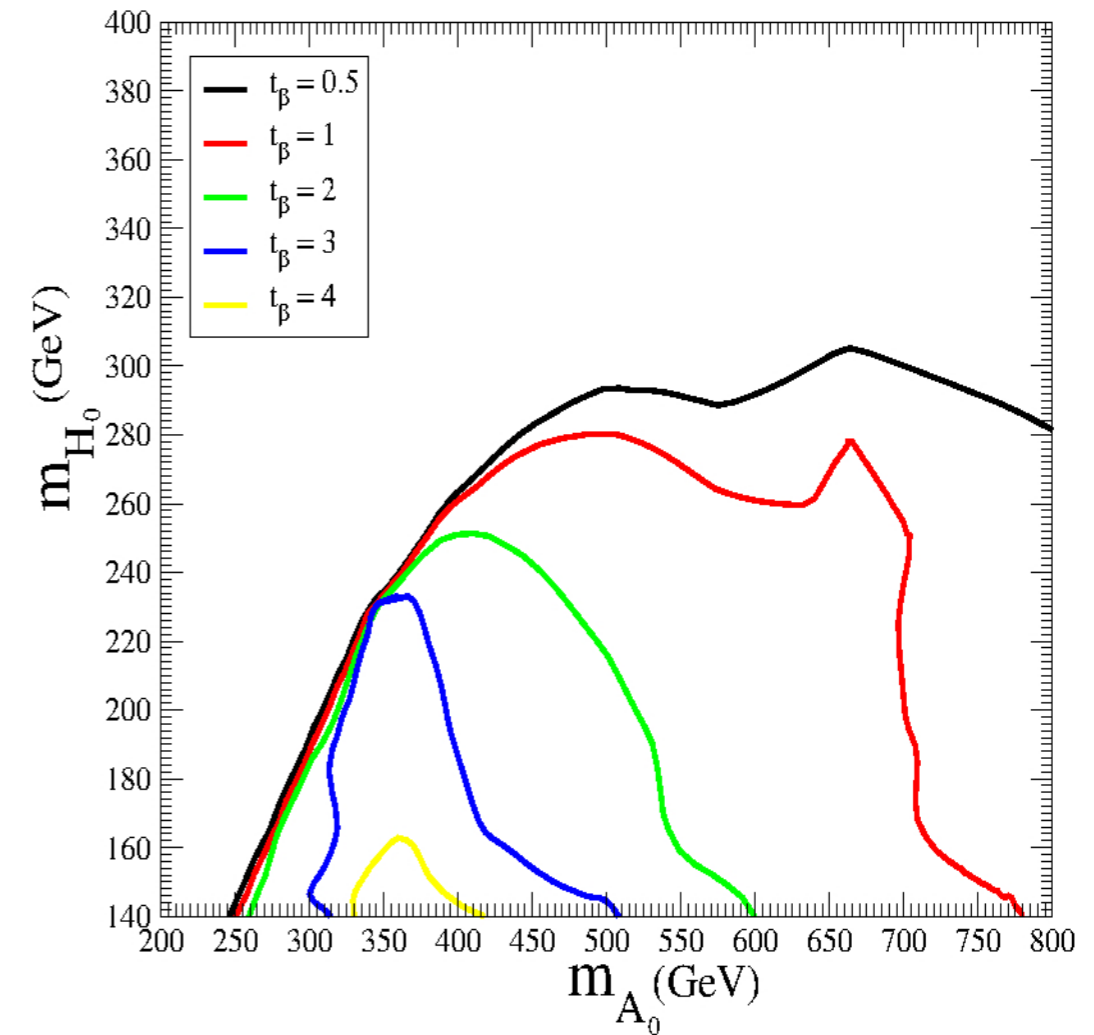
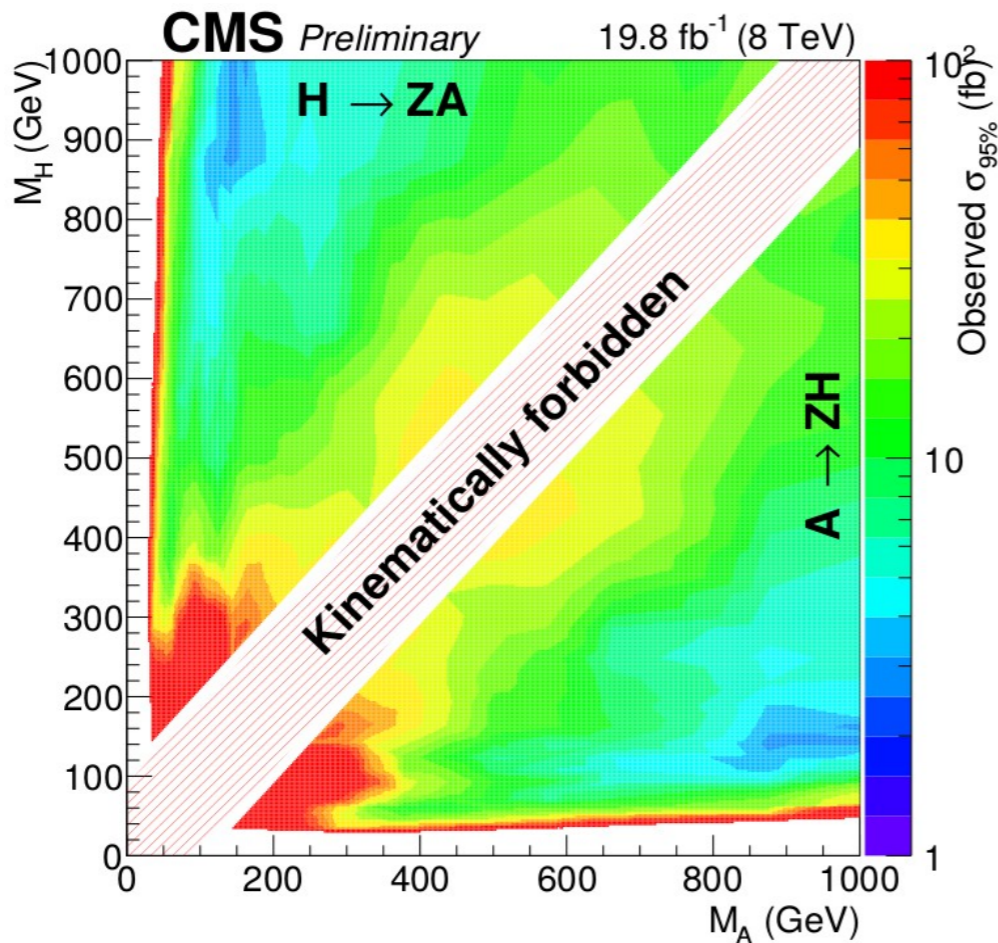
In the most general case 14 parameters are necessary to describe the scalar sector in a 2HDM. However, only 6 free parameters remain once the so-called  $Z_2$  symmetry is imposed to suppress flavor changing neutral currents, in agreement with experimental observations, and the values of the mass of the recently discovered Higgs boson (125 GeV) and the electroweak vacuum expectation value (246 GeV) are assumed. The compatibility of a 125 GeV SM-like Higgs boson with 2HDMs is possible in the so-called alignment limit. In such a limit, one of the CP-even scalars,  $h$  or  $H$ , is identified with the 125 GeV Higgs boson and the condition  $\cos(\beta - \alpha) \approx 0$  or  $\sin(\beta - \alpha) \approx 0$  is satisfied, where  $\tan \beta$  and  $\alpha$  are, respectively, the ratio of the vacuum expectation values, and the mixing angle of the two Higgs doublets. A recent theoretical study [5] has shown that, in this limit, a large mass splitting ( $>100$  GeV) between the  $A$  and  $H$  bosons would favor the electroweak phase transition that would be at the origin of the baryogenesis process in the early Universe, thus explaining the currently observed matter-antimatter asymmetry in the Universe. In such a scenario, the most frequent decay mode of the pseudoscalar  $A$  boson would be  $A \rightarrow ZH$ .

Meaningful constraints from LHC Run 1



# New Searches "Fill in the Gap"

CMS-PAS-HIG-15-001



*G. Dorsch, S. Huber, K. Mimasu, J. M. No, 1507.xxxxx*

Meaningful constraints from LHC Run 1

# Conclusions

- Extended Higgs Sectors are Archetype Scenarios for a Connection between EW Cosmology & LHC physics

EW PHASE  
TRANSITION

2HDM:

**EW Phase Transition  
Signature**

$$\underline{A_0 \rightarrow H_0 Z} \rightarrow \bar{b}b \ell\ell$$

$$\underline{A_0 \rightarrow H_0 Z} \rightarrow W^+W^- \ell\ell \rightarrow 4\ell + 2\nu$$

- For these “Hierarchical” 2HDM, Standard LHC searches become less efficient
- Promising Prospects for  $A_0 \rightarrow H_0 Z$  Searches at LHC14

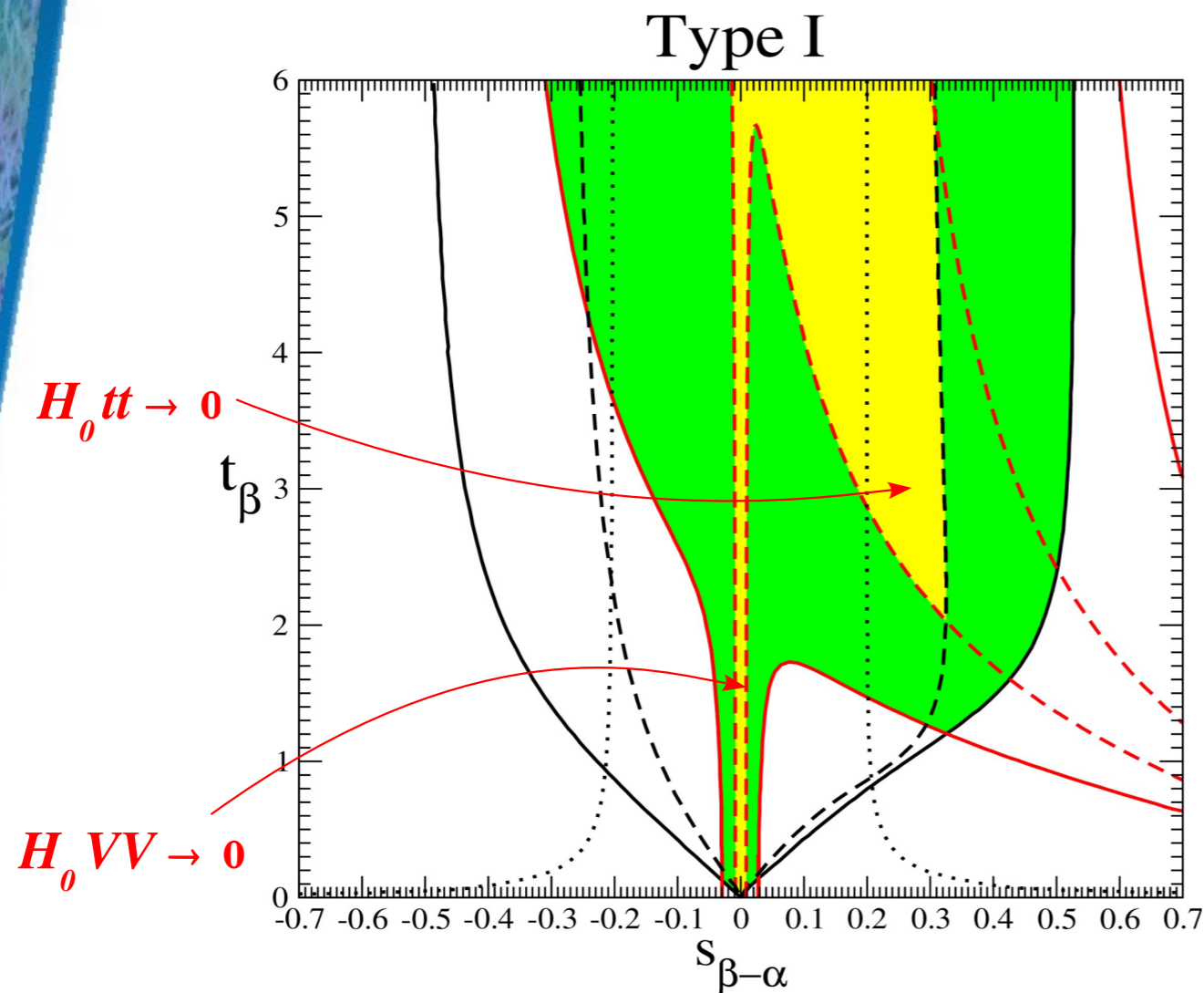




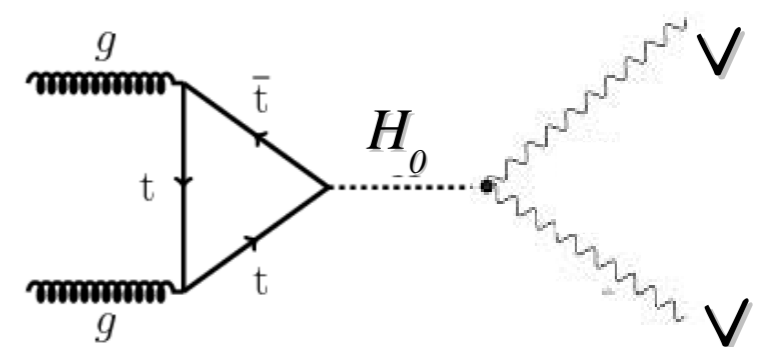
**Back-up**

# Remarks on "Standard" 2HDM LHC Searches

EW Phase Transition LHC Searches fill region where standard 2HDM LHC searches are not as sensitive



( $m_{H_0} = 200 \text{ GeV}$ )



+

Measurements of light Higgs couplings