

# A Fresh Look at the Sterile Neutrino Production from the Frozen-In Scalar

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In cooperation with Michael A. Schmidt

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# Content

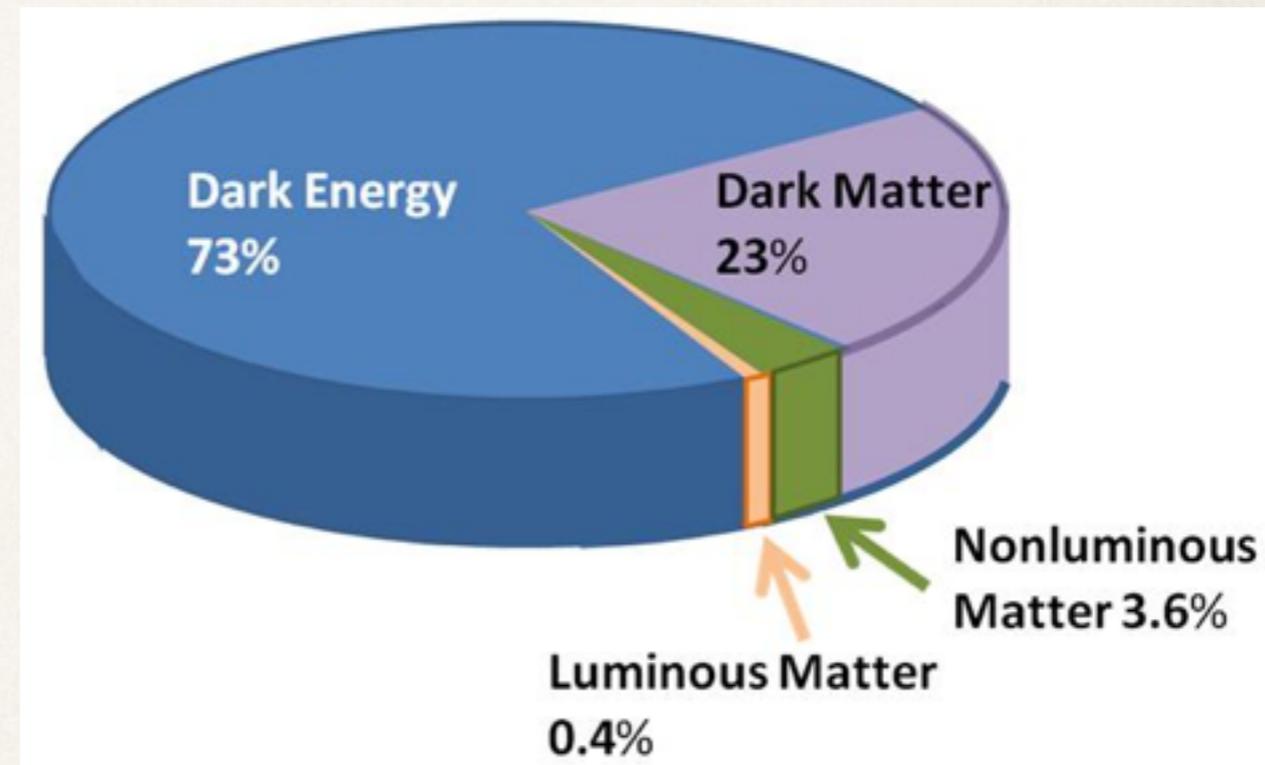
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- ❖ Motivation
- ❖ Dark Matter Production
- ❖ Free Streaming Horizon of the Dark Matter
- ❖ Conclusion

# Motivation

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- ❖ The Standard Model of Particle Physics is so successful, but yet still not complete —Hierarchy problem, flavor problem, neutrino mass, **dark matter**, etc.
- ❖ 23% of the Universe is governed by dark matter. **Whose characteristic is still unknown.** — WIMP, axion, sterile neutrino, etc.
- ❖ A good candidate that one can obtain from a minimal extension of the SM is **the keV sterile neutrino** which can serve as the good candidate for the warm dark matter.



# Possible Ways for the sterile neutrino production.

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- ❖ Dodelson-Widrow (DW) Mechanism (Non-Resonant Oscillation) [S. Dodelson and L.M. Widrow, '94]
- ❖ Shi-Fuller Mechanism (Resonant Oscillation) [X.-D. Shi and G.M. Fuller, '98]
- ❖ Inflaton decays [M. Shaposhnikov and I. Tkachev, '06; F. Bezrukov and D. Gorbunov, '09]
- ❖ Thermal Production by new gauge interactions and subsequent dilution by production of entropy [F. Bezrukov, H. Hettmansperger and M. Lindner, '09; M. Nemevsek, G. Senjanovic and Y. Zhang, '12]
- ❖ The decay of the scalar field [recent helpful reference, A. Merle and M. Totzauer, '15]
  - which is in-thermal equilibrium with the thermal plasma [A. Kusenko, '06; K. Petraki and A. Kusenko, '07; M. Frigerio and C.E. Yaguna, '14]
  - which is out-of-thermal equilibrium with the thermal plasma [A. Merle, V. Niro and D. Schmidt, '13; AA and M. A. Schmidt, '14]

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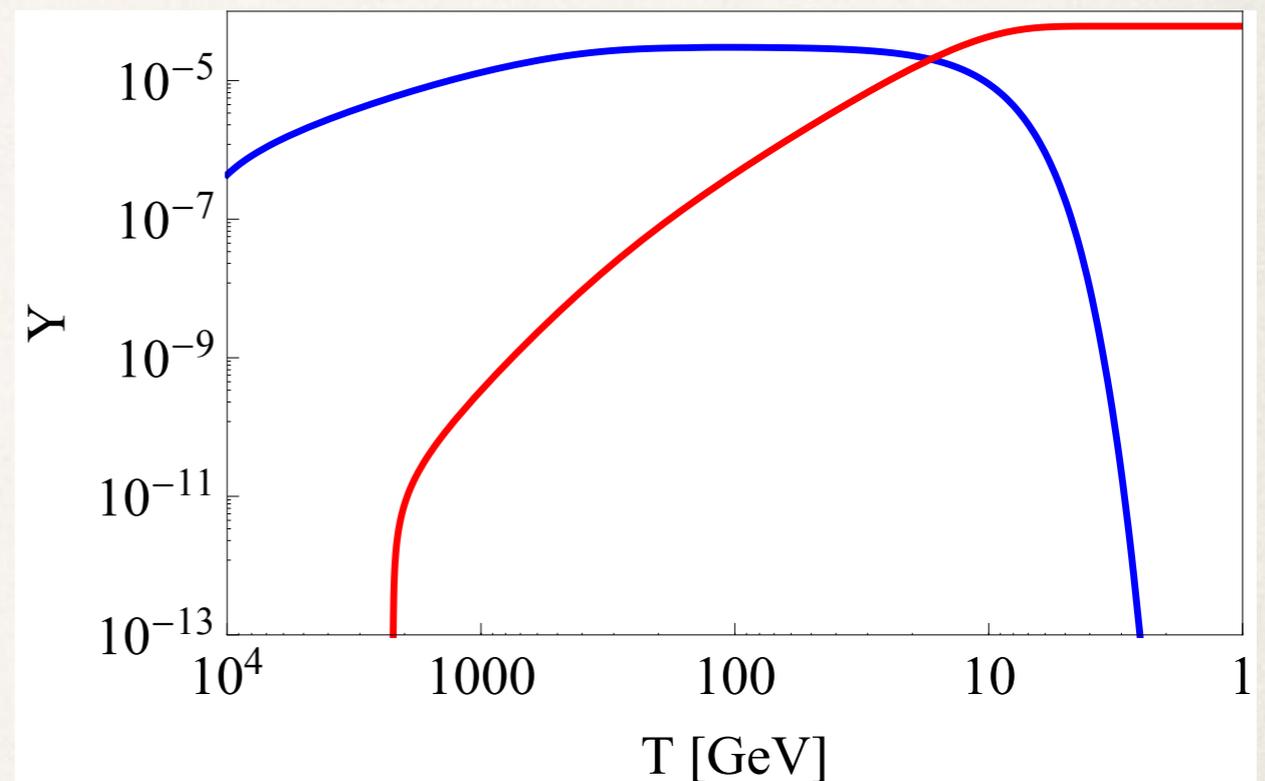
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# Sterile neutrino production from the out-of-thermal equilibrium decay of the scalar field — Frozen-In Scalar

- ❖ This is first proposed by A. Merle, V. Niro and D. Schmidt, '13 — they have investigated the region where the scalar mass being heavier than the SM Higgs.
- ❖ The Mechanism works in two steps:
  - 1. The scalar field gets produced first through its very weakly interactions with the SM fields via the Higgs portal.
  - 2. The sterile neutrino gets produced from the decay of the scalar field.

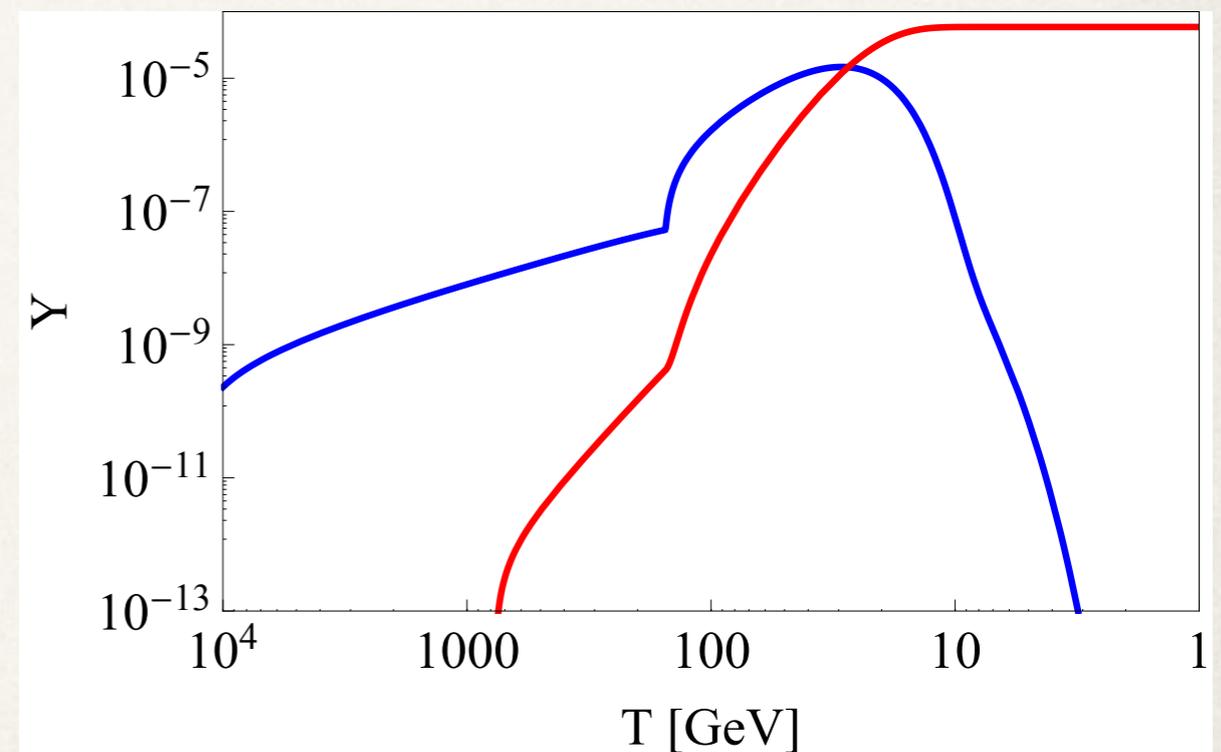


(a)  $m_\sigma = 500$  GeV

# Sterile neutrino production from the out-of-thermal equilibrium decay of the scalar field — Frozen-In Scalar

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- ❖ Our work —
- ❖ Investigating also the region where the scalar field is lighter than the SM Higgs.
- ❖ Investigating in detail the different contributions to the production.
- ❖ Investigating the possibility to produce the 7.1 keV sterile neutrino — explaining 3.55 keV Line.
- ❖ Investigating the channel of the SM Higgs decaying into the frozen-scalar field.



(b)  $m_\sigma = 60$  GeV

# Model

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- ❖ Introducing one light SM singlet fermion and one scalar singlet
- ❖ Introducing a discrete  $Z_4$  symmetry
- ❖ The transformation properties of the fields are

$$L \rightarrow iL \quad E^c \rightarrow -iE^c \quad N \rightarrow -iN \quad H \rightarrow H \quad \phi \rightarrow -\phi .$$

# Model

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- ❖ Lagrangian density

$$-\mathcal{L} = y_E L H E^C + y_{LN} L H N + \frac{1}{2} y_N \phi N^2 + \frac{y_\nu}{\Lambda^2} L L H H \phi + \text{h.c.}$$

- ❖ The scalars obtain VEVs

$$H = \begin{pmatrix} G^+ \\ v + \frac{1}{\sqrt{2}} (h + iG^0) \end{pmatrix} \quad \phi = v_\phi + \sigma$$

- ❖ Neutrino Mass

$$m_\nu = y_\nu \frac{v_\phi v^2}{\Lambda^2} - \frac{y_{LN} y_{LN}^T}{y_N} \frac{v^2}{v_\phi} .$$

# Model

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- ❖ The scalar potential

$$V = -\mu_H^2 H^\dagger H + \frac{\lambda_H}{6} (H^\dagger H)^2 + \frac{\lambda_{H\phi}}{2} H^\dagger H \phi^2 - \frac{1}{2} \mu_\phi^2 \phi^2 + \frac{\lambda_\phi}{24} \phi^4$$

- ❖ Scalar Masses

$$m_h^2 = 2\mu_H^2 = \frac{2\lambda_H}{3} v^2 \quad \text{and} \quad m_\sigma^2 = 2\mu_\phi^2 = \frac{\lambda_\phi}{3} v_\phi^2 .$$

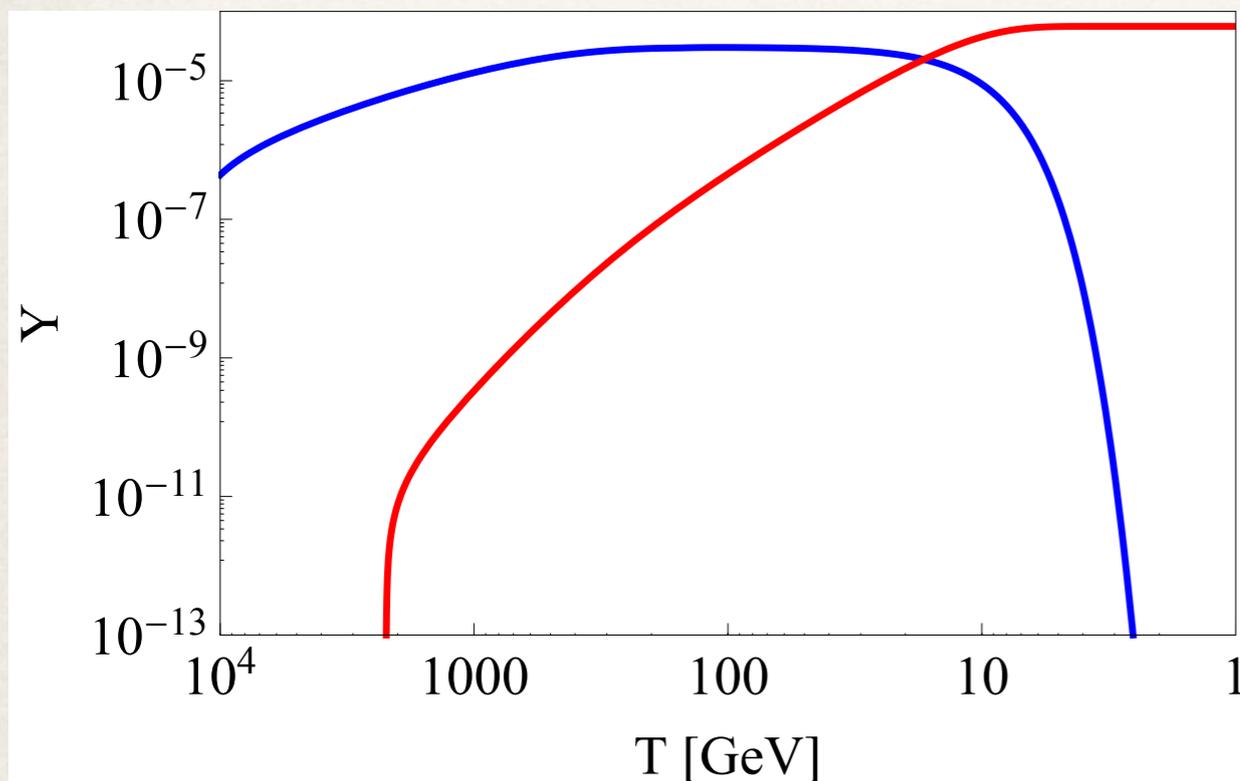
- ❖ The Relevant couplings leading to the production

$$\Delta V = \lambda_{H\phi} \frac{h^2 \sigma^2}{4} + \sqrt{2} \lambda_{H\phi} v \frac{h \sigma^2}{2}$$

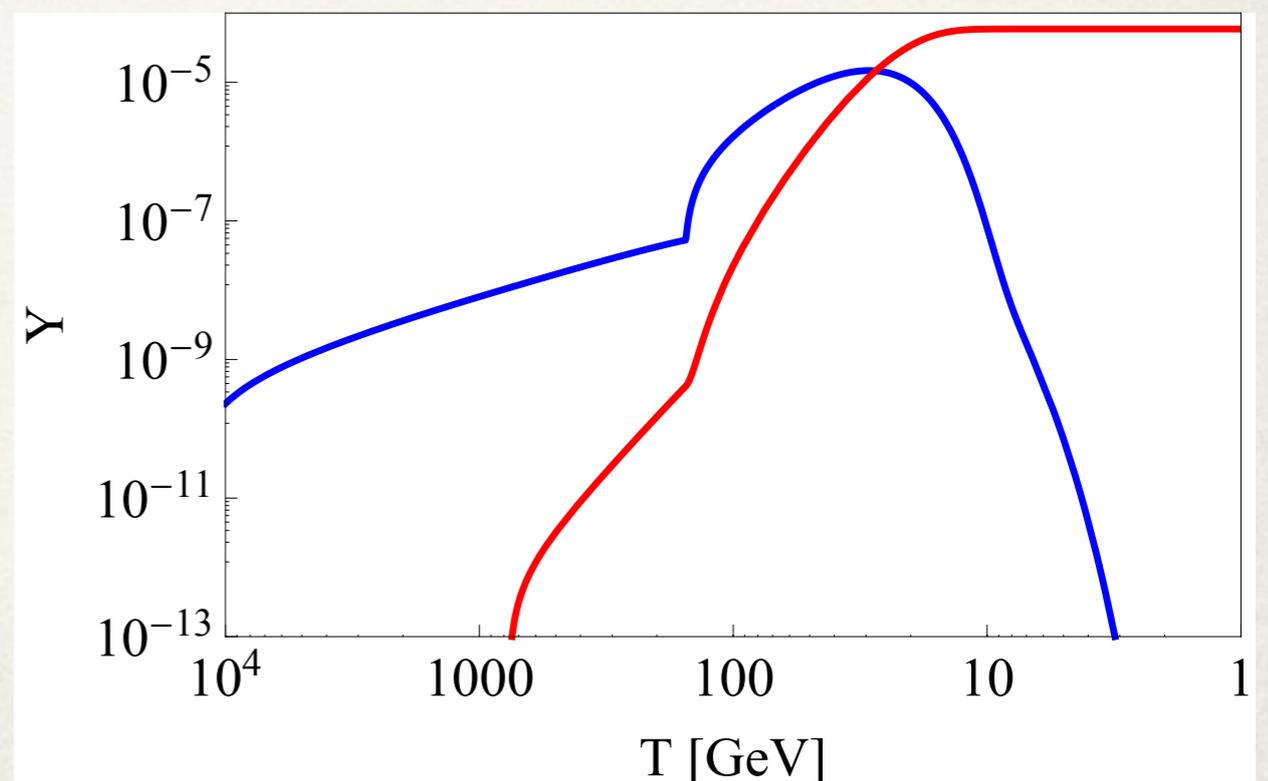
# Sterile neutrino production via the decay of the frozen-in scalar

The relevant terms in the production:

$$\Delta V = \lambda_{H\phi} \frac{h^2 \sigma^2}{4} + \sqrt{2} \lambda_{H\phi} v \frac{h \sigma^2}{2}$$



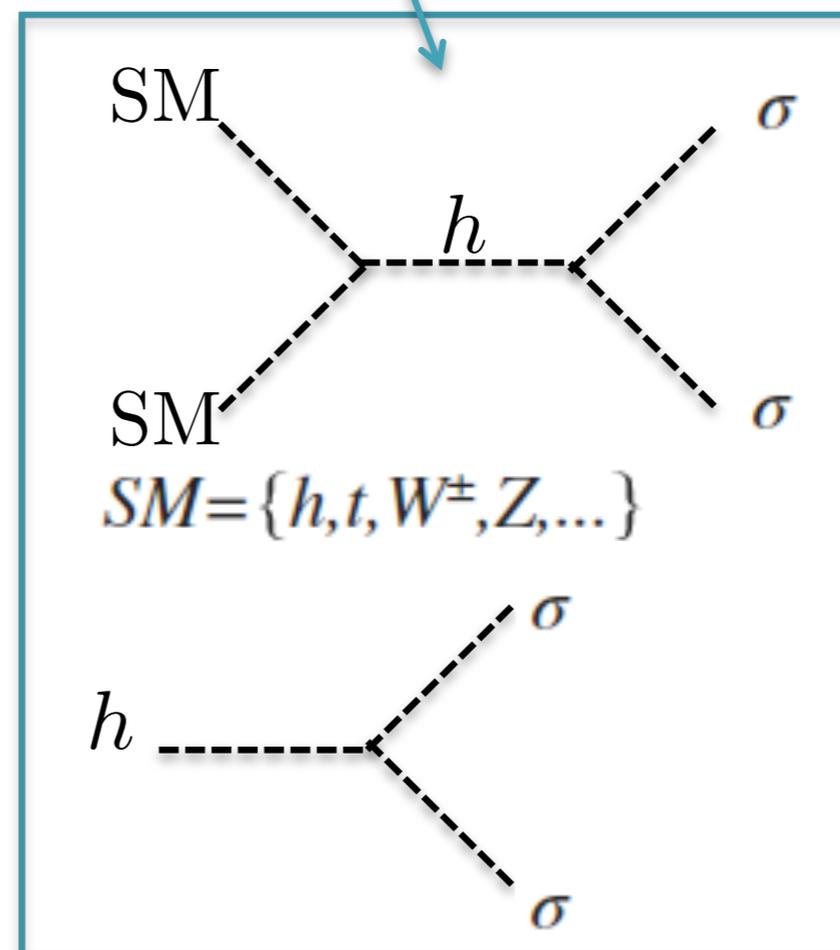
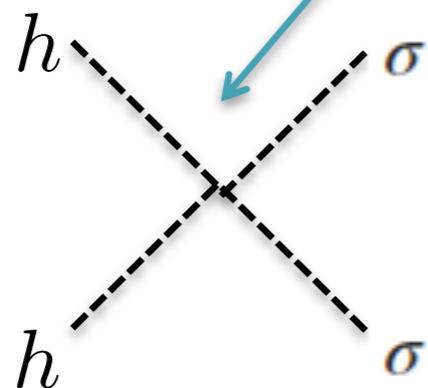
(a)  $m_\sigma = 500$  GeV



(b)  $m_\sigma = 60$  GeV

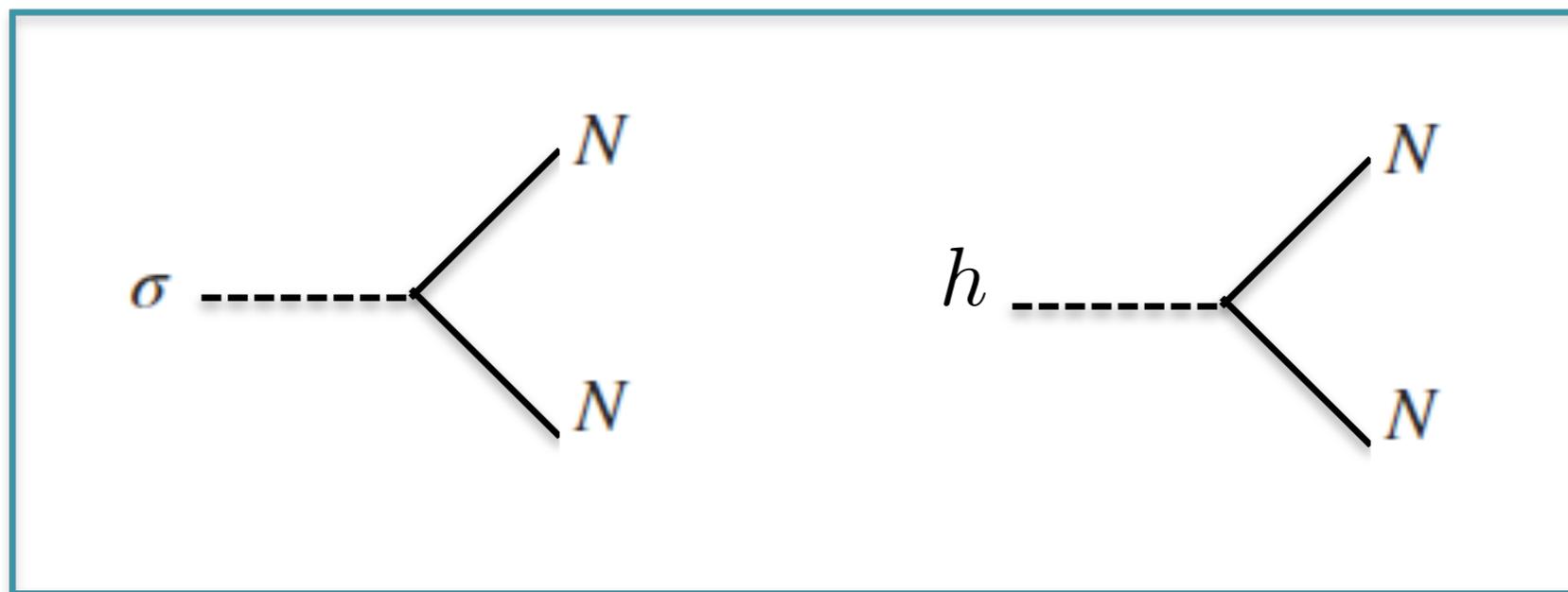
# Scalar Productions

$$\Delta V = \lambda_{H\phi} \frac{h^2 \sigma^2}{4} + \sqrt{2} \lambda_{H\phi} v \frac{h \sigma^2}{2}.$$



# Sterile Neutrino Productions

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$$\Gamma(\sigma \rightarrow NN) \gg \Gamma(h \rightarrow NN)$$

# Boltzmann Equations

$$\frac{dY_\sigma}{dT} = \frac{dY_\sigma^A}{dT} + \frac{dY_\sigma^D}{dT} + \frac{dY_\sigma^{\text{HD}}}{dT}$$

$$\frac{dY_N}{dT} = \frac{dY_N^D}{dT} + \frac{dY_N^{\text{HD}}}{dT}$$

$$Y_{\sigma,N} \equiv n_{\sigma,N}/s$$

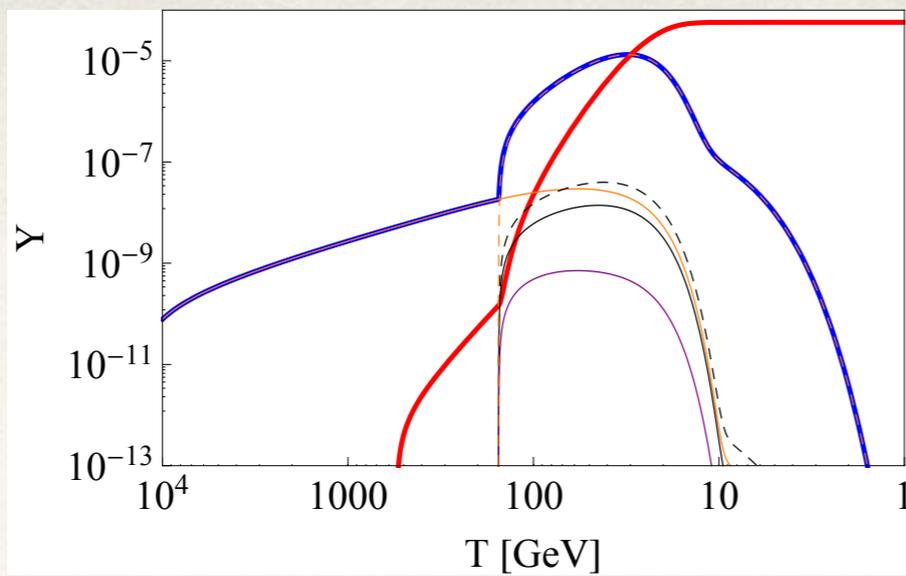
$$\frac{dY_\sigma^A}{dT} = \sqrt{\frac{\pi}{45G_N}} \sqrt{g_*(T)} \sum_{i=h,W,Z,t,b,c,\tau} \langle \sigma v(\sigma\sigma \rightarrow ii) \rangle (Y_\sigma(T)^2 - Y_\sigma^{\text{eq}}(T)^2)$$

$$\frac{dY_\sigma^D}{dT} = -\frac{1}{2} \frac{dY_N^D}{dT}$$

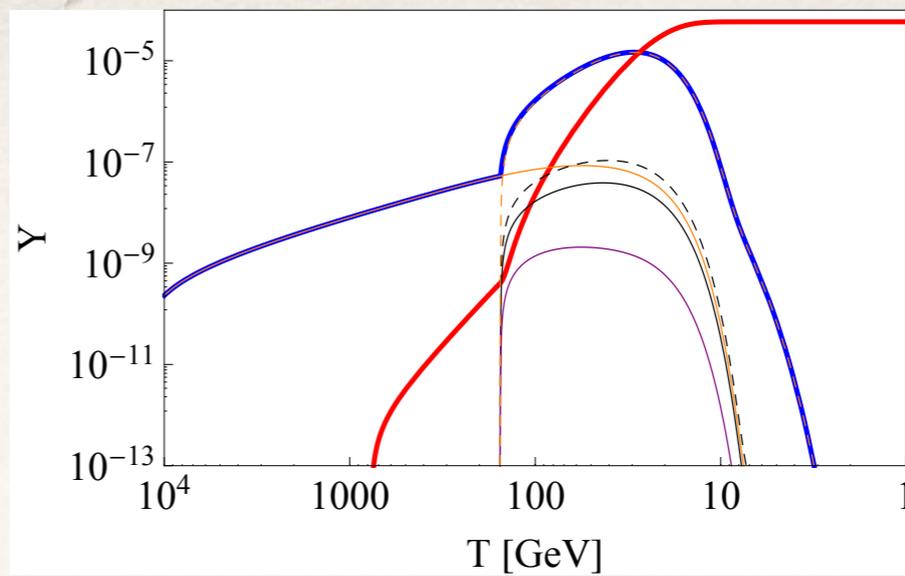
$$\frac{dY_N^D}{dT} = -\sqrt{\frac{45}{\pi^3 G_N}} \frac{1}{T^3} \frac{1}{\sqrt{g_{\text{eff}}(T)}} \langle \Gamma(\sigma \rightarrow NN) \rangle \left( Y_\sigma(T) - \left( \frac{Y_N(T)}{Y_N^{\text{eq}}(T)} \right)^2 Y_\sigma^{\text{eq}}(T) \right)$$

$$\frac{dY_\sigma^{\text{HD}}}{dT} = -\sqrt{\frac{45}{\pi^3 G_N}} \frac{1}{T^3} \frac{1}{\sqrt{g_{\text{eff}}(T)}} \langle \Gamma(H \rightarrow \sigma\sigma) \rangle \left( 1 - \left( \frac{Y_\sigma(T)}{Y_\sigma^{\text{eq}}(T)} \right)^2 \right) Y_h^{\text{eq}}(T)$$

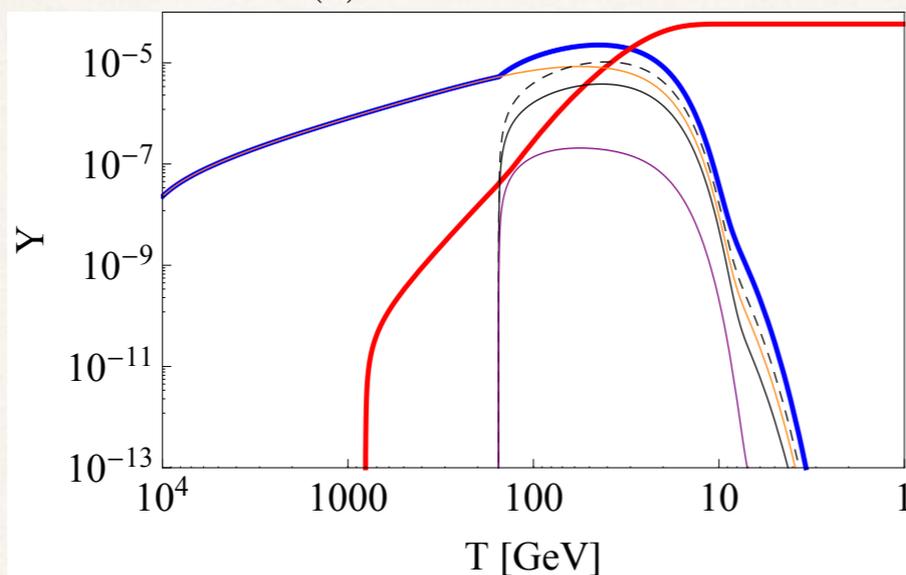
$$\frac{dY_N^{\text{HD}}}{dT} = -\sqrt{\frac{45}{\pi^3 G_N}} \frac{1}{T^3} \frac{1}{\sqrt{g_{\text{eff}}(T)}} \langle \Gamma(H \rightarrow NN) \rangle \left( 1 - \left( \frac{Y_N(T)}{Y_N^{\text{eq}}(T)} \right)^2 \right) Y_h^{\text{eq}}(T) .$$



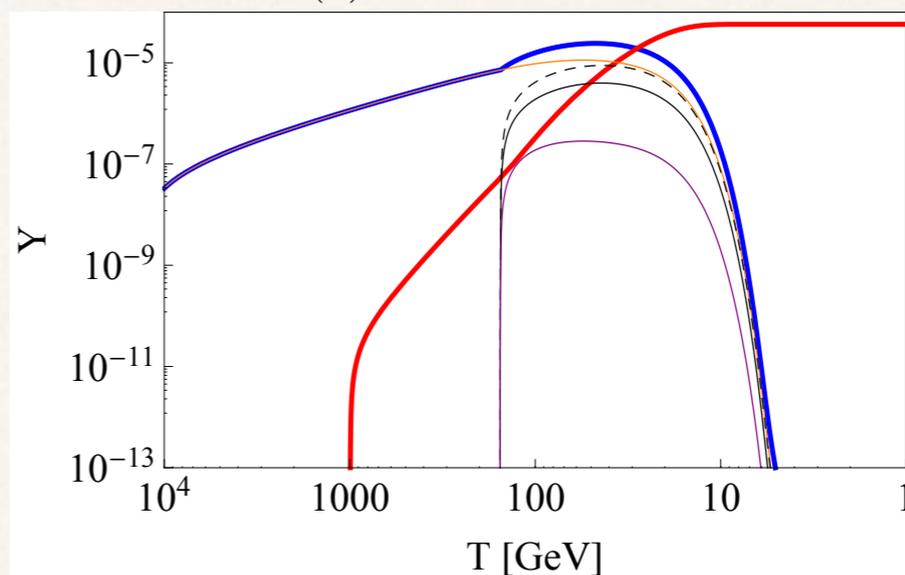
(a)  $m_\sigma = 30\text{GeV}$



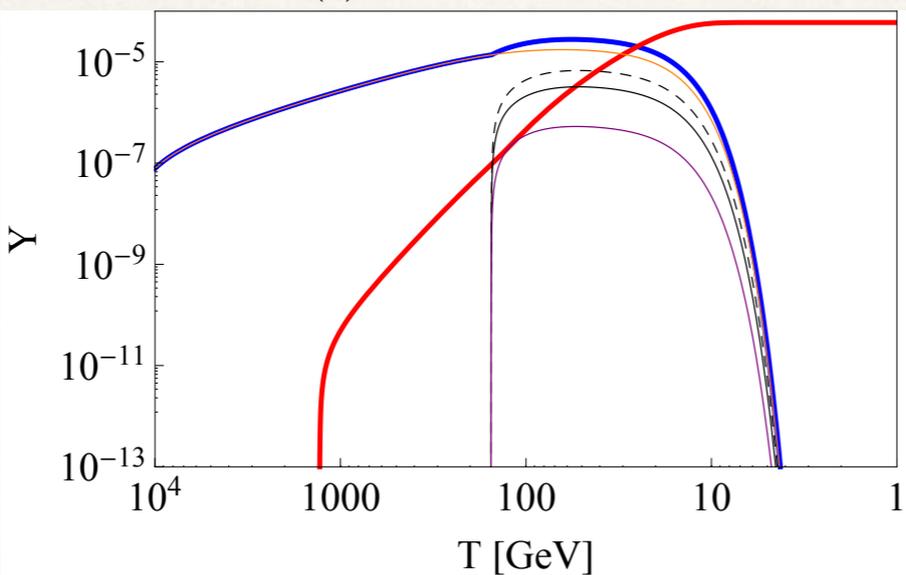
(b)  $m_\sigma = 60\text{GeV}$



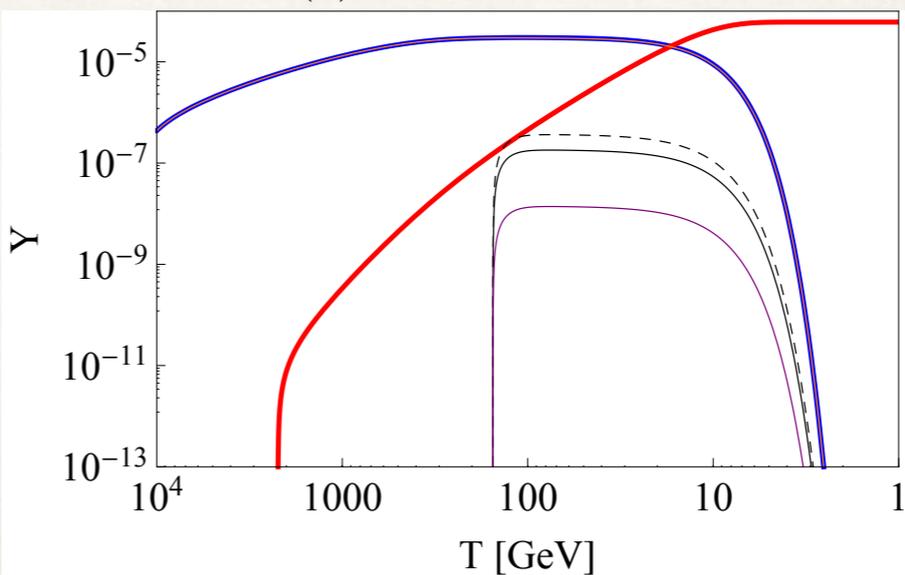
(c)  $m_\sigma = 65\text{GeV}$



(d)  $m_\sigma = 100\text{GeV}$

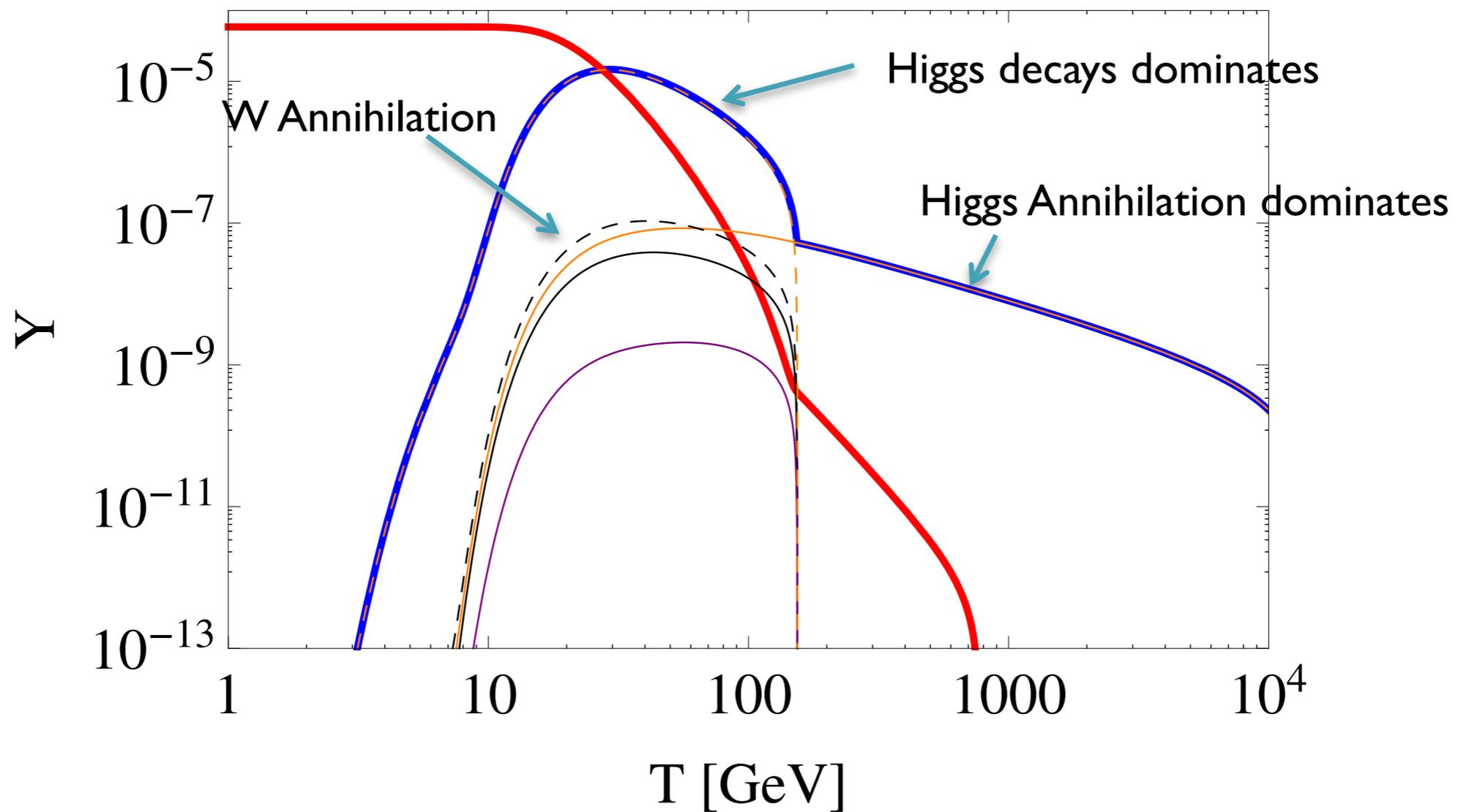


(e)  $m_\sigma = 170\text{GeV}$



(f)  $m_\sigma = 500\text{GeV}$

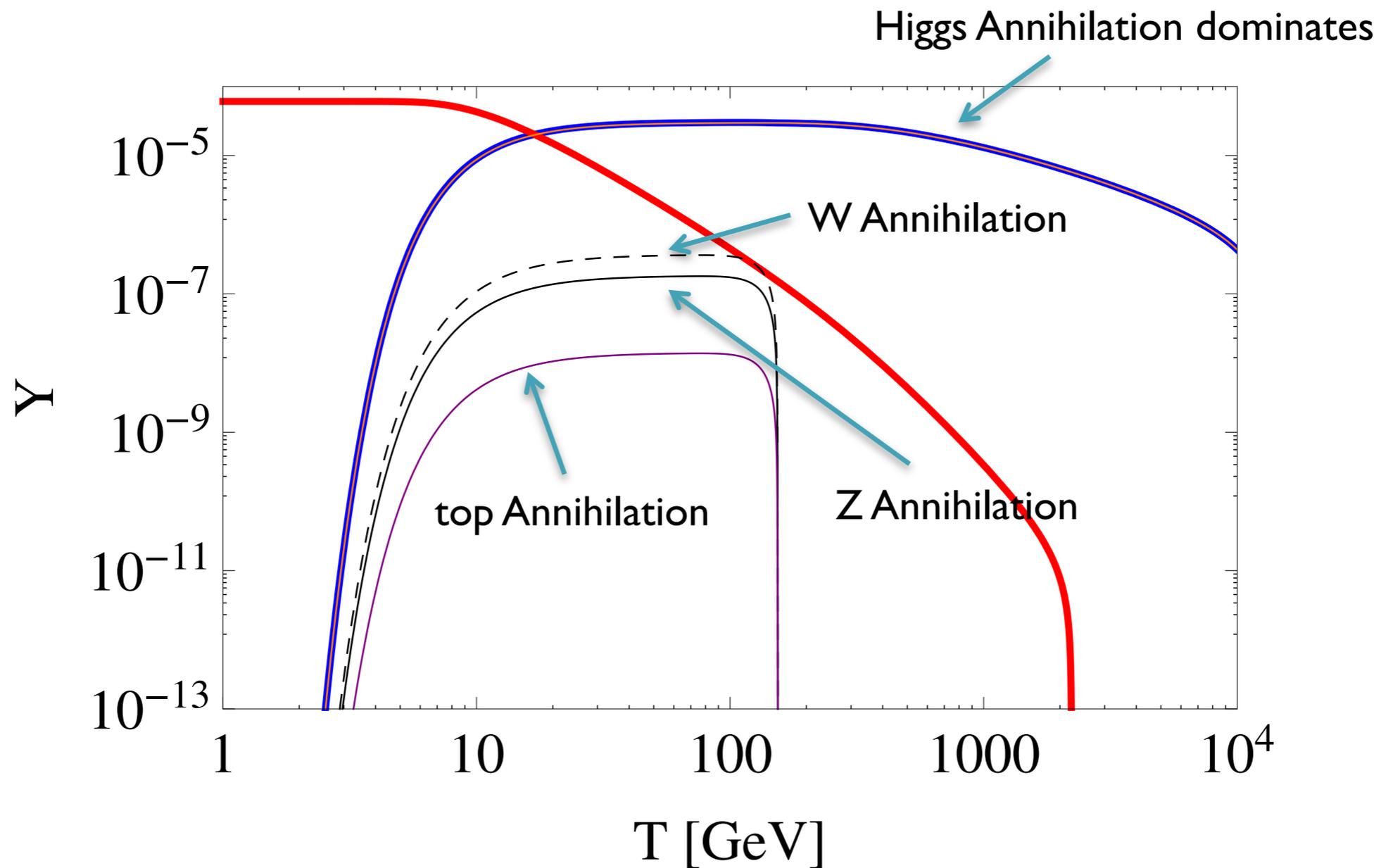
# Dark Matter Production



(b)  $m_\sigma = 60\text{GeV}$

Higgs decay dominates for  $m_\sigma < m_H/2$

# Dark Matter Production



(f)  $m_\sigma = 500\text{GeV}$

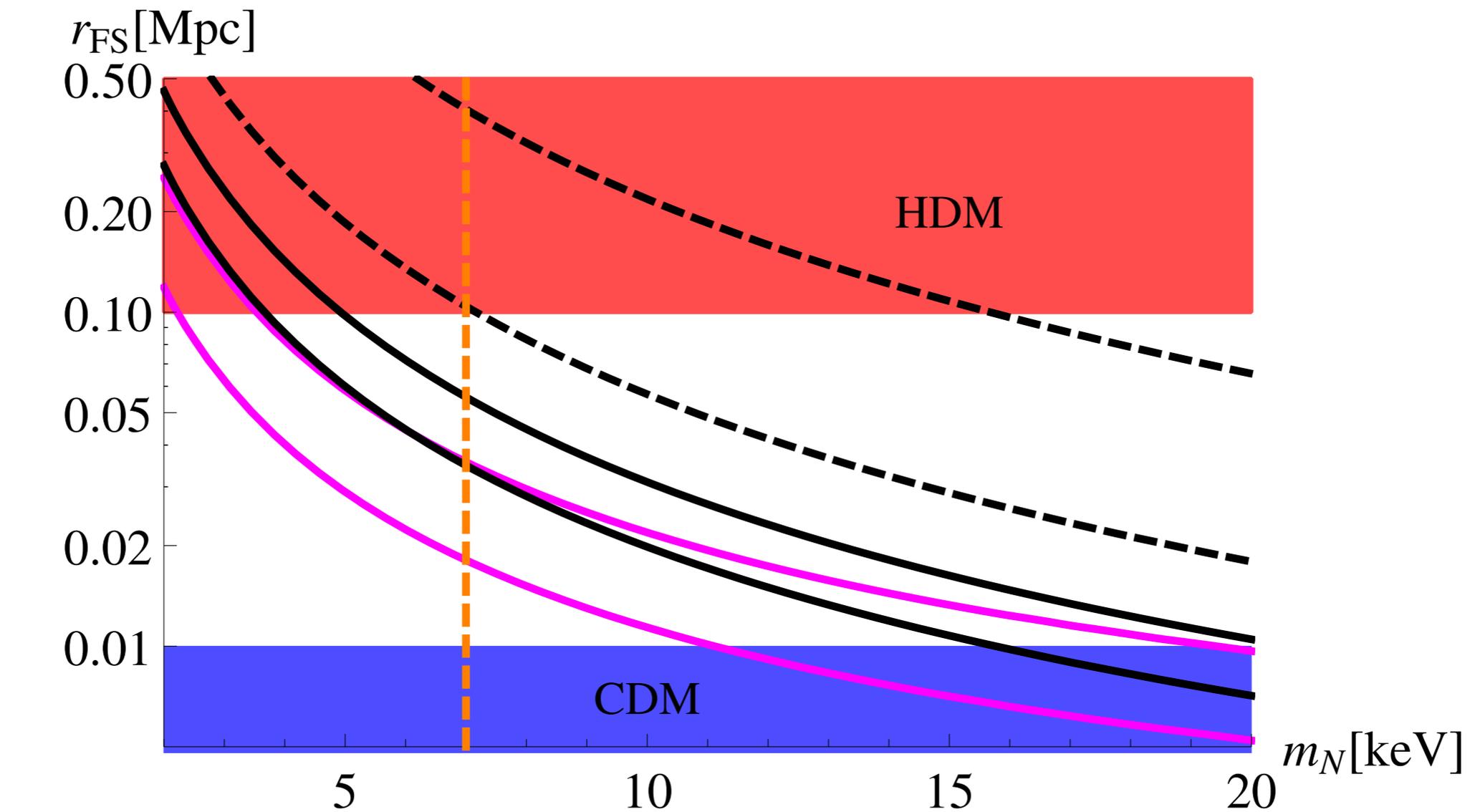
# Free Streaming Horizon

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- ❖ An indicator whether the keV sterile neutrino dark matter is hot, cold or warm dark matter is the free streaming horizon, which is the co-moving mean distance which a collisionless gravitationally unbounded particle can travel.

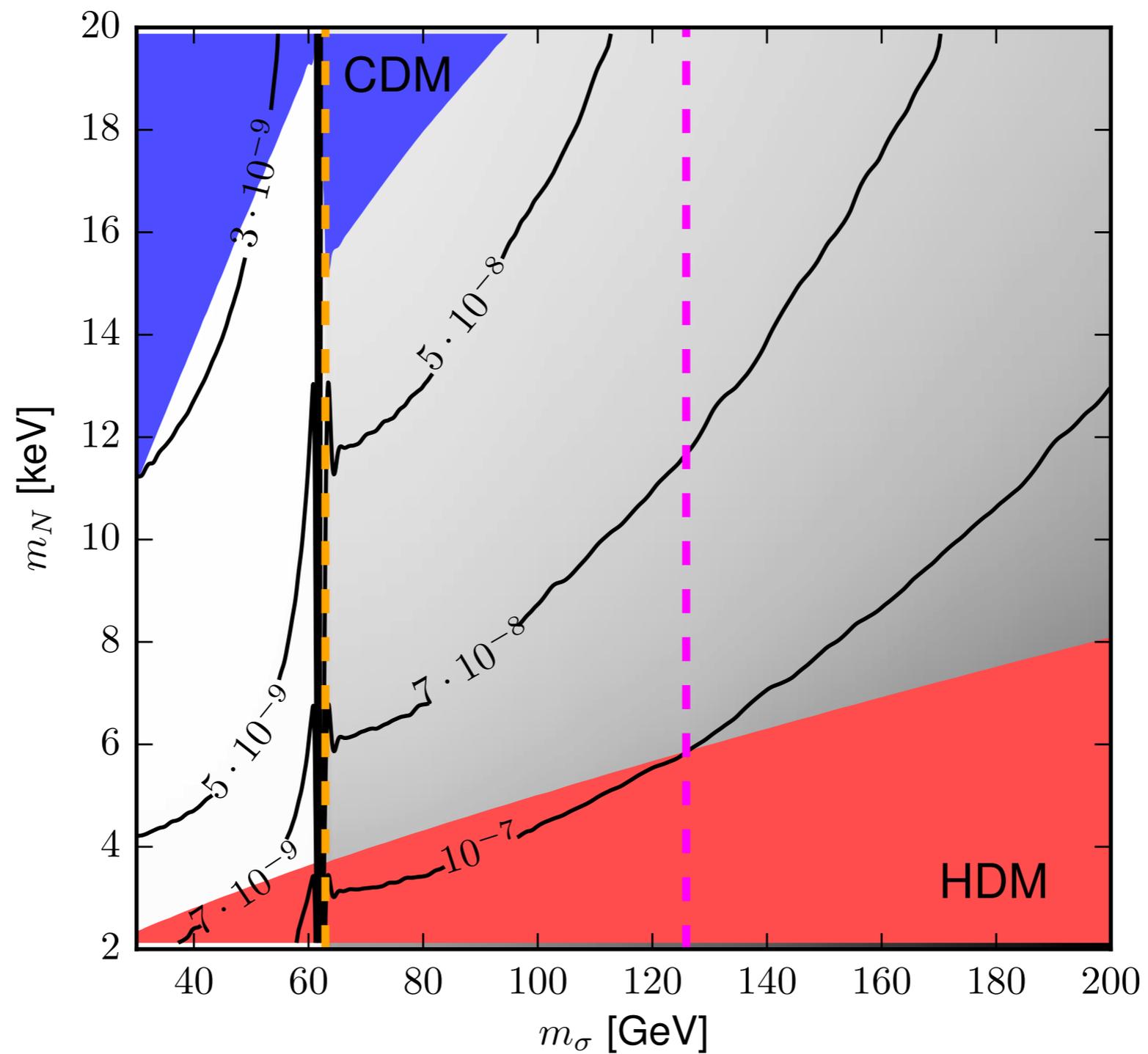
$$r_{\text{FS}} = \int_{t_{\text{in}}}^{t_0} \frac{\langle v(t) \rangle}{a(t)} dt$$

# Free Streaming Horizon



$$m_\sigma = \{30\text{GeV}, 60\text{GeV}, 65\text{GeV}, 100\text{GeV}, 170\text{GeV}, 500\text{GeV}\}$$

# Free Streaming Horizon



# Conclusion

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- ❖ We have investigated the sterile neutrino dark matter production from the frozen-in scalar whose mass is either above or below the SM Higgs mass.
- ❖ Different contributions to the production have been investigated.
- ❖ In case, the frozen-In scalar is lighter than the half of the SM Higgs mass, the SM Higgs can decay into the frozen-in scalar as well, thus providing the additional channel for the production.
- ❖ The lighter the scalar is, the colder the sterile neutrino will be.