Rotating generalization of the Buda-Lund hydrodynamical model

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Introduction	Observables	Summary
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#### Motivation

- Heavy Ion Physics
- Hydrodynamics
  Rotation
- Buda-Lund model [1]

<sup>[1]</sup> T. Csörgő and B. Lörstad, Bose-Einstein correlations for three-dimensionally expanding, cylindrically symmetric, finite systems, Phys.Rev. C54, 1390 (1996)

Buda-Lund model	Observables	Summary
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#### Definitions of the model

- ▶ Relativistic extension of a non-relativistic exact solution [2]
- Observables calculated from source function parametrized with quantities taken at the freeze-out
- Source function:

$$S(x,p)\mathrm{d}^4 x = rac{g_s}{(2\pi)^3}rac{p^\mu\mathrm{d}^4\Sigma_\mu(x)}{B(x,p)+s_q}$$

Maxwell-Jüttner distribution:

$$B(x,p) = \exp\left[\frac{p^{\nu}u_{\nu}(x)}{T(x)} - \frac{\mu(x)}{T(x)}\right]$$

► Cooper-Frye prefactor:  $p^{\mu} \mathrm{d}^4 \Sigma_{\mu}(x) = p^{\mu} u_{\mu} H(\tau) \mathrm{d}^4 x$ 

<sup>[2]</sup> M. Csanád, T. Csörgő, and B. Lörstad, Buda-Lund hydro model for ellipsoidally symmetric fireballs and the elliptic flow at RHIC, Nucl.Phys. A742, 80 (2004)

Buda-Lund model	Observables	Summary
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### The rotating velocity field

New parameter  $\omega$  describing rotation.  $\omega = 0$  is the original model.

▶ The four-velocity based on the non-relativistic solution [3]

$$\mathbf{v}_{\text{rot}} = \begin{pmatrix} \mathbf{r}_{x} \frac{\dot{X}}{X} + \mathbf{r}_{z} \frac{\omega}{Z} \frac{X + Z}{2} \\ \dot{Y} \\ \mathbf{r}_{y} \frac{\dot{Y}}{Y} \\ \mathbf{r}_{z} \frac{\dot{Z}}{Z} - \mathbf{r}_{x} \frac{\omega}{X} \frac{X + Z}{2} \end{pmatrix} \rightarrow u^{\mu} = \begin{pmatrix} \mathbf{G} \\ \mathbf{r}_{x} \frac{\dot{X}}{X} + \mathbf{r}_{z} \frac{\omega}{Z} \frac{X + Z}{2} \\ \dot{Y} \\ \mathbf{r}_{y} \frac{\dot{Y}}{Y} \\ \mathbf{r}_{z} \frac{\dot{Z}}{Z} - \mathbf{r}_{x} \frac{\omega}{X} \frac{X + Z}{2} \end{pmatrix}$$

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$$u_{\mu}u^{\mu} = 1 \Rightarrow G$$

• Normalization convention to provide v < 1

<sup>[3]</sup> M. I. Nagy and T. Csörgő, An analytic hydrodynamical model of rotating 3D expansion in heavy-ion collisions, arXiv:1512.00888 [nucl-th].

Buda-Lund model	Observables	Summary
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#### Temperature and chemical potential

Scaling variable for ellipsodial symmetry

$$s = rac{r_x^2}{2X^2} + rac{r_y^2}{2Y^2} + rac{r_z^2}{2Z^2} \; ,$$

Chemical potential

$$\frac{\mu(x)}{T(x)}=\frac{\mu_0}{T_0}-b^2s ,$$

Temperature

$$\frac{1}{T(x)} = \frac{1}{T_0} \left( 1 + a^2 s \right) \left( 1 + d^2 \frac{(\tau - \tau_0)^2}{2\Delta \tau^2} \right)$$

► 15 free parameters:  $\tau_0$ ,  $\Delta \tau$ , X, Y, Z,  $\dot{X}$ ,  $\dot{Y}$ ,  $\dot{Z}$ ,  $\omega$ ,  $\vartheta$ ,  $\mu_0$ ,  $T_0$ ,  $b^2$ ,  $a^2$  and  $d^2$ 

Buda-Lund model	Observables	Summary
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#### Saddle-point integration

Sharp distribution

$$S_0(x,p) := rac{H( au)}{B(x,p) + s_q}$$

Approximation with Gaussian

$$S(x,p) \approx \frac{g_s}{(2\pi)^3} \rho_{\mu} u^{\mu}(x_s) S_0(x_s,p) \exp\left\{-\frac{1}{2} R_{\mu\nu}^{-2} (x-x_s)^{\mu} (x-x_s)^{\nu}\right\}$$

The covariance matrix

$$R_{\mu\nu}^{-2} = \partial_{\mu}\partial_{\nu}\left[-\ln(S_0(x_s, p))\right]$$

	Observables	

### Definitions of observables

Invariant momentum distribution

$$N_1(\mathbf{p}) = \int S(x,p) \mathrm{d}^4 x$$

Elliptic flow

$$v_{2} = \frac{\int_{0}^{2\pi} N_{1}(p_{z}, p_{t}, \varphi) \cos(2\varphi) \mathrm{d}\varphi}{\int_{0}^{2\pi} N_{1}(p_{z}, p_{t}, \varphi) \mathrm{d}\varphi}$$

Bose-Einstein correlation and HBT radii

$$C(\mathbf{q}_{osl},p) = 1 + \lambda_* \exp\left(-\sum_{i,j} R_{ij}^2(p) q_i q_j\right)$$
  $i,j = o, s, l$ 

	Observables ○●○○○○	

Elliptic flow



The elliptic flow as a function of transverse momentum, at  $p_z = 0$  for different  $\vartheta$  and  $\omega$  parameters

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	Observables 000000	

#### Elliptic flow



The elliptic flow as a function of pseudorapidity, at  $p_t = 300 \,\text{MeV}$  for different  $\vartheta$  and  $\omega$  parameters

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#### Elliptic flow



The elliptic flow as a function of  $\vartheta$  and  $\omega$ , at  $p_t = 300 \text{ MeV}, p_z = 0$ 

	Observables 0000●0	



The HBT radii  $(R_{ij}^2 \text{ [fm}^2))$  without rotation, at  $p_z = 0$  as a function of transverse mass  $(m_t \text{ [MeV]})$  and the azimuth angle of the momentum  $(\varphi \text{ [rad]})$ . The angle of the rotated ellipsoid  $\vartheta = 0$ .

	Observables 0000●0	



The HBT radii ( $R_{ij}^2$  [fm<sup>2</sup>]) without rotation, at  $p_z = 0$  as a function of transverse mass ( $m_t$  [MeV]) and the azimuth angle of the momentum ( $\varphi$  [rad]). The angle of the rotated ellipsoid  $\vartheta = 0.05$ .

	Observables 0000●0	



The HBT radii ( $R_{ij}^2$  [fm<sup>2</sup>]) with rotation ( $\omega = 0.01$ ), at  $p_z = 0$  as a function of transverse mass ( $m_t$  [MeV]) and the azimuth angle of the momentum ( $\varphi$  [rad]). The angle of the rotated ellipsoid  $\vartheta = 0$ .

	Observables 0000●0	



The HBT radii ( $R_{ij}^2$  [fm<sup>2</sup>]) with rotation ( $\omega = 0.01$ ), at  $p_z = 0$  as a function of transverse mass ( $m_t$  [MeV]) and the azimuth angle of the momentum ( $\varphi$  [rad]). The angle of the rotated ellipsoid  $\vartheta = 0.05$ .

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	Observables 000000	Summary ●

#### Summary:

- Extension of non-relativistic rotating solution to a relativistic parametrization via the Buda-Lund model
- Generalization of the model taking into account rotation
- Calculation of observables using saddle-point approximation

Outlook:

- Comparison with experimental data
- Other flow components  $(v_1, v_3 \dots)$
- Analytic approximations

Buda-Lund model	Observables	Summary
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# Thank you for your attention!