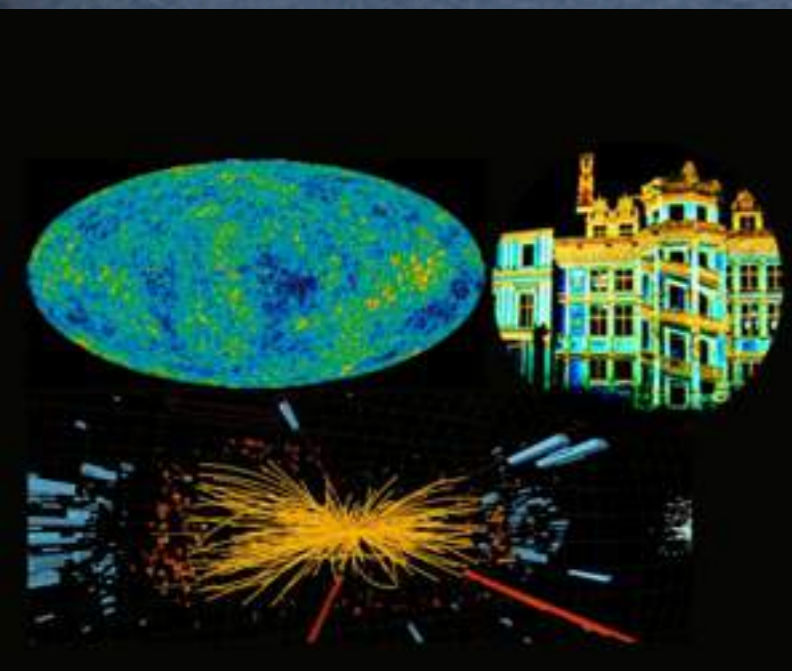




Theoretical Models for the Neutrino Mass and Mixing Pattern

Steve King
2nd June 2016

28th Rencontres de Blois
Particle Physics and Cosmology





Neutrino Mass & Mixing



Three things we have learned from the neutrino experiments

1. Neutrinos have tiny masses $\ll m_e$
2. Neutrinos mix a lot (unlike the quarks)
3. First (and so far only) new physics Beyond the Standard Model

Lepton Mixing Angles

$$\theta_{12} = 34^\circ \pm 1^\circ$$

$$\theta_{23} = 42^\circ \pm 3^\circ$$

$$\theta_{13} = 8.4^\circ \pm 0.3^\circ$$

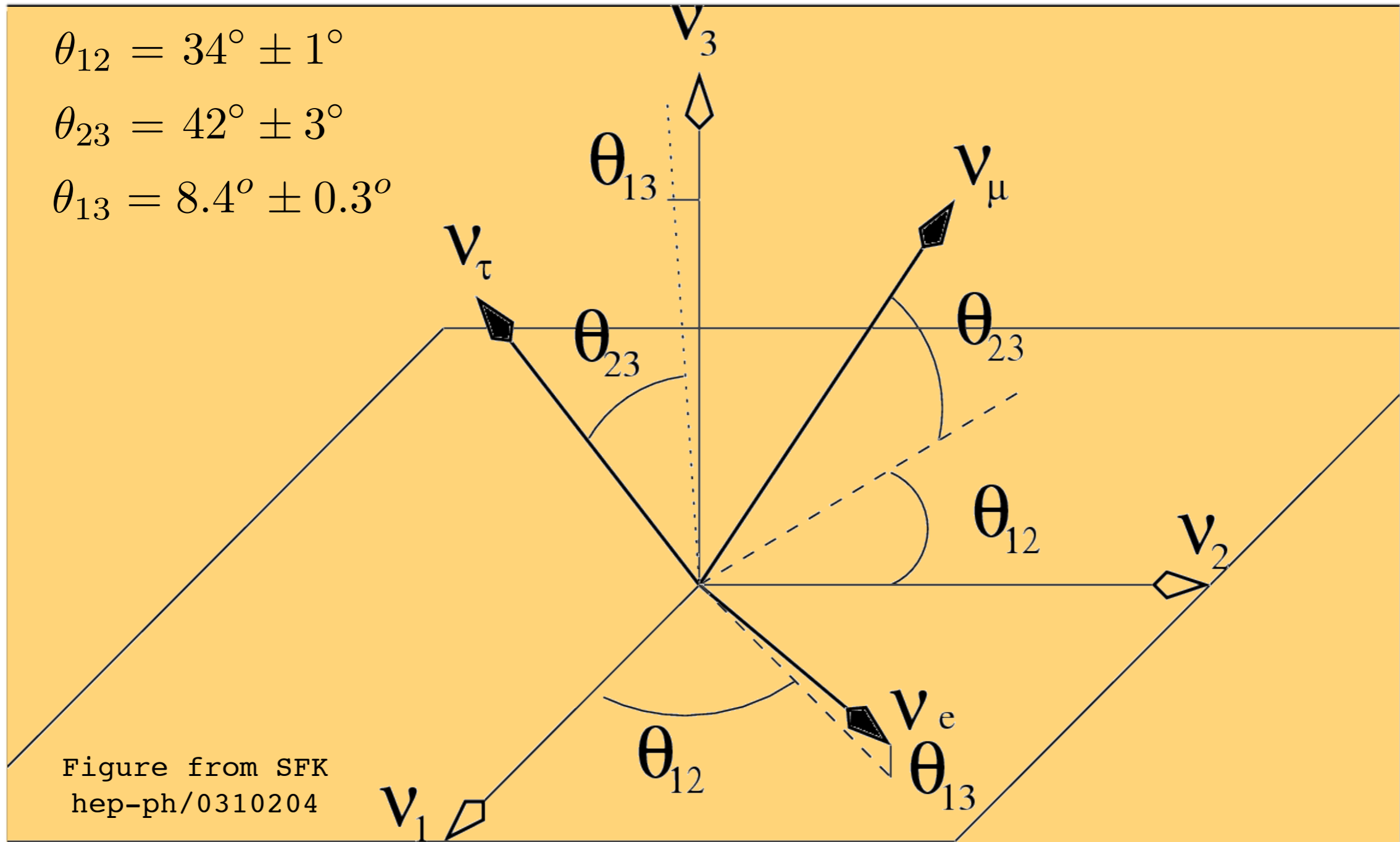
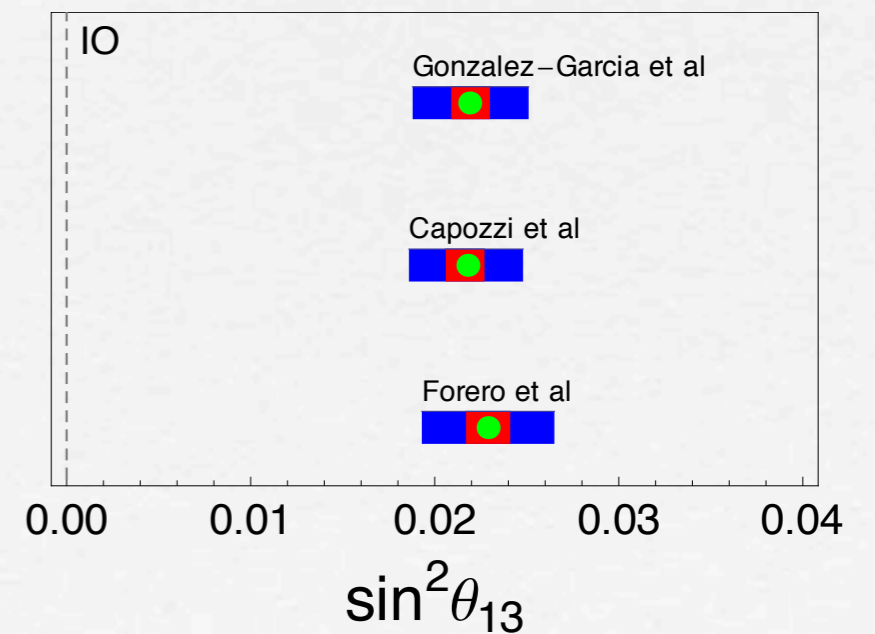
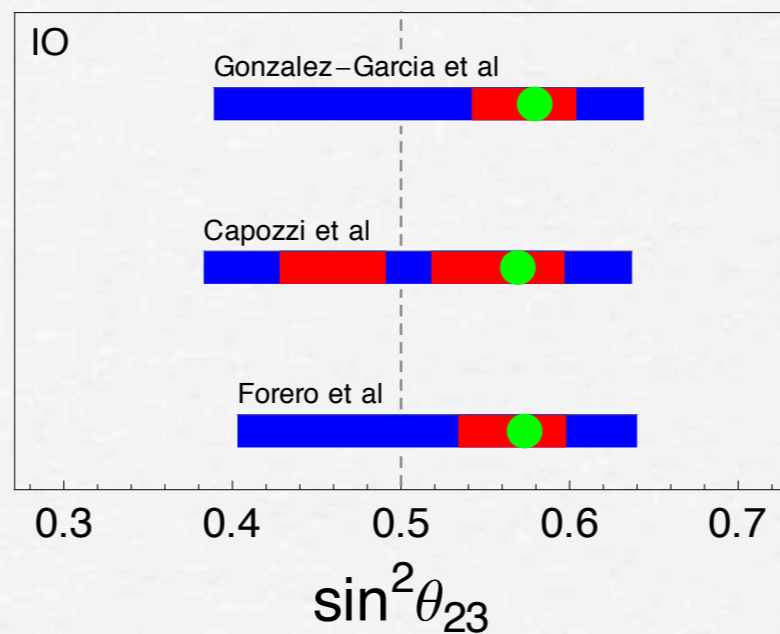
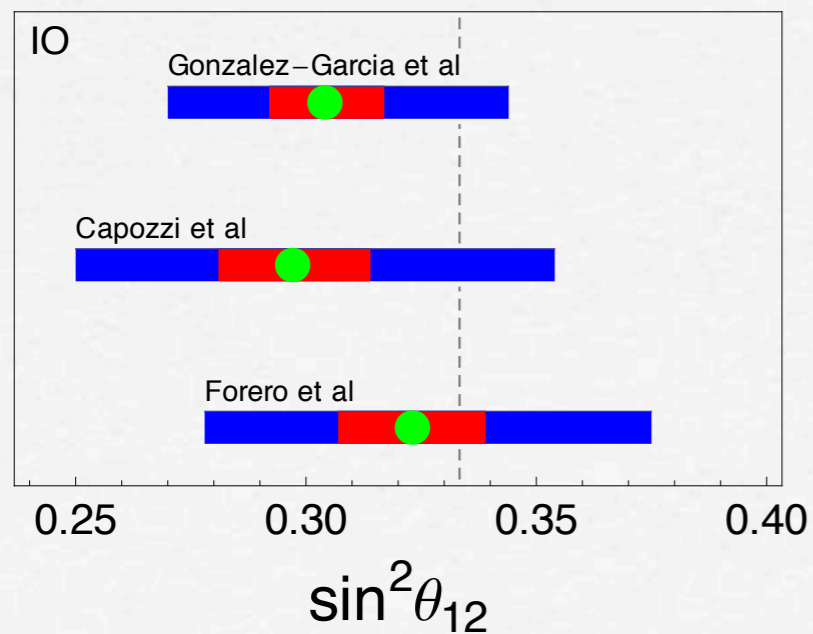
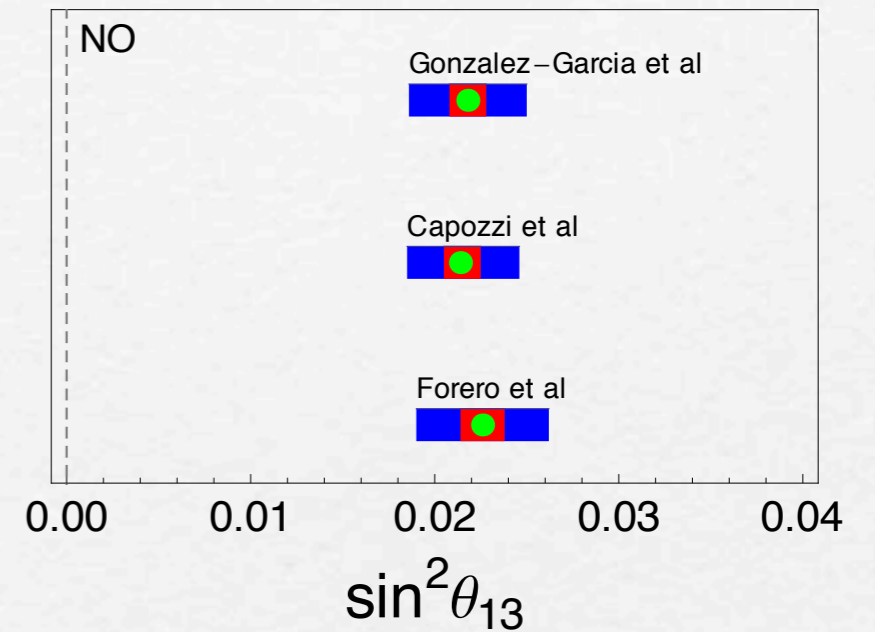
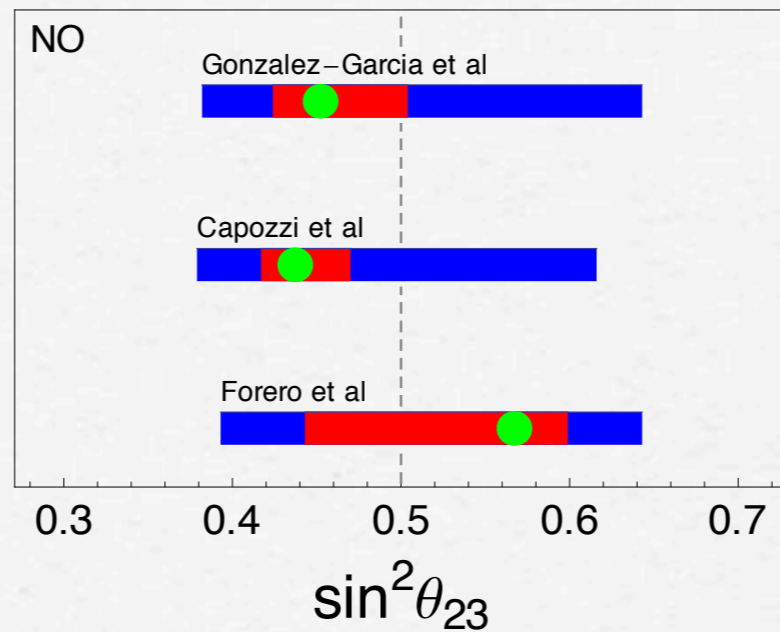
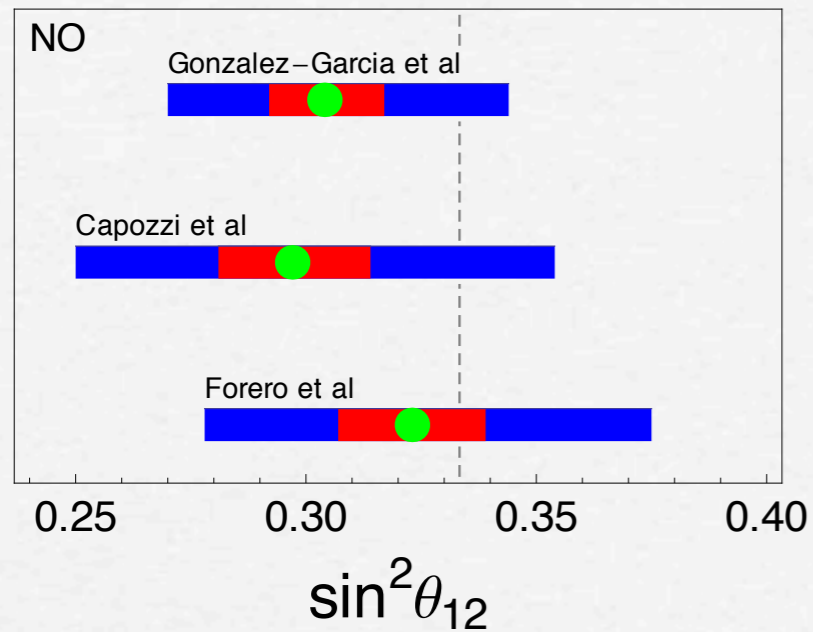


Figure from SFK
hep-ph/0310204

Gonzalez-Garcia et al = Gonzalez-Garcia, Maltoni, Salvado, Schwetz
 Fogli et al = Capozzi, Fogli, Lisi, Marrone, Montanino, Palazzo
 Forero et al = Forero, Tortola, Valle

Lepton Mixing

Plots from updated review by SFK and Luhn

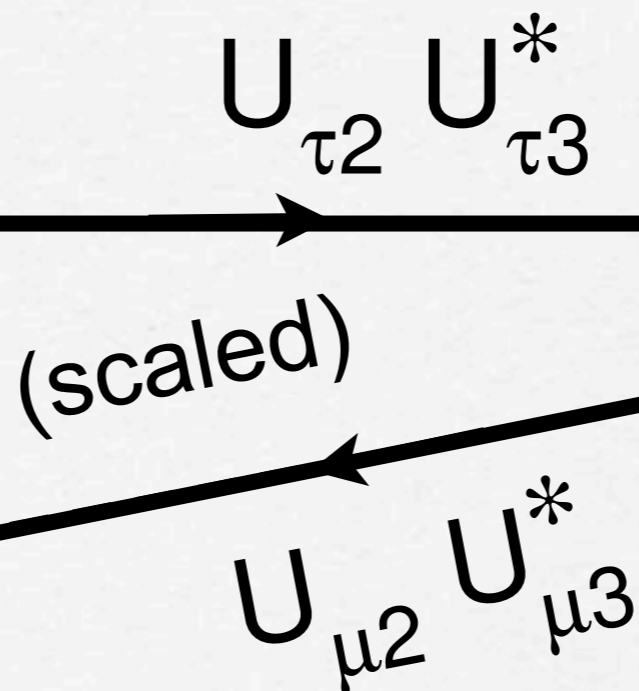
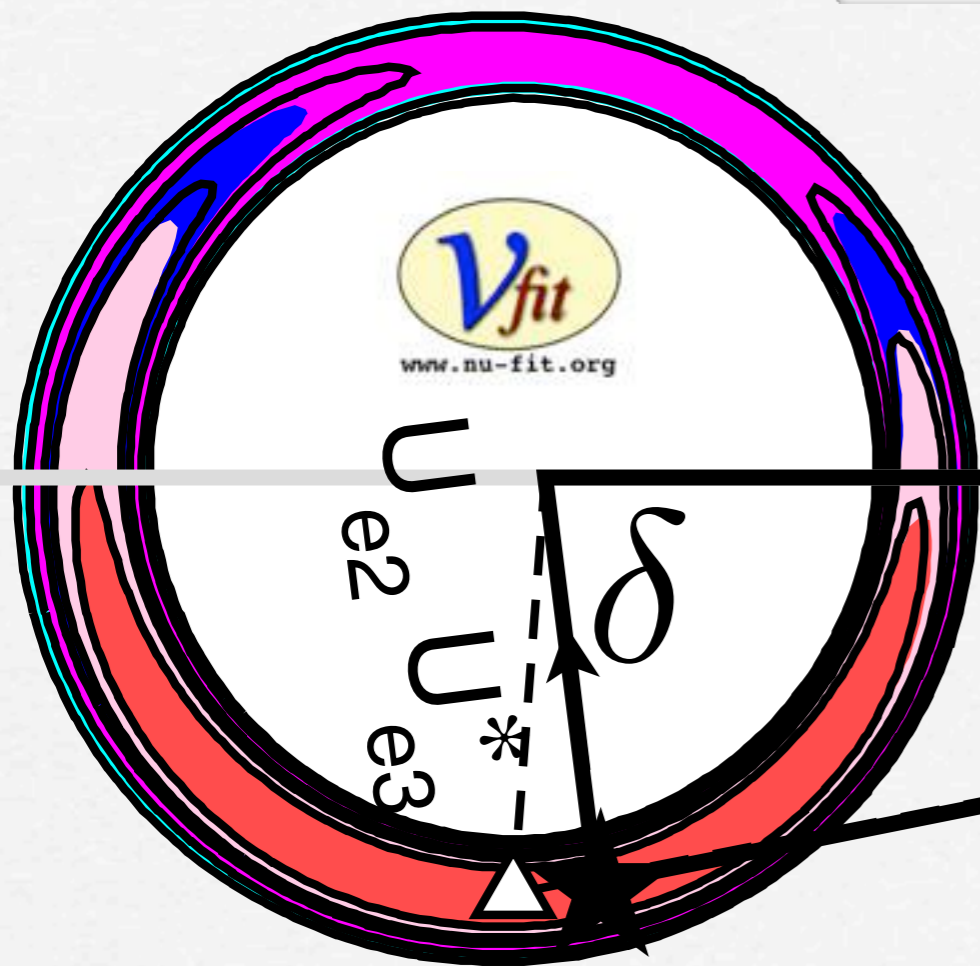


CP phase

$$\delta \sim -90^\circ?$$

$$\frac{U_{e2}U_{e3}^*}{U_{\tau2}U_{\tau3}^*} \approx -r e^{i\delta}$$

$$\frac{U_{\mu2}U_{\mu3}^*}{U_{\tau2}U_{\tau3}^*} \approx -1 + r e^{i\delta}$$



Scaled triangle independent of s, a to LO

arXiv:0710.0530

$$\sin \theta_{12} = \frac{1}{\sqrt{3}}(1 + s),$$

$$\sin \theta_{23} = \frac{1}{\sqrt{2}}(1 + a),$$

$$\sin \theta_{13} = \frac{r}{\sqrt{2}}$$

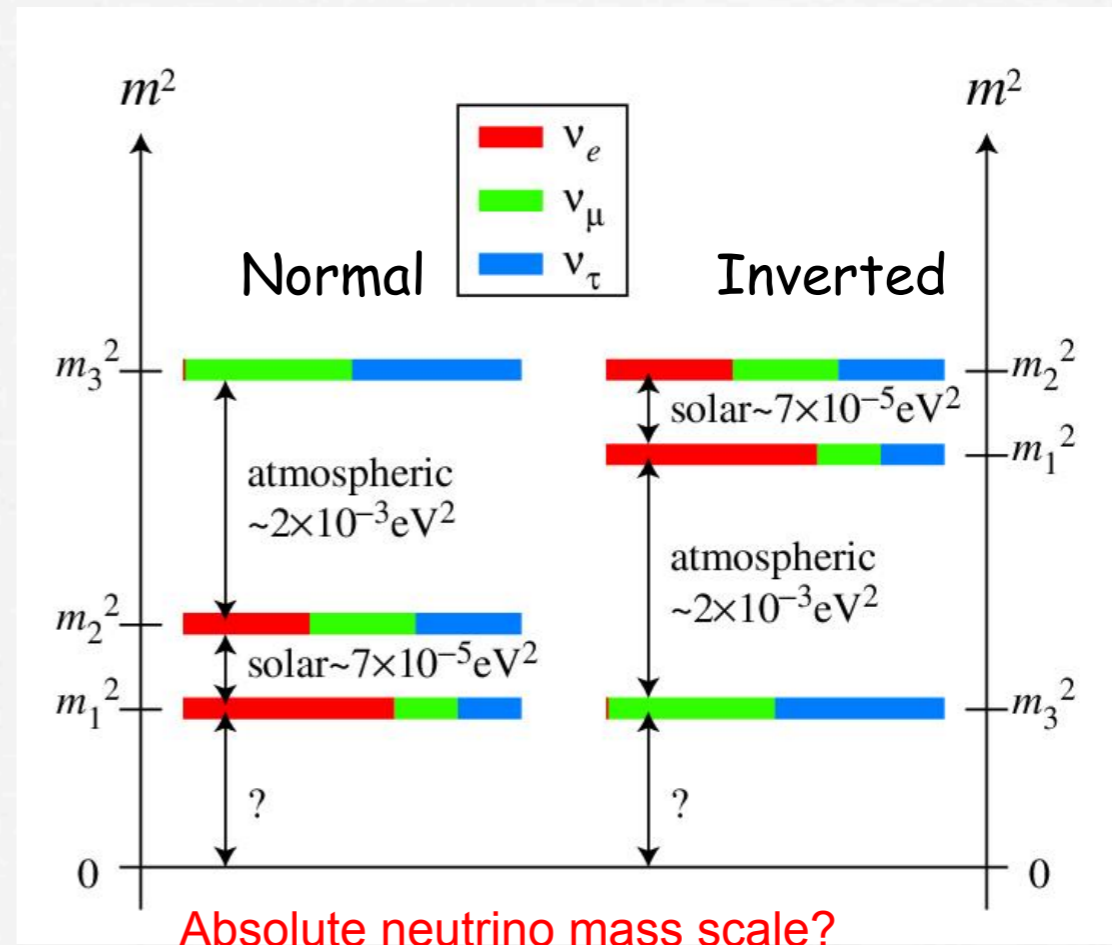
$s = \text{solar}$

$a = \text{atmospheric}$

$r = \text{reactor}$

What is the Origin of Small Neutrino Masses?

Masses?



Majorana

$$m^\nu \bar{\nu}_L \nu_L^c$$

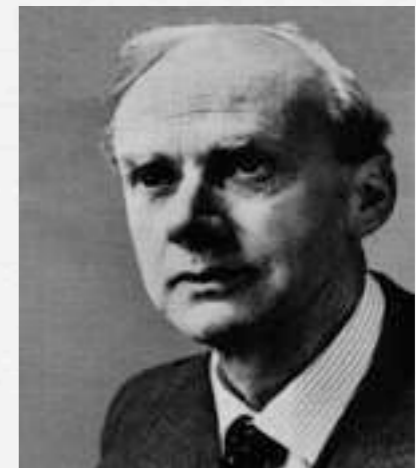
Violates L

or

$$m^\nu \bar{\nu}_L \nu_R$$

Conserves L

Dirac



The Electron Mass



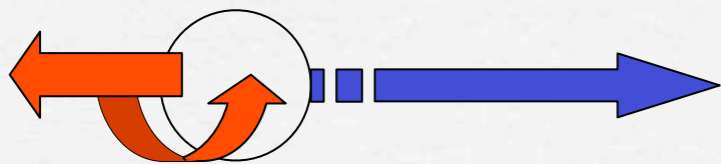
Left-handed
electron

Right-handed
electron

Neutrino Mass

Left-handed
neutrino

ν_L

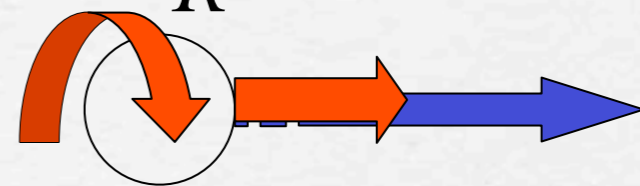


Dirac

$$m_{LR}^{\nu} \bar{\nu}_L \nu_R$$

Right-handed
neutrino

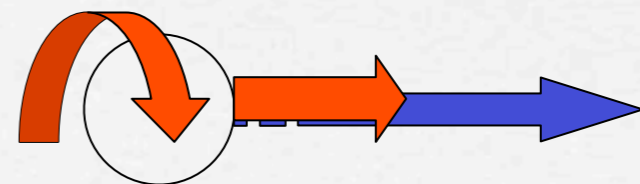
ν_R



Majorana

$$m_{LL}^{\nu} \bar{\nu}_L \nu_L^c$$

ν_L^c

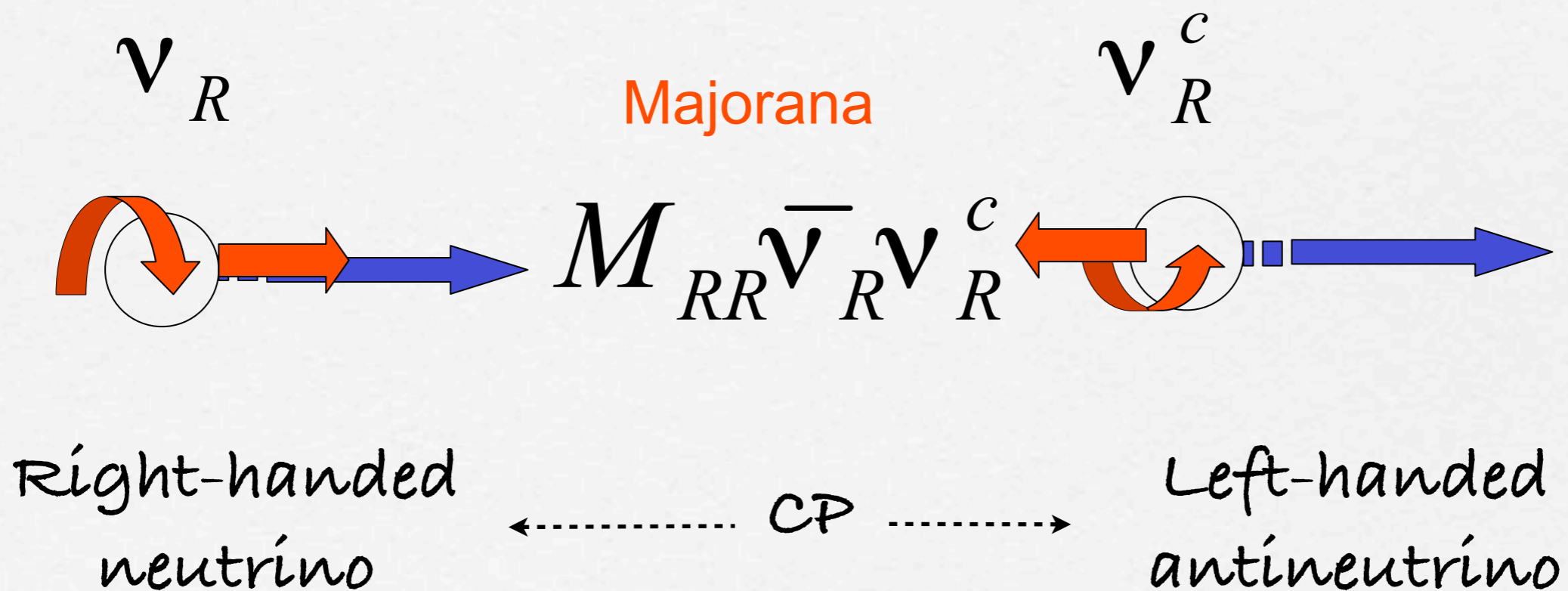


Left-handed
neutrino

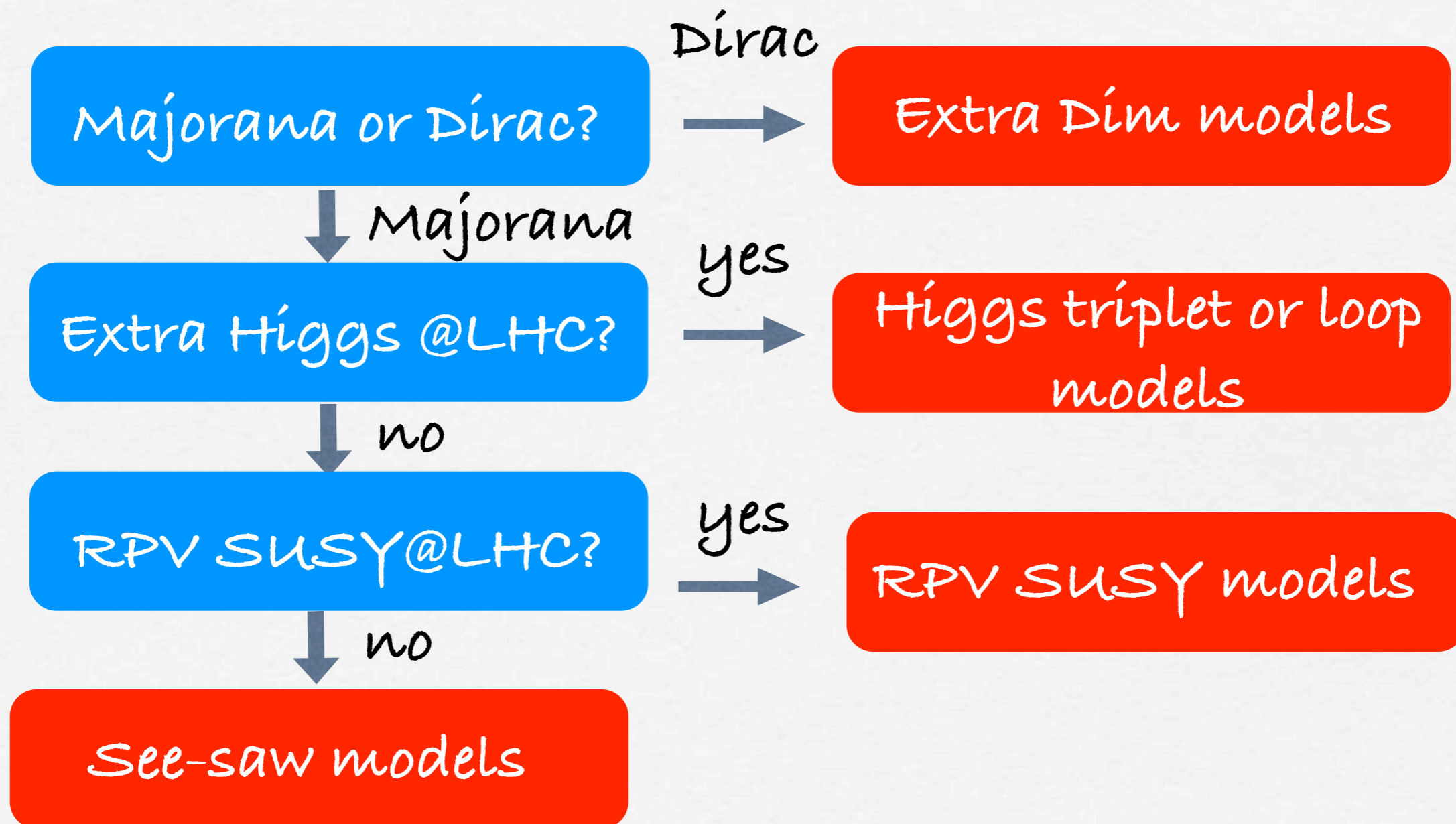
← CP →

Right-handed
antineutrino

Right-handed neutrino mass

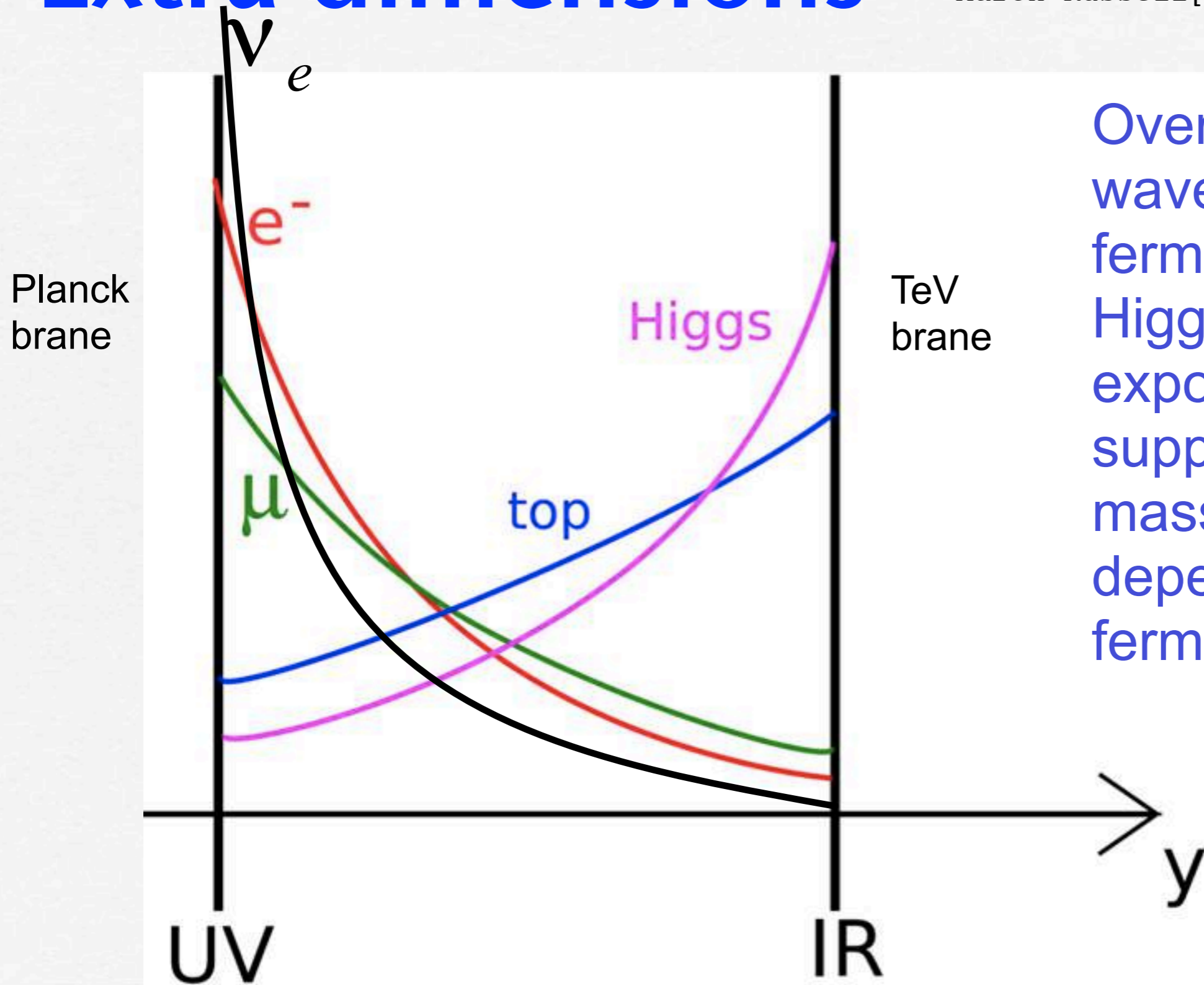


Roadmap of neutrino mass models



Extra dimensions

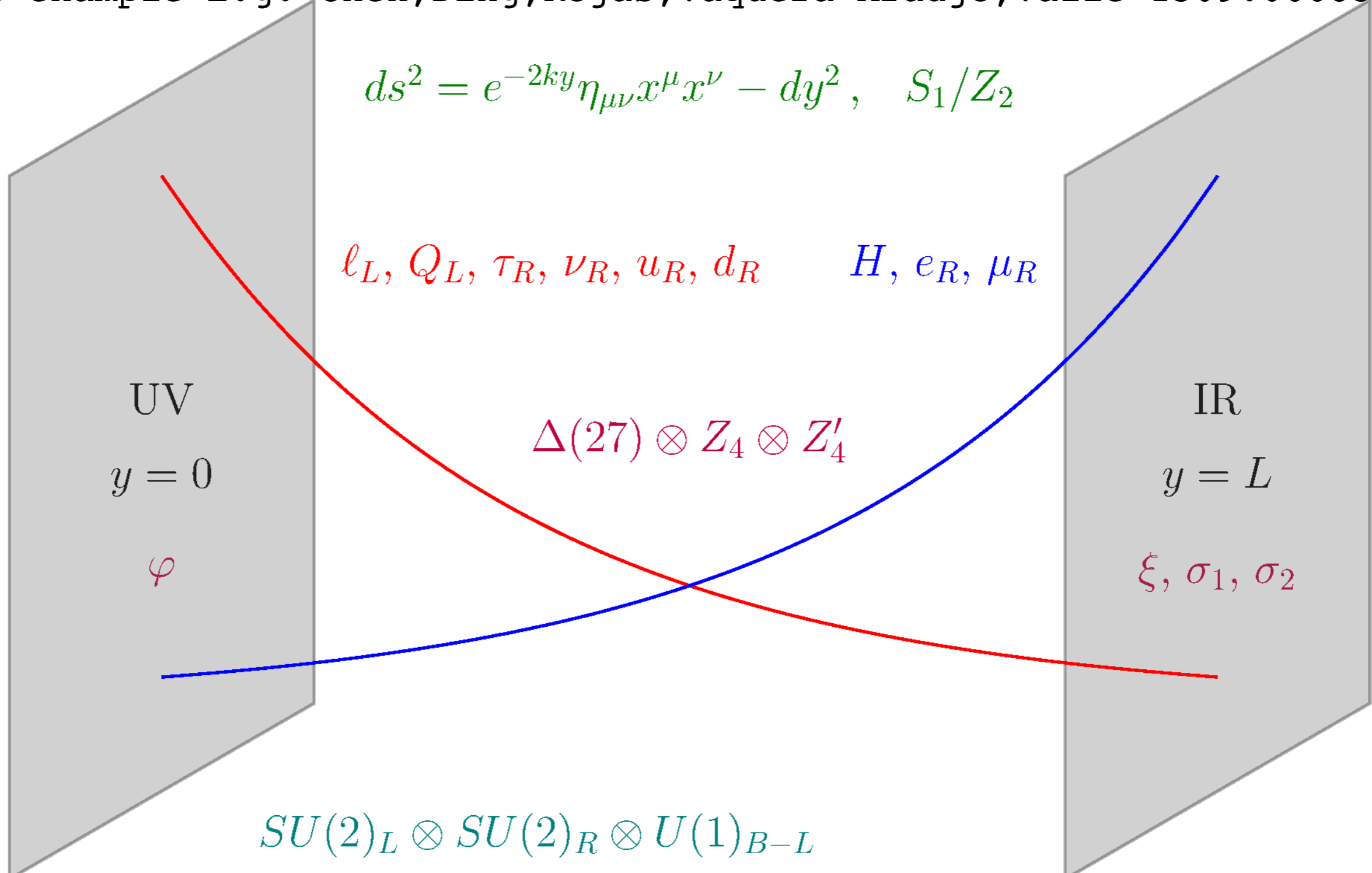
Arkani-Hamed, Dimopoulos, Dvali,
March-Russell [hep-ph/9811448]...



Overlap
wavefunction of
fermions with
Higgs gives
exponentially
suppressed Dirac
masses,
depending on the
fermion profiles

Recent example E.g. Chen, Ding, Rojas, Vaquera-Araujo, Valle 1509.06683

$$ds^2 = e^{-2ky} \eta_{\mu\nu} x^\mu x^\nu - dy^2, \quad S_1/Z_2$$



UV
 $y = 0$
 φ

$\ell_L, Q_L, \tau_R, \nu_R, u_R, d_R$ H, e_R, μ_R

$\Delta(27) \otimes Z_4 \otimes Z'_4$

IR
 $y = L$
 ξ, σ_1, σ_2

$SU(2)_L \otimes SU(2)_R \otimes U(1)_{B-L}$

$SU(2)_L \otimes U(1)_Y$

$SU(2)_D \otimes U(1)_{B-L}$

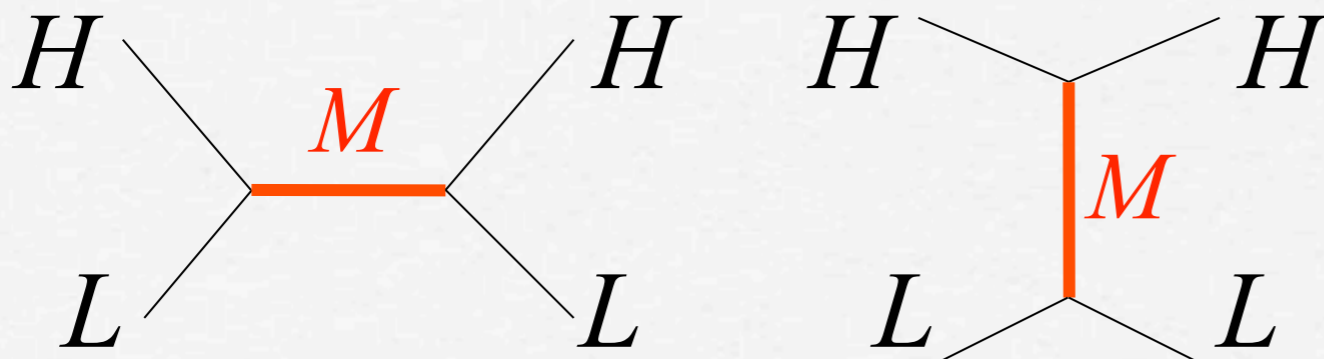
Majorana Neutrino Mass

Renormalisable $\Delta L = 2$ operator $\lambda_\nu LL\Delta$ where Δ is light Higgs triplet with $VEV < 8\text{GeV}$ from ρ parameter

Non-renormalisable $\Delta L = 2$ operator $\frac{\lambda_\nu}{M} LLHH = \frac{\lambda_\nu}{M} \langle H^0 \rangle^2 \bar{\nu}_{eL} \nu_{eL}^c$ Weinberg

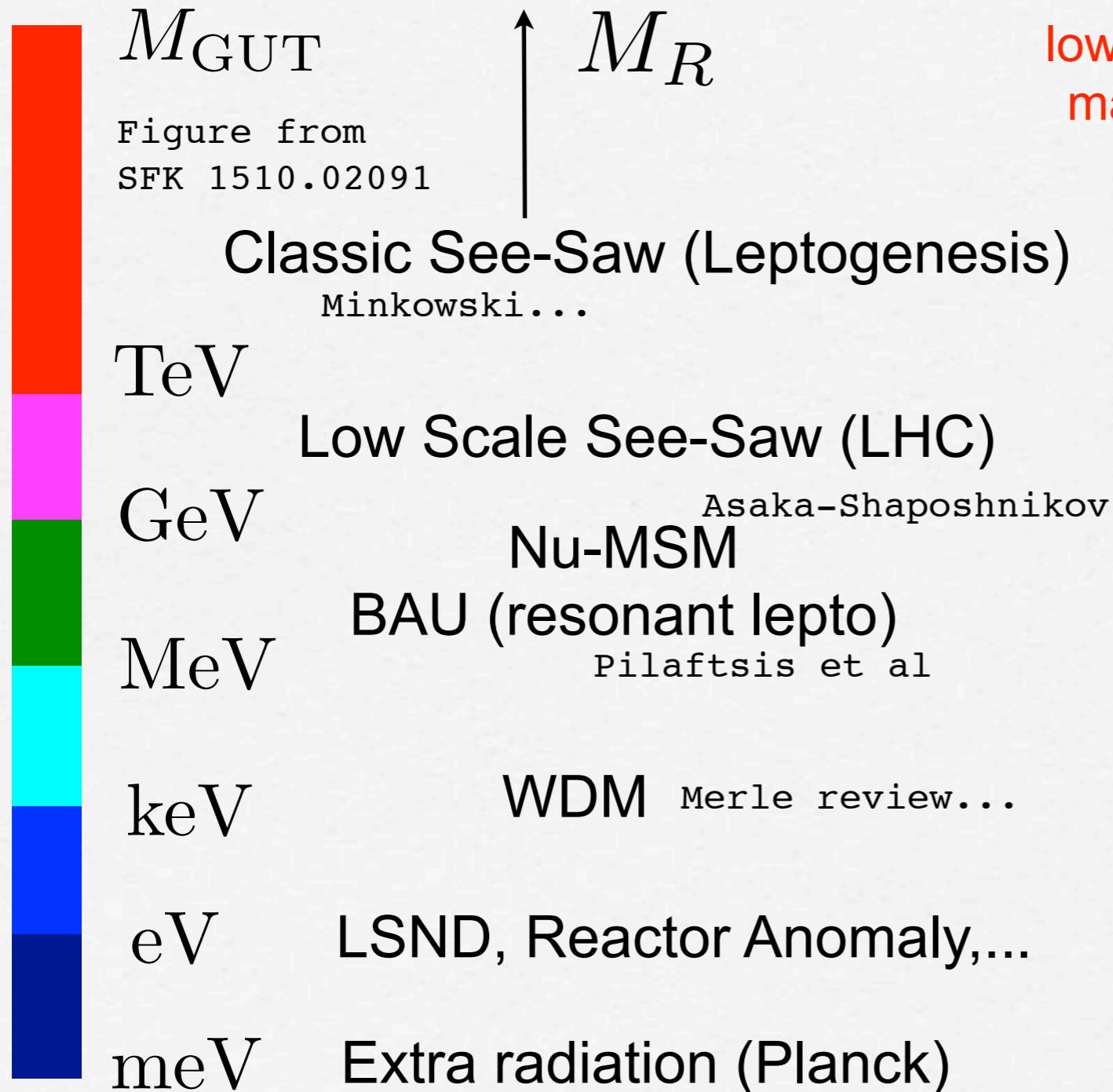
This is nice because it gives naturally small Majorana neutrino masses $m_{LL} \sim \langle H^0 \rangle^2 / M$ where M is some high energy scale

The high mass scale can be associated with some heavy particle of mass M being exchanged (can be singlet or triplet)



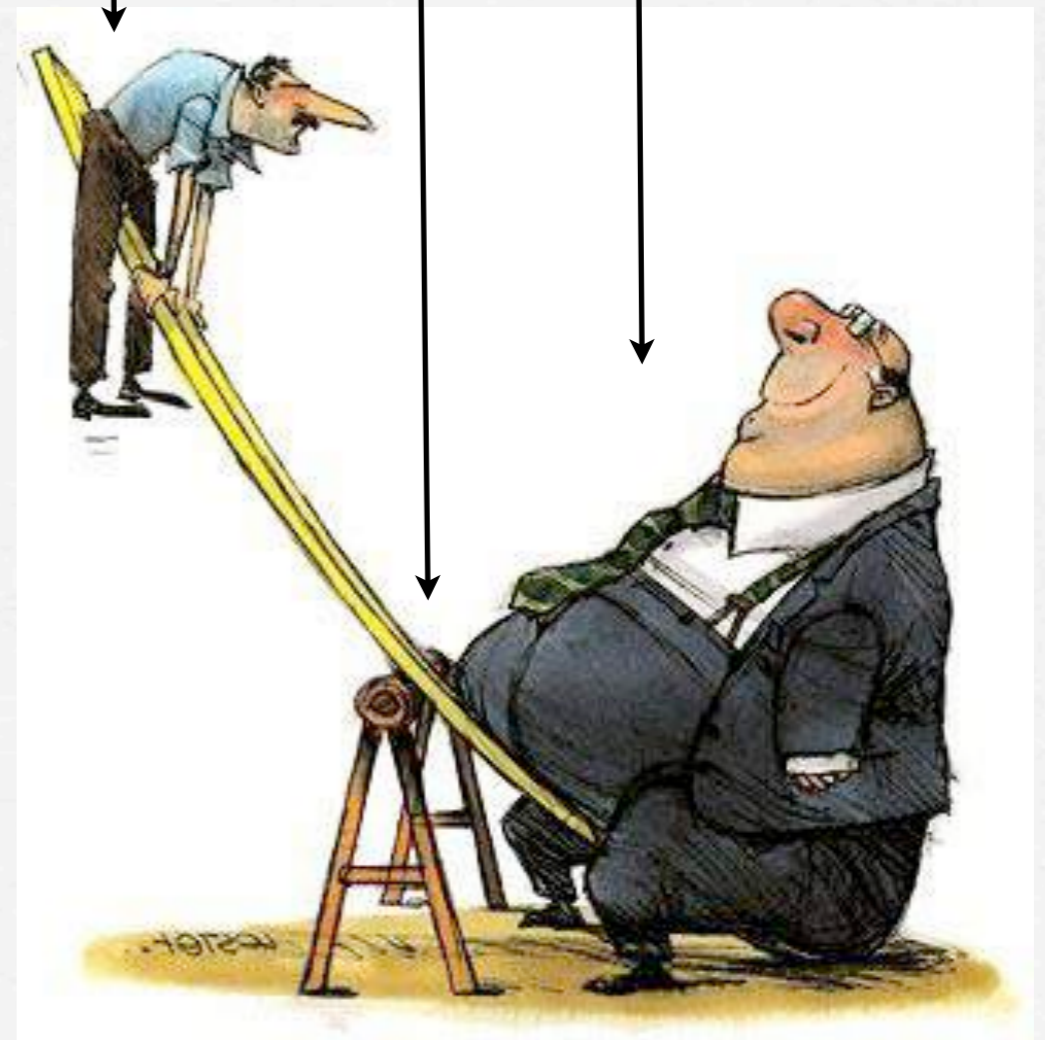
e.g. see-saw with RH neutrinos...introduces a new RH mass spectrum

The RH neutrino Spectrum



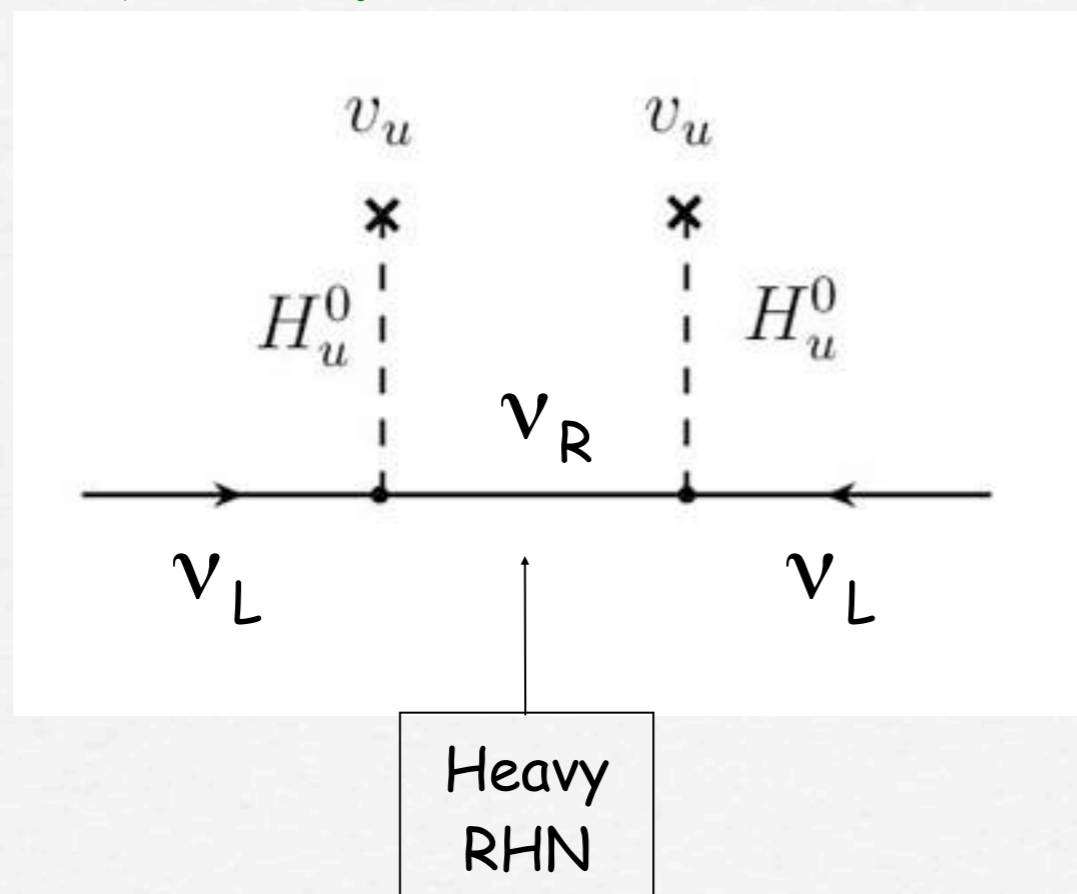
low energy nu mass matrix

$$m^\nu = -m_D \frac{1}{M_R} m_D^T$$



Type I see-saw mechanism

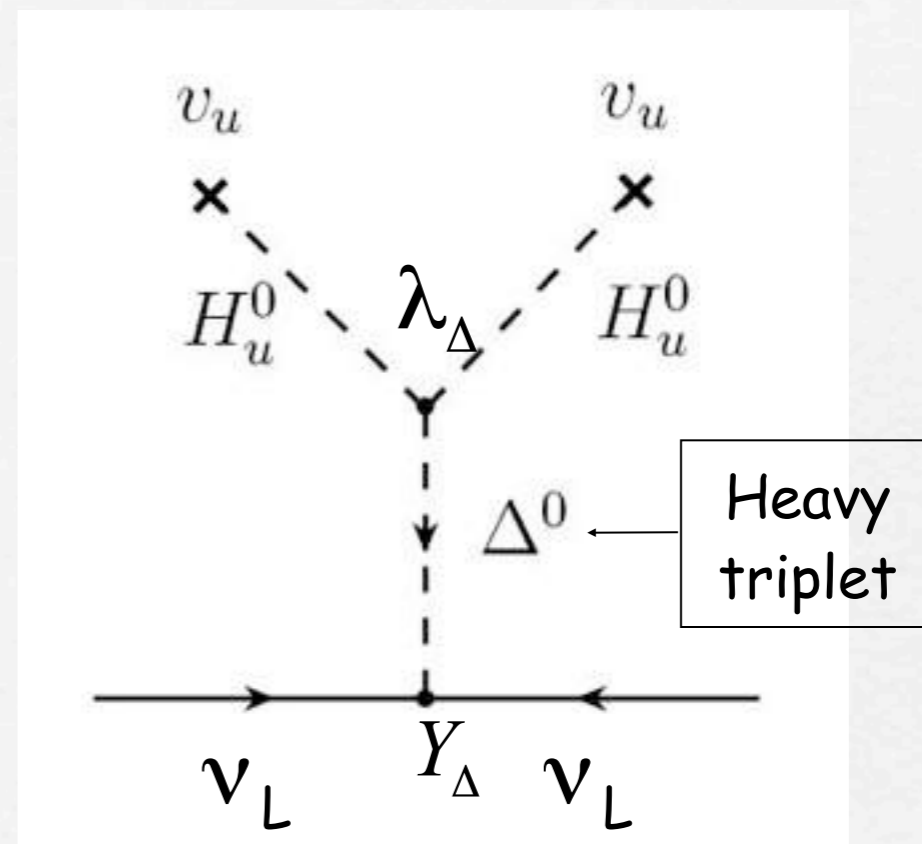
P. Minkowski (1977); Glashow; Yanagida;
Gell-Mann, Ramond, Slansky;
Mohapatra, Senjanovic; Schechter and Valle



$$m_{LL}^I \approx -m_{LR} M_{RR}^{-1} m_{LR}^T$$

Type II see-saw mechanism (SUSY)

Lazarides, Magg, Mohapatra, Senjanovic,
Shafi, Wetterich, Schechter and Valle



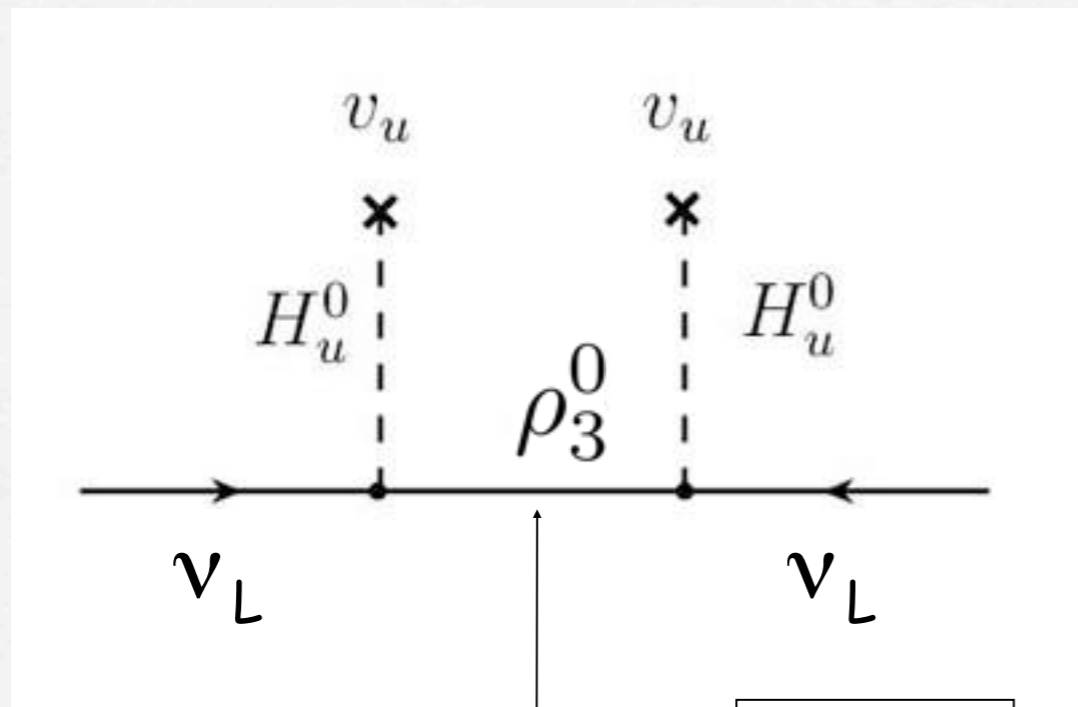
$$m_{LL}^{II} \approx \lambda_\Delta Y_\Delta \frac{v_u^2}{M_\Delta}$$

Type III see-saw mechanism

Foot, Lew, He, Joshi; Ma...

Supersymmetric adjoint SU(5)

Perez et al; Cooper, SFK, Luhn,...



SU(2)_L fermion triplet

$$M_\rho \rho \rho$$

$$m_{LL}^{III} \approx -m_{LR} M_\rho^{-1} m_{LR}^T$$

Type III

See-saw mechanisms with extra singlets S

Inverse see-saw

Wyler, Wolfenstein; Mohapatra, Valle

$$\begin{pmatrix} 0 & M_D & 0 \\ M_D^T & 0 & M \\ 0 & M^T & \mu \end{pmatrix} \quad M_\nu = M_D M^{T^{-1}} \mu M^{-1} M_D^T$$

$M \approx \text{TeV} \rightarrow \text{LHC}$

Linear see-saw

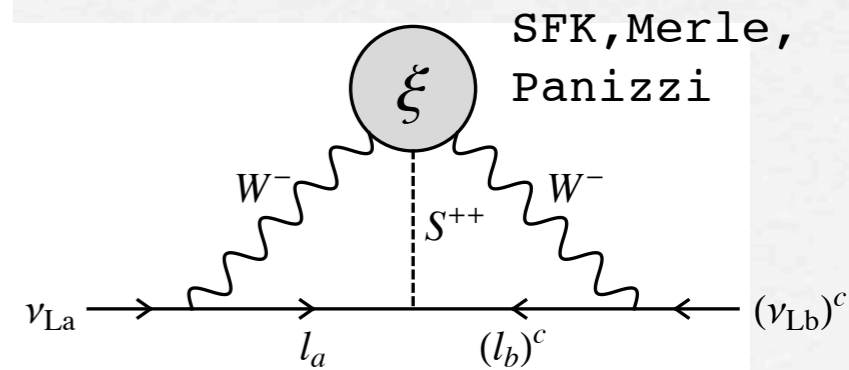
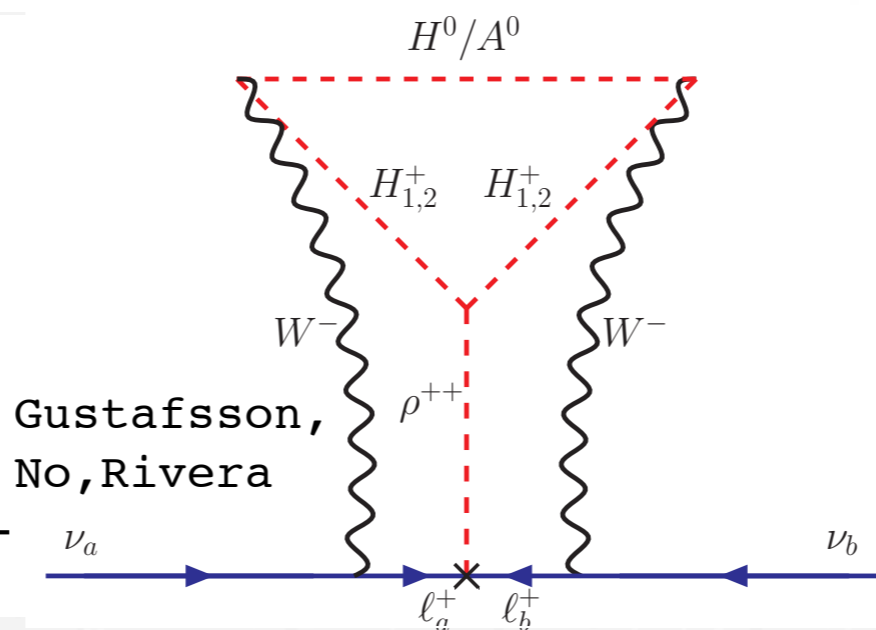
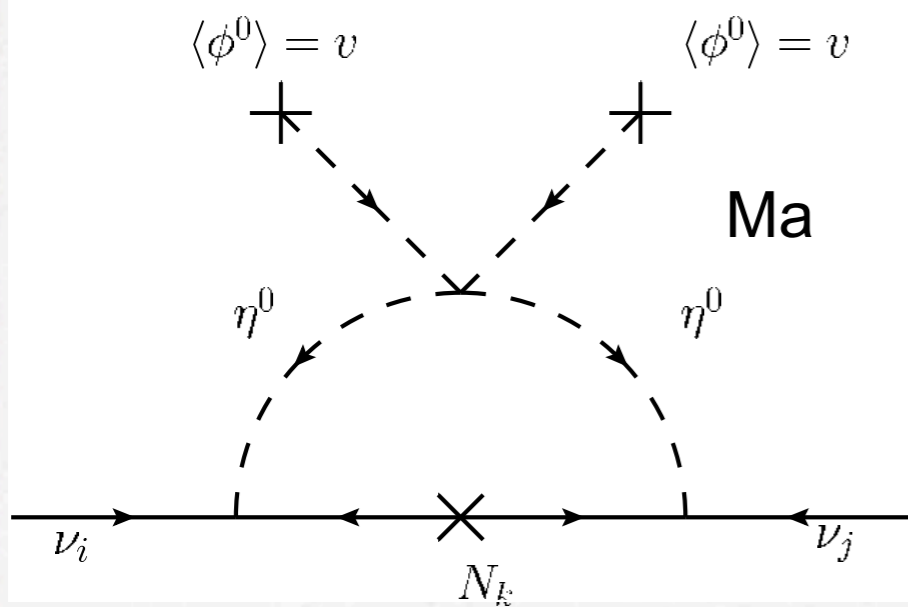
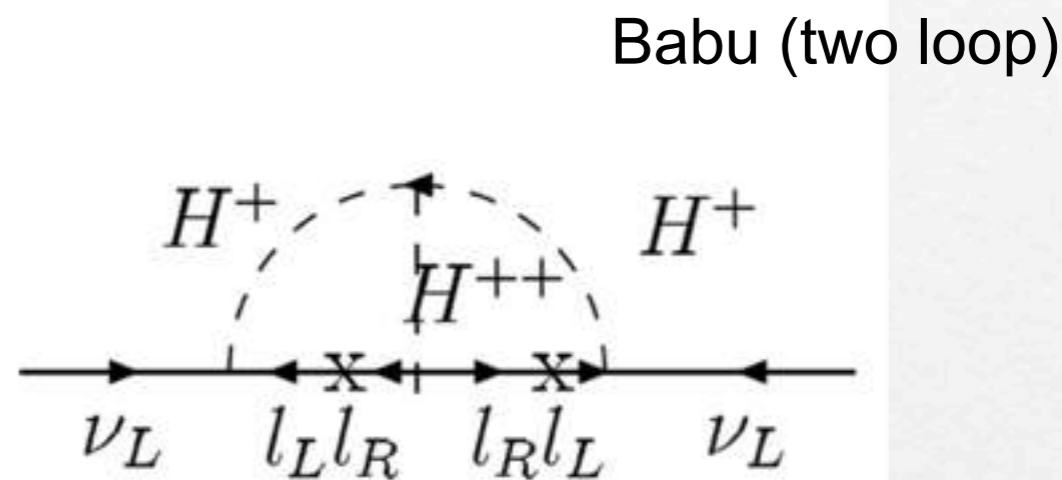
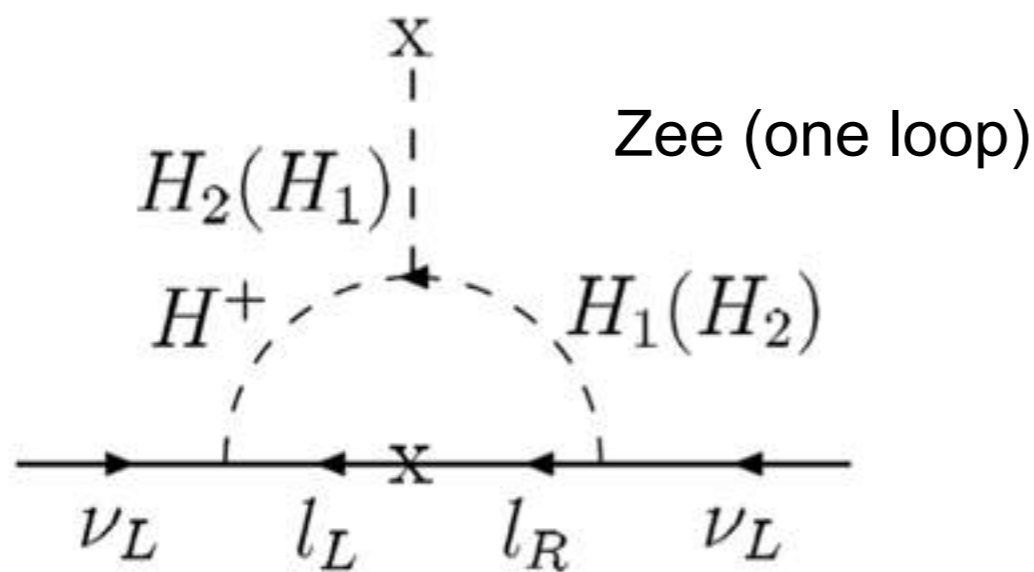
Malinsky, Romao, Valle

$$\begin{pmatrix} 0 & M_D & M_L \\ M_D^T & 0 & M \\ M_L^T & M^T & 0 \end{pmatrix}$$

$$M_\nu = M_D (M_L M^{-1})^T + (M_L M^{-1}) M_D^T$$

LFV predictions

Loop Models of Neutrino Mass



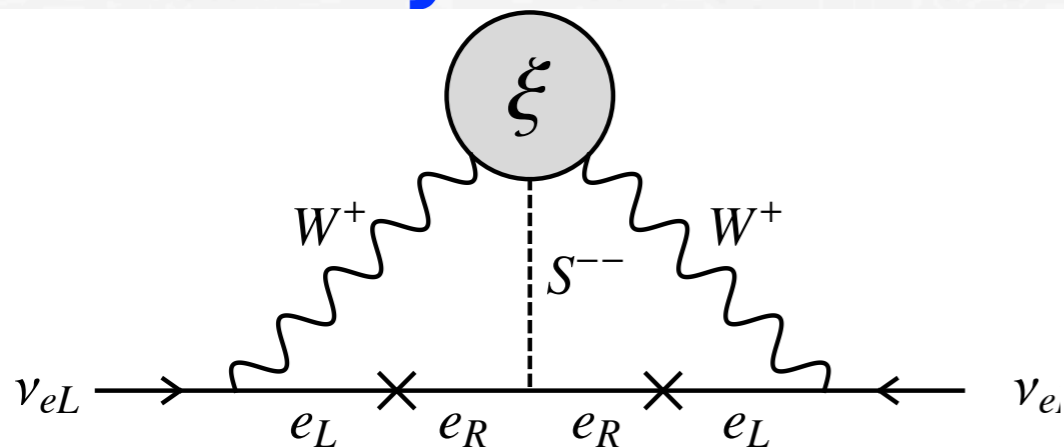
Scotogenic model

Cocktail model

Doubly charged singlet scalar

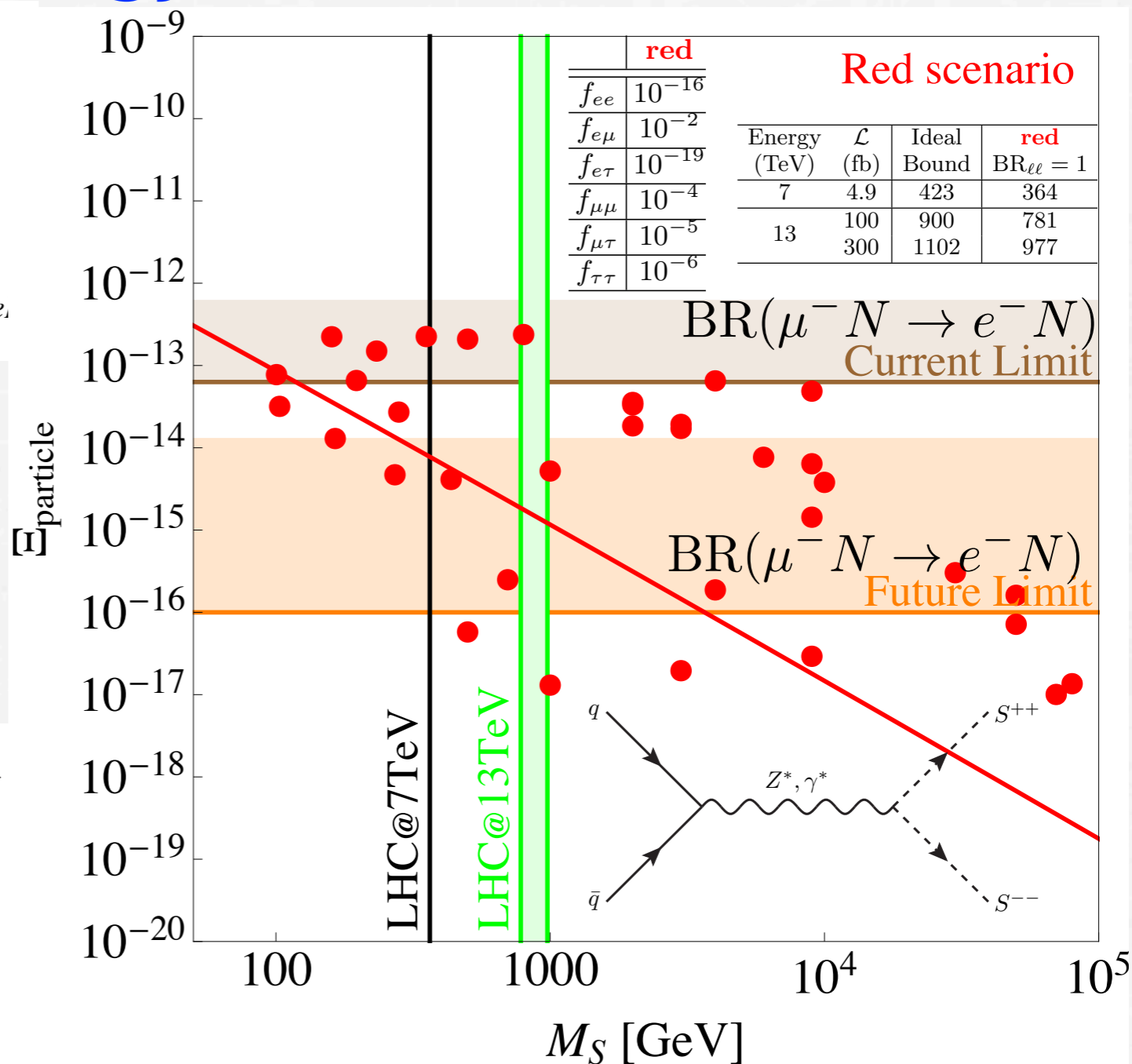
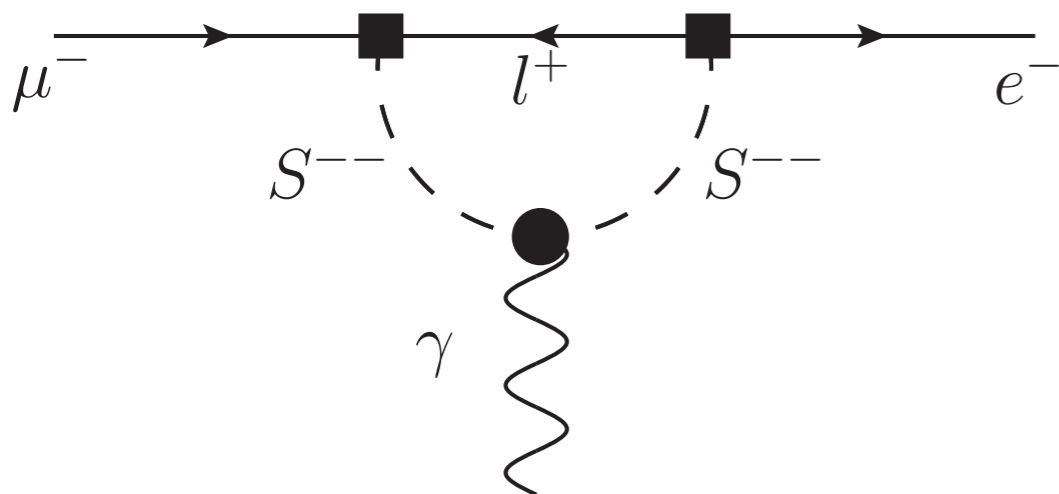
A Loop Model: Complementarity of the Intensity and the Energy Frontiers

Geib, SFK, Merle, No, Panizzi, 1512.04391



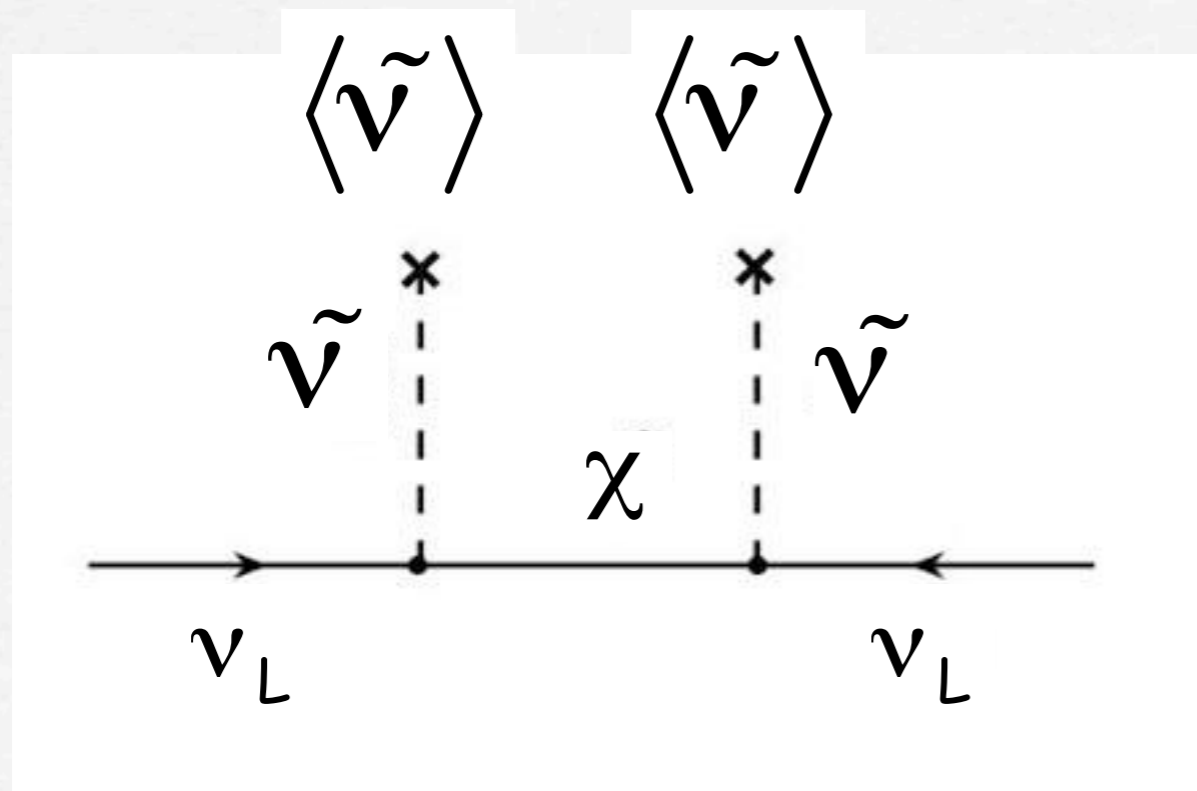
SFK, Merle, Panizzi, 1406.4137

$$(m_{ab}^\nu) \sim \begin{pmatrix} m_e^2 f_{ee} & m_e m_\mu f_{e\mu} & m_e m_\tau f_{e\tau} \\ m_e m_\mu f_{e\mu} & m_\mu^2 f_{\mu\mu} & m_\mu m_\tau f_{\mu\tau} \\ m_e m_\tau f_{e\tau} & m_\mu m_\tau f_{\mu\tau} & m_\tau^2 f_{\tau\tau} \end{pmatrix}$$



R-Parity Violating SUSY

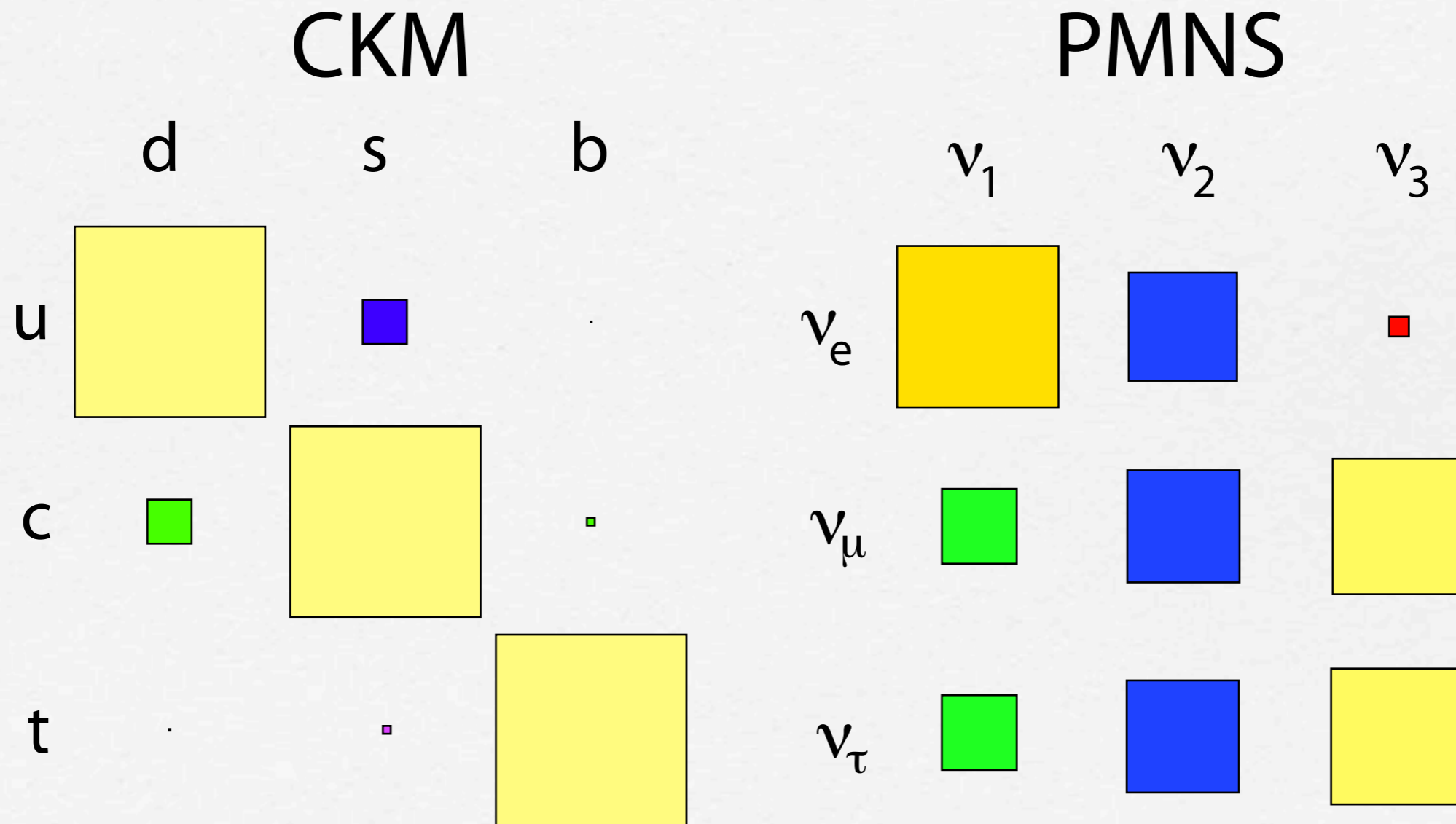
- Majorana masses can be generated via RPV SUSY
- Scalar partners of lepton doublets (slepton doublets) have same quantum numbers as Higgs doublets
- If R-parity is violated then sneutrinos may get (small) VEVs inducing a mixing between neutrinos and neutralinos χ



Hirsch, Diaz, Porod, Romao, Valle
[hep-ph/0004115]...

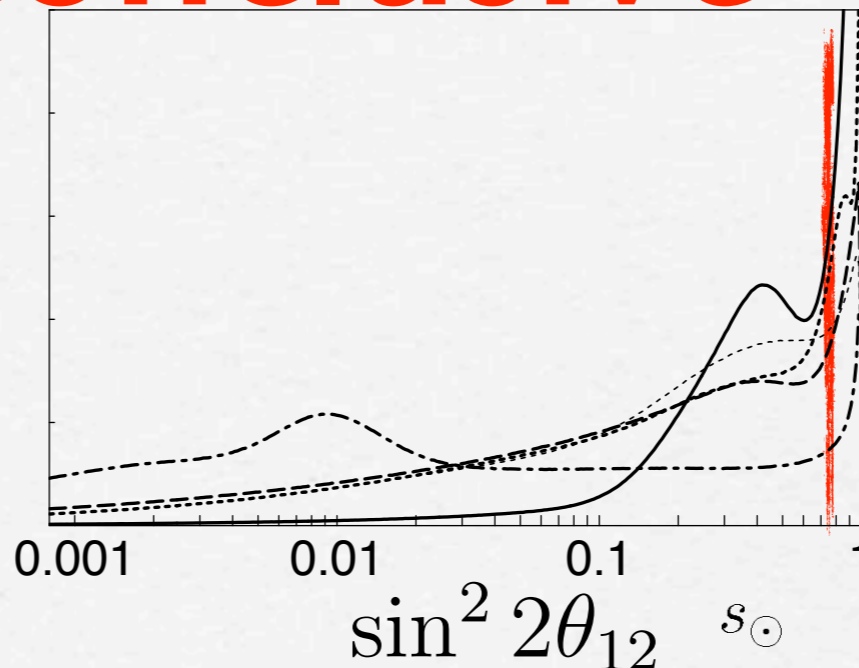
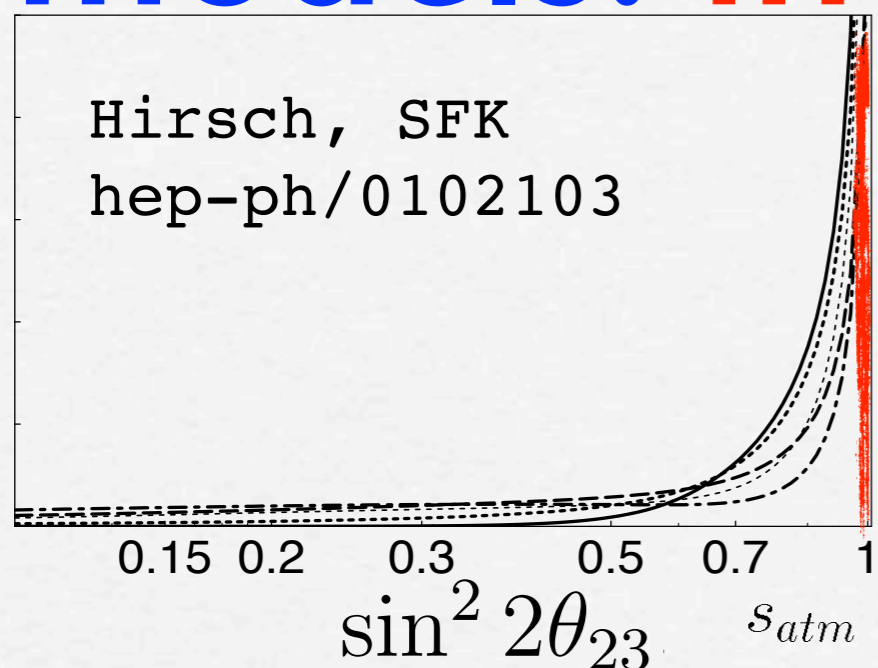
$$m_{LL}^{\nu} \approx \frac{\langle \tilde{\nu} \rangle^2}{M_{\chi}} \approx \frac{\text{MeV}^2}{\text{TeV}} \approx eV$$

What is the origin of Large Lepton Mixing?



#1: Anarchy vs U(1) models: inconclusive

Hall, Murayama, Weiner 9911341;
Hirsch, SFK 0102103;
Altarelli, Feruglio, Masina
0210342; ...

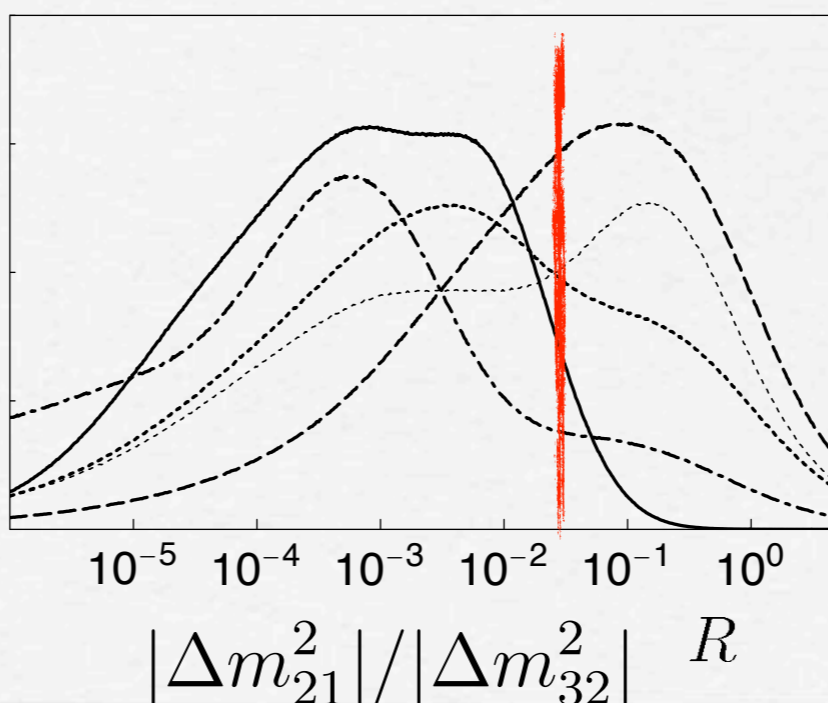
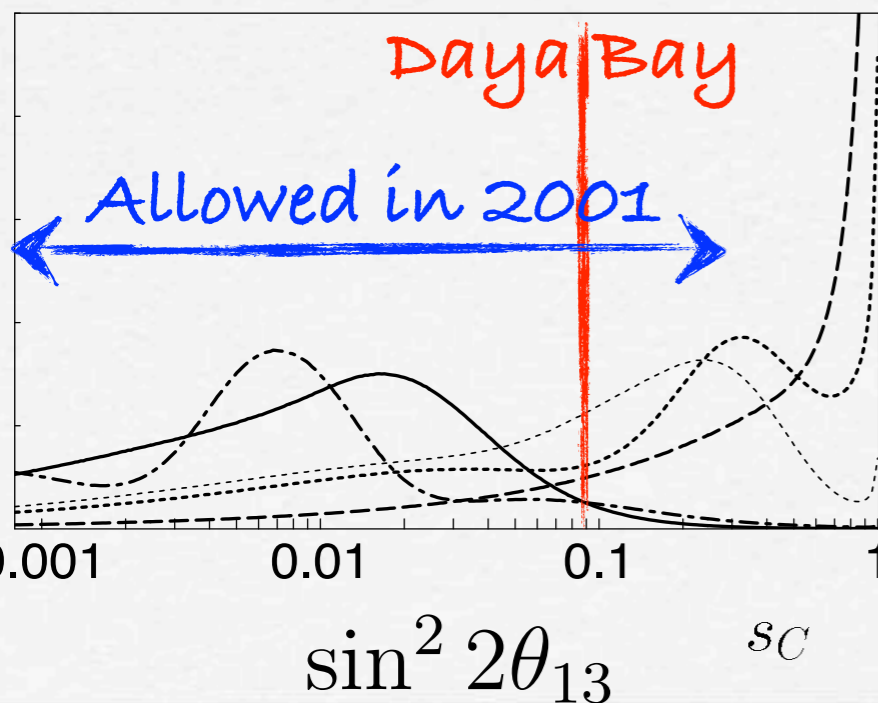


$$m_{LL}^{anarchy} \sim \begin{pmatrix} - & - & - \\ - & - & - \\ - & - & - \end{pmatrix}$$

dashes

$$m_{LL}^{FC1} \sim \begin{pmatrix} \lambda^4 & \lambda^2 & \lambda^2 \\ \lambda^2 & 1 & 1 \\ \lambda^2 & 1 & 1 \end{pmatrix}$$

full



$$m_{LL}^{FC2} \sim \begin{pmatrix} \lambda^7 & \lambda^5 & \lambda^5 \\ \lambda^5 & \lambda^3 & \lambda^3 \\ \lambda^5 & \lambda^3 & \lambda^3 \end{pmatrix}$$

dot-dashes

$$m_{LL}^{FC3} \sim \begin{pmatrix} \lambda & 1 & 1 \\ 1 & \lambda^{-1} & \lambda^{-1} \\ 1 & \lambda^{-1} & \lambda^{-1} \end{pmatrix}$$

thick dots

#2: Flavour Symmetry (non-Abelian)

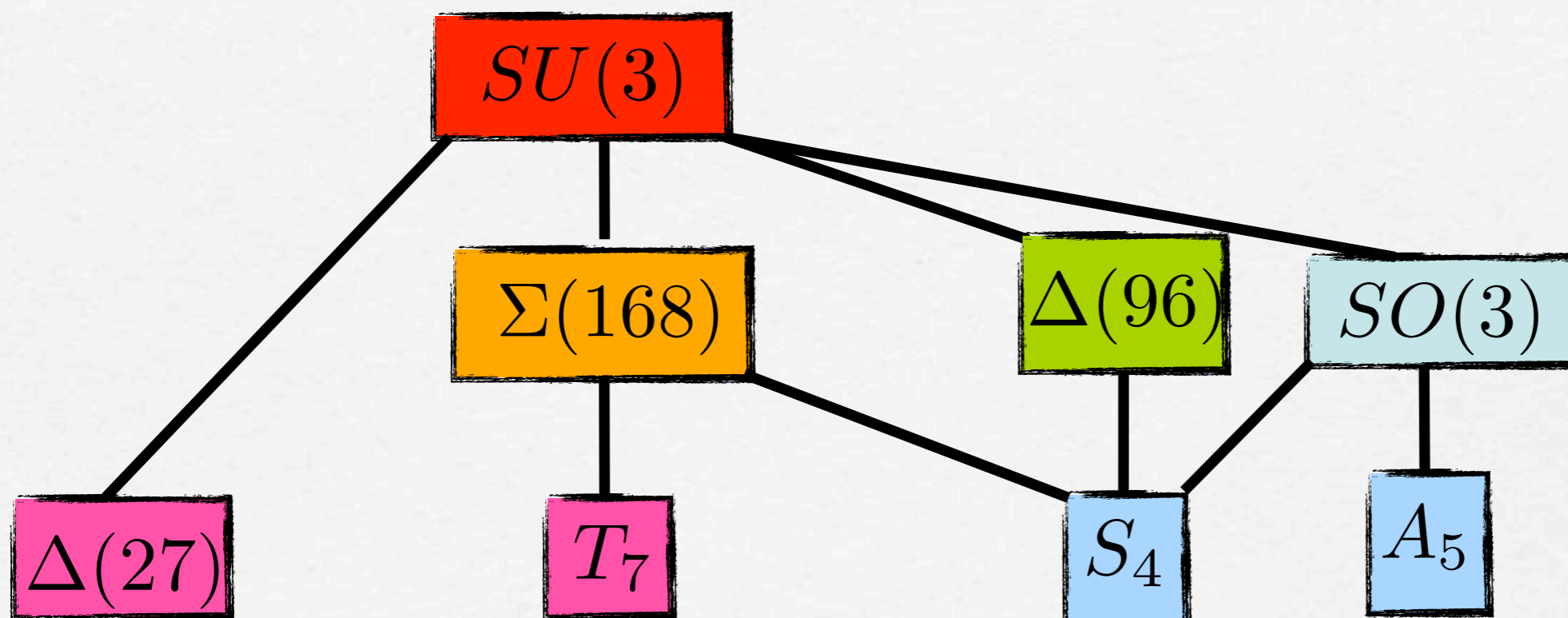
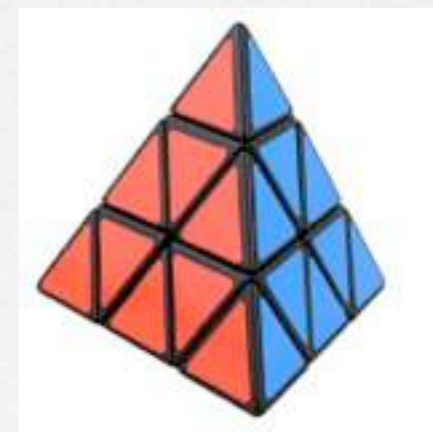
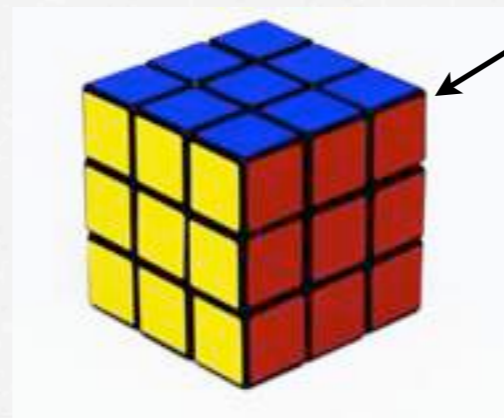
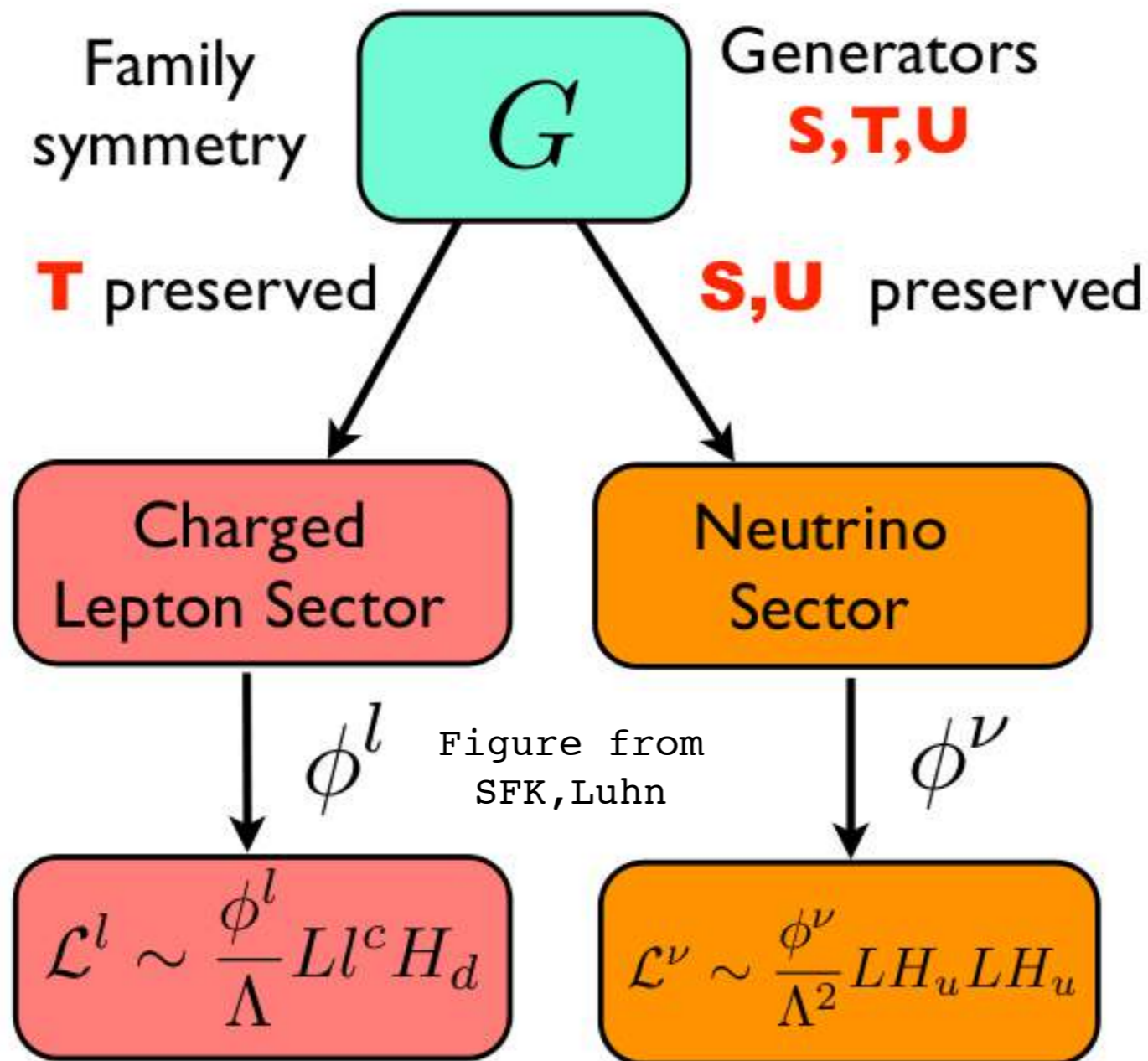


Figure from
SFK 1510.02091



Direct Models



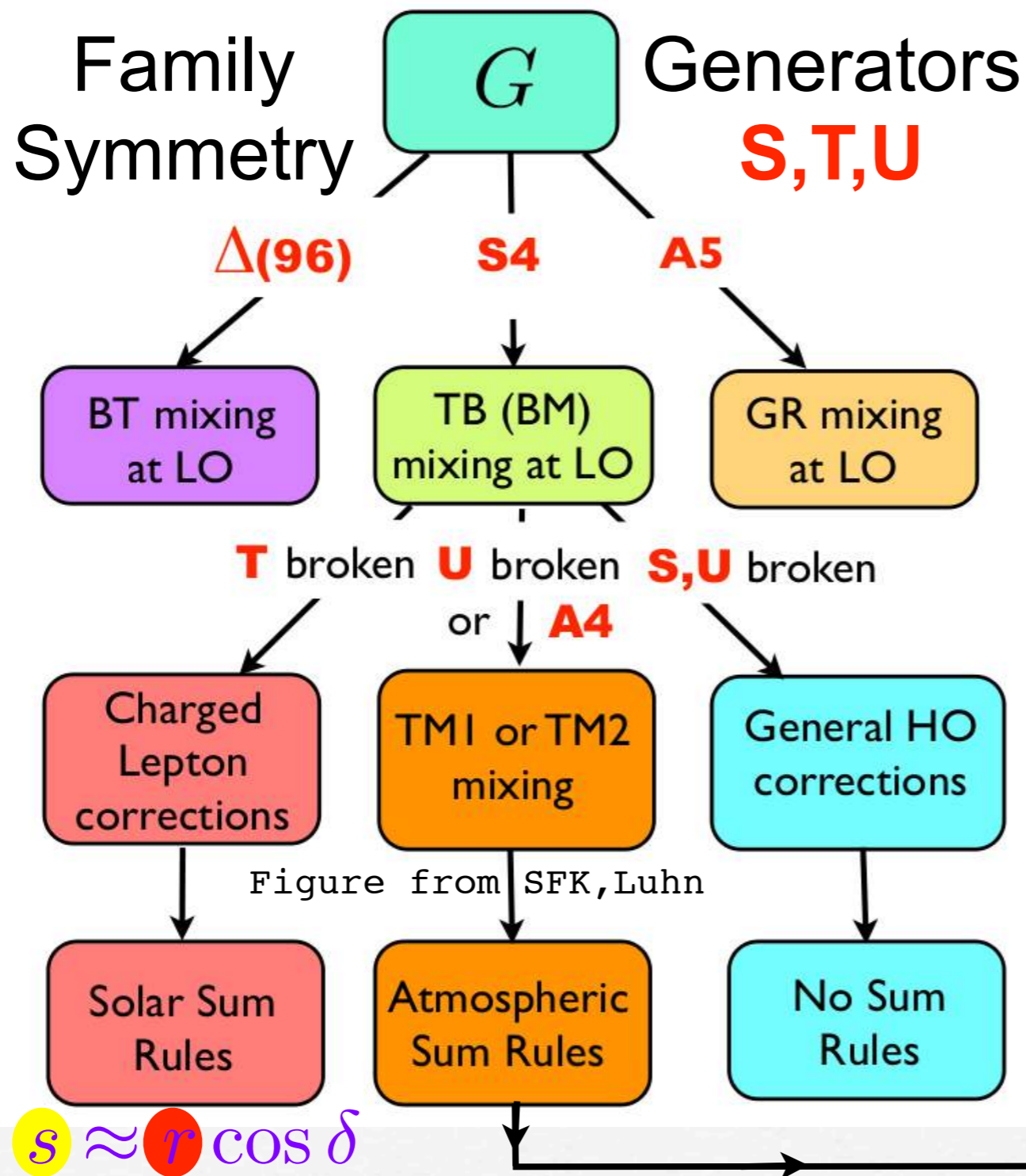
$$\Delta(6n^2)$$

Luhn et al

is the only viable symmetry class - predicts zero Dirac CPV

Holthausen, Lim, Lindner;
SFK, Neder, Stuart;
Lavoura, Ludl;
Fonseca, Grimus

Semi-Direct Models



Klein symmetry and T are **partly** preserved as subgroups of some family symmetry

$$U_{BM} = \begin{pmatrix} \frac{1}{\sqrt{2}} & \frac{1}{\sqrt{2}} & 0 \\ -\frac{1}{2} & \frac{1}{2} & \frac{1}{\sqrt{2}} \\ \frac{1}{2} & -\frac{1}{2} & \frac{1}{\sqrt{2}} \end{pmatrix} P$$

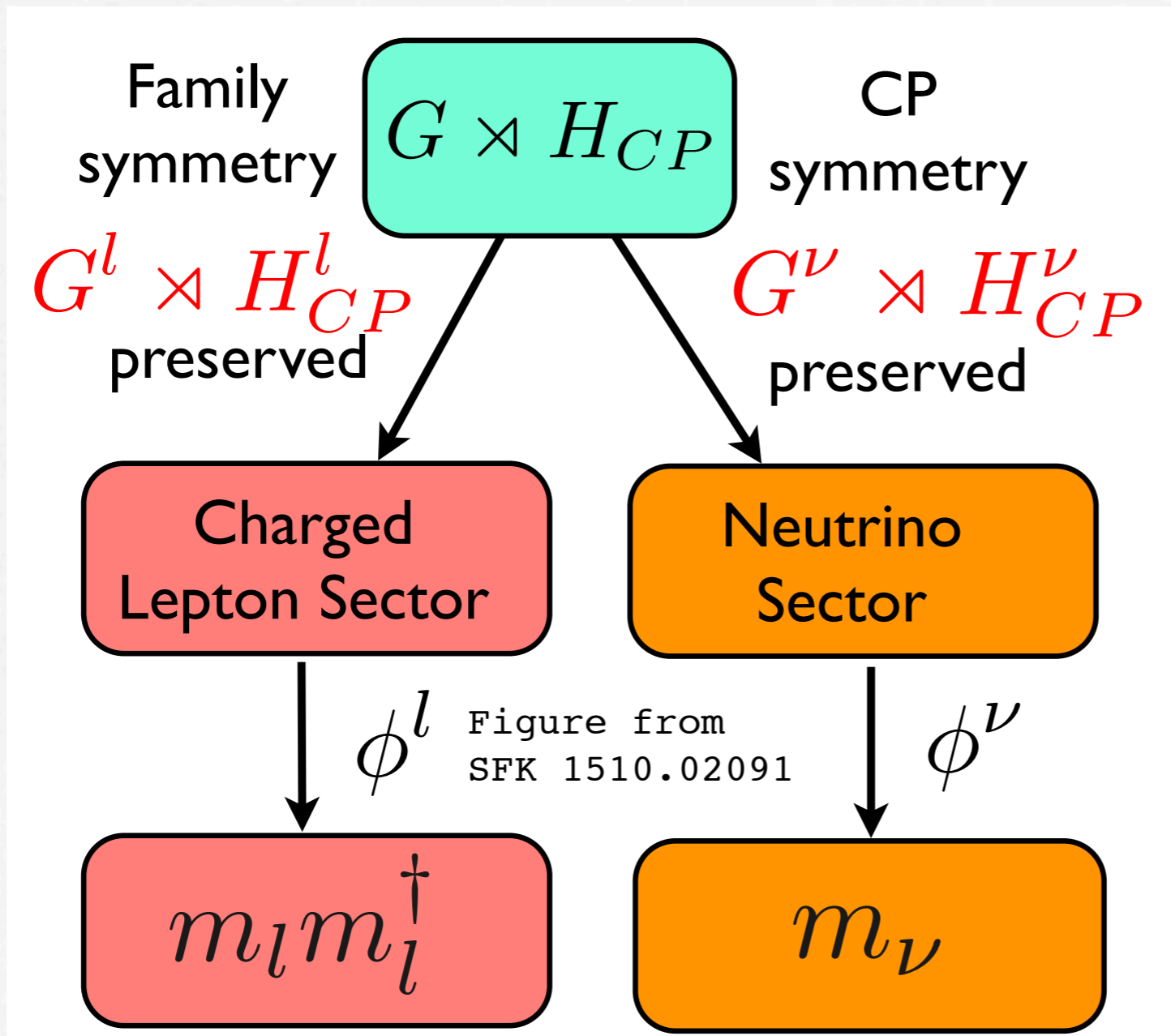
$$U_{TB} = \begin{pmatrix} \sqrt{\frac{2}{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix} P$$

TM1 TM2

$a \approx r \cos \delta$ $a \approx -\frac{r}{2} \cos \delta$

SFK, Antusch, ...

CP violation



Feruglio, Hagedorn;
 Holthausen, Lindner Schmidt;
 Ding, SFK, Luhn, Stuart;
 Nishi, Xing; Hagedorn, Meroni,
 Molinaro; Ding, SFK, Neder;
 Branco, SFK, Varzielas,
 Chen, ...

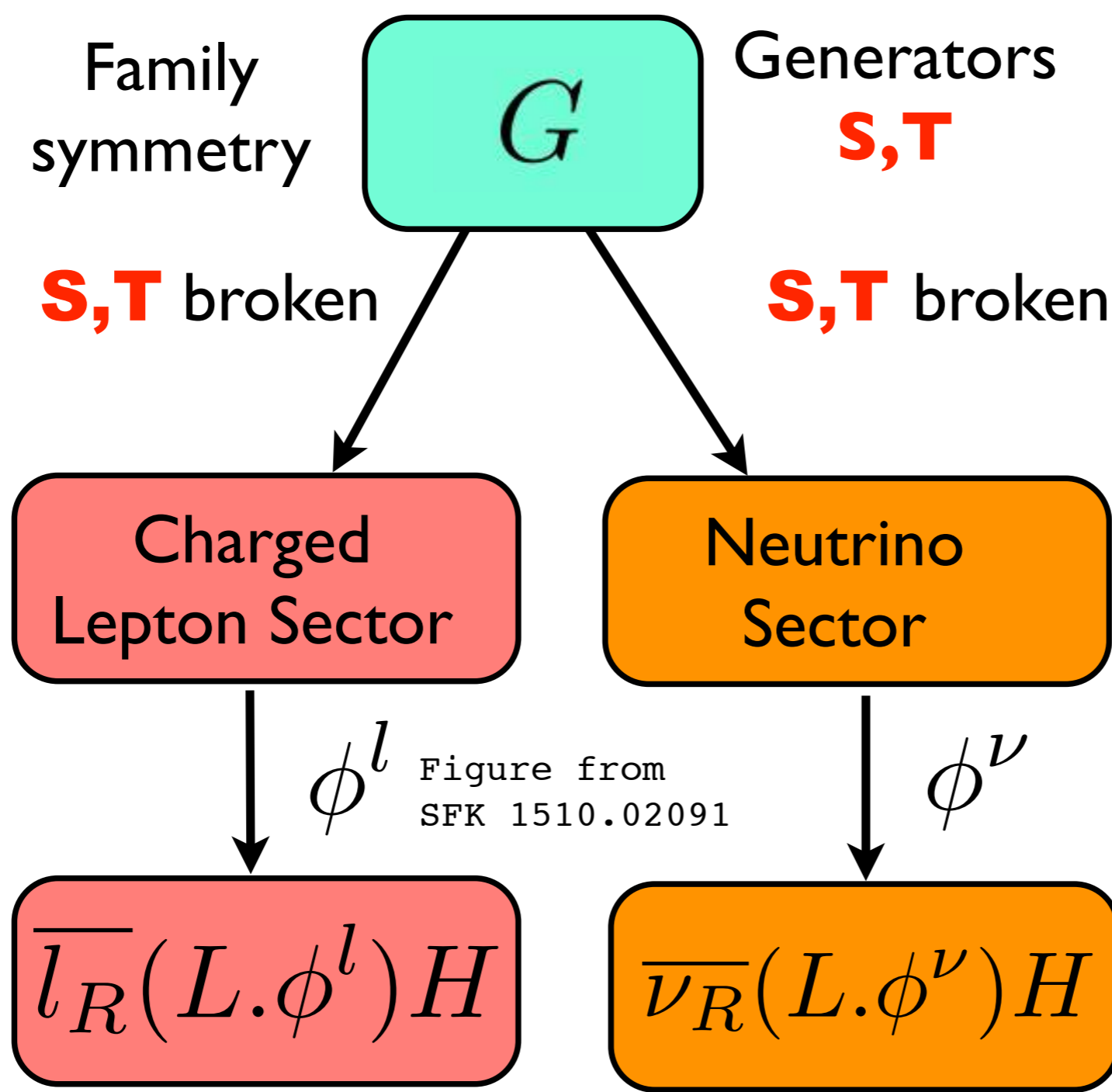
E.g. in semi-direct
 models where

$$G^\nu \sim Z_2^S$$

predicts maximal
 Dirac Phase

$$\delta = \pm \pi/2$$

Indirect Models



SFK, G. Ross, de Medeiros
Varzielas, Antusch, Luhn,
Bjorkeroth, ...

Predictions follow
from flavon vacuum
alignment

What is the minimal natural model of neutrino mass and mixing?

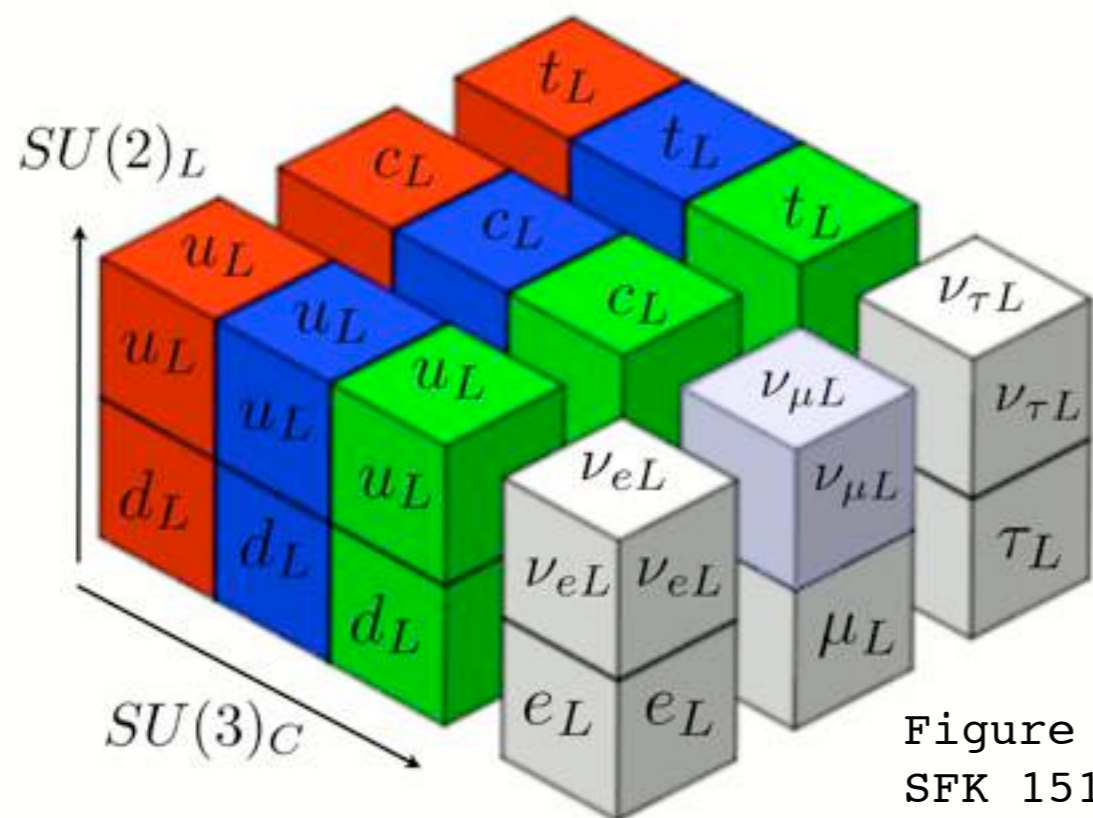
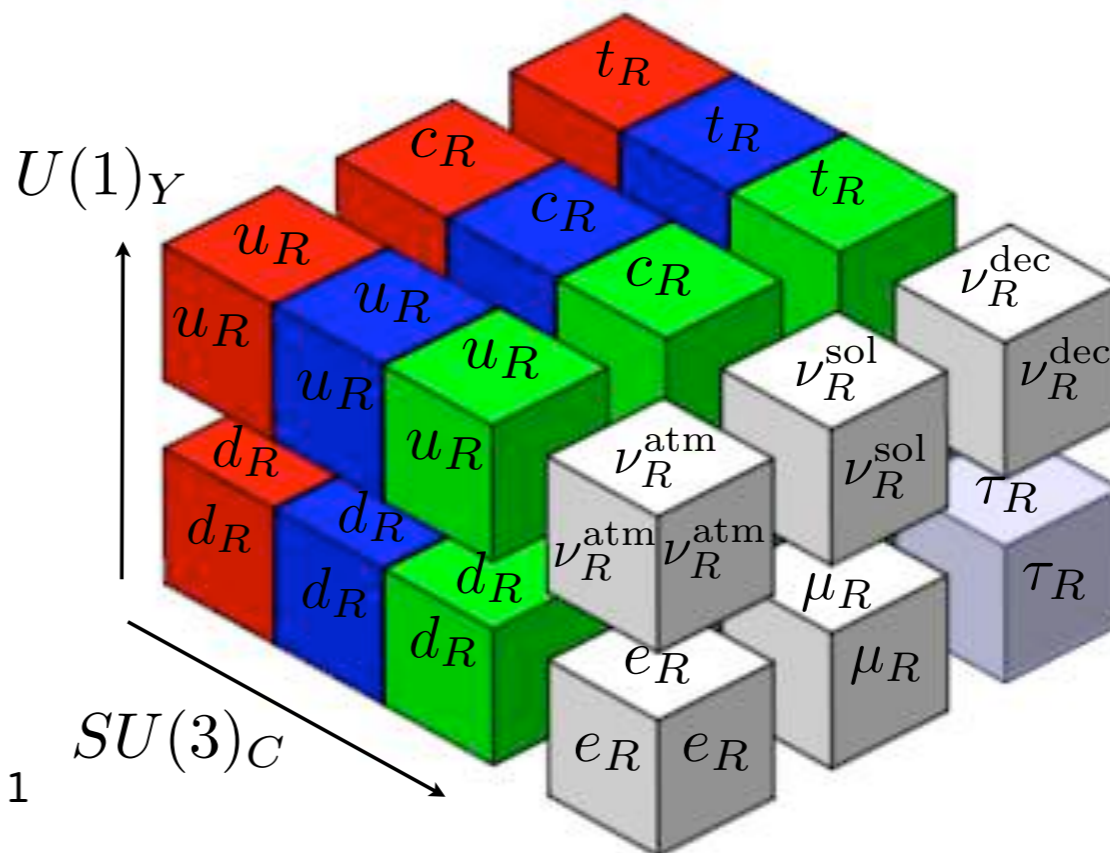


Figure from
SFK 1510.02091



Sequential Dominance

Heavy Majorana

Dirac

$$M_R = \begin{pmatrix} M_{\text{atm}} & 0 & 0 \\ 0 & M_{\text{sol}} & 0 \\ 0 & 0 & [M_{\text{dec}}] \end{pmatrix}$$

$$m^D = \begin{pmatrix} m_{e,\text{atm}}^D & m_{e,\text{sol}}^D & [m_{e,\text{dec}}^D] \\ m_{\mu,\text{atm}}^D & m_{\mu,\text{sol}}^D & [m_{\mu,\text{dec}}^D] \\ m_{\tau,\text{atm}}^D & m_{\tau,\text{sol}}^D & [m_{\tau,\text{dec}}^D] \end{pmatrix} \equiv \left(m_{\text{atm}}^D \quad m_{\text{sol}}^D \quad [m_{\text{dec}}^D] \right)$$

Light Majorana $m^\nu = m^D M_R^{-1} m^{D T}$

has three contributions, one from each RH neutrino.

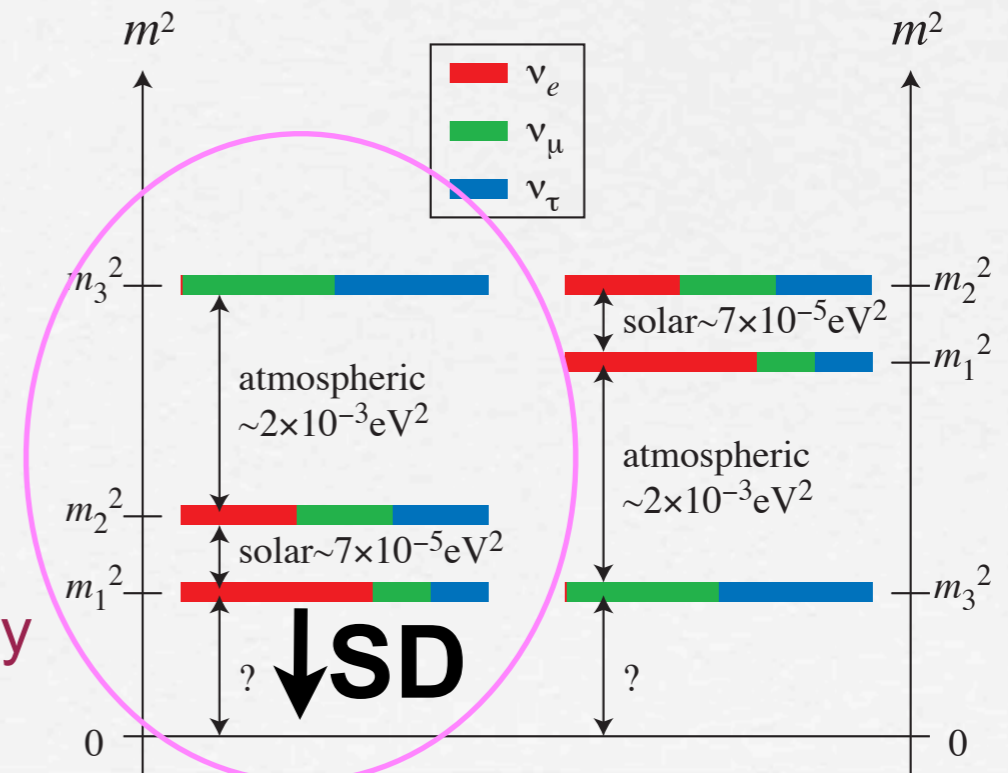
Sequential dominance assumption:

$$\frac{m_{\text{atm}}^D m_{\text{atm}}^{D T}}{M_{\text{atm}}} \gg \frac{m_{\text{sol}}^D m_{\text{sol}}^{D T}}{M_{\text{sol}}} \left[\gg \frac{m_{\text{dec}}^D m_{\text{dec}}^{D T}}{M_{\text{dec}}} \right]$$

Responsible for
atm nu mass
 $m_3 \sim 50 \text{ meV}$

Responsible for
sol nu mass
 $m_2 \sim 9 \text{ meV}$

Predicts strong
normal hierarchy
 $m_1 \ll m_2$



Littlest Seesaw Model

2RHN

$$\frac{\phi_{\text{atm}}}{\Lambda} LH N_{\text{atm}}^c + \frac{\phi_{\text{sol}}}{\Lambda} LH N_{\text{sol}}^c + M_{\text{atm}} N_{\text{atm}}^c N_{\text{atm}}^c + M_{\text{sol}} N_{\text{sol}}^c N_{\text{sol}}^c$$

Messenger scale

Flavon vevs

$$\langle \phi_{\text{atm}} \rangle = v_{\text{atm}} \begin{pmatrix} 0 \\ 1 \\ 1 \end{pmatrix} \quad \langle \phi_{\text{sol}} \rangle = v_{\text{sol}} \begin{pmatrix} 1 \\ 3 \\ 1 \end{pmatrix}$$

Constrained
Sequential
Dominance
"CSD3"

$$m^\nu = \frac{v^2}{\Lambda^2} \left(\frac{\langle \phi_{\text{atm}} \rangle \langle \phi_{\text{atm}} \rangle^T}{M_{\text{atm}}} + \frac{\langle \phi_{\text{sol}} \rangle \langle \phi_{\text{sol}} \rangle^T}{M_{\text{sol}}} \right)$$

After
seesaw

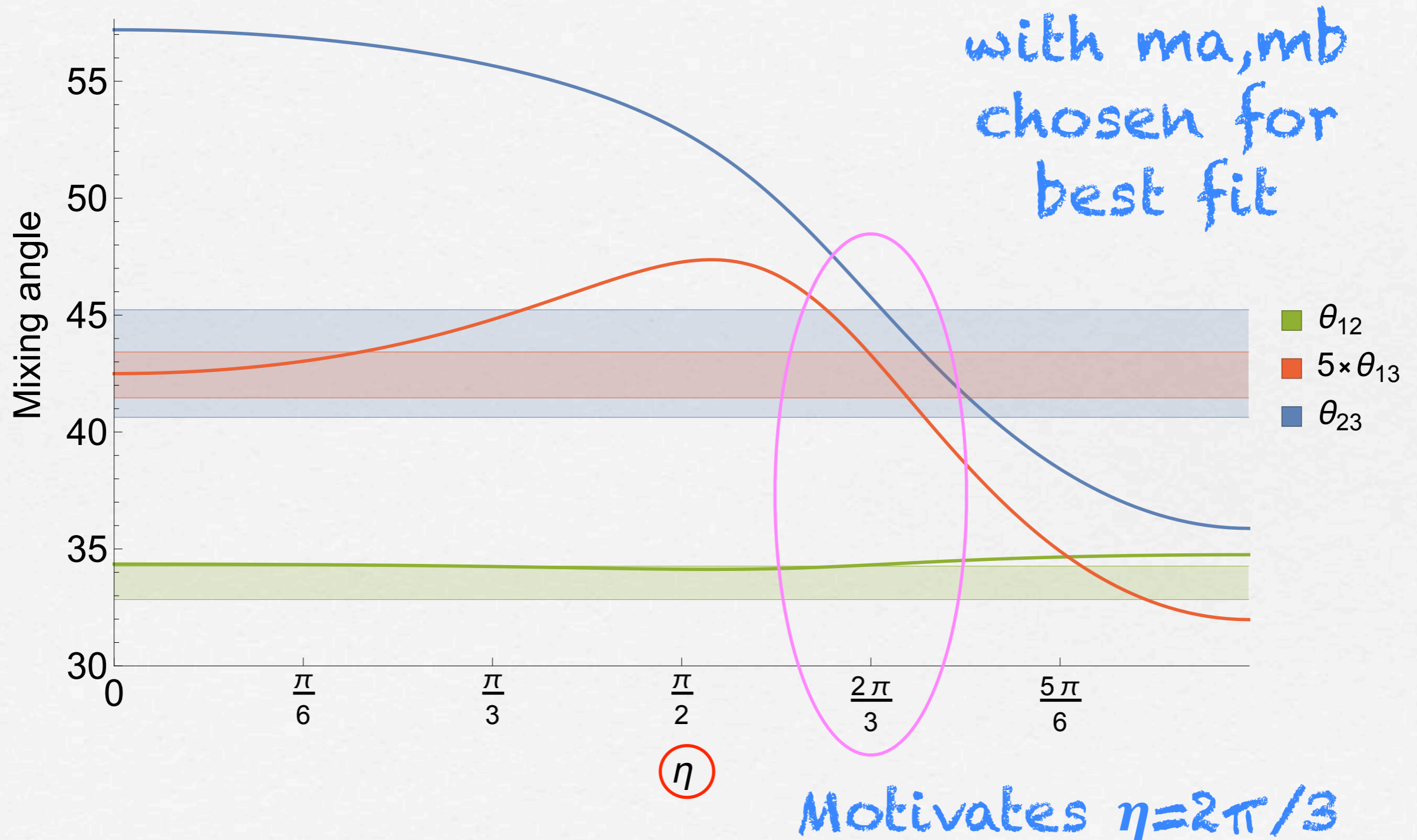
Only one phase

$$m^\nu = m_a \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 1 \\ 0 & 1 & 1 \end{pmatrix} + m_b e^{i\eta} \begin{pmatrix} 1 & 3 & 1 \\ 3 & 9 & 3 \\ 1 & 3 & 1 \end{pmatrix}$$

TM1 mixing

Littlest seesaw = smallest number of parameters: m_a, m_b, η

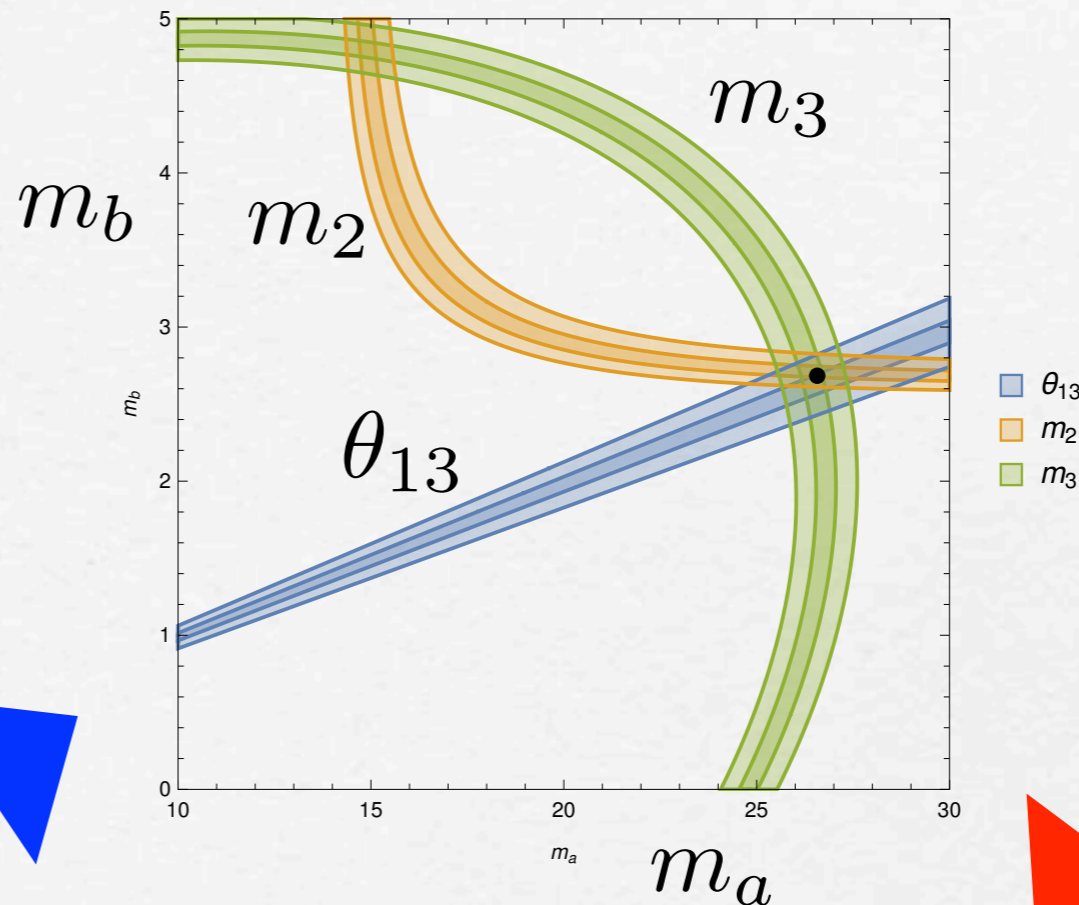
Littlest Seesaw Model



Littlest Seesaw Model

Prouse et al
(to appear)

$$m^\nu = m_a \begin{pmatrix} 0 & 0 & 0 \\ 0 & 1 & 1 \\ 0 & 1 & 1 \end{pmatrix} + m_b e^{i2\pi/3} \begin{pmatrix} 1 & 3 & 1 \\ 3 & 9 & 3 \\ 1 & 3 & 1 \end{pmatrix}$$



Only two parameters now: m_a, m_b

Very predictive:
7 outputs!!

m_a (meV)	m_b (meV)	η (rad)	θ_{12} ($^\circ$)	θ_{13} ($^\circ$)	θ_{23} ($^\circ$)	δ_{CP} ($^\circ$)	m_1 (meV)	m_2 (meV)	m_3 (meV)
26.57	2.684	$\frac{2\pi}{3}$	34.3	8.67	45.8	-86.7	0	8.59	49.8
Value	from	[25]	$33.48^{+0.78}_{-0.75}$	$8.50^{+0.20}_{-0.21}$	$42.3^{+3.0}_{-1.6}$	-54^{+39}_{-70}	0	8.66 ± 0.10	49.57 ± 0.47

Conclusions

- *What is the Origin of Small Neutrino Masses?*

Roadmap of possibilities e.g. extra dims, loops, seesaw models,...

- *What is the origin of Large Lepton Mixing?*

Flavour symmetry is really the only way to make progress, but how to implement it? Direct, semi-direct, indirect, CP violation?

- *What is the minimal natural model of neutrino mass and mixing?*

Littlest Seesaw: 2RHN, CSD3, 2 input parameters, 7 predictions!

Thank YOU