## Rare decays at LHCb

### Eluned Smith<sup>1</sup> on behalf on the LHCb collaboration

<sup>1</sup>Imperial College London

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## The LHCb detector

The LHCb detector is a single arm spectrometer which covers the forward region at LHC.



 $\Delta p/p \sim 0.4\%$  at 5 GeV,  $\sigma_{IP} = 20 \ \mu m$  for high  $p_T$  tracks.  $\pi/K$  separation:  $\epsilon_K \sim 90\%$ , 5%  $\pi \rightarrow K$  mis-id.  $\pi/\mu$  separation:  $\epsilon_\mu \sim 97\%$ , 1-3%  $\pi \rightarrow K$  mis-id.



## Rare decays: FCNC processes

Decays mediated via Flavour Changing Neutral Currents (FCNC) occur at loop order and are suppressed in the SM.



New particles can affect the decay rates or the angular distributions of the final state particles.





## Rare decays: outline

#### Test of the Minimal Flavour Violation hypothesis

1) Measuring  $V_{td}/V_{ts}$  using the  $b o d(s)\ell\ell$  transitions  $B^+ o \pi^+({\cal K}^+)\mu^+\mu^-$  [LHCb JHEP 10 (2015) 034]

#### Sensitivity to Wilson coefficients

**2)** Very rare decays  $B_s^0 \rightarrow \mu^+ \mu^-$  and  $B^0 \rightarrow \mu^+ \mu^-$  [ $C_{10}, C_S, C_P$ ] [Nature 522 (2015) pp. 68-72]

**3)** Angular analyses using  $b \to s\ell\ell$  transitions [ $C_7, C_9, C_{10}$ ] :

• 
$$B^0 
ightarrow K^{*0} \ (
ightarrow K^+ \pi^-) \mu^+ \mu^-$$
 decays [LHCb, JHEP 02 (2016) 104]

• 
$$B^0_s o \Phi(K^+K^-) \mu^+\mu^-$$
 decays [LHCb, JHEP 09 (2015) 179]

#### Lepton universality

4) Branching fraction measurements of  $b \to s\ell\ell$  transitions  $B^+ \to K^+ \mu^+ \mu^-$  and  $B^+ \to K^+ e^+ e^-$  [LHCb, PRL113 (2014) 151601]



# Minimal flavour violation

Comparing the CKM elements obtained via loop and tree level processes tests the MFV hypothesis that NP flavour structure = SM flavour structure.

$$\underbrace{\mathbb{A}}_{l} > \frac{4.4 \text{ TeV}}{|V_{ti}^* V_{tj}| / |C_{ij}|^{1/2}} \sim \begin{cases} 1.3 \times 10^4 \text{ TeV} \times |c_{sd}|^{1/2} \\ 5.1 \times 10^2 \text{ TeV} \times |c_{bd}|^{1/2} \\ 1.1 \times 10^2 \text{ TeV} \times |c_{bs}|^{1/2} \\ 1.1 \times 10^2 \text{ TeV} \times |c_{bs}|^{1/2} \\ \text{Lower mass limit on NP if a generic flavour structure is assumed for NF (i.e. NOT the MFV hypothesis)} \end{cases}$$







As top quarks don't hadronise, Vtd and Vts cannot be measured directly with tree diagrams, -> use unitarity constraints from CKM matrix

A. Bazavov et. al. [arXiv:1602.03560]





# Calculating $V_{td}$ ( $V_{ts}$ ) with $B^+ \rightarrow \pi^+$ ( $K^+$ ) $\mu^+\mu^-$ decays





## Wilson coefficients



$$B^0_s o \mu^+ \mu^-$$
 and  $B^0 o \mu^+ \mu^-$  [ ${\cal C}_{10}, {\cal C}_s, {\cal C}_p$ ]

LHCb and CMS data combined:











 $B^0 \to K^{*0} [\to K^+ \pi^-] \mu^+ \mu^-$  angular analysis  $[\mathcal{C}_7, \mathcal{C}_9, \mathcal{C}_{10}]$ Angular decay fully described by the dilepton mass  $(q^2)$  and the angles  $\cos(\theta)_{\parallel} \cos(\theta)_k$  and  $\phi$ :





 $B^0 \rightarrow K^{*0} [\rightarrow K^+ \pi^-] \mu^+ \mu^-$  angular analysis: Results



Use  $B^0 \to K^{*0} J/\psi$  as control channel.



 $B^0 
ightarrow K^{*0} [
ightarrow K^+ \pi^- ] \mu^+ \mu^-$  angular analysis: Results









#### Form factor free observables

Can construct ratios of angular observables where form-factors cancel, giving clean theoretical predictions:

$$P_5' = S_5/\sqrt{F_L(1-F_L)}$$

 $P'_5$  plot: Bins 4/5 = local SM tension of 2.8 and 3.0 $\sigma$ . Global tension= 3.4 $\sigma$ , assuming tension due to shift in Wilson coeff.  $\mathcal{R}e(C_9)$ 



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$$B^0_s o \phi[ o K^+K^-]\mu^+\mu^-$$
 [ $\mathcal{C}_7, \mathcal{C}_9, \mathcal{C}_{10}$ ]

Equivalent process of  $B^0 \to K^{*0} \mu^+ \mu^-$  for  $B^0_s$  mesons.

Angular variables consistent with the SM.  $P_5'$  cannot be measured as  $B_s^0 \rightarrow \phi \mu^+ \mu^-$  not self-tagging.

In bin 1<  $q^2$  <6 GeV/ $c^2$  the data is 3.3 $\sigma$  from the SM prediction.





## Global fits



An example of a fit to many results from  $b \rightarrow s\ell\ell$  transitions. 3-4 $\sigma$  tension with SM observed in the  $C_9$  Wilson coefficient.



# $B^+ ightarrow K^+ \mu^+ \mu^-$ and $B^+ ightarrow K^+ \; e^+ e^-$ decays

## Lepton universality

The quantity:  $R_{k} = \frac{\int \Gamma(B^{+} \rightarrow K^{+} \mu^{+} \mu^{-})/dq^{2}.dq^{2}}{\int \Gamma(B^{+} \rightarrow K^{+} e^{+} e^{-})/dq^{2}.dq^{2}}$ differs from unity only due to phase space.

Theoretically clean as matrix elements cancel.



Experimentally challenging due to electrons.

Deficit of  $b \rightarrow s \mu^+ \mu^-$  transitions, tension in same direction as measurements previously discussed



# Conclusions

- Flavour observables in rare decays allow for NP searches and can place many strong constraints on NP models.
- Some tensions with the SM observed, particular within  $b \rightarrow s\ell\ell$  transitions.
- Many rare decay analyses performed with LHC Run 1 data and many more results to come using the Run 1 and first Run 2 datasets.



# Back-up slides





# Global fits $b \rightarrow s \ell \ell$ decays

# Decays included in fit in *Altmannshofer et al.* with tension $> 1.9 \sigma$ in a Wilson Coefficient.

Decay	obs.	$q^2$ bin	SM pred.	measurement		pull
$\bar{B}^0\to \bar{K}^{*0}\mu^+\mu^-$	$F_L$	[2, 4.3]	$0.81\pm0.02$	$0.26\pm0.19$	ATLAS	+2.9
$\bar{B}^0\to \bar{K}^{*0}\mu^+\mu^-$	$F_L$	[4, 6]	$0.74\pm0.04$	$0.61\pm0.06$	LHCb	+1.9
$\bar{B}^0\to \bar{K}^{*0}\mu^+\mu^-$	$S_5$	[4, 6]	$-0.33\pm0.03$	$-0.15\pm0.08$	LHCb	-2.2
$\bar{B}^0\to \bar{K}^{*0}\mu^+\mu^-$	$P_5'$	[1.1, 6]	$-0.44\pm0.08$	$-0.05\pm0.11$	LHCb	-2. <mark>9</mark>
$\bar{B}^0\to \bar{K}^{*0}\mu^+\mu^-$	$P_5'$	[4, 6]	$-0.77\pm0.06$	$-0.30\pm0.16$	LHCb	-2.8
$B^- \to K^{*-} \mu^+ \mu^-$	$10^7 \frac{d \mathrm{BR}}{dq^2}$	[4, 6]	$0.54\pm0.08$	$0.26\pm0.10$	LHCb	+2.1
$\bar{B}^0\to \bar{K}^0\mu^+\mu^-$	$10^8 \frac{d \mathrm{BR}}{dq^2}$	[0.1,2]	$2.71\pm0.50$	$1.26\pm0.56$	LHCb	+1.9
$\bar{B}^0\to \bar{K}^0\mu^+\mu^-$	$10^8 \frac{d\mathrm{BR}}{dq^2}$	[16, 23]	$0.93 \pm 0.12$	$0.37\pm0.22$	$\operatorname{CDF}$	+2.2
$B_s \to \phi \mu^+ \mu^-$	$10^7 \frac{dBR}{dq^2}$	[1, 6]	$0.48\pm0.06$	$0.23\pm0.05$	LHCb	+3.1



# $B_s^0 \rightarrow \mu^+ \mu^-$ results from Atlas





# NP models: contribution to $C_9$ arXiv:1308.1501

# Z' models

- Z' models involve tree-level exchange of a heavy neutral boson (Z') with a flavour-changing  $b \rightarrow s$  coupling.
- Z' models allow for significant contributions to the C<sub>9</sub> vector coefficient which could accommodate tensions in b → sℓℓ measurements.
- Best fit values for  $C_9$  global fits and constraints from  $B_s$  meson mixing favour a light Z' with mass of order 1 TeV. To avoid constraints from di-lepton searches, the Z' coupling to SM fermions must be  $\sim$  an order of magnitude less than those of the SM Z.



# NP models: contribution to $C_9$ arXiv:1308.1501

### Models with partial composites

- Most well-motivated to contribute to C<sub>7</sub> coefficients, contributions could be large enough to significantly reduce tension in C<sub>9</sub> coefficient.
- To contribute *C*<sub>9</sub> would require several cancellations in the V-A coefficient *C*<sub>10</sub> and properties of muons would be strongly dependent on chirality.

### MSSM

- Contributions to *C*<sub>9</sub> from *Z* penguins, charged Higgs (box and loops), Higgsinos (box and loops) all negligible.
- MSSM can affect C<sub>7</sub> which could soften tensions.

