

Beyond standard neutrinos

Pedro A N Machado



physics with
Beyond standard neutrinos

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Beyond standard neutrinos

uses of

Pedro A N Machado



Importance of SM for BSM

Mixing parameters

nu-fit.org

$\sin^2\theta_{12}$ @ 4%

$\sin^2\theta_{13}$ @ 6%

$\sin^2\theta_{23}$ @ 10%-ish

Mass squared differences

Δm^2_{21} @ 3%

$|\Delta m^2_{atm}|$ @ 2%

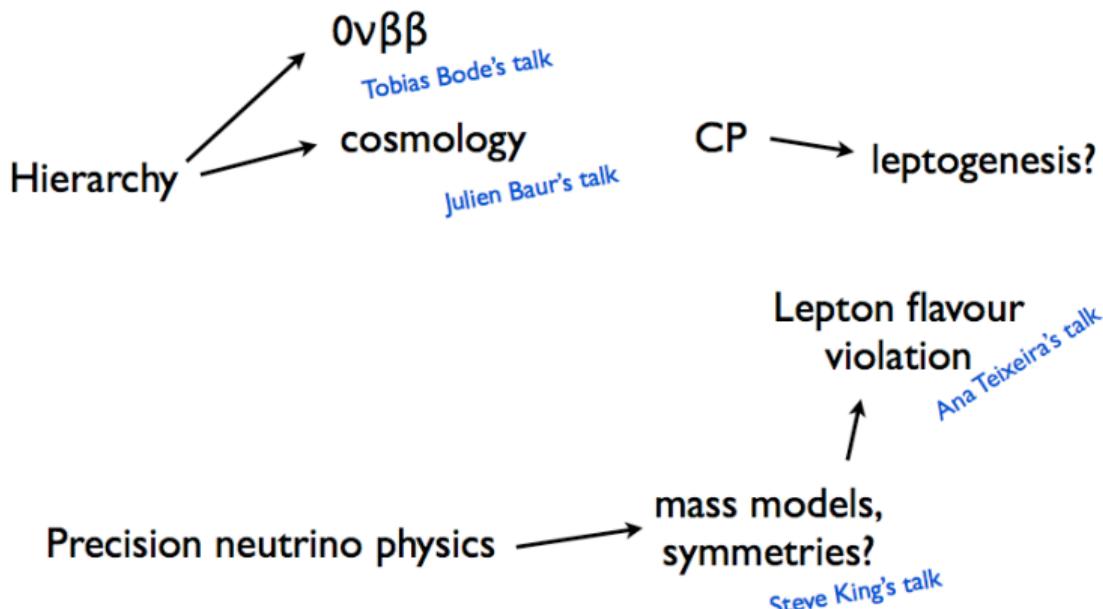
Kate Scholberg's talk

Hierarchy unknown

CP phase unknown

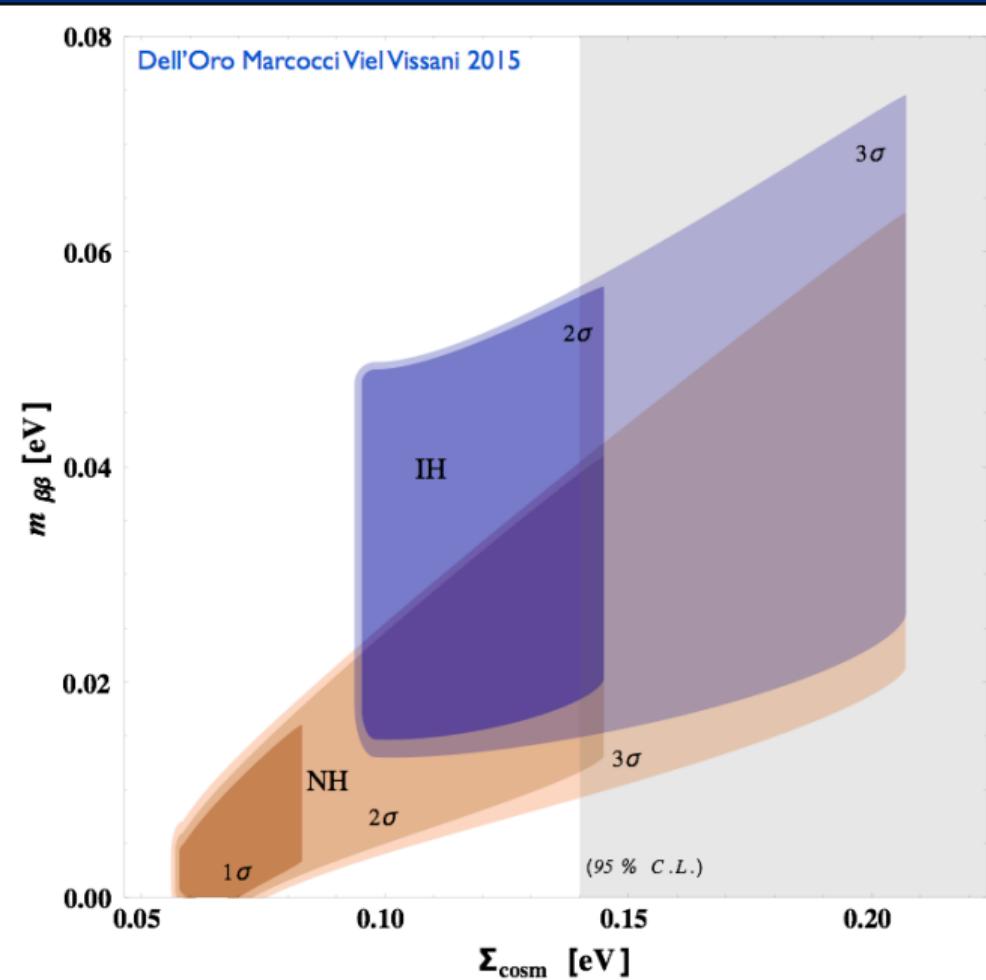
Mass scale unknown

Importance of SM for BSM

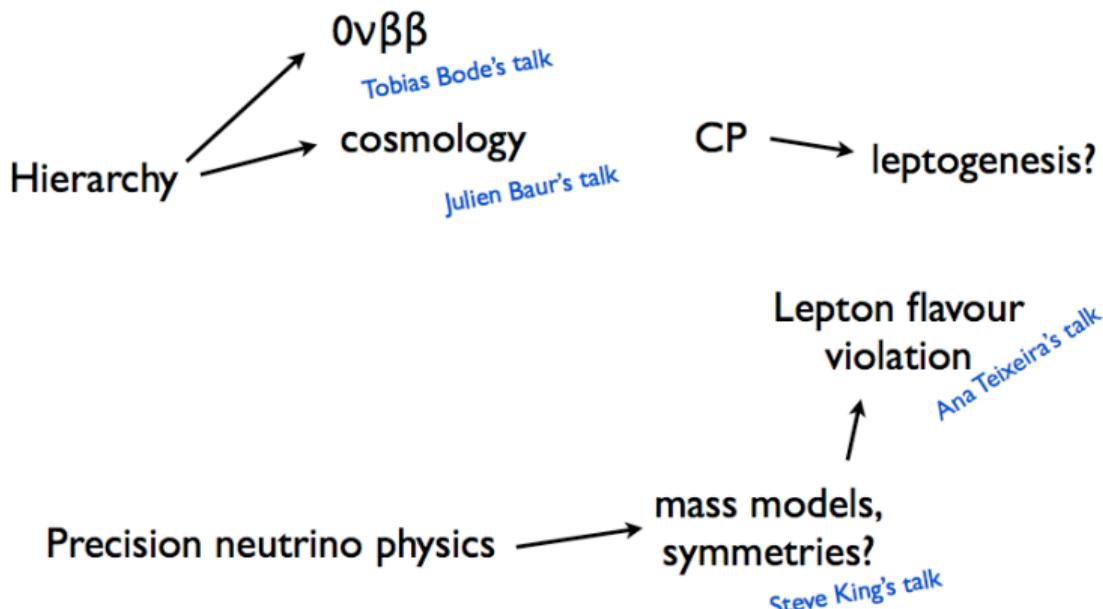


Hierarch

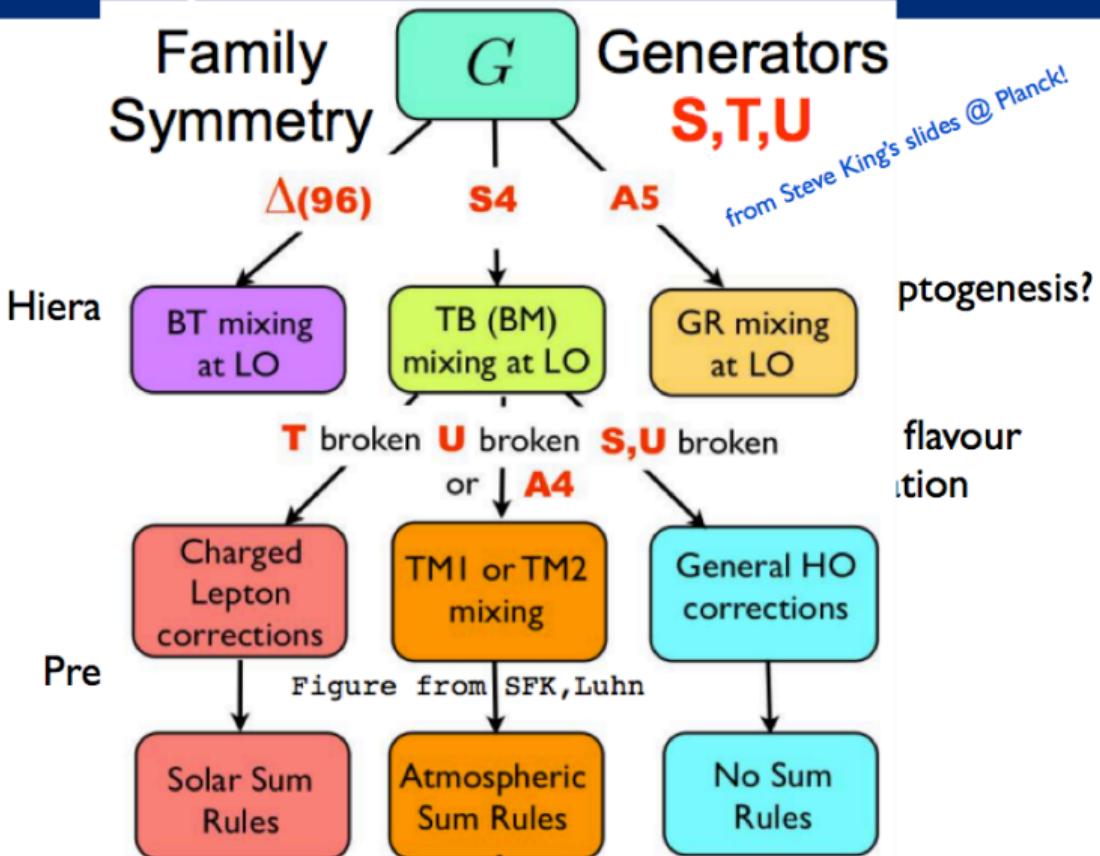
Precisi



Importance of SM for BSM



Importance of SM for BSM



What kind of BSM can be probed by neutrino experiments?

Can neutrinos probe reasonable new physics
beyond the reach of collider experiments?

Beyond standard physics with neutrinos

Non standard interactions

$$\mathcal{L}_{\text{NSI}} = -2\sqrt{2}G_F \varepsilon_{\alpha\beta}^{fP} (\bar{\nu}_\alpha \gamma^\mu \nu_\beta) (\bar{f} \gamma_\mu P f)$$

↓ ?

$$\frac{1}{\Lambda^2} (\bar{L}_\alpha \gamma^\rho L_\beta) (\bar{L}_\gamma \gamma_\rho L_\delta)$$

Flavour changing in charged lepton sector

Too big to be true

...

Flavour below G_F scale

Flavour remnants below G_F can show up
in neutrino oscillation experiments

Babu Friedland PANM Mocioiu, to appear soon

Gauge B-L of 3rd family

	ϕ_1	ϕ_2	s
$SU(2)_L$	2	2	1
$U(1)_Y$	+1	+1	0
$U(1)_{B-L}^{(3)}$	+1/3	0	+1/3

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breaks EW and new symmetry

Z-X mixing:

$$Z_\mu \simeq -s_w B_\mu + c_w W_\mu^3 - s_X X_\mu^0,$$

$$X_\mu \simeq s_X (s_w B_\mu - c_w W_\mu^3) + X_\mu^0,$$

$$s_X \equiv \frac{2}{3} \frac{g_X}{\sqrt{g^2 + g'^2}} \frac{v_1^2}{v^2}.$$

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Flavour changing neutral currents in quark sector:

$$\frac{g_X}{3} \bar{\mathbf{Q}}_L \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \gamma^\mu \mathbf{Q}_L X_\mu \longrightarrow \frac{g_X}{3} \bar{\mathbf{u}}_L \begin{pmatrix} V_{ub}^2 & V_{ub}V_{cb} & V_{ub} \\ V_{ub}V_{cb} & V_{cb}^2 & V_{cb} \\ V_{ub} & V_{cb} & 1 \end{pmatrix} \gamma^\mu \mathbf{u}_L X_\mu + \dots$$

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No FCNC (at tree level) in charged lepton sector!

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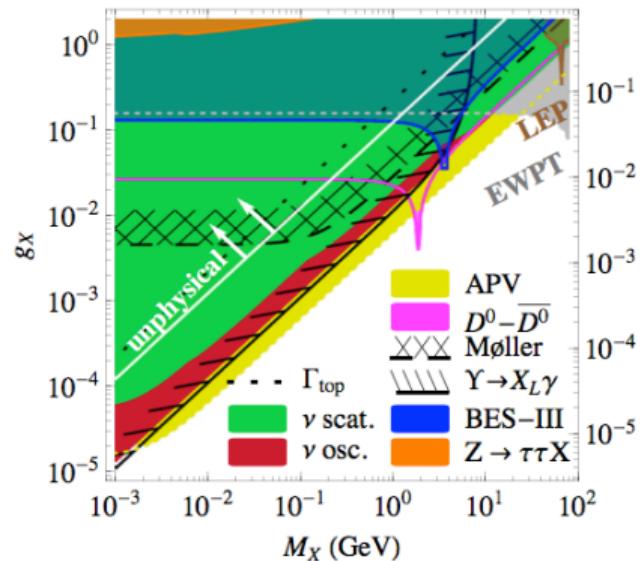
Neutrinos may experience large non standard matter effects

$$V_X \propto \text{diag}(0, 0, \varepsilon_{\tau\tau})$$

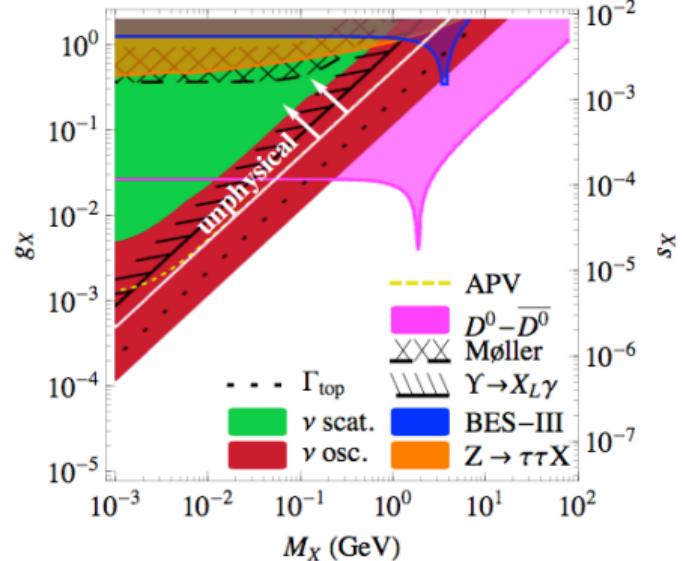
Flavour below G_F scale

Flavour remnants below G_F can show up
in neutrino oscillation experiments

$$\tan\beta = v_2/v_1 = 0.5$$



$$\tan\beta = v_2/v_1 = 10$$

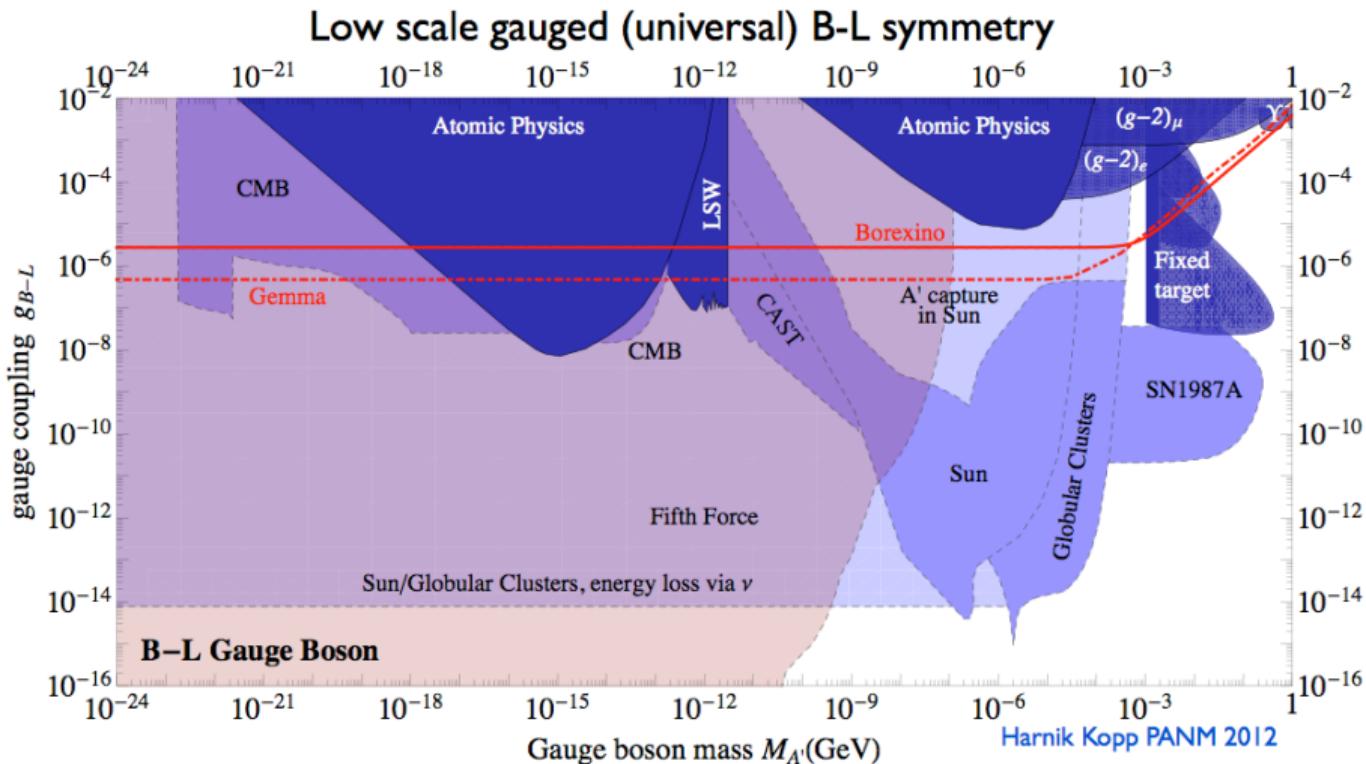


Babu Friedland PANM Mocioiu, to appear soon

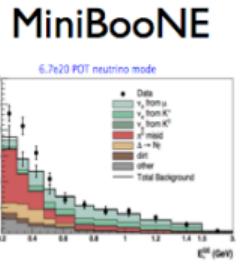
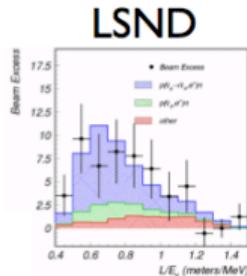
see also Farzan 2015, Farzan Shoemaker 2015

pedro.machado@uam.es

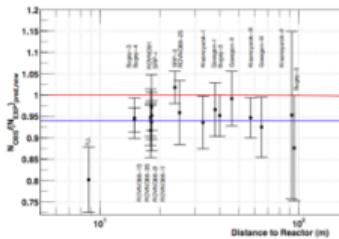
Light weakly coupled physics



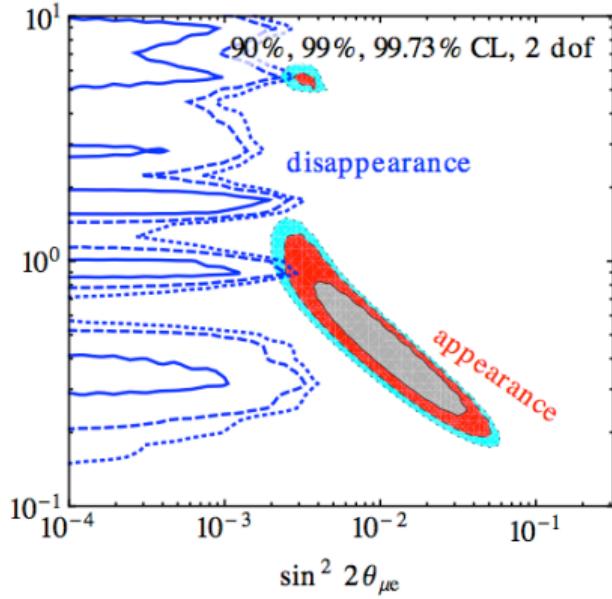
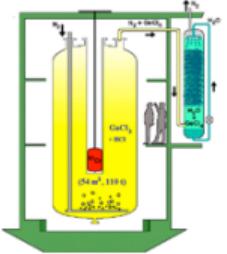
Sterile neutrinos



Reactor anomaly

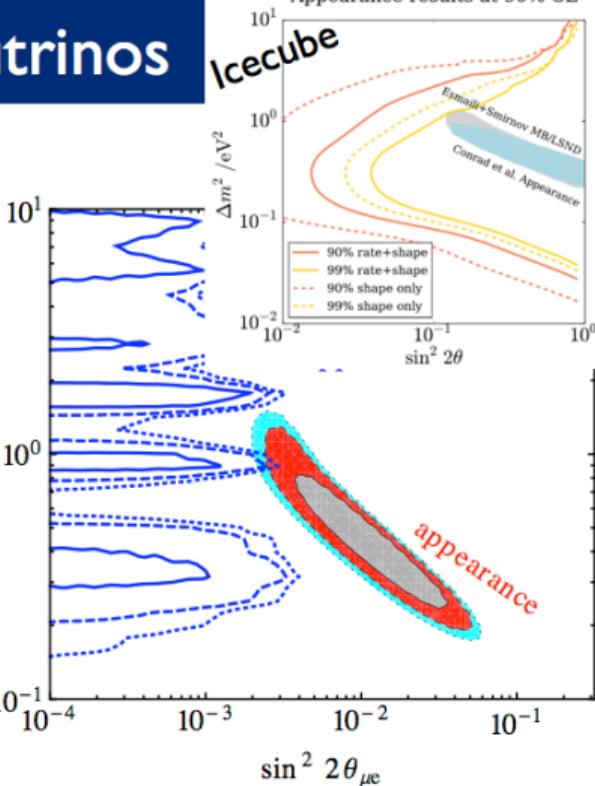
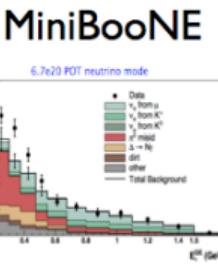
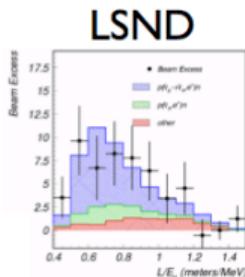


Gallium anomaly

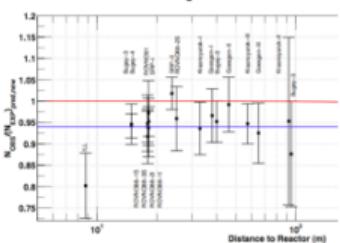


Kopp, PANM, Maltoni, Schwetz 2013

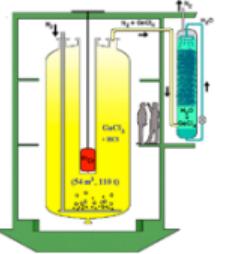
Sterile neutrinos



Reactor
anomaly



Gallium
anomaly



Kopp, PANM, Maltoni, Schwetz 2013

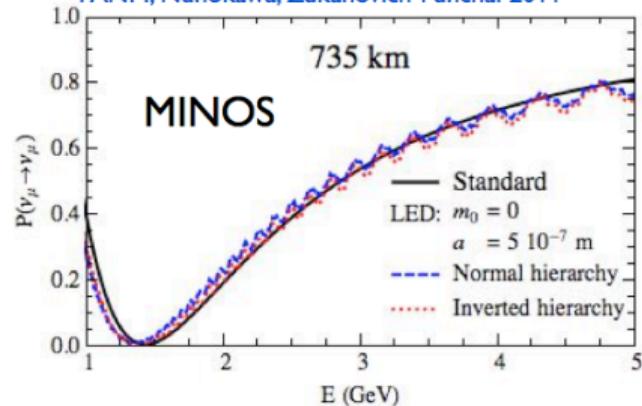
Neutrino masses

Neutrino masses from extra dimensions

Arkani-Hamed, Dimopoulos, March-Russel, Dienes, Dudas,
Gherghetta, Dvali, Smirnov, Barbieri, Creminelli, Strumia, ...

$$m_{\alpha\beta}^D = h_{\alpha\beta} \langle H \rangle M_{Pl(4+\delta)} / M_{Pl}$$

PANM, Nunokawa, Zukanovich-Funchal 2011



Could explain the reactor anomaly

PANM, Nunokawa, dos Santos, Zukanovich-Funchal 2011

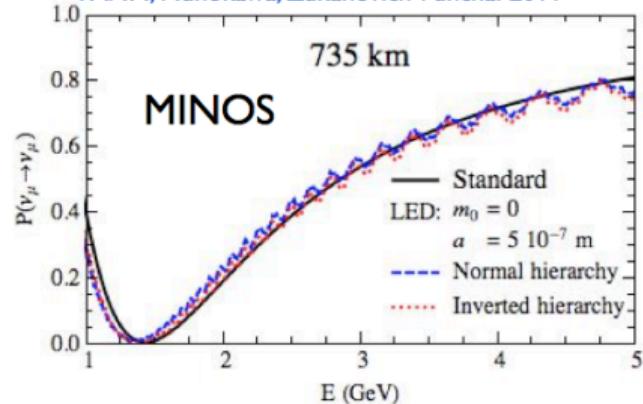
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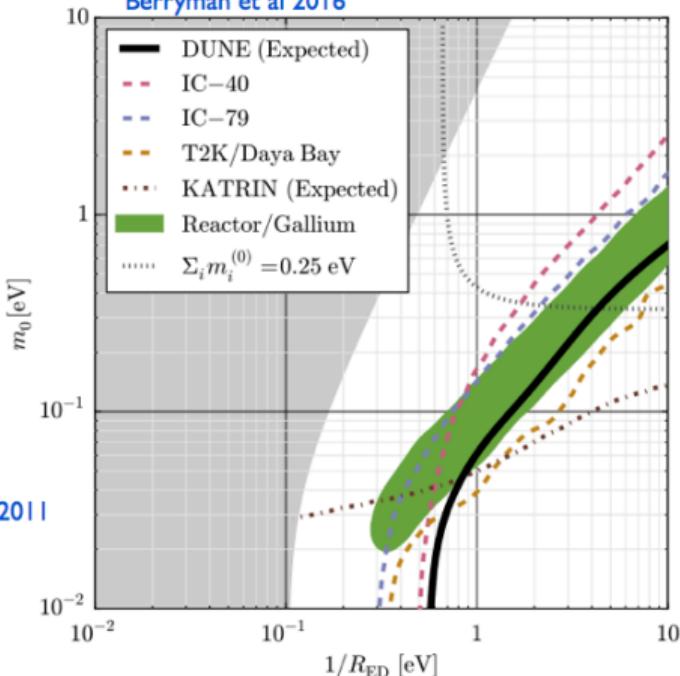
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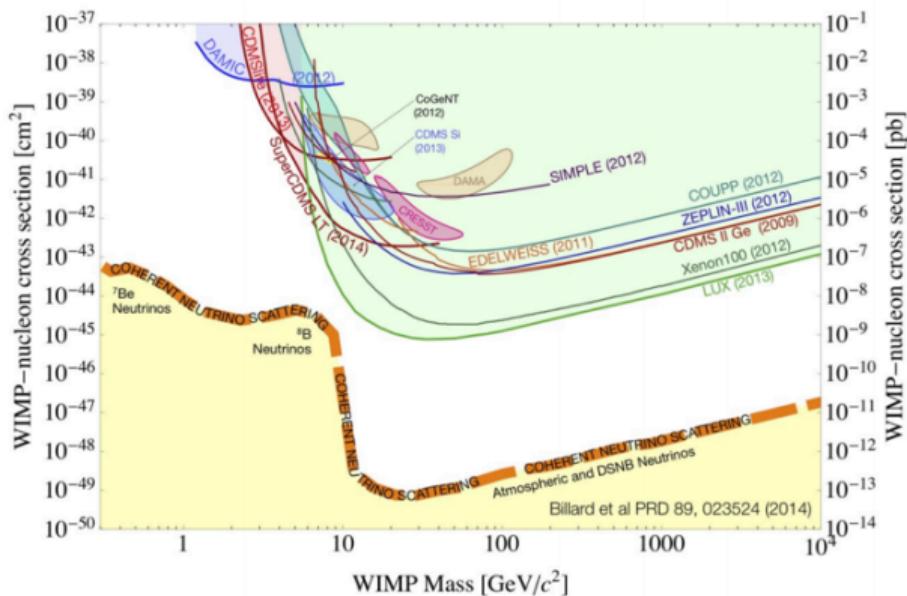
PANM, Nunokawa, dos Santos, Zukanovich-Funchal 2011

Berryman et al 2016



Beyond standard uses of neutrinos

DM direct detection experiments



Neutrino floor
DM-nucleus cross section for which the solar neutrino background hinders the experimental sensitivity

$$\frac{d\sigma_{\nu N}}{dE_R} = \frac{G_F^2}{4\pi} Q_v^2 m_N \left(1 - \frac{m_N E_R}{2E_\nu^2}\right) F^2(E_R)$$

$$Q_v = N - (1 - 4 \sin^2 \theta_W) Z$$

DM direct detection experiments

signal \longleftrightarrow **background**

see Adam Ritz's talk for
DM in neutrino experiments

What can be learned with
solar neutrinos in dark matter detectors?

Boehm, Cerdeno, Fairbairn, Jubb, PANM, Vincent 2016

Standard solar model
(opacity, metallicity, neutrinos)

BSM physics
(light mediators)

Standard model
(Coherent ν -N, $\sin^2\theta_w$)

Pospelov 2011, Harnik Kopp PANM 2012, Baudis et al 2014, Ruppin et al 2014, Billard et al 2015, Schumann et al 2015, Bilmis et al 2015, Dutta et al 2016, Strigari 2016, ...

Blois 2016

pedro.machado@uam.es

Neutrinos in DM experiments

Experiment	ϵ (ton-year)	$E_{th,n}$ (keV)	$E_{th,o}$ (keV)	E_{max} (keV)	$R(pp)$	$R(^8B)$	$R(CNO)$
G2-Ge	0.25	0.35	0.05	50	–	[62 – 85]	[0 – 3]
G2-Si	0.025	0.35	0.05	50	–	[3 – 3]	0
G2-Xe	25	3.0	2.0	30	[2104 – 2167]	[0 – 64]	0
Future-Xe	200	2.0	1.0	30	[17339 – 17846]	[520 – 10094]	0
Future-Ar	150	2.0	1.0	30	[14232 – 14649]	[6638 – 12354]	0
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Super-CDMS-like: 50 (5) kg Ge (Si) x 5 years

LZ-like: 5 ton Xe (fiducial) x 5 years

Darwin-like: 200 ton x y

Dark matter sensitivity of multi-ton
liquid xenon detectors

Marc Schumann,^{a,*} Laura Baudis,^b Lukas Bütkofer,^a
Alexander Kish,^b Marco Selvi^c

With an exposure of 200 t x y and assuming detector parameters which have been already demonstrated experimentally, spin-independent cross sections as low as $2.5 \times 10^{-49} \text{ cm}^2$ can be probed for WIMP masses around $40 \text{ GeV}/c^2$. Additional improvements in terms of background

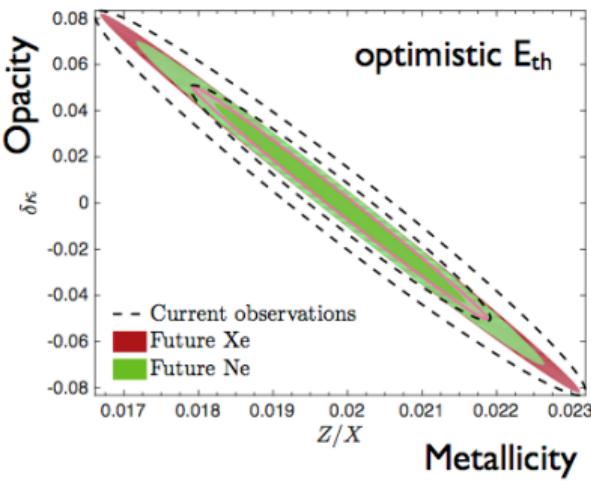
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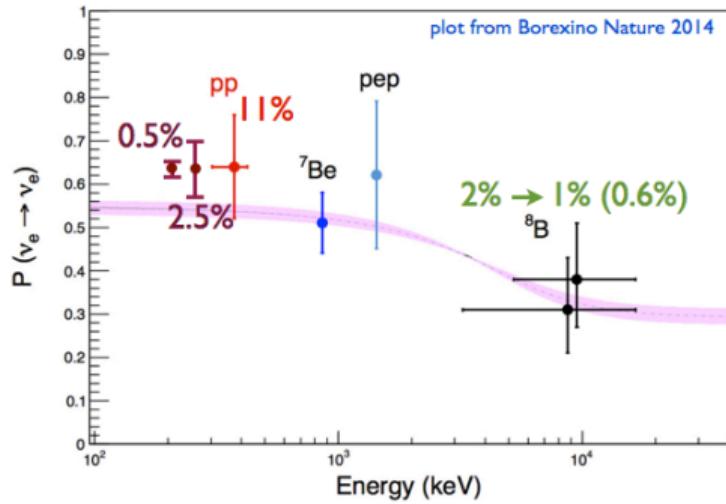
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Solar properties

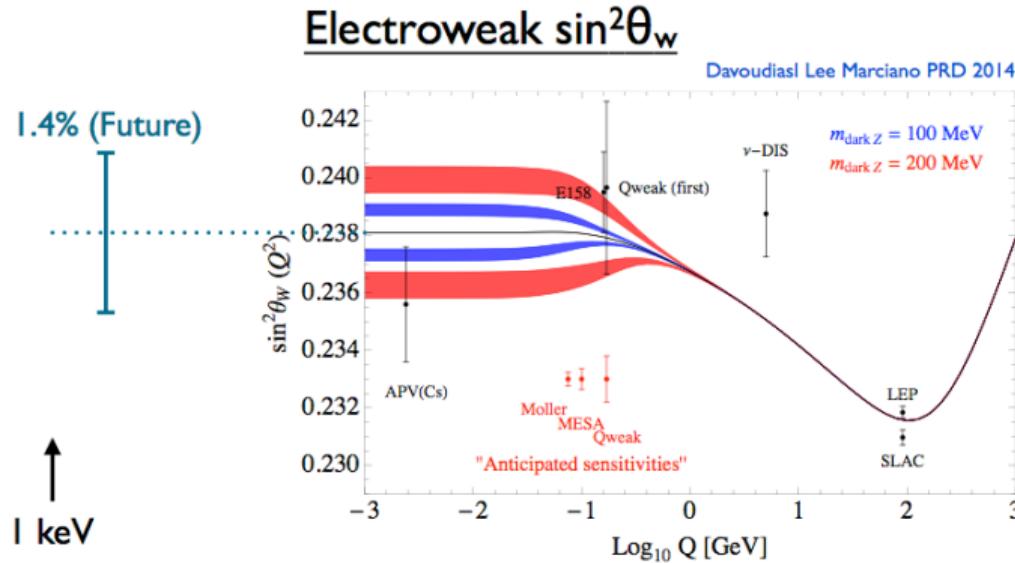


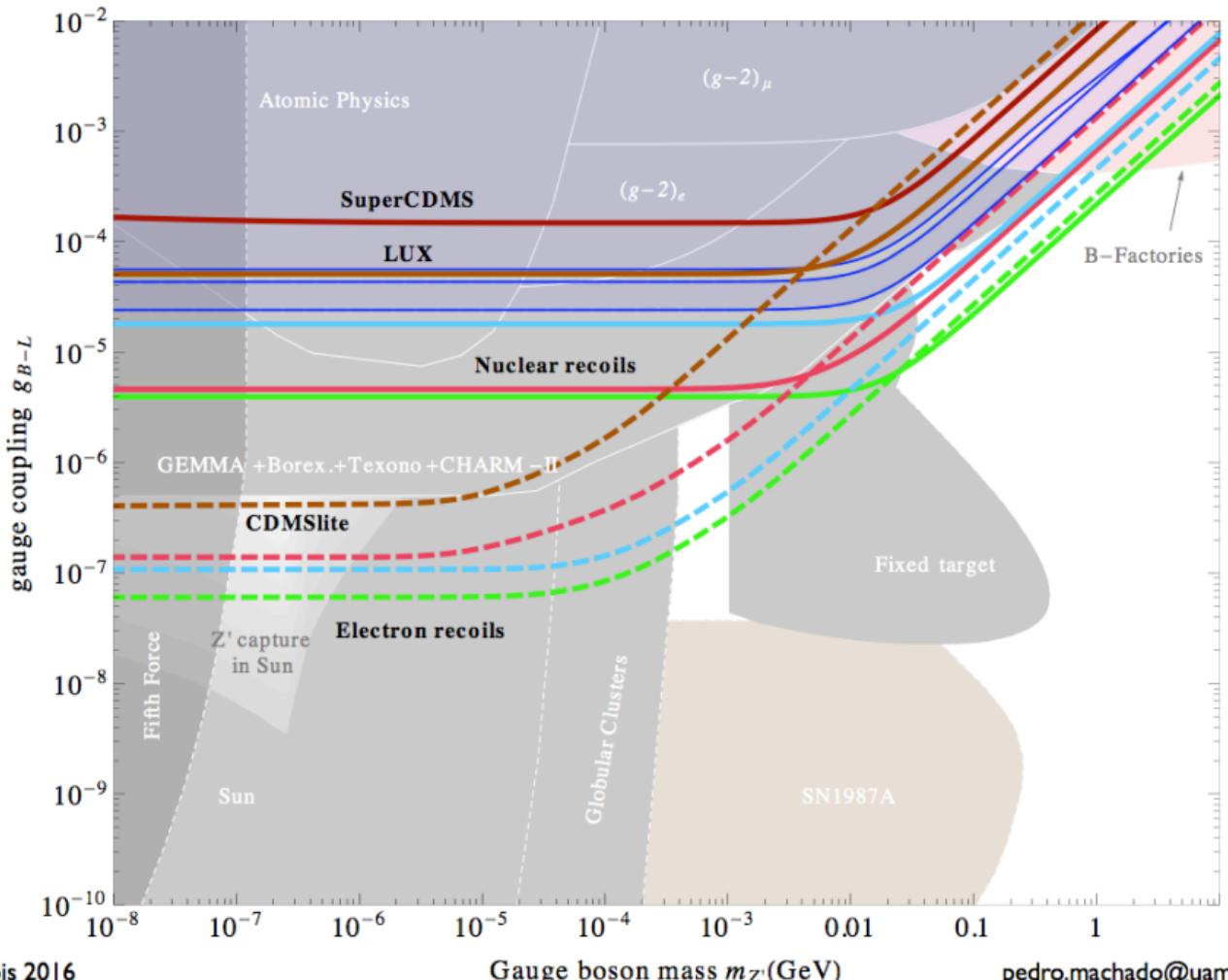
Solar neutrino fluxes



Neutrinos in DM experiments

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Conclusions

Neutrinos experiments can be sensitive to
physics beyond the reach of collider experiment

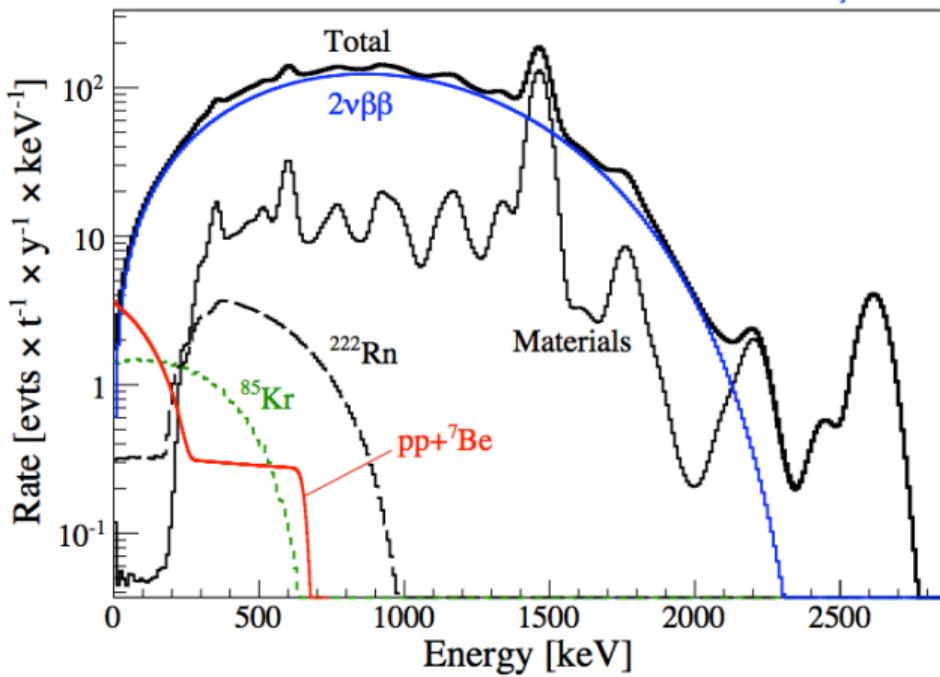
- Low scale flavour models
- Light mediator physics
- Sterile neutrinos ...

Neutrinos may become an important probe in
dark matter experiments

- Solar model parameters
- Neutrino fluxes ($p\bar{p}$, 8B)
- $\sin^2\theta_w$ at very low scales

Backgrounds

Baudis et al JCAP 2014



Backup

$$\mathcal{L}_{yuk}^q = \overline{\mathbf{Q}}_L \begin{pmatrix} y_{11}^u \tilde{\phi}_1 & y_{12}^u \tilde{\phi}_1 & y_{13}^u \tilde{\phi}_2 \\ y_{21}^u \tilde{\phi}_1 & y_{22}^u \tilde{\phi}_1 & y_{23}^u \tilde{\phi}_2 \\ 0 & 0 & y_{33}^u \tilde{\phi}_1 \end{pmatrix} \mathbf{u}_R + \overline{\mathbf{Q}}_L \begin{pmatrix} y_{11}^d \phi_1 & y_{12}^d \phi_1 & 0 \\ y_{21}^d \phi_1 & y_{22}^d \phi_1 & 0 \\ y_{31}^d \phi_2 & y_{32}^d \phi_2 & y_{33}^d \phi_1 \end{pmatrix} \mathbf{d}_R + \text{h.c.}$$

We choose to generate the CKM mixing on the up sector

$$R_{12}^{uL} \cdot M_u \cdot R_{12}^{uR\dagger} \approx \begin{pmatrix} m_u & 0 & V_{ub}^0 m_t \\ 0 & m_c & V_{cb}^0 m_t \\ 0 & 0 & m_t \end{pmatrix} \quad \text{and} \quad R_{12}^{dL} \cdot M_d \cdot R_{12}^{dR\dagger} \approx \begin{pmatrix} m_d^0 & 0 & 0 \\ 0 & m_s^0 & 0 \\ 0 & 0 & m_b^0 \end{pmatrix}$$

$$\mathcal{L}_{X-\text{FCNC}} = \frac{g_X}{3} \overline{\mathbf{Q}}_L \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \gamma^\mu X_\mu \mathbf{Q}_L$$

$$\mathcal{L}_{X-\text{FCNC}} \approx \frac{g_X}{3} \overline{\mathbf{u}}_L \begin{pmatrix} V_{ub}^2 & V_{ub} V_{cb} & V_{ub} \\ V_{ub} V_{cb} & V_{cb}^2 & V_{cb} \\ V_{ub} & V_{cb} & 1 \end{pmatrix} \gamma^\mu X_\mu \mathbf{u}_L + \frac{g_X}{3} \overline{\mathbf{d}}_L \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \gamma^\mu X_\mu \mathbf{d}_L$$