

# ***Type Ia Supernovae meet Maximum Likelihood Estimators***

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J. Trøst Nielsen, AG and S. Sarkar, arXiv:1506.01354

J. Trøst Nielsen, arXiv:1508.07850

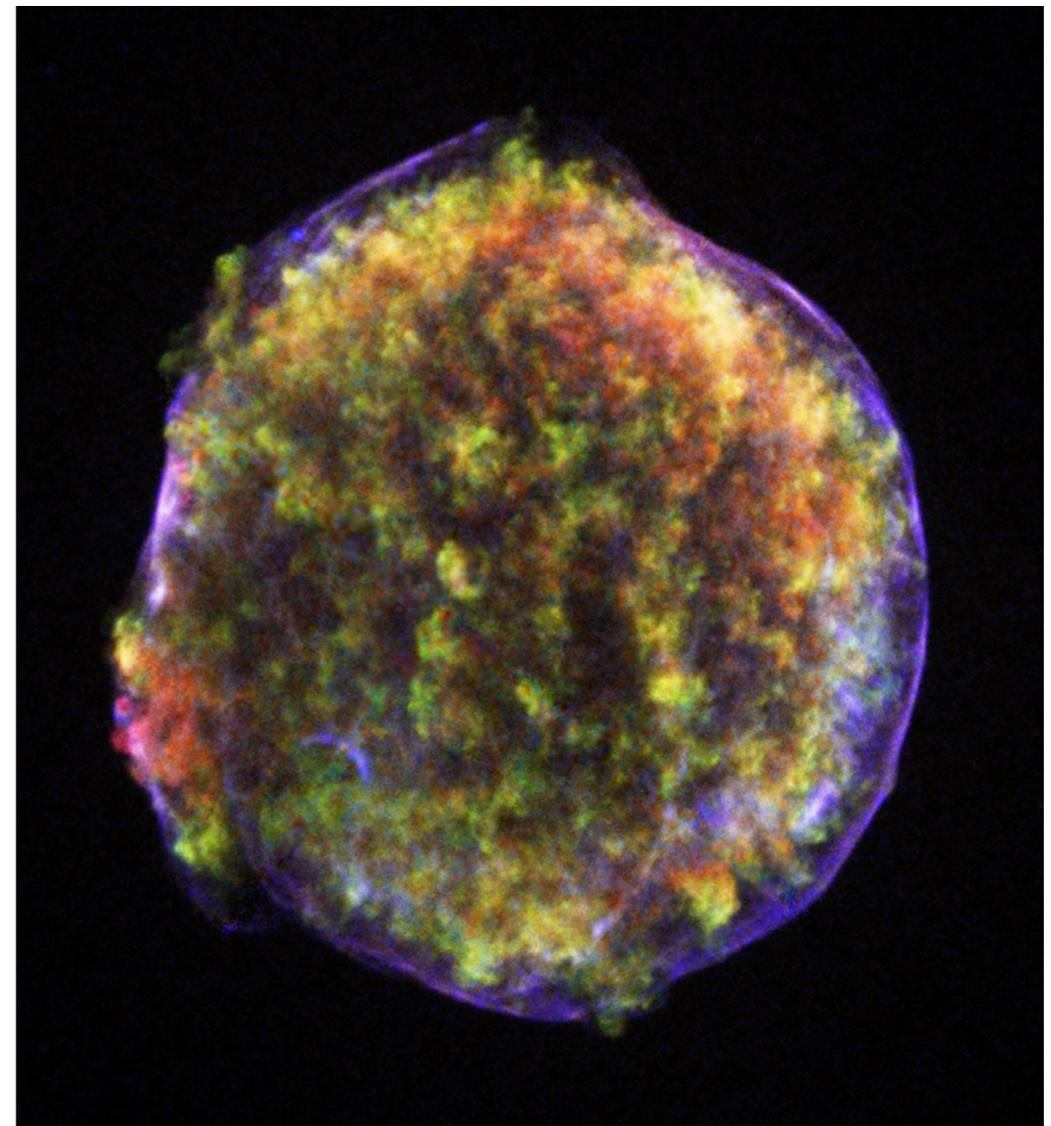
XXVIII<sup>th</sup> Rencontres de Blois - Particle Physics and Cosmology

Blois, 29 May - 3 June, 2016

# Type Ia Supernovae



Suzuki *et al*, 1105.3470



**SN 1572 (Tycho)**  
NASA/CXC/Rutgers/Warren & Hughes



# Type Ia Supernovae

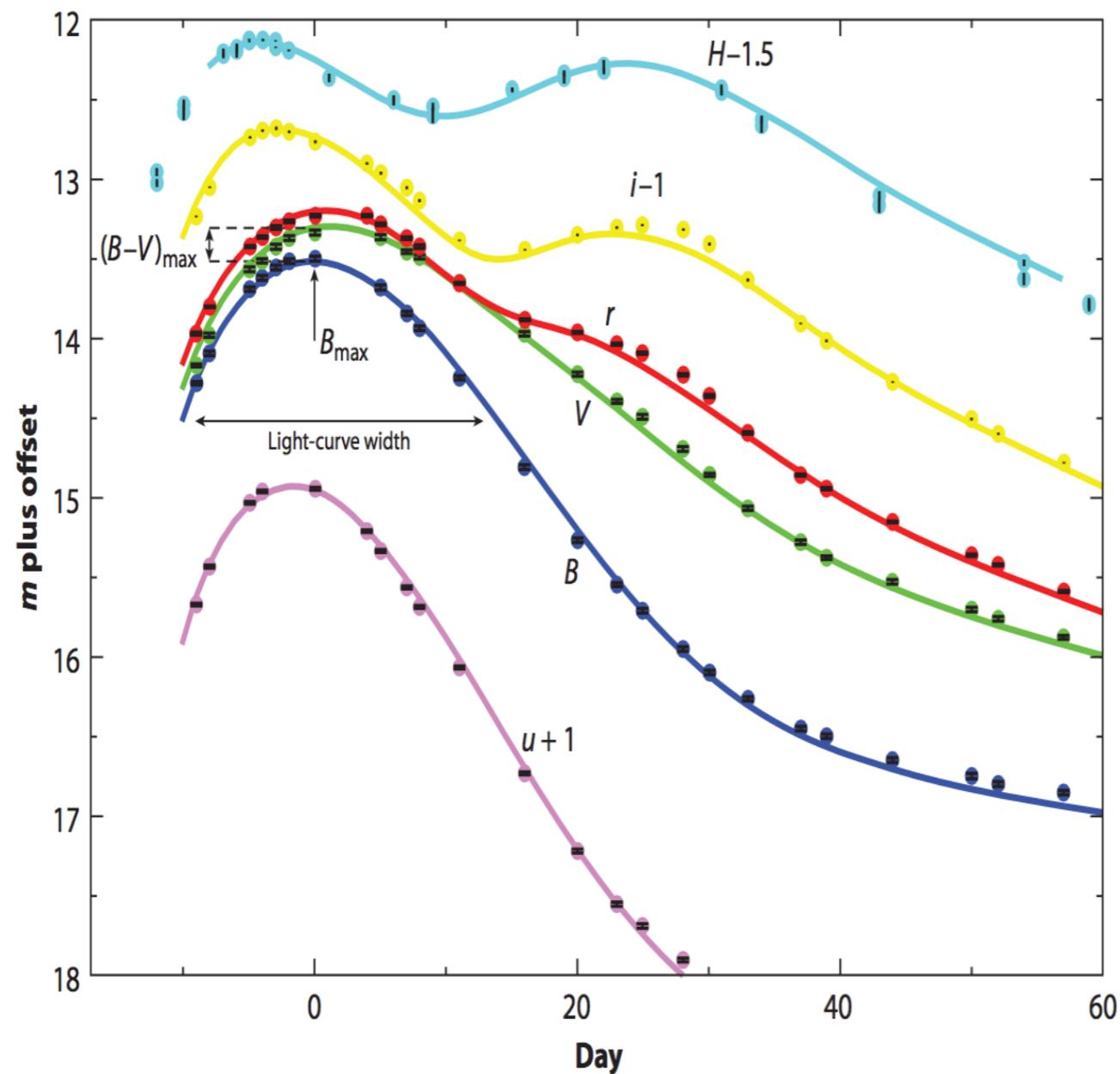
Diagram illustrating the classification of Supernovae (SN) based on their chemical composition and explosion mechanism:

- SN** (Supernova) branches into:
  - no H** (Hydrogen):
    - Si** (Silicon): **Type Ia** (Thermonuclear)
    - no Si**:
      - He** (Helium): **Type Ib**
      - no He**: **Type Ic**
  - H** (Hydrogen): **Type II**

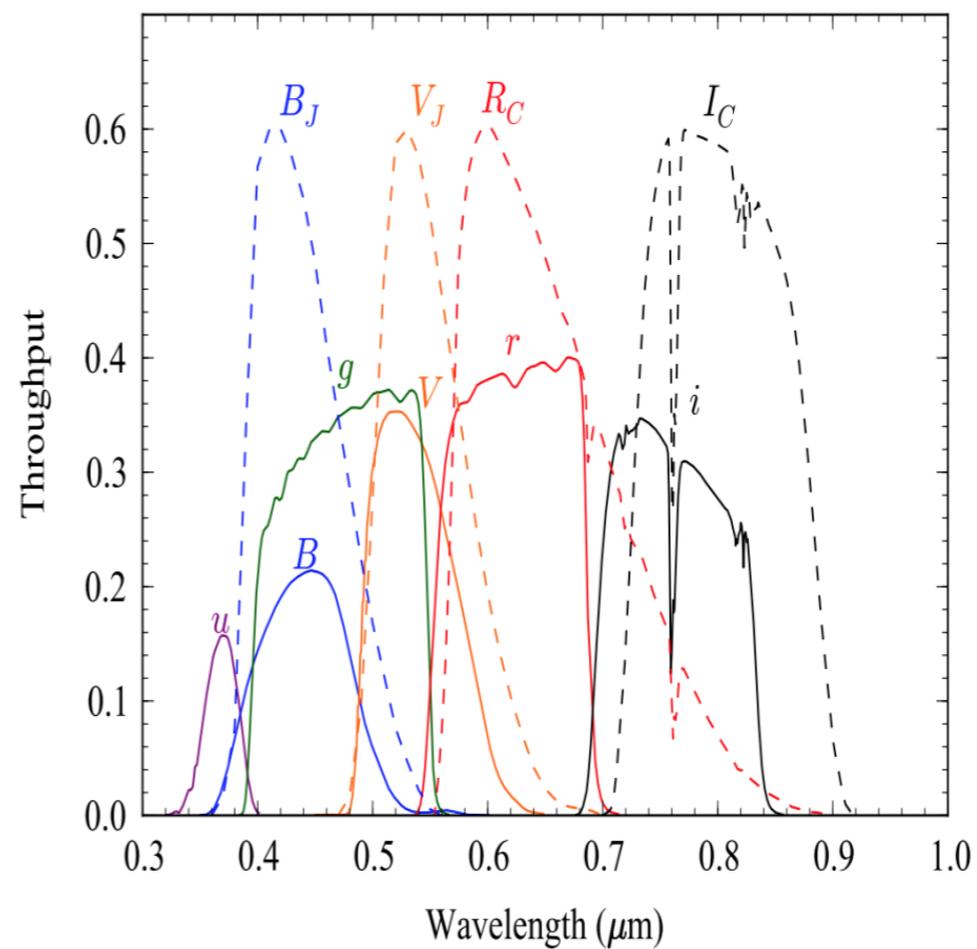
The **Type Ib**, **Type Ic**, and **Type II** categories are collectively labeled as **Core Collapse**. A red 'X' is drawn over the Core Collapse section of the diagram.



# Type Ia Supernovae



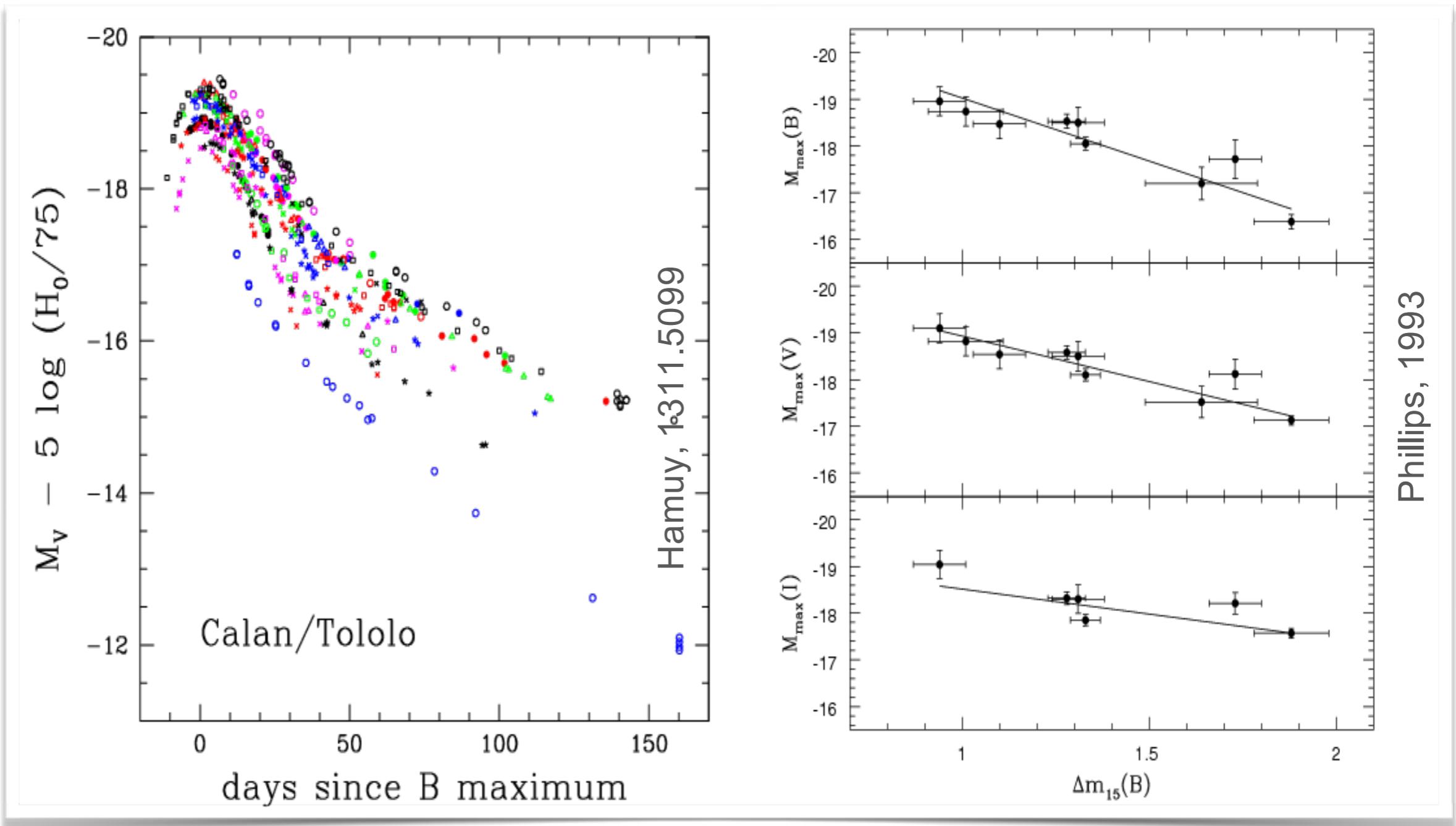
$$m = -2.5 \log(F/F_{\text{ref}})$$



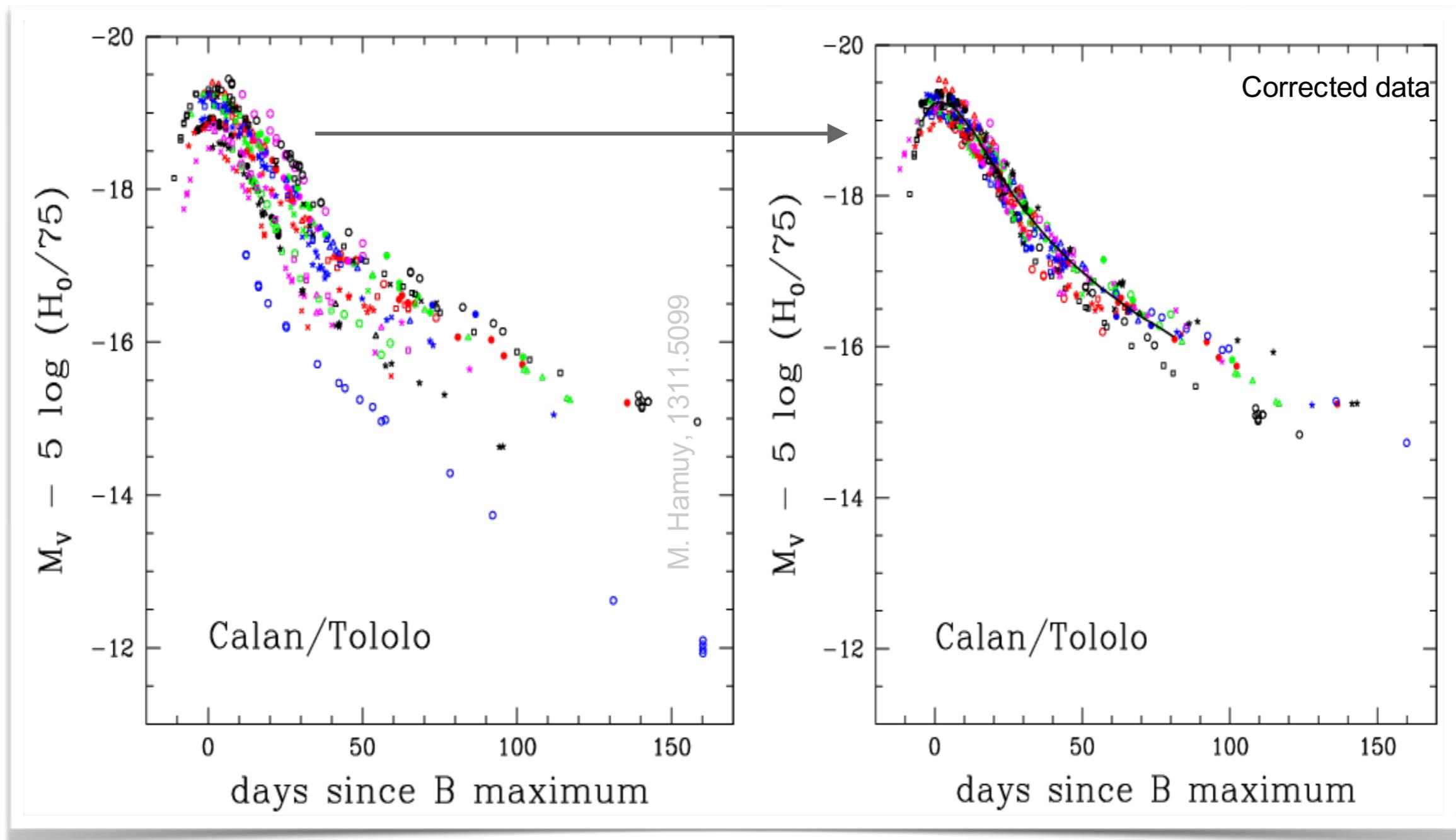
Goobar & Leibundgut, 1102.1431



# Type Ia Supernovae - Standard Candles?



# Type Ia Supernovae - Standardizable Candles



# Type Ia Supernovae - Standardizable Candles

- We take as input to our analysis SN Ia standardises using the SALT2 LC fitting procedure [Guy et al, 2007, 2010]

SALT 2 parameters						Betoule <i>et al.</i> , 1401.4064
Name	$z_{\text{cmb}}$	$m_B^*$	$X_1$	$C$	$M_{\text{stellar}}$	?
03D1ar	0.002	$23.941 \pm 0.033$	$-0.945 \pm 0.209$	$0.266 \pm 0.035$	$10.1 \pm 0.5$	?
03D1au	0.503	$23.002 \pm 0.088$	$1.273 \pm 0.150$	$-0.012 \pm 0.030$	$9.5 \pm 0.1$	?
03D1aw	0.581	$23.574 \pm 0.090$	$0.974 \pm 0.274$	$-0.025 \pm 0.037$	$9.2 \pm 0.1$	?
03D1ax	0.495	$22.960 \pm 0.088$	$-0.729 \pm 0.102$	$-0.100 \pm 0.030$	$11.6 \pm 0.1$	?
03D1bp	0.346	$22.398 \pm 0.087$	$-1.155 \pm 0.113$	$-0.041 \pm 0.027$	$10.8 \pm 0.1$	?
03D1co	0.678	$24.078 \pm 0.098$	$0.619 \pm 0.404$	$-0.039 \pm 0.067$	$8.6 \pm 0.3$	?
03D1dt	0.611	$23.285 \pm 0.093$	$-1.162 \pm 1.641$	$-0.095 \pm 0.050$	$9.7 \pm 0.1$	
03D1ew	0.866	$24.354 \pm 0.106$	$0.376 \pm 0.348$	$-0.063 \pm 0.068$	$8.5 \pm 0.8$	
03D1fc	0.331	$21.861 \pm 0.086$	$0.650 \pm 0.119$	$-0.018 \pm 0.024$	$10.4 \pm 0.0$	
03D1fq	0.799	$24.510 \pm 0.102$	$-1.057 \pm 0.407$	$-0.056 \pm 0.065$	$10.7 \pm 0.1$	
03D3aw	0.450	$22.667 \pm 0.092$	$0.810 \pm 0.232$	$-0.086 \pm 0.038$	$10.7 \pm 0.0$	
03D3ay	0.371	$22.273 \pm 0.091$	$0.570 \pm 0.198$	$-0.054 \pm 0.033$	$10.2 \pm 0.1$	
03D3ba	0.292	$21.961 \pm 0.093$	$0.761 \pm 0.173$	$0.116 \pm 0.035$	$10.2 \pm 0.1$	
03D3bl	0.356	$22.927 \pm 0.087$	$0.056 \pm 0.193$	$0.205 \pm 0.030$	$10.8 \pm 0.1$	

$$\mu_B = m_B^* - M + \alpha X_1 - \beta C$$


# Distance moduli as probes of cosmology

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$$F = \frac{L}{4\pi d_L^2}$$

$$d_L = (1+z) \frac{d_H}{\sqrt{\Omega_k}} \sinh \left( \sqrt{\Omega_k} \int_0^z \frac{H_0 dz'}{H(z')} \right)$$

$$\mu_C = m - M = -2.5 \log \frac{F/F_{\text{ref}}}{L/L_{\text{ref}}} = 5 \log \frac{d_L}{10\text{pc}}$$

$$\chi^2 = \sum_{\text{objects}} \frac{(\mu_B - 5 \log_{10}(d_L(\theta, z)/10\text{pc}))^2}{\sigma^2(\mu_B) + \sigma_{\text{int}}^2}$$

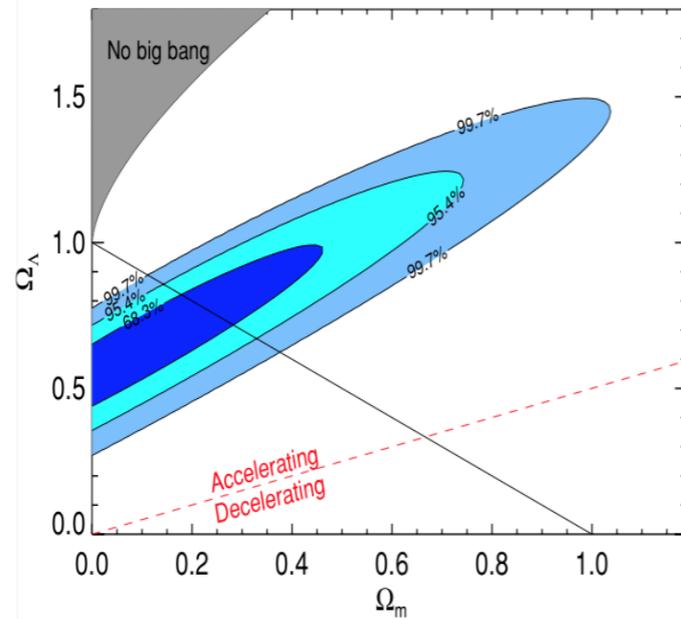
Astier et al., 2006



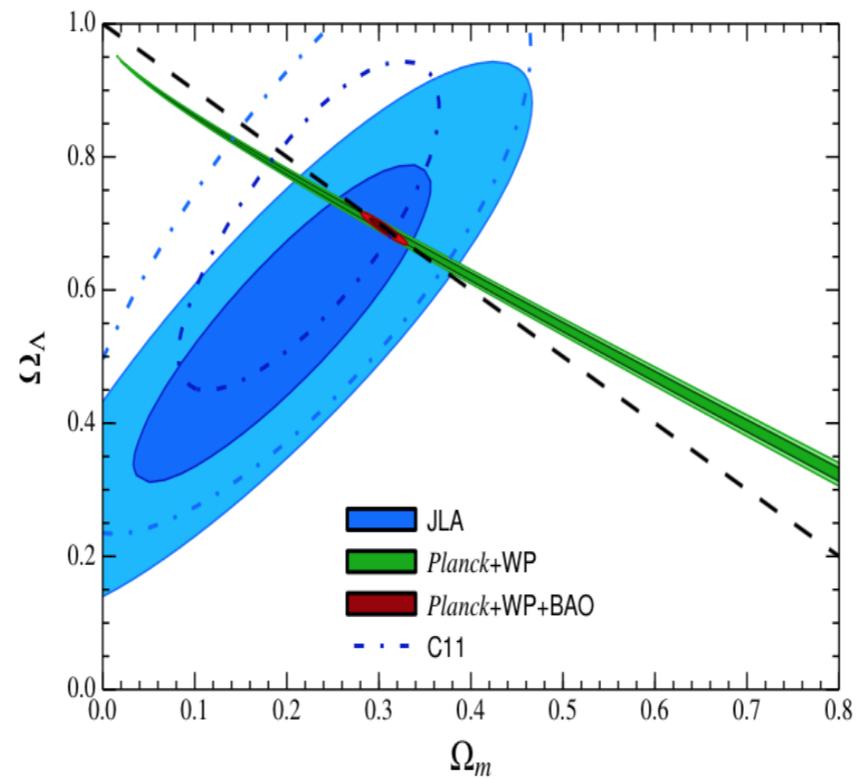
# Supernova Cosmology: the conventional wisdom

“SN data alone require cosmic acceleration at >99.999% confidence, including systematic effects”

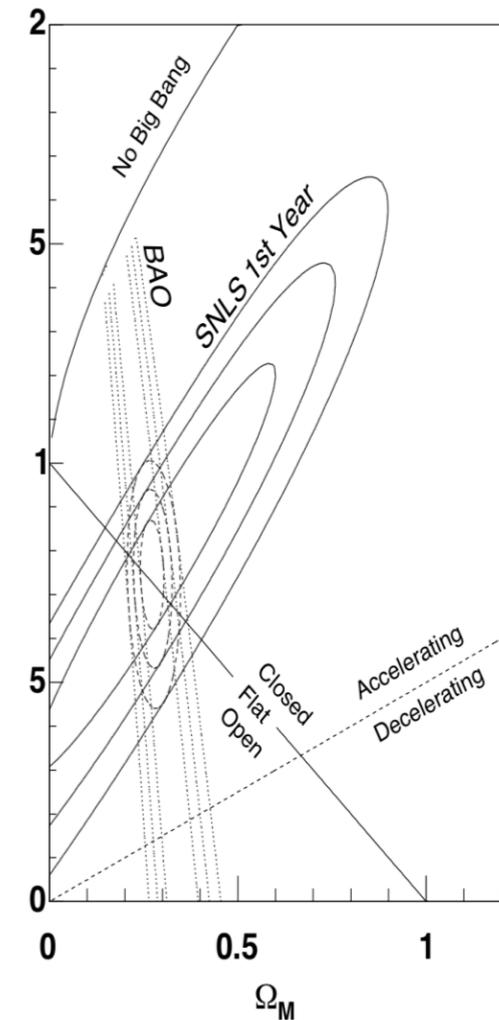
Conley *et al*, 2011



Betoule *et al*, 2014



Astier *et al*, 2006



# A critical appraisal of the standard constrained $\chi^2$

$$\mu_B = m_B^{\max} + \alpha \cdot x_1 - \beta \cdot c + \delta \cdot P(m_*^{\text{true}} < m_*^{\text{threshold}}) - M_B,$$

where  $M_B$  is the absolute  $B$ -band magnitude of a SN Ia with  $x_1 = 0$ ,  $c = 0$  and  $P(m_*^{\text{true}} < m_*^{\text{threshold}}) = 0$ . The parameters  $\alpha$ ,  $\beta$ ,  $\delta$  and  $M_B$  are nuisance parameters that are fitted simultaneously with the cosmological parameters. The SN Ia

## 4.5. Systematic errors

In this paper, we follow the systematics analysis we presented in Amanullah et al. (2010). Systematic errors that directly affect supernova distance measurements (calibration, and galactic extinction, for example) are treated as nuisance parameters to be fit simultaneously with the cosmology. Minimizing over these nuisance parameters gives additional terms to add to the distance modulus covariance matrix

$$U_{ij} = \sum_{\epsilon} \frac{d\mu_i(\alpha, \beta)}{d\epsilon} \frac{d\mu_j(\alpha, \beta)}{d\epsilon} \sigma_{\epsilon}^2, \quad (5)$$

where the sum is over each of these distance systematic errors in the analysis. (Although the distance modulus depends on  $\delta$  as well as  $\alpha$  and  $\beta$ , the derivatives with respect to the zero-points do not.) In this analysis,  $\alpha$  and  $\beta$  have little interaction with cosmological parameters. When computing cosmologi-

## 4.4. Fitting the Cosmology

Following Amanullah et al. (2010), the best-fit cosmology is determined by minimizing

$$\chi_{\text{stat}}^2 = \sum_{\text{SNe}} \frac{[\mu_B(\alpha, \beta, \delta, M_B) - \mu(z; \Omega_m, \Omega_w, w)]^2}{\sigma_{\text{lc}}^2 + \sigma_{\text{ext}}^2 + \sigma_{\text{sample}}^2}. \quad (4)$$

A detailed discussion of the terms in this equation can be found in Amanullah et al. (2010). We only comment on the final term in the denominator,  $\sigma_{\text{sample}}^2$ , which is computed by setting the reduced  $\chi^2$  of each sample to unity. This term was referred to as “ $\sigma_{\text{systematic}}^2$ ” in Kowalski et al. (2008); Amanullah et al. (2010). We note that  $\sigma_{\text{sample}}^2$  includes intrinsic dispersion as well as sample-dependent effects. This term effectively further deweights samples with poorer-quality data that has sources of error which have not been accounted for. As noted in Amanullah et al. (2010), this may occasionally deweight an otherwise well-measured supernova.



# A critical appraisal of the standard constrained $\chi^2$

$$\mu_B = m_B^{\max} + \alpha \cdot x_1 - \beta \cdot c + \delta \cdot P(m_*^{\text{true}} < m_*^{\text{threshold}}) - M_B,$$

## C. EMPIRICAL ADJUSTMENT OF UNCERTAINTIES

The propagated uncertainties are underestimates, as they do not account for the pixel-pixel covariance introduced by warping, sub-sampling, stacking, and convolution of the images. In order to empirically determine by how much the uncertainties are underestimated, we measure the flux  $f_r$  and its uncertainty  $\sigma_r$  at random positions in a given difference image in exactly the same way we measure the SN flux. We calculate the weighted mean  $\bar{f}_r$  of these flux measurements. In order to guard against reduction and image artifacts, we apply a  $3\sigma$  cut to the normalized flux distribution  $(f_r - \bar{f}_r)/\sigma_r$ , rather than cutting on the underestimated errors,  $\sigma_r$ , for the following reason: let's assume that all uncertainties are underestimated by the same factor  $s_r$ . If we nominally apply a N-sigma cut using these underestimated uncertainties, we effectively apply an  $N/s_r$ -cut, e.g. for a nominal 3-sigma cut and  $s_r = 1.5$ , the real cut-off is at 2-sigma. In order to avoid this, we determine the normalized flux distribution  $(f_r - \bar{f}_r)/\sigma_r$ , which has a standard deviation of  $s_r$ . The true 3-sigma outliers can then be identified and removed by doing a 3-sigma cut on the normalized flux distribution. Note that the standard deviation  $s_r$  is equivalent to the square-root of the chi-square distribution

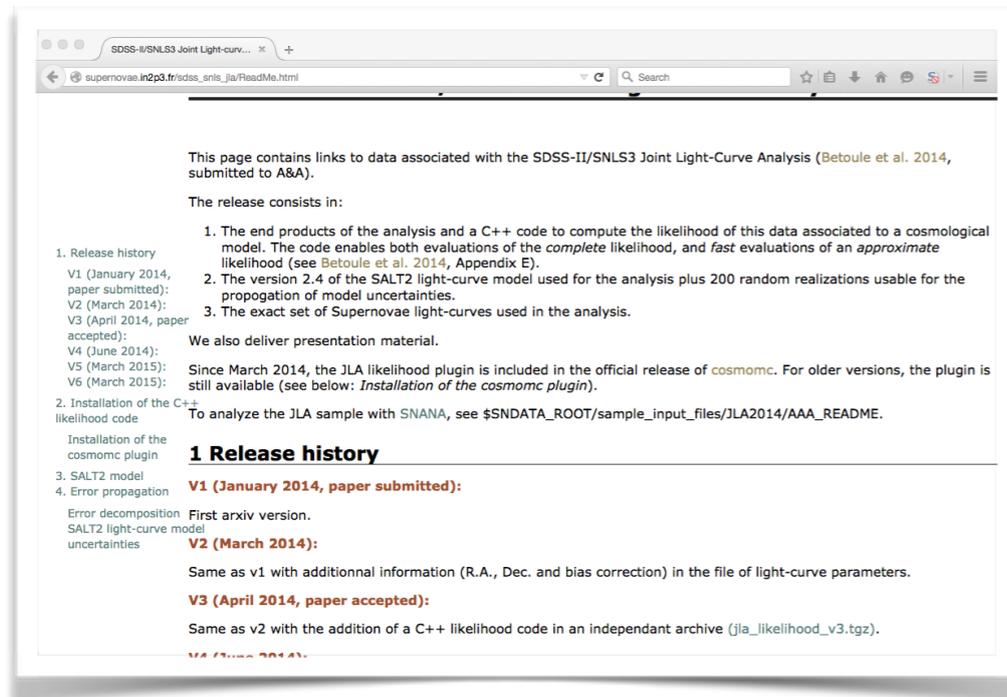
$$s_r = \sqrt{\chi_r^2} = \frac{1}{N-2} \sum \left( \frac{f_r - \bar{f}_r}{\sigma_r} \right)^2 \quad (\text{C1})$$

We multiply all uncertainties by the factor  $s_r$  in order to empirically correct the uncertainties. **We find that it is imperative to employ this robust way of determining  $s_r$  for the method to work correctly. The fact that the reduced chi-square of the baseline flux of the SN light curves peaks at 1.0 validates our method (see §5.3.3).**

~~In addition, for a given difference image,  $\bar{f}_r$  is an estimate of the bias in the flux measurements.~~ The values of  $\bar{f}_r$  are in general very small, much smaller than the typical uncertainties. Nevertheless, we adjust all fluxes by this value.



# Data - JLA catalogue



This page contains links to data associated with the SDSS-II/SNLS3 Joint Light-Curve Analysis (Betoule et al. 2014, submitted to A&A).

The release consists in:

1. The end products of the analysis and a C++ code to compute the likelihood of this data associated to a cosmological model. The code enables both evaluations of the *complete likelihood*, and *fast evaluations of an approximate likelihood* (see Betoule et al. 2014, Appendix E).
2. The version 2.4 of the SALT2 light-curve model used for the analysis plus 200 random realizations usable for the propagation of model uncertainties.
3. The exact set of Supernovae light-curves used in the analysis.

We also deliver presentation material.

Since March 2014, the JLA likelihood plugin is included in the official release of *cosmomc*. For older versions, the plugin is still available (see below: *Installation of the cosmomc plugin*).

To analyze the JLA sample with *SNANA*, see `$SNDATA_ROOT/sample_input_files/JLA2014/AAA_README`.

### 1 Release history

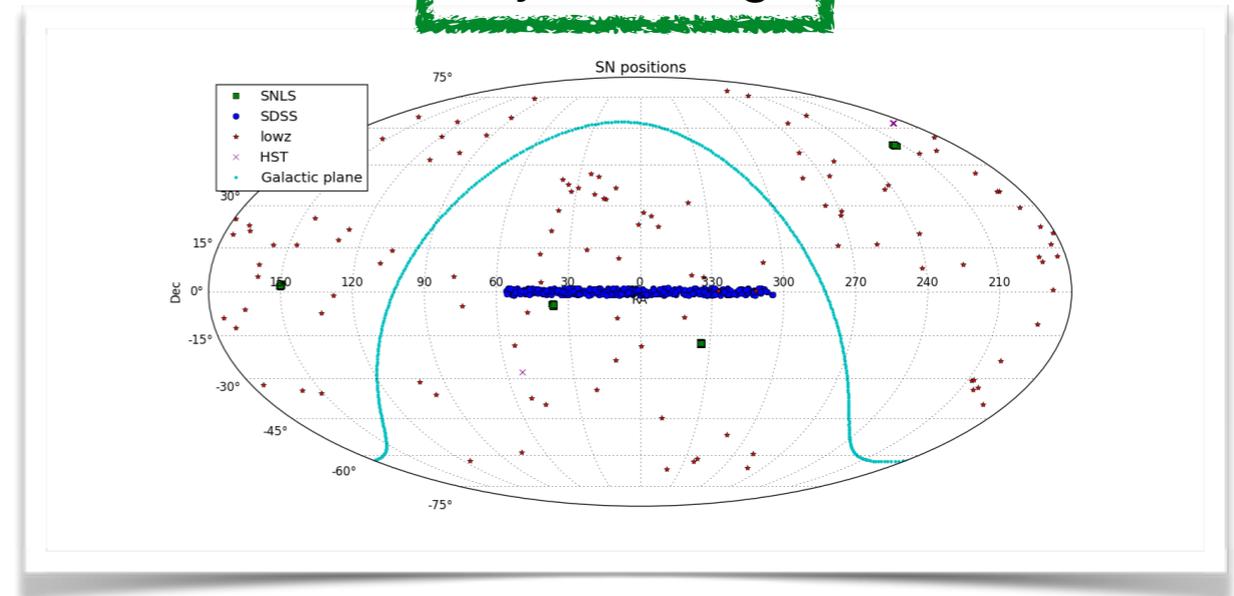
**V1 (January 2014, paper submitted):**  
First arxiv version.

**V2 (March 2014):**  
Same as v1 with additional information (R.A., Dec. and bias correction) in the file of light-curve parameters.

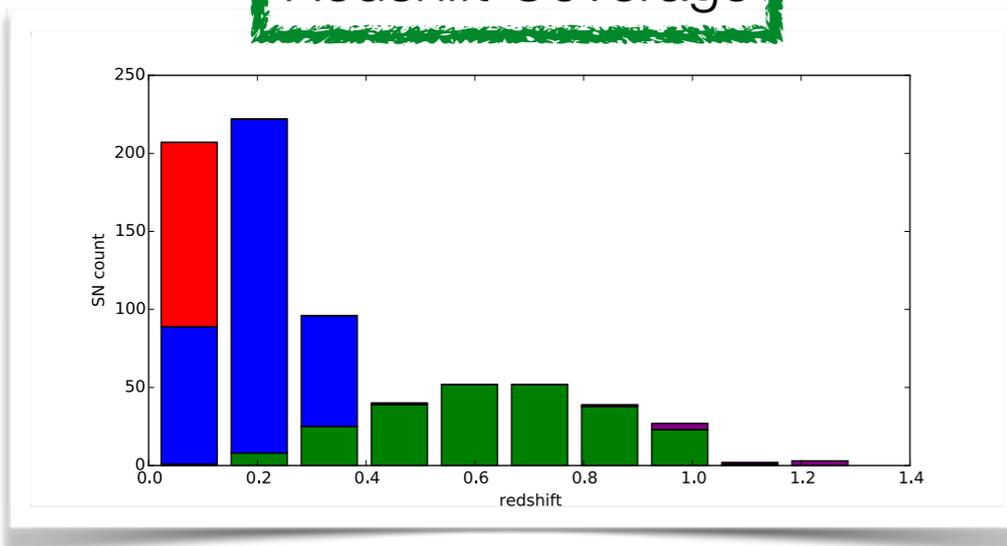
**V3 (April 2014, paper accepted):**  
Same as v2 with the addition of a C++ likelihood code in an independant archive (`jla_likelihood_v3.tgz`).

**V4 (June 2014):**

## Sky Coverage



## Redshift Coverage



Our analysis is based on (and makes use of) the public information available through from the JLA website



# *Maximum Likelihood Estimator analysis*

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***It all starts with a likelihood***

$\mathcal{L}$  = “probability density”(data|model)

## **Disclaimer**

(for my statistics-savvy friends)

I apologise for the abuse of language of defining the likelihood function as a probability density, since technically it is not normalised ... life (and talks) are often too short to discuss all the details



# Maximum Likelihood Estimator analysis

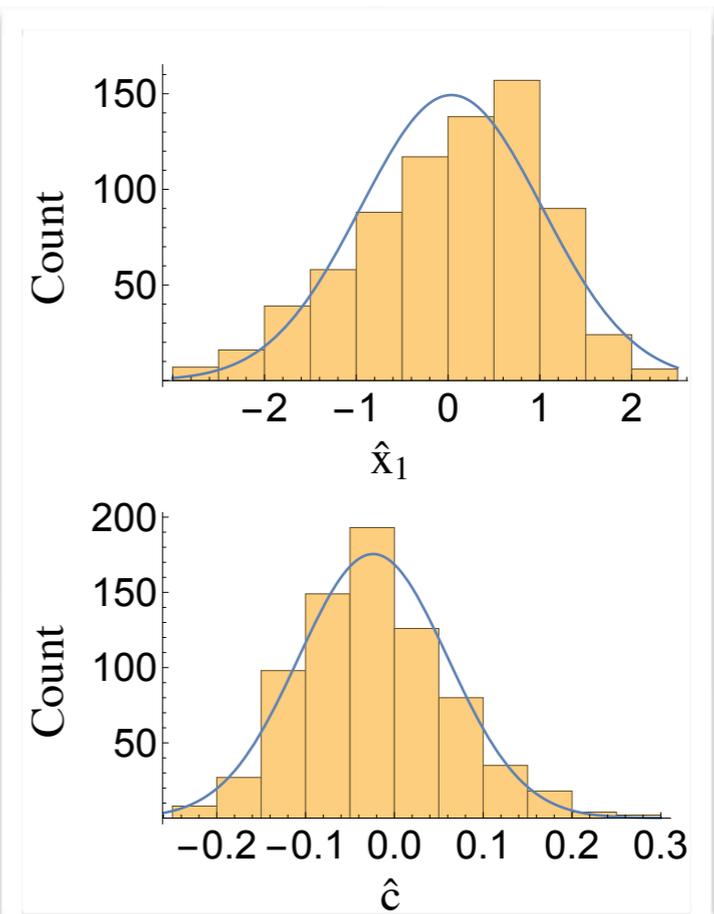
$$\begin{aligned}\mathcal{L} &= p[(\hat{m}_B^*, \hat{x}_1, \hat{c})|\theta] \\ &= \int p[(\hat{m}_B^*, \hat{x}_1, \hat{c})|(M, x_1, c), \theta] p[(M, x_1, c)|\theta] dM dx_1 dc,\end{aligned}$$

$p[(M, x_1, c)|\theta] = p(M|\theta)p(x_1|\theta)p(c|\theta)$ , where:

$$p(M|\theta) = (2\pi\sigma_{M_0}^2)^{-1/2} \exp\left\{-\frac{[(M - M_0) / \sigma_{M_0}]^2}{2}\right\},$$

$$p(x_1|\theta) = (2\pi\sigma_{x_{1,0}}^2)^{-1/2} \exp\left\{-\frac{[(x_1 - x_{1,0}) / \sigma_{x_{1,0}}]^2}{2}\right\},$$

$$p(c|\theta) = (2\pi\sigma_{c_0}^2)^{-1/2} \exp\left\{-\frac{[(c - c_0) / \sigma_{c_0}]^2}{2}\right\}.$$



# Maximum Likelihood Estimator analysis

$$p(Y|\theta) = \frac{1}{\sqrt{|2\pi\Sigma_l|}} \exp \left[ -\frac{1}{2}(Y - Y_0)\Sigma_l^{-1}(Y - Y_0)^T \right]$$

$$p(\hat{X}|X, \theta) = \frac{1}{\sqrt{|2\pi\Sigma_d|}} \exp \left[ -\frac{1}{2}(\hat{X} - X)\Sigma_d^{-1}(\hat{X} - X)^T \right]$$

$$\mathcal{L} = \frac{1}{\sqrt{|2\pi(\Sigma_d + A^T\Sigma_l A)|}} \times \exp \left( -\frac{1}{2}(\hat{Z} - Y_0 A)(\Sigma_d + A^T\Sigma_l A)^{-1}(\hat{Z} - Y_0 A)^T \right)$$

intrinsic distributions

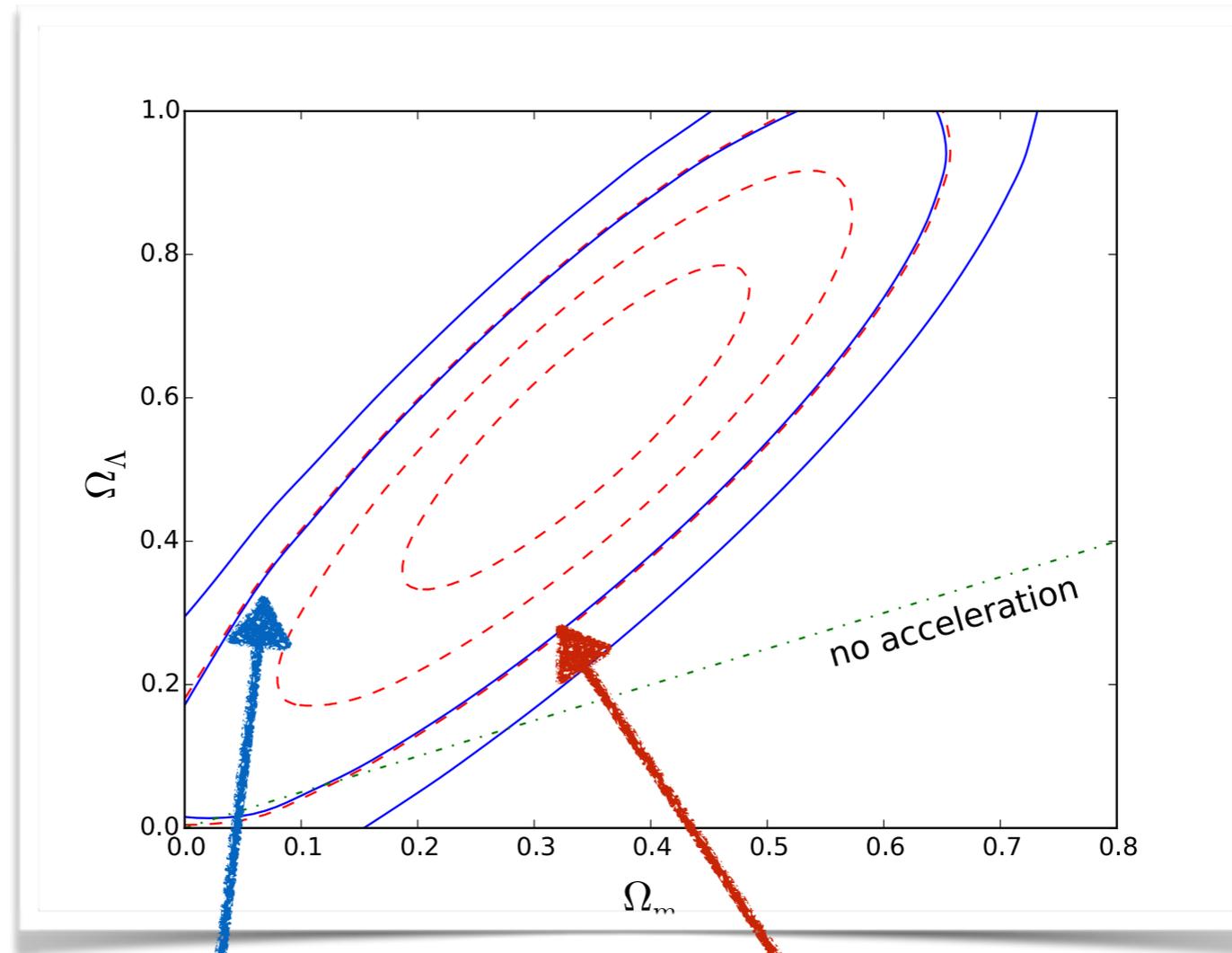
cosmology

SALT2



# MLE fit result

MLE, best fit	
$\Omega_M$	0.341
$\Omega_\Lambda$	0.569
$\alpha$	0.134
$x_0$	0.038
$\sigma_{x_0}^2$	0.931
$\beta$	3.058
$c_0$	-0.016
$\sigma_{c_0}^2$	0.071
$M_0$	-19.05
$\sigma_{M_0}^2$	0.108

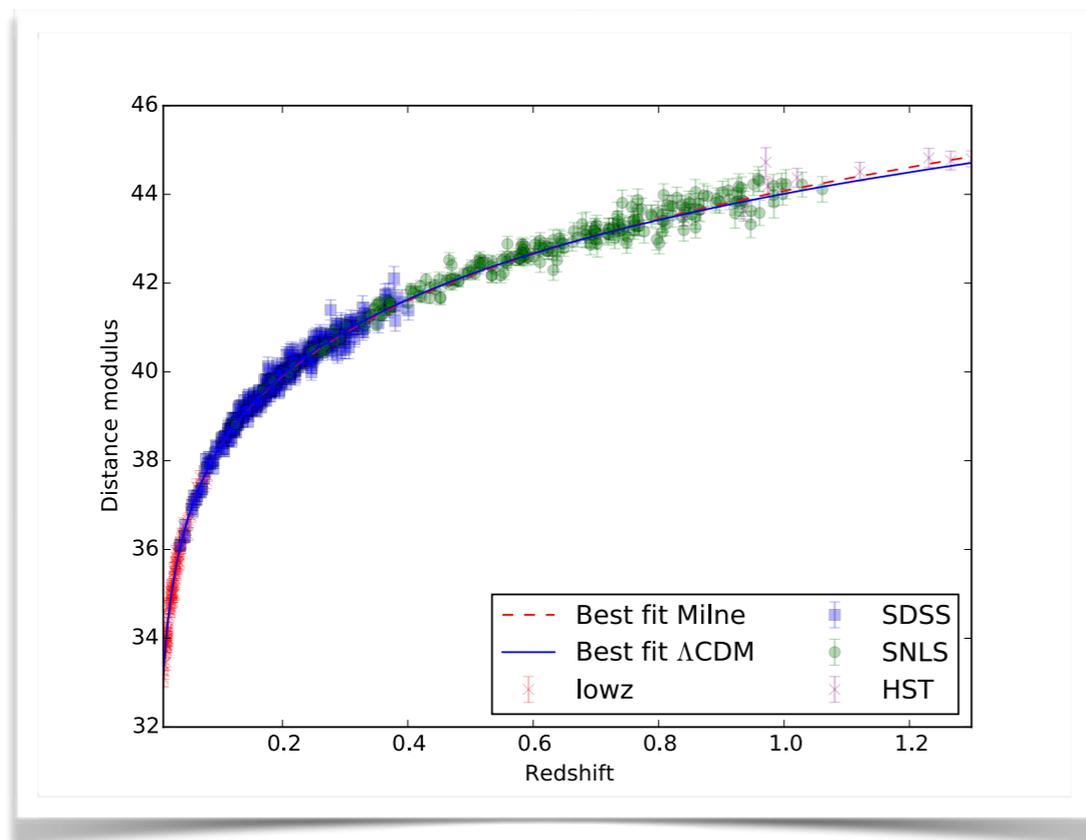


Projected 10D  
confidence regions

2D confidence regions

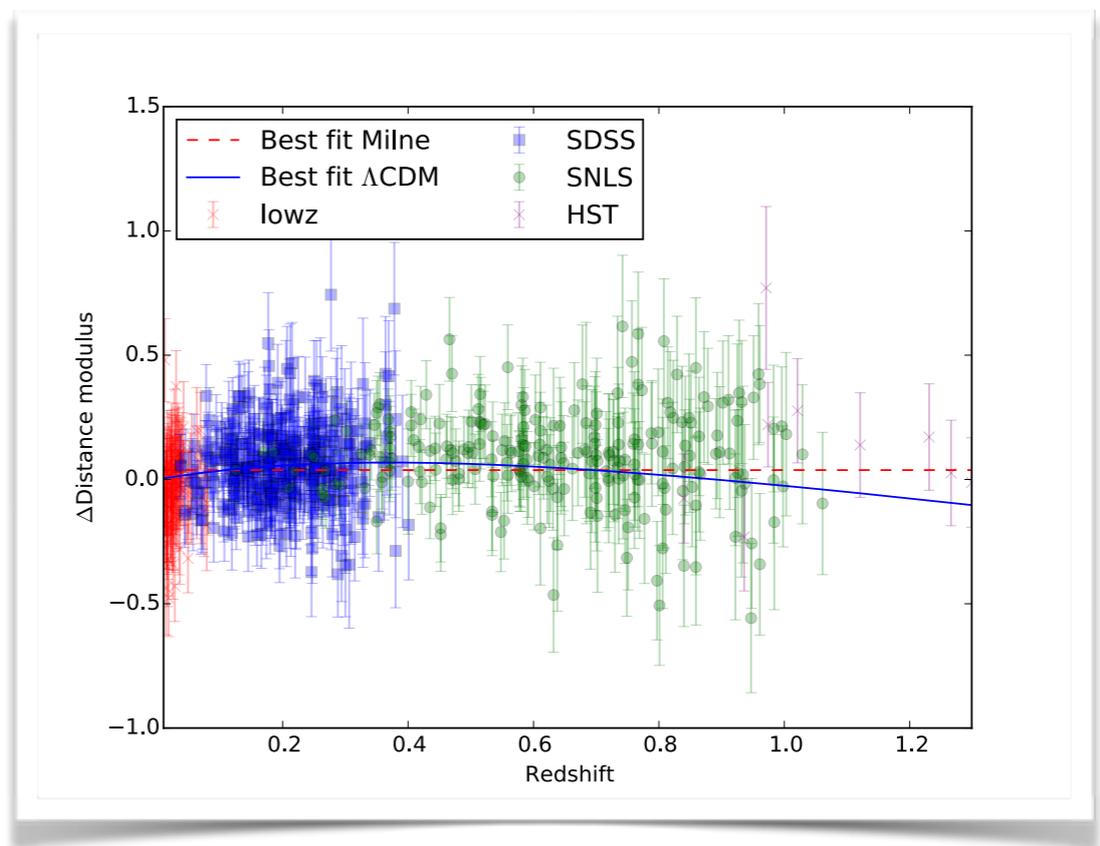


# MLE fit result - does it look like a good fit?



Looks like a good fit when looking at the Hubble diagram

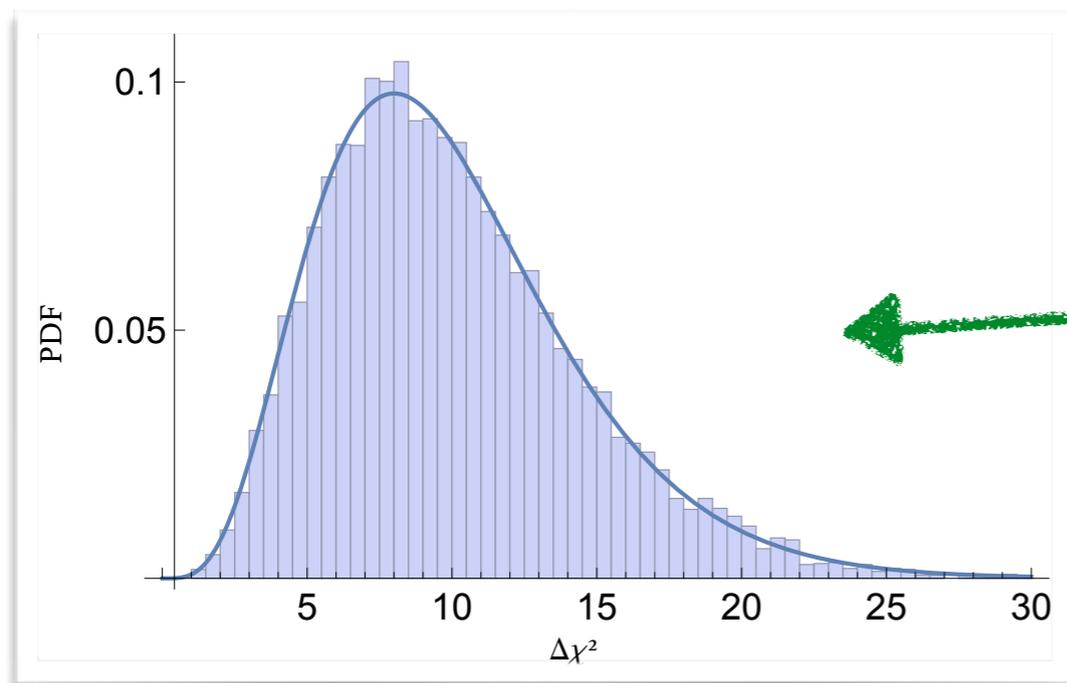
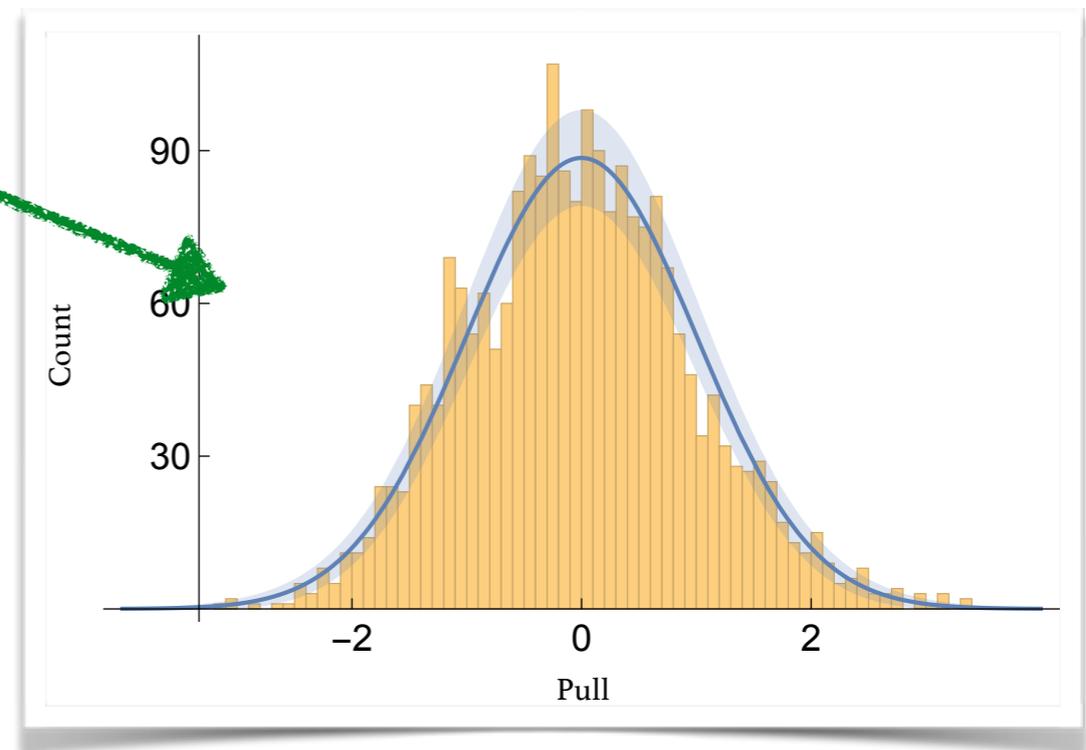
... but in all honesty ...  
any (half-decent) fit would ...



# MLE fit result - Is it a good fit?

## Pulls distribution

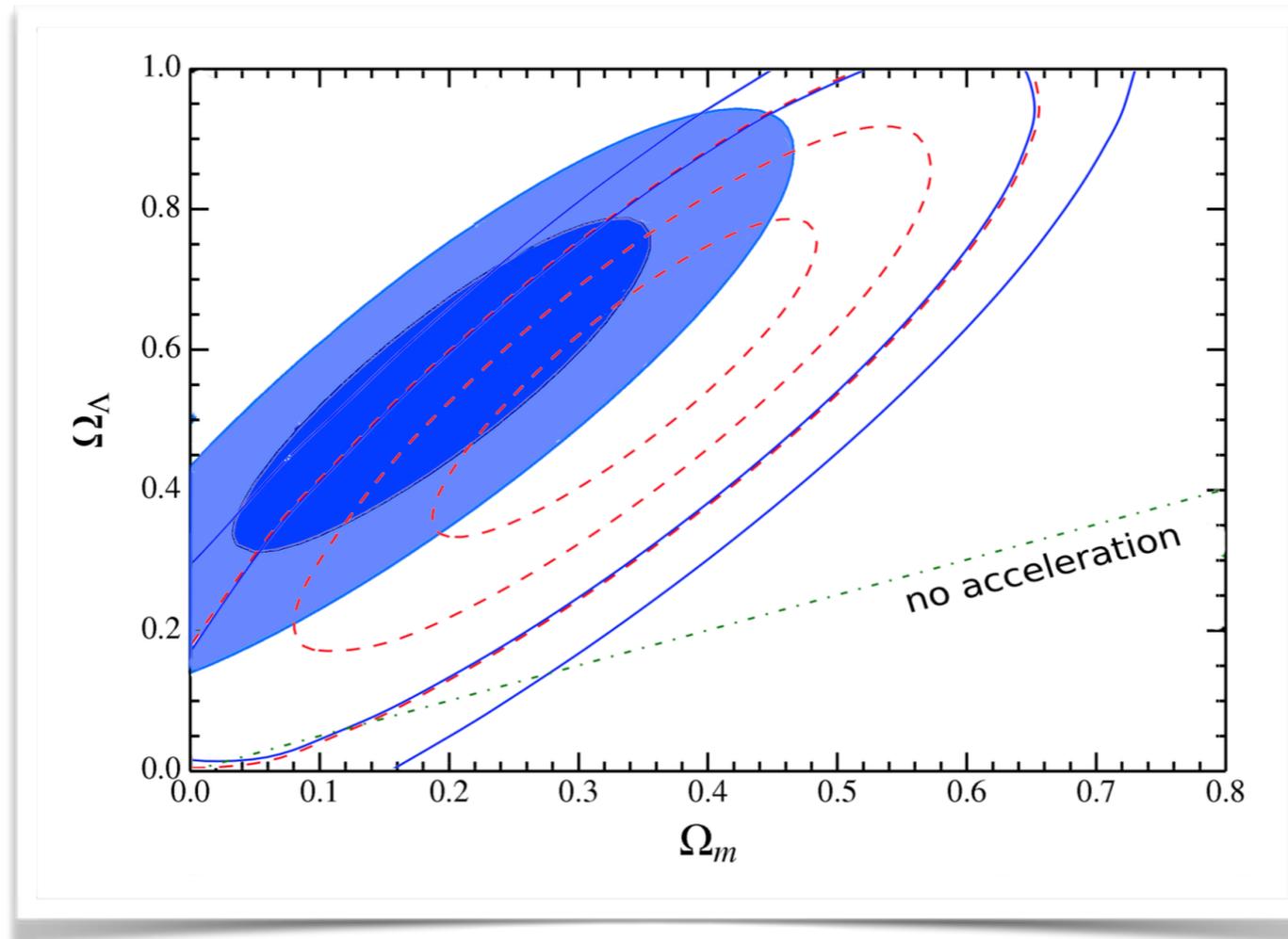
$$(\Sigma_d + A^T \Sigma_l A)^{-1/2} (\hat{Z} - Y_0 A)$$



**Do the best fit and its uncertainties satisfy Wilks' theorem?**



# Comparison to JLA results



Size of “ellipses” comparable to JLA  
Systematic shift towards larger values of  $\Omega_m$



# Recent developments - Unity & BAHAMAS

- More recently two extensive analyses based on Bayesian frameworks have appeared.

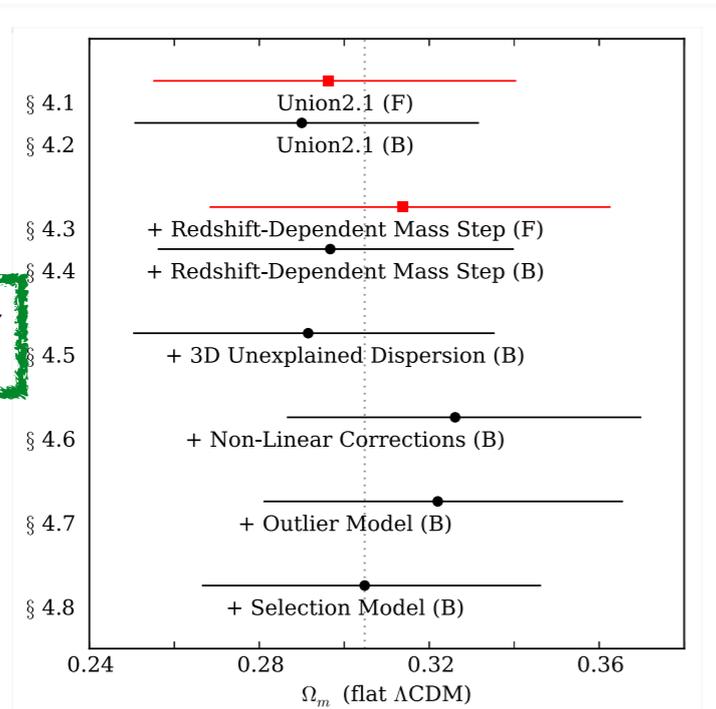
[UNITY, Rubin et al., arXiv:1507.01602]

[BAHAMAS, H. Shariff et al., arXiv:1510.05954]

- The frameworks they develop allow for the systematic study of effects that have been, up to now, disregarded (or poorly treated) in SN cosmology studies, *i.e.* a modelling of the redshift dependence of the color and stretch corrections

BAHAMAS

UNITY



	JLA SNIa only			JLA SNIa + Planck 2015		
	Baseline	$z$ -Linear color Corr	$z$ -Jump color Corr	Baseline	$z$ -Linear color Corr	$z$ -Jump color Corr
<i>Baseline Model parameters</i>						
$\Omega_m$	$0.340 \pm 0.101$	$0.362 \pm 0.094$	$0.429 \pm 0.097$	$0.399 \pm 0.027$	$0.420 \pm 0.031$	$0.425 \pm 0.025$
$\Omega_\Lambda$	$0.542 \pm 0.157$	$0.557 \pm 0.145$	$0.632 \pm 0.155$	$0.625 \pm 0.020$	$0.609 \pm 0.025$	$0.604 \pm 0.019$
$\Omega_\kappa$	$0.119 \pm 0.249$	$0.081 \pm 0.230$	$-0.061 \pm 0.244$	$-0.024 \pm 0.008$	$-0.028 \pm 0.008$	$-0.029 \pm 0.007$
$\alpha$	$0.137 \pm 0.006$	$0.136 \pm 0.006$	$0.136 \pm 0.006$	$0.137 \pm 0.006$	$0.135 \pm 0.007$	$0.136 \pm 0.006$
$\beta$	$3.058 \pm 0.085$	n/a	n/a	$3.068 \pm 0.097$	n/a	n/a
<i>Redshift evolution of color correction parameters</i>						
$\beta_0$	n/a	$3.211 \pm 0.120$	$3.137 \pm 0.092$	n/a	$3.219 \pm 0.119$	$3.136 \pm 0.096$
$\beta_1$	n/a	$-0.622 \pm 0.342$	n/a	n/a	$-0.732 \pm 0.360$	n/a
$\Delta\beta$	n/a	n/a	$-1.120 \pm 0.240$	n/a	n/a	$-1.145 \pm 0.243$
$z_t$	n/a	n/a	$0.662 \pm 0.055$	n/a	n/a	$0.670 \pm 0.056$
<i>Intrinsic magnitude and residual dispersion parameters</i>						
$M_0$	$-19.140 \pm 0.022$	$-19.140 \pm 0.020$	$-19.144 \pm 0.021$	$-19.140 \pm 0.018$	$-19.138 \pm 0.018$	$-19.140 \pm 0.016$
$\sigma_{\text{res}}$	$0.104 \pm 0.005$	$0.104 \pm 0.005$	$0.103 \pm 0.005$	$0.105 \pm 0.005$	$0.105 \pm 0.004$	$0.103 \pm 0.005$



# Conclusions & Outlook

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- Type Ia Supernovae are a crucial probe of cosmic acceleration
- The precision of present and upcoming (*i.e.* DES ...) data needs to be matched by a comparably sound and accurate statistical analysis
- Maximum Likelihood Estimators methods provide such a framework, forcing one to make explicit the assumption made in the modelling of the data
- When MLE techniques are applied to the data from the data in the JLA catalogue marginal evidence ( $\sim 3$  sigma) for cosmic acceleration is supported by SN data alone
- A better modelling of SN standardisation corrections is crucial to improve the quality of the analysis



# *Backup Slides*



# Best fit correlations

## Results

