

Results from the VES experiment

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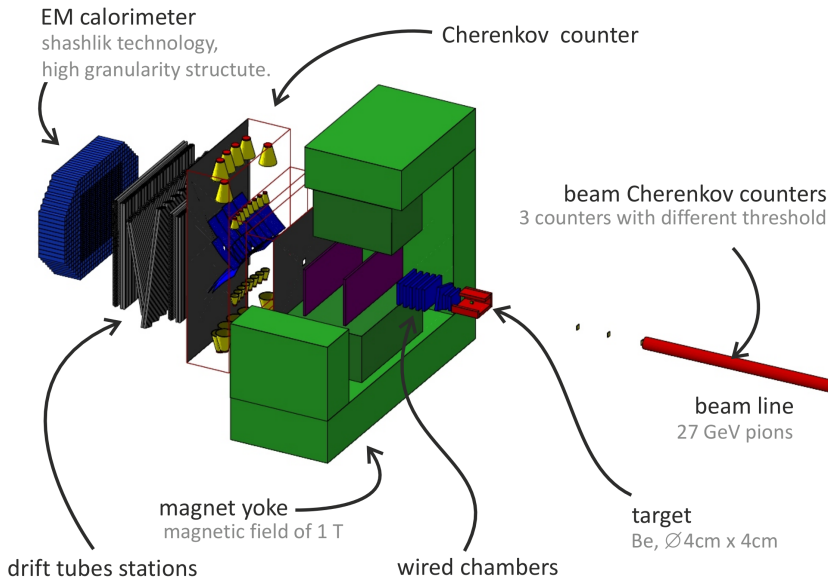
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5 september 2016

- The VES detector
- $\pi^- Be \rightarrow \pi^- \pi^- \pi^+ Be$ and $\pi^- Be \rightarrow \pi^- \pi^0 \pi^0 Be$ at $p_{\pi^-} = 29$ GeV
 - Introduction to 3π analysis
 - Specific features of PWA
 - Isospin relations
 - Rank=1 applicability for diffraction on nucleus targets ?
 - Preliminary Mass-independent PWA
 - Example of preliminary Mass-dependent analysis
 - Prospectives of increased statistics for 3π states
- Summary, outlook



The VES apparatus - scheme



VES apparatus - main features and upgrade

- Secondary beam from U-70 PS: $p_{beam} = 29 \text{ GeV}$ ($\approx 97 \% \pi^-$, $\approx 1.7 \% K^-$)
- Be target ($\approx 0.1 \lambda$)
- Magnetic spectrometer $p_{kick} = 0.56 \text{ GeV}$
- Tracking system
- EM - calorimetry
- Multicell Cherenkov Counter (28 Cells) - PID for secondary tracks
- Minimum-bias trigger (Beam Killer)

The huge upgrade of VES detector - finished in 2011

- **Newly built EMC:** 1215 small (38x38 mm) + 368 big (76x76 mm) Shashlyk counters
- **New large area trackers:** Drift tubes
- **New FE electronics** for tracking detectors
- **New fast DAQ:** $4 \cdot 10^4 / 9s$ cycle, dead time = $22 \mu s$
- **New trigger** Flexible \rightarrow Single - track events (first time at VES !)
- **New Detector Control System**
- **New Beam Momentum Spectrometer**
- **New off-line software** reconstruction, GEANT4 based MC, developed PWA

The main existing experimental results on 3π - systems:

- Existence and properties of $a_1(1260)$ - meson, parameters of $a_2(1320)$ and $\pi_2(1670)$ - ACCMOR (1980-1981)
- Study of $\pi(1800)$ - hybrid candidate VES, 1991-2005
- Existence and properties of spin-exotic $\pi_1(1600)$ - meson, VES \leftrightarrow E852 (1992-2005), COMPASS (2009)
- First analysis of $\pi^-\pi^-\pi^+$ and $\pi^-\pi^0\pi^0$ with same technique E852 (2005)
- Evidence for new narrow state $a_1(1420)$ in 3π COMPASS (2013-2015) and $\pi^-\pi^0\pi^0$ VES (2013)
- Chiral dynamics and radiative decays $a_2(1320)$ and $\pi_2(1670) \rightarrow \pi\gamma$ COMPASS (2011-2014)
- Analysis of 3π with unlimited rank of spin-density matrix and $\pi^-\pi^0\pi^0$ using rank=1 VES (2013-2016)

Still high-disputed issues:

- Spin-exotic $\pi_1(1600) \rightarrow 3\pi$
- Further study of $a_1(1420)$ (tail of rescattered $a_1(1260) \rightarrow K\bar{K}\pi$?)
- Excited states $\pi_2(1880)$, $\pi_2(2100)$, $a_1(1900)$, $a_2(1700)$

Possible model-dependence of the results. What model is correct ?

Example: unlimited rank of spin-density matrix model vs. rank=1 model.

Brief introduction to mass-independent PWA

Mass-independent PWA events density:

$$\mathcal{I}(m, t', \tau) = \sum_{\epsilon=\pm 1} \sum_{r=1}^{N_r} \left| \sum_i T_{i,r}^\epsilon(m, t') \bar{\psi}_i^\epsilon(m, \tau) \right|^2 + FLAT$$

Normalized decay amplitude $\bar{\psi}_i^\epsilon(m, \tau) = \psi_i^\epsilon(m, \tau) / \sqrt{\int |\psi_i^\epsilon(m, \tau')|^2 d\rho(\tau')}$

- $\psi_i^\epsilon(\tau)$ enumerated by its quantum numbers i (coherent) ϵ (incoherent)
- $i, \epsilon = J^{PC} M^\epsilon$ [isobar] πL
- $\psi(\tau)$ has no free parameters \rightarrow this makes fits very quick

Transition amplitudes $T_{ir}^\epsilon(m, t')$ fitted independently in each (m, t') - bin

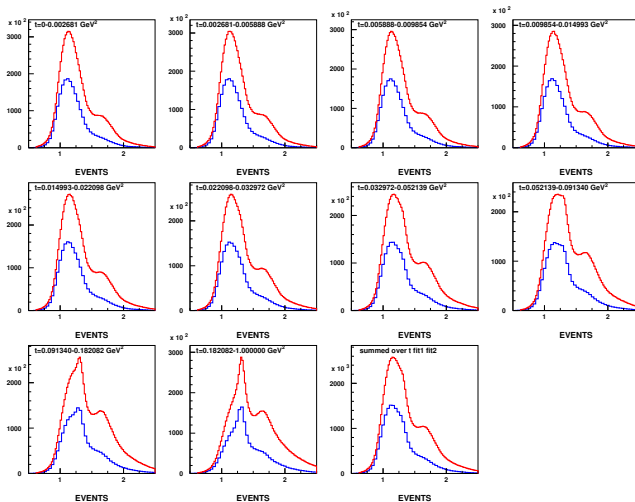
Events density expressed through spin-density matrix $\rho_{i,j}^\epsilon$:

$$\mathcal{I}(m, t', \tau) = \sum_{\epsilon=\pm 1} \sum_{i,j} \rho_{i,j}^\epsilon(m, t') \bar{\psi}_i^\epsilon(m, \tau) \bar{\psi}_j^{\epsilon*}(m, \tau), \quad \rho_{i,j}^\epsilon = \sum_{r=1}^{N_r} T_{i,r}^\epsilon T_{j,r}^{\epsilon*}$$

Diagonal element $\rho_{i,i}^\epsilon$ - number of events associated with a given decay amplitude

Events input to PWA $\pi^- \pi^- \pi^+$ (red) $\pi^- \pi^0 \pi^0$ (blue)

The current data was taken end of 2012



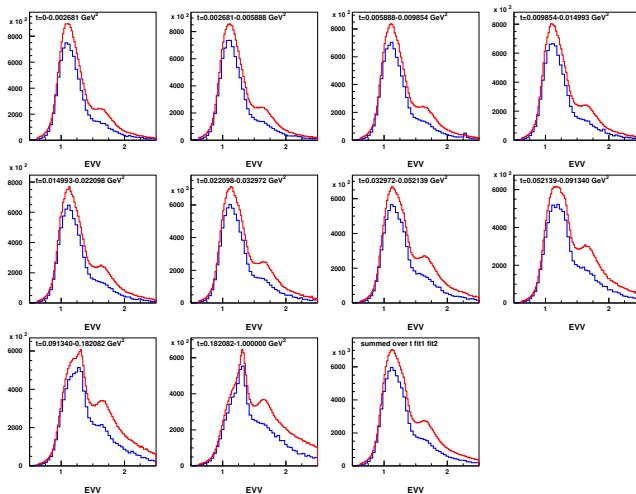
$$N(\pi^- \pi^- \pi^+) = 41\,000\,000$$

$$N(\pi^- \pi^0 \pi^0) = 20\,500\,000$$

$$8\,200\,000 \quad (t' = 0.1 - 1.0 \text{ GeV}^2)$$

$$3\,800\,000 \quad (t' = 0.1 - 1.0 \text{ GeV}^2)$$

Acc-Corrected events $\pi^-\pi^-\pi^+$ (red) $\pi^-\pi^0\pi^0$ (blue)



Isospin relations between $\pi^- \pi^- \pi^+$ and $\pi^- \pi^0 \pi^0$

Decay amplitudes for $\rho\pi$ ($I=1$ of 2π isobar) are connected if isospin $I(3\pi) = 1$:

$$\sqrt{\frac{1}{2}} \left(\frac{1}{\sqrt{2}} [(\pi_{(1)}^- \pi^+) \pi_{(2)}^- + (\pi_{(2)}^- \pi^+) \pi_{(1)}^-] \right) \leftrightarrow -\sqrt{\frac{1}{2}} \left(\frac{1}{\sqrt{2}} [(\pi^- \pi_{(1)}^0) \pi_{(2)}^0 + (\pi^- \pi_{(2)}^0) \pi_{(1)}^0] \right)$$

Decay amplitudes for $f\pi$ ($I=0$ of 2π isobar) channels are always connected :

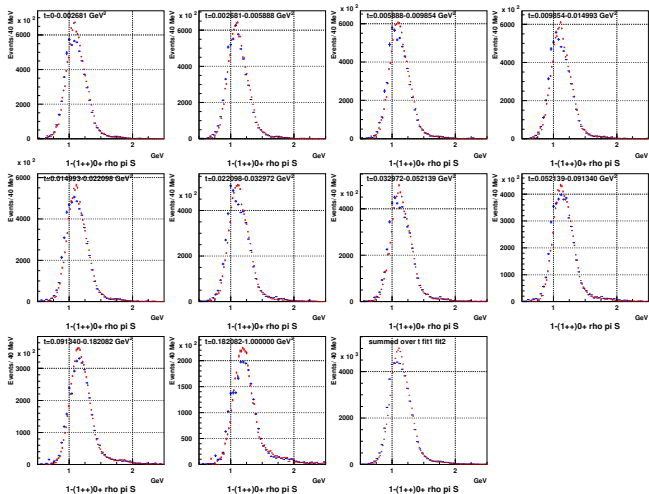
$$\sqrt{\frac{2}{3}} \left(\frac{1}{\sqrt{2}} [(\pi_{(1)}^- \pi^+) \pi_{(2)}^- + (\pi_{(2)}^- \pi^+) \pi_{(1)}^-] \right) \leftrightarrow -\sqrt{\frac{1}{3}} \left((\pi_{(1)}^0 \pi_{(2)}^0) \pi^- \right)$$

Connection means same production amplitudes.

$Br = N(\pi^- \pi^0 \pi^0) / N(\pi^- \pi^- \pi^+)$ is just ratio of integrals of squared decay amplitudes

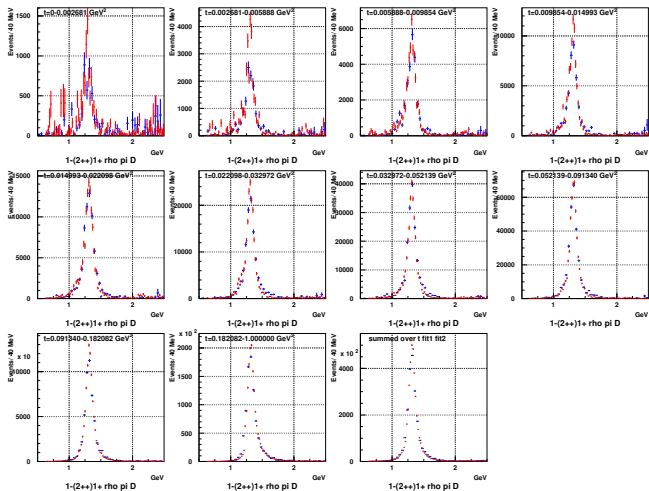
- 1) $\rho\pi$: $Br=1$ as decay amplitudes have same structure up to $m(\pi^{+/-}) \rightarrow m(\pi^0)$.
- 2) $f\pi^-$: Br comes from different isobar decay and bose-symmetrization. Br can differ from 0.5 in case of big cross-term interference on $\pi^- \pi^- \pi^+$ Dalitz-plot.
- 3) All corresponding relative phases in $\pi^- \pi^0 \pi^0$ and $\pi^- \pi^- \pi^+$ are equal - in case of appropriate choosing of directions of "spin analyzers" (π^- - direction in 2π CMS for both systems)

$1^{++}0^+ \rho(770)\pi S$ in $\pi^- \pi^- \pi^+$ (red) $\pi^- \pi^0 \pi^0$ (blue)



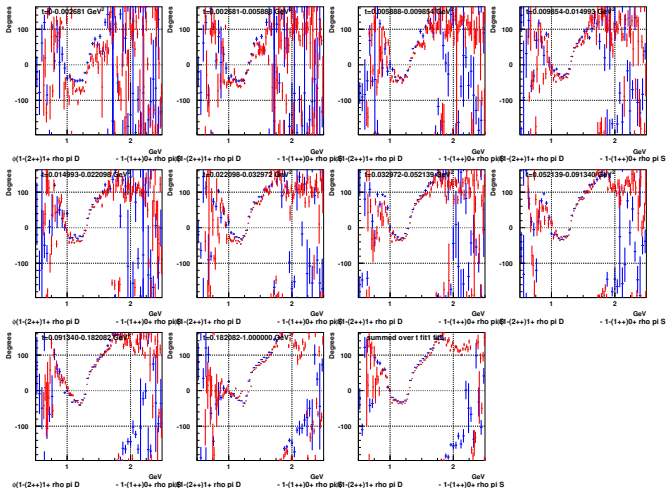
Isospin relations for intensity - approx. satisfied in each (m, t) - bin

$2^{++}1^+\rho(770)\pi D$ in $\pi^-\pi^-\pi^+$ (red) $\pi^-\pi^0\pi^0$ (blue)



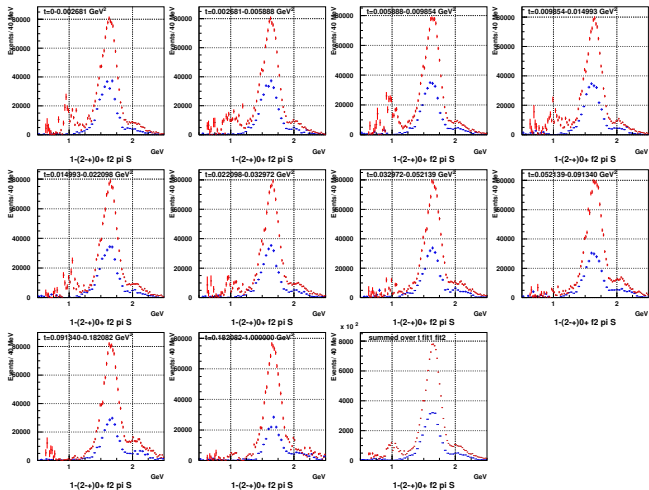
Isospin relations for intensity - approx. satisfied in each (m, t) - bin

Phase of $2^{++}1^+\rho(770)\pi D$ $\pi^-\pi^-\pi^+$ (red) $\pi^-\pi^0\pi^0$ (blue)



Isospin relations for relative phase - approx. satisfied in each (m, t) - bin

$2^{-+}0^{+}f_2\pi S$ in $\pi^{-}\pi^{-}\pi^{+}$ (red) $\pi^{-}\pi^0\pi^0$ (blue)



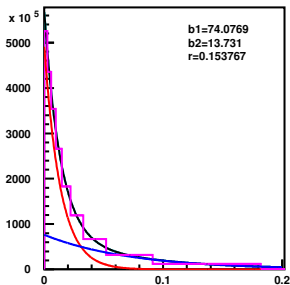
Isospin relations for intensity - approx. satisfied in each (m, t) - bin

$$\frac{dN}{dt} \text{ for } 1^{++}0^+ \rho\pi S, 2^{++}1^+ \rho\pi D, 2^{-+}0^+ f_2\pi S$$

t' -dependence of partial-wave intensities, integrated in broad $m(\pi^-\pi^-\pi^+)$ regions

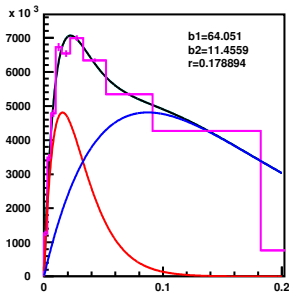
Fit function $\frac{dN}{dt} = x t'^M (e^{-b_c t'} + r e^{-b_i t'})$

m=1.2-1.4



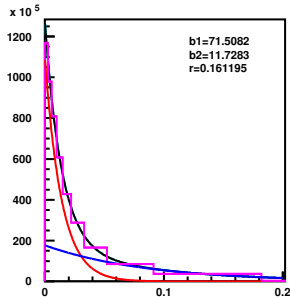
1-(1⁺⁺)0⁺ rho pi S

m=1.2-1.4



1-(2⁺⁺)1⁺ rho pi D

m=1.6-1.8



1-(2⁻⁺)0⁺ f₂ pi S

Toy formula for adding coherent+incoherent processes

Adding coherent+incoherent density matrices:

- example : 2x2 matrix; $M=0$ and $M=1$ states
- phase is assumed to be the same for 2 processes
- if b_c, b_i are universal and $r_1 = r_2 \rightarrow Coh = 1$
- in general, different t^M factors are cancelled

$$\begin{pmatrix} e^{-b_c t} & \sqrt{r} e^{-b_c t} \\ t e^{-b_c t} & \end{pmatrix} + \begin{pmatrix} r_1 e^{-b_i t} & \sqrt{r_1 r_2} \sqrt{t} e^{-b_i t} \\ t \cdot r_2 e^{-b_i t} & \end{pmatrix}$$

$$Coh = \frac{\sqrt{r} (e^{-b_c t} + \sqrt{r_1 r_2} e^{-b_i t})}{\sqrt{(e^{-b_c t} + r_1 e^{-b_i t})(e^{-b_c t} + r_2 e^{-b_i t})}}$$

Brief introduction to mass-dependent PWA

Mass-independent PWA events density:

$$\mathcal{I}(m, t', \tau) = \sum_{\epsilon=\pm 1} \sum_{r=1}^{N_r} \left| \sum_i T_{i,r}^{\epsilon}(m, t') \bar{\psi}_i^{\epsilon}(m, \tau) \right|^2$$

The model for global $m(3\pi)$ - dependence

$$\mathcal{I}(m, t', \tau) = \sum_{\epsilon} \sum_r \left| \sum_i \sum_k C_{ikr}^{\epsilon} BW_{ik}(m, t', \zeta) \sqrt{\int |\psi_i^{\epsilon}(\tau', m)|^2 d\rho(\tau')} \bar{\psi}_i^{\epsilon}(\tau, m) \right|^2$$

Comparing mass-indep and global mass-dep :

$$\rho_{i,j}^{\epsilon}(m, t') = \sum_r \sum_{k,l} C_{ikr}^{\epsilon} C_{jlr}^{\epsilon*} BW_{ik}(m, t', \zeta) BW_{jl}^*(m, t', \zeta) \sqrt{\int |\psi_i^{\epsilon}(m, \tau')|^2 d\rho(\tau')} \sqrt{\int |\psi_j^{\epsilon}(m, \tau')|^2 d\rho(\tau')}$$

Mass-independent intensities normalized to number of observed events \rightarrow
mass-dependent intensities contain opening phase-space factors

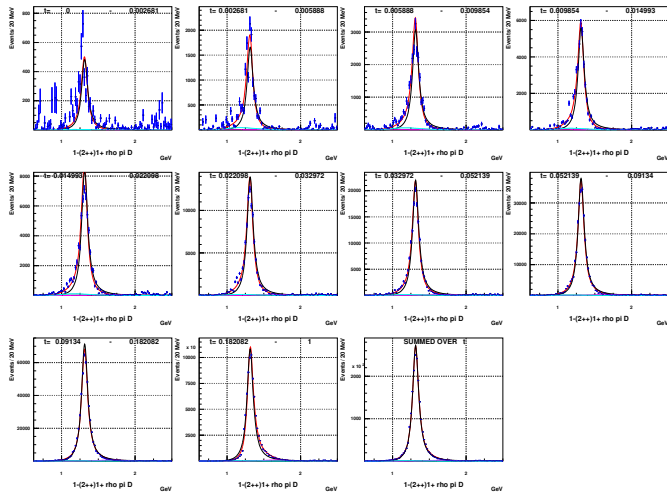
The current mass-dependent fit

The sub-set of partial waves fitted mass-dependently:

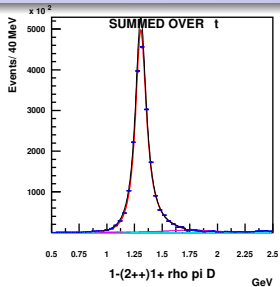
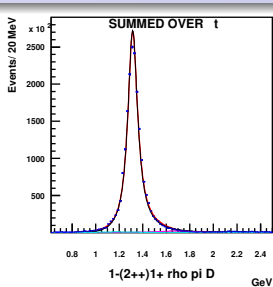
- $1^{++}0^+ \rho(770)\pi S$
- $2^{++}1^+ \rho(770)\pi D$
- $2^{++}1^+ f_2(1270)\pi P$
- $0^{-+}0^+ f_0(980)\pi S$
- $1^{++}0^+ f_0(980)\pi P$
- $2^{-+}0^+ f_2(1270)\pi S$
- $4^{++}1^+ \rho(770)\pi G$
- $4^{++}1^+ f_2(1270)\pi F$

The following mass-dependent fit curves and following numbers are NOT to be considered as final results !

$2^{++}1^+\rho^0(770)\pi^-D$ m-dep fit in 10 t' bins for $\pi^-\pi^-\pi^+$

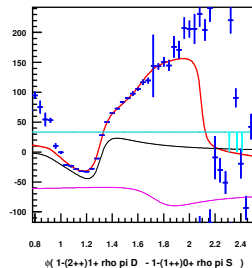


$2^{++}1^+\rho(770)\pi D$ intens and phase (m) - combined t



$$\pi^- \pi^- \pi^+$$

$$M = 1.319 \quad \Gamma = 0.103$$

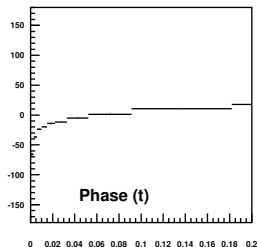
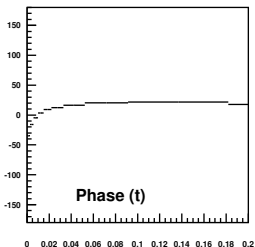
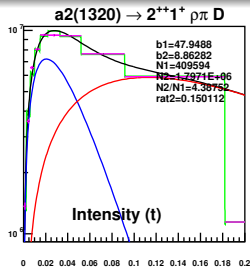
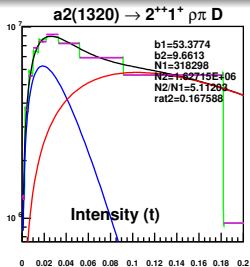


$$\pi^- \pi^0 \pi^0$$

$$M = 1.315 \quad \Gamma = 0.124$$

$a_2(1320)$

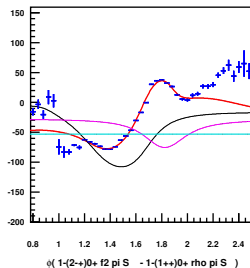
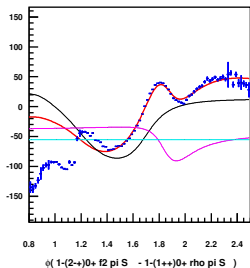
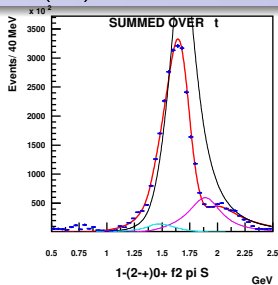
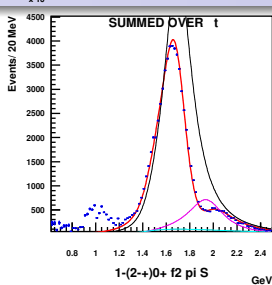
$a_2(1320)$ intens and prod. phase (t)



$$\pi^- \pi^- \pi^+$$

$$\pi^- \pi^0 \pi^0$$

$2^{-+}0^{+}f_2\pi^{-}S$ intens and phase (m)

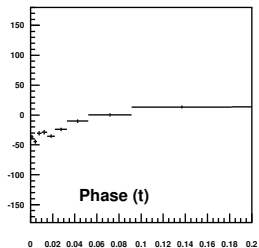
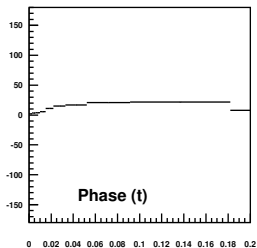
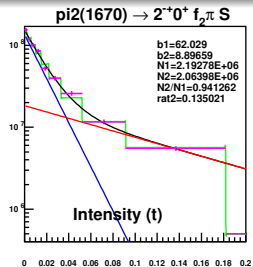
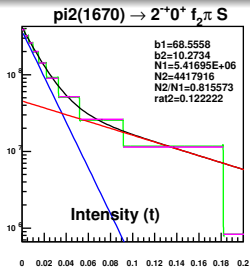


$\pi_2(1670)$
 $\pi_2(1880)$

$M = 1.690 \quad \Gamma = 0.308$
 $M = 1.940 \quad \Gamma = 0.380$

$M = 1.678 \quad \Gamma = 0.319$
 $M = 1.890 \quad \Gamma = 0.370$

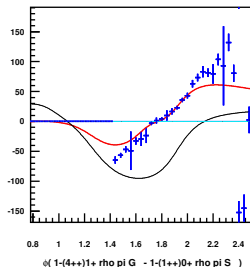
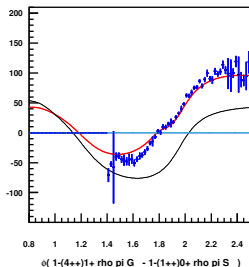
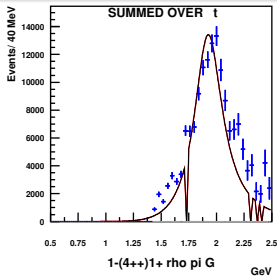
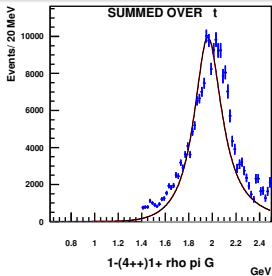
$\pi_2(1670)$ intens and prod-phase (t)



$\pi^- \pi^- \pi^+$

$\pi^- \pi^0 \pi^0$

$4^{++}1^+\rho\pi G$ intens and phase (m)



$a_4(2050)$

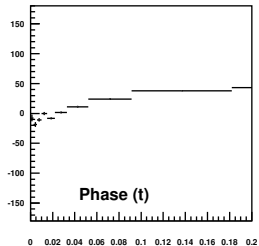
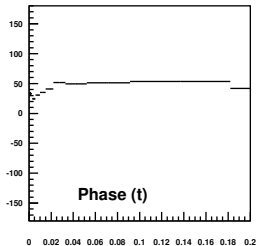
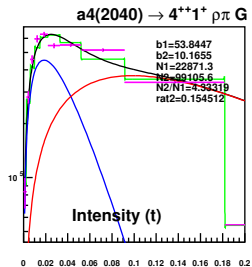
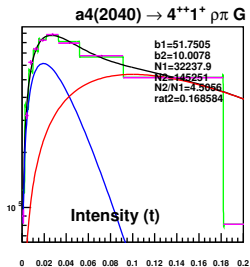
$$\pi^- \pi^- \pi^+$$

$$M = 1.950 \quad \Gamma = 0.294$$

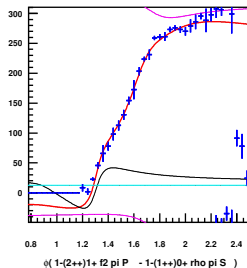
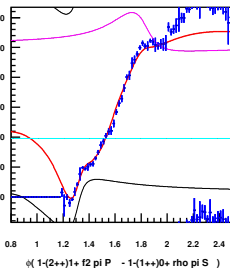
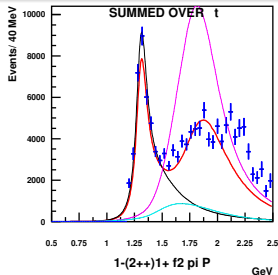
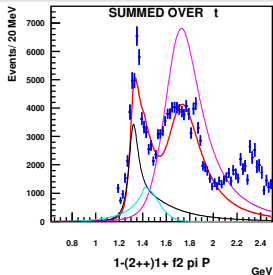
$$\pi^- \pi^0 \pi^0$$

$$M = 1.910 \quad \Gamma = 0.317$$

$a_4(2050)$ intens and phase (t)



$2^{++}1^+ f_2 \pi^- P$ intens and phase (m)

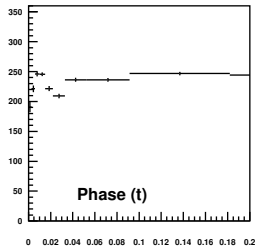
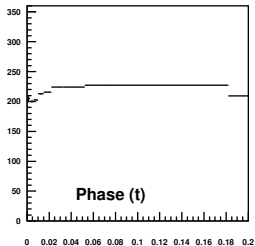
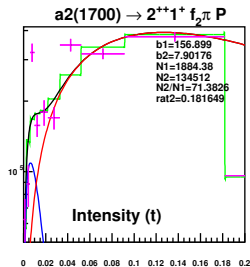
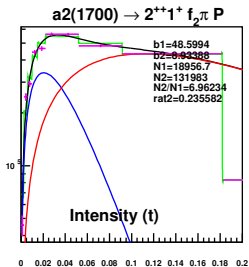


$a_2(1700)$

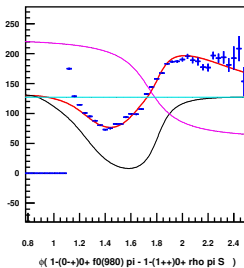
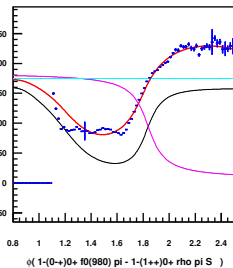
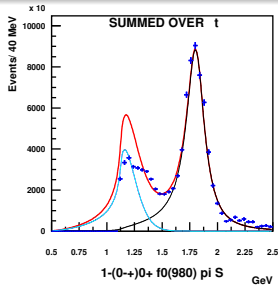
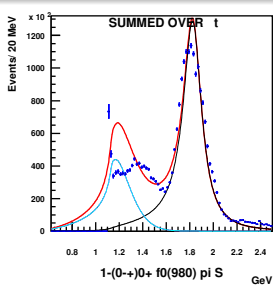
$\pi^- \pi^- \pi^+$
 $M = 1.690 \quad \Gamma = 0.430$

$\pi^- \pi^0 \pi^0$
 $M = 1.780 \quad \Gamma = 0.550$

$a_2(1700)$ intens and phase (t)



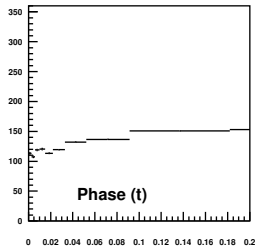
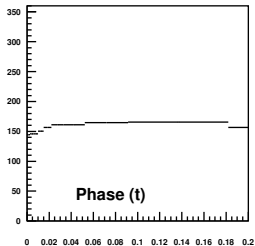
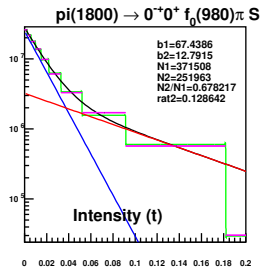
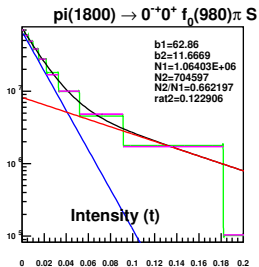
$0^{-+}0^{+}f_0(980)\pi S$ intens and phase (m)



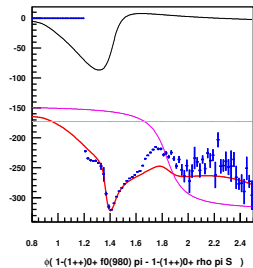
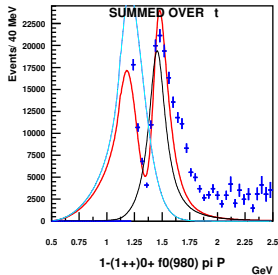
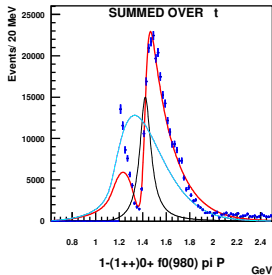
$$M = 1.820 \quad \Gamma = 0.191$$

$$M = 1.800 \quad \Gamma = 0.199$$

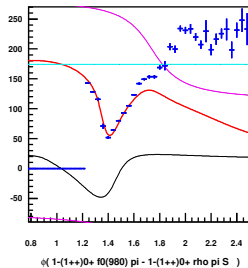
$\pi(1800)$ intens and phase (t)



$1^{++}0^+ f_0(980) \pi P$ intens and phase (m)

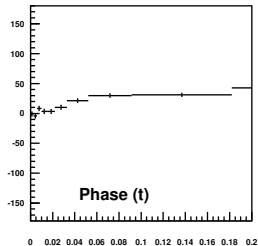
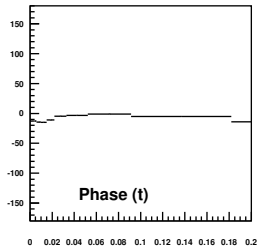
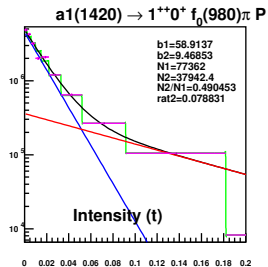
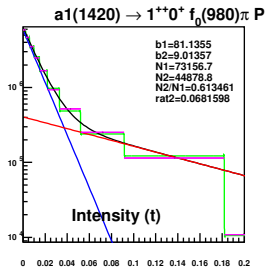


$$M = 1.419 \quad \Gamma = 0.115$$



$$M = 1.450 \quad \Gamma = 0.188$$

$a_1(1420)$ intens and phase (t)



- The mass independent PWA, using rank=1 fit model and 10 t' intervals is first applied for both $\pi^-\pi^-\pi^+$ and $\pi^-\pi^0\pi^0$ systems, diffractively produced on Be nucleus
- The preliminary mass-dependent PWA results are first compared between $\pi^-\pi^-\pi^+$ and $\pi^-\pi^0\pi^0$ systems

- Further development of the mass independent PWA
 - Comparison of rank=1 vs. unlimited rank density-matrix model
 - Partial coherence model ?
 - Free-isobarred analysis for both $\pi^- \pi^- \pi^+$ and $\pi^- \pi^0 \pi^0$
- Further development of the mass-dependent PWA
 - Choosing of the best mass-dependent wave sub-set
 - Choosing of resonant and background parametrization (t' -dependent background shape ?)
 - Full study of local minima and spurious solutions in the mass-dep fit