Neutrinos Theory and Phenomenology

Joachim Kopp EPS Conference on High Energy Physics | Venice, Italy | July 11th, 2017









Apologies for not being here in person

The reason are these little fellows, born yesterday









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In this Talk

- Phenomenoloy for Long Baseline Experiments
- Oscillation Anomalies and Sterile Neutrinos
- **Meutrinos and Dark Matter**
- Neutrinoless Double Beta Decay
- Understanding Neutrino Mass and Mixing







- Astrophysical Neutrinos

 talk by Maarten de Jong
- Supernova Neutrinos
 talk by Maarten de Jong
- Meutrino Cosmology
 - → talks by François Bouchet, Antonio Riotto
- Neutrinos as a probe for Dark Matter
 talks by Yonit Hochberg, Maarten de Jong







Phenomenology for Long Baseline Experiments























LHC



- Matrix Elements
 - o NⁿLO calculation
 - **o** MC simulations (Pythia, ...)
- Model Building
- 🗹 Global Fits

Neutrinos

- Beam FluxesCross Sections
- Oscillation Physics (SM + beyond)
 - **o** Simulations (GENIE, GLoBES, ...)
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Optimization of Future Experiments

- **o** What are the most important physics goals?
- SM precision measurements ↔ BSM searches
- **o** Boundary conditions (available sites, budget, politics, ...)

Mow could BSM physics affect precision measurements?







Long Baseline Phenomenology









- **M** Fits to oscillation data fraught with correlations and degeneracies
- Can often be resolved in global fits

NuFit 3.0 (Esteban Gonzalez-Garcia Maltoni Martinez Schwetz), 2016







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Global Fits



Global Fits

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Fits to oscillation data fraught with correlations and degeneracies

Can often be resolved in global fits

★ O(few %) precision
 ★ no sensitivity to mass ordering, δ_{CP}, octant of θ₂₃
 ★ ... yet!
 ★ First hints will probably emerge from global fits

NuFit 3.0 (Esteban Gonzalez-Garcia Maltoni Martinez Schwetz), 2016







Oscillation Anomalies





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Anomalies in Short Baseline Oscillations







\mathbf{V} LSND / MiniBooNE: anomalous $\nu_{\mu} \rightarrow \nu_{e}$ oscillations











 \mathbf{V} LSND / MiniBooNE: anomalous $\nu_{\mu} \rightarrow \nu_{e}$ oscillations







 $\mathbf{V} \sqcup \mathbf{LSND} / \mathbf{MiniBooNE}$: anomalous $\nu_{\mu} \to \nu_{e}$ oscillations

 \mathbf{M} Reactor & Gallium Experiments: anomalous ν_e disappearance









 \mathbf{V} LSND / MiniBooNE: anomalous $\nu_{\mu} \rightarrow \nu_{e}$ oscillations

 \mathbf{V} Reactor & Gallium Experiments: anomalous ν_e disappearance







"I felt a great disturbance in the force."

....

 \checkmark Promote mixing matrix to 4×4

Solution channels are related:

$$\begin{aligned} P_{\nu_e \to \nu_e} \simeq 1 - 2|U_{e4}|^2 (1 - |U_{e4}|^2) \\ P_{\nu_\mu \to \nu_\mu} \simeq 1 - 2|U_{\mu4}|^2 (1 - |U_{\mu4}|^2) \\ P_{\nu_\mu \to \nu_e} \simeq 2|U_{e4}|^2 |U_{\mu4}|^2 \end{aligned}$$
(for $4\pi E/\Delta m_{41}^2 \ll L \ll 4\pi E/\Delta m_{31}^2$)

Models can be over-constrained.







Global Fit in 3+1 Model



Dentler Hernandez JK Machado Maltoni Martinez Schwetz, in preparation see also works by Collin Argüelles Conrad Shaevitz, <u>1607.00011</u>, Gariazzo Giunti Laveder Li, <u>1703.00860</u>







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severe tension ($p < 10^{-4}$)

 \star scrutinize anomalies for unknown systematics (need 4 independent effects!)

 \star scrutinize also null results.



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Standard picture: ν_s production via oscillation at T \ge MeV Constrained by N_{eff} and Σm_v







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- Entropy production at T < MeV</p>
 - ν_s diluted Fuller Kishimoto Kusenko, <u>1110.6479</u>; Ho Scherrer, <u>1212.1689</u>







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 - o production suppressed by thermal potential

Hannestad et al. <u>1310.5926</u>; Dasgupta JK, <u>1310.6337</u>







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 \mathbf{V}_s properties change in late phase transition

Bezrukov Chudaykin Gorbunov, <u>1705.02184</u>







Flux Measurement by Daya Bay









Reactor fuel composition evolves with time ("burnup")









Reactor fuel composition evolves with time ("burnup")

Measure inverse β decay rate per isotope







 \checkmark

Reactor fuel composition evolves with time ("burnup") Measure inverse β decay rate *per isotope*






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- Measure inverse β decay rate per isotope
- **Markov Full analysis:**
 - O Compare fit with free ²³⁵U, ²³⁸U, ²³⁹Pu, ²⁴¹Pu fluxes to fit with fixed fluxes + $\sin^2 2\theta_{14}$

$$\Delta \chi^2 = 7.9$$







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- o Include uncertainties in fixed fluxes?







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Fluxes within errors + $\sin^2 2\theta_{14}$, Δm_{41}^2 : p = 0.18 Fluxes free : p = 0.73 $\Delta \chi^2$ (sterile neutrino vs. free fluxes) : p = 0.006















Neutrinos and Dark Matter





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Leading candidate for Warm Dark Matter

- Improved small scale structure
- o x-ray line signature
- Production through oscillations challenged by e.g. Lyman-α data
- Promising alternative production mechanisms
 - o Decays of heavy particles
 - o High-T freeze-in









DM Annihilation in the Sun









DM Annihilation in the Sun









Coherent forward scattering of neutrinos on DM

- o analogous to SM matter effects ("MSW effect")
- **o** Requires huge DM number density
- Fuzzy Dark Matter
 - **o** scalar or vector, $m < 10^{-20} \, \mathrm{eV}$
 - o Compton wave length $\sim \mathrm{pc}$
 - o Interesting for small scale structure

Krnjaic Machado Necib, <u>1705.06740</u> Brdar JK Liu Prass Wang, <u>1705.09455</u>







Neutrino – DM Interactions





o an Limits from Long-Baseline Experiments Re





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"You underestimate the power of the dark side."

Matrix Elements for Neutrinoless Double Beta Decay





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 $\Gamma \propto G_F^4 |\mathcal{M}|^2 \left| \sum U_{ej}^2 m_j \right|^2 p_e^2$









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Nuclear Matrix Elements cannot be measured independently Nuclear Shell Model

- o good for small nuclei
- o fails for heavy nuclei, does not capture multiparticle excitation
- Quasiparticle Random Phase Approximation (QRPA)
 - **o** quasiparticle state = linear combination of single nucleon states
 - o large set of states
 - o but less accurate description of nucleon nucleon correlations
- \mathbf{V} Currently $\mathcal{O}(1)$ uncertainties

of Future:

- o Improvements to Shell Model, QRPA
- **ο** Ab Initio method (e.g. in χPT)





Engel Menéndez, <u>1610.06548</u>



Nuclear Theory for 0v2β Decay



Understanding Neutrino Mass and Mixing





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Seesaw Models

Right-handed singlet: (type-l seesaw)



Scalar triplet: (type-II seesaw)



Fermion triplet: (type-III seesaw)



$$m_{\nu} = Y_{\Sigma}^T \frac{1}{M_{\Sigma}} Y_{\Sigma} v^2$$

Minkowski; Gellman, Ramon, Slansky; Yanagida; Glashow; Mohapatra, Senjanovic

 $m_{\nu} = Y_N^T \frac{1}{M_N} Y_N v^2$

Magg, Wetterich; Lazarides, Shafi; Mohapatra, Senjanovic; Schechter, Valle

Foot, Lew, He, Joshi; Ma; Ma, Roy; T.H., Lin, Notari, Papucci, Strumia; Bajc, Nemevsek, Senjanovic; Dorsner, Fileviez-Perez;....

slide by Thomas Hambye







- Generic seesaw scale (10¹⁴ GeV) out of reach
- Opportunities in low-scale seesaw, LR symmetry, B-L models.





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Radiative Neutrino Mass Models



diagrams from Sugiyama, <u>1505.01738</u>







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Radiative Neutrino Mass Models



diagrams from Sugiyama, 1505.01738







Radiative Neutrino Mass Models





Mased on symmetries

- o continuous ↔ discrete
- o Abelian ↔ non-Abelian
- o with/without CP symmetry
- Can reproduce observed masses / mixing angles
- Predictive power in specific models

Hagedorn, <u>1705.00684</u>







Summary









Summary

- More that the second second
- Meed to get to the bottom of oscillation anomalies
- **Mathematical Content of Series and DM physics closely intertwined**
- Ov2β decay needs accurate matrix elements
- **Meutrino mass & flavor models need experimental input**







"Always in motion is the future"

Thank you!





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