

LBL experiments: a new window on light sterile neutrinos



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Outline

Introduction

Interference effects mediated by sterile neutrinos

LBL constraints on sterile ν_s : present

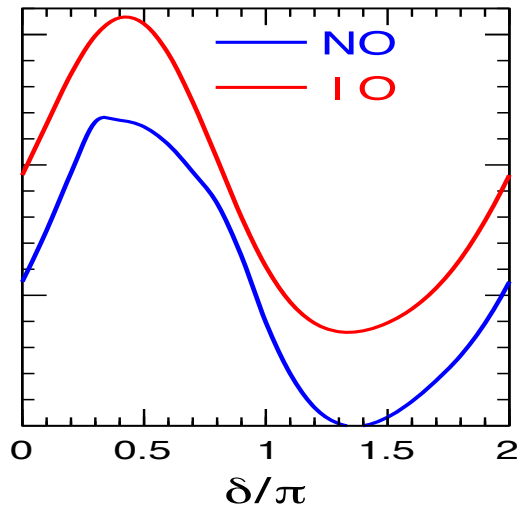
LBL constraints on sterile ν_s : future

Conclusions

Introduction



It is timely to pose a new question



Capozzi, Di Valentino, Lisi,
Marrone, Melchiorri, A.P,
PRD 95, 096014 (2017)

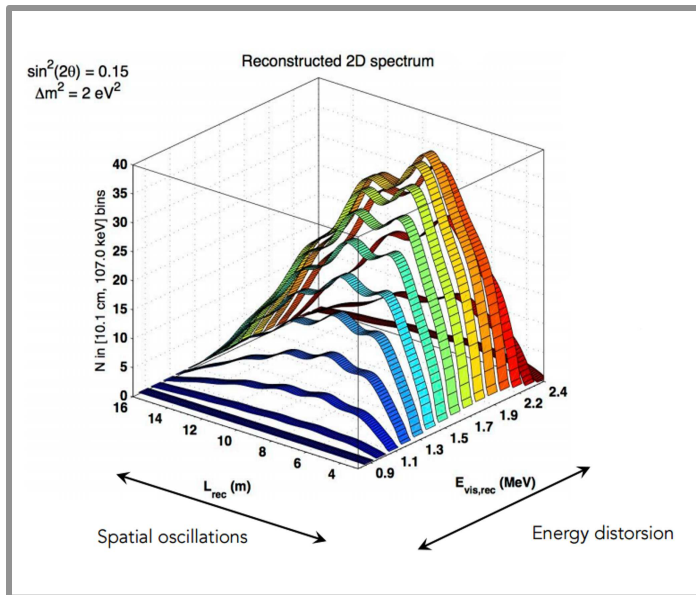
**LBL experiments start
to be sensitive to the
CP violating phase δ**

**Can sterile neutrinos generate new observable CP
violating effects at LBL?**

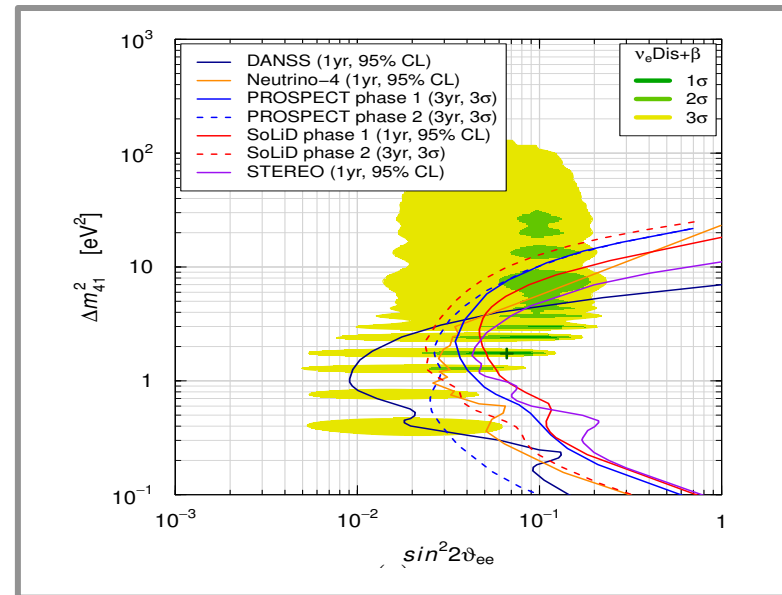
Question basically ignored in the past

A discovery can come only at SBL experiments

SOX experiment @ LNGS



Gariazzo et al., 1703.00860



Observing the oscillation pattern (in energy and/or space)

However ...

SBL have an intrinsic limitation

At SBL atm/sol oscillations are negligible

$$\frac{L}{E} \sim \frac{m}{\text{MeV}}$$

$$\begin{aligned} \Delta_{12} &\simeq 0 \\ \Delta_{13} &\simeq 0 \end{aligned}$$

$$\Delta_{ij} = \frac{\Delta m_{ij}^2 L}{4E}$$

Impossible to observe phenomena of interference between the new frequency ($\Delta_{14} \sim 1$) and atm/sol ones

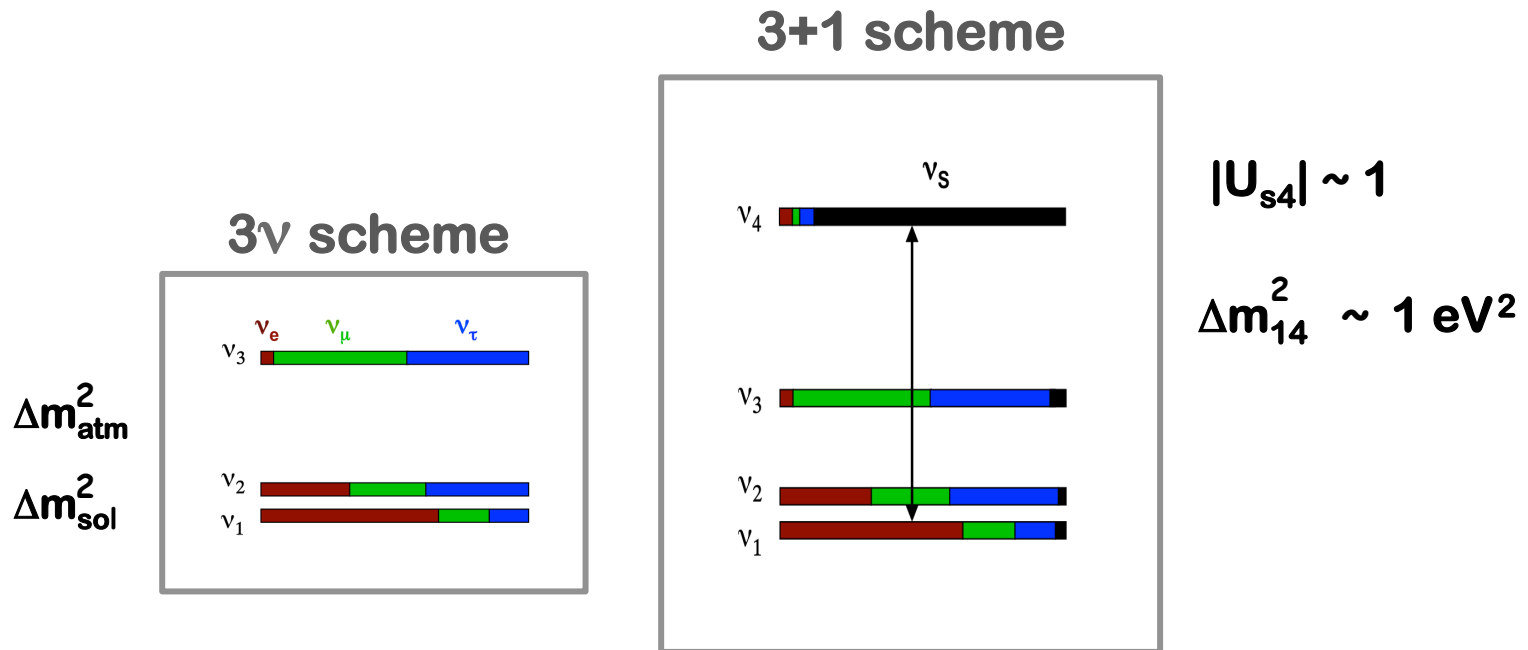
But we have LBL, which are sensitive interferometers



Interference effects mediated by sterile ν s

N. Klop & A.P., PRD 91 073017 (2015)
arXiv: 1412.7524

How to enlarge the 3-flavor scheme



At LBL the effective 2-flavor SBL description is no more valid and calculations should be done in the 3+1 (or 3+N_s) scheme

Mixing Matrix in the 3+1 scheme

$$U = \tilde{R}_{34} R_{24} \tilde{R}_{14} R_{23} \underbrace{\tilde{R}_{13} R_{12}}_{3\nu}$$

$$R_{ij} = \begin{bmatrix} c_{ij} & s_{ij} \\ -s_{ij} & c_{ij} \end{bmatrix}$$

$$\tilde{R}_{ij} = \begin{bmatrix} c_{ij} & \tilde{s}_{ij} \\ -\tilde{s}_{ij}^* & c_{ij} \end{bmatrix}$$

$$\begin{aligned} s_{ij} &= \sin \theta_{ij} \\ c_{ij} &= \cos \theta_{ij} \\ \tilde{s}_{ij} &= s_{ij} e^{-i\delta_{ij}} \end{aligned}$$

$$3\nu \begin{cases} 3 \text{ mixing angles} \\ 1 \text{ Dirac phase} \\ 2 \text{ Majorana phases} \end{cases}$$

$$3+1 \begin{cases} 6 \\ 3 \\ 3 \end{cases}$$

$$3+N \begin{cases} 3+3N \\ 1+2N \\ 2+N \end{cases}$$

In general, we have additional sources of CPV

LBL transition probability in 3-flavor

$$P_{\nu_\mu \rightarrow \nu_e}^{3\nu} = P^{\text{ATM}} + P^{\text{SOL}} + P^{\text{INT}}$$

in vacuum:

$$P^{\text{ATM}} = 4s_{23}^2 s_{13}^2 \sin^2 \Delta$$

$$P^{\text{SOL}} = 4c_{12}^2 c_{23}^2 s_{12}^2 (\alpha \Delta)^2$$

$$P^{\text{INT}} = 8s_{23}s_{13}c_{12}c_{23}s_{12}(\alpha \Delta) \sin \Delta \cos(\Delta + \delta_{CP})$$

$$\Delta = \frac{\Delta m_{31}^2 L}{4E}, \quad \alpha = \frac{\Delta m_{21}^2}{\Delta m_{31}^2}$$

$$\Delta \sim \pi/2$$

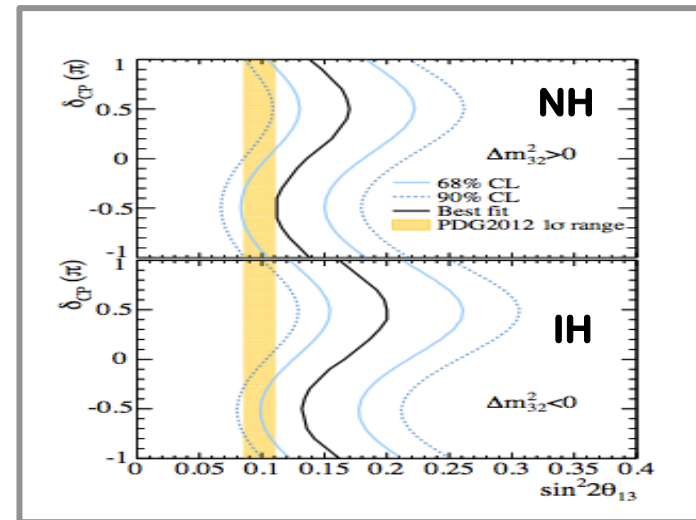
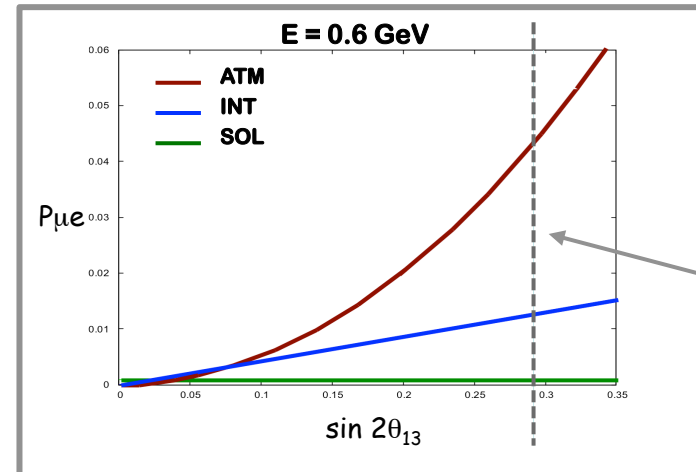
$$\alpha \sim 0.03$$

P^{ATM} leading $\rightarrow \theta_{13} > 0$

P^{INT} subleading \rightarrow dependency on δ

P^{SOL} negligible

Matter effects break
NH-IH degeneracy



A new interference term in the 3+1 scheme

N. Klop & A.P., PRD (2015)

- $\Delta_{14} \gg 1$: fast oscillations are averaged out

- But interference of Δ_{14} & Δ_{13} survives and is observable

$$P_{\mu e}^{4\nu} \simeq P^{\text{ATM}} + P_{\text{I}}^{\text{INT}} + P_{\text{II}}^{\text{INT}}$$

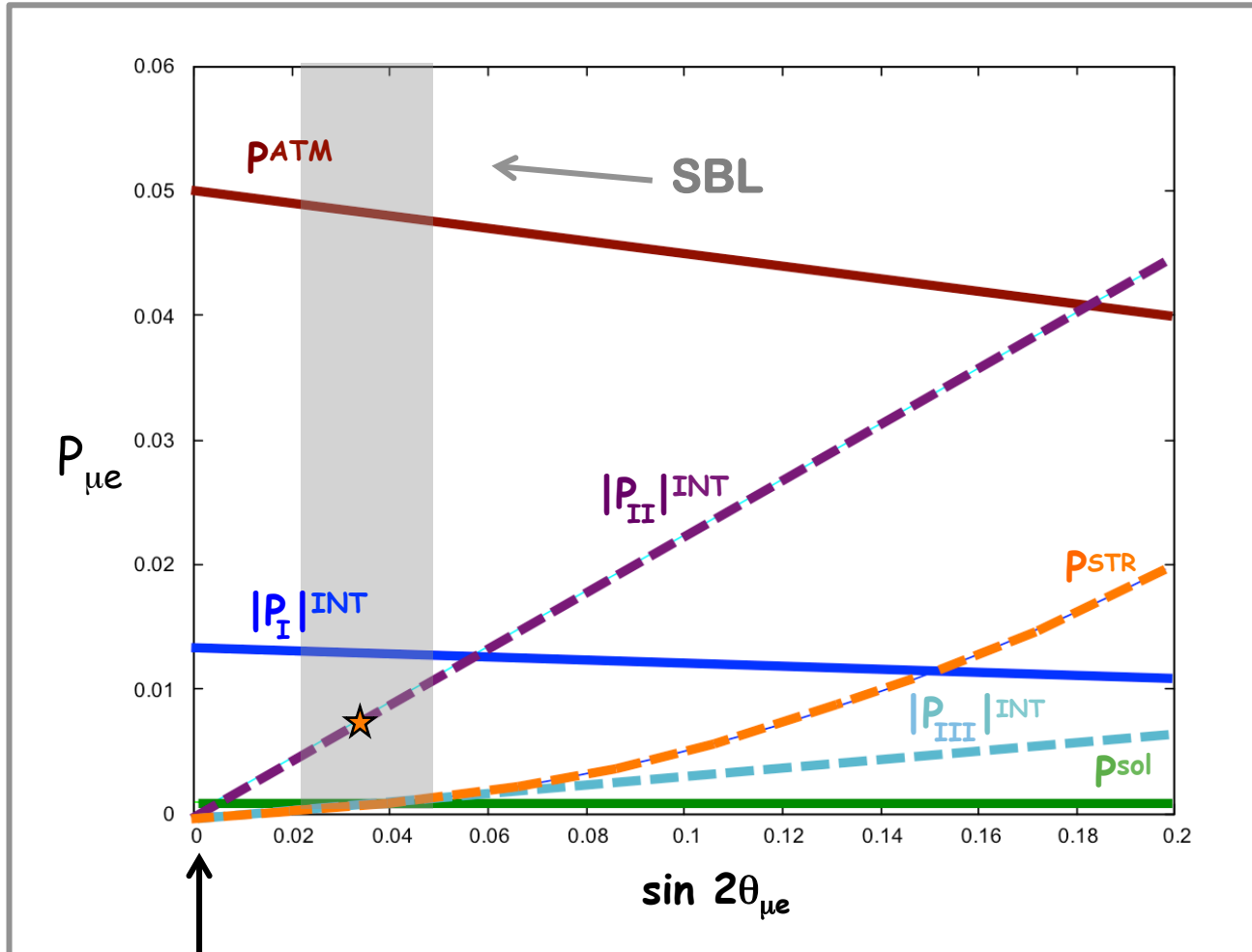
$$\begin{aligned} S_{13} \sim S_{14} \sim S_{24} &\sim 0.15 \sim \epsilon \\ \alpha = \delta m^2 / \Delta m^2 &\sim 0.03 \sim \epsilon^2 \end{aligned}$$

$$\begin{cases} P^{\text{ATM}} \simeq 4s_{23}^2 \underline{s_{13}^2} \sin^2 \Delta & \sim \epsilon^2 \\ P_{\text{I}}^{\text{INT}} \simeq 8\underline{s_{13}} \underline{s_{23}} c_{23} s_{12} c_{12} (\alpha \Delta) \sin \Delta \cos(\Delta + \delta_{13}) & \sim \epsilon^3 \\ P_{\text{II}}^{\text{INT}} \simeq 4\underline{s_{14}} \underline{s_{24}} \underline{s_{13}} \underline{s_{23}} \sin \Delta \sin(\Delta + \delta_{13} - \delta_{14}) & \sim \epsilon^3 \end{cases}$$

Sensitivity to the new CP-phase δ_{14}

Amplitude of the new interference term

N. Klop & A.P., PRD (2015)

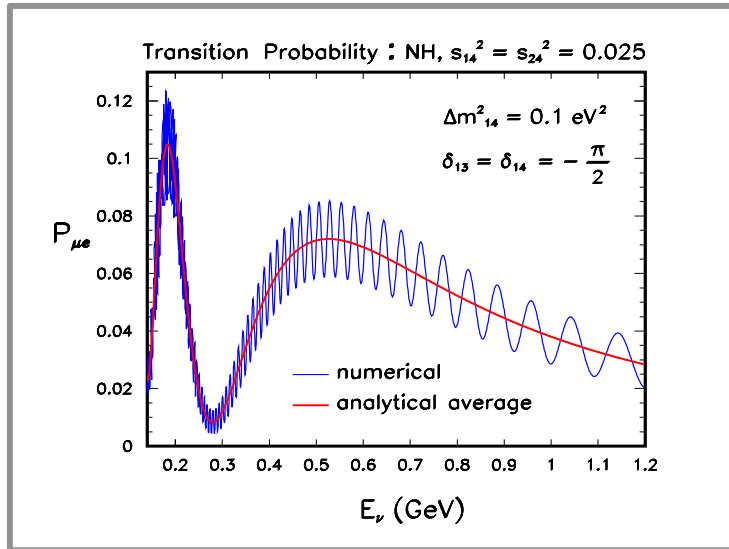


T2K
 $\theta_{13} = 9^\circ$
 $E = 0.6 \text{ GeV}$

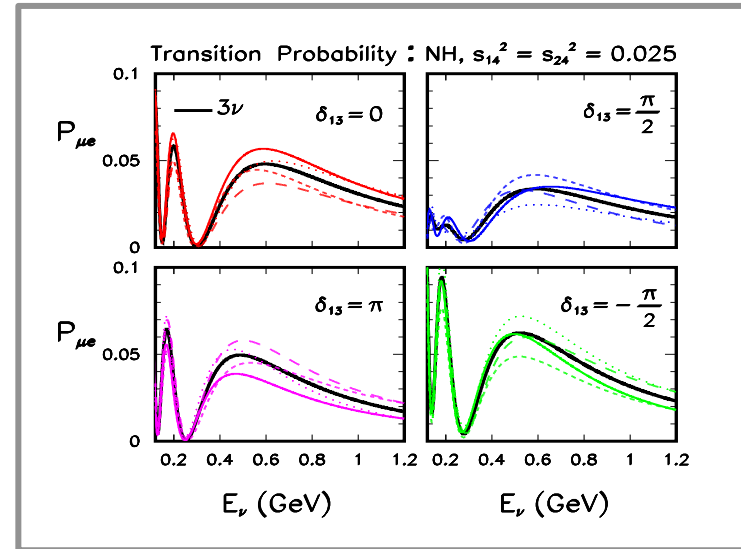
$$\sin^2 2\theta_{\mu e} = 4|U_{e4}|^2|U_{\mu 4}|^2$$

3v limit

Numerical examples of 4ν probability



The fast oscillations get averaged out due to the finite energy resolution



Different line styles



Different values of δ_{14}

The modifications induced by δ_{14} are almost as large as those induced by the standard CP-phase δ_{13}

Consequences...



LBL constraints on sterile ν s: present

A.P., PRD (RC) 91, 091301 (2015)
arXiv:1503.03966

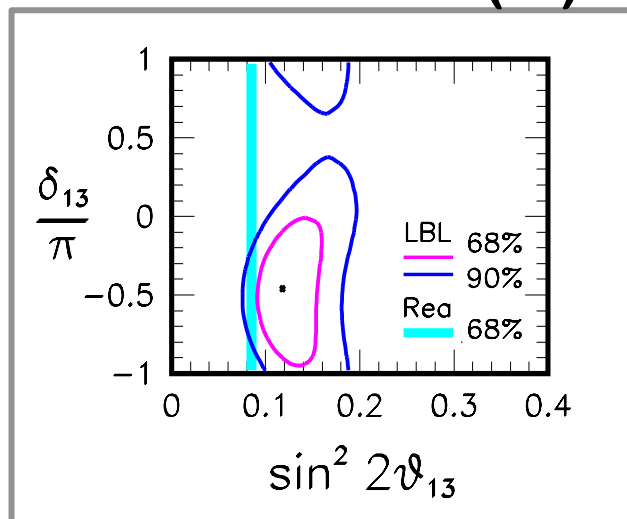
A.P., PLB 757, 142 (2016)
arXiv:1509.03148

Capozzi, Giunti, Laveder & A.P.,
PRD 95 (2017)
arXiv:1612.07764

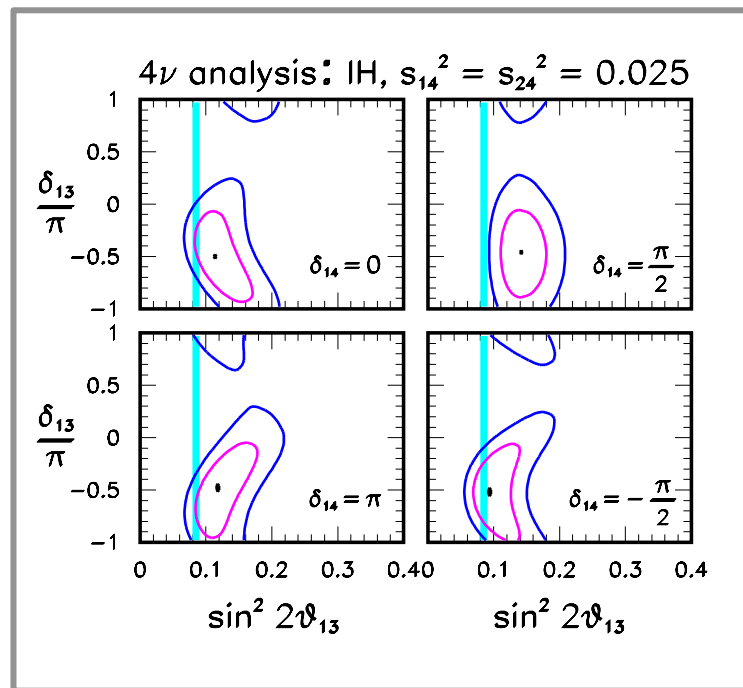
LBL constraints change in the 3+1 scheme

PLB (2016)

3 ν : T2K + NO ν A (IH)



4 ν →

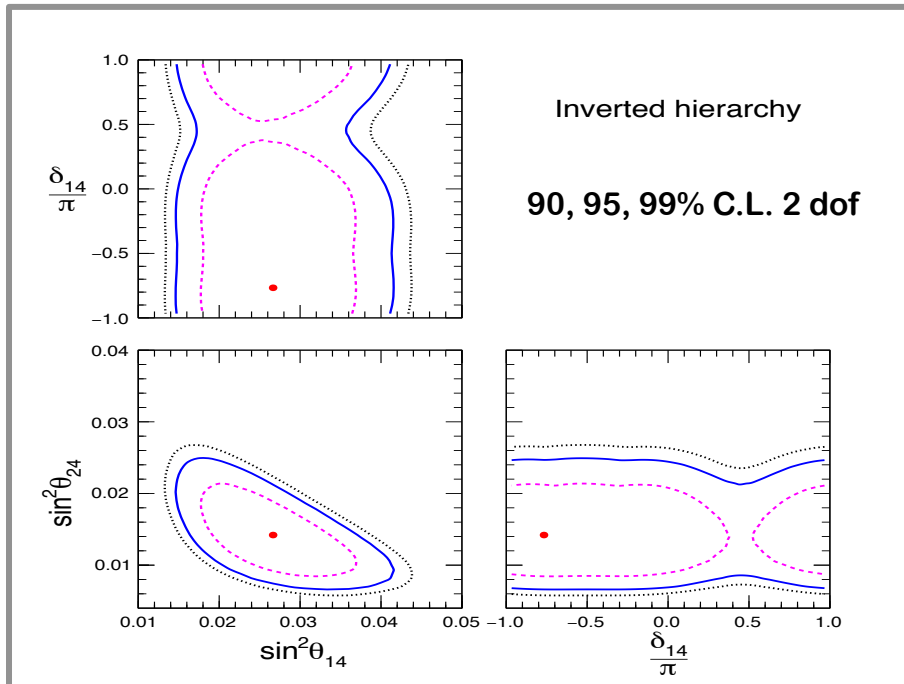


- The level of (dis-)agreement of LBL & Rea. depends on δ_{14}
- In this analysis θ_{14} and θ_{24} are fixed at the SBL best fit values
- These results call for a more refined analysis ...

Joint SBL and LBL constraints on $[\theta_{14}, \theta_{24}, \delta_{14}]$

PRD (2017)

SBL + LBL



SBL (all available data)

(Icecube and NEOS not included in this analysis)

LBL \equiv T2K + NO ν A

(Neutrino 2016 data)

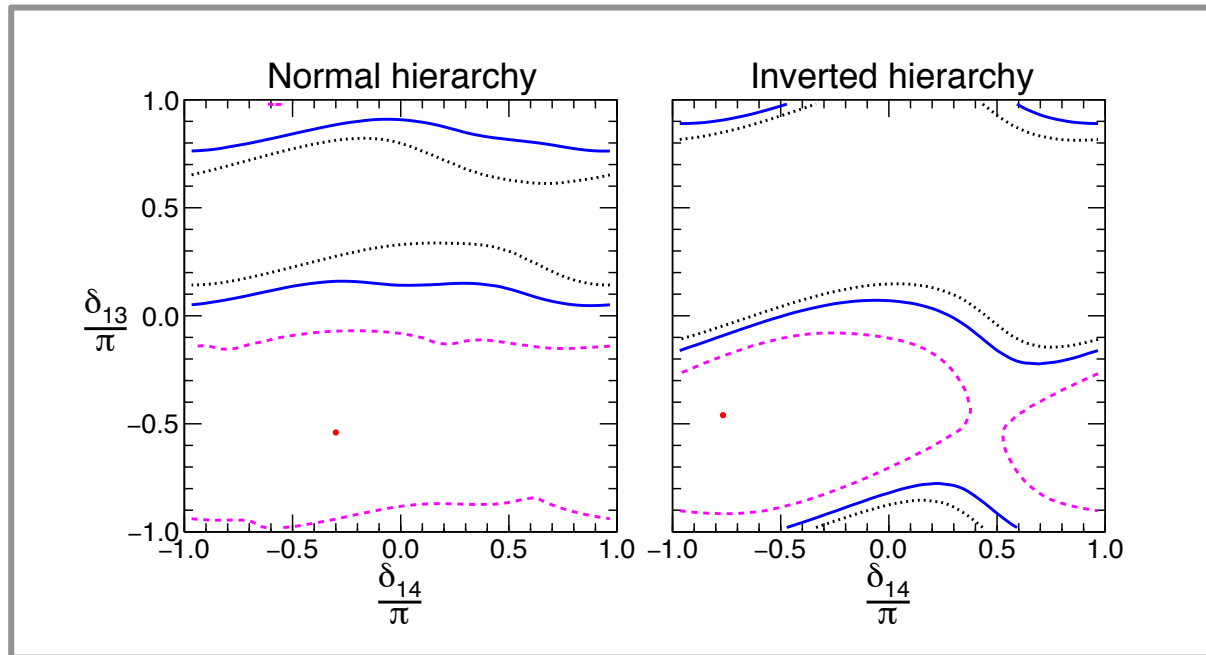
- $[\theta_{14}, \theta_{24}]$ determined by SBL experiments

- δ_{14} constrained by LBL experiments

Constraints on the two CP-phases

PRD (2017)

SBL + LBL

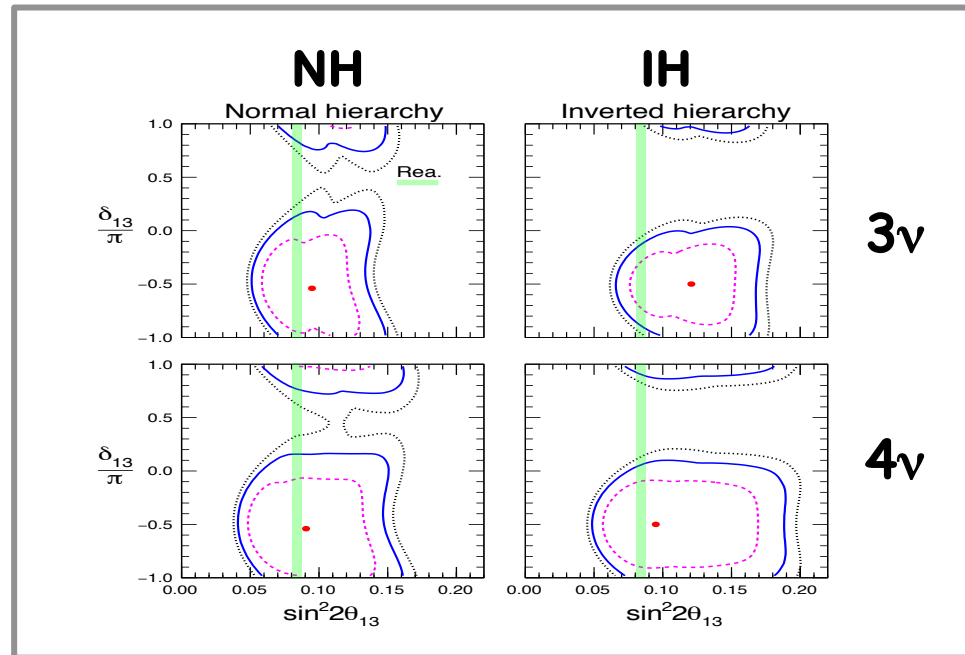


- δ_{13} is more constrained than δ_{14}
- Best fit values: $\delta_{13} \sim \delta_{14} \sim -\pi/2$
- This information cannot be extracted from SBL alone !

Impact on the standard parameters [θ_{13}, δ_{13}]

PRD (2017)

SBL + LBL

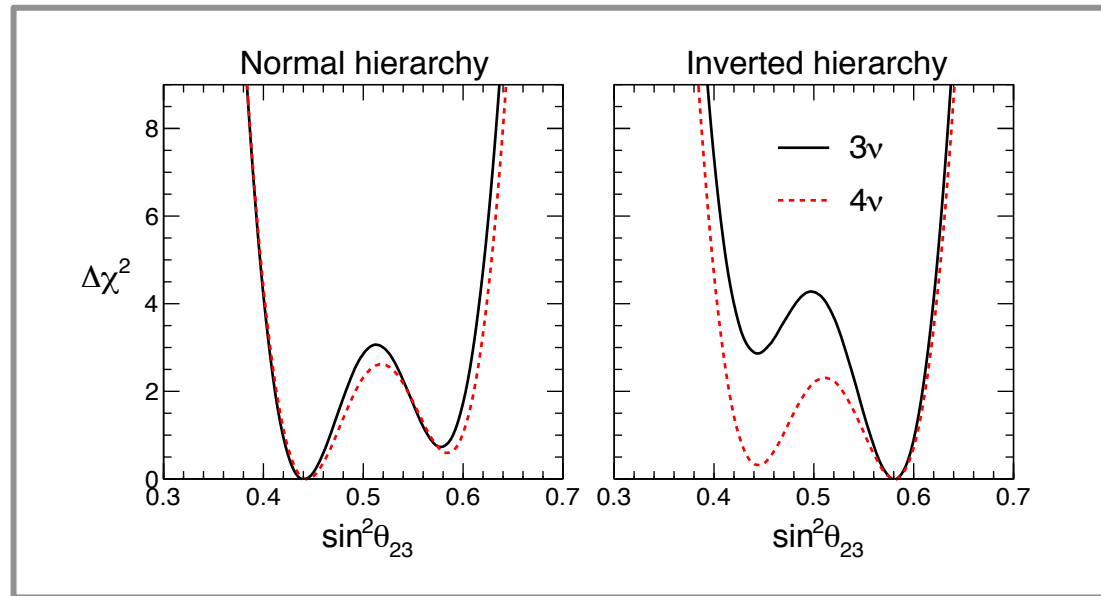


- Allowed range for θ_{13} from LBL alone gets enlarged
- Values preferred for $\delta_{13} \equiv \delta$ basically unaltered
- Mismatch (in IH) of LBL and Reactors decreases in 3+1

Impact of sterile neutrinos on θ_{23}

PRD (2017)

SBL + LBL



Indication for non-maximal θ_{23} persists in 3+1 scheme

Preference for θ_{23} octant disappears in 3+1 scheme

Octant fragility seems to be a general feature (see later)



Looking to the future

Agarwalla, Chatterjee, Dasgupta, A.P.,
arXiv: 1601.05995 (JHEP 2016)

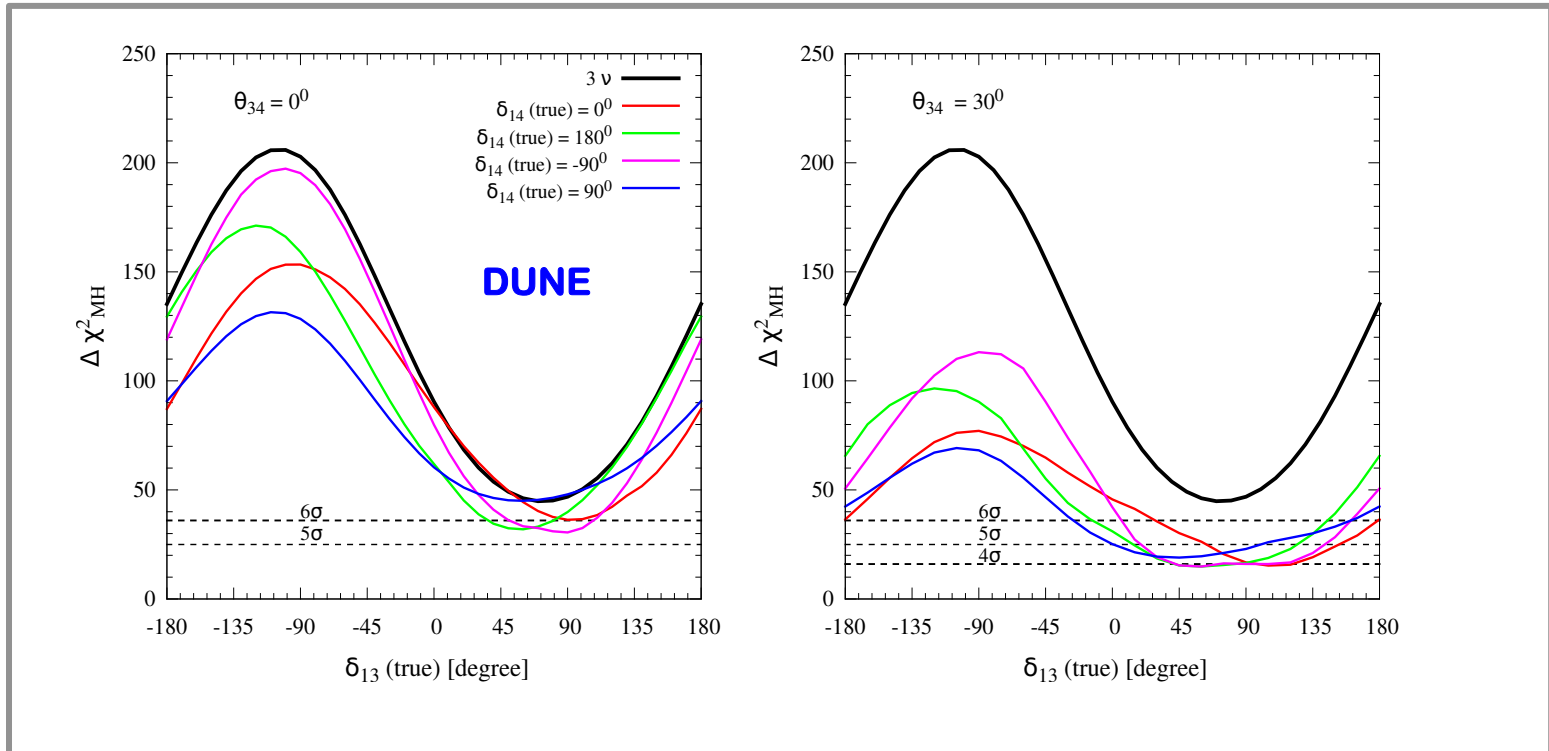
Agarwalla, Chatterjee, A.P.,
arXiv: 1603.03759 (JHEP 2016)

Agarwalla, Chatterjee, A.P.,
arXiv: 1607.01745 (PLB 2016)

Agarwalla, Chatterjee, A.P.,
arXiv: 1605.04299 (PRL 2017)

Discovery potential of mass hierarchy

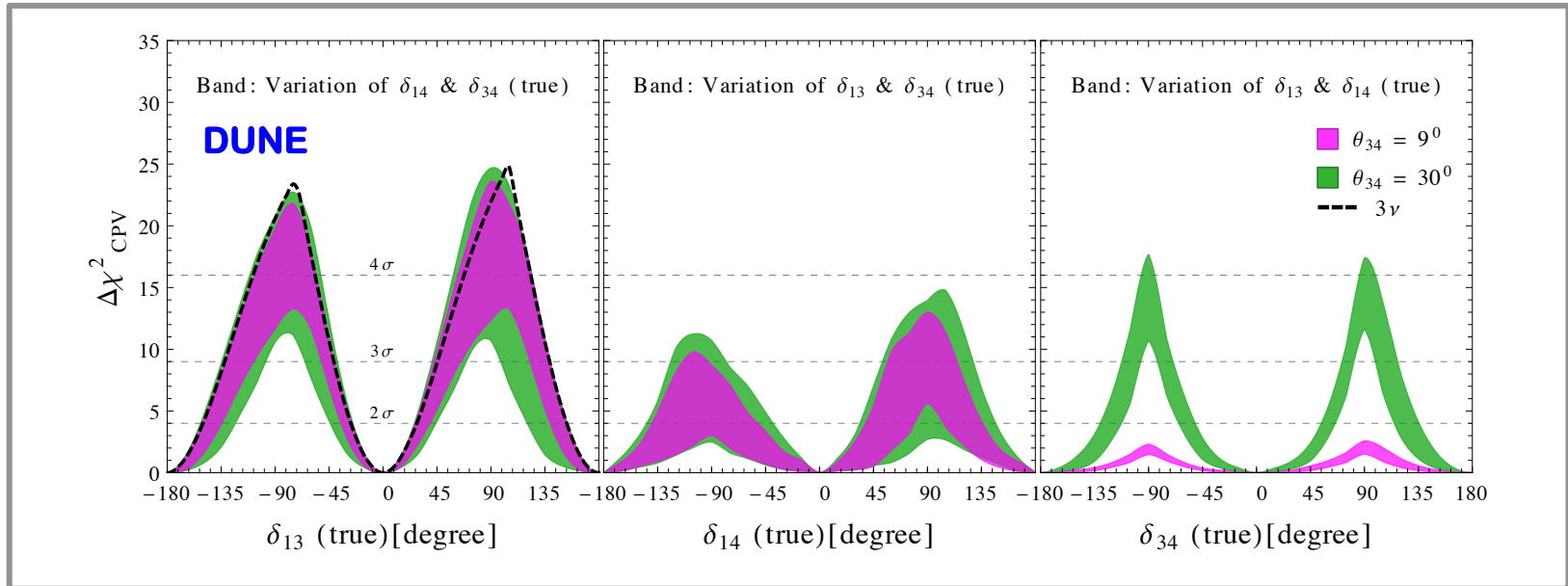
JHEP 2016



Degradation of sensitivity but 4 σ level preserved

CPV discovery potential

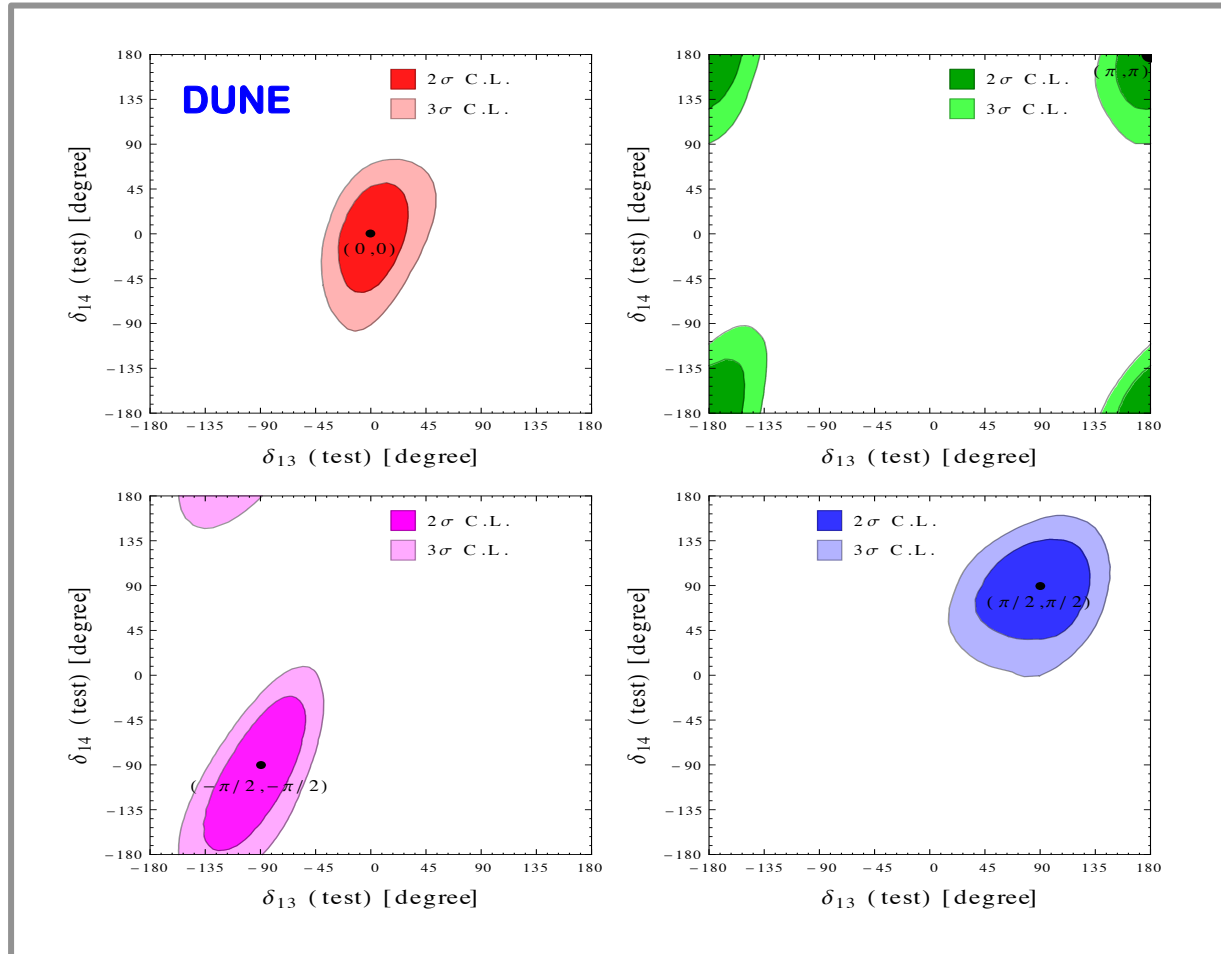
JHEP 2016



- Sensitivity to CPV induced by δ_{13} reduced in 3+1 scheme
- Potential sensitivity also to the new CP-phases δ_{14} e δ_{34}
- Clear hierarchy in the sensitivity: $\delta_{13} > \delta_{14} > \delta_{34}$ for $\theta_{14} = \theta_{24} = \theta_{34} = 9^\circ$

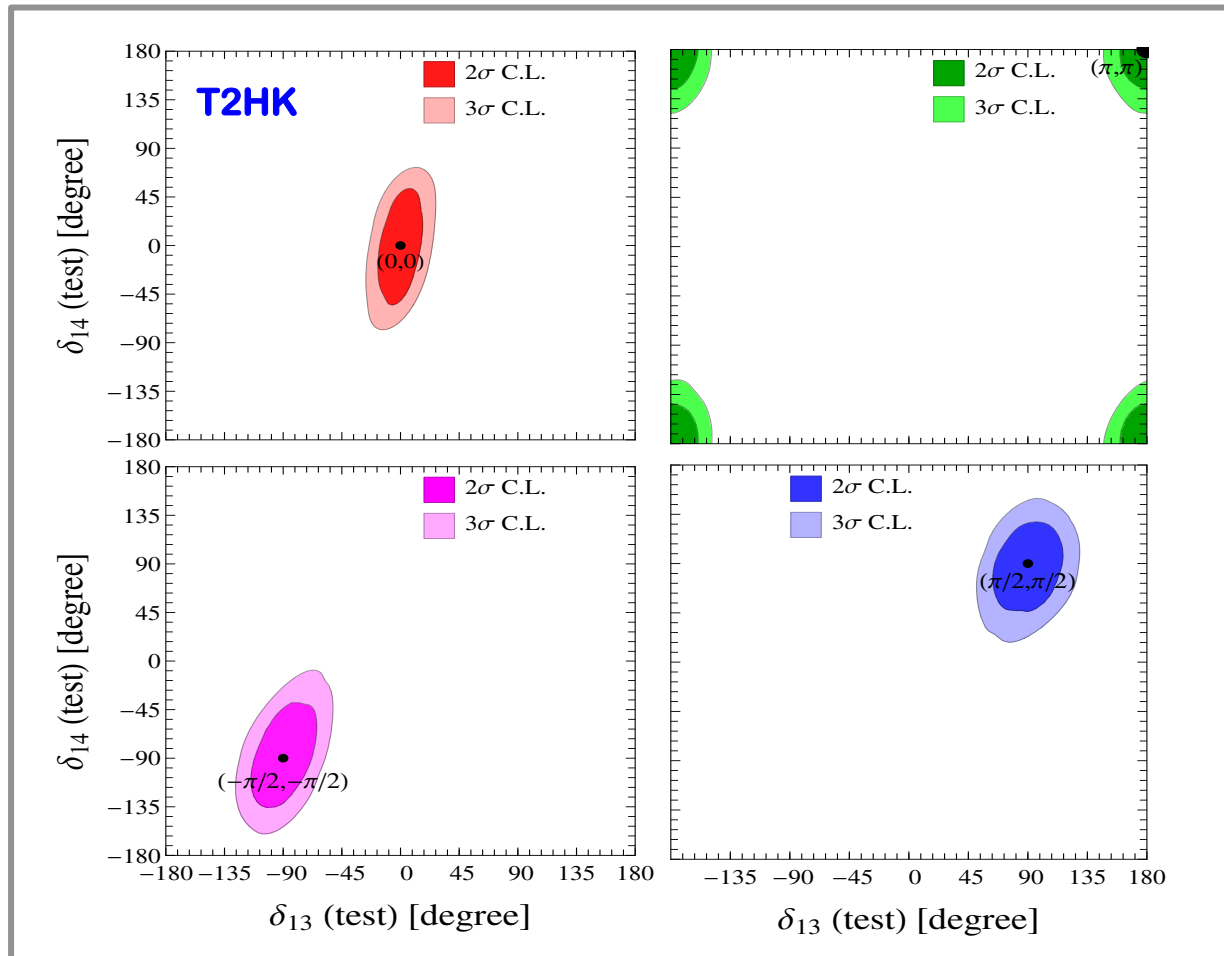
Reconstruction of the CP phases in DUNE

JHEP 2016



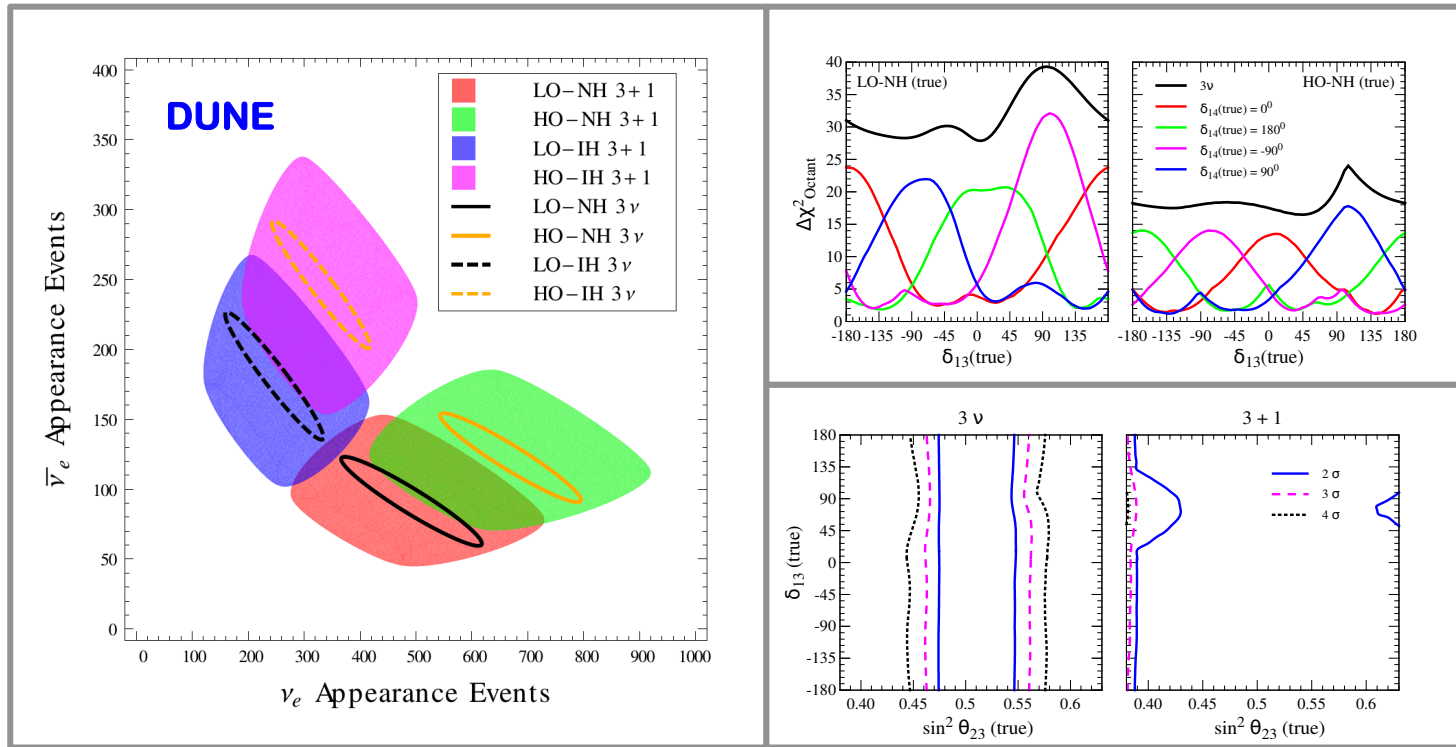
Reconstruction of the CP phases in T2HK

Preliminary plot realized by S.S. Chatterjee



Octant of θ_{23} in danger with a sterile neutrino

PRL 2017



Distinct ellipses (3ν) become overlapping blobs (3+1)

For unfavorable combinations of δ_{13} & δ_{14} sensitivity is lost

Conclusions

- **Sterile neutrinos are sources of additional CPV**
- **Full exploration of new CPV possible only with LBL**
- **LBL experiments complementary to the SBL ones**



**Looking ahead to
the next mooring!**

**Thank you
for your
attention!**

Back up slides

CPV and averaged oscillations

$$A_{\alpha\beta}^{\text{CP}} \equiv P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta)$$

$$A_{\alpha\beta}^{\text{CP}} = -16 J_{\alpha\beta}^{12} \sin \Delta_{21} \underbrace{\sin \Delta_{13} \sin \Delta_{32}}$$

if

$$\Delta \equiv \Delta_{13} \simeq \Delta_{23} \gg 1$$

osc. averaged out by finite E resol.

→

$$\langle \sin^2 \Delta \rangle = 1/2$$

It can be:

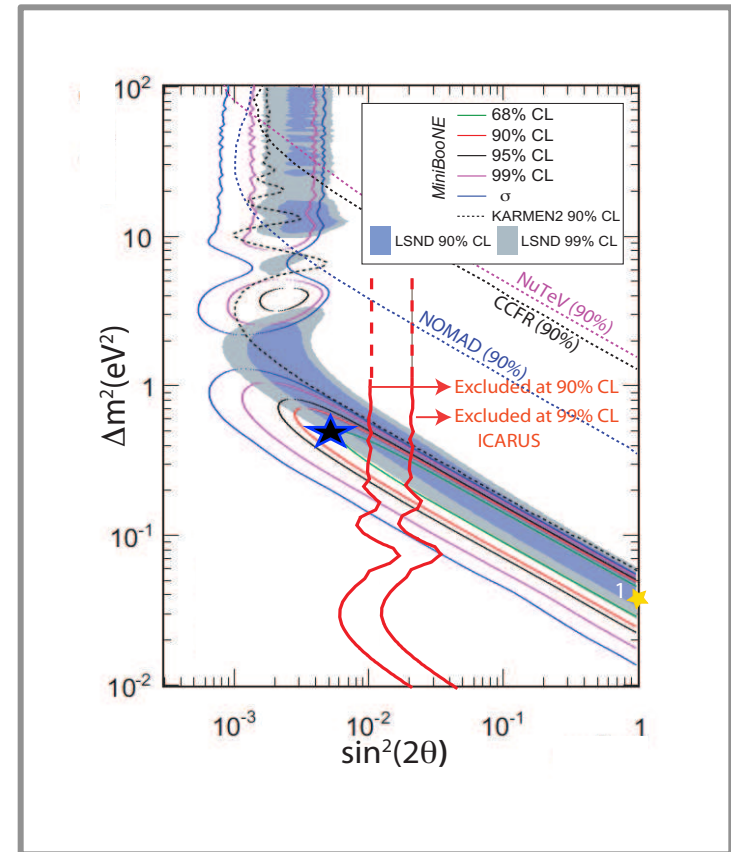
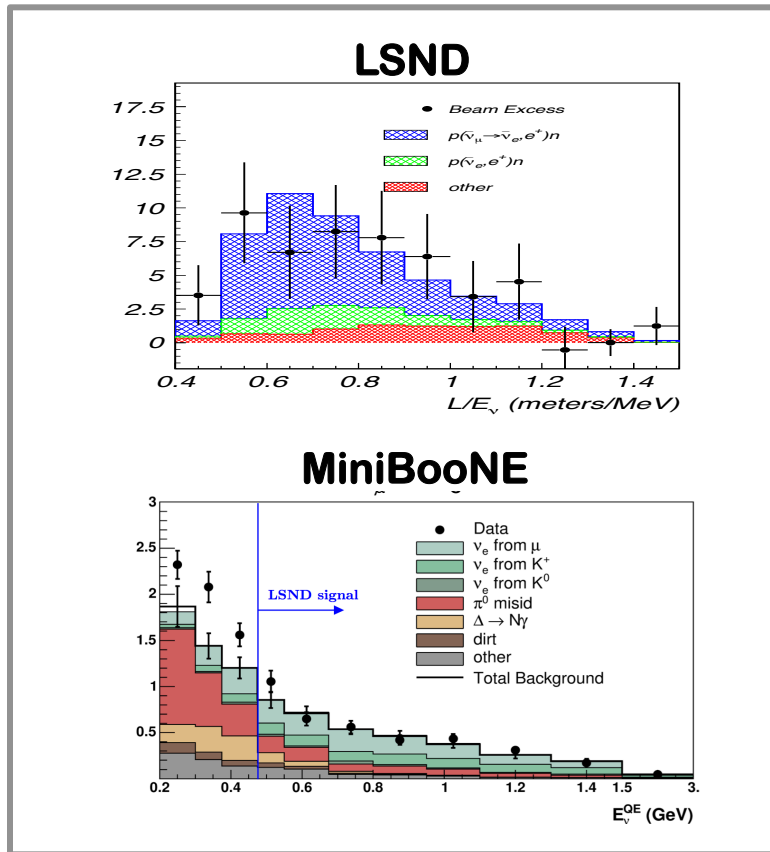
$$A_{\alpha\beta}^{\text{CP}} \neq 0$$

(if $\sin \delta \neq 0$)

The bottom line is that if one of the three ν_i is ∞ far from the other two ones this does not erase CPV
(relevant for the 4 ν case)

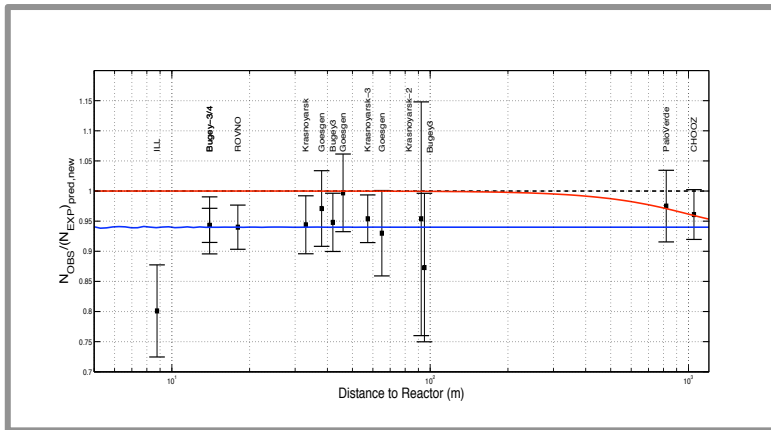
The SBL accelerator anomalies

(unexplained ν_e appearance in a ν_μ beam)

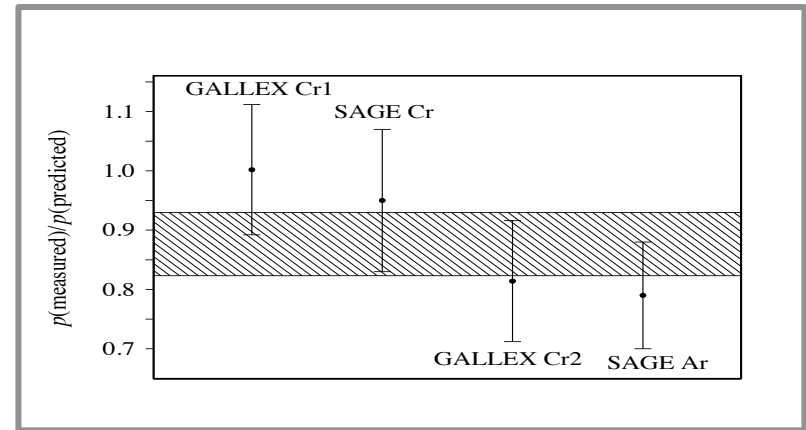


The reactor and gallium anomalies

(unexplained ν_e disappearance)



Mention et al. arXiv:1101:2755 [hep-ex]

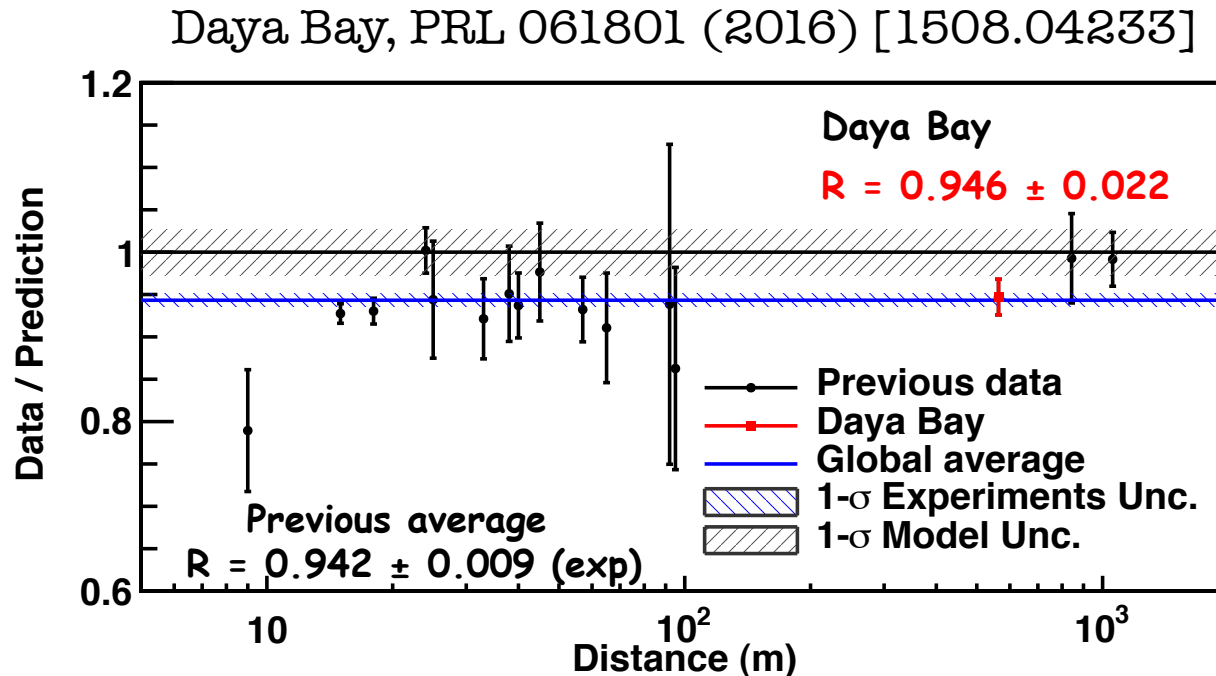


SAGE coll., PRC 73 (2006) 045805

Warning: both are mere normalization issues

The culprit may be hidden in unknown systematics

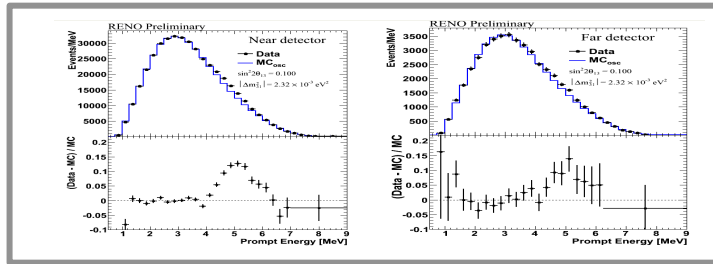
New-generation detectors confirm deficit



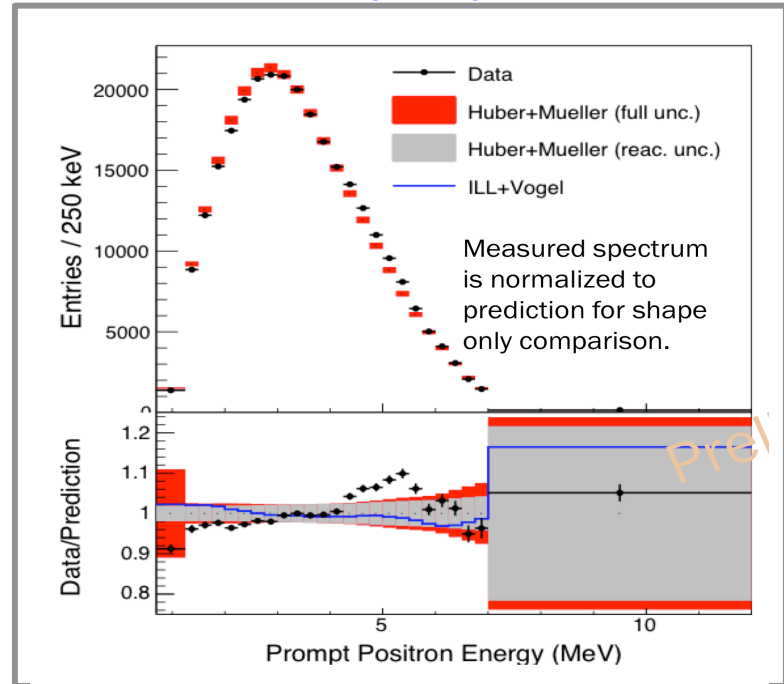
However, the same detectors give us a warning ...

Understanding of reactor spectrum is incomplete

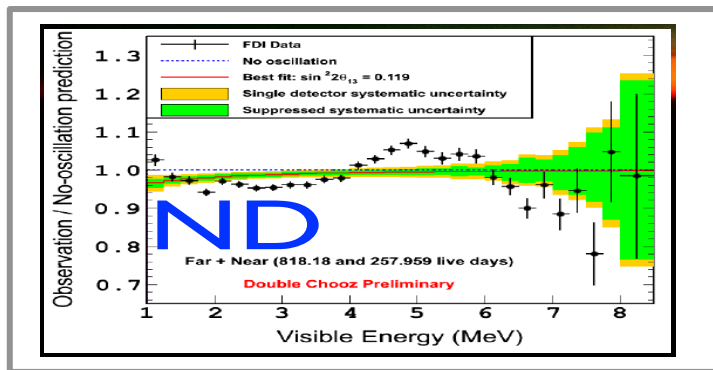
RENO



Daya Bay



Double-CHOOZ



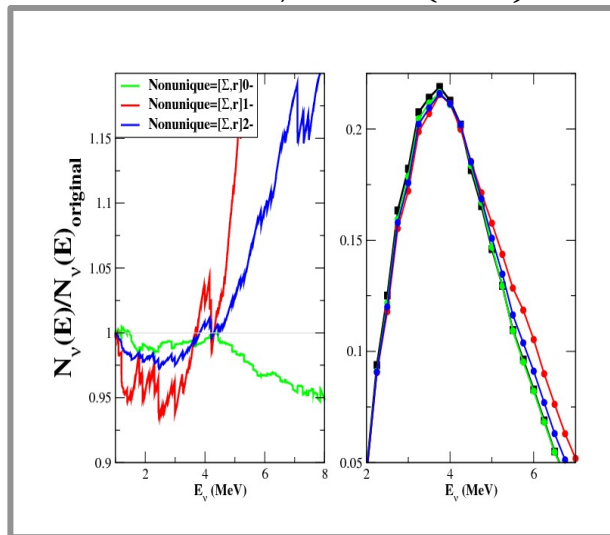
Bump/shoulder at 5 MeV observed in all the three experiments

Found both a near & far sites: not imputable to new osc. physics

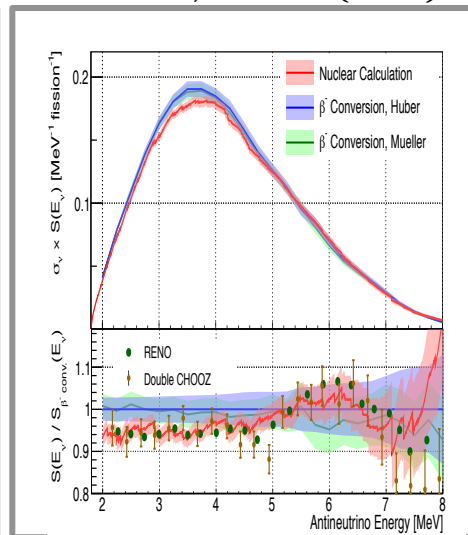
θ_{13} extraction is unaffected (based on near/far comparison)

5 MeV bump is under active investigation

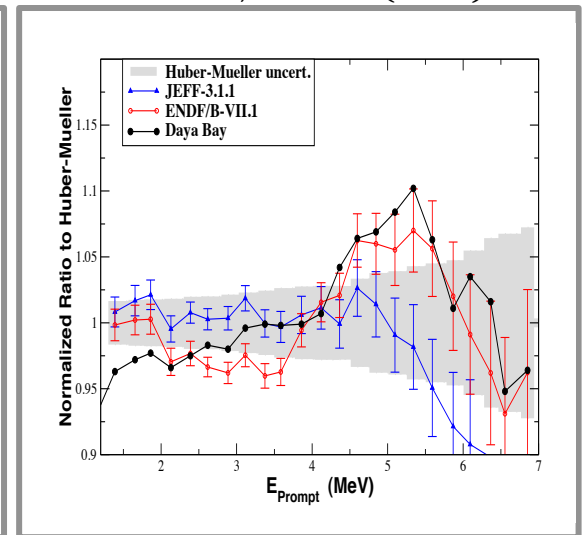
Hayes et al.
PRL 112, 202501 (2014)



Dwyer and Langford
PRL 114, 012502 (2015)



Hayes et al.
PRD 92, 033015 (2015)



- Systematics in reactor spectra not entirely under control
- Dissimilar results with two different nuclear databases
- Normalization & spectral shape issues not necessarily related
- New SBL experiments needed to shed light on both issues

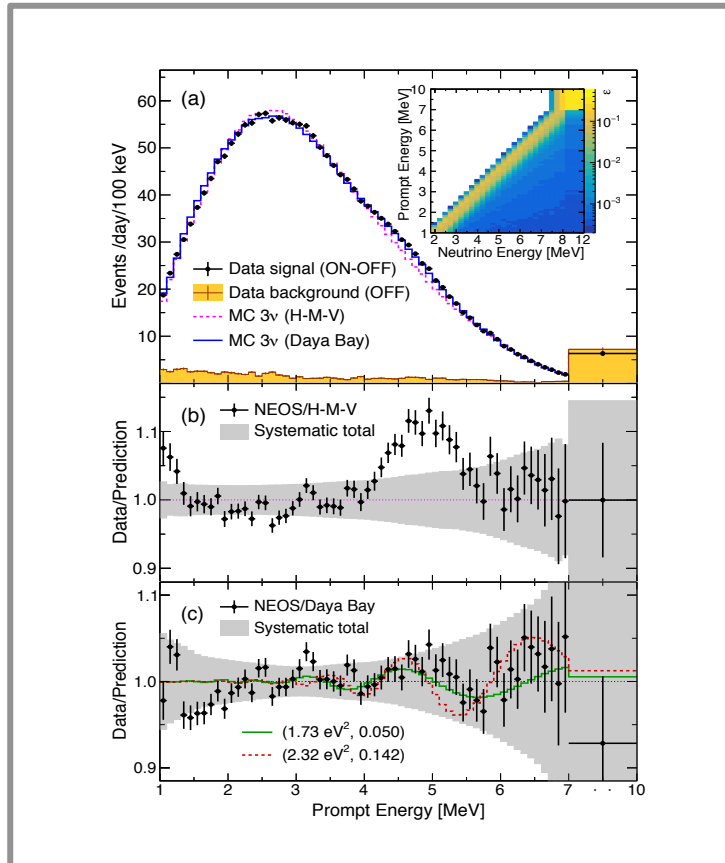
NEOS

Hanbit Nuclear Power
Complex, Korea

Detector: 1 ton
Gd-loaded liquid
scintillator 24 m from
the reactor core

Daya-Bay absolute
spectrum used as
a normalization

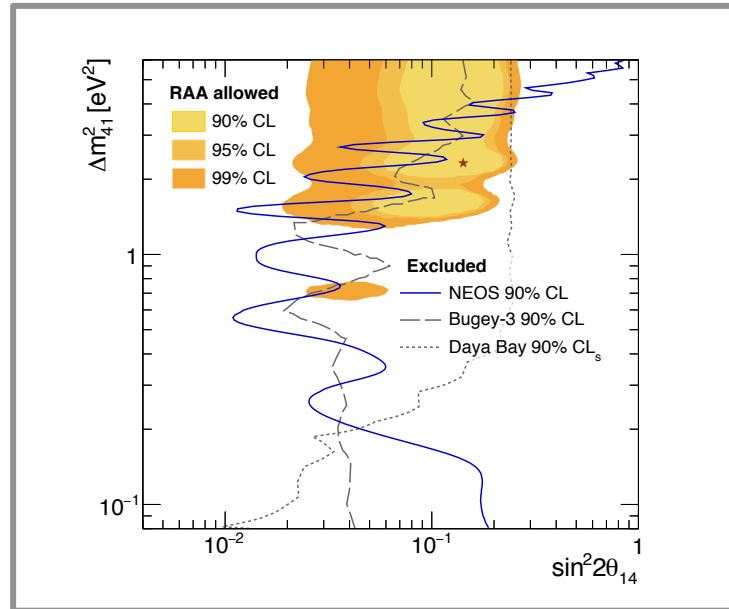
Oscillating pattern
visible after normalization



NEOS arXiv:1610:05134

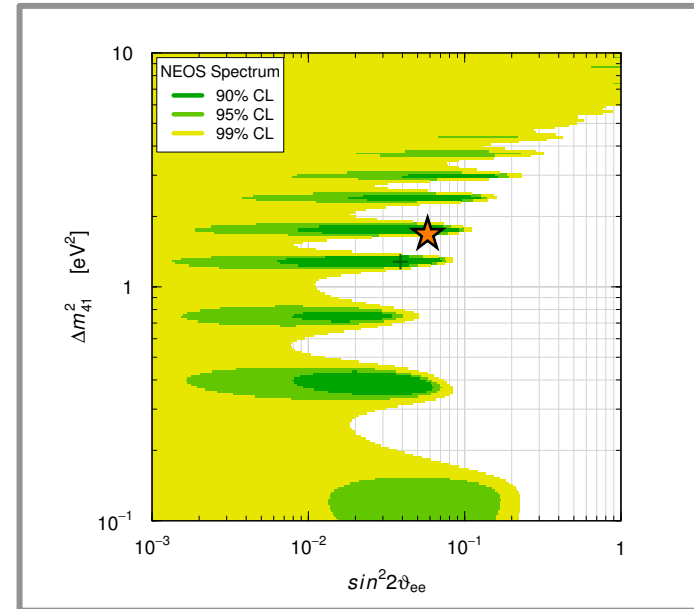
Two different perspectives

negative view



NEOS, arXiv:1610:05134

positive view



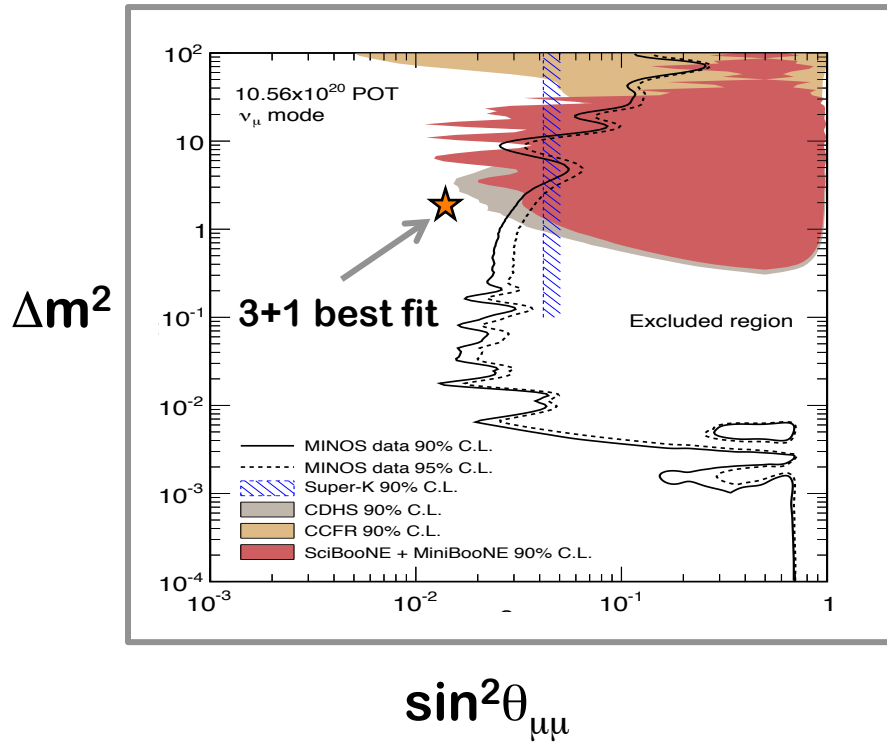
Gariazzo et al., arXiv: 1703.00860

Best fit: $\Delta m^2 = 1.73 \text{ eV}^2$ $\sin^2 2\theta = 0.05$

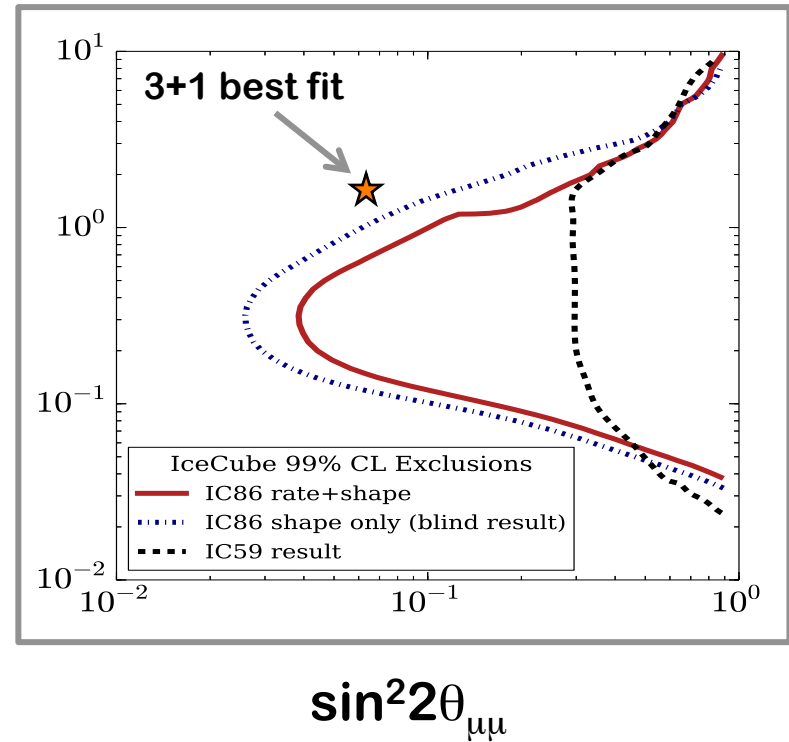
$$\chi_{\text{no osc}}^2 - \chi_{\text{min}}^2 = 6.5 \quad > \mathbf{95\% \text{ CL indication!}}$$

No anomaly in ν_μ disappearance

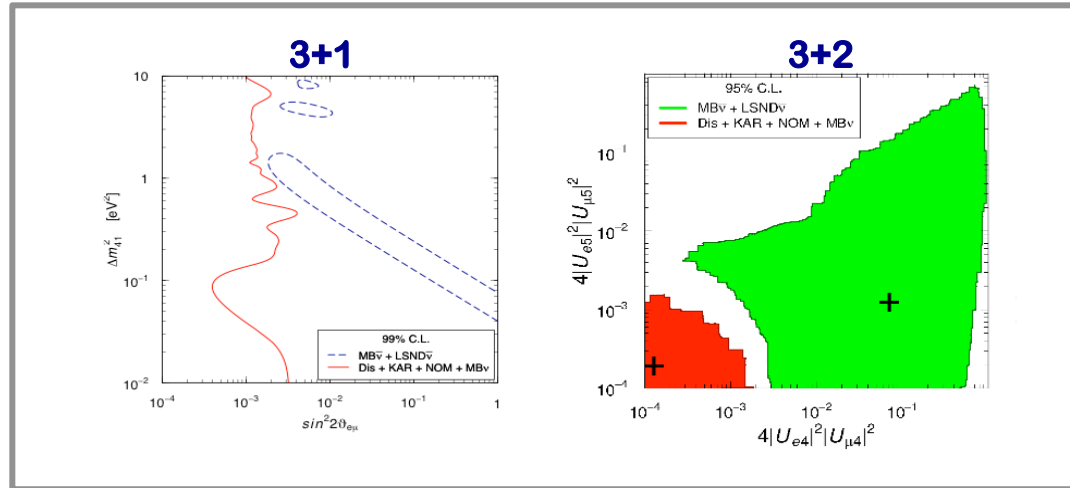
SBL & MINOS (NC)



IceCube



Tension in all ν_s models



Giunti
&
Laveder

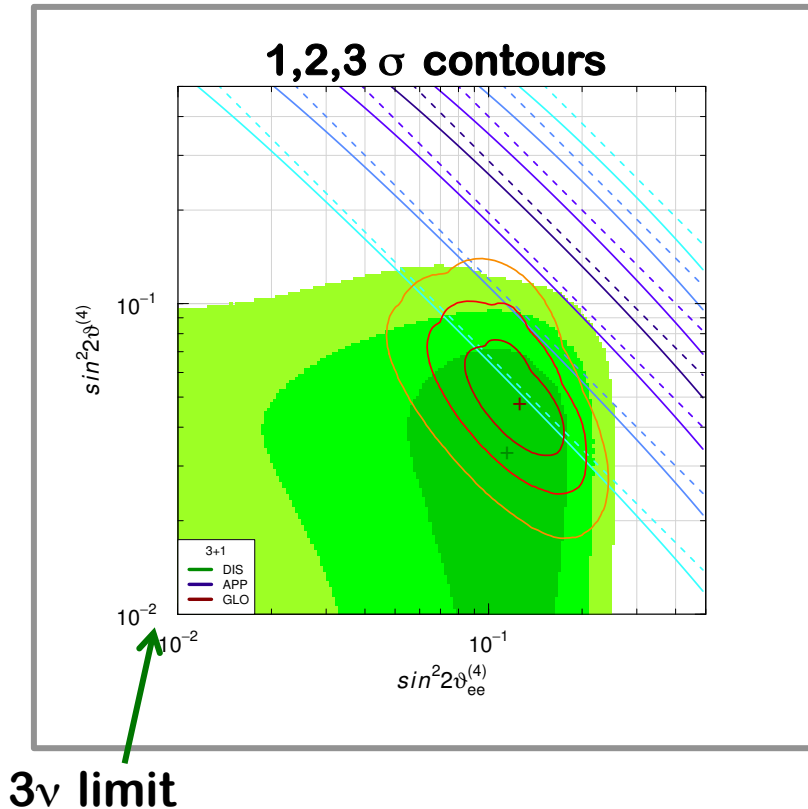
arXiv:1107.1452

$\nu_\mu \rightarrow \nu_e$ **positive**
 $\nu_e \rightarrow \nu_e$ **positive**
 $\nu_\mu \rightarrow \nu_\mu$ **negative**

$|U_{e4}||U_{\mu4}| > 0$
 $|U_{e4}| > 0$
 $|U_{\mu4}| \sim 0$

$$\sin^2 2\theta_{e\mu} \simeq \frac{1}{4} \sin^2 2\theta_{ee} \sin^2 2\theta_{\mu\mu} \simeq 4|U_{e4}|^2|U_{\mu4}|^2$$

An “undecidable” problem



App. & Dis. barely overlap at 2σ level

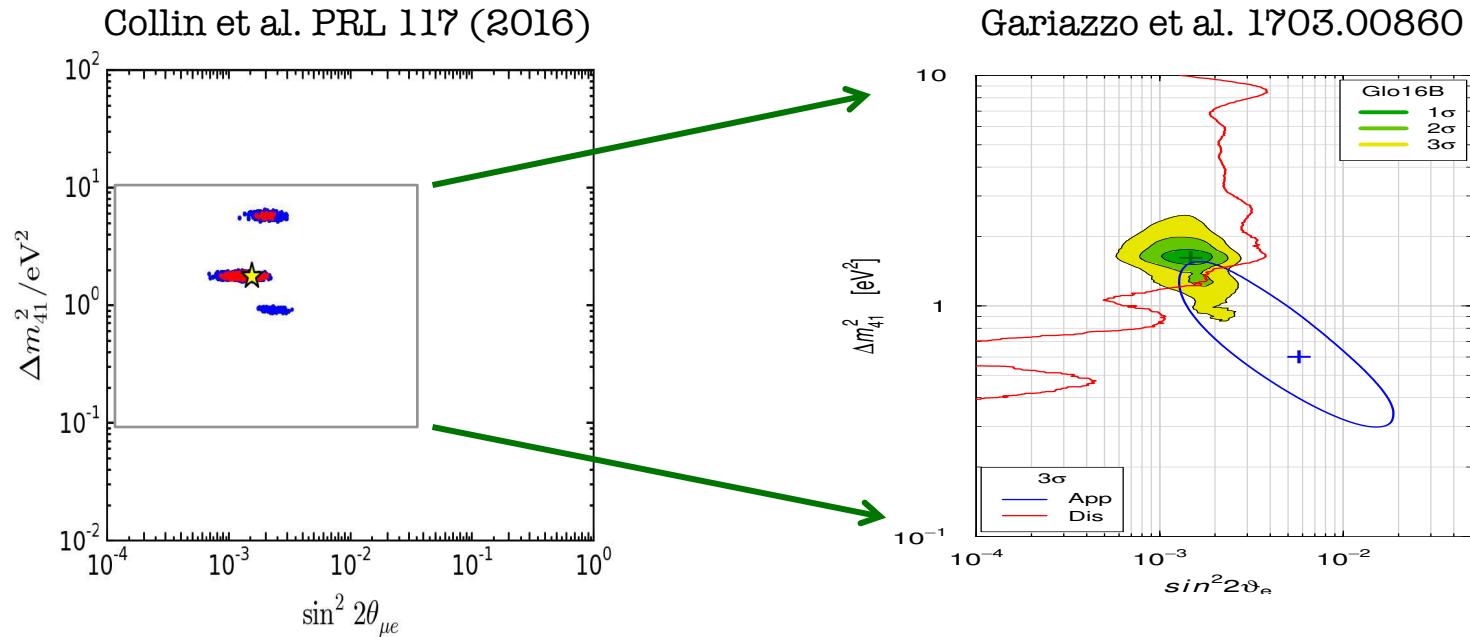
But their combination gives a 6σ improvement with respect to the 3ν case

Difficult to take a decision on sterile ν s!

Only new more sensitive experiments can decide

Figure from Giunti & Zavanin, arXiv:1508:03172
(tension slightly increased after NEOS, MINOS, IceCube)

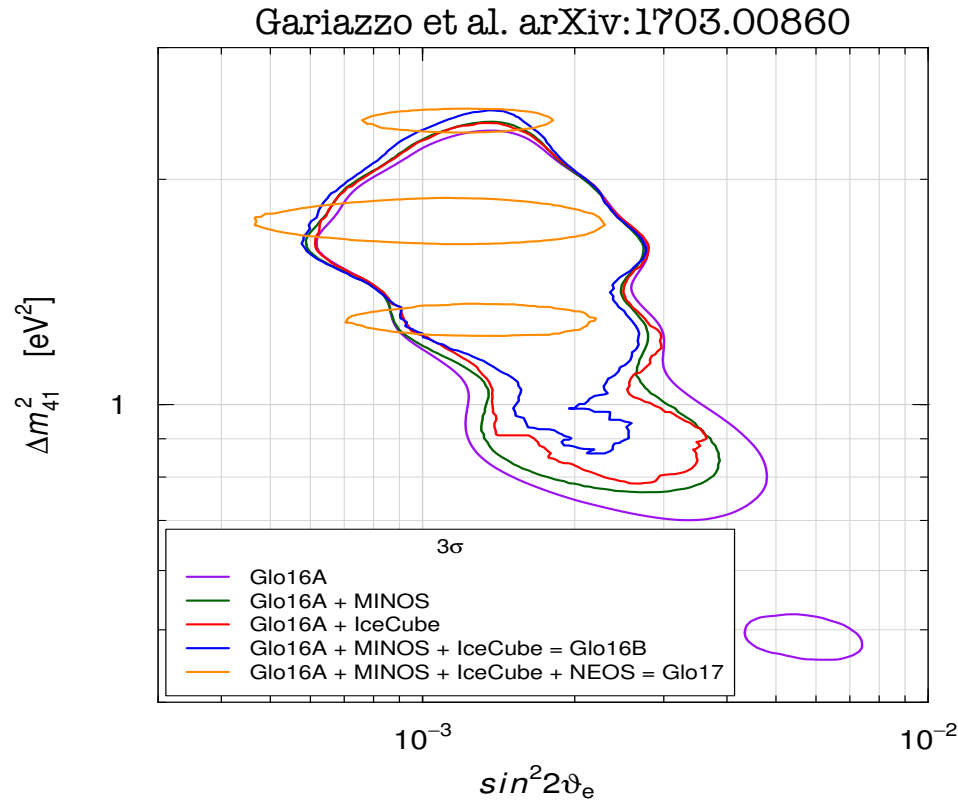
If one accepts to live with the tension



Both analyses include IceCube data

Similar best fit points around $\Delta m^2 \sim 1.7 \text{ eV}^2$

Impact of the latest measurements



NEOS selects a subregion of the region allowed by all the other data : very intriguing!

The SBL race for the light sterile neutrino

Gariazzo et al., arXiv: 1703.00860

