

Expectations for the muon $g-2$ in simplified models with dark matter

Enrico Maria Sessolo

Fakultät für Physik, TU Dortmund, Germany
National Centre for Nuclear Research, Warsaw, Poland

EPS-HEP 2017
Venice, July 7 2017

Based on:

K.Kowalska, EMS

1707.00753

Muon g-2

- **Measured at Brookhaven:**

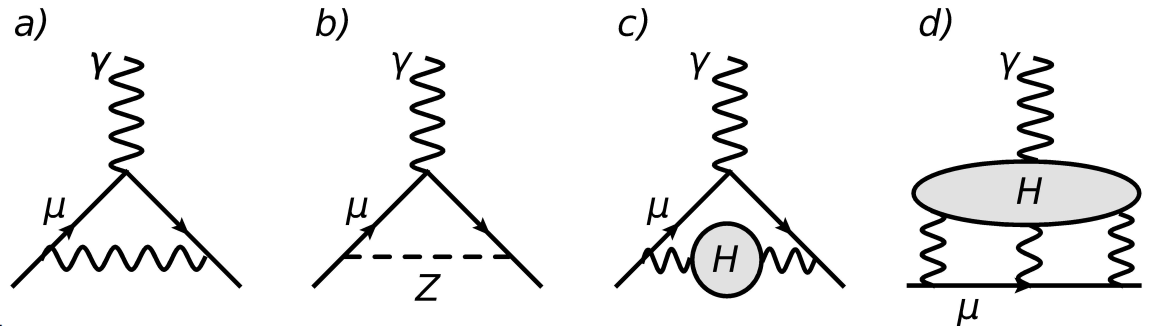
$$a_{\mu}^{\text{exp}} = 116592089(63) \times 10^{-11}$$

Bennet *et al.*, hep-ex/0602035 (PRD 2006)

- **Discrepancy with SM:**

$$a_{\mu}^{\text{SM}} = 116591802(49) \times 10^{-11}$$

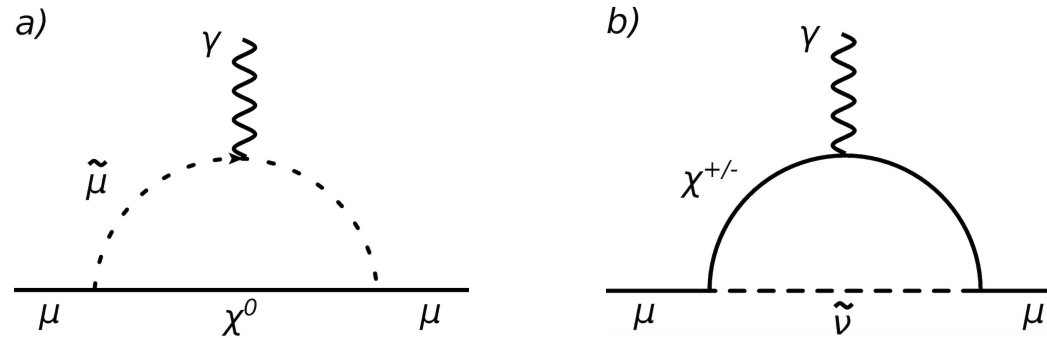
$$\delta(g-2)_{\mu} = (28.7 \pm 8.0) \times 10^{-10} \quad \mathbf{3.5 \sigma}$$



New upcoming experiments will test the anomaly...
Fermilab, J-PARC, 7σ!!

Muon $g-2$

Explanations of the anomaly abundant in the literature...
e.g. SUSY:



Interesting assumption in MSSM:

Same **BSM physics** is responsible for $\delta(g - 2)_\mu$ and $\Omega h^2 \approx 0.12$
(charginos, sleptons / neutralino DM)

We will apply the same assumption to generic frameworks...

Simplified models with dark matter approach to the muon g-2

- The dark matter interacts with the muons through renormalizable couplings
- Interactions are CP conserving and invariant under the SM gauge group, $SU(2) \times U(1)$
- Each model satisfies the constraints from perturbativity and unitarity
- The measurement of the relic abundance is an active constraint on the parameter space.

Fermion-scalar-fermion:

$$\mathcal{L} = g_s \bar{\psi}_E \psi_\mu \phi_S + ig_p \bar{\psi}_E \gamma^5 \psi_\mu \phi_S + \text{h.c.}$$

Fermion-vector-fermion:

$$\mathcal{L} = g_v Z'_\nu \bar{\psi}_E \gamma^\nu \psi_\mu + g_a Z'_\nu \bar{\psi}_E \gamma^\nu \gamma^5 \psi_\mu + \text{h.c.}$$

 **This work!**

MINIMAL ASSUMPTIONS:
 Z_2 symmetry + vector-like new fermions

Investigated models

Models with a real neutral scalar SU(2) singlet

- Fermion SU(2) singlets
- Fermion SU(2) doublets
- Mixing singlet and doublet fermions

Models with a complex scalar SU(2) singlet

- Fermion SU(2) singlets
- Fermion SU(2) doublets
- Mixing singlet and doublet fermions

Models with an SU(2) scalar doublet

- Fermion SU(2) singlets
- Fermion SU(2) doublets
- Fermion SU(2) triplets
- Fermion adjoint triplets
- Mixing cases

Model 1: Scalar R singlet / fermion singlets

Lagrangian:

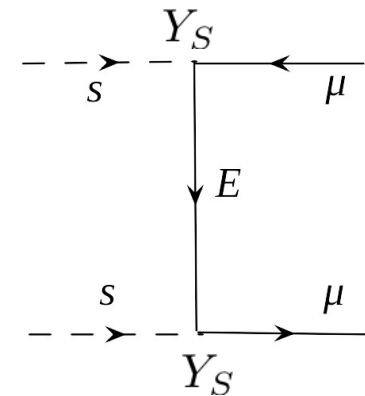
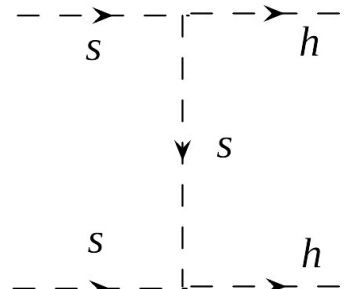
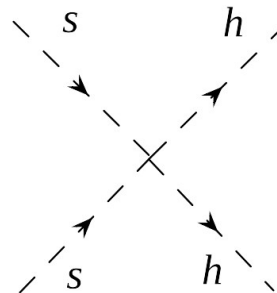
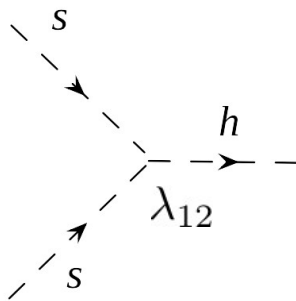
$$V = -\mu^2 \phi^\dagger \phi + \frac{\lambda}{2} (\phi^\dagger \phi)^2 + \frac{\mu_s^2}{2} s^2 + \frac{\lambda_s}{2} s^4 + \lambda_{12} s^2 \phi^\dagger \phi \quad (\text{s inert field, } \phi \text{ SM Higgs})$$

$$\mathcal{L} \supset -Y_S E e_R^* s - M_E E' E + \text{h.c.}$$

VL heavy lepton

muon

Relic abundance can be easily obtained:



Model 1: Scalar R singlet / fermion singlets

Lagrangian:

$$V = -\mu^2 \phi^\dagger \phi + \frac{\lambda}{2} (\phi^\dagger \phi)^2 + \frac{\mu_s^2}{2} s^2 + \frac{\lambda_s}{2} s^4 + \lambda_{12} s^2 \phi^\dagger \phi \quad (\text{s inert field, } \phi \text{ SM Higgs})$$

$$\mathcal{L} \supset -Y_S E e_R^* s - M_E E' E + \text{h.c.}$$

VL heavy lepton

muon

Muon g-2 contributions:

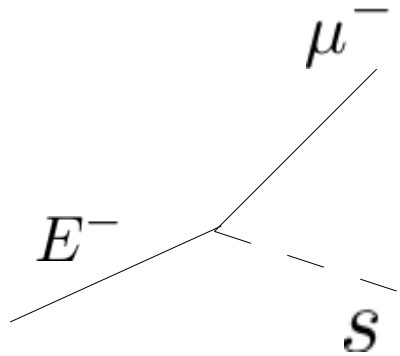
$$\delta(g-2)_\mu = \frac{1}{16\pi^2} \sum_{S^0, E^\pm} \left[\frac{m_\mu^2}{m_S^2} (|c_L|^2 + |c_R|^2) \mathcal{F}_1(r) + 2 \frac{m_\mu m_E}{m_S^2} \Re(c_L c_R^*) \mathcal{F}_2(r) \right]$$

Only piece in Model 1

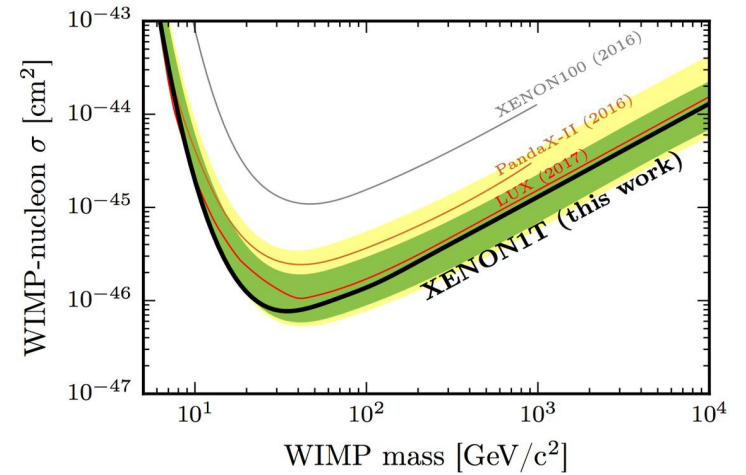
Constraints

**ATLAS, CMS: 2 lepton + missing E_T
Mono-jet**

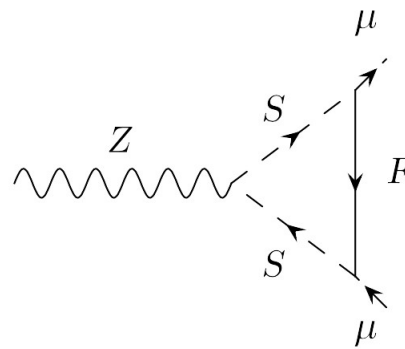
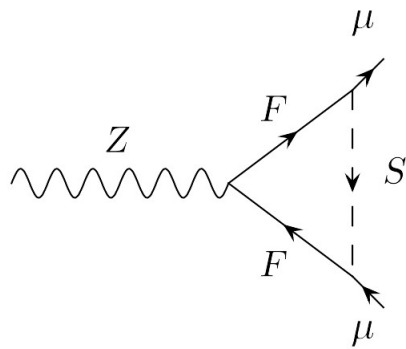
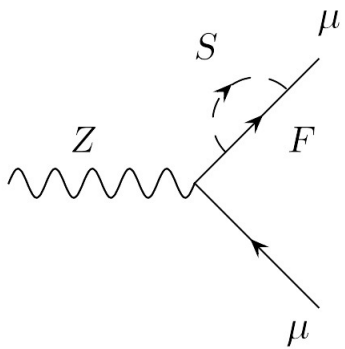
MadGraph, PYTHIA, DELPHES3



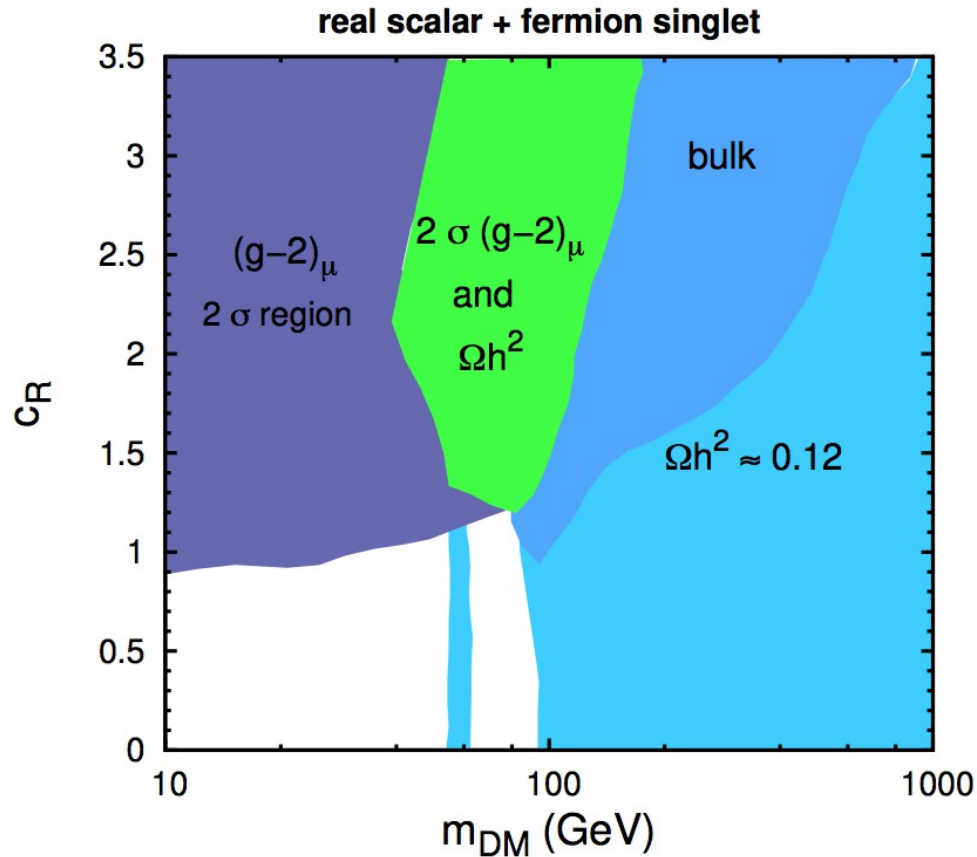
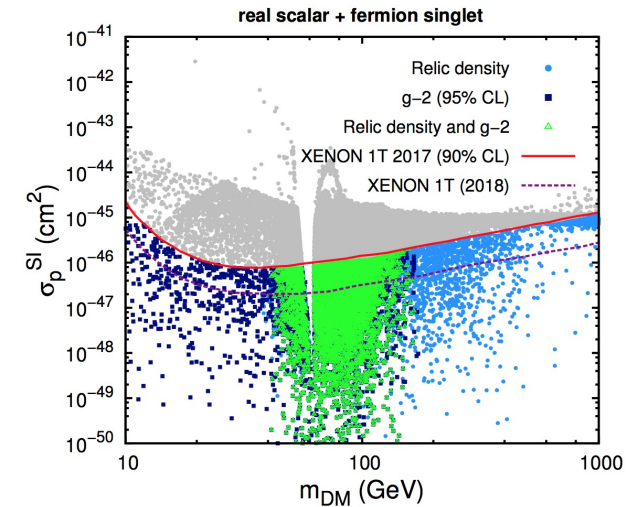
DIRECT DM SEARCHES:



PRECISION TESTS at LEP:



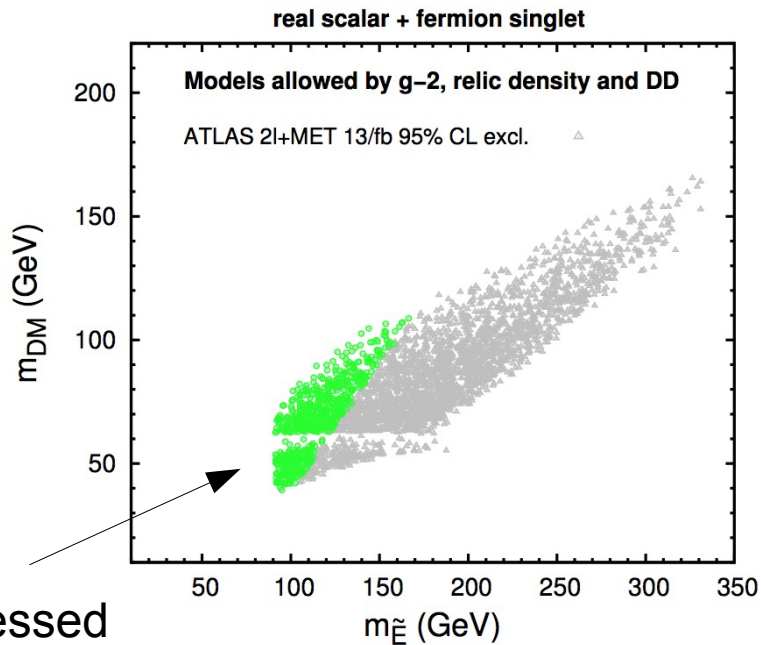
Model 1 results



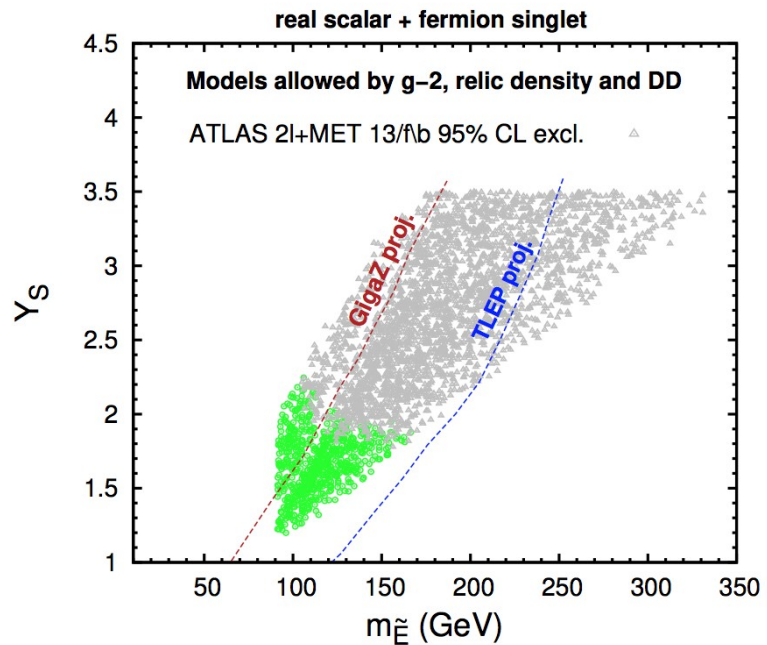
- **New Yukawa $Y_s > 1 - 1.5$**
- **m_{DM} range: 40 GeV – 180 GeV**
- **DM dominant mechanism: t -channel fermion exchange**

Model 1 results

AFTER LHC CONSTRAINTS:

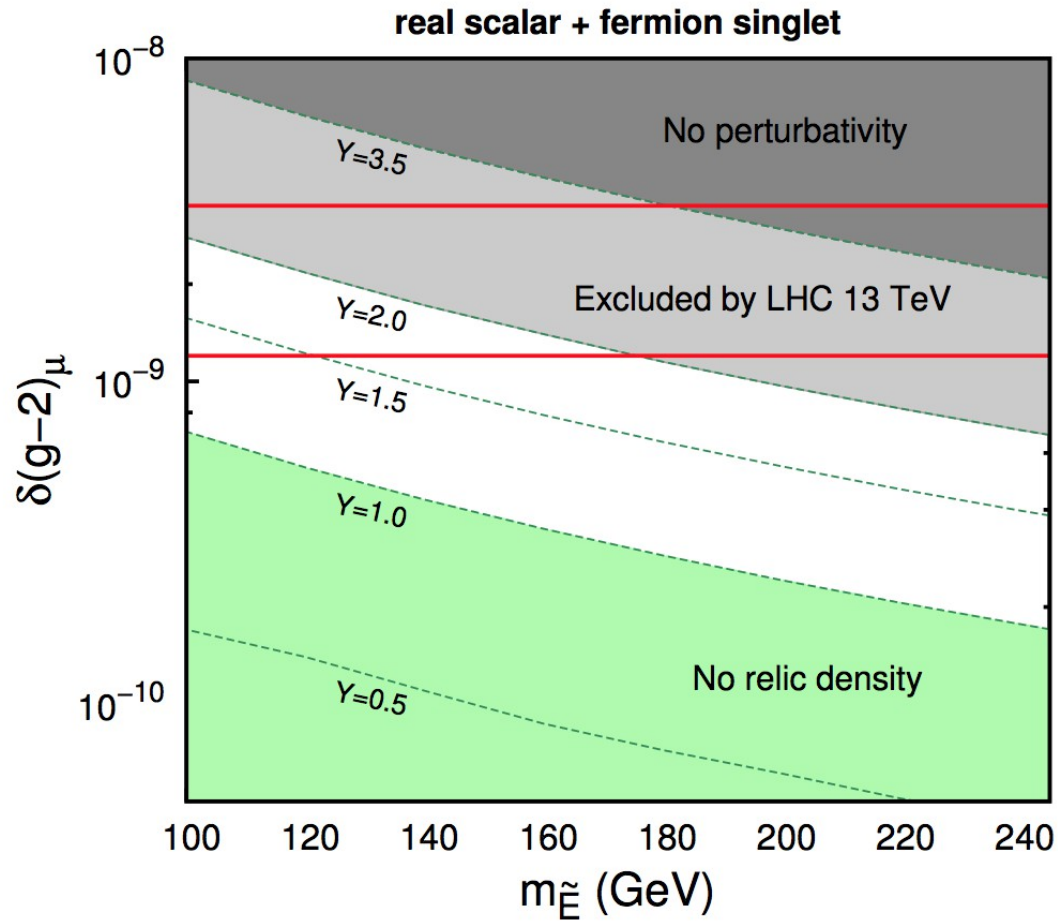


Compressed
Region survives



**Surviving parameter space tested by
future precision experiments GigaZ, TLEP**

Model 1 results

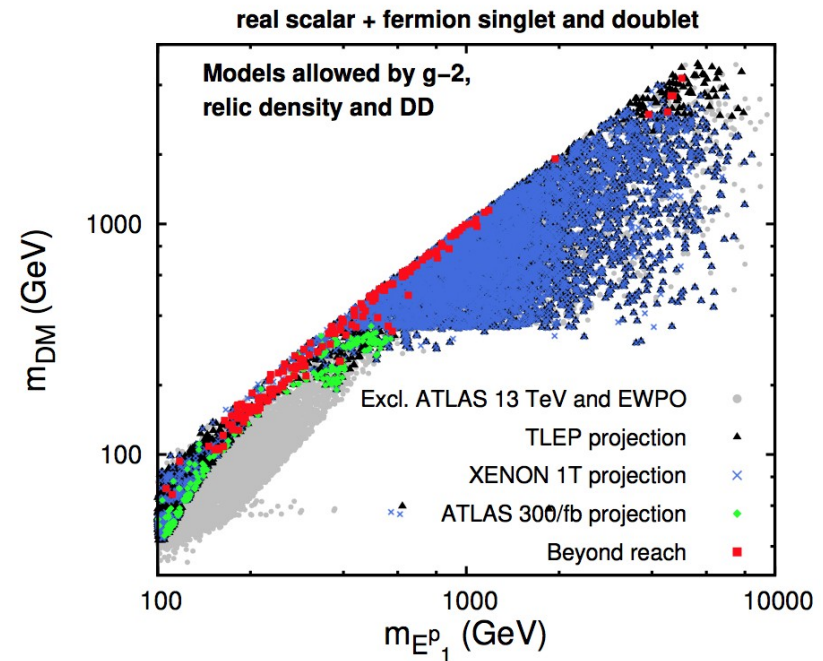
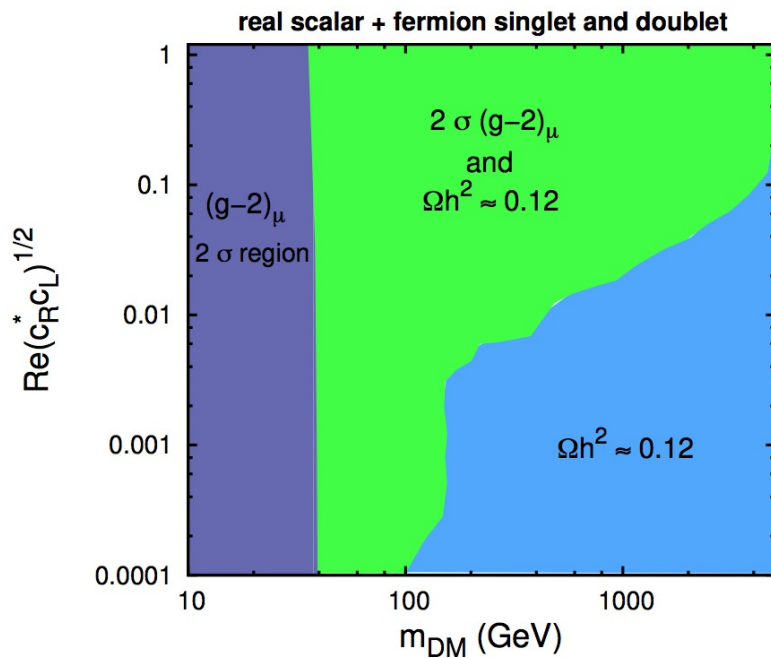


Predictive: interplay of bounds from LHC and the relic density

Mixing singlet/doublet fermions

$$\delta(g-2)_\mu = \frac{1}{16\pi^2} \sum_{S^0, E^\pm} \left[\frac{m_\mu^2}{m_S^2} (|c_L|^2 + |c_R|^2) \mathcal{F}_1(r) + 2 \frac{m_\mu m_E}{m_S^2} \Re(c_L c_R^*) \mathcal{F}_2(r) \right]$$

g-2 boost



- Large mass range allowed
- Tested by combination of LHC, PRECISION, DM DIRECT DETECTION

Scalar SU(2) doublet cases

Each case differs:

F singlet:

$$\mathcal{L} \supset -Y_S \left(\nu_L S^- E' + e_L \frac{S^{0*} E'}{\sqrt{2}} \right)$$

F doublet:

$$\mathcal{L} \supset -Y_D \left(N_1 S^- e_R^* + \frac{E_1 S^{0*}}{\sqrt{2}} e_R^* \right) + \text{h.c.} \quad (\text{charged scalar contribution})$$

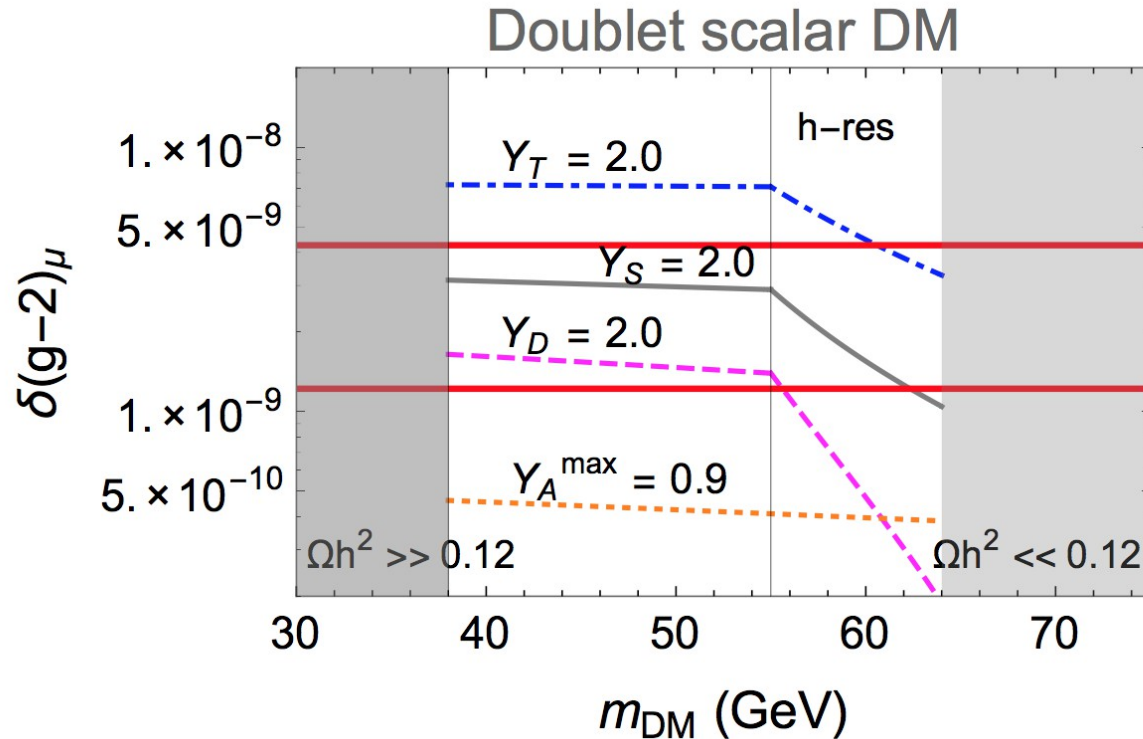
F triplet:

$$\mathcal{L} \supset -Y_T \left(\frac{\Psi^+ S^{0*}}{2} + \Psi^{++} S^- \right) e_L - Y_T \left(\frac{\Psi^0 S^{0*}}{\sqrt{2}} + \frac{\Psi^+ S^-}{\sqrt{2}} \right) \nu_L + \text{h.c.} \quad (\text{Doubly charged fermion contribution})$$

F adjoint triplet:

$$\mathcal{L} \supset -Y_A \left(\frac{\Psi^0 S^+}{\sqrt{2}} + \frac{\Psi^+ S^0}{\sqrt{2}} \right) e_L - Y_A \left(\frac{\Psi^0 S^0}{2} - \Psi^- S^+ \right) \nu_L + \text{h.c.} \quad (\text{Majorana fermions})$$

Scalar SU(2) doublet cases

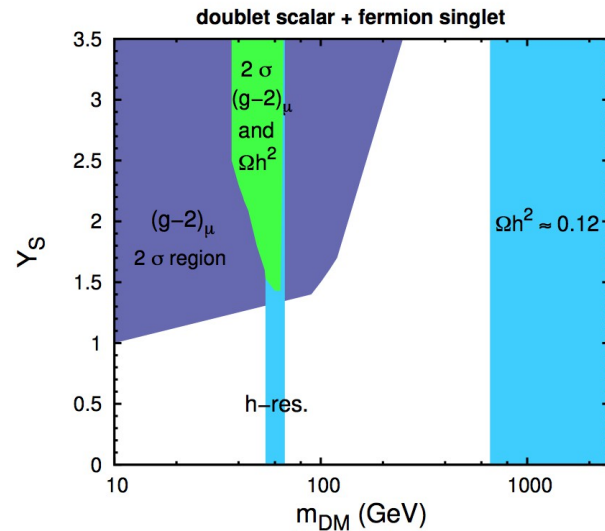


Connection g-2 / DM relic abundance:

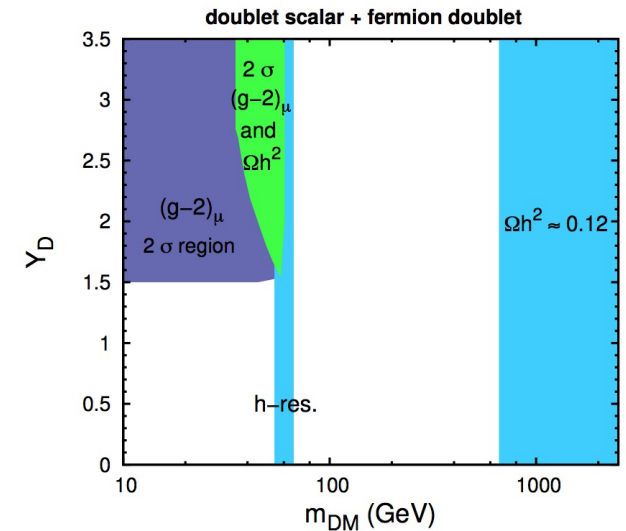
- F triplet fits BLN only in Higgs resonance
- F doublet fits only outside of Higgs resonance
- F adjoint triplet does not fit BLN

Scalar SU(2) doublet cases

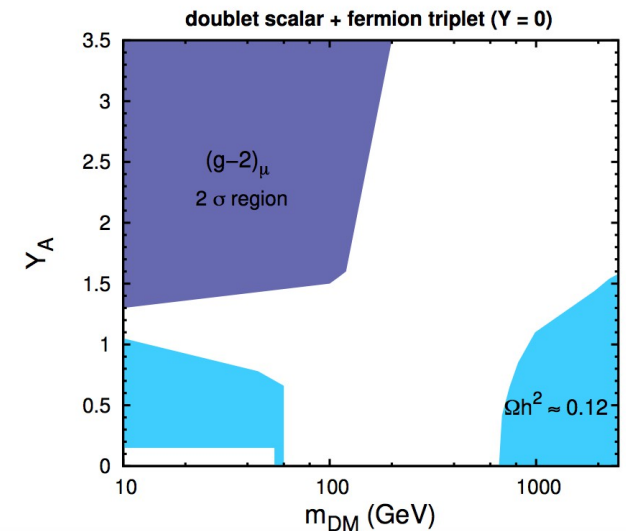
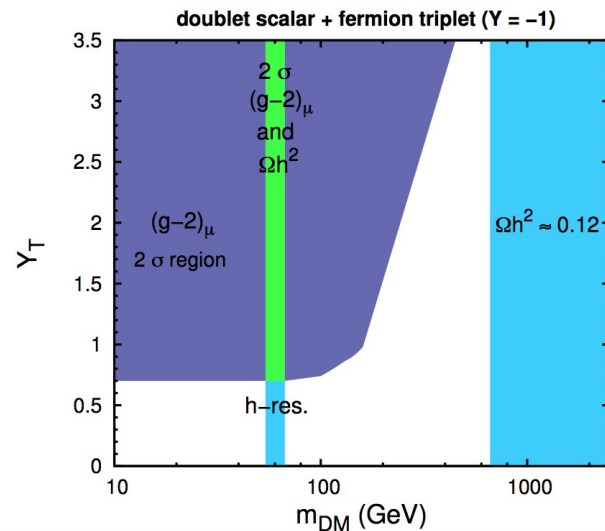
Favored parameter space differs in each case:



(a)



(b)



Summary

- **Test hypothesis that, if $g-2$ is confirmed, the same physics is subject to $\Omega h^2 \sim 0.12$.**
- **Renormalizable $SU(2) \times U(1)$ invariant extensions of the SM with different gauge structure**
- **LHC and precision tests strongly constrain the p.s., and make some of these models stand out from one another**
- **If $g-2$ confirmed, new $e^+ e^-$ collider can probe even minimal models**