

# Quasi-sterile neutrinos at long-baseline oscillation experiments

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Long-baseline neutrino experiments can search for various new particles.

This talk: a theory that predicts unusual signals in the far detector.

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New gauge symmetry acting on quarks only:  $U(1)_B$   
charges proportional to baryon number

Fields	spin	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$	$U(1)_B$
$q_L^i = (u_L^i, d_L^i)^\top$	1/2	3	2	+1/6	+1/3
$u_R^i, d_R^i$	1/2	3	1	+2/3, -1/3	+1/3

*Theoretical requirements:*

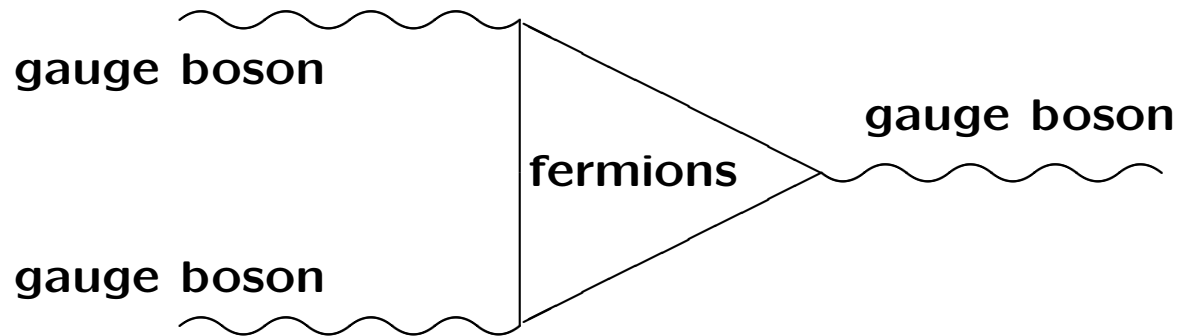
- $U(1)_B$  must be spontaneously broken.  
Simple choice: a new scalar field  $\phi$  acquires a VEV.
- All  $U(1)_B$  gauge anomalies must cancel.  
⇒ Some new fermions (“anomalons”) must be vectorlike with respect to  $SU(3)_c \times SU(2)_W \times U(1)_Y$ , and chiral with respect to the new gauge group.

# Gauge anomaly cancellation

W. Bardeen, 1969, ...

Gauge symmetries may be broken by quantum effects.

Cure: sums over fermion triangle diagrams must vanish.



*Standard Model – anomalies cancel within each fermion generation:*

$$[SU(3)_c]^2 U(1)_Y: \quad 2(1/6) + (-2/3) + (1/3) = 0$$

$$[SU(2)_W]^2 U(1)_Y: \quad 3(1/6) + (-1/2) = 0$$

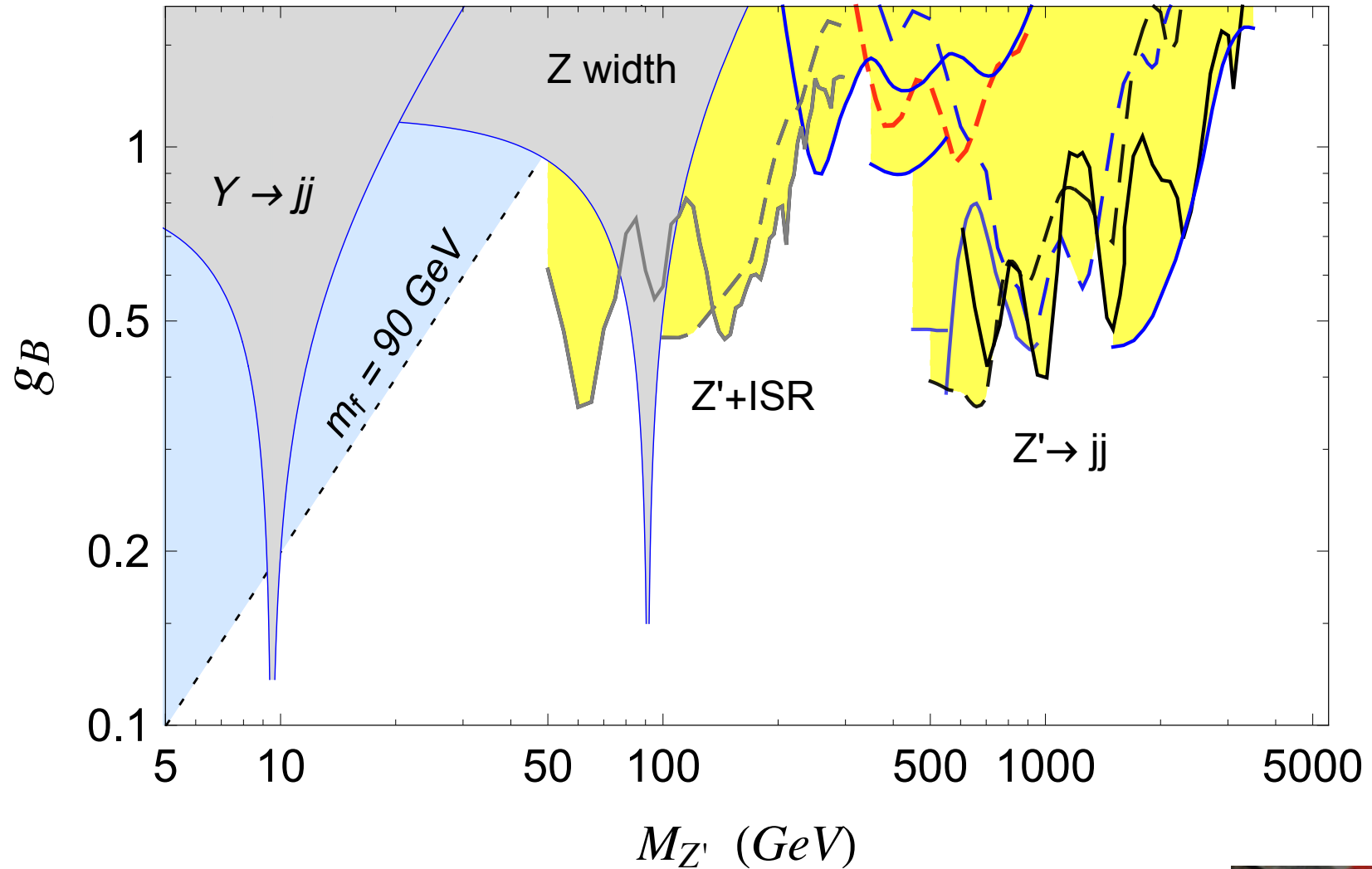
$$[U(1)_Y]^3: \quad 3 \left[ 2(1/6)^3 + (-2/3)^3 + (1/3)^3 \right] + 2(-1/2)^3 + (-1)^3 = 0$$

...  $(u_L, d_L)$   $u_R$   $d_R$   $(\nu_L, e_L)$   $e_R$

## A complete, renormalizable model:

Fields	spin	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$	$U(1)_B$
$q_L^i = (u_L^i, d_L^i)^\top$	1/2	3	2	+1/6	+1/3
$u_R^i, d_R^i$	1/2	3	1	+2/3, -1/3	+1/3
$L_L = (L_L^N, L_L^E)^\top$	1/2	1	2	-1/2	-3/2
$L_R = (L_R^N, L_R^E)^\top$	1/2	1	2	-1/2	+3/2
$E_L$	1/2	1	1	-1	+3/2
$E_R$	1/2	1	1	-1	-3/2
$N_L$	1/2	1	1	0	+3/2
$N_{*R}$	1/2	1	1	0	-3/2
$N_{0R}$	1/2	1	1	0	0
$\phi$	0	1	1	0	+3
$\phi'$	0	1	1	0	+3/2

“Baryonic”  $Z'_B$ : same coupling ( $g_B$ ) to all six quark flavors.



$$\mathcal{L}_q = \frac{g_B}{2} Z'_\mu \sum_q \left( \frac{1}{3} \bar{q}_L \gamma^\mu q_L + \frac{1}{3} \bar{q}_R \gamma^\mu q_R \right)$$

with Felix Yu:  
update to 1306.2629

Limits on anomalon masses impose the constraint  
on  $g_B$  at low mass B.A. Dobrescu, C. Frugiuele, 1404.3947



**Yukawa interactions of neutral fermions:**

$$- (\lambda_N \langle \phi \rangle \bar{N}_{*R} + \lambda'_N \langle \phi' \rangle \bar{N}_{0R}) N_L + \text{H.c.}$$

**This leads to a Dirac mass for a fermion  $N$  with**

$$N_R \equiv N_{*R} \cos \alpha + N_{0R} \sin \alpha$$

**The state orthogonal to  $N_R$ ,**

$$\nu_R \equiv -N_{*R} \sin \alpha + N_{0R} \cos \alpha$$

**remains a massless Weyl fermion.**

**This is a “quasi-sterile” neutrino because its interactions are mediated only by the  $Z'_B$  boson.**

If the SM neutrinos acquire Majorana masses as usual from dimension-5 operators

$$\frac{c_{ij}}{M} H H \bar{\ell}_L^i \ell_L^j ,$$

then the dimension-5 operator

$$\frac{c'_i}{M} \phi' H \bar{\ell}_L^i N_{*R}$$

is also likely to be present. This leads to mixing between SM neutrinos and the  $\nu_R$  and  $N$  fermions.

The mixing of the SM  $\nu_\mu$  with the massless  $\nu_R$  is restricted by various measurements of neutrino properties:  $\theta \lesssim O(0.2)$

For the future DUNE experiment,  $L \approx 1300$  km and  $E_\nu$  is typically in the 1–4 GeV range, so that  $P(\nu_\mu \rightarrow \nu_R) \approx \theta^2 \sim O(10^{-2})$ .

The quasi-sterile neutrino  $\nu_R$  and the heavy Dirac state  $N$  have the following interactions with the  $Z'_B$  boson:

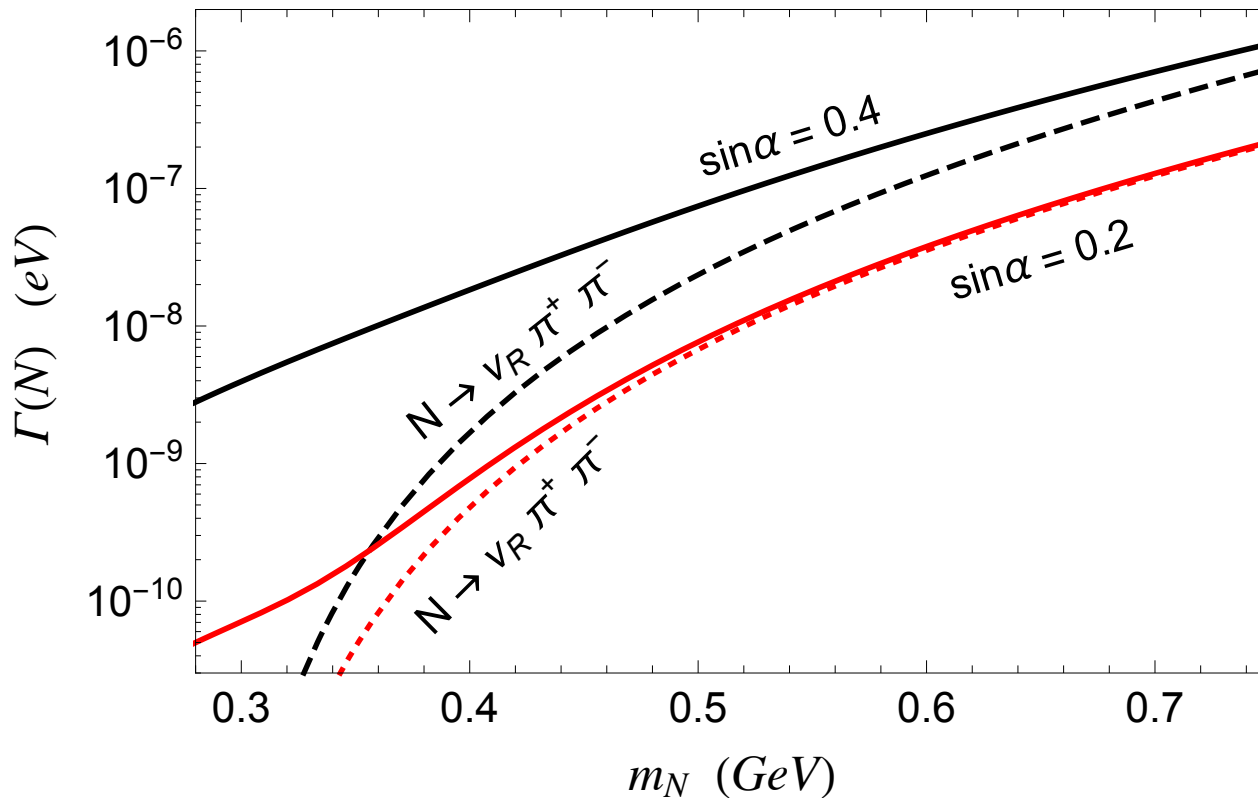
$$-\frac{3g_B}{4} Z'_{B\mu} \left[ -\sin\alpha \cos\alpha (\bar{N}_R \gamma^\mu \nu_R + \text{H.c}) \right. \\ \left. + \sin^2\alpha \bar{\nu}_R \gamma^\mu \nu_R + \cos^2\alpha \bar{N}_R \gamma^\mu N_R \right]$$

The first term here allows the up-scattering of  $\nu_R$  in  $N_R$ , when the energy of the incoming neutrino is large enough.

$\sim 100$   $N$  fermions may be produced inside DUNE.



Total decay width,  $\Gamma(N)$ , of the  $N$  fermion (solid lines) and partial width for the  $N \rightarrow \nu_R \pi^+ \pi^-$  decay (dashed lines):

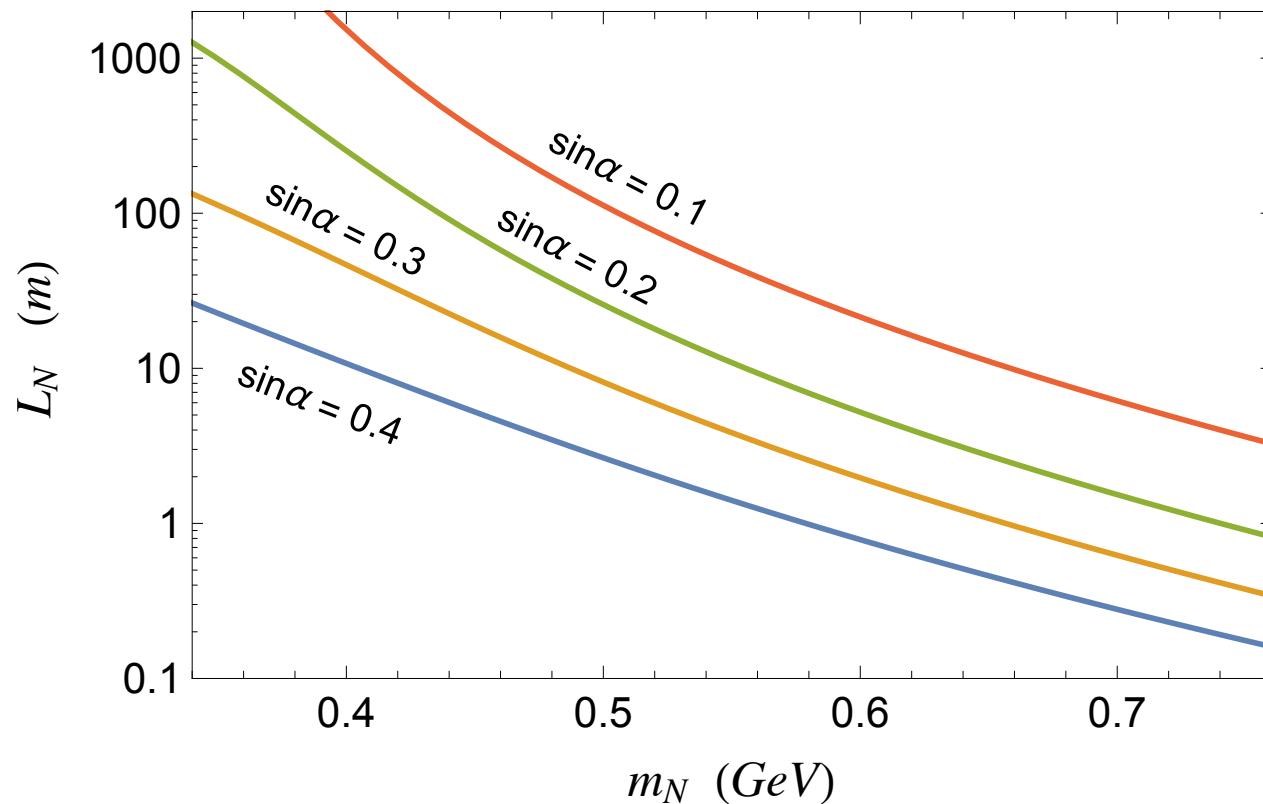


The difference between the total width and the  $N \rightarrow \nu_R \pi^+ \pi^-$  width is due to the invisible  $N \rightarrow 3\nu_R$  decay.

Parameters used here are  $M_{Z'} = 50$  GeV,  $g_B = 0.5$ .

Decay length of the  $N$  fermion in the rest frame,  $L_N$ , is proportional to  $M_{Z'}/g_B$ ,

For  $M_{Z'} = 50$  GeV, and  $g_B = 0.5$ :



## Conclusions

- Quasi-sterile neutrinos have interactions only through nonstandard bosons.
- A quasi-sterile neutrino may be produced in long-baseline neutrino oscillation experiments.
- Interaction of a quasi-sterile neutrino  $\nu_R$  with a  $Z'$  boson and a new fermion  $N$  of GeV-scale mass may lead to  $N$  production in the far detector.
- $N \rightarrow \pi^+ \pi^- \nu_R$  decay would give a spectacular signal.

**DUNE will probe the laws of nature in a new regime → many searches for new particles will be possible.**