

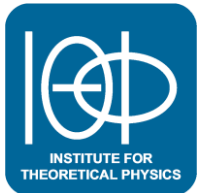
Broken boost invariance in the Glasma via finite nuclei thickness

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FWF

Der Wissenschaftsfonds.



Introduction

Various heavy-ion collision experiments:

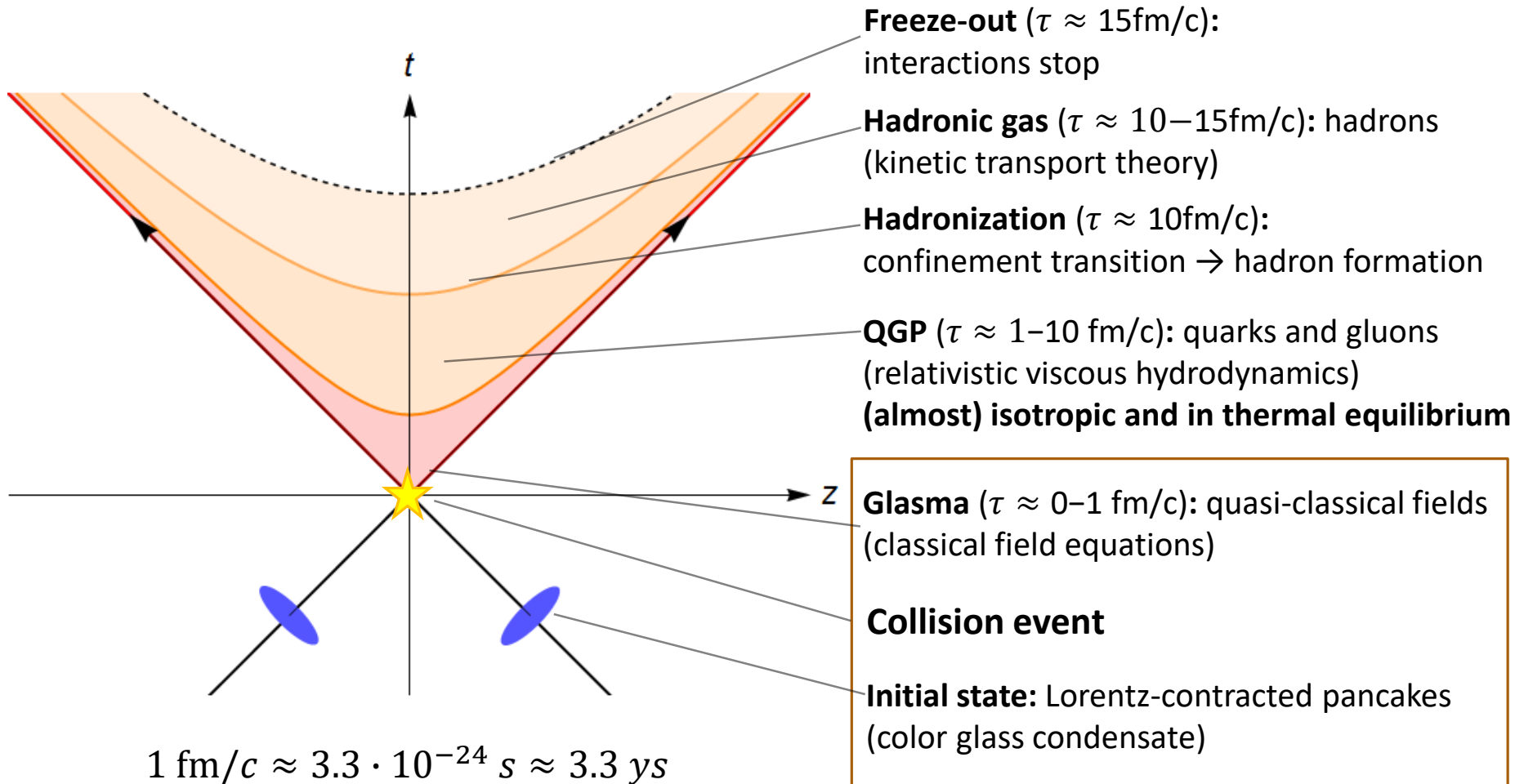
- LHC (ALICE) @ CERN: Pb+Pb with **~5.5 TeV** per nucleon pair. ($\gamma \approx 2700$)
- RHIC @ BNL: Au+Au with **~200 GeV** per nucleon pair. ($\gamma \approx 100$)
- RHIC beam energy scan: **~7.7 – 62.4 GeV** ($\gamma \approx 4 – 30$)

Early stages (few fm/c):

- Color glass condensate / Glasma: classical fields
- Finite longitudinal extent of the nuclei?

Goal: Simulate heavy-ion collisions in the color glass condensate (CGC) framework with finite nucleus thickness. Possible with colored particle-in-cell (CPIC) method.

Stages of a heavy-ion collision



scope of this project

Color glass condensate

Nuclei at ultrarelativistic speeds can be described by **classical effective theory** in the color glass condensate (CGC) framework.

[Gelis, Iancu, Jalilian-Marian, Venugopalan, Ann.Rev.Nucl.Part.Sci.60:463-489,2010]

Large gluon occupation numbers \rightarrow coherent, classical gluon field

Split degrees of freedom into ...

- Hard partons = classical color charges
 - Soft gluons = classical gauge field
- generates..
- Static field configuration due to time dilation.
 - Collision of two such classical fields creates the **Glasma**.

[Gelis, Int.J.Mod.Phys. A28 (2013) 1330001]

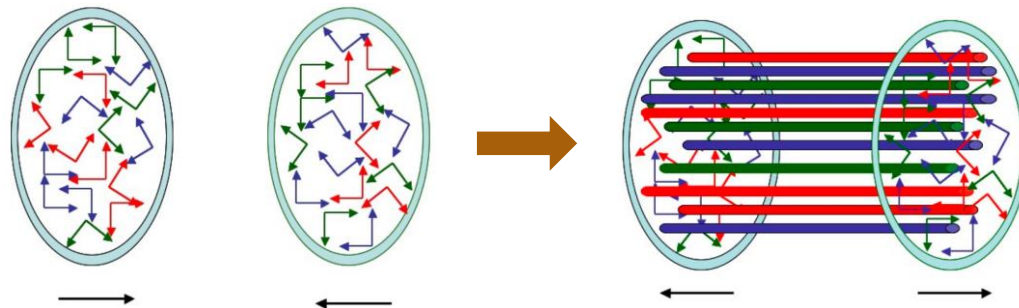
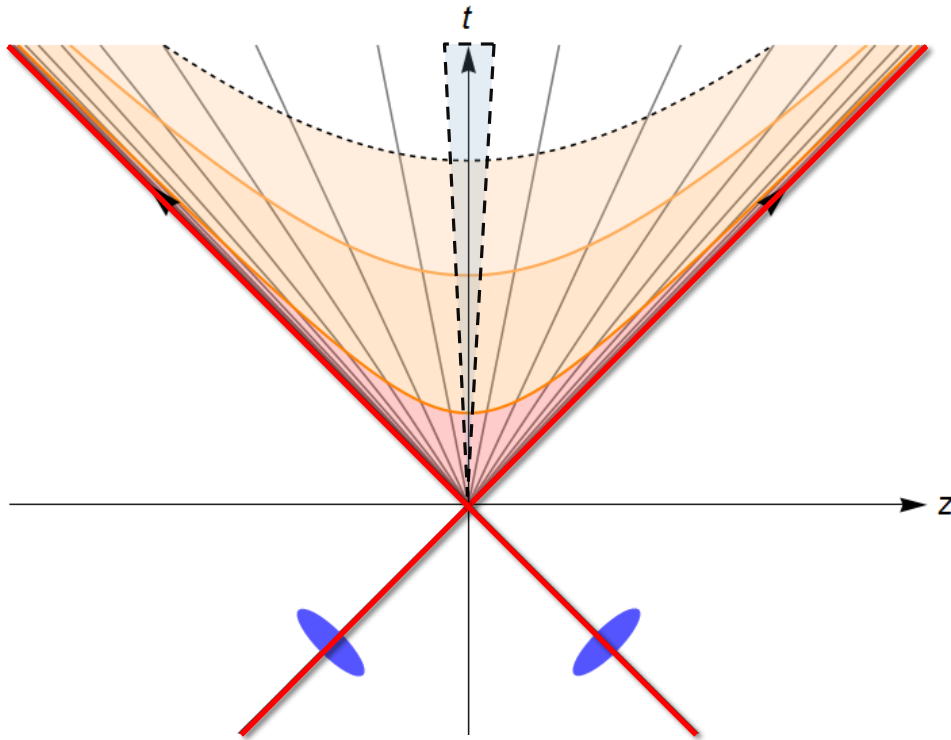


Figure from L. McLerran:
Proceedings of ISMD08, p.3-18 (2008)

Boost-invariant collision

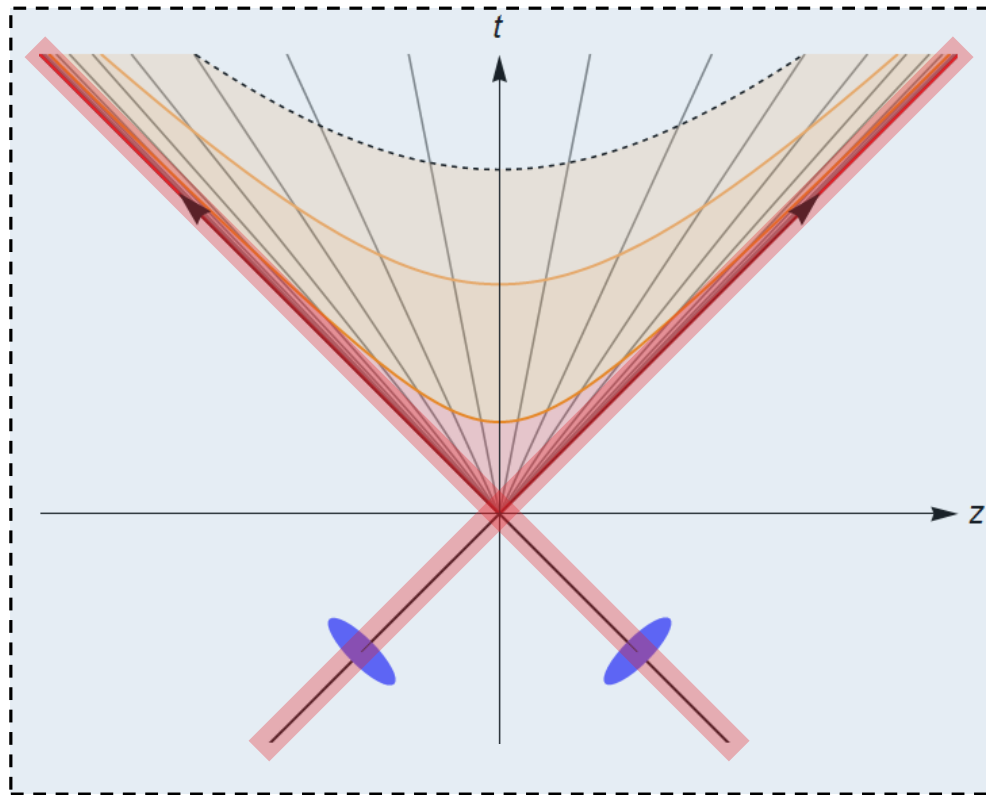


- Color currents of the nuclei restricted to the light cone and infinitely thin
- Point-like collision region
- Analytical solutions exist for everything except the forward light cone. Glasma initial conditions on the boundary.
- Fields in the forward light cone are independent of rapidity η . Reduction from 3+1 to 2+1.
- Need to solve **2D+1 source-free Yang-Mills** equations in the forward light cone with Glasma initial conditions on the boundary of the light cone:

$$D_\mu F^{\mu\nu}(\tau, x_T) = 0$$

Description in τ, η (comoving) frame

Broken boost invariance



Description in t, z (laboratory) frame

- Extended color currents
- Extended collision region
- Fields depend on rapidity
- Need to solve **full 3D+1 Yang-Mills equations with currents!**

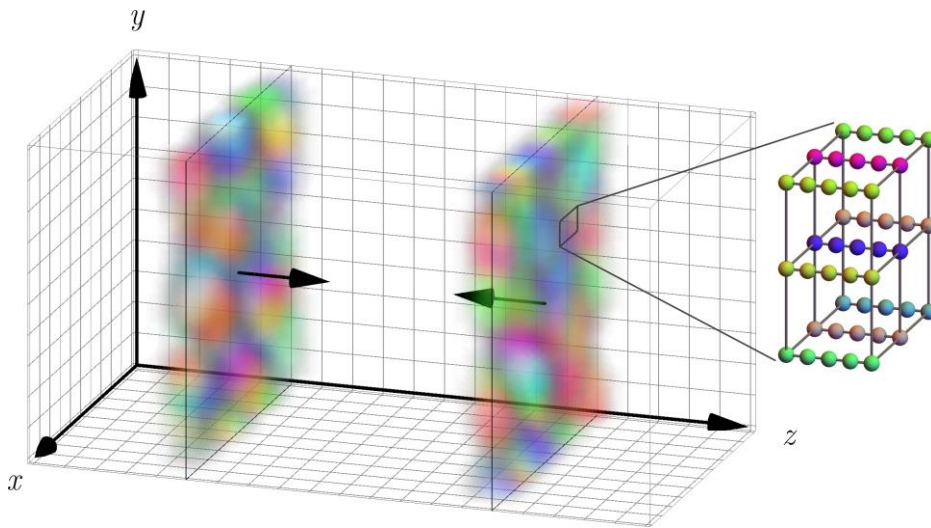
$$\begin{aligned} D_\mu F^{\mu\nu}(t, z, x_T) &= J^\nu \\ D_\mu J^\mu(t, z, x_T) &= 0 \end{aligned}$$

Colored particle-in-cell (CPIC) provides a framework to numerically solve the field and current equations on a lattice.

Colored particle-in-cell (CPIC) method

Generalization of the particle-in-cell (PIC) method for Abelian plasmas to non-Abelian gauge theories.

[A. Dumitru, Y. Nara, M. Strickland: Phys.Rev.D75:025016 (2007)]



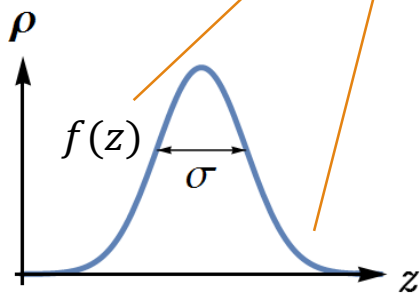
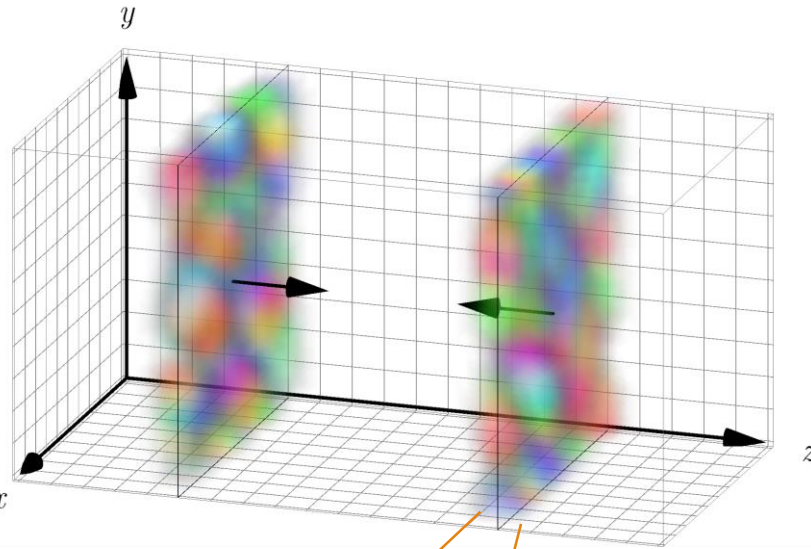
Heavy-ion collision in the lab frame with CPIC:

- Fields: discretized on a lattice
- **Color charge densities: replace with a large number of colored particles.**
- Current J_μ on the grid generated by particle movement \rightarrow field equations

$$\begin{aligned} D_\mu F^{\mu\nu}(t, z, x_T) &= J^\nu \\ D_\mu J^\mu(t, z, x_T) &= 0 \end{aligned}$$

Non-Abelian gauge theory \rightarrow parallel transport of charges

Initial conditions



longitudinal profile function $f(z)$ (Gaussian)

Color charges in the nuclei are randomly distributed

Models needed to fix the exact distribution

Simple model: **2D McLerran-Venugopalan (MV)**

[McLerran, Venugopalan: PRDD49 (1994) 3352-3355]

- Infinite transverse extent (periodic b.c.)
- **New:** embed 2D density into 3D space

$$\langle \hat{\rho}^a(\mathbf{x}_T) \hat{\rho}^b(\mathbf{x}'_T) \rangle = g^2 \mu^2 \delta^{(2)}(\mathbf{x}_T - \mathbf{x}'_T) \delta^{ab}$$

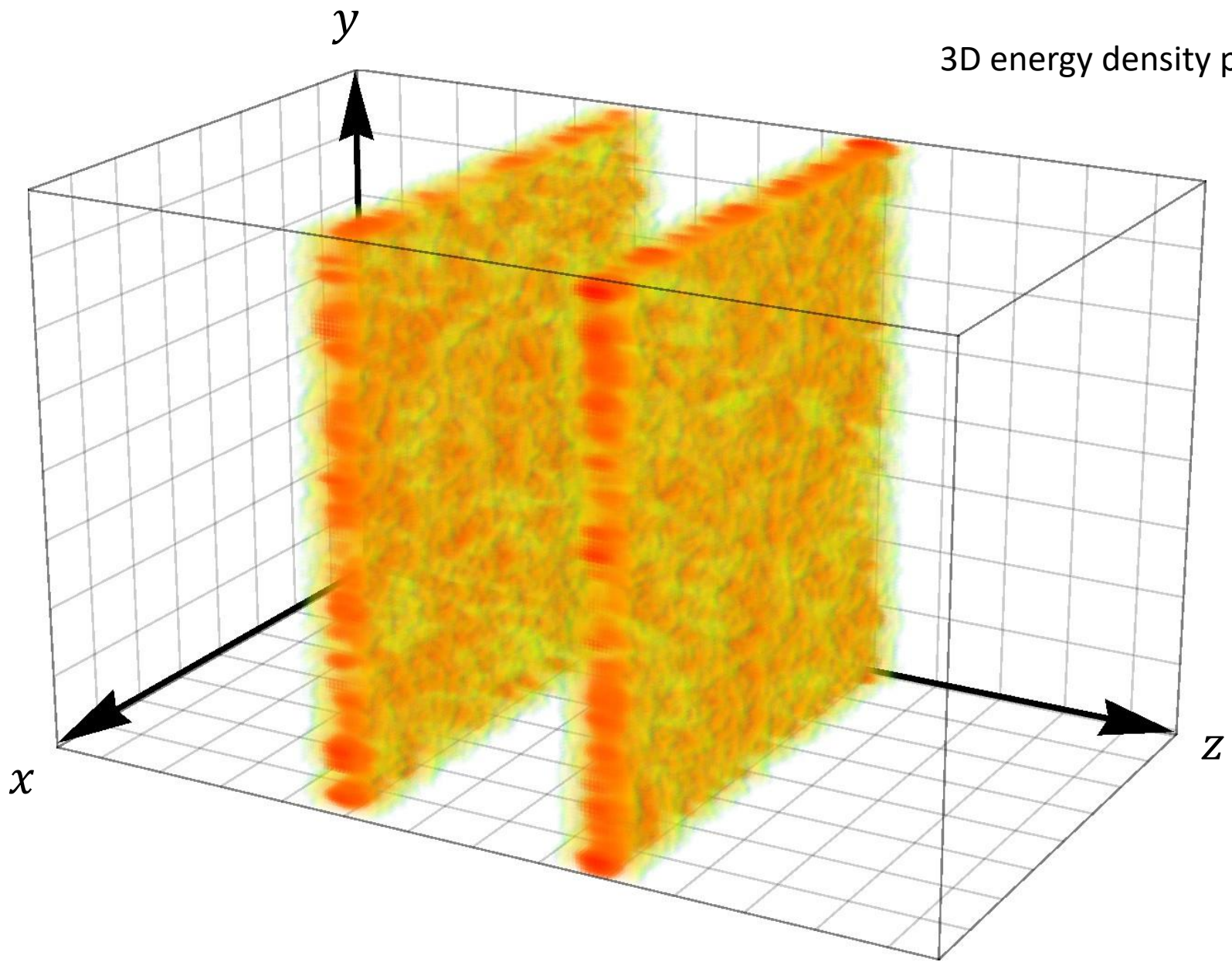
$$\rho(t, z, \mathbf{x}_T) = f(z - t) \hat{\rho}(\mathbf{x}_T)$$

- “Trivial” longitudinal structure
- Regulation of IR and UV modes required

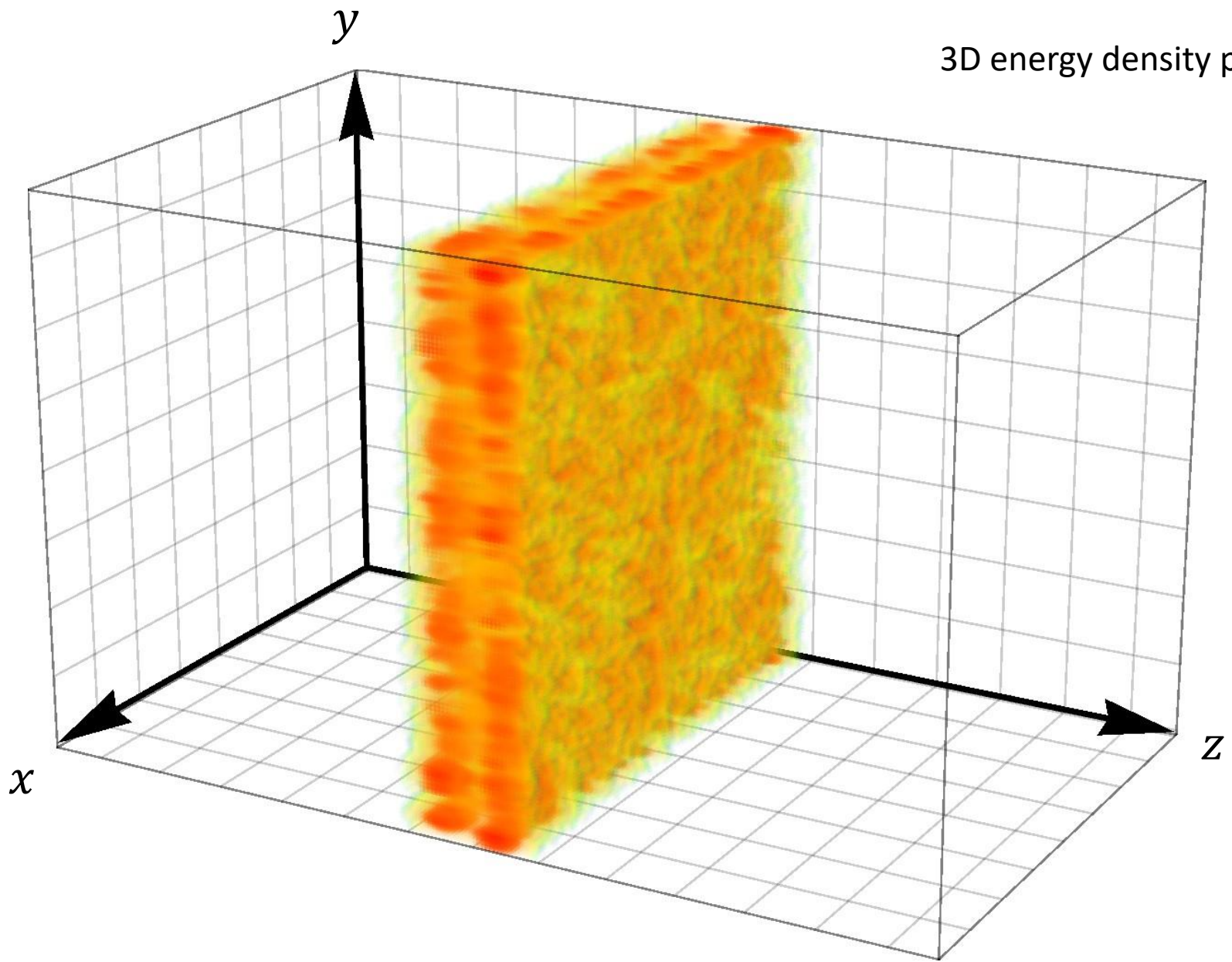
MV parameter $\mu \approx 0.5 \text{ GeV (Au)}$

Numerical results

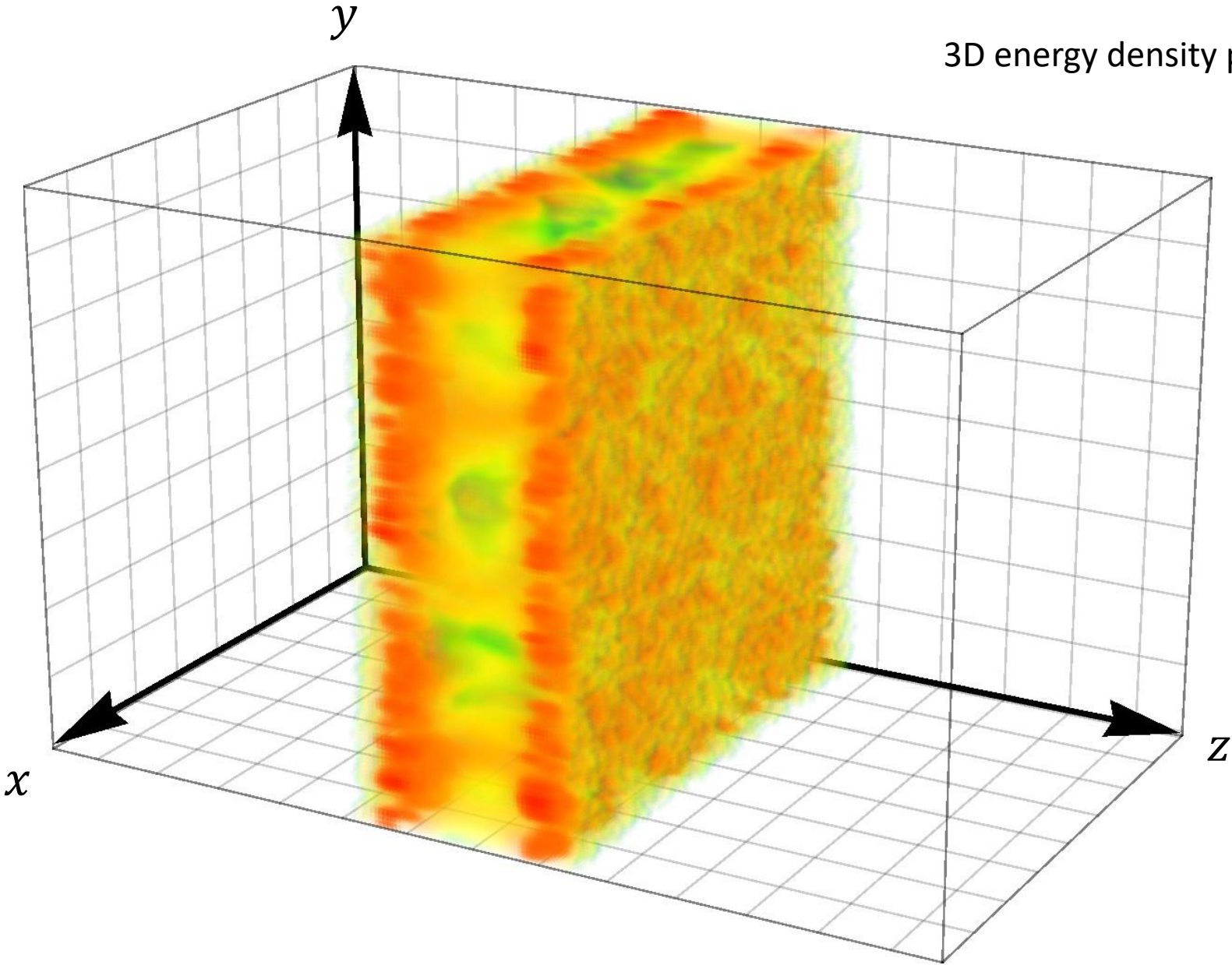
3D energy density plot



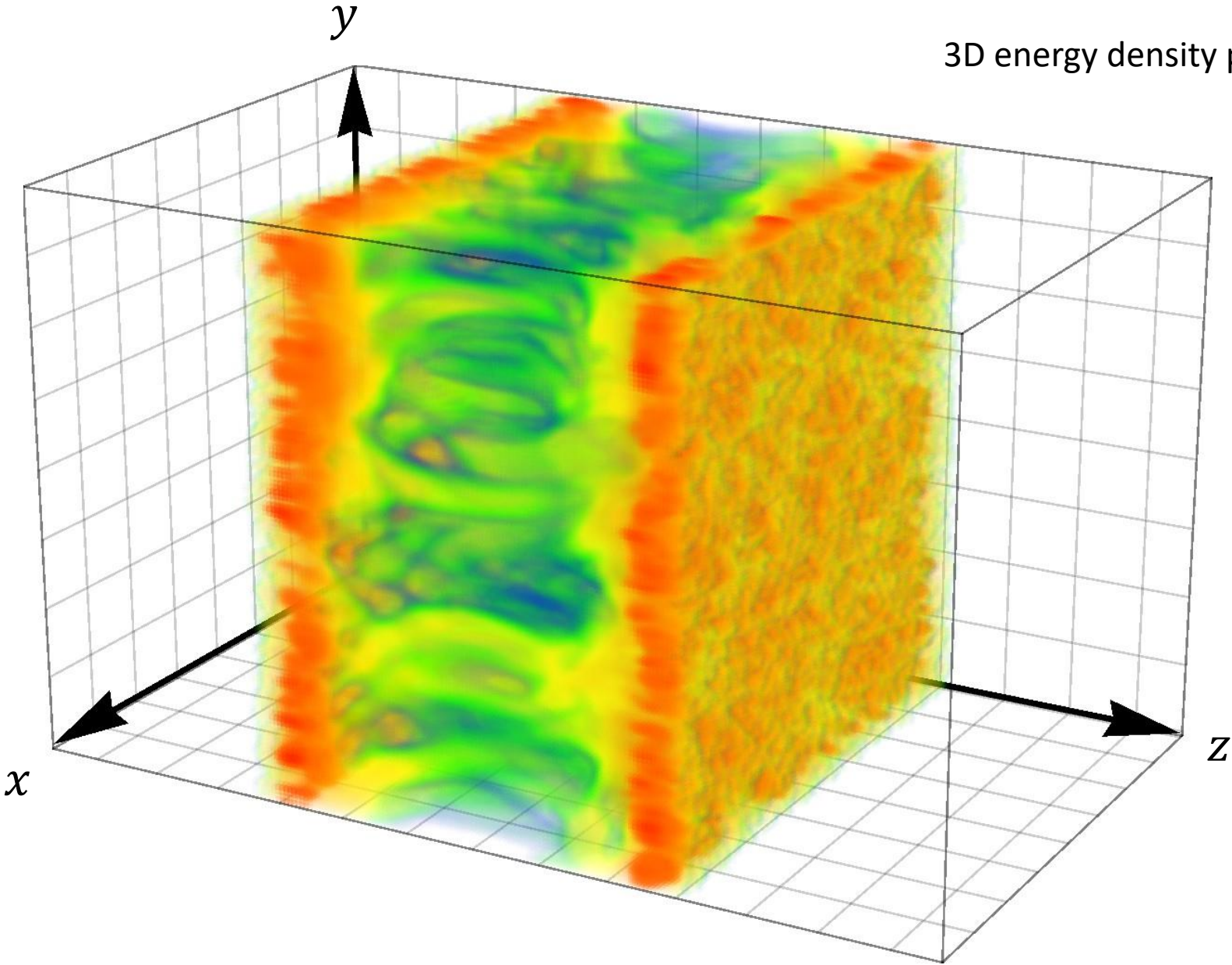
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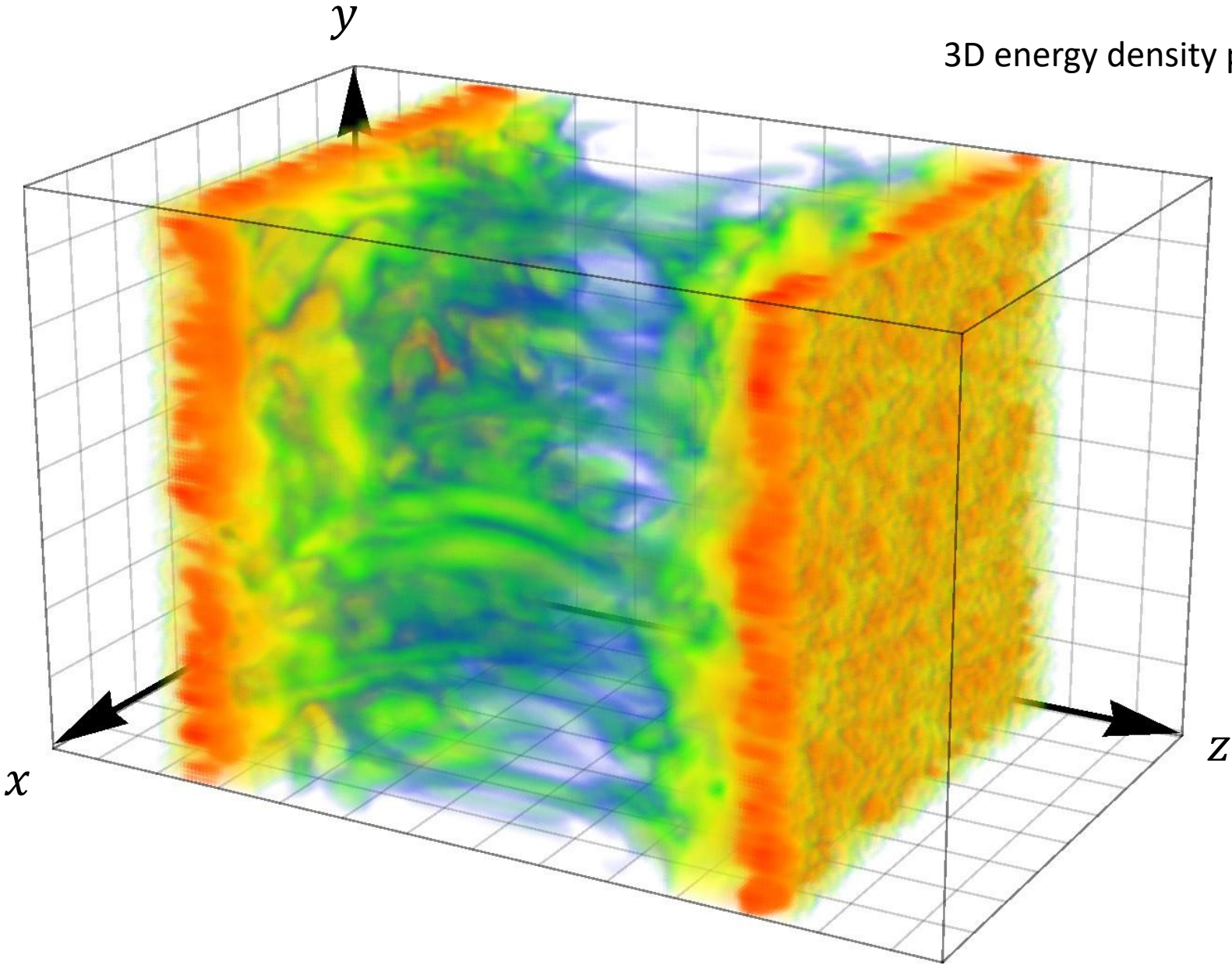
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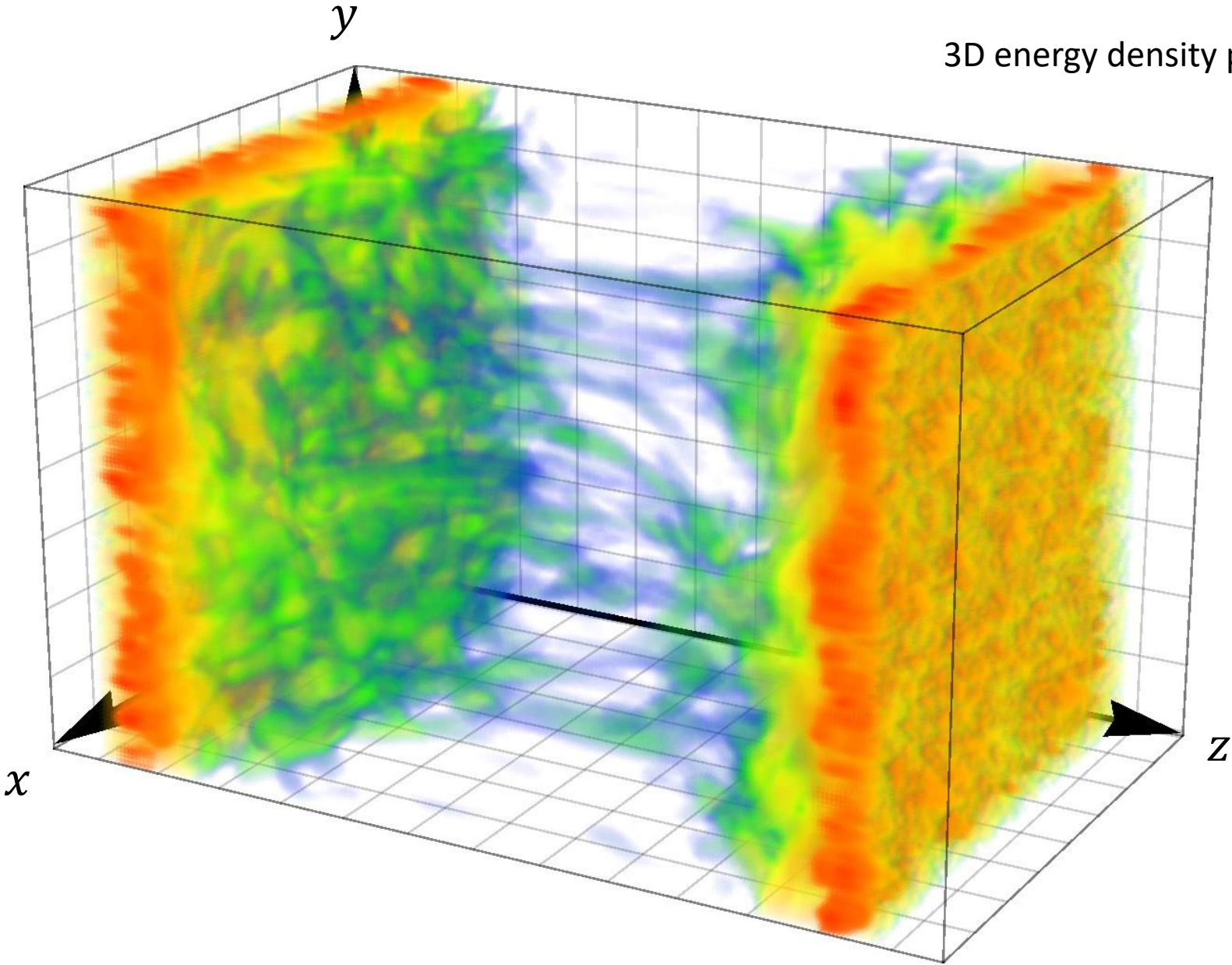
3D energy density plot



3D energy density plot



3D energy density plot



Observables

Main observable: energy-momentum tensor $T^{\mu\nu}(x)$

- Compute $T^{\mu\nu}(x)$ from electric and magnetic fields $E_i^a(x), B_i^a(x)$
- Average over configurations and integrate over transverse plane

$$\langle T^{\mu\nu} \rangle = \begin{pmatrix} \langle \varepsilon \rangle & 0 & 0 & \langle S_L \rangle \\ 0 & \langle p_T \rangle & 0 & 0 \\ 0 & 0 & \langle p_T \rangle & 0 \\ \langle S_L \rangle & 0 & 0 & \langle p_L \rangle \end{pmatrix}$$
$$\begin{aligned} \langle \varepsilon \rangle &= \frac{1}{2} \langle E_T^2 + B_T^2 + E_L^2 + B_L^2 \rangle \\ \langle p_T \rangle &= \frac{1}{2} \langle E_L^2 + B_L^2 \rangle \\ \langle p_L \rangle &= \frac{1}{2} \langle E_T^2 + B_T^2 - E_L^2 - B_L^2 \rangle \\ \langle S_L \rangle &= \left\langle (\vec{E}^a \times \vec{B}^a)_L \right\rangle \end{aligned}$$

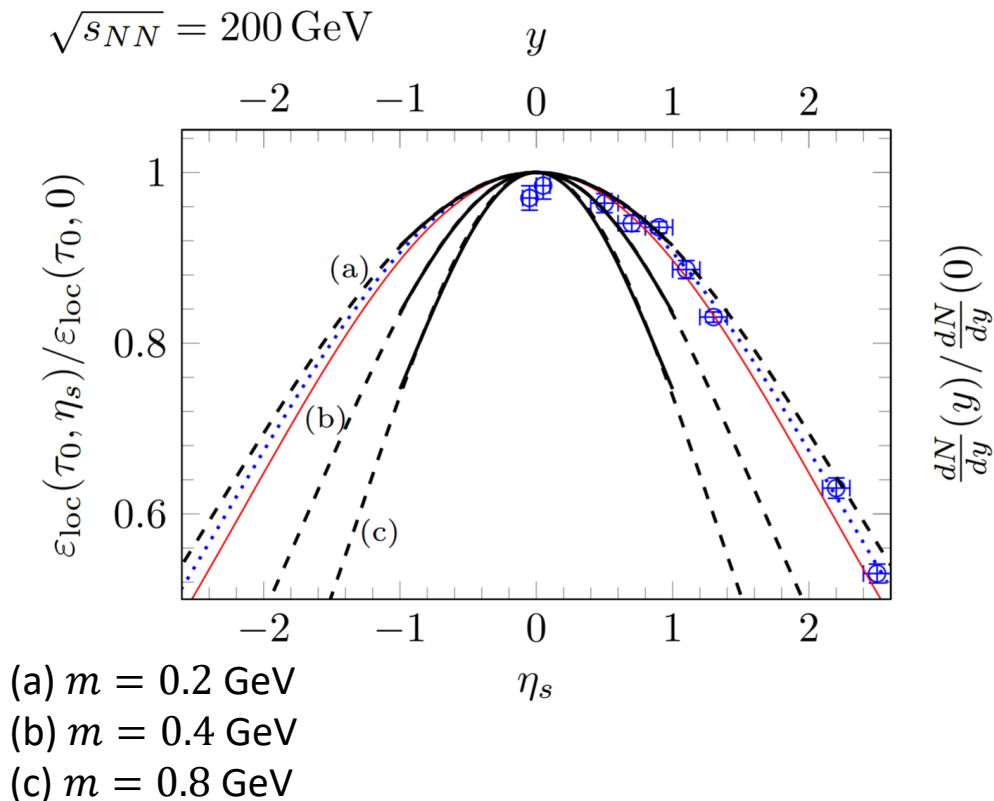
- Diagonalize, obtain local rest-frame energy density

$$\langle \varepsilon_{\text{loc}} \rangle = \frac{1}{2} \left(\langle \varepsilon \rangle - \langle p_L \rangle + \sqrt{(\langle \varepsilon \rangle + \langle p_L \rangle)^2 - 4 \langle S_L \rangle^2} \right)$$

Rapidity profiles

Plot (space-time) rapidity profile of local rest-frame energy density

Compare to measured **rapidity profile of particle multiplicity (RHIC)** and **Landau model** prediction

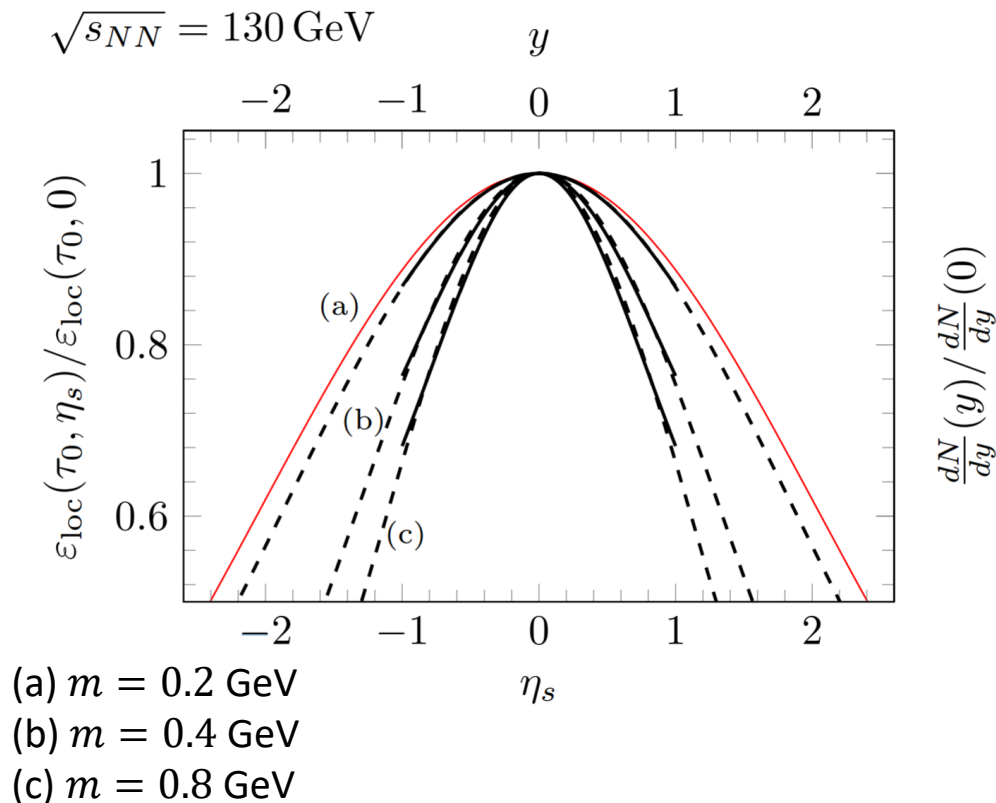


- Simulation data in interval $\eta_s \in (-1, 1)$ at $\tau = 1 \text{ fm/c}$
- Fit to Gaussian profile (dashed)
- Dependency on thickness (or rather \sqrt{s})
- Strong dependency on IR regulator, but $m = 0.2 \text{ GeV}$ gives realistic shape
- However: no hydrodynamic expansion included
- Limited rapidity interval

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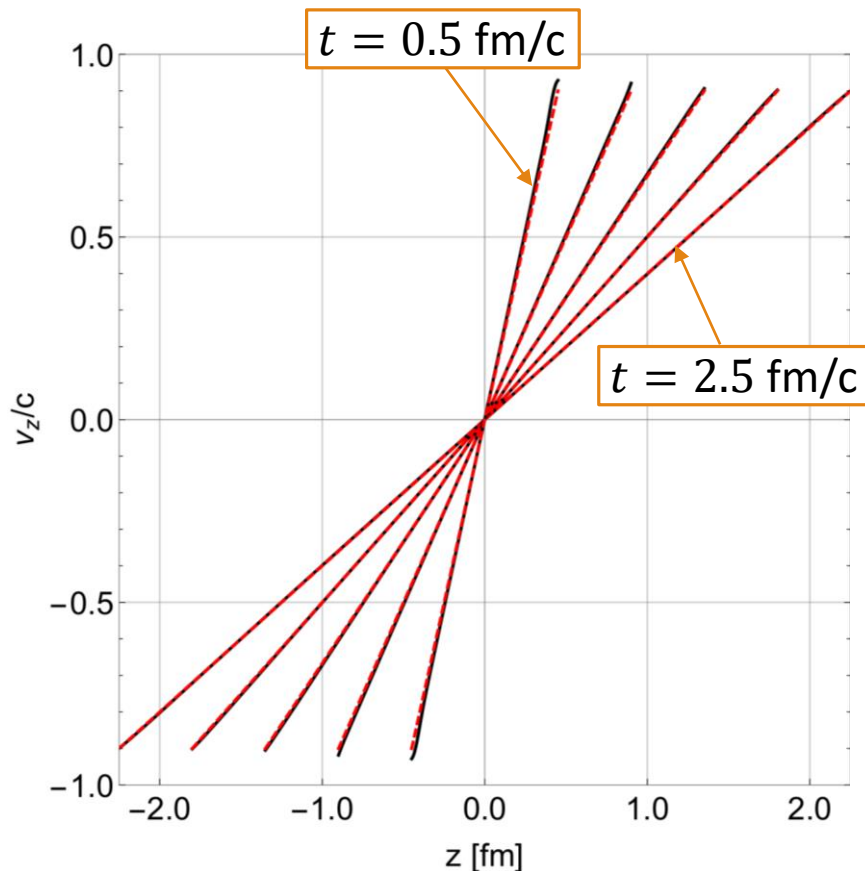


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Longitudinal velocity v_z

How fast is the Glasma flowing in the longitudinal direction?

Is the use of the τ, η frame justified?



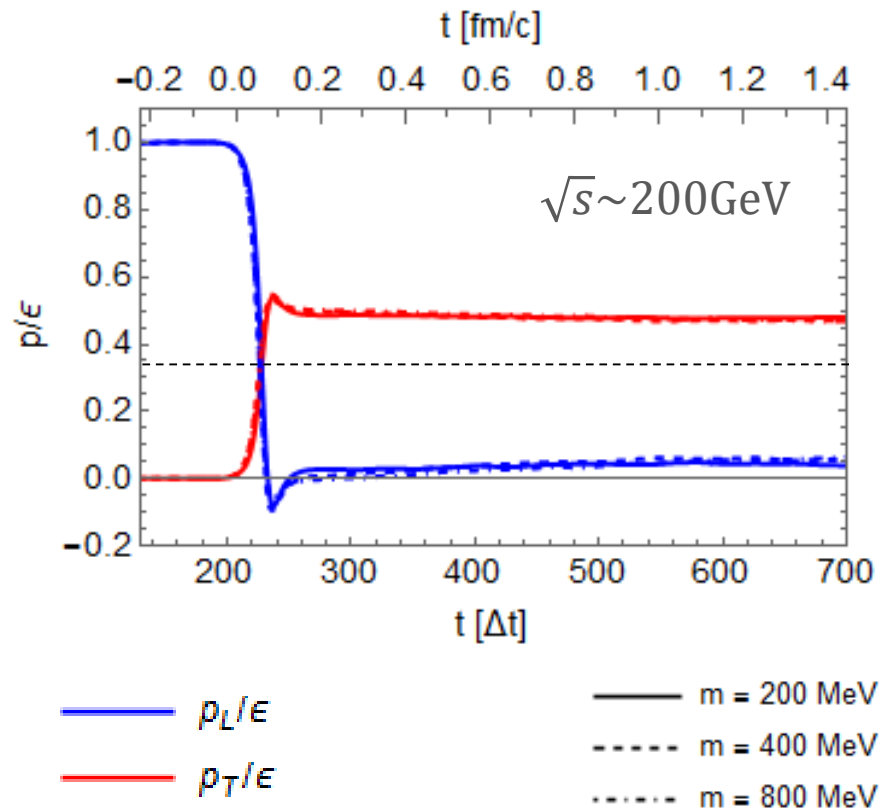
- Compute v_z from energy-momentum tensor (**black**)
- Compare to free-streaming case $v_z = z/t$ (**red, dashed**)
- Glasma fields depend on η , but still free-streaming!
- Local rest-frame $\cong \tau, \eta$ frame

Pressure anisotropy

Look at **pressure to energy density ratio** at mid-rapidity to investigate pressure anisotropy of the Glasma fields

$$\frac{\langle p_L \rangle}{\langle \epsilon \rangle}, \frac{\langle p_T \rangle}{\langle \epsilon \rangle} \text{ at } \eta = 0$$

- Strong anisotropy, no isotropization
- Similar to boost-invariant case
- No dependence on IR regulator m
- Initially negative longitudinal pressure is “hidden”

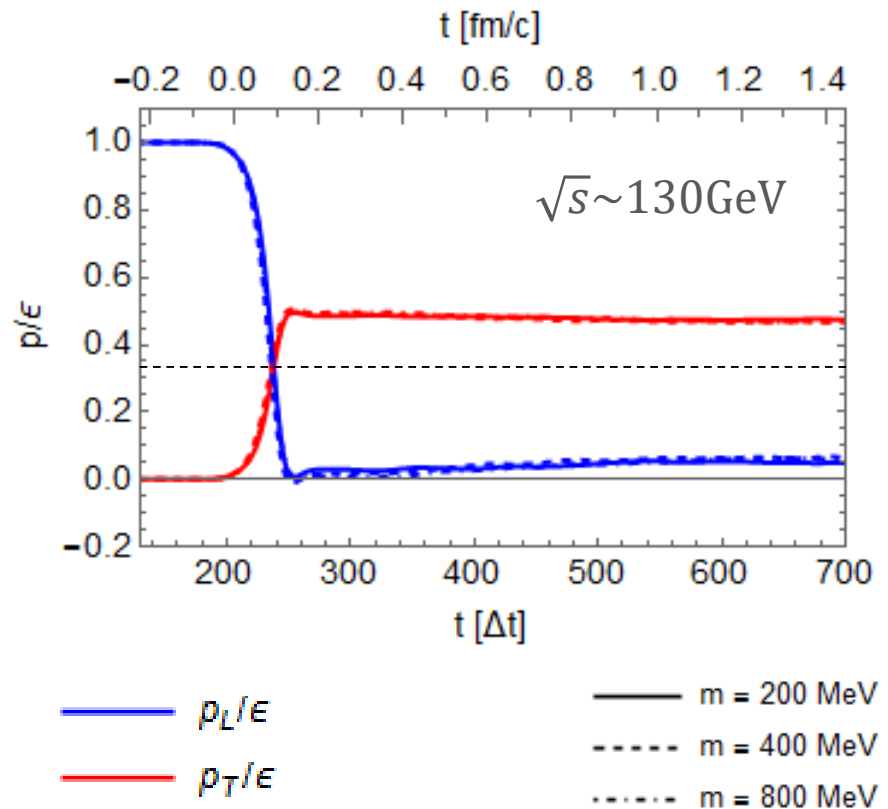


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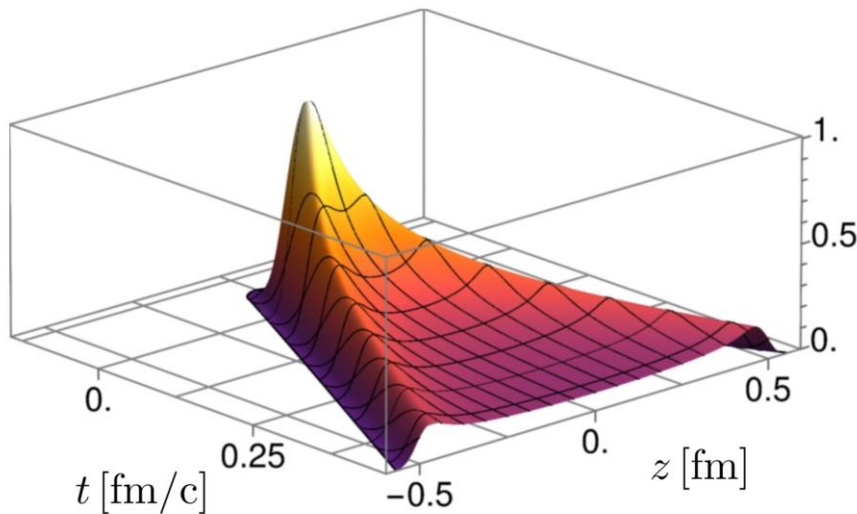


Transverse pressure distribution

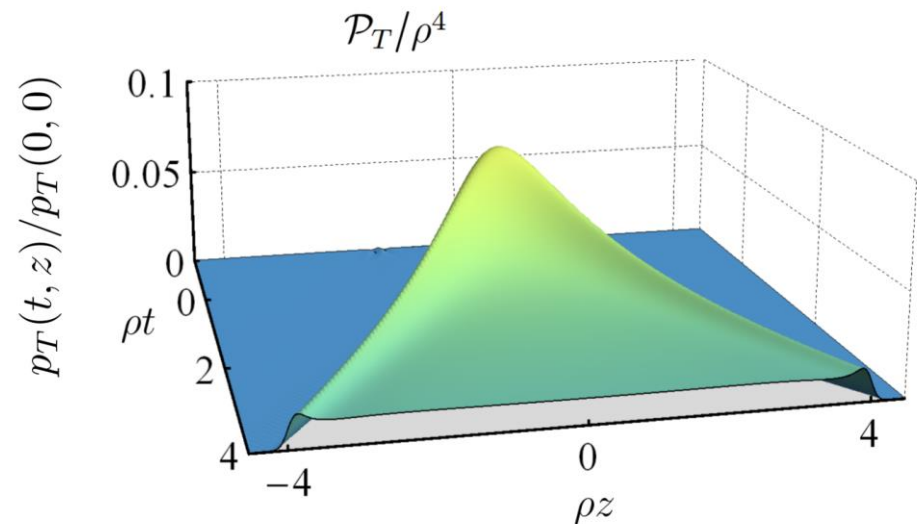
Transverse pressure $p_T(x)$ generated by longitudinal fields

$$\langle p_T \rangle = \frac{1}{2} \langle E_L^2 + B_L^2 \rangle$$

Boost-invariant case: initial conditions at $\tau = 0$ for longitudinal **E** and **B** fields, i.e. constant p_T along the boundary of the forward light cone



3+1 Yang-Mills



Holographic model

[Casalderrey-Solana et al., PRL (2013) 181601]

Conclusions and outlook

- Genuine 3D picture of heavy-ion collisions in the CGC framework
- Finite thickness breaks boost invariance → Gaussian rapidity profiles
- Strong pressure anisotropy, no isotropization
- Free-streaming flow

Future:

- Initial conditions with random (complicated) longitudinal structure
- Numerical improvements to study higher energies (LHC)
- JIMWLK evolved initial conditions

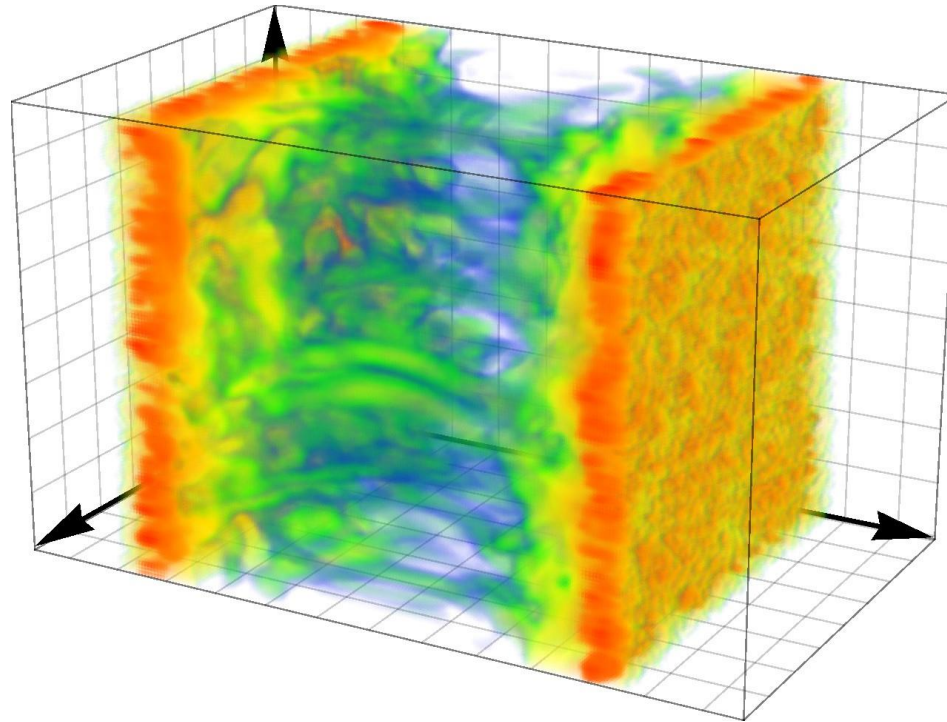
arXiv:1605.07184
arXiv:1703.00017

open source!

github.com/openpixi/openpixi (Java)

gitlab.com/monolithu/pyglasma3d (Python, Cython)

Thank you for your attention!



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Backup slides

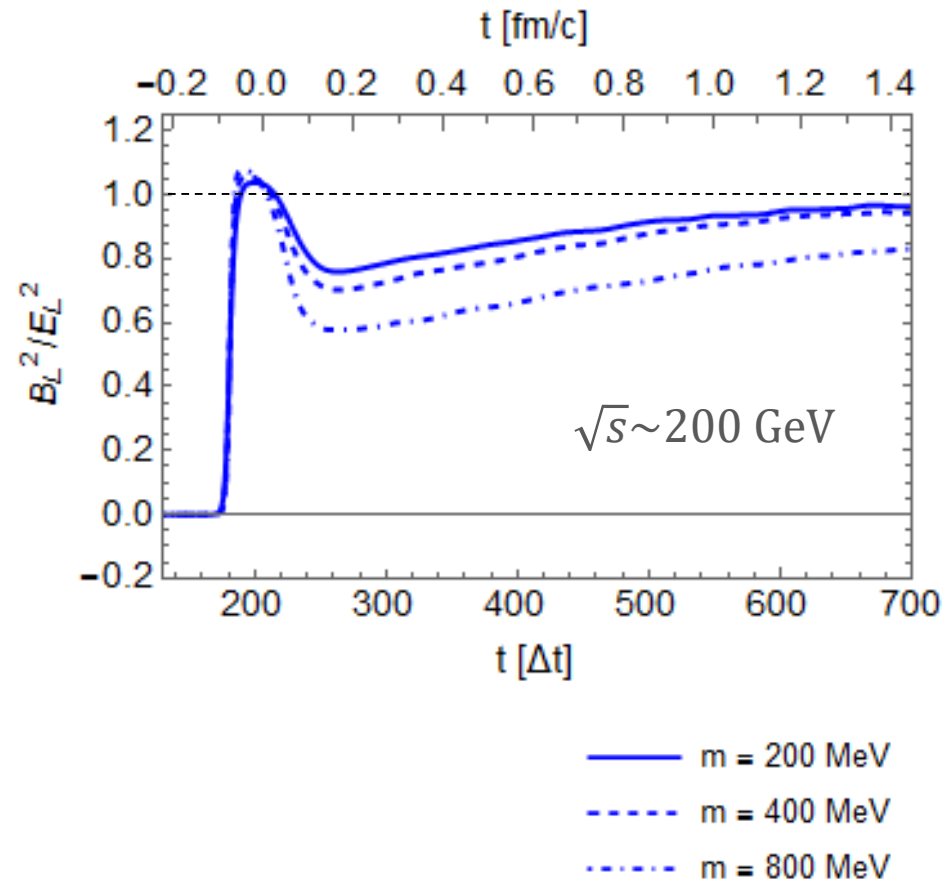
Chromo-magnetic and electric fields

Study creation and evolution of longitudinal fields

- Boost-invariant case:
 $\langle E_L^2 \rangle \sim \langle B_L^2 \rangle$
- Here: magnetic and electric fields are not created equally
- Strong dependence on IR regulator m and thickness

Large m , thick nuclei: low ratio
Small m , thin nuclei: ratio ~ 1

Artefact of initial conditions?



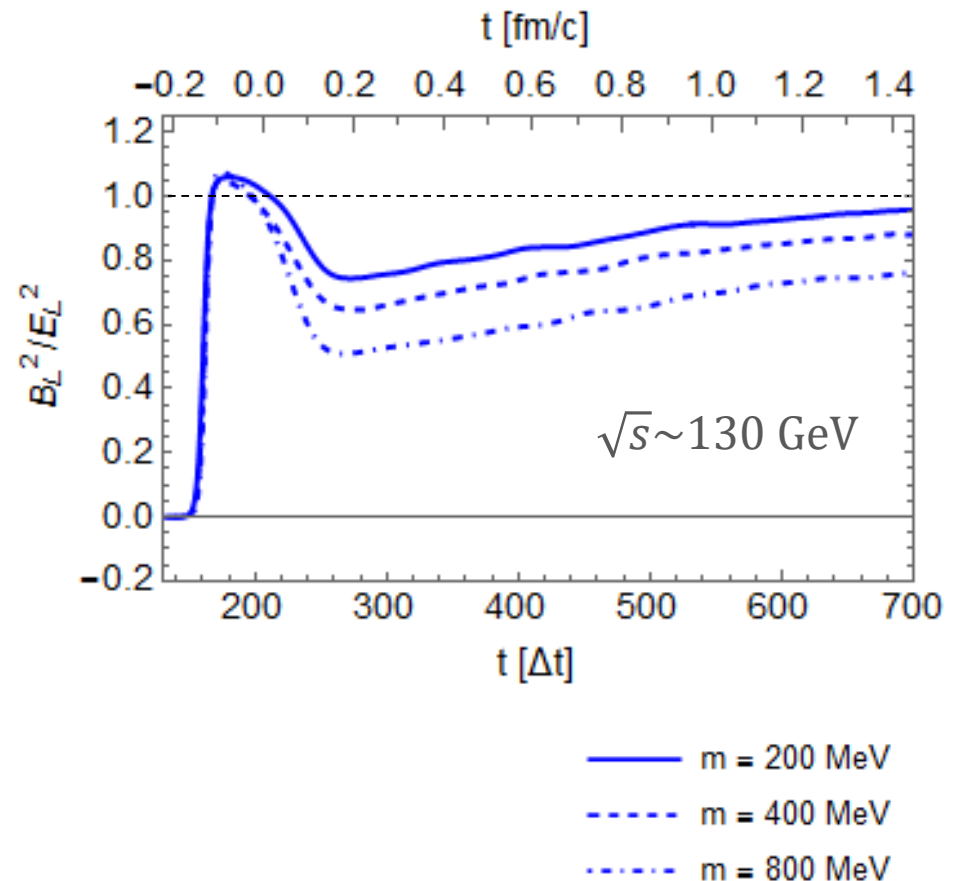
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Longitudinal structure (1)

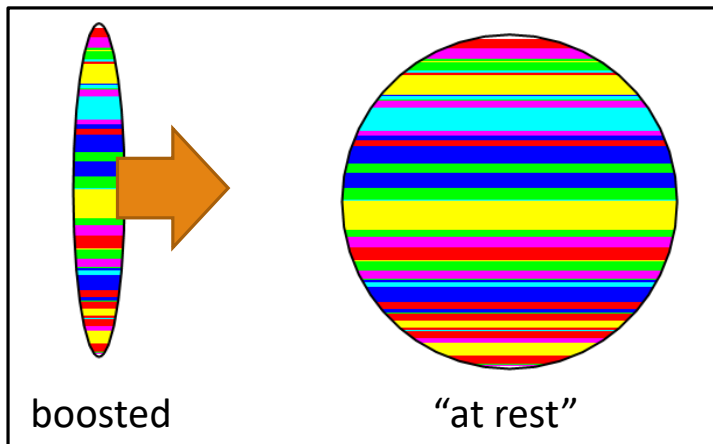
Initial conditions are still missing **random longitudinal structure**:

- Color charges (or fields) of the nuclei are correlated in the longitudinal direction (figure on the left)
- Realistic nuclei should consist of independent “sheets” (figure on the right) [Fukushima, PRD 77 (2008) 074005]

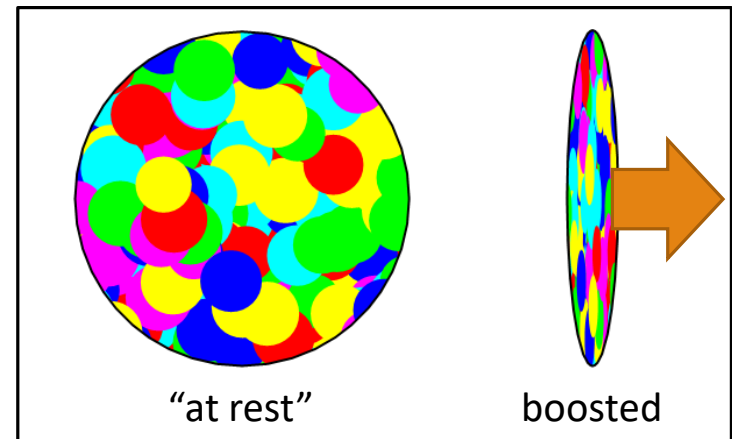
Consequences

- Higher energy density in the Glasma
- MV model: Higher Q_s for fixed μ [Lappi, Eur. Phys. J. C55 (2008) 285]
- Effect on ratio of longitudinal electric to magnetic field?

Current implementation



Longitudinal randomness



Longitudinal structure (2)

Wilson line expectation value $\langle \text{tr}(V) \rangle$ of a single nucleus is sensitive to longitudinal structure.

Embedded 2D MV-model:

$$\langle \hat{\rho}^a(\mathbf{x}_T) \hat{\rho}^b(\mathbf{x}'_T) \rangle = g^2 \mu^2 \delta^{(2)}(\mathbf{x}_T - \mathbf{x}'_T) \delta^{ab}$$

$$\rho(t, z, \mathbf{x}_T) = f(z - t) \hat{\rho}(\mathbf{x}_T)$$

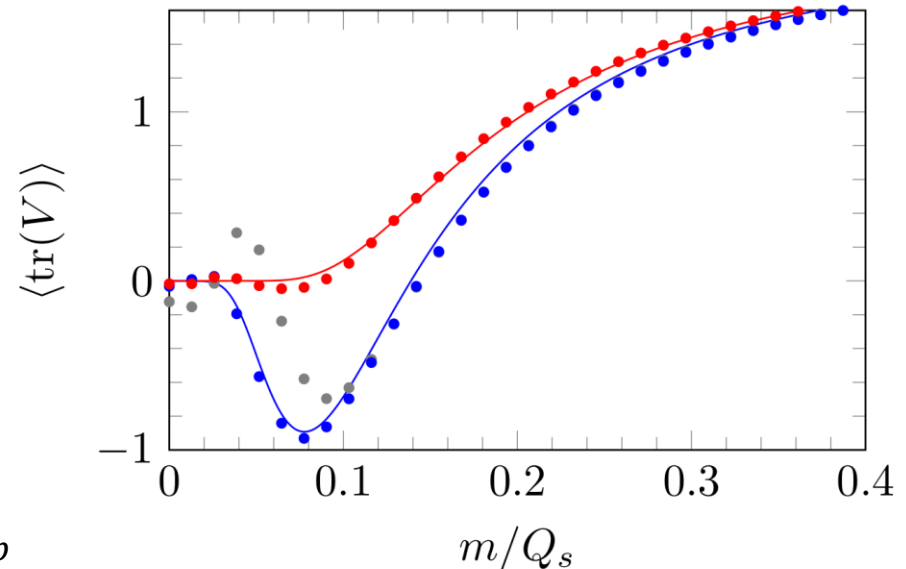
3D MV-model:

(with random longitudinal structure)

$$\langle \rho^a(t, \mathbf{x}) \rho^b(t, \mathbf{x}') \rangle = g^2 \mu^2 f(z) \delta^{(3)}(\mathbf{x} - \mathbf{x}') \delta^{ab}$$

$f(z)$... longitudinal profile function

Introducing independent “sheets” in longitudinal direction
[Fukushima, PRD 77 (2008) 074005]



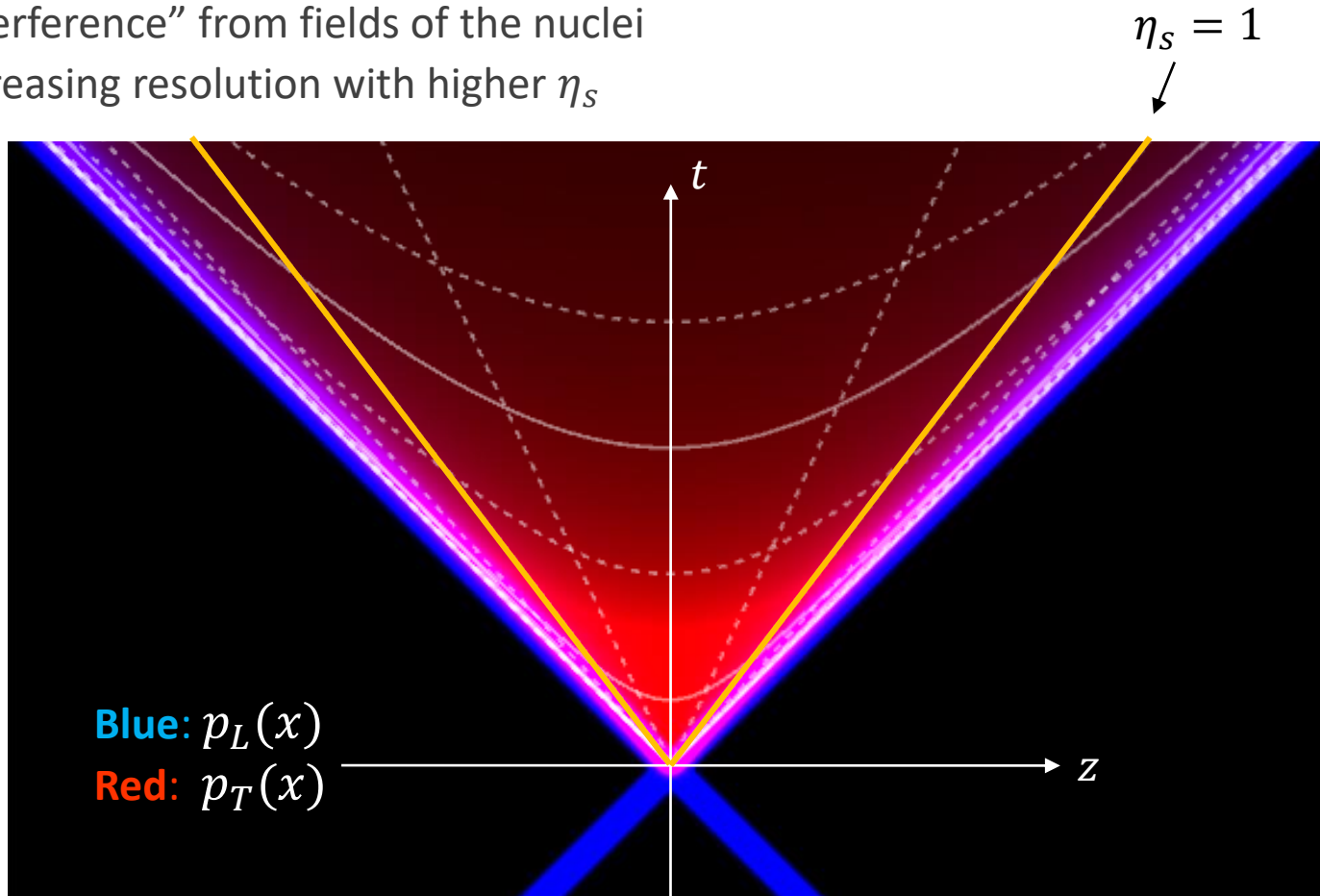
Lines: analytical result
Dots: numerical result

Blue: 2D MV-model
Red: 3D MV-model
Gray: intermediate

Extending the rapidity interval (1)

Limited rapidity interval due to ..

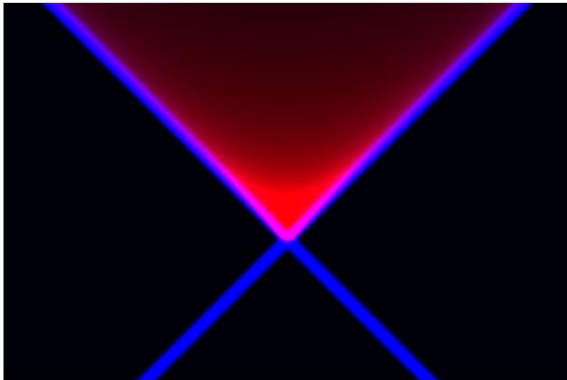
- longitudinal simulation box length / simulation time
- “interference” from fields of the nuclei
- decreasing resolution with higher η_s



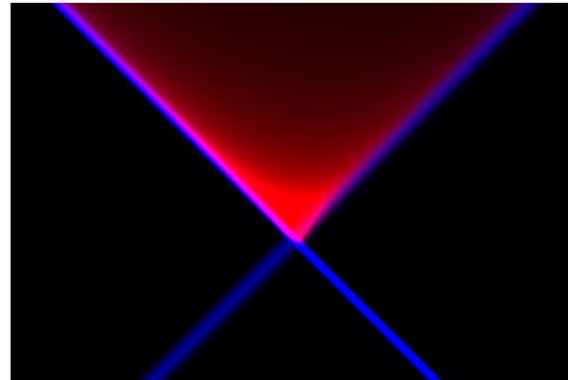
Extending the rapidity interval (2)

Idea:

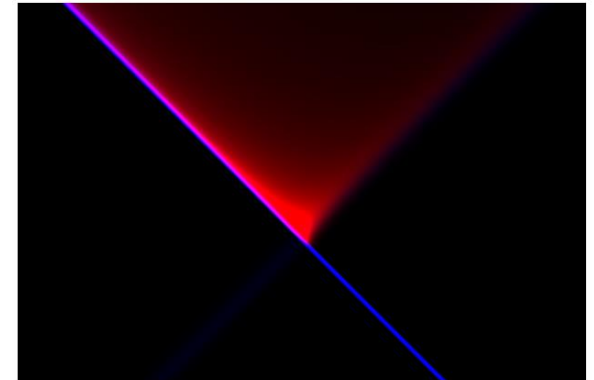
- Simulate collisions in boosted frame (asymmetric collision), parametrize with η'
- Record data only in “safe” interval $\eta_s \in (-1,1)$
- Boost back to center of mass frame



(a) $\eta' = 0$



(b) $\eta' = 0.5$



(c) $\eta' = 1$

Extending the rapidity interval (3)

Result:

- Rapidity profiles can be extended to higher ranges of η_s
- Non-trivial test of the boost invariance of the code

