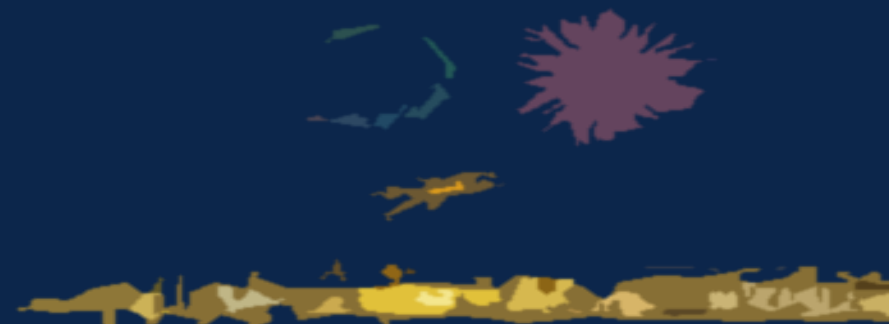




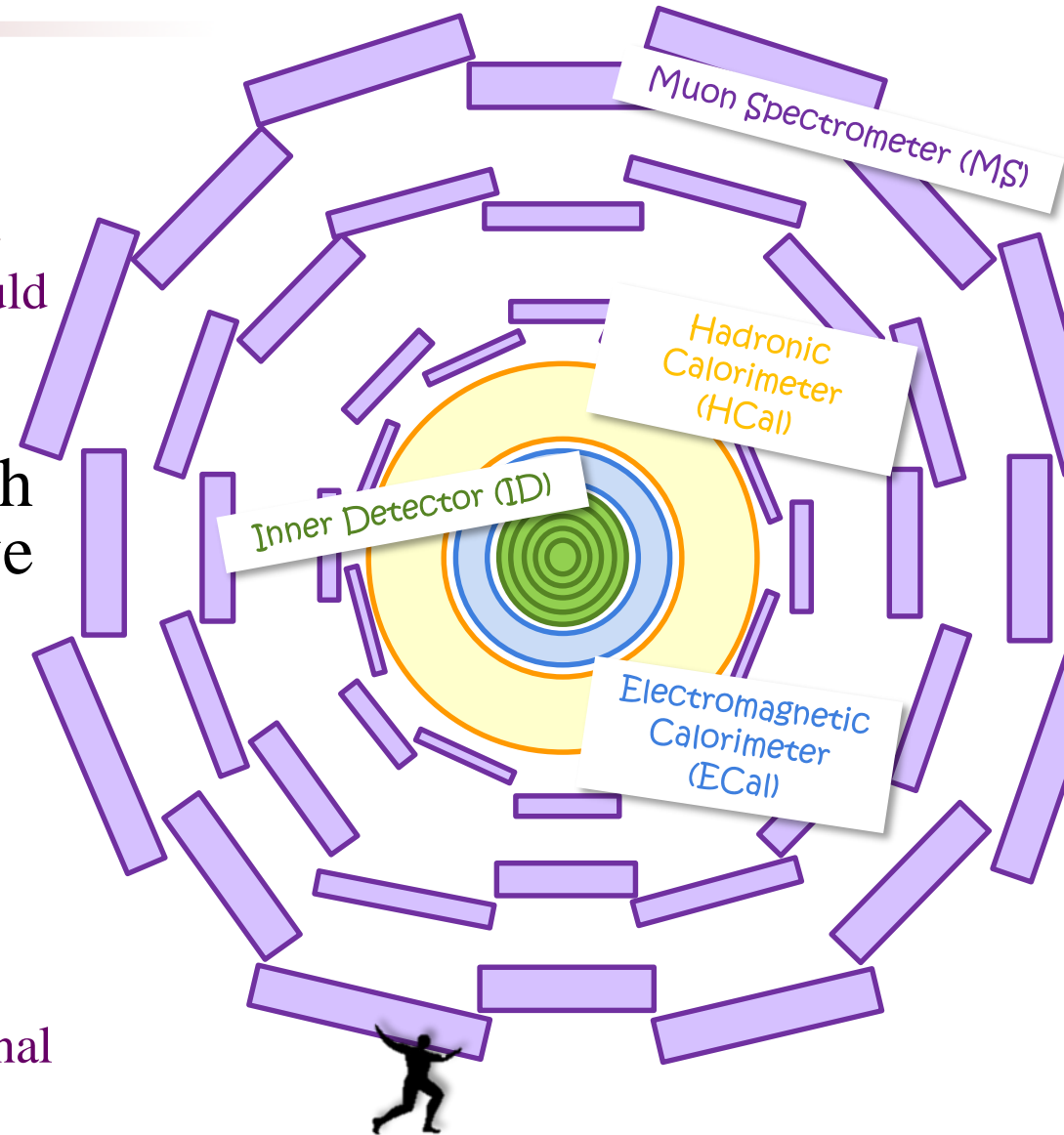
# Search for New Physics through the Reconstruction of Challenging Signatures and for Long-lived particles with the ATLAS detector

*Nora Pettersson (UMass) on behalf of  
The ATLAS Experiment*



# Introduction

- Many Beyond the Standard Model (BSM) theories predict atypical and unique signatures
  - ◀ These events often pose challenging to reconstruct
    - ◀ Detector responses to these signals are not what one would expect
- To this category belongs Long Lived Particles (LLP) of a wide range of lifetimes, particles with abnormal values of electrical charge and massive stable charged particles
- For example – an LLP decaying at a significant distance from the Interaction Point (IP) is strenuous for the tracking system
  - ◀ Standard tracking algorithms assume the particle trajectories arises from the IP
    - ◀ Need dedicated track reconstruction methods to gain signal significance for such signatures



# In this Talk...

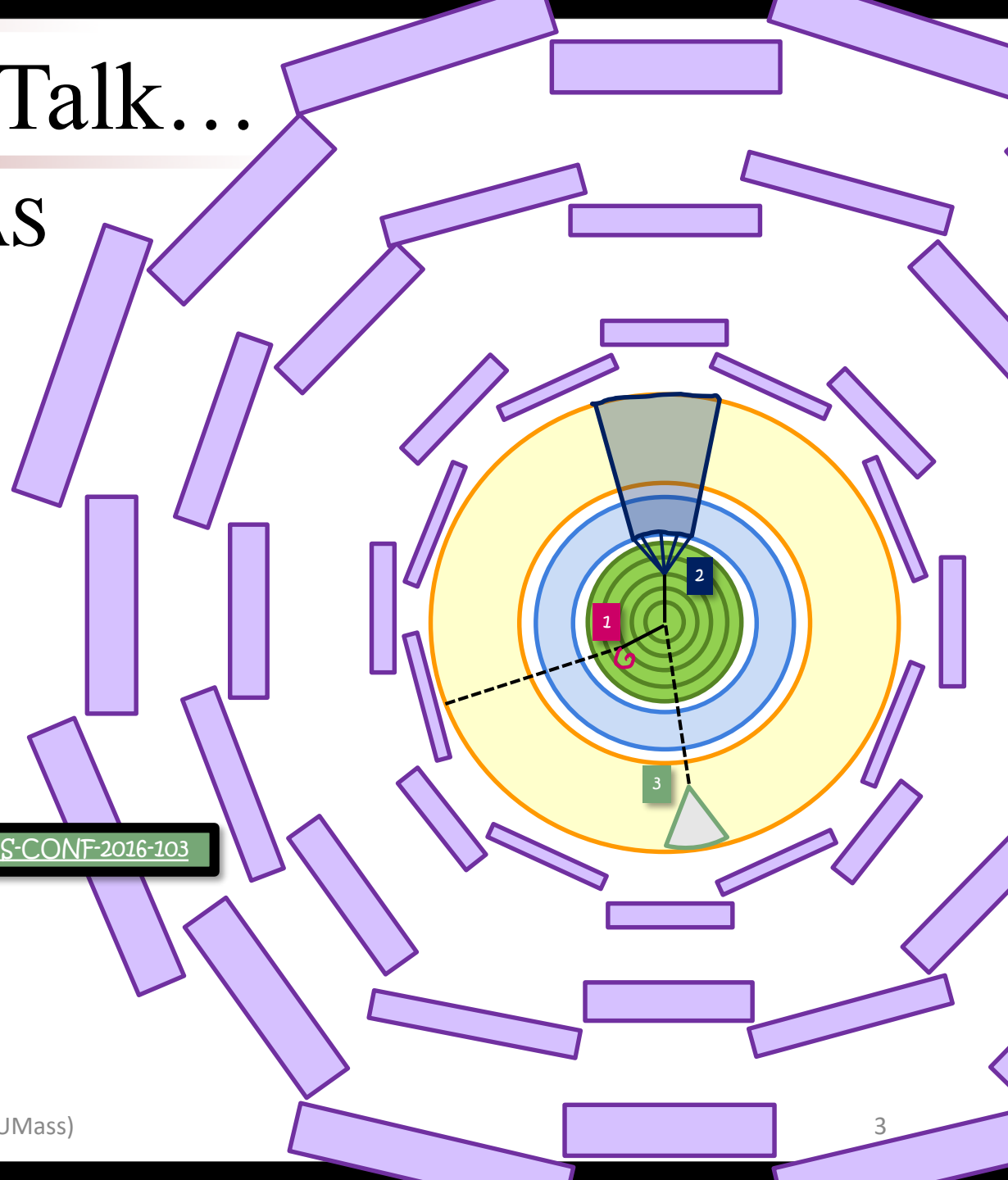
- Highlight three analyses from ATLAS

◀ All using unique and dedicated, specialised methods of signal event selection and reconstruction

**1** Disappearing-track [ATLAS-CONF-2017-017](#)

**2** Displaced Vertex [ATLAS-CONF-2017-026](#)

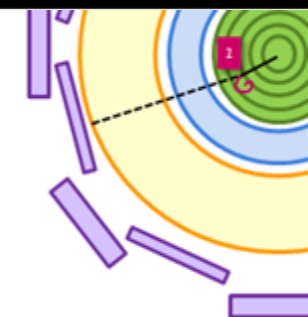
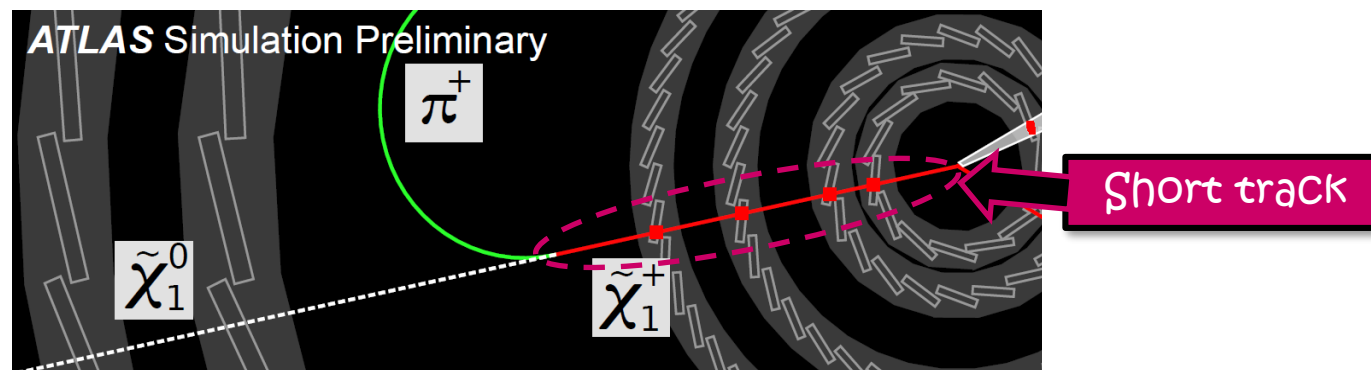
**3** Decays within the Calorimeter [ATLAS-CONF-2016-103](#)



# Disappearing Track\*

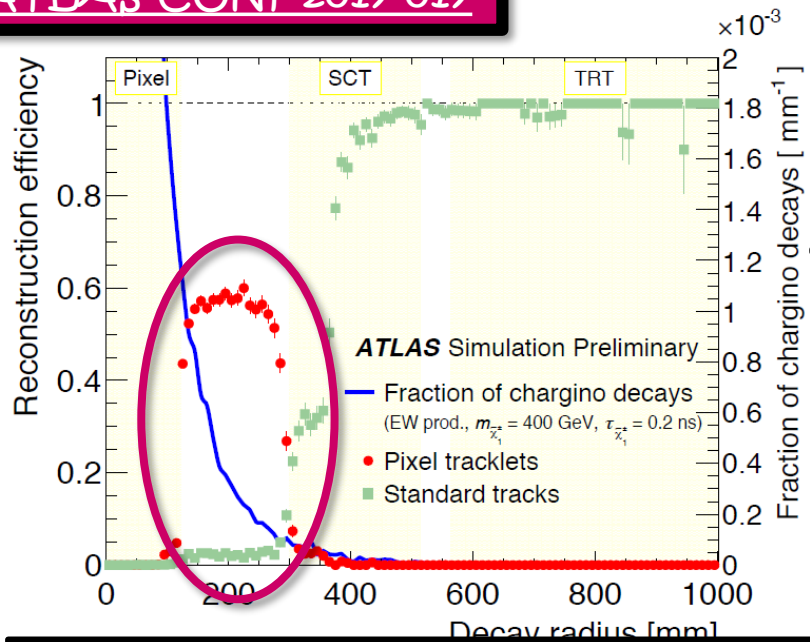
\*Track = Charged Particle Trajectory

Event topology: high transverse momentum ( $p_T$ ), large missing transverse energy ( $E_T^{\text{miss}}$ ), high  $p_T$  jets and a short track



- Assume SUSY model where  $\tilde{\chi}_1^\pm$  (NLSP) is nearly mass-degenerate with  $\tilde{\chi}_1^0$  (LSP) – Long-lived Chargino decays:  $\tilde{\chi}_1^+ \rightarrow \tilde{\chi}_1^0 \pi^+$  (soft)
  - ◀ Common to wino LSP scenarios – vital to a large portion of SUSY dark matter searches
- Disappearing track – How does one reconstruct short tracks?
- Track reconstruction is done in two steps
  - ◀ Standard algorithms– e.g. to find mainly the primary tracks
    - ◀ Requiring at least seven measurements (hits) in the silicon detector layers
  - ◀ A second pass of the tracking is run
    - ◀ Only using leftover measurements from the first pass
    - ◀ The hit requirement is significantly looser and aimed at short tracks: *at least four hits in the pixel layers*

# Disappearing Track

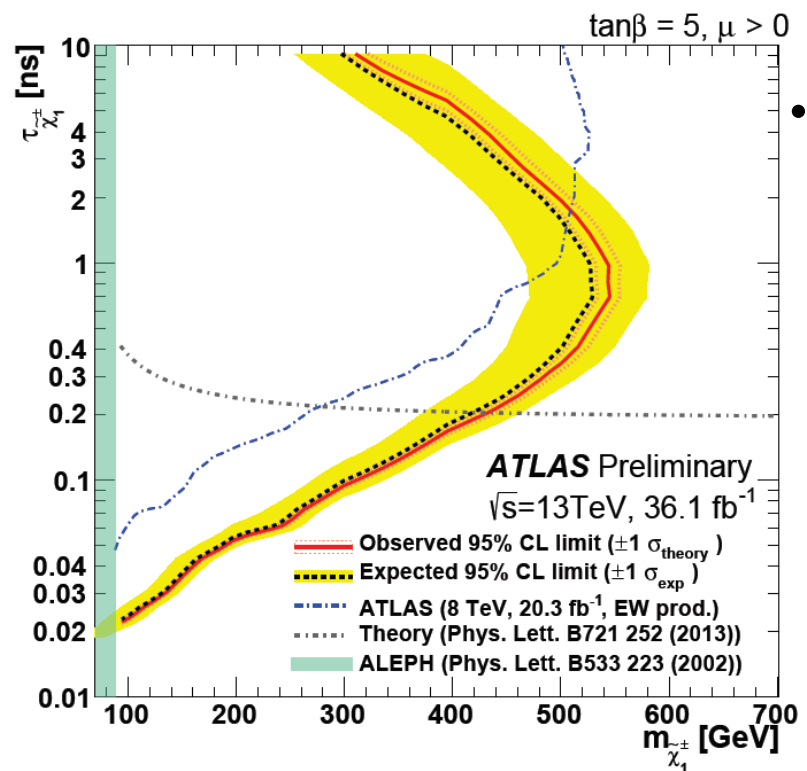


- Pixel tracklets effectively increase the efficiency at decay radius  $< 300 \text{ mm}$  from near 0% up to 60%
- Remaining inefficiency is due to LLP decaying before obtaining 4 pixel hits

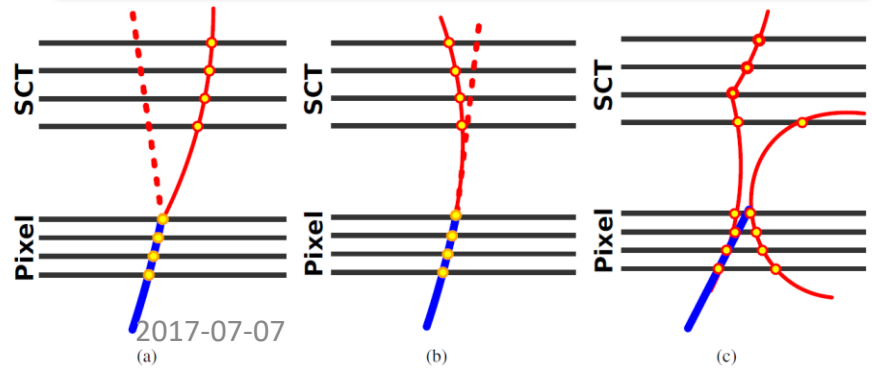
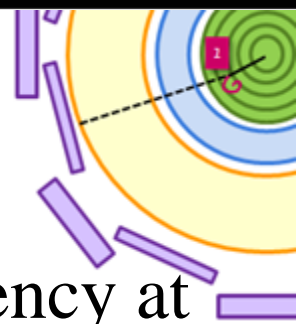
Red solid (dotted) shows neutral (charged) particles; thick blue represents the reconstructed tracklet

**Background:**

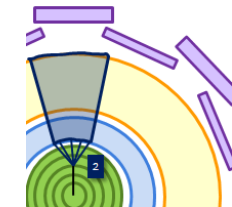
- Hard scattering of a hadron
- Lepton emitting a hard photon
- Random combination of hits



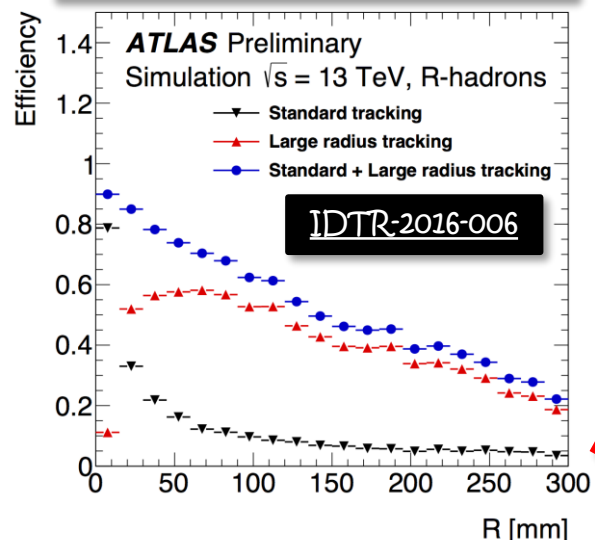
- Exclusion limits are presented as a function of the LLP lifetime and mass
- Improved sensitivity to shorter lifetimes is visible comparing previous ATLAS results to the new set limits



# Displaced Vertex

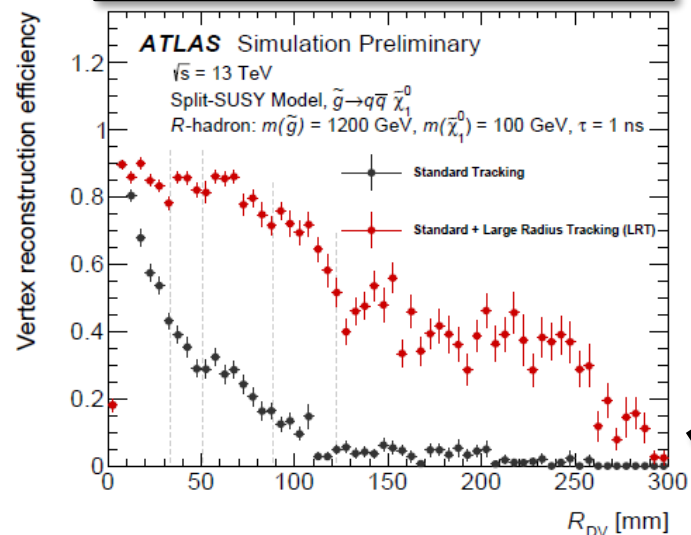


Track Reconstruction Efficiency



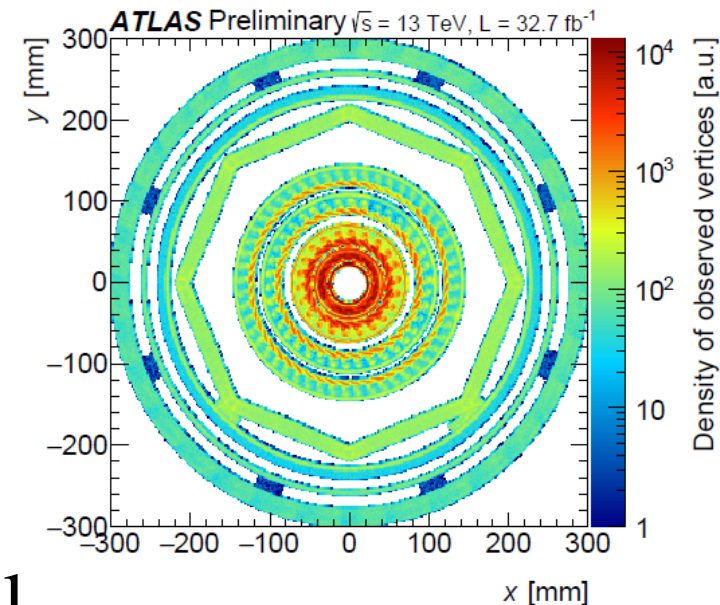
- R-hadrons decaying within the ID can be reconstructed as a Displaced Vertex (DV)
  - ◀ Long-lived Gluinos are present in Mini-split SUSY models
- Low track-reconstruction efficiency relying only on standard tracking
  - ◀ Strict cuts on transverse ( $d_0$ ) and longitudinal ( $z_0$ ) impact parameters with respect to the IP
    - ◀ Where as tracks from a displaced decay typical have significant large impact parameters
  - ◀ Yields low track reconstruction efficiency means low vertex reconstruction efficiency for displaced decays
- A dedicated second pass of the tracking is ran on leftover hits from the standard track - much like for the disappearing-track analysis - but with *looser cut on  $d_0$  and  $z_0$*  !
  - ◀ Greatly increase the signal vertex reconstruction efficiency
  - ◀ Analysis sensitive to decay length up to 300 mm

Vertex Reconstruction Efficiency



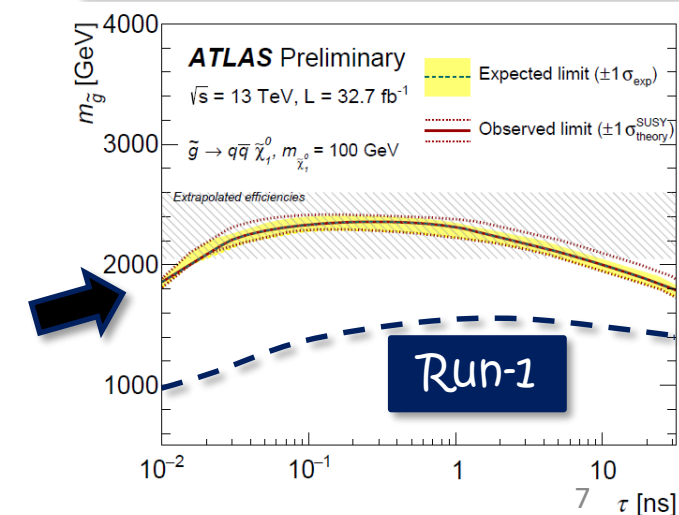
# Displaced Vertex

- Signal regions defined for massive displaced vertices with large track multiplicity
  - ◀ No SM background to the analysis
- Hadronic interactions with the material are the largest background - typically large impact parameters as well
  - ◀ Create a 3D map of the dense material region of the ID
    - ◀ Veto vertices in these regions
- Detector effects such as mis-reconstructions and combinatorics remains
  - ◀ Fully data-driven background estimations
- The increase in energy and luminosity have greatly improved the analysis exclusions set for R-hadrons

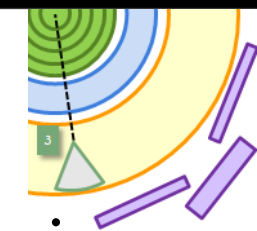


## Material Veto:

Projection of all vertices in data vetoed by the material map in the x-y plane of the ATLAS ID



# Decays within the Calorimeter

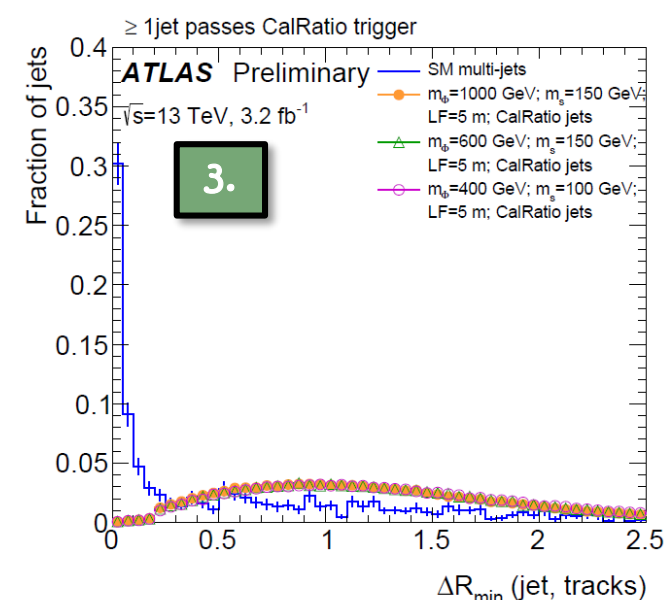
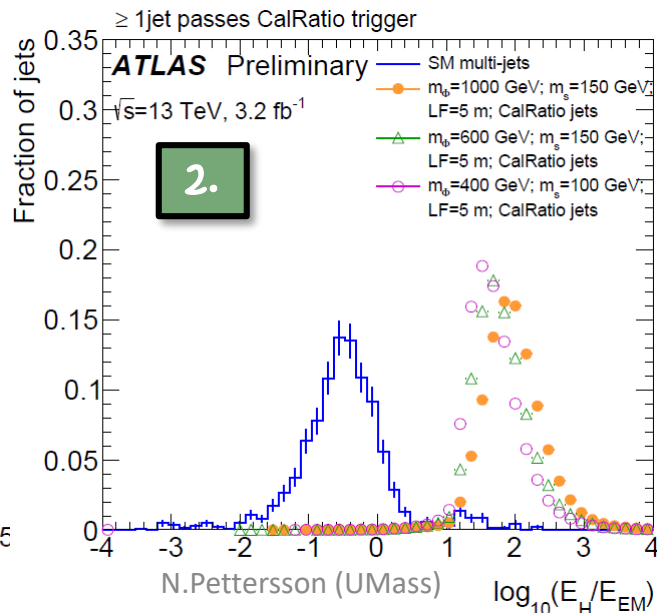
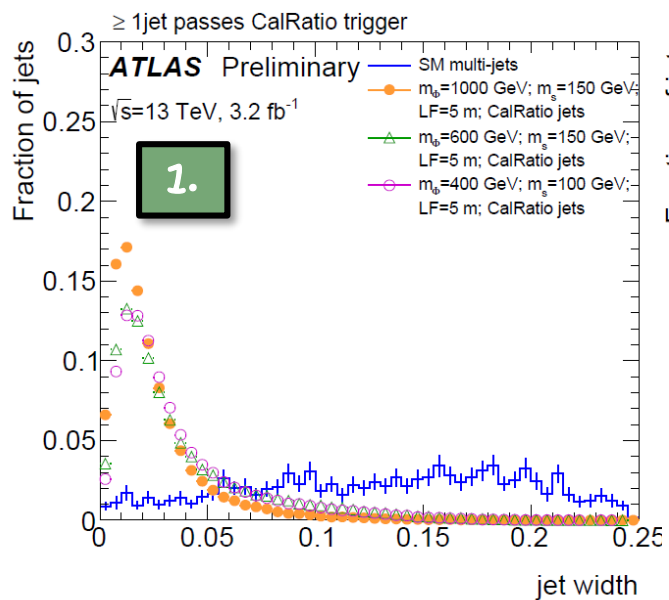


- Weakly coupled Hidden Sector (HS) may communicate with the SM via neutral LLP that decays within the ATLAS detector

◀ Decays within the HCal give an unique signature of two atypical jets

◀ Decays within the calorimeter volume:

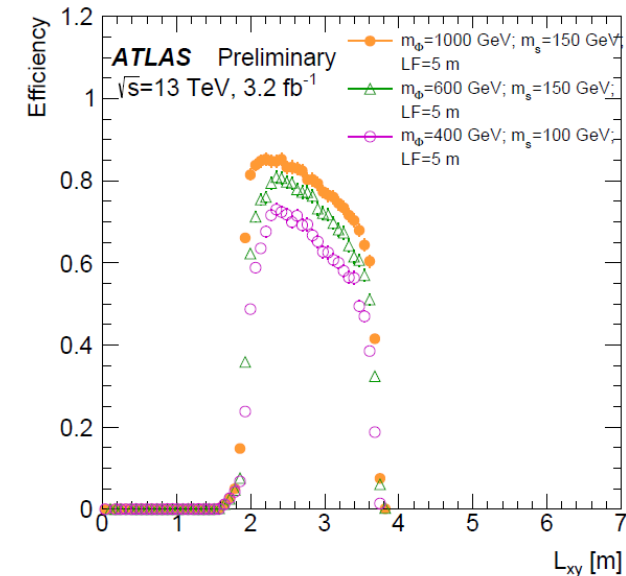
1. Produce *narrower cone* than typical SM jets originating from the IP (SM multi-jets)
2. Ratio of energy in the HCal ( $E_H$ ) to the energy in the ECal ( $E_E$ ) is large ( $\log_{10}(E_H/E_E)$  *CalRatio*)
3. No associated ID tracks to the jet – large  $\Delta R_{\min}$  between the jet direction and the tracks





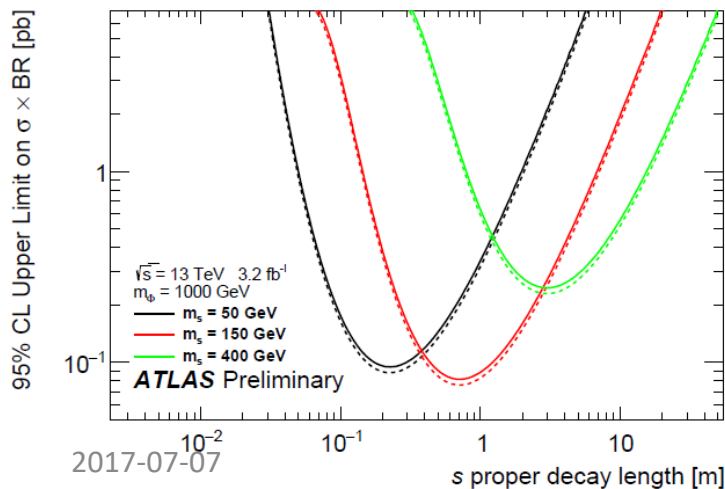
# Decays within the Calorimeter

- Dedicated *CalRatio* trigger developed on the topology of neutral LLP within the HCal
  - ◀ Narrow jets using cone radius  $R=0.2$  (standard is  $R=0.4$ )
  - ◀  $\log_{10}(E_H/E_{EM}) > 1.2$ 
    - ◀ Small deposits are expected in the ECal
    - ◀ LLP decays at the outer edge of the ECal or contribution from pile-up
  - ◀ Requiring that no tracks are within the jet cone



### CalRatio Trigger Efficiency:

High efficiency (~70%) is obtained to trigger on decays within the calorimeter volume and for obvious reasons drops significantly outside the acceptance



- Boosted Decision tree (BDT) is used to select signal based on kinematics of the jets
  - ◀ Background purely data-driven ABCD method
  - ◀ Acceptance significantly improved by relaxing of  $E_T^{\text{miss}}$  and using of the BDT for high mass LLP
- Exclusion are set for few mass points of a heavy scalar boson
  - ◀  $<1$  TeV and a wide range of decay lengths from 0.05 mm to 16 m

# Conclusion

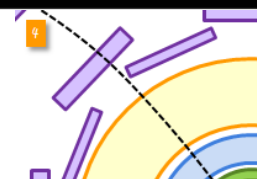
- Several interesting atypical searches have been conducted by the ATLAS Experiment – all using their unique techniques and methods for reconstruction, trigger, and object selections
  - ◀ Three of these analyses have been reviewed focusing on the technical aspects
    - ◀ Searches using signatures of Disappearing Track, Displaced Vertex, Decays within the HCal, and Highly Ionising Particles
- Many more interesting signatures
  - ◀ Charged LLP decays to a light neutral and a heavy charged particle
    - ◀ Results in a kinked track – opposite scenario of the disappearing track
  - ◀ Neutral LLP decaying to neutral stable particles and a photon
    - ◀ Gives a signature of a late photon
  - ◀ More generic final states with Displaced Vertices
  - ◀ Different types of HIP
    - ◀ Stopped in the calorimeter, particle ID track, slow heavy ionising muon
  - ◀ Stopped Particles
    - ◀ Very, very, very slow particles – decays in a different bunch crossing

Kinked Track:
<a href="#"><u>Phys. Rev. D88 112006 (2013)</u></a>
Late Photon:
<a href="#"><u>Phys. Rev. D90 112005 (2014)</u></a>
Displaced Vertices:
<a href="#"><u>Phys. Rev. D92 072004 (2015)</u></a>
HIP:
<a href="#"><u>PLB 760 (2016) 647</u></a>
Stopped Particles:
<a href="#"><u>Phys. Rev. D88 112003 (2013)</u></a>

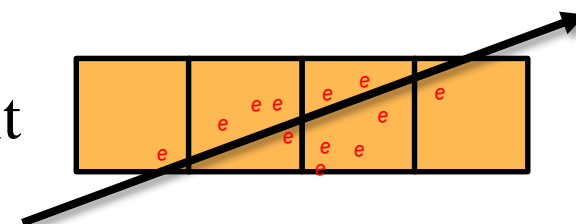
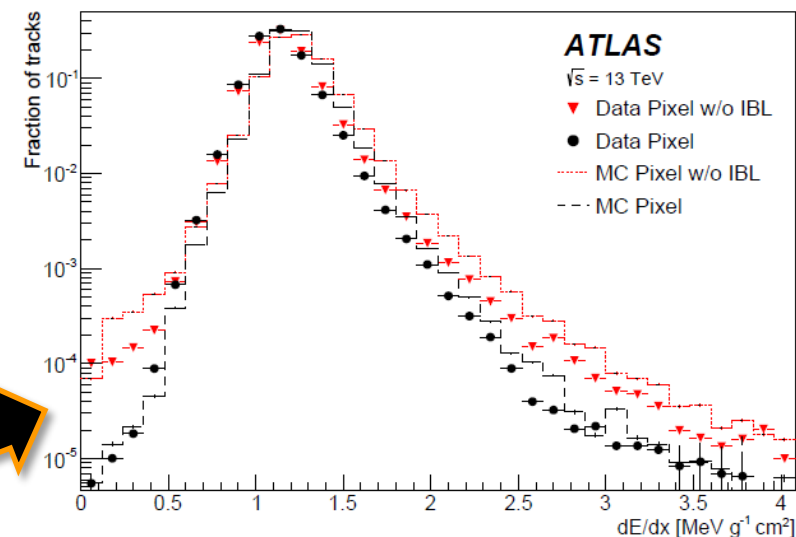
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# BACKUP

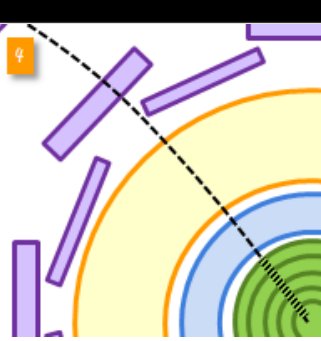
# Metastable Heavy Charge Particles



- LLP with a larger unit of charge than expected by SM particles
  - ◀ Measured by the energy loss in the pixel subsystem  $dE/dx$ 
    - ◀ Targeting models predicting R-hadrons of a wide range of lifetimes
- Additional pixel layer in after run-1: Insertable B-Layer (IBL)
  - ◀ Improves  $dE/dx$  measurement capability
    - ◀ Reduces the tail of measurement and increasing the number of tracks with valid ionisation measurement
- A charged particle usually crosses several pixels and the deposits in adjoining pixels are combined into clusters
  - ◀ All pixels in a cluster are used to calculate the total charge
- And *finally the  $dE/dx$*  is measured as the charge in the cluster per unit length of the track in the sensor
  - ◀ Using all clusters associated to a track
- One or two of the highest  $dE/dx$  measurements are removed to reduce tails of the distribution

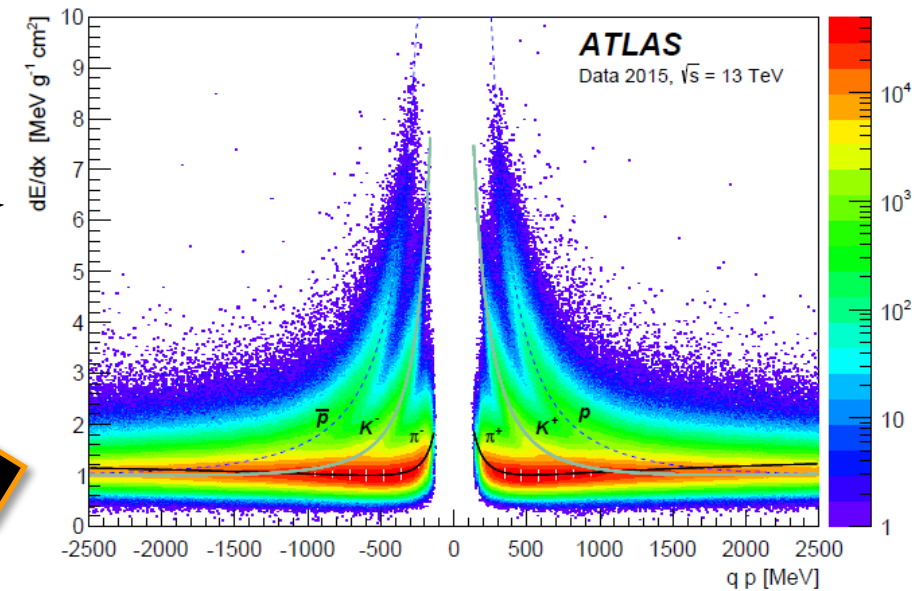


# Metastable Heavy Charge Particles

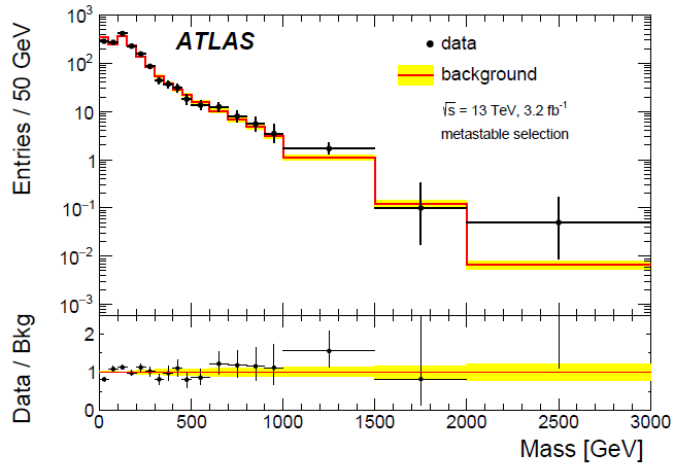
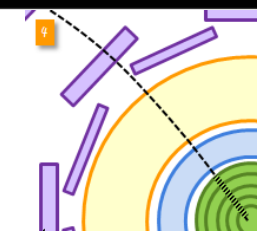


Average charge released:  
20000 (16000) e in pixels (IBL) and  
the threshold is set at 3500 (2500) e  
The signal Time over Threshold (ToT):  
is proportional to the ionisation charge  
With a dynamic range of 8.5 (1.5)  
times the average charge released by a  
Minimum Ionising Particle (MIP) for  
the pixel (IBL)

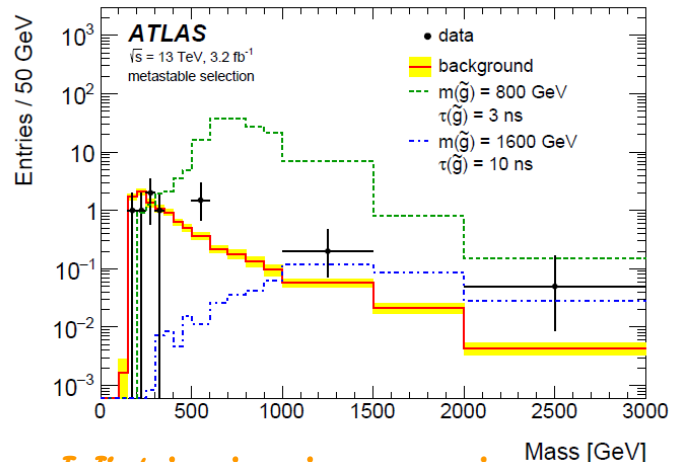
- Mass fit is performed on the energy loss measurement and the momentum of the reconstructed track
  - ◀ The fit used is a parametric Bethe-Bloch distribution
- $dE/dx$  versus  $qp$  in collected minbias events
  - ◀ Mass fits of the pions, kaons, and protons



# Metastable Heavy Charge Particles



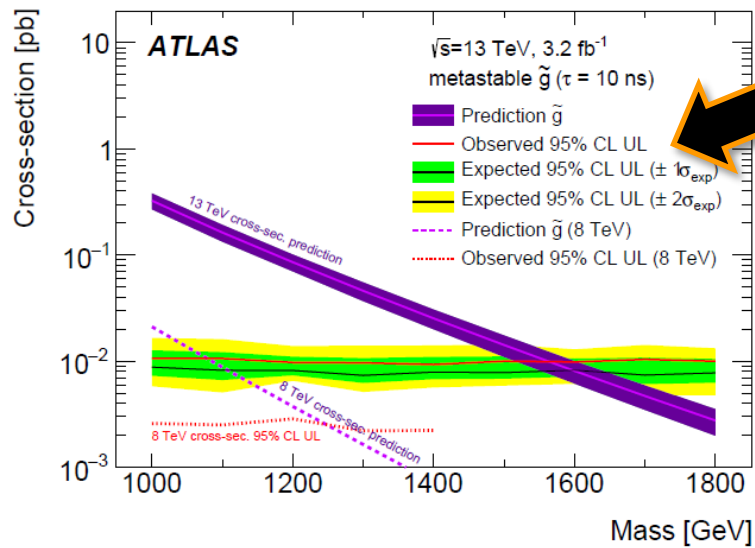
Without ionisation requirement



With ionisation requirement

- In addition to requirements on high track  $p_T$ , isolation and large  $E_T^{miss}$  - an  $\eta$  dependent ionisation cut is demanded

◀ Greatly reduce the background



- Exclusions are set on the masses for lifetimes of 10 ns or greater which are the strongest yet
- The increase in luminosity, the additional pixel layer and improved analysis techniques have yielded significant extended exclusion with respect to Run-1

