



Stochastic GW searches and Cosmology with GWs

Giancarlo Cella (INFN Pisa) for the LIGO-VIRGO Collaboration

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Stochastic background in a nutshell

A stochastic background can be seen as

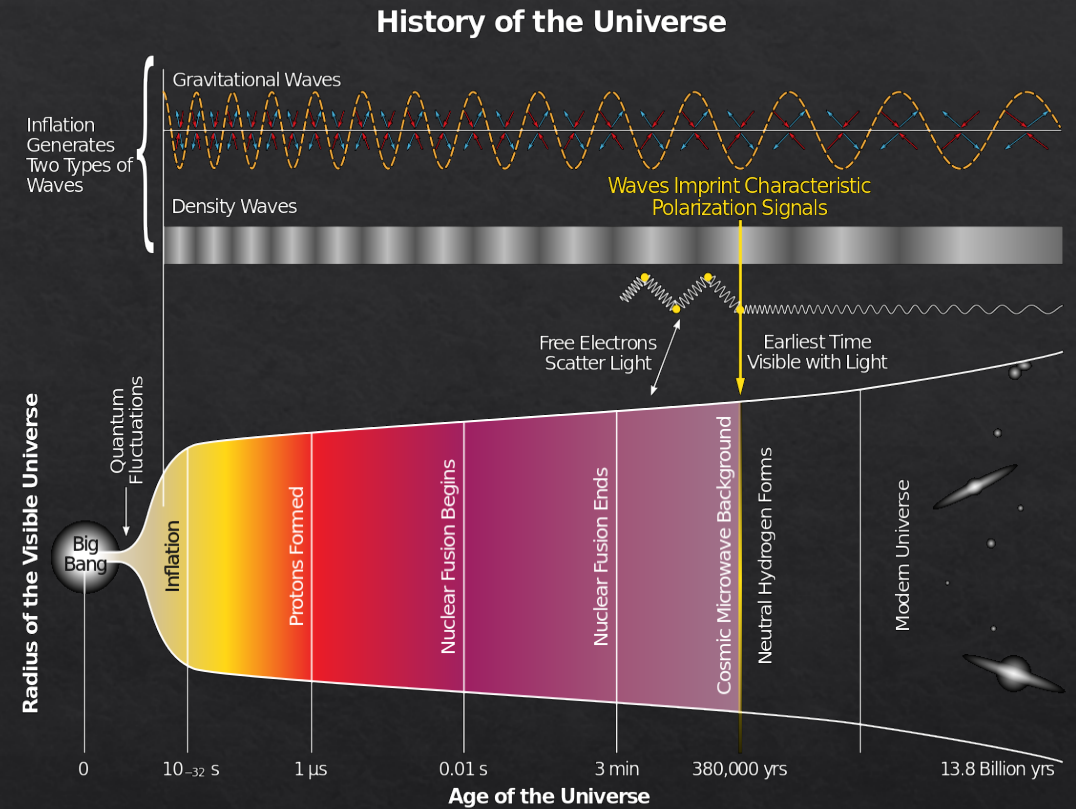
- a GW field which evolves from an initially random configuration
- the result of a superposition of many uncorrelated and unresolved sources

Two different kinds:

- **Cosmological:**
signature of the early Universe
inflation, cosmic strings, phase transitions...
- **Astrophysical:**
sources since the beginning of stellar activity
compact binaries, supernovae, rotating NSs, core-collapse to NSs or BHs, supermassive BHs...

Typical «first approximations» :

- 1) Gaussian, because sum of many contributions
- 2) Stationary, because physical time scales much larger than observational ones
- 3) Isotropic (at least for cosmological backgrounds)



If these are true, SB is completely described by its power spectrum

What we observe

Uncorrelated

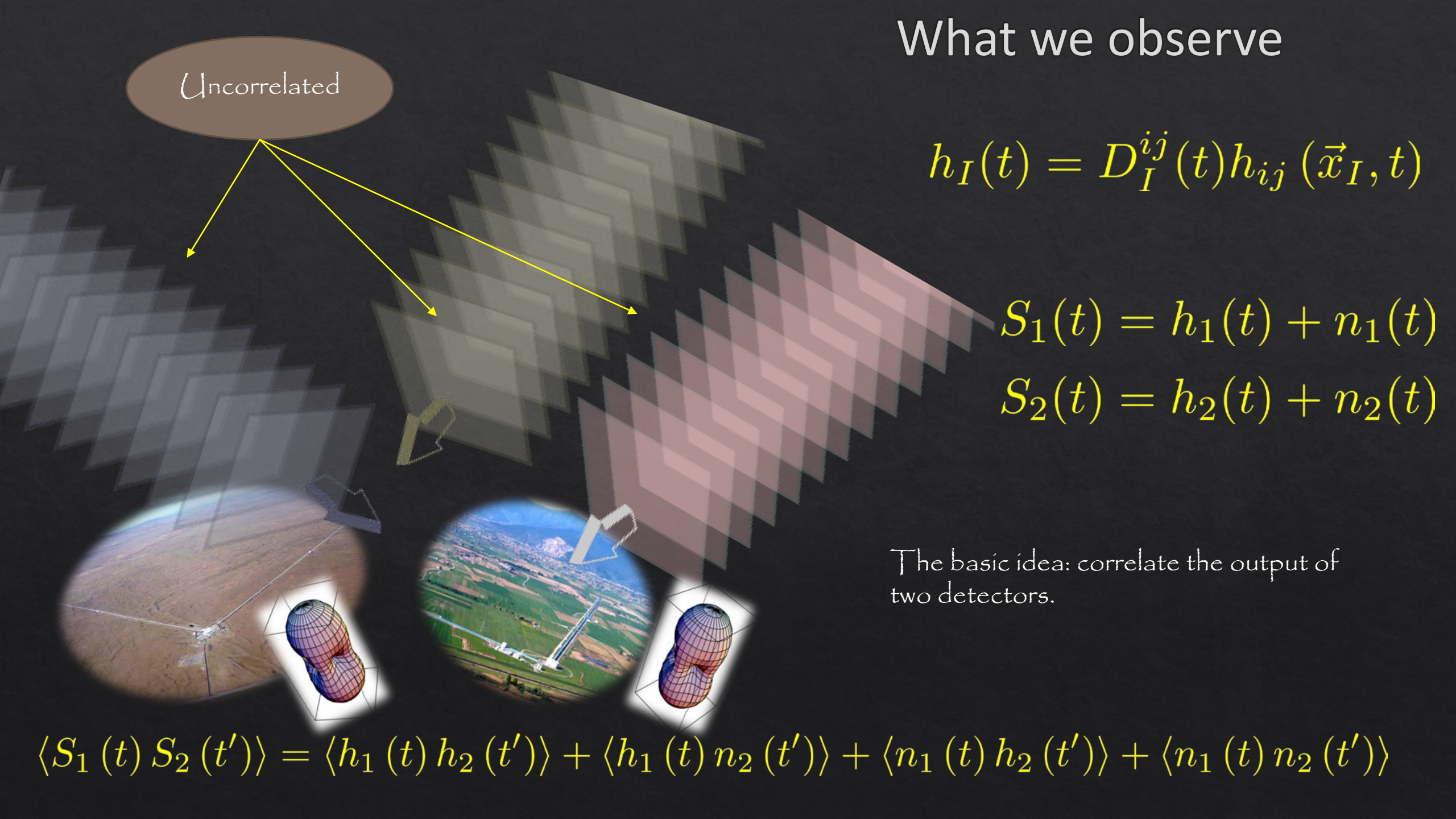
$$h_I(t) = D_I^{ij}(t) h_{ij}(\vec{x}_I, t)$$

$$S_1(t) = h_1(t) + n_1(t)$$

$$S_2(t) = h_2(t) + n_2(t)$$

The basic idea: correlate the output of two detectors.

$$\langle S_1(t) S_2(t') \rangle = \langle h_1(t) h_2(t') \rangle + \langle h_1(t) n_2(t') \rangle + \langle n_1(t) h_2(t') \rangle + \langle n_1(t) n_2(t') \rangle$$



What we observe

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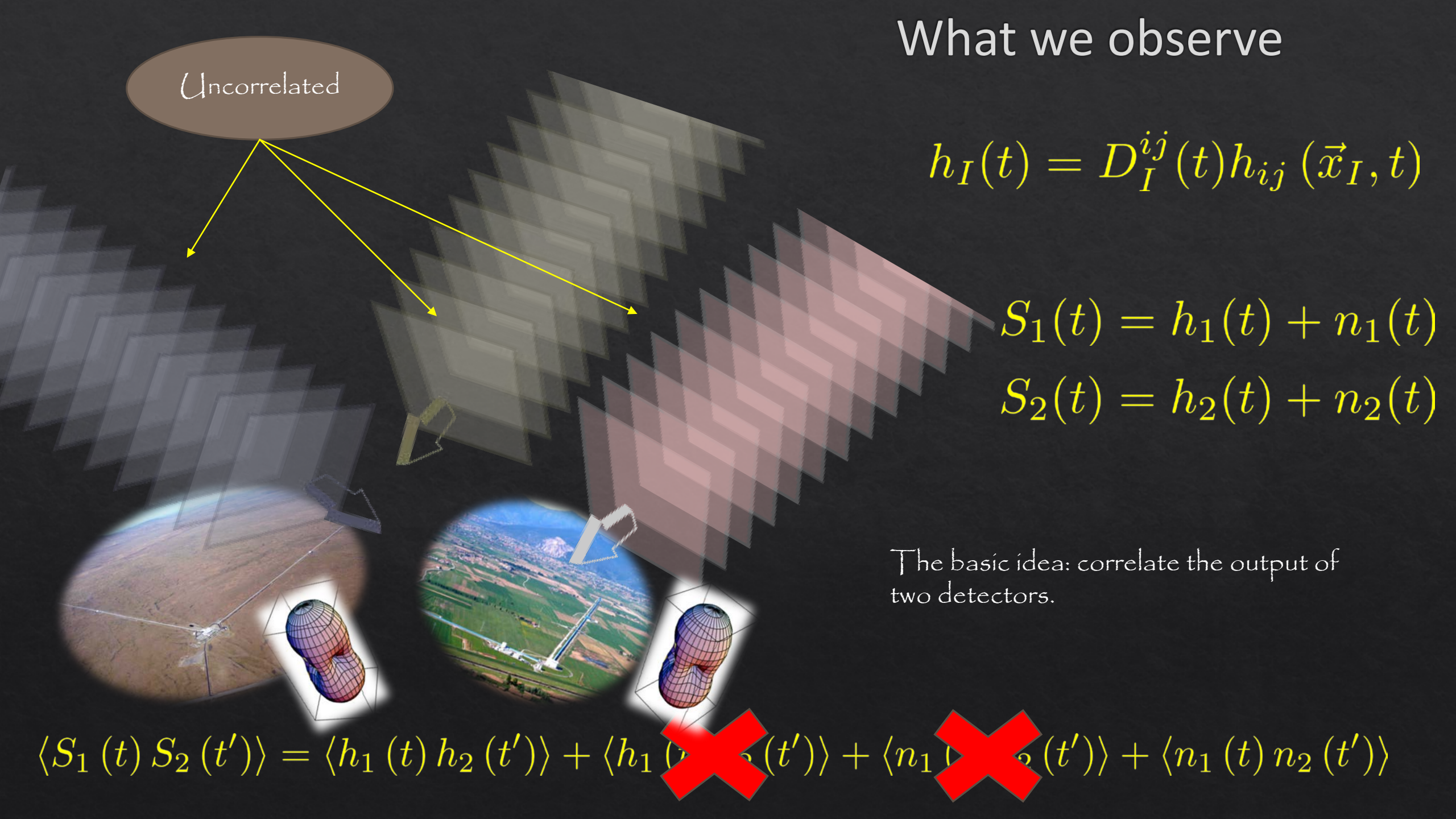
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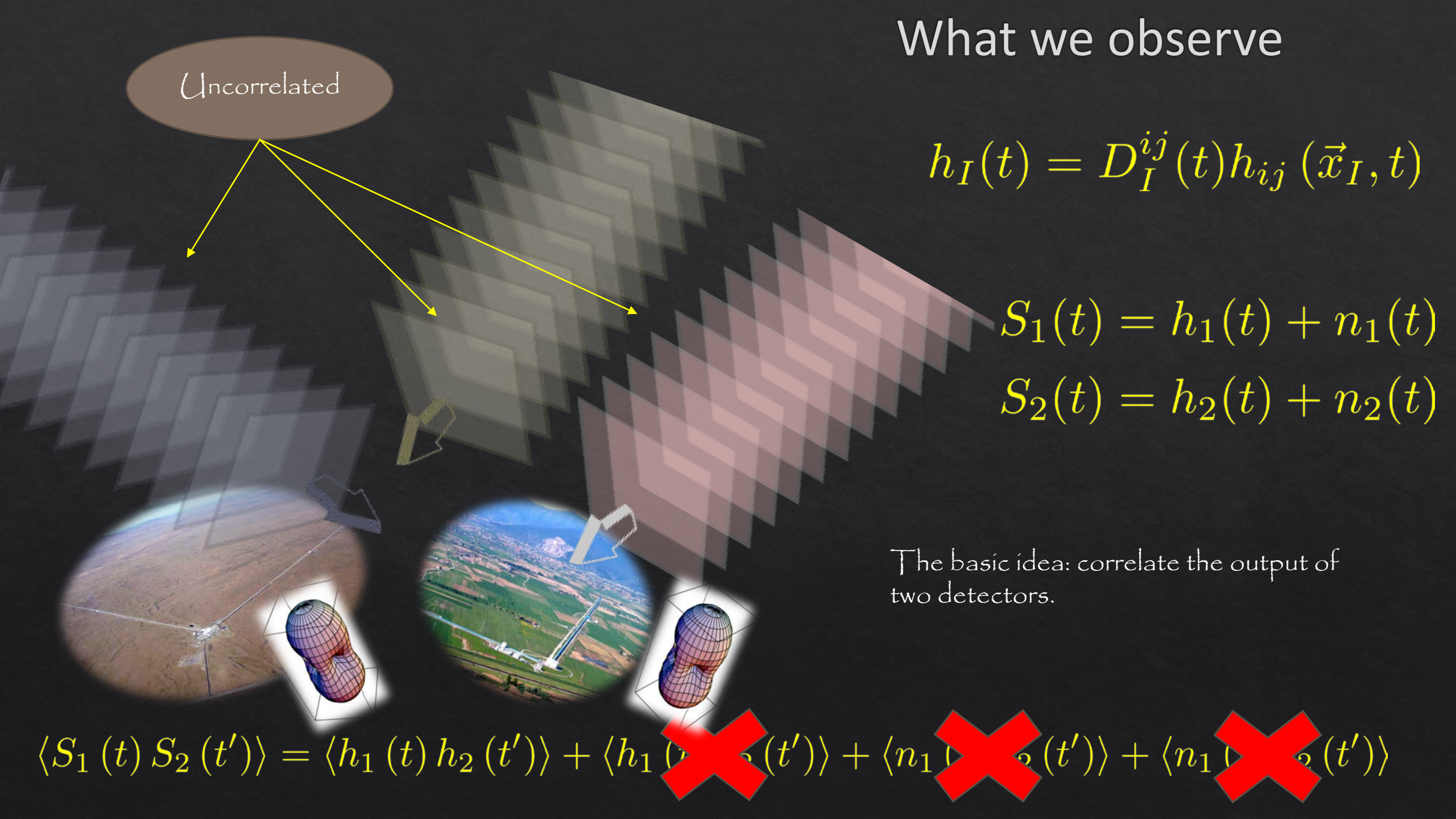
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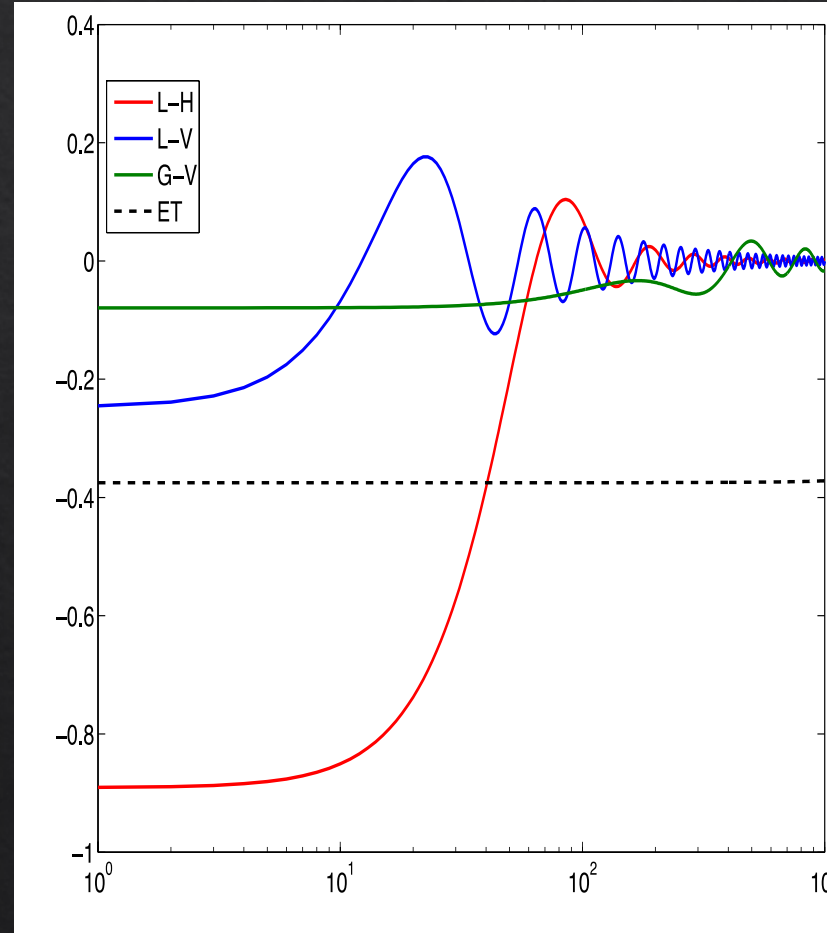
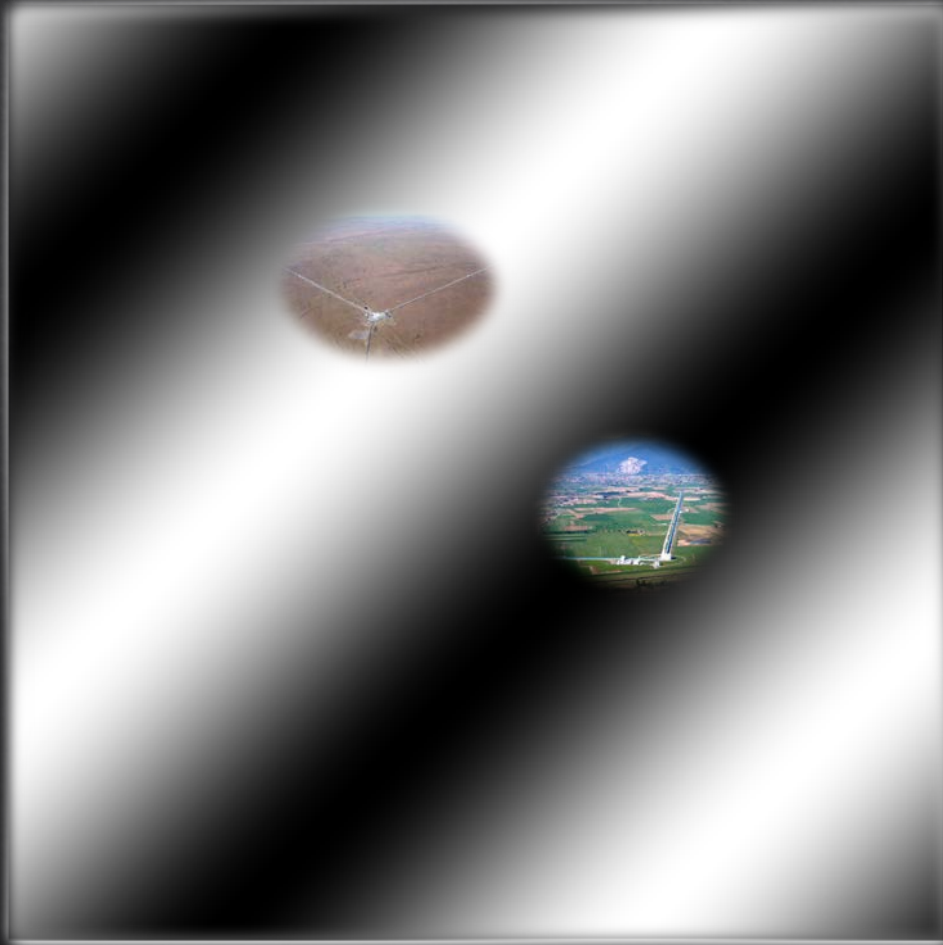
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SNR & overlap reduction function (a.k.a. coherence)

$$Y = \lambda \int df \frac{\tilde{h}_1^*(f) \tilde{h}_2(f) \gamma_{12}(f) S_h(f)}{S_{n,1}(f) S_{n,2}(f)} \quad \text{SNR}_Y^2 := \frac{\mu_Y^2}{\sigma_Y^2} = 2T \int_0^\infty S_h^2(f) \frac{\gamma_{12}^2(f)}{S_{n,1}(f) S_{n,2}(f)} df$$

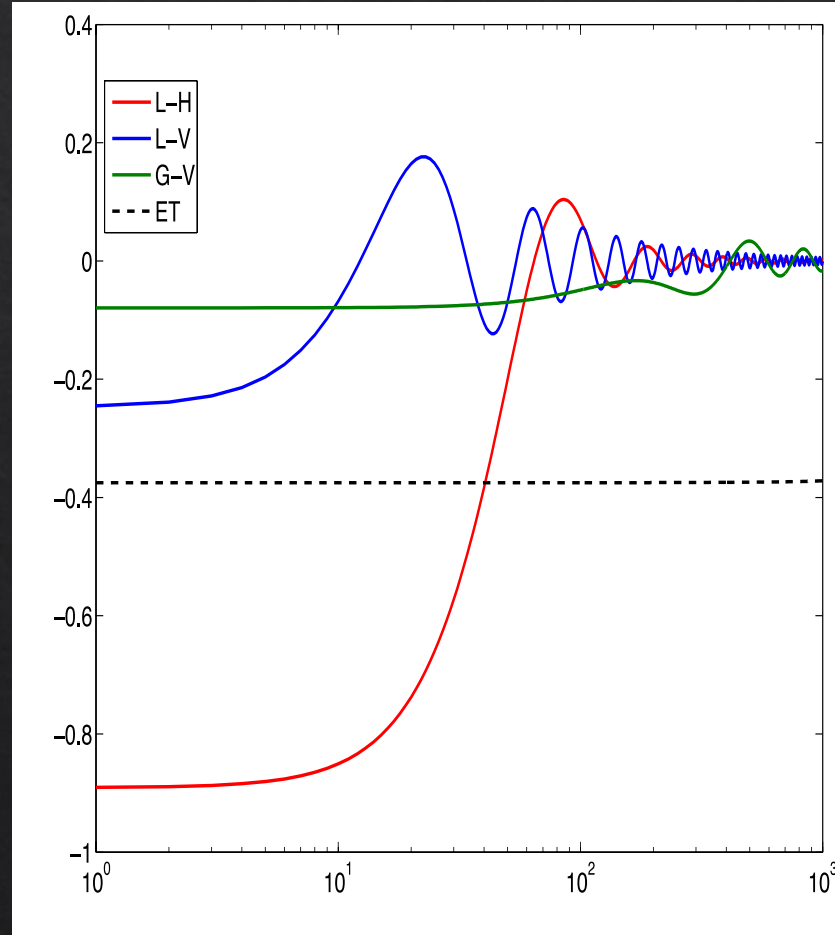
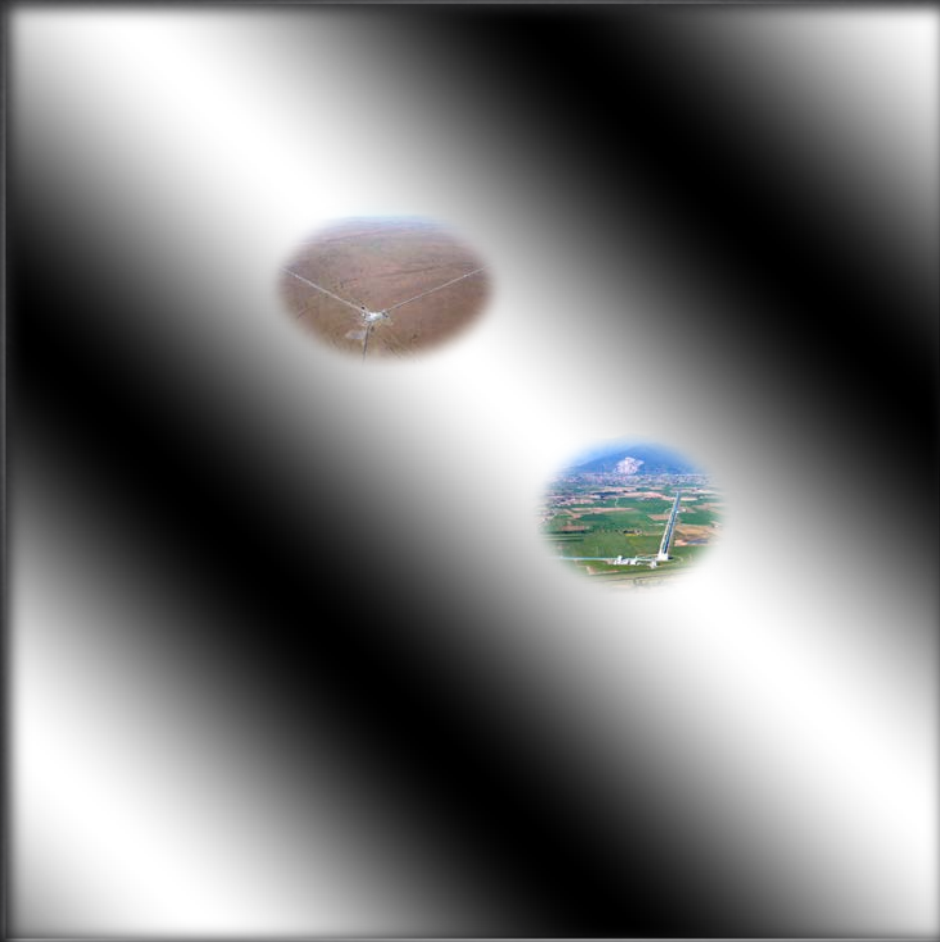


Optimal statistic takes into account:

- The spectral distribution of the noise;
- The spectral distribution of the signal;
- The loss of coherence due to the fact that different detectors are coupled to a set of GW in different ways.

Overlap reduction function (a.k.a. coherence)

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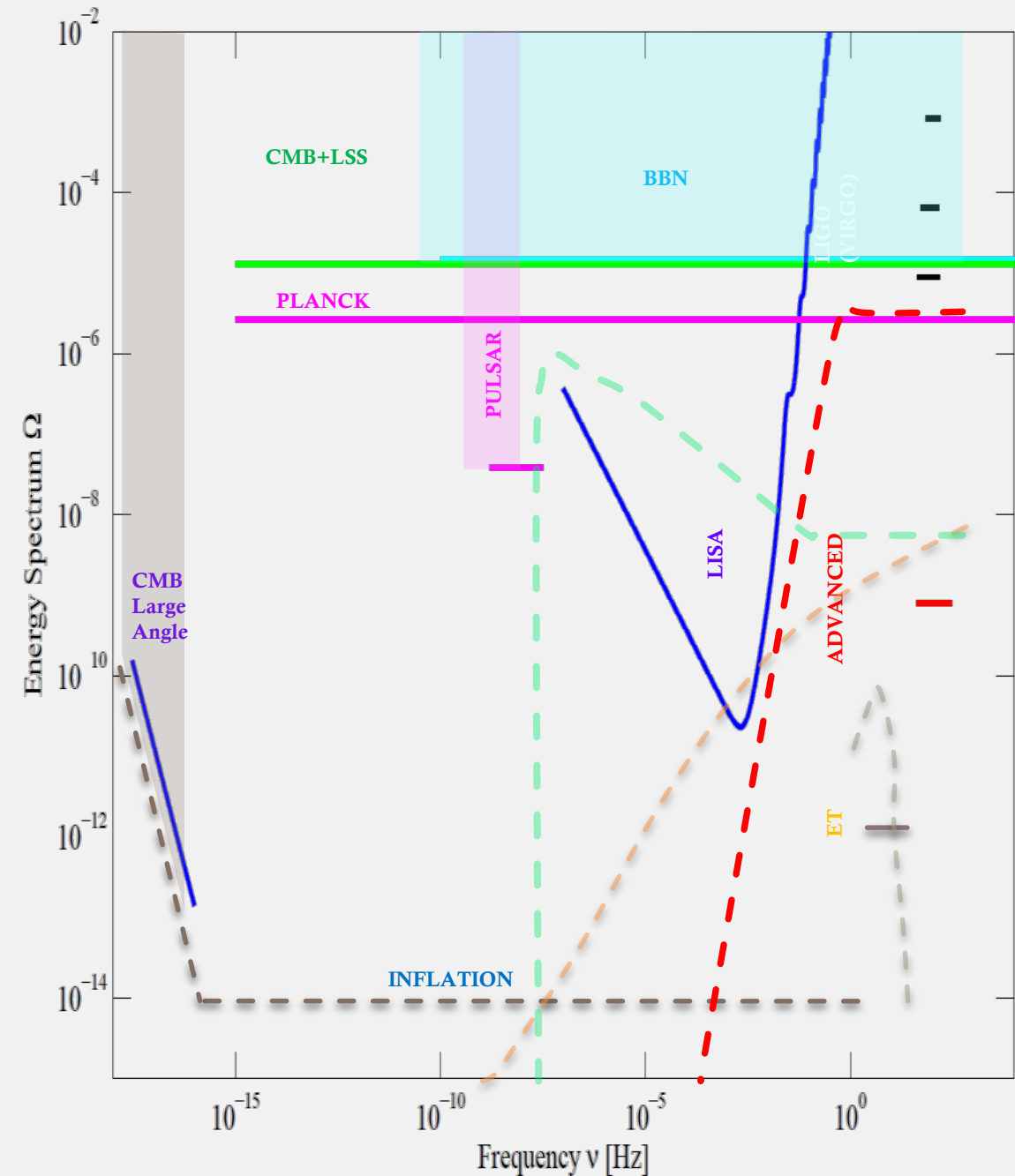
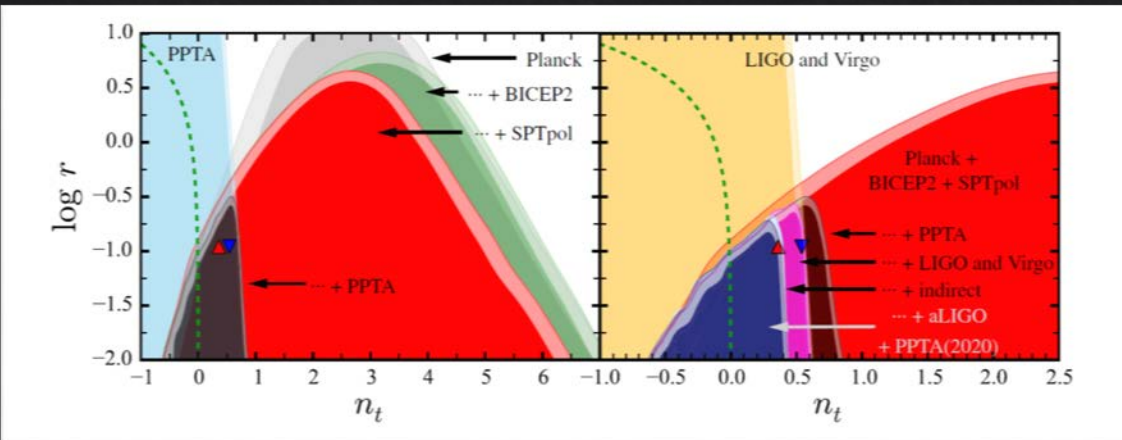
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Upper limits & expected sensitivities

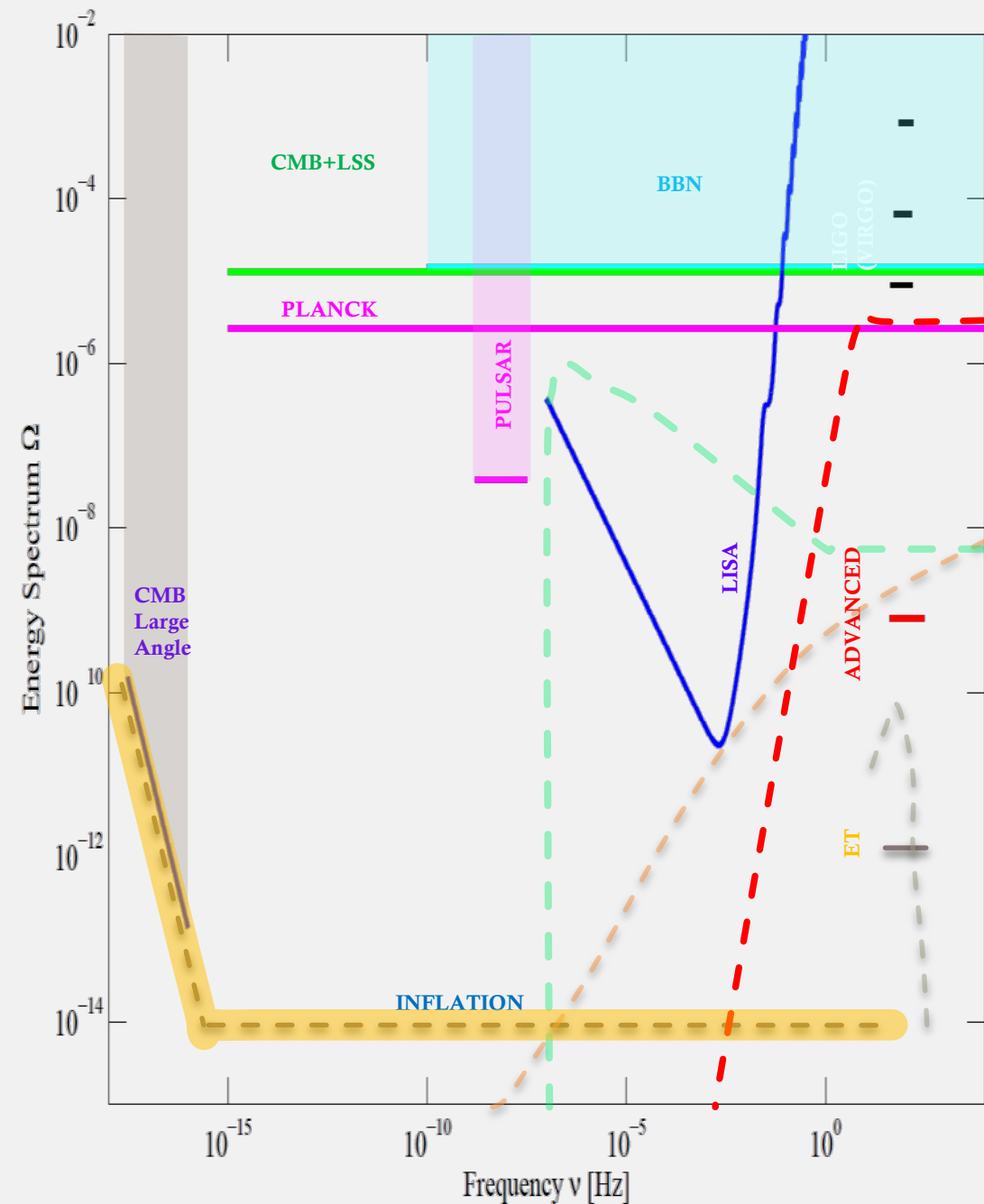
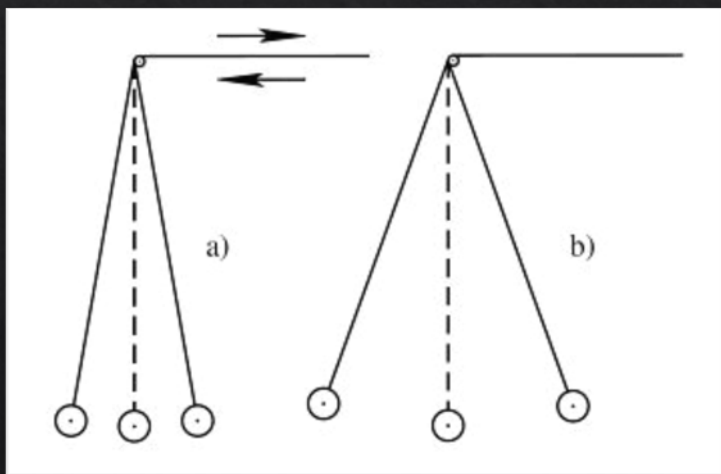
- **Big-Bang Nucleosynthesis model and observations** constrain the total GW energy at the time of BBN (integral bound)
- Similar bound from **CMB observations**
- Too much GW gives too much **large angle anisotropy** by the Sachs Wolfe effect
- Signal from **millisecond pulsars** works as a (big) interferometer

Complementary constraint in very different frequency regions: see [Phys. Rev. X 6, 011035 \(2016\)](#)



Inflation

- Parametric amplification of vacuum quantum fluctuations
- Standard inflationary models are weakly dependent on frequency
- Tight bound: CMB large angle anisotropy
- Out of reach of advanced detectors by 5 orders of magnitude

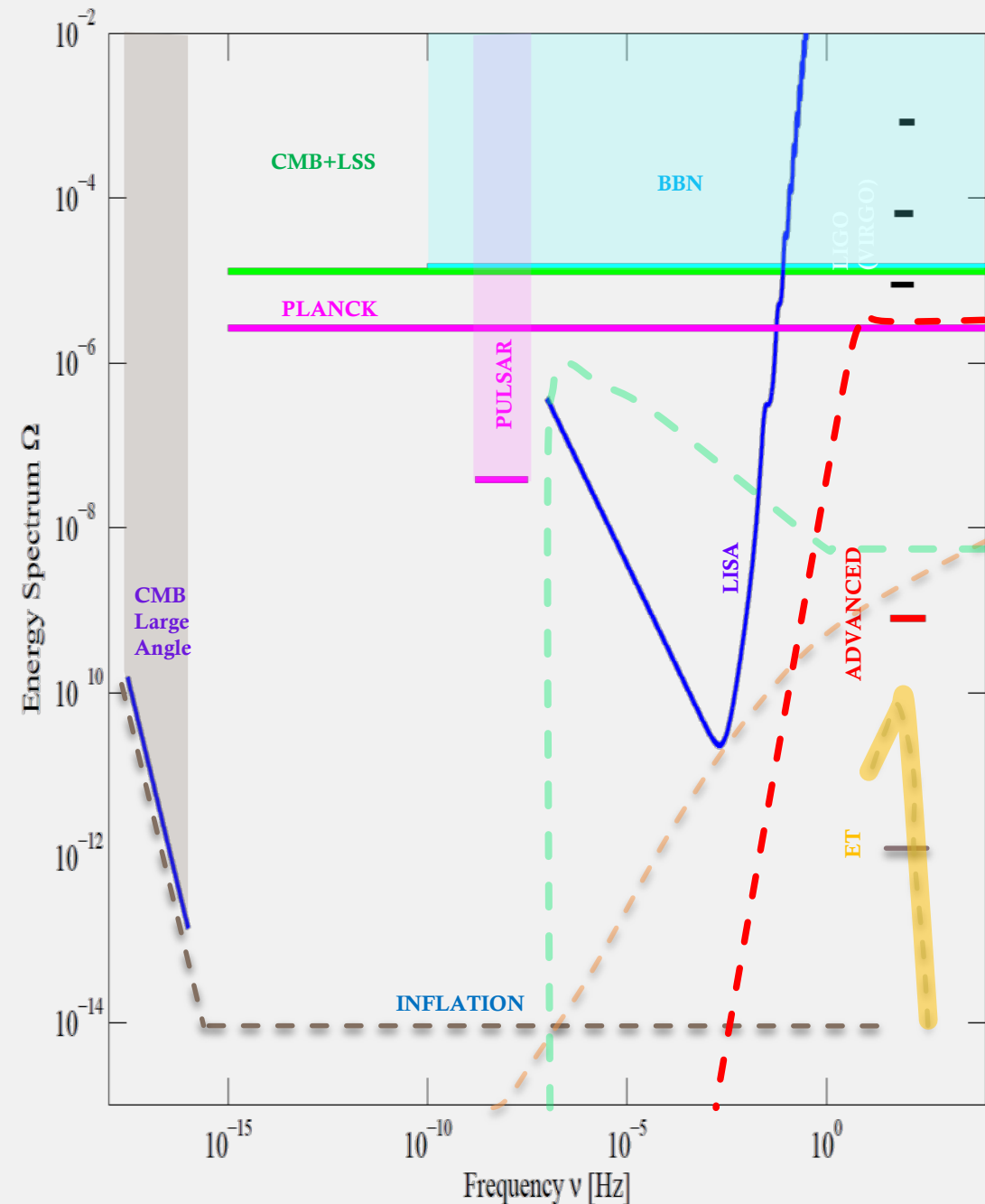


Resonant preheating



- During a resonant preheating phase at the end of the inflation, inflaton energy can be transferred efficiently to other particles
- Can produce a significant GW background
- Spectrum peak depends on energy scale (here 10 GeV, higher frequencies for a larger scale)

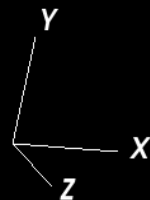
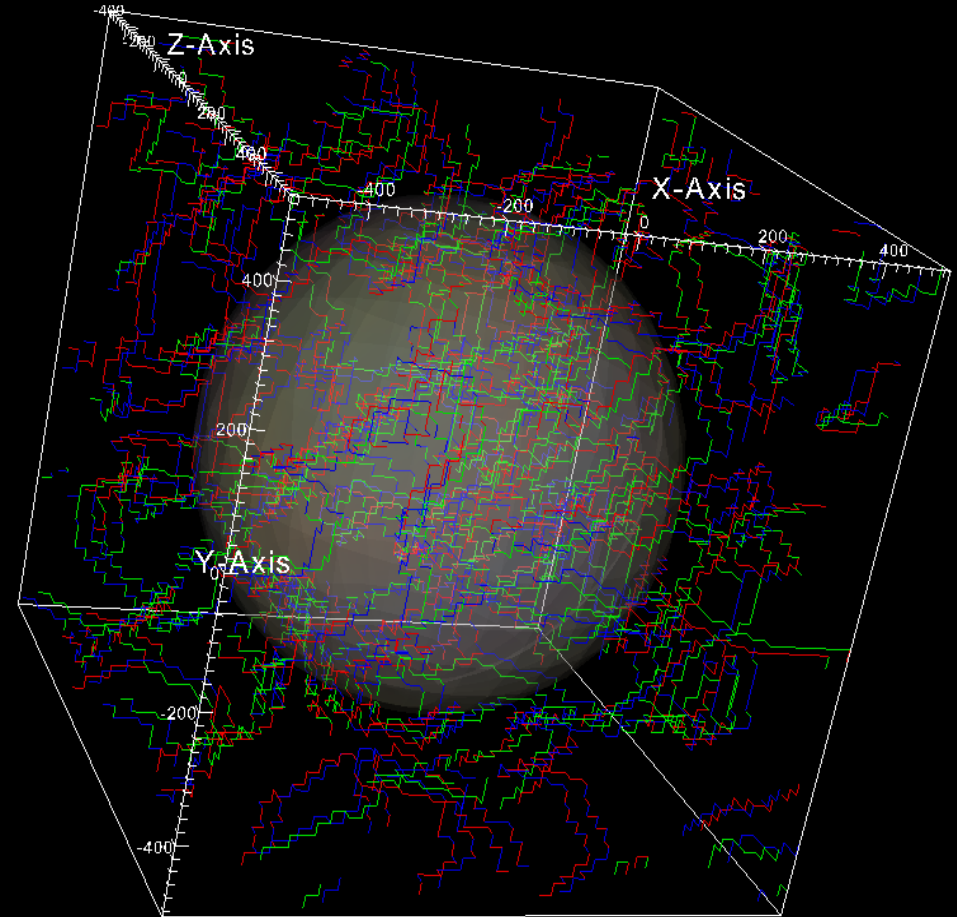
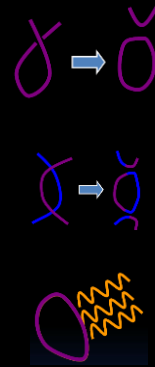
Easter & Lim, JCAP 0604, 010 (2006)
Easter et al., PLR 99, 221301 (2007)
Easter, Nucl. Phys. Proc. Suppl. 194, 33 (2009)



Cosmic (super)string models

- Dynamical network:
 - Strings entering in the horizon
 - Interconnection: loops generation
 - Radiation (GW and other fields): loop destruction
- Most efficient emission mechanisms: cusps and kinks
- Integrating over the whole universe leads to a GW background
- Large parameter space

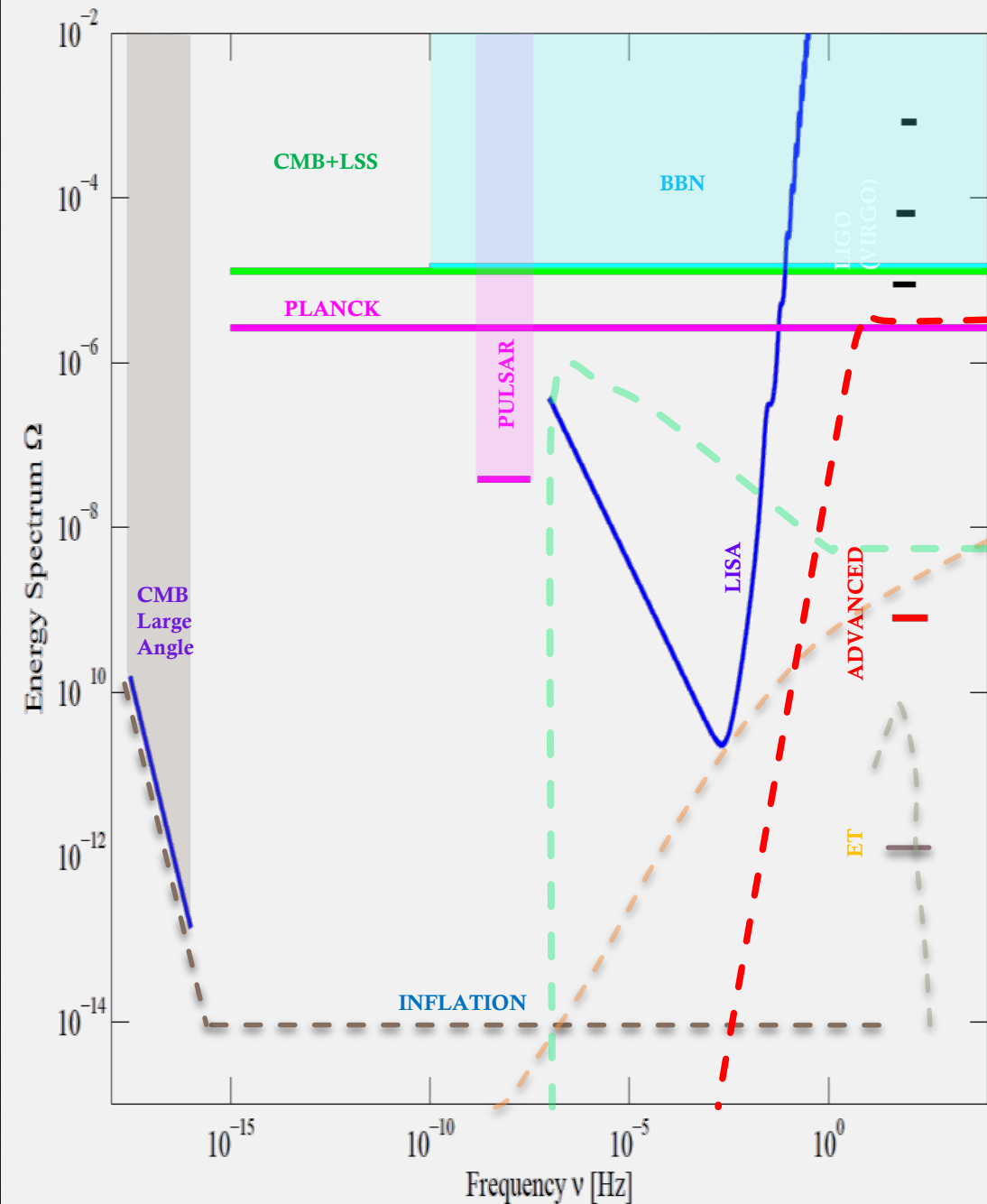
Damour & Vilenkin, PRL 85, 3761 (2000)
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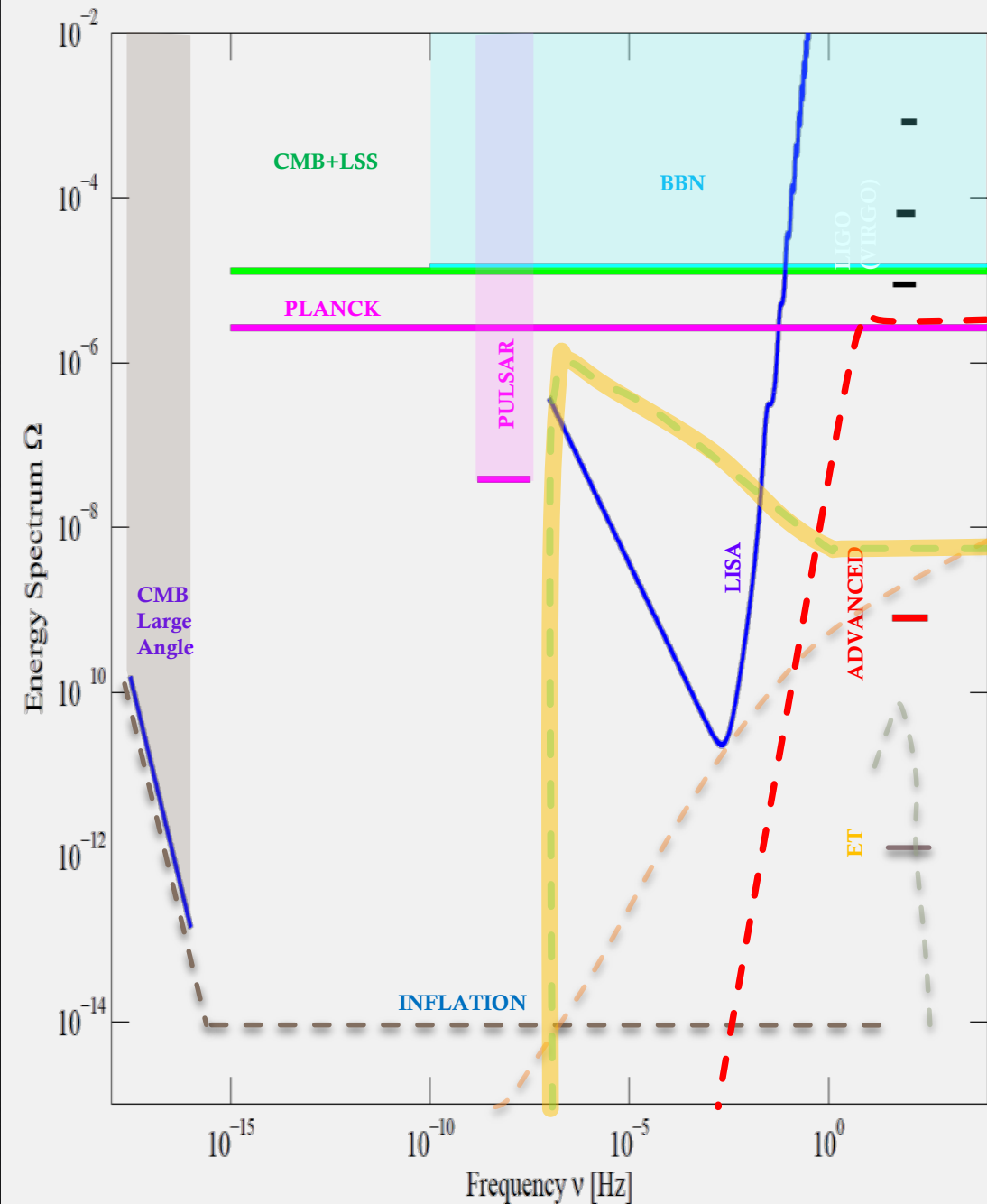
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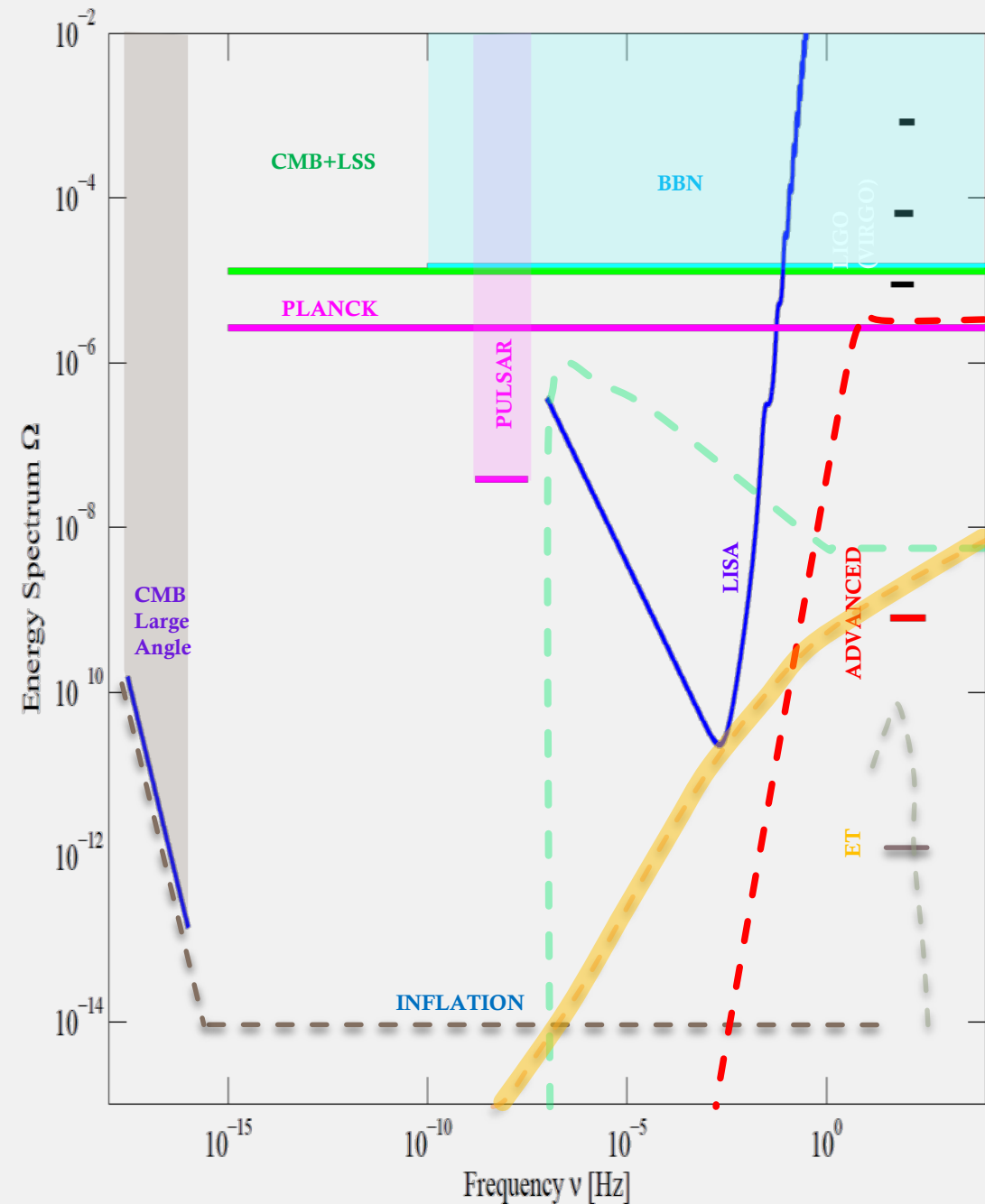
Damour & Vilenkin, PRL 85, 3761 (2000)
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Axion-based inflation models

- Models include axion-gauge couplings
- Gauge back-reaction on the inflaton extends inflation
- The late inflationary phase increases GW production at high frequencies

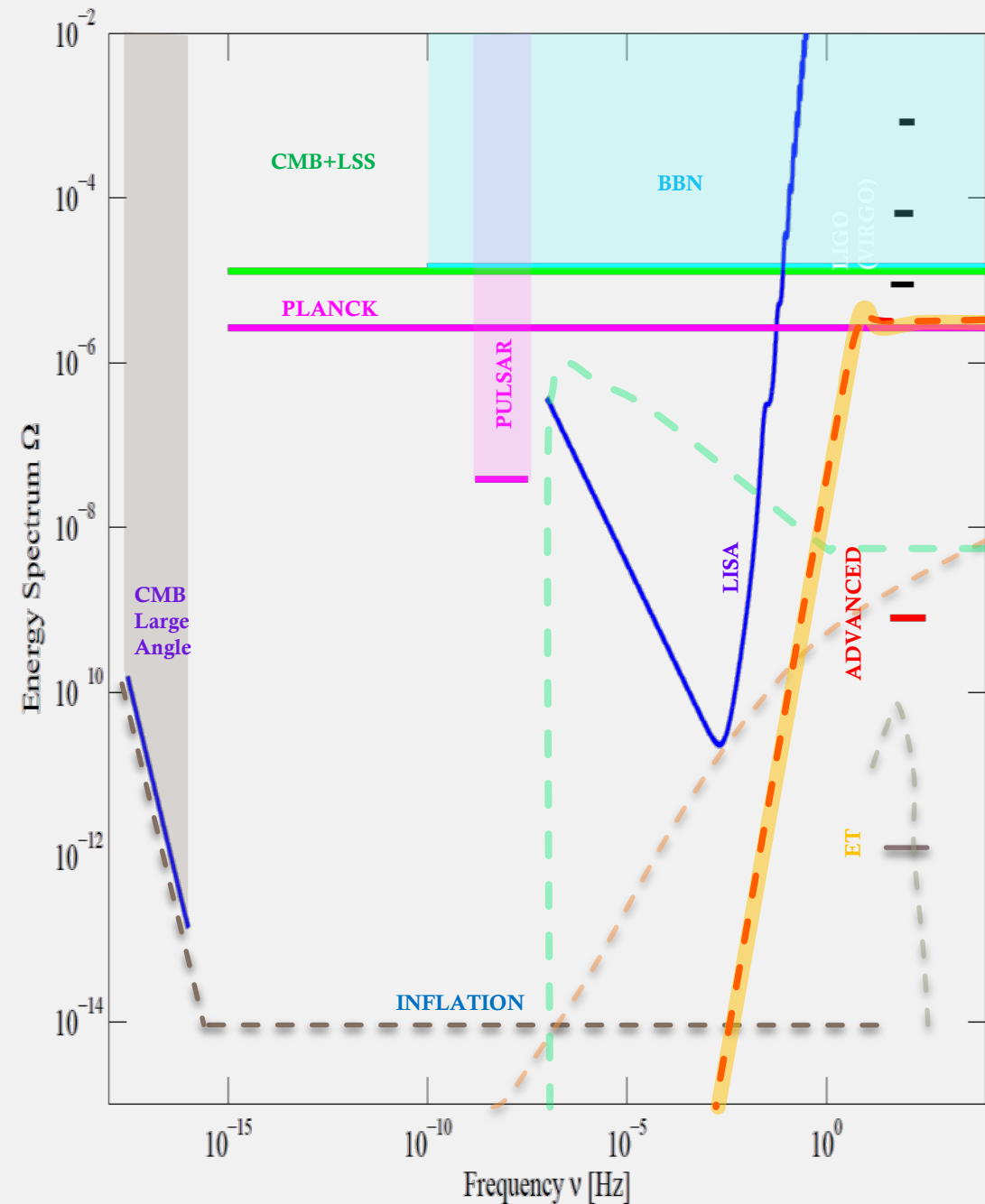
Barnaby, Pajer and Peloso - Phys. Rev. D 85, 023525



Pre-BB models

- Alternative cosmologies
- Evade the CMB large angle anisotropy bound
- Evade the BBN-CMB integral bound
- Can be significant at Virgo/LIGO frequencies

Gasperini & Veneziano, Phys. Rep. 373, 1 (2003)
Buonanno et al., PRD 55, 3330 (1997)



Isotropic upper limits

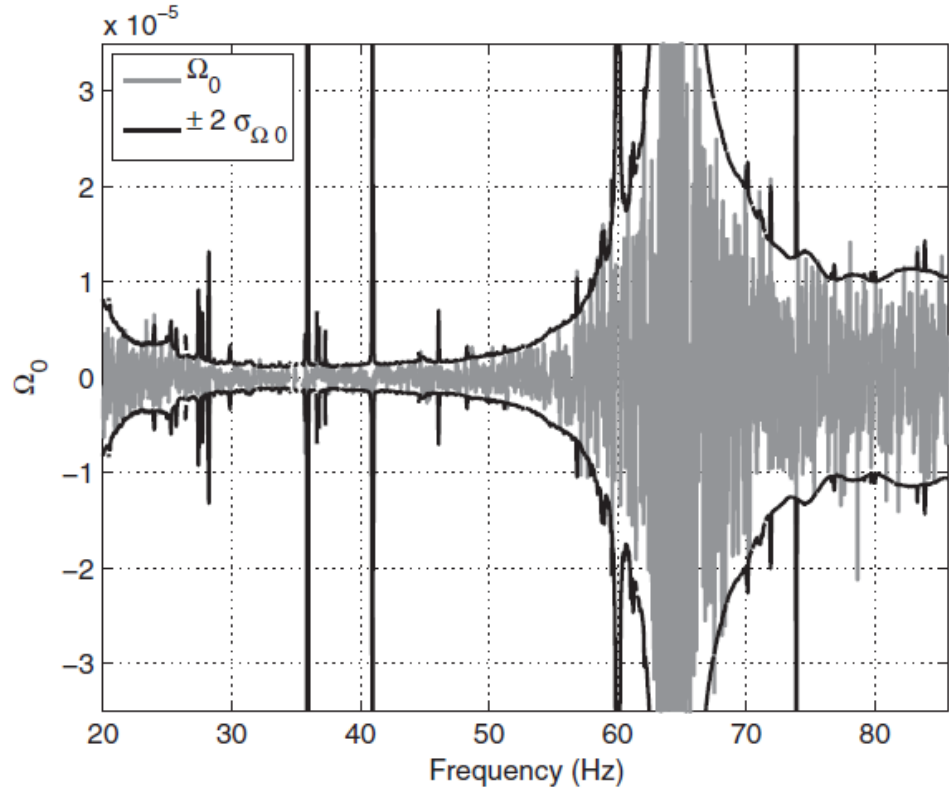


FIG. 1. We show the estimator for Ω_0 in each frequency bin, along with $\pm 2\sigma$ error bars, in the frequency band that contains 99% of the sensitivity for $\alpha = 0$. The loss of sensitivity at around 65 Hz is due to a zero in the overlap reduction function. There are several lines associated with known instrumental artifacts which do not lead to excess cross-correlation. The data are consistent with Gaussian noise, as described in the Results section.

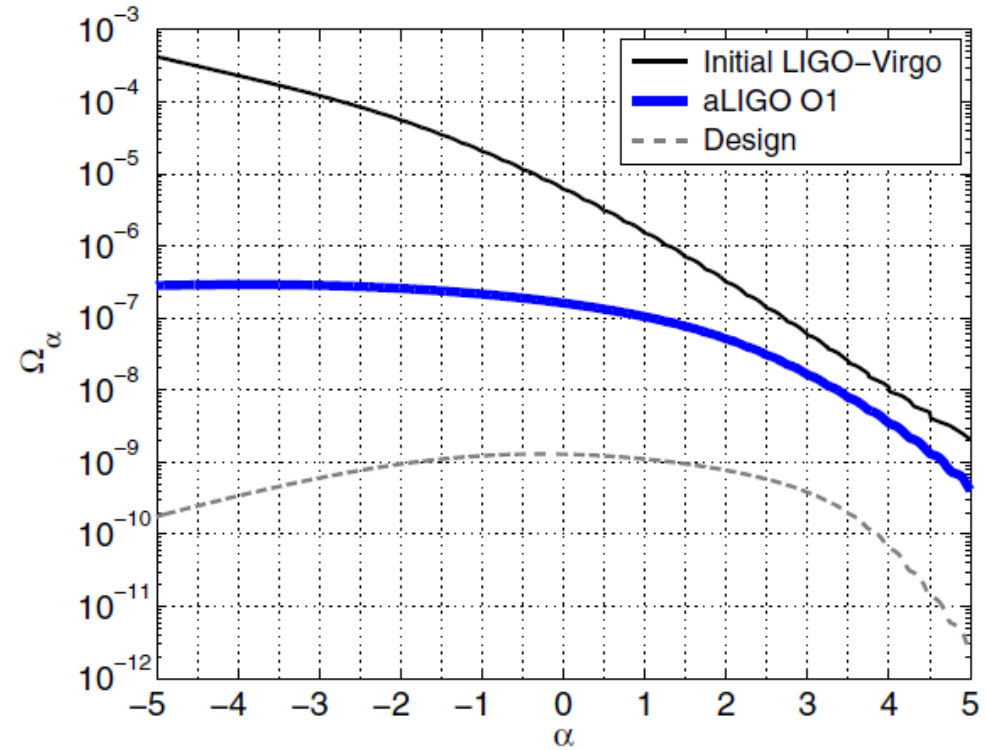


TABLE I. The frequency bands with 99% of the sensitivity are shown, along with the point estimate and standard deviation for the amplitude of the background, and 95% confidence level upper limits using O1 data for three values of the spectral index, $\alpha = 0, 2/3, 3$. We also show the previous upper limits using Initial LIGO-Virgo data.

Spectral index α	Frequency band with 99% sensitivity	Amplitude Ω_α	95% C.L. upper limit	Previous limits [36]
0	20–85.8 Hz	$(4.4 \pm 5.9) \times 10^{-8}$	1.7×10^{-7}	5.6×10^{-6}
2/3	20–98.2 Hz	$(3.5 \pm 4.4) \times 10^{-8}$	1.3×10^{-7}	–
3	20–305 Hz	$(3.7 \pm 6.5) \times 10^{-9}$	1.7×10^{-8}	7.6×10^{-8}

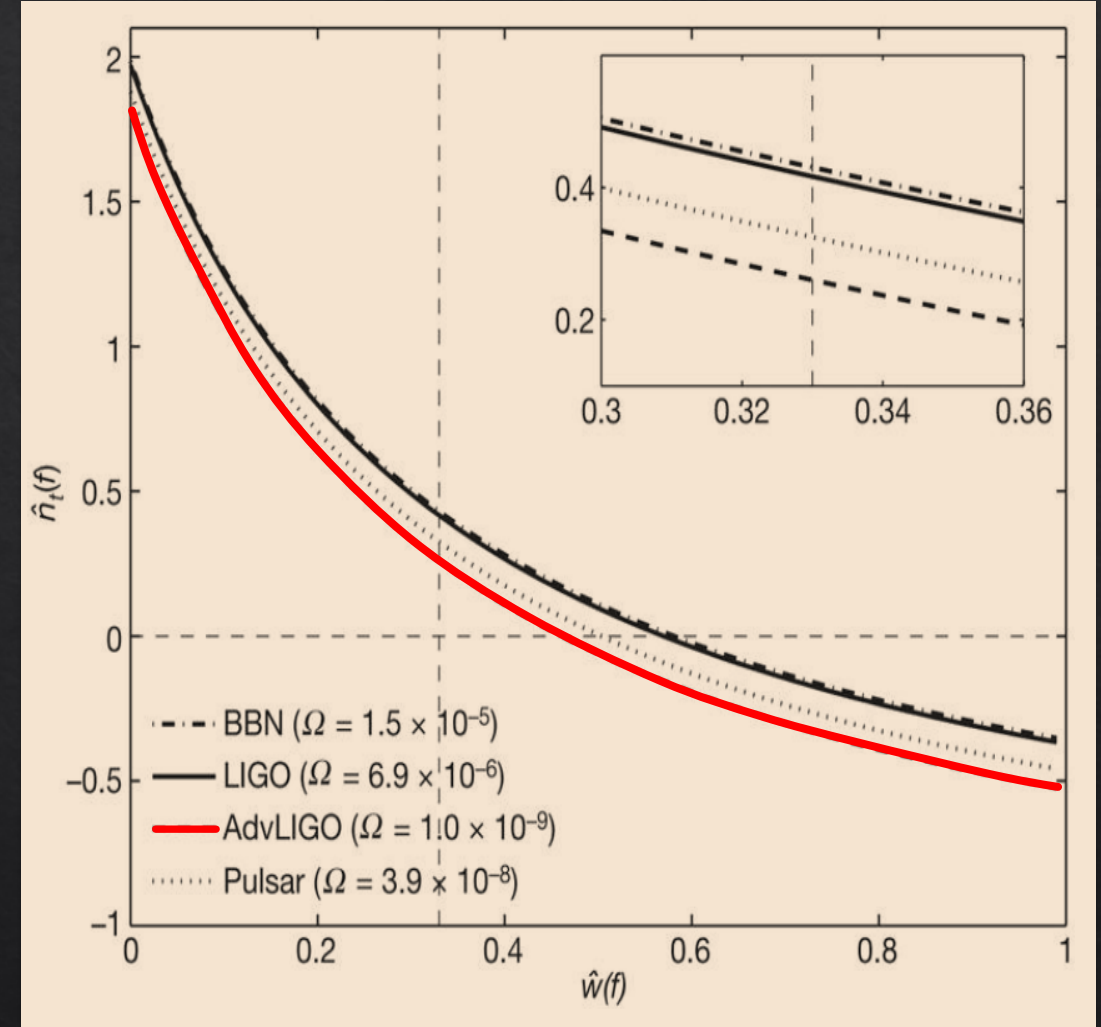
Constraint on early universe state equation

Spectrum parameterization:

$$\Omega_{GW}(f) = A f^{\hat{\alpha}(f)} f^{\hat{n}(f)} r$$

$$\hat{\alpha}(f) = 2 \frac{3\hat{w}(f) - 1}{3\hat{w}(f) + 1}$$

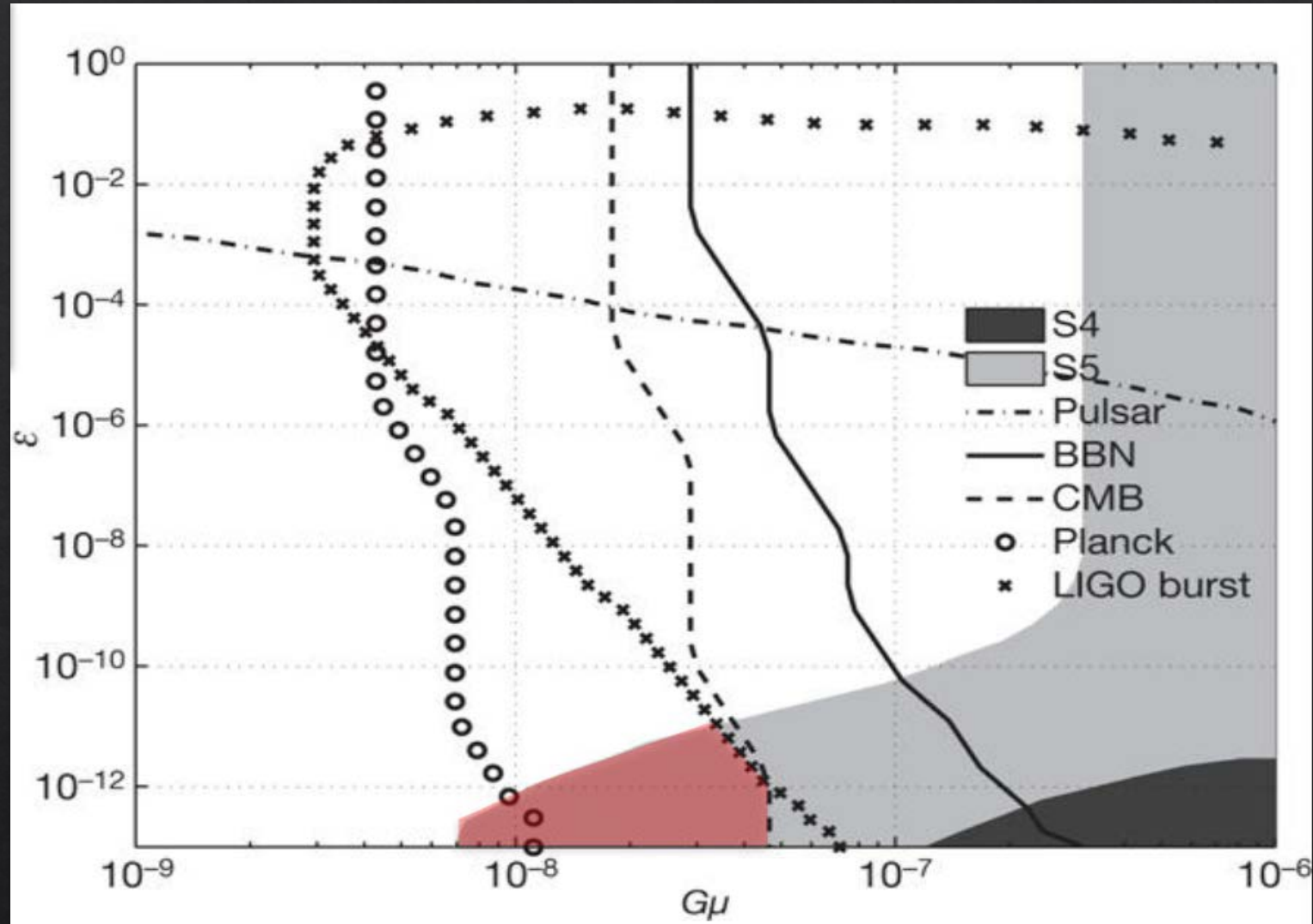
- $\hat{n}_t(f)$ effective tensor tilt parameter
- r ratio of tensor and scalar perturbation amplitudes (here $r = 0.1$)
- $\hat{w}(f)$ equation of state parameter



Constraint on cosmic strings models

- Network of cosmic strings parameterized by:
- String tension μ
- Reconnection probability p
- Loop size (parameterized by ε)
- $G\mu < 10^{-6}$ (CMB observations)
- ε unconstrained
- $10^{-4} < p < 1$ ($p = 10^{-3}$ here)

Region excluded: entire plane will be probed by advanced detectors.



Constraints on pre-Big-Bang models

Spectrum:

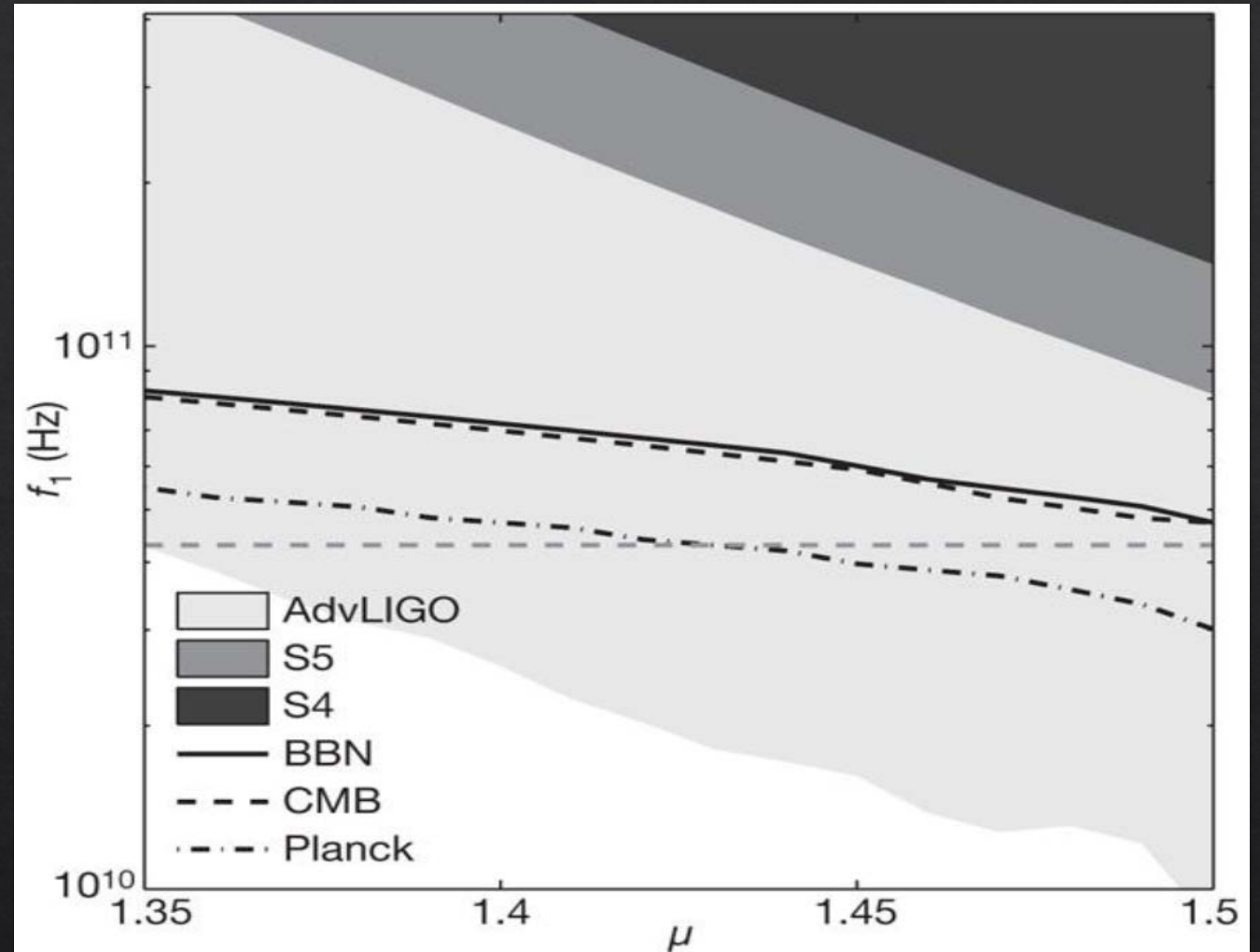
$$\Omega_{GW}(f) \propto f^3$$

below f_s , $f_s = 30$ Hz here

$$\Omega_{GW}(f) \propto f^{3-2\mu}$$

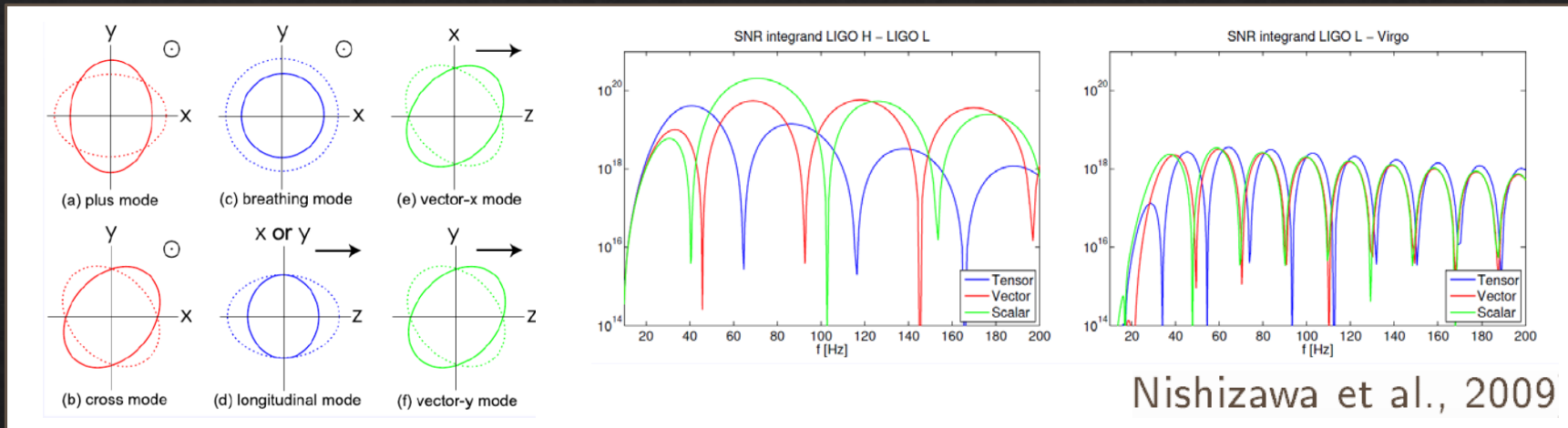
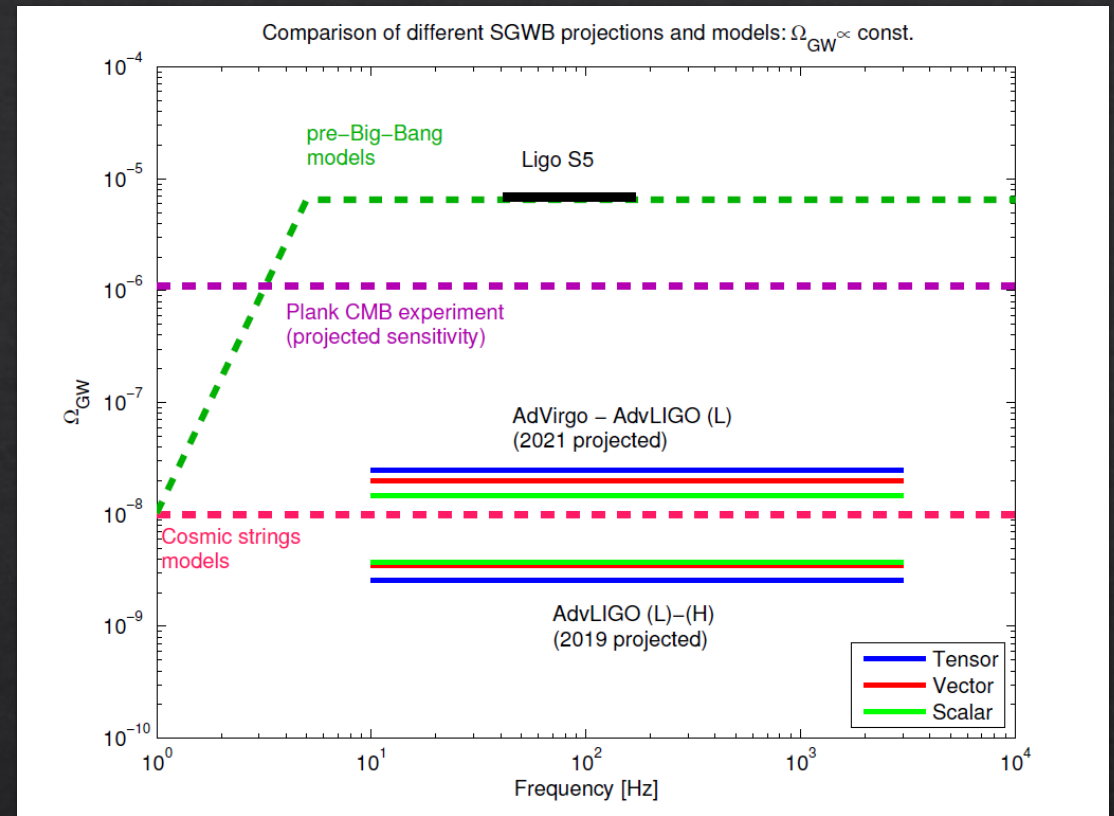
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f_1 : cut off frequency
(a factor 10 from 4.3×10^{10})

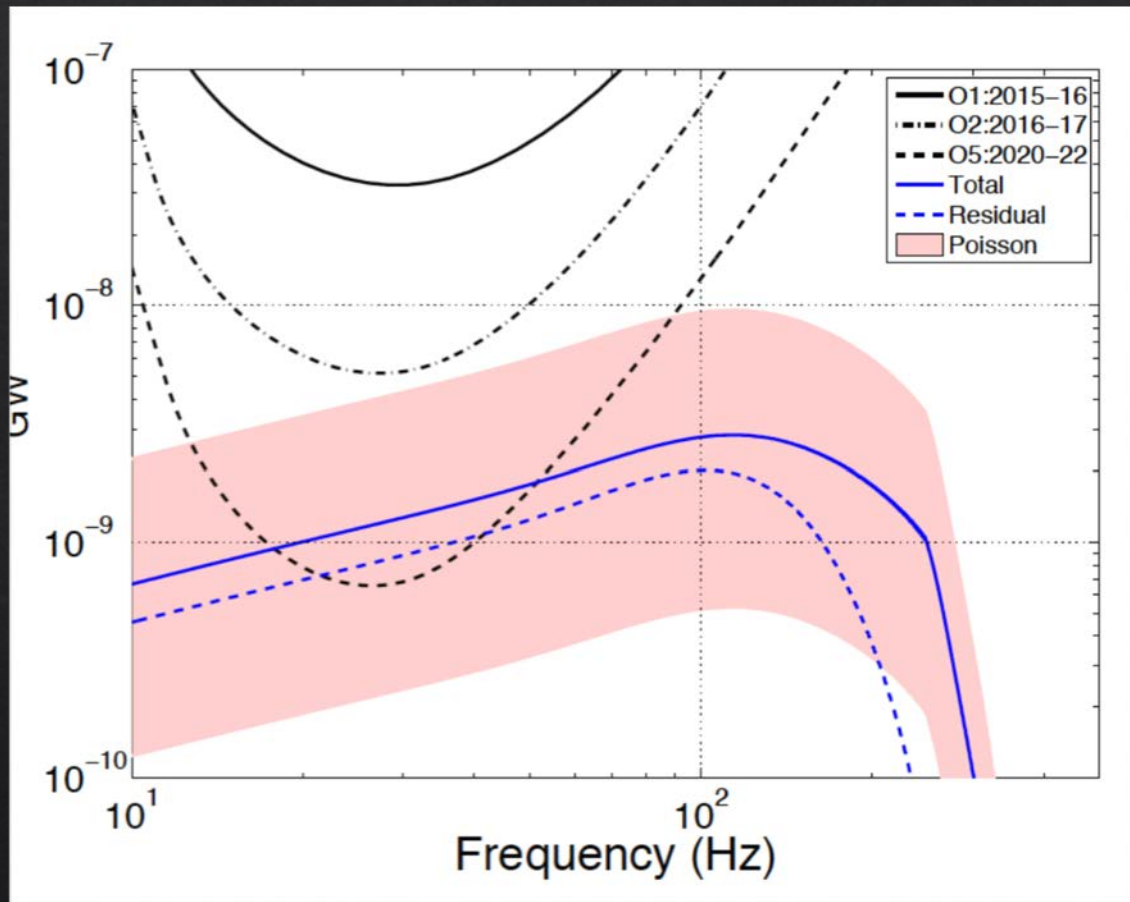


Non standard polarizations

- Looking at alternative theories of gravity
- Indirect constraints on scalar modes
 - Binary pulsars
 - WMAP
- Need many detector pairs to disentangle different contributions



Conclusions



- ◇ Stochastic background not yet measured
- ◇ Direct detection not excluded with some model in the advanced era, however
 - ◇ Astrophysical models can be detected also;
 - ◇ A cosmological contribution must be disentangled.
- ◇ Several interesting constraints on models start to be available, and will improve;
- ◇ In principle, possibility of testing alternate gravity models.
- ◇ Not discussed here:
 - ◇ Directional searches (maps of anisotropies)