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# Breaking the scale invariance of primordial spectrum in Horava-Lifshitz cosmology

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Work with

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Phys. Rev. D, to appear; arXiv: 0907.1549

# Quantum gravity & Cosmology

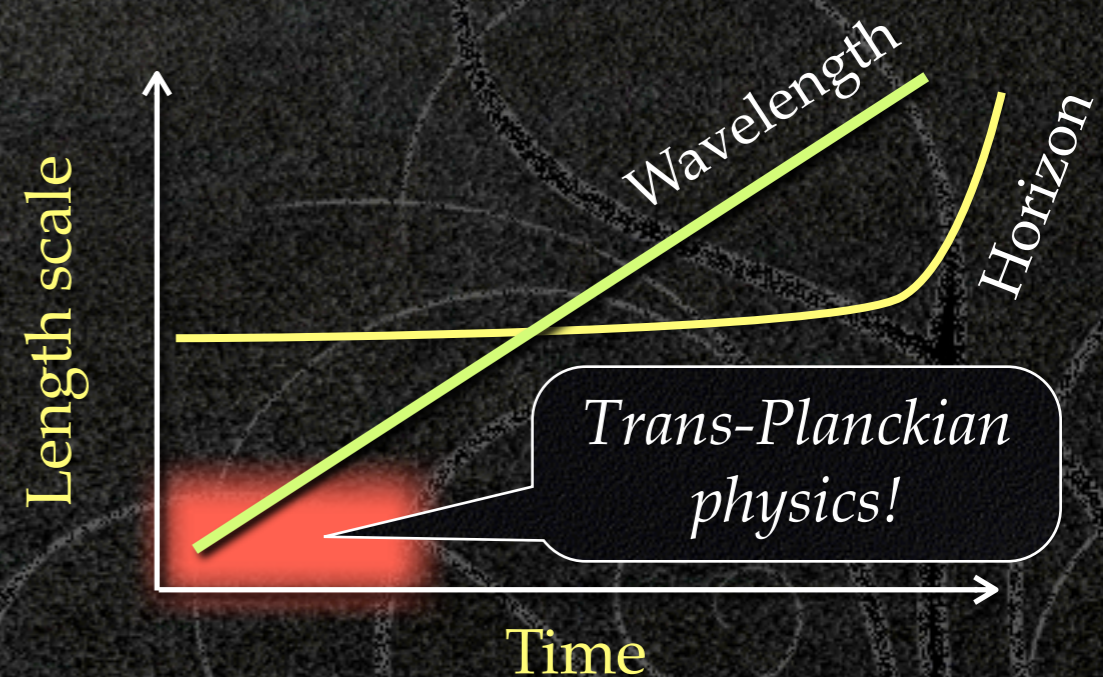


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- **Cosmological window on quantum gravity phase**
  - Fluctuations are stretched by the expansion of the universe
  - **Can probe high energy physics with cosmological perturbations**



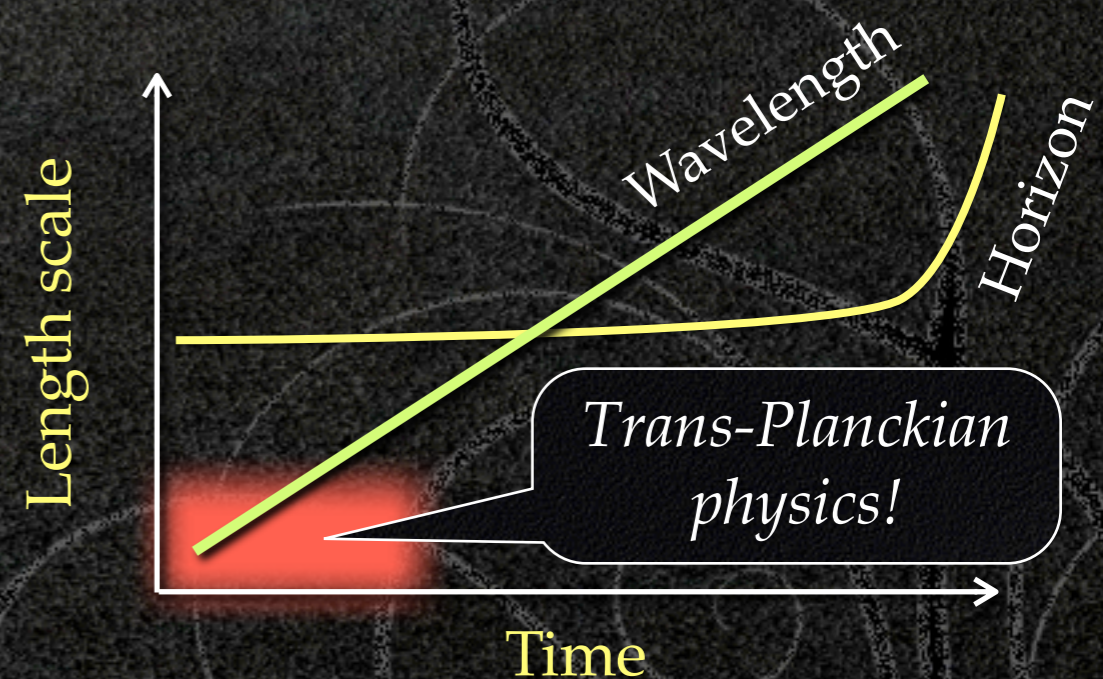
# Quantum gravity & Cosmology

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*Why don't we study cosmological perturbations in Horava's gravity?*

- Fluctuations are stretched by the expansion of the universe
- **Can probe high energy physics with cosmological perturbations**

Gao, Wang, Brandenberger & Riotto;  
Wang & Maartens; Koh;  
T.K., Urakawa & Yamaguchi, 2009; .....



# Dispersion relation modified in the UV

## Essence of Horava's gravity

Horava, 2009

- Action is invariant under the Lifshitz-like anisotropic scaling:

$$t \rightarrow \ell^z t, \quad \vec{x} \rightarrow \ell \vec{x} \quad \text{with } \mathbf{z} = \mathbf{3} \text{ in the UV}$$

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- Consider a mode equation (for GWs and a scalar field fluctuation)

$$\varphi_k'' + \left[ k_{\text{eff}}^2(\eta) - \frac{a''}{a} \right] \varphi_k = 0$$

- w/ dispersion relation  $\omega^2 = k_{\text{eff}}^2(\eta) / a^2$



# Scale-invariant spectrum from $z = 3$

Kiritsis, Kofinas, 2009;  
Mukohyama, 2009

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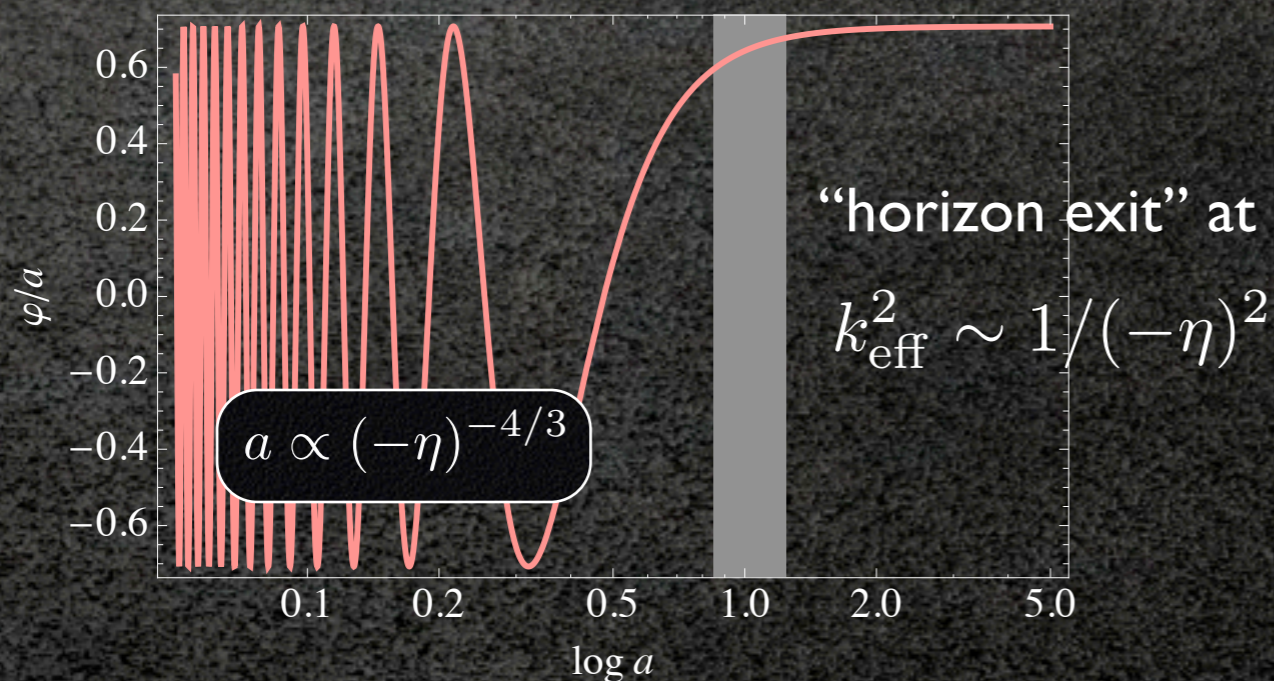
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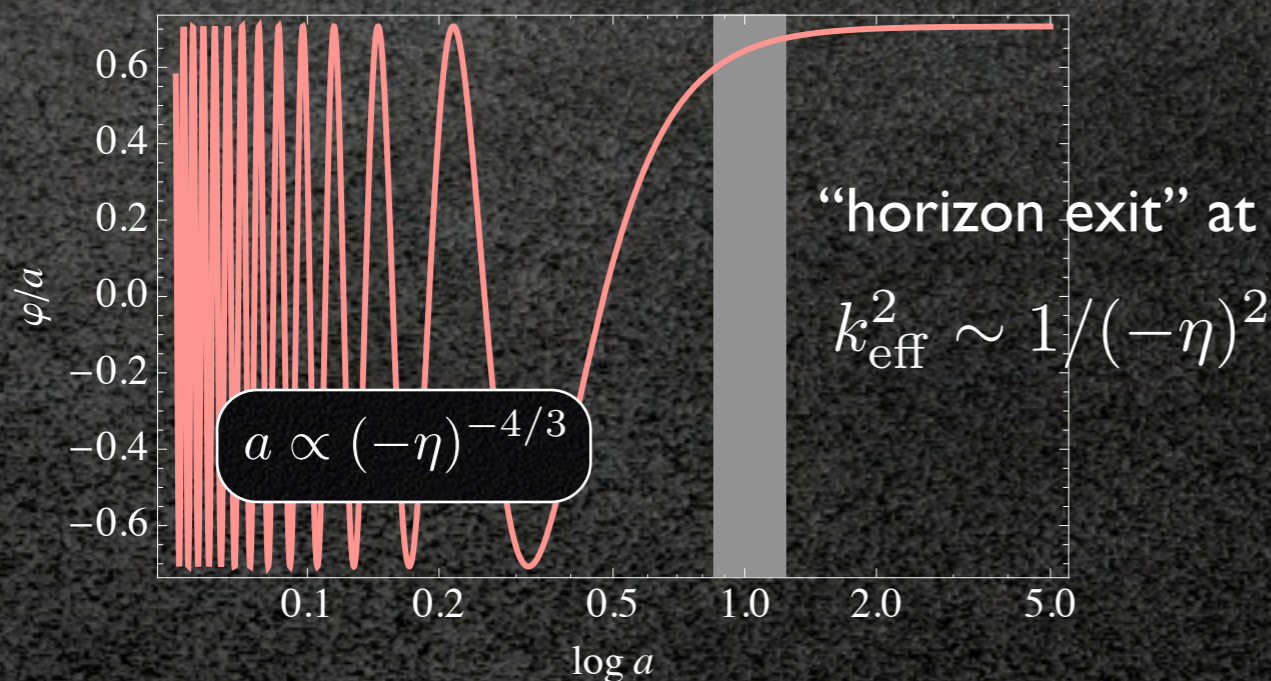
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Independent of the  
cosmic expansion rate,  $H$

**Exactly** scale-invariant even  
for non-de Sitter background

# How to break scale invariance

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**z=3 UV**

**z=1 IR**

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**Spectral tilt**

- $k^5$ : parity violating term for GWs      Takahashi and Soda, 2009

Non-standard dispersion relation  
implied by Horava's quantum gravity



# Non-standard dispersion relation implied by Horava's quantum gravity

This work:

Evaluate the spectral index  
using the **uniform approximation**

K. Yamamoto, T.K., and G. Nakamura  
Phys. Rev. D, to appear; arXiv: 0907.1549

# Uniform approximation

Olver, 1997; Martin, Schwarz, 2003; Habib *et al.* 2002; 2004

- Based essentially on the WKB approximation
- Can fit the exact behavior of the solution accurately
- Suited to calculate the  **$k$ -dependence** of the solution
  - **w/o slow-roll assumptions**

# Uniform approximation

Olver, 1997; Martin, Schwarz, 2003; Habib *et al.* 2002; 2004

- Mode equation  $\varphi_k'' = \left[ g(\eta) - \frac{1}{4\eta^2} \right] \varphi_k$   $g(\eta) := -k_{\text{eff}}^2(\eta) + \frac{a''}{a} + \frac{1}{4\eta^2}$

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- admits the approximate solution in terms of **Airy functions**

$$\varphi_k = \sqrt{\frac{\pi}{2}} \left[ \frac{\zeta(\eta)}{g(\eta)} \right]^{1/4} [\text{Ai}(\zeta) - i \text{Bi}(\zeta)]$$

- where  $\zeta(\eta) := \pm \left[ \pm \frac{3}{2} \int_{\bar{\eta}}^{\eta} \sqrt{\pm g(\eta')} d\eta' \right]^{2/3}$   $\begin{array}{l} + \text{ for } \eta > \bar{\eta} \\ - \text{ for } \eta < \bar{\eta} \end{array}$

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- The solution satisfies the **Wronskian normalization condition** and has the asymptotic **early-time** behavior:

$$\varphi_k \approx \exp\left[-i \int_{\bar{\eta}}^{\eta} k_{\text{eff}}(\eta') d\eta'\right] / \sqrt{2k_{\text{eff}}}$$

# Uniform approximation contd.

- Asymptotic **late-time** behavior:  $\varphi_k \approx -\frac{i}{\sqrt{2}} [g(\eta)]^{-1/4} \exp \left[ \int_{\bar{\eta}}^{\eta} \sqrt{g(\eta')} d\eta' \right]$
- Uniform approximation formula for the **spectral index**:

$$n_T = 3 + \lim_{\eta \rightarrow 0^-} 2 \frac{d}{d \ln k} \int_{\bar{\eta}(k)}^{\eta} \sqrt{g(\eta')} d\eta'$$

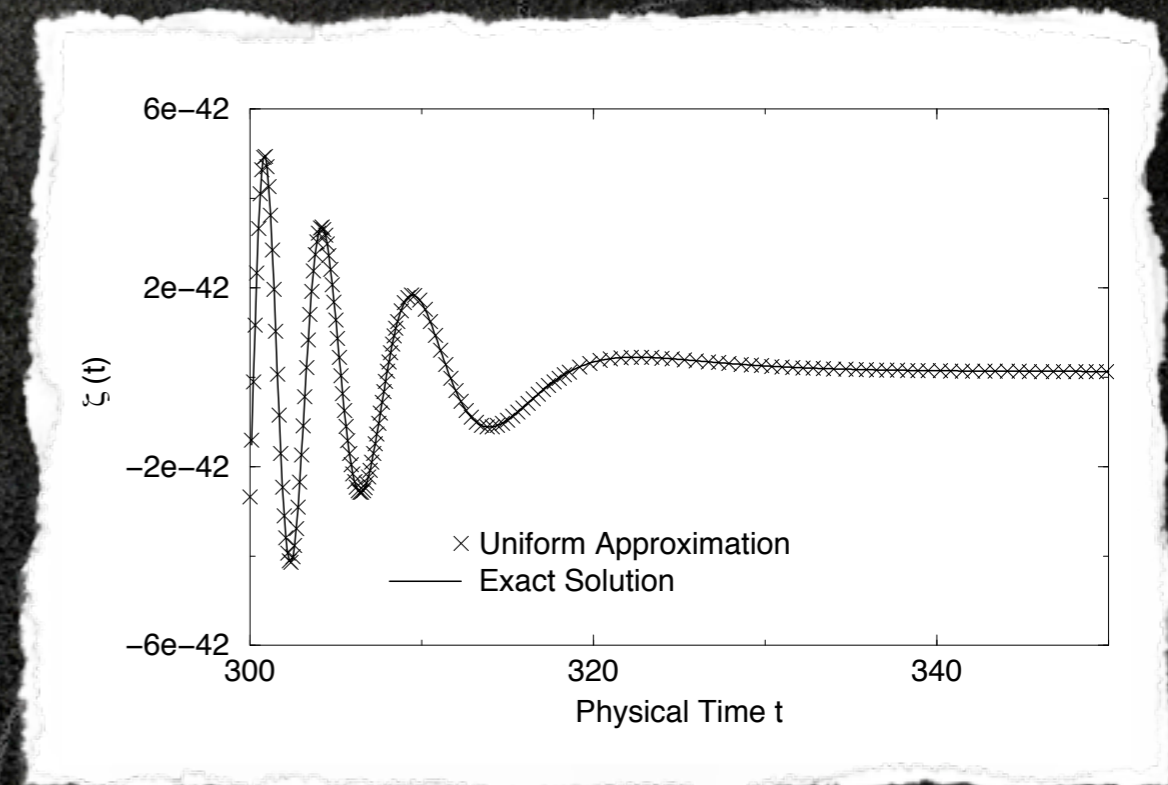
# Uniform approximation contd.

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- So far, the uniform approximation has been used only for inflation models with the standard dispersion relation,  $\omega^2 = k^2/a^2$

Habib *et al.* 2002



# Application I

- Power-law inflation  $a \propto t^{1+n}$  with  $k_{\text{eff}}^2 = k^2 \left( \frac{k^2}{M^2 a^2} \right)^\mu$



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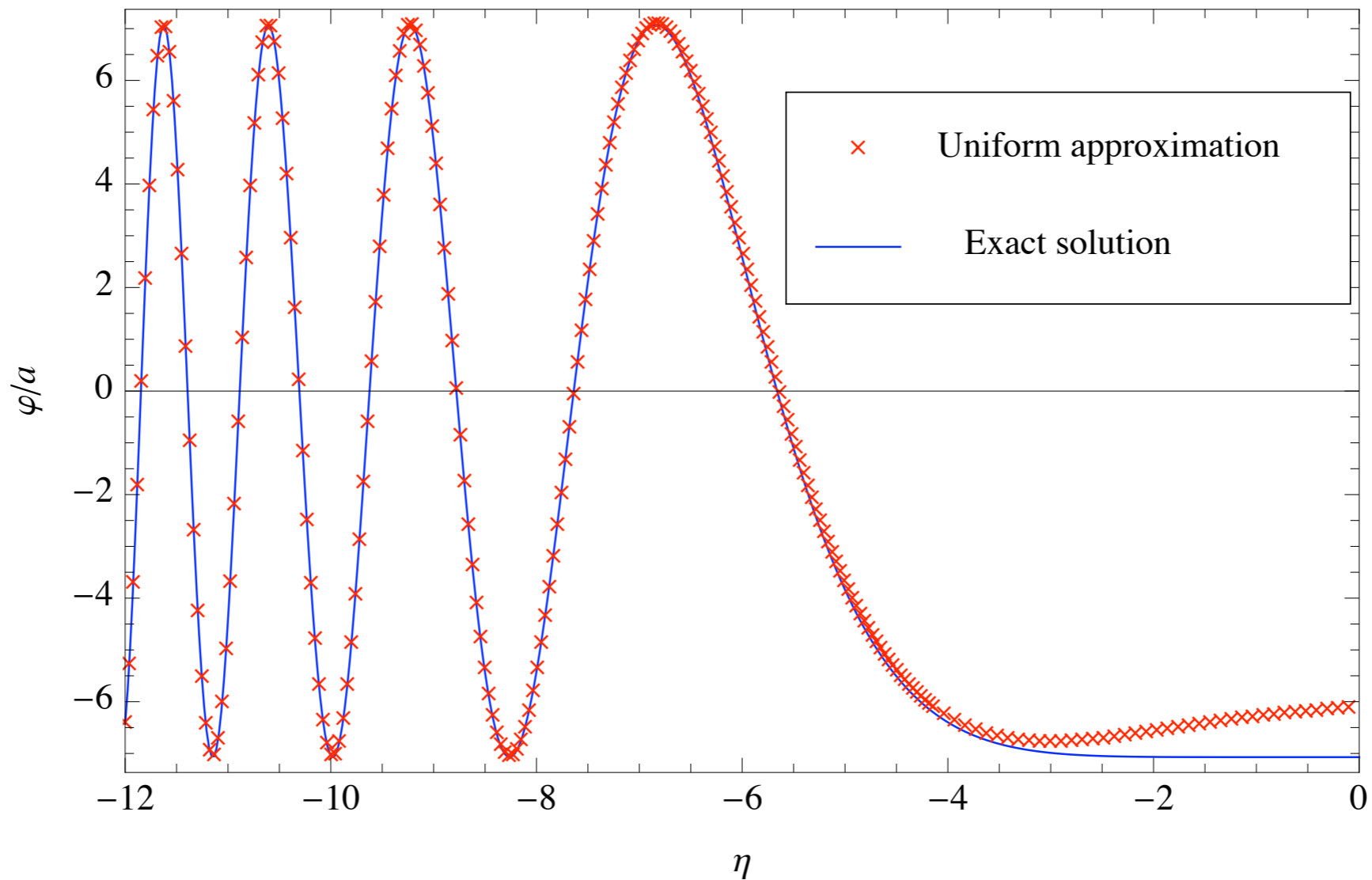
# Application I

- Power-law inflation  $a \propto t^{1+n}$  with  $k_{\text{eff}}^2 = k^2 \left( \frac{k^2}{M^2 a^2} \right)^\mu$
- Can be compared with the exact solution
- **k-independent** small correction to the amplitude:  $|1 - \mathcal{C}| = \mathcal{O}(0.1)$

$$|\varphi_{\text{uniform}}| = \mathcal{C} |\varphi_{\text{exact}}|$$
$$\mathcal{C} := \sqrt{2\pi} e^{-\nu} \nu^{\nu-1/2} / \Gamma(\nu)$$
$$\nu := \frac{1 + 3n/2}{\mu + n(1 + \mu)}$$

- **Exact result reproduced for the spectral index**

$$n_T = \frac{\mu - 2}{\mu + n(1 + \mu)}$$



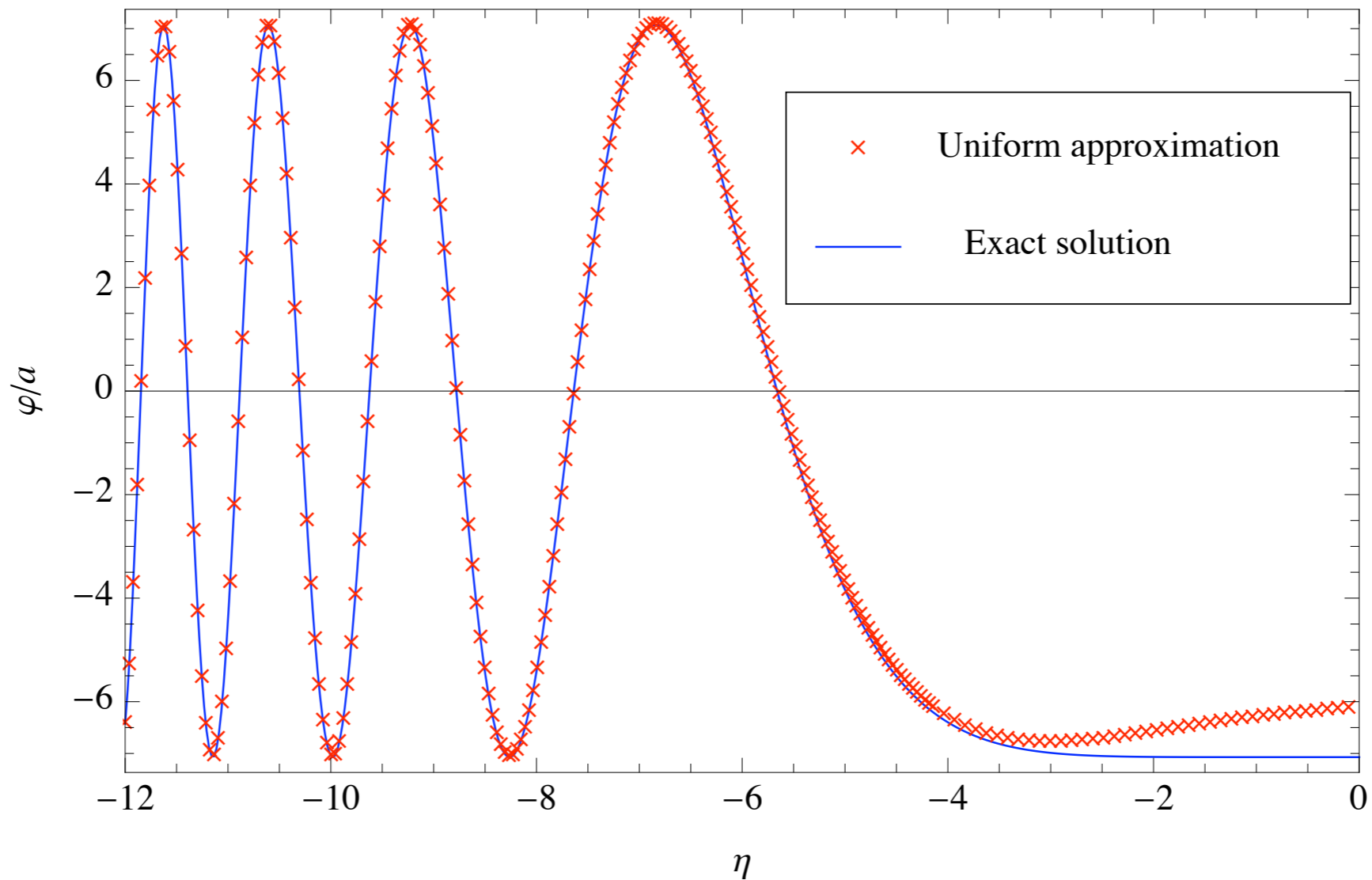
$$\mu + n(1 + \mu)$$

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$\mu$

$\mathcal{O}(0.1)$



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Scale-invariant for

$$\mu = 2 \quad (k_{\text{eff}}^2 \propto k^6)$$

Consistent with Mukohyama 2009; Cai and Zhang, 2009

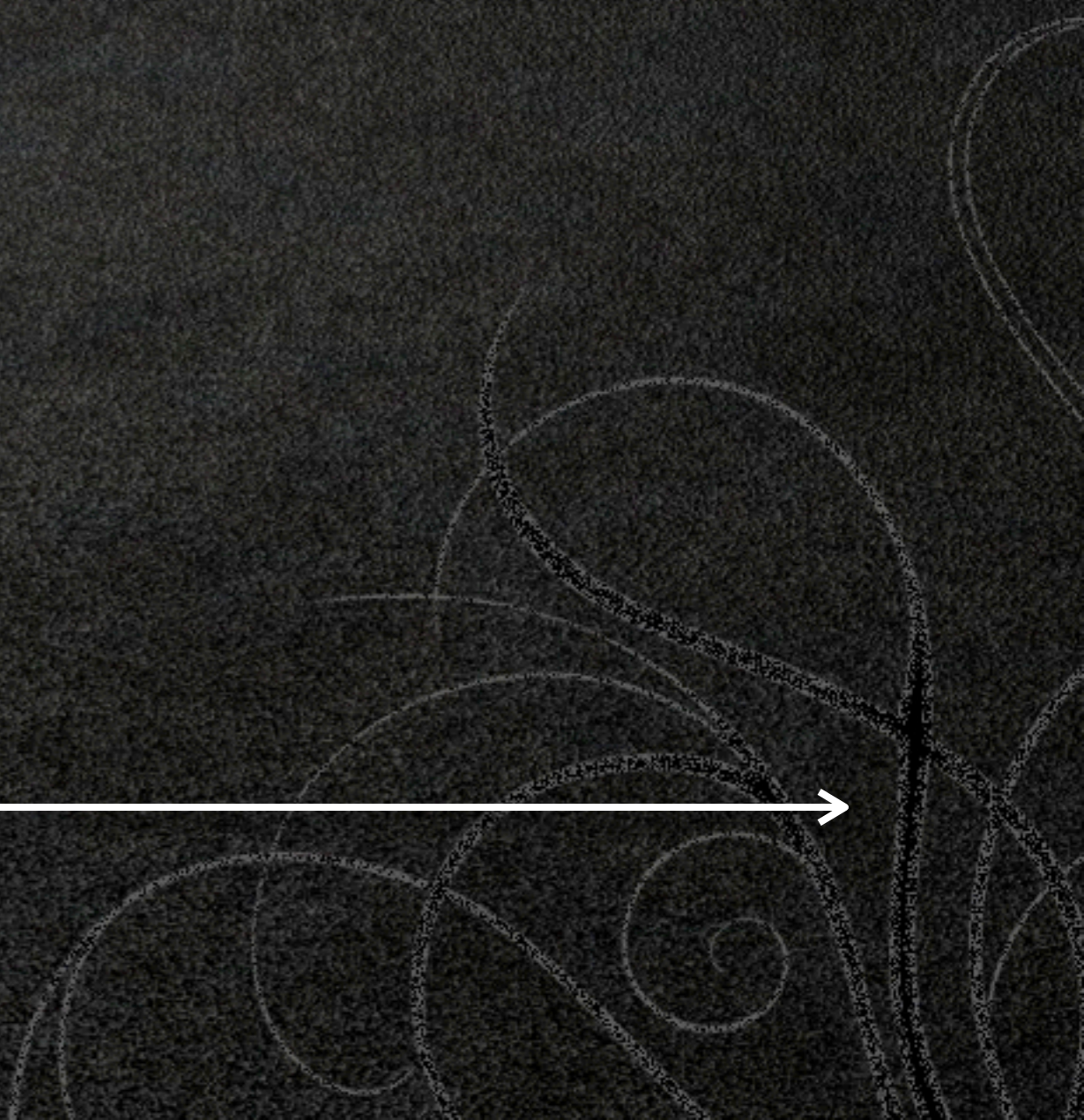
# Application 2

- Horava-Lifshitz cosmology, power-law inflation  $a \propto t^{1+n}$

$$k_{\text{eff}}^2(\eta) = k^2 \left( \frac{k^4}{M^4 a^4} + \alpha_1 \frac{k^3}{M^3 a^3} + \alpha_2 \frac{k^2}{M^2 a^2} + 1 \right)$$

$\log[(\text{length})^{-2}]$

$\log a$



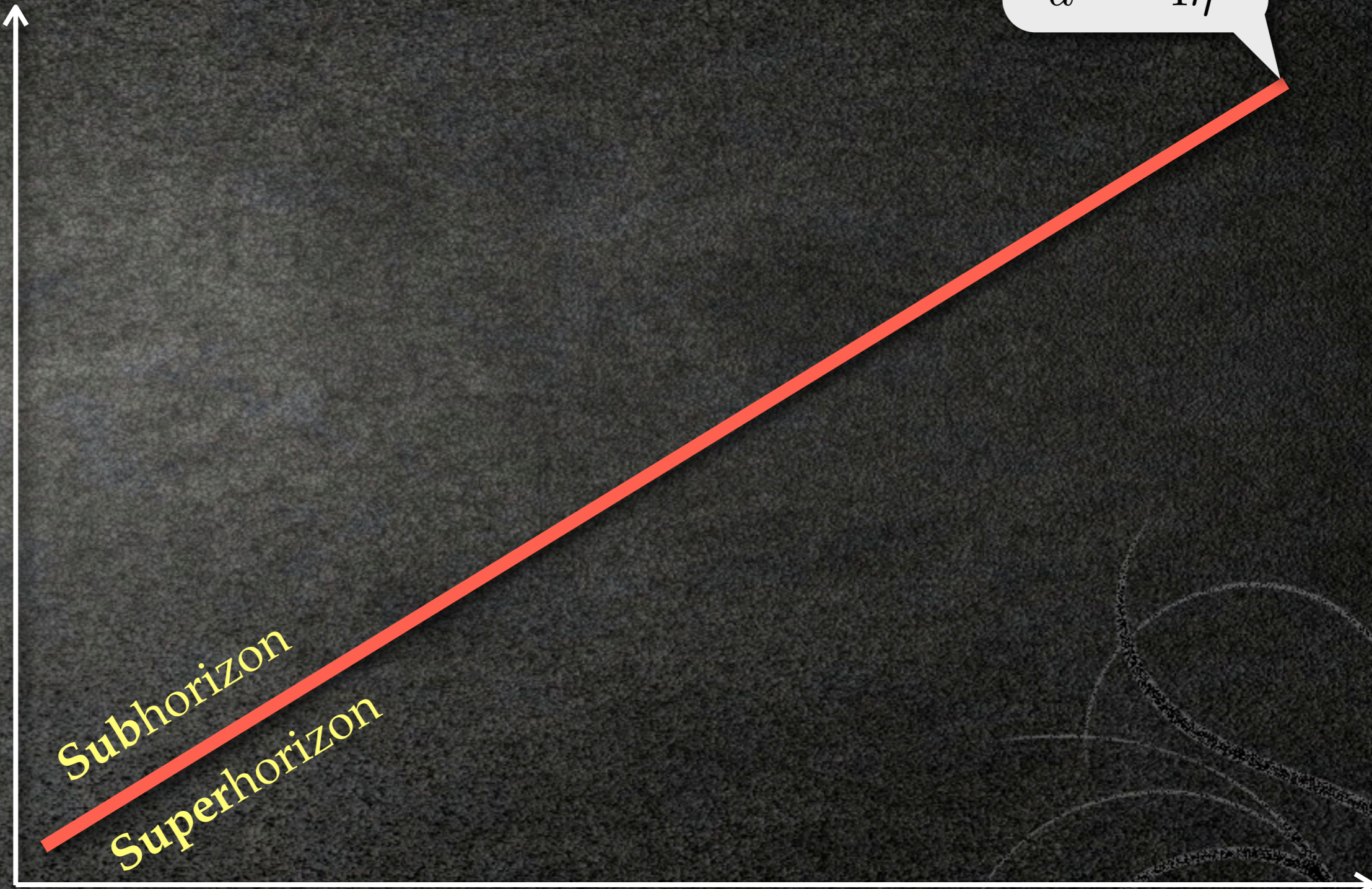
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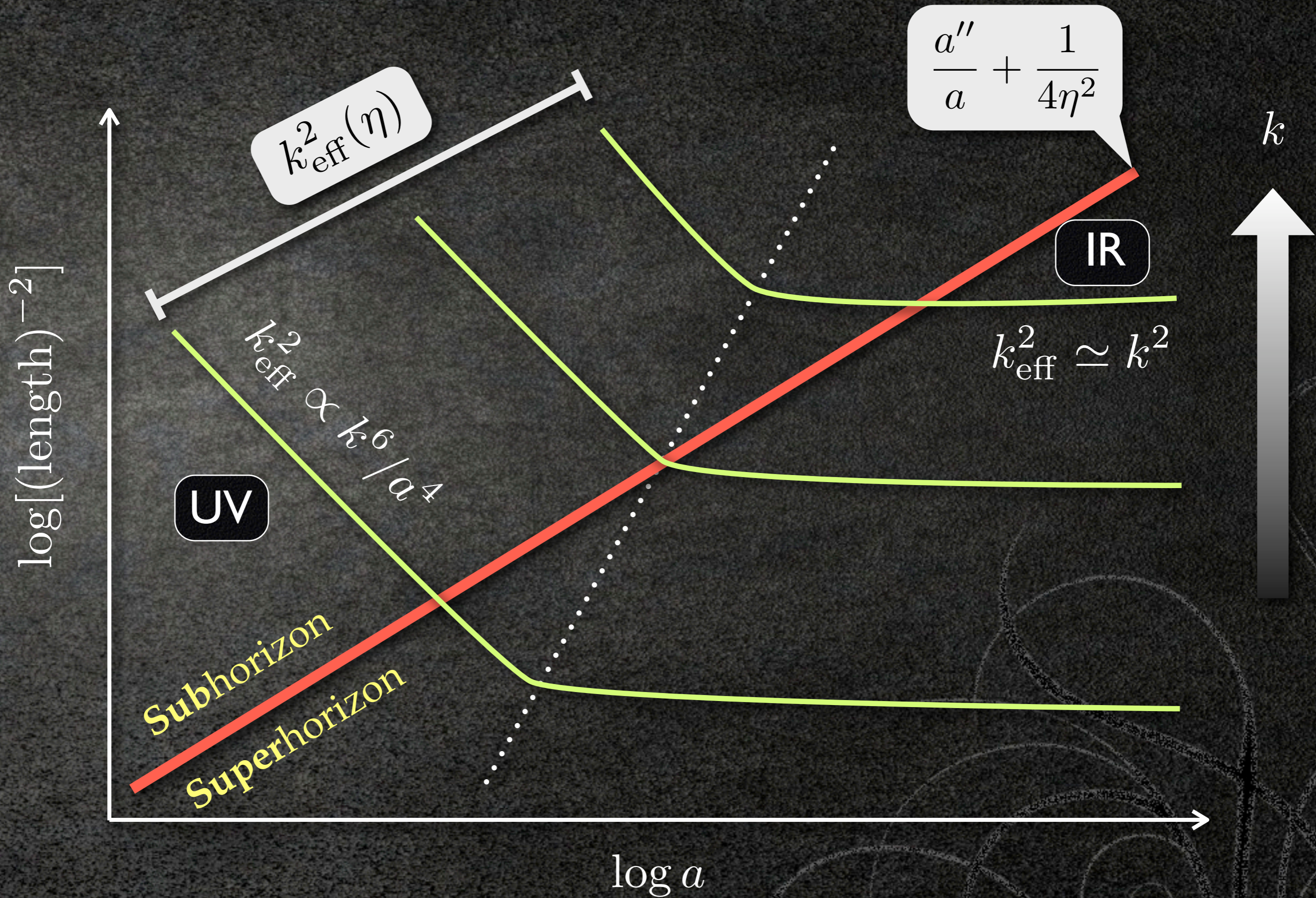
Subhorizon

Superhorizon

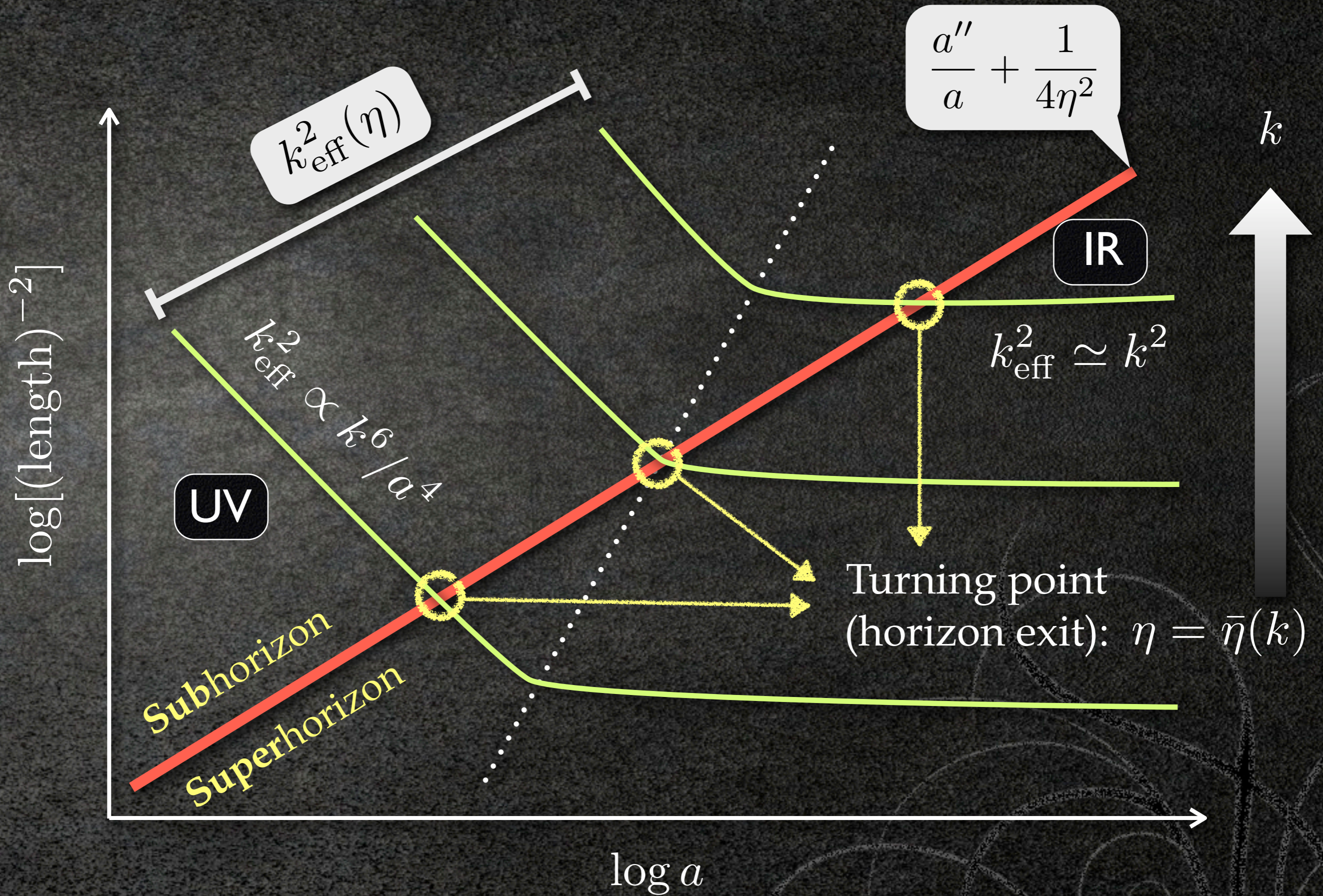
$\log a$

$$\frac{a''}{a} + \frac{1}{4\eta^2}$$









# Application 2

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- Analytic approximate results for limiting cases:

- **Horizon exit in the UV limit** (large  $q$ )

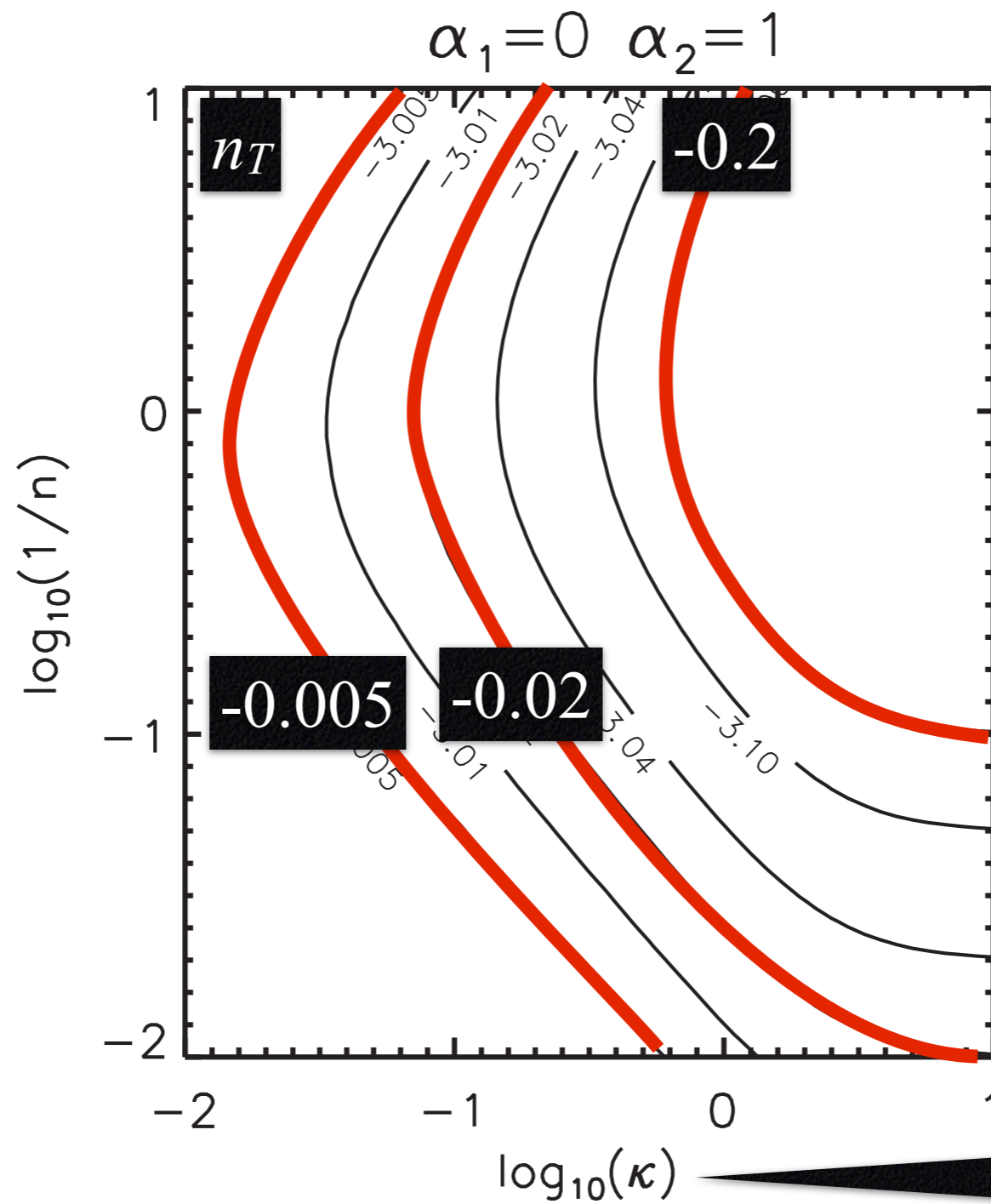
$$n_T \simeq - \left( \frac{\alpha_1}{q} - \frac{\alpha_1^2 - 2\alpha_2}{2q^2} \right) \frac{1}{2 + 3n}$$

- **Horizon exit in the IR limit** (small  $q$ )

$$n_T \simeq -\frac{2}{n} + \left( \alpha_2 q^2 + \frac{3}{2} \alpha_1 q^3 \right) \frac{2 + 3n}{n^2} \quad \text{where} \quad q := \frac{k}{M a(\bar{\eta})}$$

- Generic cases: numerical evaluation of uniform approximation formula for  $n_T$

# Numerical example



$$\kappa \propto k^{1/(1+n)}$$

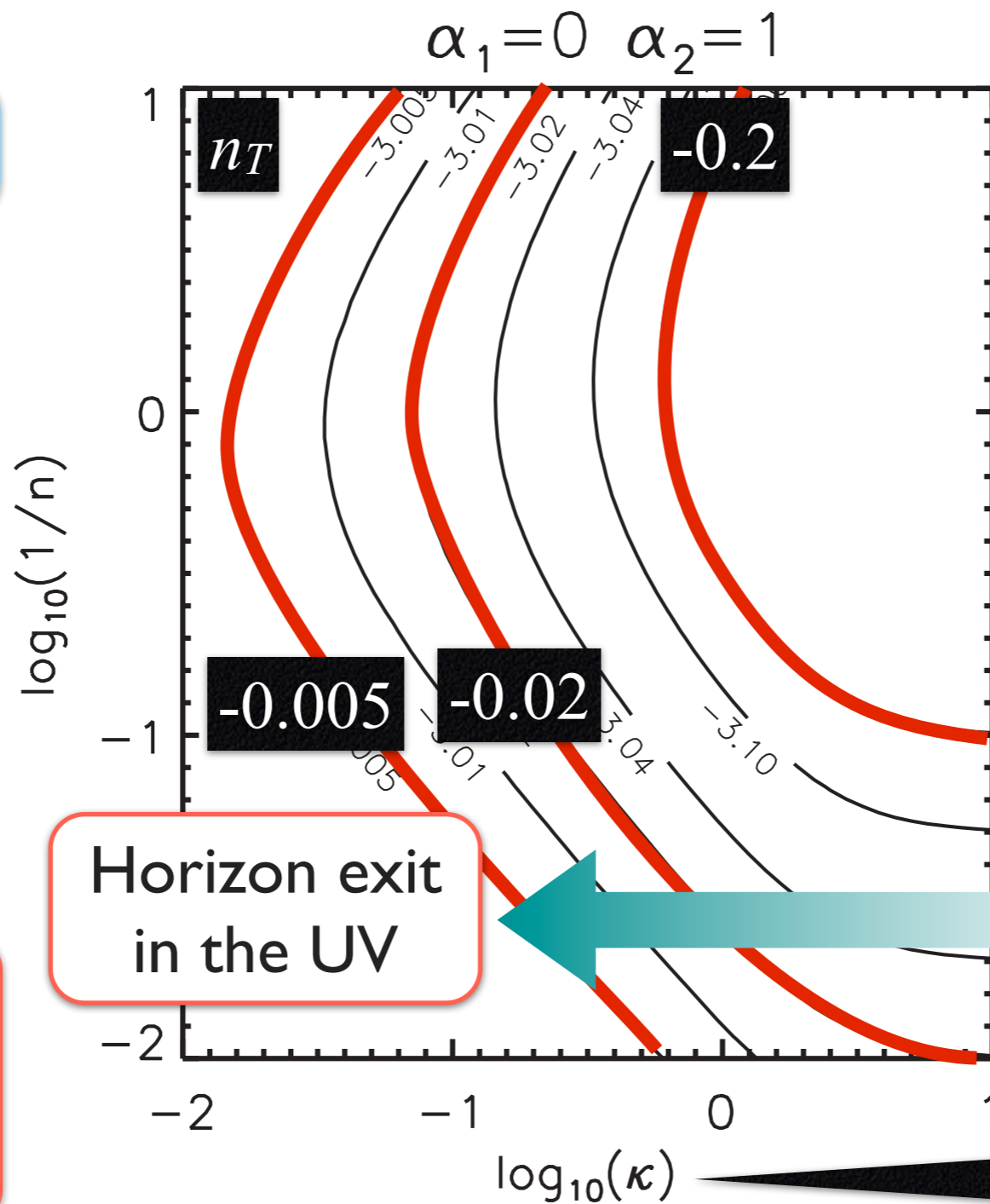
# Numerical example

Non-slow-roll



Slow-roll limit

$$-\dot{H}/H^2 \ll 1$$



Horizon exit  
in the UV

Horizon exit  
in the IR

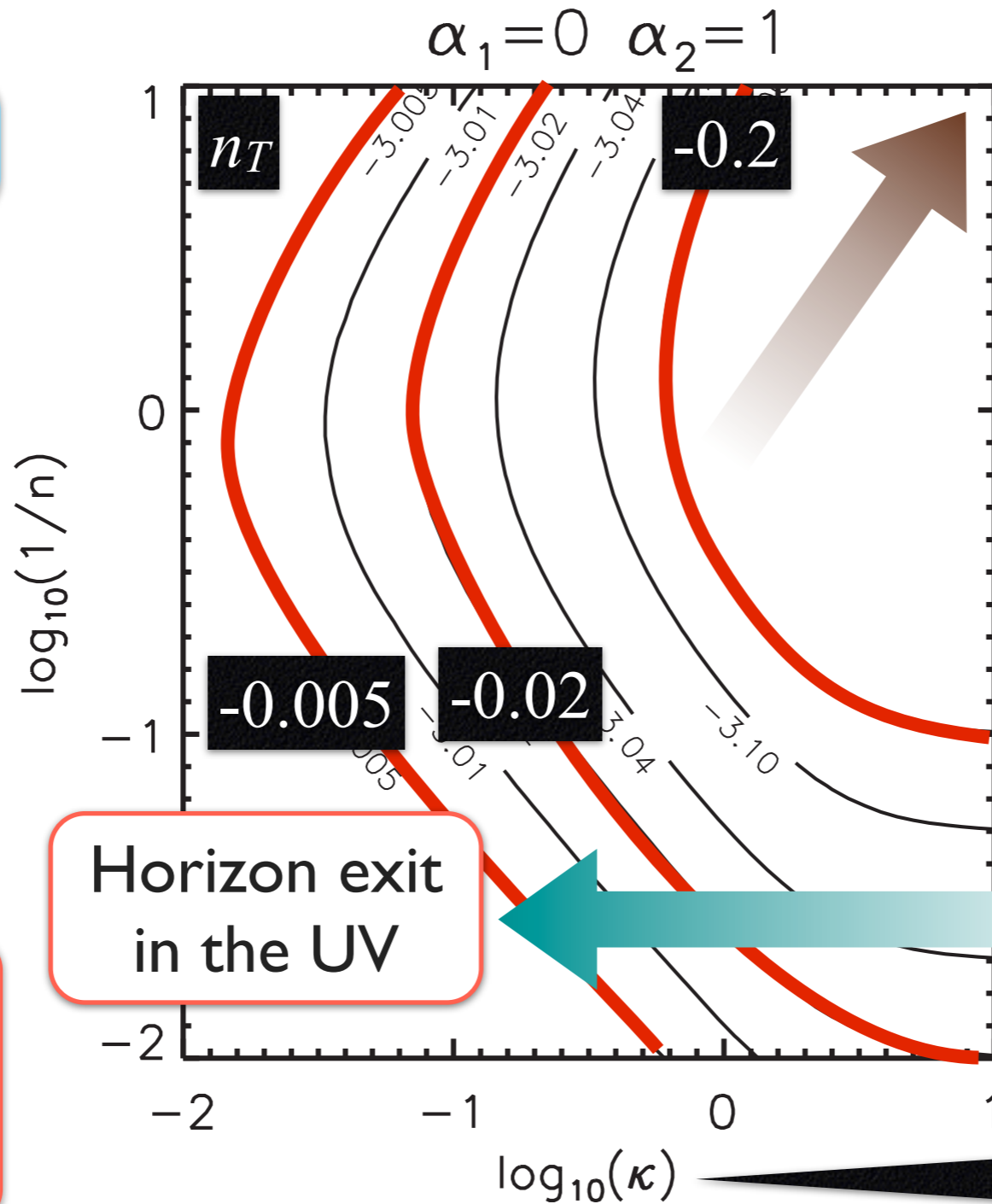
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Large deviation  
from  $n_T = 0$

Horizon exit  
in the IR

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# Conclusions

- In **Horava's gravity**, anisotropic scaling in the UV implies **modified dispersion relation** in the early universe
- **New mechanism to flatten the primordial spectrum:**

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- **New mechanism to flatten the primordial spectrum:**

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- Subdominant terms in the modified dispersion relation give rise to small **tilt** of the spectrum
- Spectral index can be computed accurately using the **uniform approximation**

***Thank you!***





# Appendix



# What is neglected in uniform approximation

- The following quantity is neglected as a first approximation:

$$\psi(\zeta) := [4g(\eta)g''(\eta) - 5g'^2(\eta)] \frac{\zeta}{16g^3(\eta)} + \frac{\zeta q(\eta)}{g(\eta)} + \frac{5}{16\zeta^2}$$

where  $g(\eta) := -k_{\text{eff}}^2 + \frac{a''}{a} + \frac{1}{4\eta^2}$        $q(\eta) := -\frac{1}{4\eta^2}$

$$\zeta(\eta) := \pm \left[ \pm \frac{3}{2} \int_{\bar{\eta}}^{\eta} \sqrt{\pm g(\eta')} d\eta' \right]^{2/3}$$

- Next-order expression is more involved.

Olver, 1997; Martin, Schwarz, 2003;  
Habib *et al.* 2002; 2004