Photon-photon interactions with PYTHIA 8: Current status and future plans FCC-ee Mini-Workshop: "Physics Behind Precision"

Ilkka Helenius in collaboration with Torbjörn Sjöstrand

Lund University Department of Astronomy and Theoretical Physics

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Outline

1 Introduction

- \triangleright Event generation in PYTHIA 8
- 2 Photon-photon interactions
	- \triangleright PDFs for photons
	- \blacktriangleright Initial state radiation with photon beams
	- \triangleright Beam remnant handling
	- \blacktriangleright Results
- **3** Future plans
	- \triangleright Photon flux from electron beams

4 Summary

Hard process

 \blacktriangleright Collinear factorization

$$
d\sigma^{p+p\to k+l} = \sum_{i,j} f_i(x_1, Q^2) \otimes f_j(x_2, Q^2) \otimes d\hat{\sigma}^{i+j\to k+l}
$$

 \triangleright Scale evolution for PDFs given by DGLAP equation

Initial state radiation (ISR)

- \triangleright Backwards evolution, trace back splittings before the hard process
- \triangleright Splitting probability from DGLAP (Conditional probability)

$$
\mathrm{d}\mathcal{P}_{a\leftarrow b} = \frac{\mathrm{d}f_b}{f_b} = \frac{\mathrm{d}Q^2}{Q^2} \frac{x' f_a(x', Q^2)}{x f_b(x, Q^2)} \frac{\alpha_s}{2\pi} P_{a\rightarrow bc}(z) \,\mathrm{d}z
$$

Final state radiation (FSR)

- \blacktriangleright Forward evolution, find splittings after the hard process
- \triangleright Splitting probability from DGLAP

$$
d\mathcal{P}_{a \to bc} = \frac{dQ^2}{Q^2} \frac{\alpha_s}{2\pi} P_{a \to bc}(z) dz,
$$

Multiparton interactions (MPI)

- ► Number of MPIs regulated by screening parameter p_{T0} ($\sim 2 \text{ GeV/c}$)
- Interleaved generation with ISR and FSR using p_T as a common evolution scale equipped with Sudakov factors

Multiparton interactions (MPI)

- ► Number of MPIs regulated by screening parameter p_{T0} ($\sim 2 \text{ GeV/c}$)
- Interleaved generation with ISR and FSR using p_T as a common evolution scale equipped with Sudakov factors
	- \Rightarrow Further emissions from partons generated by MPIs

Beam Remnants

- Add required number of partons to conserve flavour and colour
- Add primordial k_T for the partons
- \triangleright Fix the momentum of remnant partons to obtain total momentum conservation

Hadronization with Lund string model

- \triangleright Connect partons with colour strings
- Allow colour strings to be reconnected (Colour reconnection)

Hadronization with Lund string model

- \triangleright Connect partons with colour strings
- ▶ Allow colour strings to be reconnected (Colour reconnection)
- \blacktriangleright Let the strings decay and form hadrons
- \blacktriangleright Decays to stable hadrons

Photon-photon collisions

Motivation

- \triangleright Give access to many interesting processes and test QCD factorization
- ► Background for future e^+e^- colliders, higher rate with higher luminosity
- \triangleright Aim for a new robust model exploiting PYTHIA 8 developments

Current status (PYTHIA 8.215)

- \triangleright Can generate hard processes in resolved $\gamma\gamma$ interaction with real photons
- \triangleright Can generate ISR and FSR with photon beams
- \blacktriangleright Beam remnants with primordial k_T can be added
- \blacktriangleright Events can be hadronized

Resolved photons

- \triangleright Partonic structure of resolved photons described with PDFs
- \triangleright PDFs obtained via global QCD analysis

DGLAP equations for photons

Additional term due to $\gamma \rightarrow q\bar{q}$ splittings

$$
\frac{\partial f_i^{\gamma}(x, Q^2)}{\partial \log(Q^2)} = \frac{\alpha_{em}}{2\pi} e_i^2 P_{i\gamma}(x) + \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{dz}{z} P_{ij}(z) f_j(x/z, Q^2)
$$

where $P_{i\gamma}(x)=3(x^2+(1-x)^2)$ for quarks, 0 for gluons (LO)

 \triangleright Solution has two components:

$$
f_i^{\gamma}(x,Q^2) = f_i^{\gamma,\text{pl}}(x,Q^2) + f_i^{\gamma,\text{had}}(x,Q^2)
$$

Hadron-like part need non-perturbative input which is fixed by data

$$
f_i^{\gamma, \text{had}}(x, Q_0^2) = N_i x^{a_i} (1-x)^{b_i}
$$

Photon PDFs

 \triangleright Several groups have performed photon PDF analyses

- \triangleright Reasonable agreement between the data and the fits
- \triangleright Currenty we are using PDFs from CJKL analysis [PRD 68 014010 (2003)]
	- \triangleright Leading order analysis, consistent with LO MC
	- \triangleright Provides a parametrization for the PDFs
	- \triangleright Provides point-like and hadron-like parts separately

ISR with photon beams

Different DGLAP evolution

 \triangleright Need to add a term corresponding to $\gamma \rightarrow q\bar{q}$ splitting

$$
d\mathcal{P}_{a\leftarrow b} = \frac{dQ^2}{Q^2} \frac{x' f_a^{\gamma}(x', Q^2)}{xf_b^{\gamma}(x, Q^2)} \frac{\alpha_s}{2\pi} P_{a\rightarrow bc}(z) dz + \frac{dQ^2}{Q^2} \frac{\alpha_{EM}}{2\pi} \frac{e_b^2 P_{\gamma \rightarrow bc}(x)}{f_b^{\gamma}(x, Q^2)}
$$

- \triangleright Possibility to end up to the original photon during the evolution ⇒ No further emissions
	- ⇒ No need for beam remnants

ISR vs PDFs

Backwards evolution should give same number of partons as PDF evolution

 \blacktriangleright The PDFs integrated over relevant region of x

 \blacktriangleright Number of particles produced below Q^2 from ISR algorithm

- \triangleright Qualitatively similar behaviour but also some differences due to different inputs
- \triangleright CJKL analysis uses ACOT(χ) scheme to deal with heavy quark masses ⇒ Some differences in scale evolution

Beam remnants

Valence content of resolved photon

- \triangleright Two "valence" quarks (q \bar{q} pair), flavors can fluctuate
- \triangleright Decompose PDFs to valence and sea parts

$$
f_i^\gamma(x,Q^2) = f_{i,\text{val}}^\gamma(x,Q^2) + f_{i,\text{sea}}^\gamma(x,Q^2)
$$

- \triangleright Decide whether the parton from beam is a valence quark
	- If yes: Other valence quark is the corresponding (anti)quark
	- \triangleright If not: Sample valence flavour using PDFs

A problem:

- \blacktriangleright Decomposition not provided
- \triangleright CJKL: Valence content for hadron-like part

 $f_{i,\text{val}}^{\gamma} = f_{i,\text{val}}^{\gamma,\text{had}} + f_i^{\gamma,\text{pl}}$ i

Beam remnants

- \triangleright Add minimal number of partons to conserve flavour, e.g. if
	- \triangleright a valence quark from beam \Rightarrow Add the other valence quark
	- \triangleright a gluon from beam \Rightarrow Add the valence quarks
	- \triangleright a sea quark from beam \Rightarrow Add the valence quarks + one companion quark

Need invariant mass for remnants

- \blacktriangleright Condition: $W_{\text{rem}} > W_1 + W_2$
- \blacktriangleright Definitive limit when interaction between two valence quarks:

$$
\sqrt{s(1-x_1)(1-x_2)} > m_1 + m_2
$$

- \Rightarrow Reject hard processes and ISR violating this condition
- Also primordial k_T increases the invariant mass

Results

Charged particle p_T spectrum

- \blacktriangleright $\gamma\gamma$ generated with
	- \triangleright CJKL PDFs for photons
	- \triangleright ISR and FSR
	- \blacktriangleright Beam remnants with primordial k_T
- \triangleright MPIs are not included

Comparison to p+p

- \triangleright Cross section smaller due to EM-coupling $(\alpha_{em}^2 \sim 10^{-4})$ Cross section smaller due to
EM-coupling $(\alpha_{em}^2 \sim 10^{-4})$
Harder spectra due to larger
number of large-x partons
- \blacktriangleright Harder spectra due to larger

 \triangleright p_T spectra of the hard $2 \rightarrow 2$ scattering showed for comparison (dashed)

Future plans

Soft processes and MPI

- \blacktriangleright Include also soft QCD processes
- For MPIs need a model for total $\gamma\gamma$ cross section

Photon emissions from electrons

- \blacktriangleright Needed to get the rate of $\gamma\gamma$ interactions in $\mathrm{e^{+}e^{-}}$
- \blacktriangleright Photon flux machine dependent
	- \triangleright For a circular collider (FCC-ee) bremsstrahlung dominant (EPA)

$$
x f_{\gamma}^{e}(x, Q^{2}) = \frac{\alpha_{em}}{2\pi} \log(Q^{2}/m_{e}^{2})(1 + (1 - x)^{2})
$$

- ► For linear $\mathrm{e^+e^-}$ collider with tighter bunches larger flux of photons \Rightarrow Need more flexible form for photon flux
- \triangleright Consider also photon virtuality (so far only real photons)

Once $\gamma\gamma$ is set up, easy to extend to γ p

- \triangleright Still need to make sure that remnants can be added
- \triangleright γ not vet included to public version
- \blacktriangleright Also the photon emissions from proton could be included

 \blacktriangleright $d\sigma^{\gamma\gamma} < d\sigma^{\gamma p} < d\sigma^{\text{pp}}$

Summary & Outlook

Current status

- \triangleright Can generate hard processes for resolved $\gamma\gamma$ interactions
- \blacktriangleright Parton showers can handle photon beams
- \blacktriangleright Beam remnants with primordial k_T can be added
- \blacktriangleright Events can be hadronized
	- \Rightarrow Included into PYTHIA 8.215

Future plans

- \triangleright Model the photon flux from electron beams
- \triangleright Consider virtuality of the photons
- \triangleright Add soft interactions and MPIs
- \triangleright Develop further γ p and ep machinery

Backup

Parton distribution functions (PDFs)

- \triangleright PDFs can not be calculated from first principles of QCD
- \blacktriangleright However, the Q^2 depedence is given by DGLAP evolution equations:

$$
\frac{\partial f_i(x, Q^2)}{\partial \log Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{dz}{z} P_{ij}(z) f_j(x/z, Q^2)
$$

where j runs over the parton flavours

Data for photon PDFs

► Photon structure functions can be measured in $e^- + e^+$ collisions

"Photon DIS"

- \triangleright Other electron emits a virtual photon (γ^*)
	- ⇒ This electron is measured
- \triangleright Other electron is not detected as the scattering angle is small
	- \Rightarrow Photon from this electron has small virtuality
- Also $W_{\gamma\gamma}$ need to be measured to construct kinematics
- \triangleright Data available mainly from different LEP experiments ($\mathcal{O}(200)$ points)
- \triangleright Precision and kinematic coverage more limited than for proton PDFs

 $ACOT(\chi)$ scheme for heavy quarks

DIS kinematics

- \blacktriangleright Limit for heavy quark production $W^2 = Q^2(x^{-1} - 1) > (2m_H)^2$
- In ACOT (y) scheme this is taken into account by rescaling $x \to \chi = x(1 + 4m_H^2/Q^2)$
- \blacktriangleright In CJKL the heavy quark PDFs are zero for $x > 1/(1+\frac{4m_H^2}{Q^2})$

$\gamma + \gamma$ kinematics

 \blacktriangleright Heavy quark limit not related to Q^2 but $\sqrt{s} \Rightarrow$ Undo rescaling $x \to x/(1+4m_H^2/Q^2)$