Photon-photon interactions with PYTHIA 8: Current status and future plans FCC-ee Mini-Workshop: "Physics Behind Precision"

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## Outline

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#### Hard process

Collinear factorization

$$\mathrm{d}\sigma^{p+p\to k+l} = \sum_{i,j} f_i(x_1,Q^2) \otimes f_j(x_2,Q^2) \otimes \mathrm{d}\hat{\sigma}^{i+j\to k+l}$$

Scale evolution for PDFs given by DGLAP equation



#### Initial state radiation (ISR)

- Backwards evolution, trace back splittings before the hard process
- Splitting probability from DGLAP (Conditional probability)

$$\mathrm{d}\mathcal{P}_{a\leftarrow b} = \frac{\mathrm{d}f_b}{f_b} = \frac{\mathrm{d}Q^2}{Q^2} \frac{x'f_a(x',Q^2)}{xf_b(x,Q^2)} \frac{\alpha_s}{2\pi} P_{a\to bc}(z) \,\mathrm{d}z$$



#### Final state radiation (FSR)

- Forward evolution, find splittings after the hard process
- Splitting probability from DGLAP

$$\mathrm{d}\mathcal{P}_{a\to bc} = \frac{\mathrm{d}Q^2}{Q^2} \frac{\alpha_s}{2\pi} P_{a\to bc}(z) \,\mathrm{d}z,$$



#### Multiparton interactions (MPI)

- ▶ Number of MPIs regulated by screening parameter  $p_{T0}$  (~ 2 GeV/c)
- Interleaved generation with ISR and FSR using  $p_T$  as a common evolution scale equipped with Sudakov factors



#### Multiparton interactions (MPI)

- ▶ Number of MPIs regulated by screening parameter  $p_{T0}$  (~ 2 GeV/c)
- Interleaved generation with ISR and FSR using  $p_T$  as a common evolution scale equipped with Sudakov factors
  - $\Rightarrow$  Further emissions from partons generated by MPIs



#### Beam Remnants

- Add required number of partons to conserve flavour and colour
- Add primordial  $k_T$  for the partons
- Fix the momentum of remnant partons to obtain total momentum conservation



#### Hadronization with Lund string model

- Connect partons with colour strings
- Allow colour strings to be reconnected (Colour reconnection)



#### Hadronization with Lund string model

- Connect partons with colour strings
- Allow colour strings to be reconnected (Colour reconnection)
- Let the strings decay and form hadrons
- Decays to stable hadrons

## Photon-photon collisions

#### Motivation

- Give access to many interesting processes and test QCD factorization
- ► Background for future e<sup>+</sup>e<sup>-</sup> colliders, higher rate with higher luminosity
- ► Aim for a new robust model exploiting Pythia 8 developments

## Current status (PYTHIA 8.215)

- $\blacktriangleright$  Can generate hard processes in resolved  $\gamma\gamma$  interaction with real photons
- Can generate ISR and FSR with photon beams
- Beam remnants with primordial  $k_T$  can be added
- Events can be hadronized

## Resolved photons

- ► Partonic structure of resolved photons described with PDFs
- PDFs obtained via global QCD analysis

#### DGLAP equations for photons

- Additional term due to  $\gamma 
ightarrow {
m q} {ar q}$  splittings

$$\frac{\partial f_i^{\gamma}(x,Q^2)}{\partial \log(Q^2)} = \frac{\alpha_{em}}{2\pi} e_i^2 P_{i\gamma}(x) + \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{\mathrm{d}z}{z} P_{ij}(z) f_j(x/z,Q^2)$$

where  $P_{i\gamma}(x) = 3(x^2 + (1-x)^2)$  for quarks, 0 for gluons (LO)

Solution has two components:

$$f_i^{\gamma}(x,Q^2) = f_i^{\gamma,\mathrm{pl}}(x,Q^2) + f_i^{\gamma,\mathrm{had}}(x,Q^2)$$

Hadron-like part need non-perturbative input which is fixed by data

$$f_i^{\gamma, \text{had}}(x, Q_0^2) = N_i x^{a_i} (1-x)^{b_i}$$

## Photon PDFs

Several groups have performed photon PDF analyses



- Reasonable agreement between the data and the fits
- Currenty we are using PDFs from CJKL analysis [PRD 68 014010 (2003)]
  - Leading order analysis, consistent with LO MC
  - Provides a parametrization for the PDFs
  - Provides point-like and hadron-like parts separately

## ISR with photon beams

#### Different DGLAP evolution

 $\blacktriangleright$  Need to add a term corresponding to  $\gamma \to q \bar{q}$  splitting

$$\mathrm{d}\mathcal{P}_{a\leftarrow b} = \frac{\mathrm{d}Q^2}{Q^2} \frac{x' f_a^{\gamma}(x',Q^2)}{x f_b^{\gamma}(x,Q^2)} \frac{\alpha_s}{2\pi} P_{a\to bc}(z) \,\mathrm{d}z + \frac{\mathrm{d}Q^2}{Q^2} \frac{\alpha_{EM}}{2\pi} \frac{e_b^2 P_{\gamma\to bc}(x)}{f_b^{\gamma}(x,Q^2)}$$

- Possibility to end up to the original photon during the evolution
   No further emissions
  - $\Rightarrow$  No need for beam remnants



## ISR vs PDFs

Backwards evolution should give same number of partons as PDF evolution

The PDFs integrated over relevant region of x  Number of particles produced below Q<sup>2</sup> from ISR algorithm



- Qualitatively similar behaviour but also some differences due to different inputs
- ► CJKL analysis uses ACOT( $\chi$ ) scheme to deal with heavy quark masses ⇒ Some differences in scale evolution

## Beam remnants

## Valence content of resolved photon

- ▶ Two "valence" quarks  $(q\bar{q} pair)$ , flavors can fluctuate
- Decompose PDFs to valence and sea parts

$$f_i^\gamma(x,Q^2) = f_{i,\mathrm{val}}^\gamma(x,Q^2) + f_{i,\mathrm{sea}}^\gamma(x,Q^2)$$

- Decide whether the parton from beam is a valence quark
  - ▶ If yes: Other valence quark is the corresponding (anti)quark
  - If not: Sample valence flavour using PDFs

## A problem:

- Decomposition not provided
- CJKL: Valence content for hadron-like part

 $f_{i,\mathrm{val}}^{\gamma} = f_{i,\mathrm{val}}^{\gamma,\mathrm{had}} + f_{i}^{\gamma,\mathrm{pl}}$ 



## Beam remnants

- ► Add minimal number of partons to conserve flavour, e.g. if
  - $\blacktriangleright$  a valence quark from beam  $\Rightarrow$  Add the other valence quark
  - a gluon from beam  $\Rightarrow$  Add the valence quarks
  - ► a sea quark from beam ⇒ Add the valence quarks + one companion quark

#### Need invariant mass for remnants

- Condition:  $W_{\text{rem}} > W_1 + W_2$
- Definitive limit when interaction between two valence quarks:

$$\sqrt{s(1-x_1)(1-x_2)} > m_1 + m_2$$

- ⇒ Reject hard processes and ISR violating this condition
- Also primordial k<sub>T</sub> increases the invariant mass



## Results

#### Charged particle $p_T$ spectrum

- $\gamma\gamma$  generated with
  - CJKL PDFs for photons
  - ISR and FSR
  - Beam remnants with primordial k<sub>T</sub>
- MPIs are not included

#### Comparison to p+p

- Cross section smaller due to EM-coupling  $(\alpha_{em}^2 \sim 10^{-4})$
- Harder spectra due to larger number of large-x partons



▶  $p_T$  spectra of the hard  $2 \rightarrow 2$ scattering showed for comparison (dashed)

## Future plans

## Soft processes and MPI

- Include also soft QCD processes
- $\blacktriangleright$  For MPIs need a model for total  $\gamma\gamma$  cross section

#### Photon emissions from electrons

- $\blacktriangleright$  Needed to get the rate of  $\gamma\gamma$  interactions in  $e^+e^-$
- Photon flux machine dependent
  - ► For a circular collider (FCC-ee) bremsstrahlung dominant (EPA)

$$xf_{\gamma}^{\rm e}(x,Q^2) = \frac{\alpha_{em}}{2\pi} \log(Q^2/m_{\rm e}^2)(1+(1-x)^2)$$

- For linear e<sup>+</sup>e<sup>−</sup> collider with tighter bunches larger flux of photons
   ⇒ Need more flexible form for photon flux
- Consider also photon virtuality (so far only real photons)

## Bonus: $\gamma p$ collisions

Once  $\gamma\gamma$  is set up, easy to extend to  $\gamma p$ 

- Still need to make sure that remnants can be added
- $\gamma p$  not yet included to public version
- Also the photon emissions from proton could be included





 $\blacktriangleright d\sigma^{\gamma\gamma} < d\sigma^{\gamma p} < d\sigma^{pp}$ 

## Summary & Outlook

#### Current status

- $\blacktriangleright$  Can generate hard processes for resolved  $\gamma\gamma$  interactions
- Parton showers can handle photon beams
- Beam remnants with primordial  $k_T$  can be added
- Events can be hadronized
  - $\Rightarrow$  Included into Pythia 8.215

#### Future plans

- Model the photon flux from electron beams
- Consider virtuality of the photons
- Add soft interactions and MPIs
- $\blacktriangleright$  Develop further  $\gamma p$  and ep machinery



# Backup

I. Helenius (Lund U.)

## Parton distribution functions (PDFs)

- PDFs can not be calculated from first principles of QCD
- However, the  $Q^2$  dependence is given by DGLAP evolution equations:

$$\frac{\partial f_i(x,Q^2)}{\partial \log Q^2} = \frac{\alpha_s(Q^2)}{2\pi} \sum_j \int_x^1 \frac{\mathrm{d}z}{z} P_{ij}(z) f_j(x/z,Q^2)$$

where j runs over the parton flavours



$$P_{qq}(z) = \frac{4}{3} \left[ \frac{1+z^2}{(1-z)_+} + \frac{3}{2} \delta(1-z) \right]$$

$$P_{qg}(z) = \frac{4}{3} \left[ \frac{1+(1-z)^2}{z} \right]$$

$$P_{gq}(z) = \frac{1}{2} \left[ z^2 + (1-z)^2 \right]$$

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$$P_{gq}(z) = 6 \left[ \frac{z}{(1-z)_+} + \frac{1-z}{z} + z(1-z) + \frac{11-\frac{2}{3}n_f}{12} \delta(1-z) \right]$$

## Data for photon PDFs

▶ Photon structure functions can be measured in e<sup>-</sup>+e<sup>+</sup> collisions



[Phys.Lett.B436 (1998) 403-416]

## "Photon DIS"

- Other electron emits a virtual photon  $(\gamma^*)$ 
  - $\Rightarrow$  This electron is measured
- Other electron is not detected as the scattering angle is small
  - $\Rightarrow$  Photon from this electron has small virtuality
- Also  $W_{\gamma\gamma}$  need to be measured to construct kinematics
- ▶ Data available mainly from different LEP experiments ( $\mathcal{O}(200)$  points)
- Precision and kinematic coverage more limited than for proton PDFs

## $ACOT(\chi)$ scheme for heavy quarks

## **DIS** kinematics

- ► Limit for heavy quark production  $W^2 = Q^2 (x^{-1} 1) > (2m_H)^2$
- In ACOT(χ) scheme this is taken into account by rescaling

 $x \to \chi = x(1 + 4m_H^2/Q^2)$ 

- In CJKL the heavy quark PDFs are zero for  $x>1/(1+\frac{4m_{H}^{2}}{Q^{2}})$ 



#### $\gamma + \gamma$ kinematics

► Heavy quark limit not related to  $Q^2$  but  $\sqrt{s} \Rightarrow$  Undo rescaling  $x \rightarrow x/(1 + 4m_H^2/Q^2)$