

Tools for precision electroweak calculations

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FCC-ee Mini-Workshop “Physics Behind Precision”

2–3 February 2016

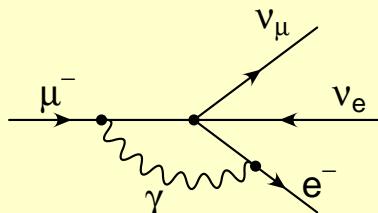
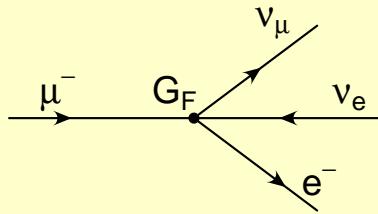
- 1. Status report of WG2 (Precision Electroweak Calculations)**
- 2. Open questions**
- 3. Tools for electroweak precision observables**

Phenomenology WG2, conveners: S. Heinemeyer and A. Freitas



- SM corrections to electroweak precision observables and related processes
 - Fixed-order multi-loop calculations
 - QED/QCD radiation and soft/collinear photon resummation
 - Monte-Carlo tools
 - Evaluation of theory uncertainties
- SM input parameters
 - Robust theoretical definition (in particular for m_t)
 - Prospects for experimental determination of m_t , m_b , α_s , $\Delta\alpha_{\text{had}}$, ...
 - Theory uncertainties in parameter extraction
- Implications for BSM physics
 - (N)MSSM
 - “Beautiful” mirrors
 - ...
 - Effective field theory

μ decay in Fermi Model

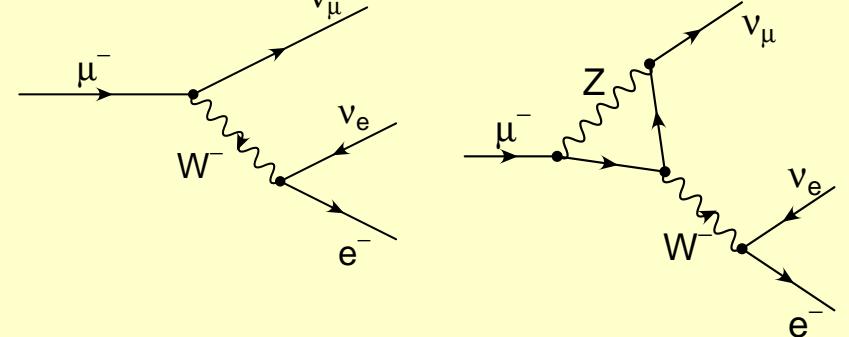
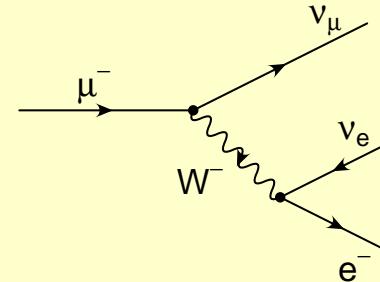


← QED corr.
(2-loop)

$$\Gamma_\mu = \frac{G_F^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) (1 + \Delta q)$$

Ritbergen, Stuart '98
Pak, Czarnecki '08

μ decay in Standard Model



$$\frac{G_F^2}{\sqrt{2}} = \frac{e^2}{8s_W^2 M_W^2} (1 + \Delta r)$$

electroweak corrections

- Deconvolution of initial-state QED radiation:

$$\sigma[e^+e^- \rightarrow f\bar{f}] = \mathcal{R}_{\text{ini}}(s, s') \otimes \sigma_{\text{hard}}(s')$$

- Subtraction of γ -exchange, $\gamma-Z$ interference, box contributions:

$$\sigma_{\text{hard}} = \sigma_Z + \sigma_\gamma + \sigma_{\gamma Z} + \sigma_{\text{box}}$$

- Z -pole contribution:

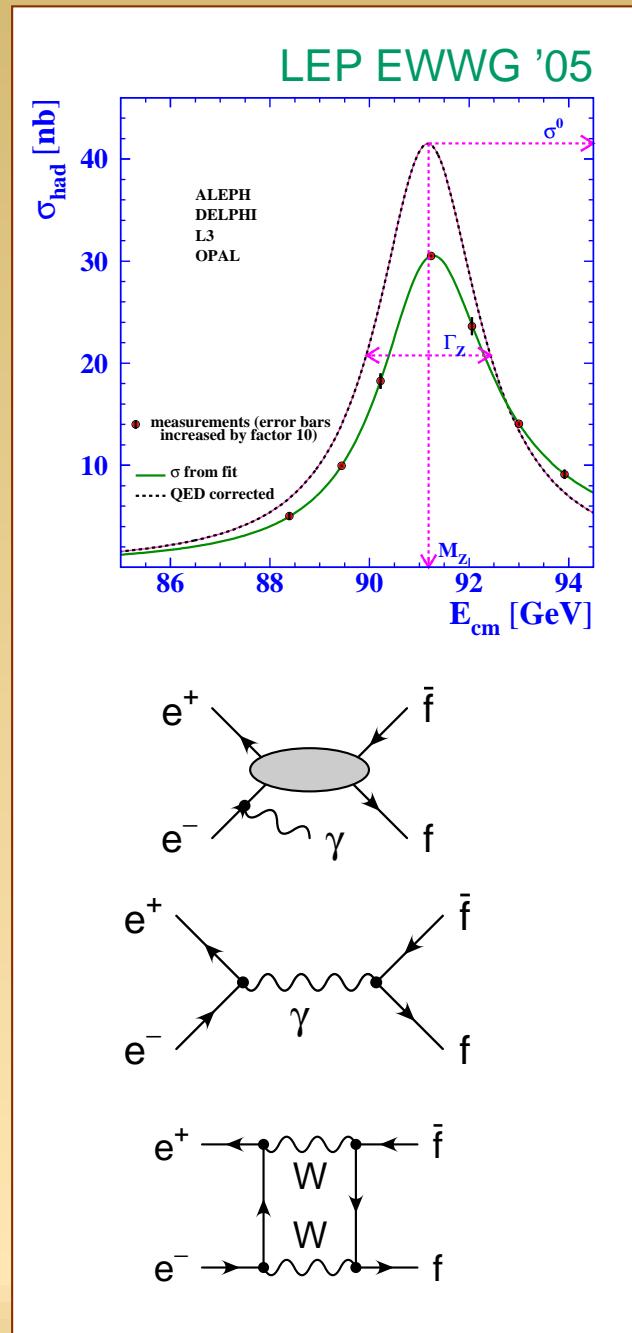
$$\sigma_Z = \frac{R}{(s - M_Z^2)^2 + M_Z^2 \Gamma_Z^2} + \sigma_{\text{non-res}}$$

- In experimental analyses:

$$\sigma \sim \frac{1}{(s - M_Z^2)^2 + s^2 \Gamma_Z^2 / M_Z^2}$$

$$\overline{M}_Z = M_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx M_Z - 34 \text{ MeV}$$

$$\overline{\Gamma}_Z = \Gamma_Z / \sqrt{1 + \Gamma_Z^2 / M_Z^2} \approx \Gamma_Z - 0.9 \text{ MeV}$$



Effective weak mixing angle:

Z -pole asymmetries:

$$A_{\text{FB}}^f \equiv \frac{\sigma(\theta < \frac{\pi}{2}) - \sigma(\theta > \frac{\pi}{2})}{\sigma(\theta < \frac{\pi}{2}) + \sigma(\theta > \frac{\pi}{2})} = \frac{3}{4} \mathcal{A}_e \mathcal{A}_f$$

$$A_{\text{LR}} \equiv \frac{\sigma(\mathcal{P}_e > 0) - \sigma(\mathcal{P}_e < 0)}{\sigma(\mathcal{P}_e > 0) + \sigma(\mathcal{P}_e < 0)} = \mathcal{A}_e$$

$$\mathcal{A}_f = 2 \frac{g_V f / g_{A f}}{1 + (g_V f / g_{A f})^2} = \frac{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f}{1 - 4|Q_f| \sin^2 \theta_{\text{eff}}^f + 8(|Q_f| \sin^2 \theta_{\text{eff}}^f)^2}$$

Most precisely measured for $f = \ell$ (also $f = b, c$)

	Experiment	Theory error	Main source
M_W	80.385 ± 0.015 MeV	4 MeV	$\alpha^3, \alpha^2 \alpha_s$
Γ_Z	2495.2 ± 2.3 MeV	0.5 MeV	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2 \alpha_s, \alpha \alpha_s^2$
σ_{had}^0	41540 ± 37 pb	6 pb	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2 \alpha_s$
$R_b \equiv \Gamma_Z^b / \Gamma_Z^{\text{had}}$	0.21629 ± 0.00066	0.00015	$\alpha_{\text{bos}}^2, \alpha^3, \alpha^2 \alpha_s$
$\sin^2 \theta_{\text{eff}}^\ell$	0.23153 ± 0.00016	4.5×10^{-5}	$\alpha^3, \alpha^2 \alpha_s$

Methods for theory error estimates:

- Parametric factors, *i. e.* factors of α, N_c, N_f, \dots
- Geometric progression, *e. g.* $\frac{\mathcal{O}(\alpha^3)}{\mathcal{O}(\alpha^2)} \sim \frac{\mathcal{O}(\alpha^2)}{\mathcal{O}(\alpha)}$
- Renormalization scale dependence (often underestimates error)
- Renormalization scheme dependence (may underestimate error)

	ILC	FCC-ee	perturb. error with 3-loop [†]	Param. error ILC*	Param. error FCC-ee**
M_W [MeV]	3–5	~ 1	1	2.6	1
Γ_Z [MeV]	~ 1	~ 0.1	$\lesssim 0.2$	0.5	0.06
R_b [10^{-5}]	15	$\lesssim 5$	5–10	< 1	< 1
$\sin^2 \theta_{\text{eff}}^\ell$ [10^{-5}]	1.3	0.6	1.5	2	2

[†] **Theory scenario:** $\mathcal{O}(\alpha \alpha_s^2)$, $\mathcal{O}(N_f \alpha^2 \alpha_s)$, $\mathcal{O}(N_f^2 \alpha^2 \alpha_s)$
 $(N_f^n = \text{at least } n \text{ closed fermion loops})$

Parametric inputs:

* **ILC:** $\delta m_t = 100$ MeV, $\delta \alpha_s = 0.001$, $\delta M_Z = 2.1$ MeV

****FCC-ee:** $\delta m_t = 50$ MeV, $\delta \alpha_s = 0.0001$, $\delta M_Z = 0.1$ MeV

also: $\delta(\Delta \alpha) \sim 5 \times 10^{-5}$

- m_t : Several sources of theory error
 $(e^+e^- \rightarrow t\bar{t}$ cross-section, scheme conversion, α_s)
→ Projected future precision: $\delta_{\text{th}} m_t \sim 50$ MeV Hoang '15
- M_W : From threshold scan of $e^+e^- \rightarrow W^{(*)}W^{(*)}$
Current state of art: full 1-loop for $e^+e^- \rightarrow WW \rightarrow 4f$ and higher-order contribution from effective field theory
→ $\delta_{\text{th}} M_W \sim 3$ MeV Actis, Beneke, Falgari, Schwinn '08
With 2-loop and resummed higher-order terms for $e^+e^- \rightarrow WW$ and $W \rightarrow ff'$ expect $\delta_{\text{th}} M_W \sim 1$ MeV
- α_s :
 $\delta_{\text{exp}} \alpha_s \sim 0.0001$ from measuring R_ℓ at FCC-ee
Theory error (with fermionic electroweak 3-loop corrections):
 $\delta_{\text{th}} R_\ell \sim 0.0015 \quad \Rightarrow \quad \delta \alpha_s^{\text{th}} \sim 0.0002$

- $\Delta\alpha_{\text{had}}$: Could be limiting factor

a) From $e^+e^- \rightarrow \text{had.}$ using dispersion relation

Current: $\delta(\Delta\alpha_{\text{had}}) \sim 10^{-4}$

Improvement to $\delta(\Delta\alpha_{\text{had}}) \sim 5 \times 10^{-5}$ likely

b) Direct determination at FCC-ee from $e^+e^- \rightarrow \mu^+\mu^-$ off the Z peak

(i.e. $A_{\text{FB}}^{\mu\mu}$ at $\sqrt{s} \sim 88$ GeV and $\sqrt{s} \sim 95$ GeV)

$\rightarrow \delta_{\text{th}}(\Delta\alpha_{\text{had}}) \sim 3 \times 10^{-5}$

Janot '15

Requires high-precision theory prediction for $e^+e^- \rightarrow \mu^+\mu^-$ including γ -exchange and box contributions:

$\rightarrow \delta_{\text{th}}(\Delta\alpha_{\text{had}}) \sim \text{few} \times 10^{-4}$ (existing results)

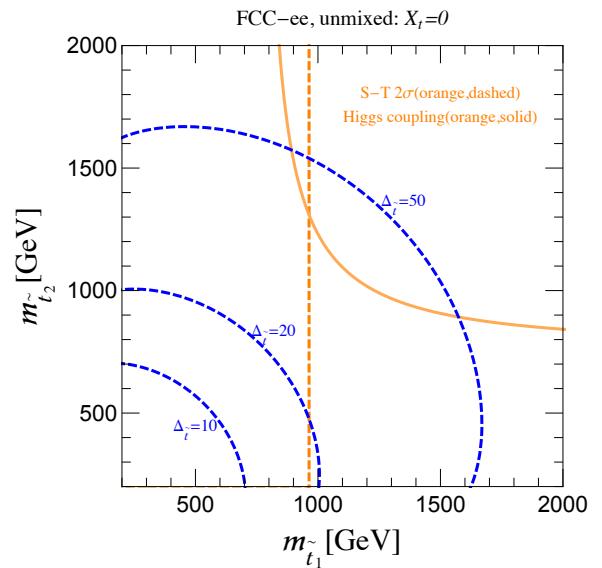
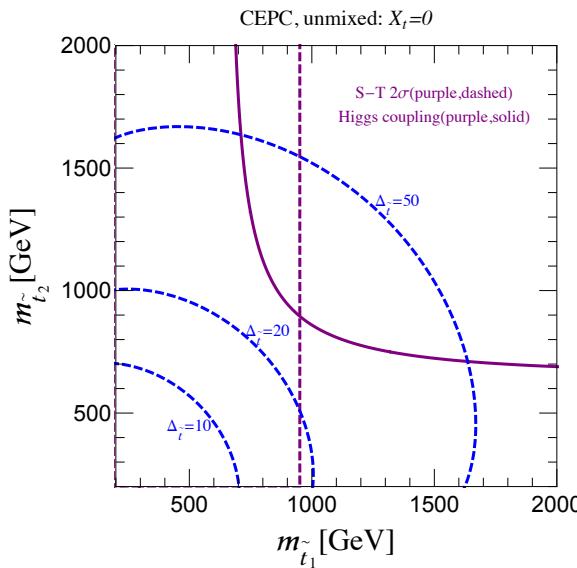
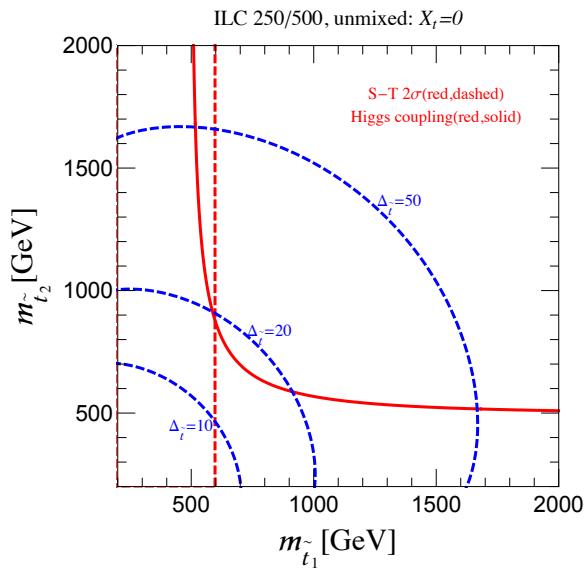
$\delta_{\text{th}}(\Delta\alpha_{\text{had}}) \sim 10^{-4}$ (available methods)

$\delta_{\text{th}}(\Delta\alpha_{\text{had}}) \sim \text{few} \times 10^{-5}$ (reasonable future developments)

Indirect constraints on MSSM:

Fan, Reece, Wang '14

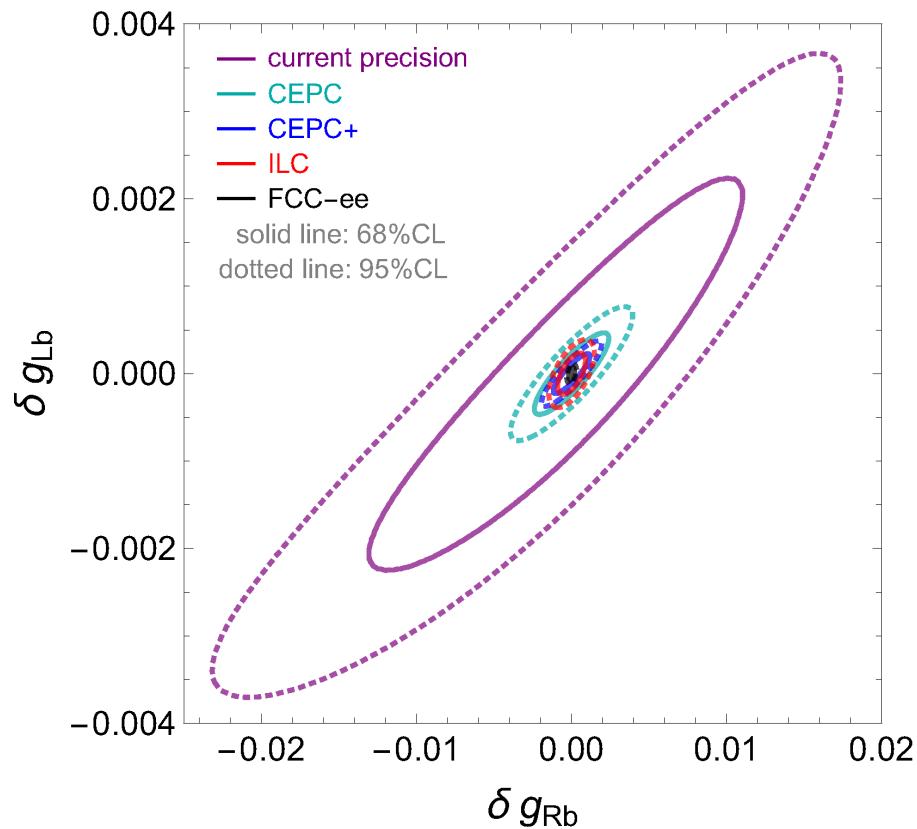
Constraints on stop masses, no mixing ($\Delta_{\tilde{t}} = \text{fine-tuning measure}$)



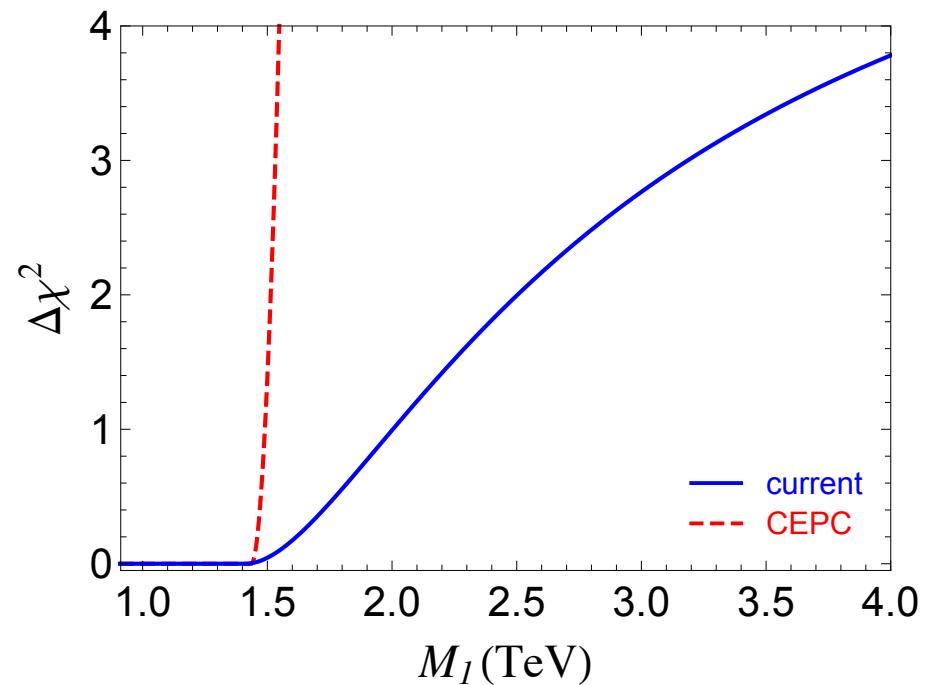
Indirect constraints on beautiful mirrors:

Modified $Z b\bar{b}$ couplings due to vector-like b partners

Choudhury, Tait, Wagner '01



Gori, Gu, Wang '15



Effective field theory: $\mathcal{L} = \sum_i \frac{c_i}{\Lambda^2} \mathcal{O}_i + \mathcal{O}(\Lambda^{-3}) \quad (\Lambda \gg M_Z)$

$$\mathcal{O}_{\phi 1} = (D_\mu \Phi)^\dagger \Phi \Phi^\dagger (D^\mu \Phi)$$

$$\mathcal{O}_{BW} = \Phi^\dagger B_{\mu\nu} W^{\mu\nu} \Phi$$

$$\mathcal{O}_{LL}^{(3)e} = (\bar{L}_L^e \sigma^a \gamma_\mu L_L^e)(\bar{L}_L^e \sigma^a \gamma^\mu L_L^e)$$

$$\mathcal{O}_R^f = i(\Phi^\dagger \overset{\leftrightarrow}{D}_\mu \Phi)(\bar{f}_R \gamma^\mu f_R)$$

$$\mathcal{O}_L^F = i(\Phi^\dagger \overset{\leftrightarrow}{D}_\mu \Phi)(\bar{F}_L \gamma^\mu F_L)$$

$$\mathcal{O}_L^{(3)F} = i(\Phi^\dagger \overset{\leftrightarrow}{D}_\mu^a \Phi)(\bar{F}_L \sigma_a \gamma^\mu F_L)$$

$$\alpha \Delta T = -\frac{v^2}{2} \frac{c_{\phi 1}}{\Lambda^2}$$

$$\alpha \Delta S = -e^2 v^2 \frac{c_{BW}}{\Lambda^2}$$

$$\Delta G_F = -\sqrt{2} \frac{c_{LL}^{(3)e}}{\Lambda^2}$$

$$f = e, \mu, \tau, b, lq$$

$$F = \binom{\nu_e}{e}, \binom{\nu_\mu}{\mu}, \binom{\nu_\tau}{\tau}, \binom{u, c}{d, s}, \binom{t}{b}$$

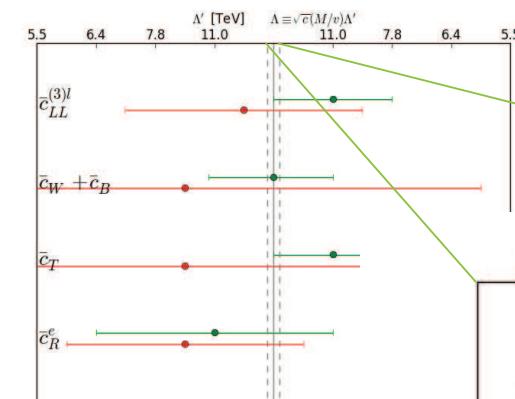
More operators than EWPOs

→ Some can be constrained by $W \rightarrow \ell\nu$, had., $e^+e^- \rightarrow W^+W^-$

EFT description (dimension-6 operators):

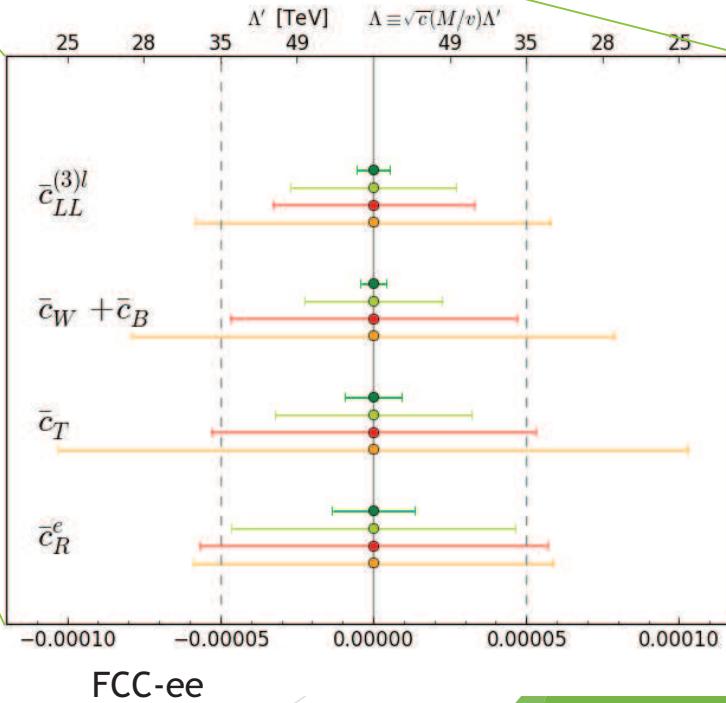
Model-independent, but more general than S/T/U parameters

FCC-ee Summary



LEP

- Dark green: One-by-one (exp. uncertainty only)
- Light green: One-by-one (exp + TH uncertainty)
- Red: Marginalised (exp. uncertainty only)
- Orange: Marginalised (exp + TH uncertainty)



FCC-ee

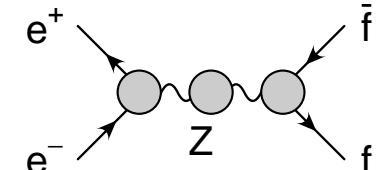
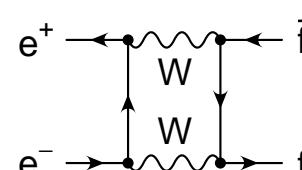
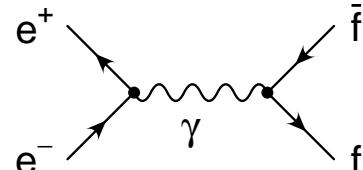
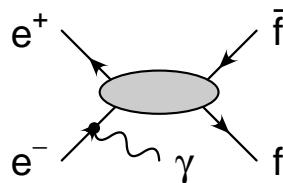
Tevong You (KCL)

Ellis, You '15

- Many methods for evaluating **theory error** and no clear statistical interpretation
→ Ideally we want $\Delta_{\text{th}} \ll \Delta_{\text{exp}}$
But: Theory error may be limiting factor for ILC/FCC-ee precision
- Order of calculations required for **EWPO**
(electroweak 3-loop + leading 4-loop/5-loop?)
- Order of calc. required for **other processes** used to extract M_W , $\Delta\alpha_{\text{had}}$, etc.
(electroweak 2-loop + leading 3-loop)
- Need consistent use of **short-distance definition** of m_t across all sectors
- **MC Tools:** Control of QED effects to sub-permille level is challenge,
but progress is underway
- Evaluation of **new physics reach** including all theory and parametric errors

“Analytical” tools for $e^+e^- \rightarrow f\bar{f}$

- State of the art: Zfitter 6.42
Older code: TOPAZ0
 - Describes true observables ($\sigma_{e^+e^- \rightarrow f\bar{f}}(s)$, etc.)
and pseudo-observables (Γ_Z , σ_{had}^0 , \mathcal{A}_f , etc.)
 - Final-state QED and QCD corrections at $\mathcal{O}(\alpha^2)$, $\mathcal{O}(\alpha\alpha_s)$, $\mathcal{O}(\alpha_s^3)$
 - Deconvolution of initial-state and initial-final QED radiation
at $\mathcal{O}(\alpha)$, $\mathcal{O}(\alpha^2 L)$ and $\mathcal{O}(\alpha^3 L^2)$ ($L \equiv \log(s/m_e^2)$)
 - Full NLO electroweak corrections for $e^+e^- \rightarrow f\bar{f}$
 - Partial $\mathcal{O}(\alpha^2)$ and higher-order electroweak corrections



Drawbacks:

- Not all available NNLO and higher-order corrections implemented (code structure makes implementation difficult)
- For consistent treatment beyond NLO, need expansion of $\mathcal{A}[e^+e^- \rightarrow \mu^+\mu^-]$ about $s_0 = M_Z^2 - iM_Z\Gamma_Z$:

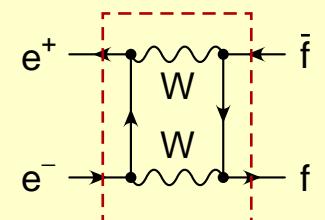
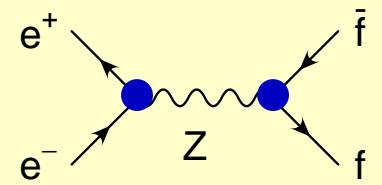
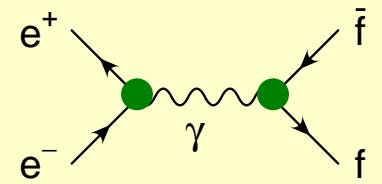
$$\mathcal{A}[e^+e^- \rightarrow f\bar{f}] = \frac{R}{s - s_0} + S + (s - s_0)T + \dots$$

$$R = g_Z^e(s_0)g_Z^f(s_0)$$

$$S = \left[\frac{1}{M_Z^2} g_\gamma^e g_\gamma^f + g_Z^e g_Z^{f'} + g_Z^{e'} g_Z^f + S_{\text{box}} \right]_{s=s_0}$$

$g_V^f(s)$: effective $Vff\bar{f}$ couplings

At NNLO: Need R at $\mathcal{O}(\alpha^2)$, S at $\mathcal{O}(\alpha)$, etc.



- State of the art: KKMC , BabaYaga

Jadach, Ward, Wąs '13
Carloni Calame et al. '12

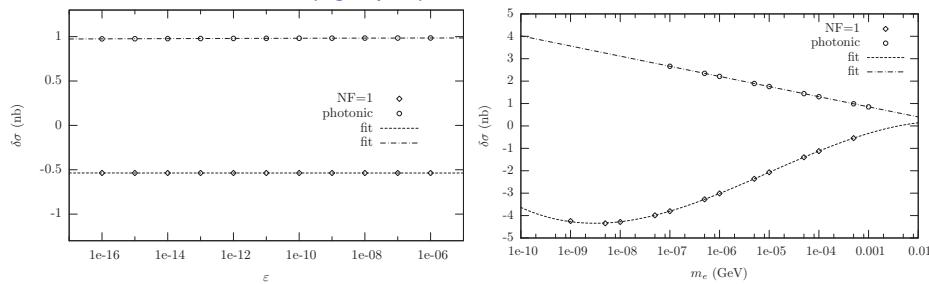
- YFS exponentiation for QED radiation, approximate NNLO QED
- currently $\mathcal{O}(0.1\%)$ precision, $\mathcal{O}(0.01\%)$ feasible in (near) future,
but more may be needed for FCC-ee

Comparison with (a subset of) NNLO

Comparison of $\sigma_{SV}^{\alpha^2}$ calculation of BabaYaga@NLO with

G. Balossini et al., NPB758 (2006) 227

- Penin (photonic): switching off the vacuum polarisation contribution in BabaYaga@NLO, as a function of the logarithm of the soft photon cut-off ([left plot](#)) and of a fictitious electron mass ([right plot](#))



- ★ differences are infrared safe, as expected
- ★ $\delta\sigma(\text{photonic})/\sigma_0 \propto \alpha^2 L$, as expected
- Numerically, for various selection criteria at the Φ and B factories

$$\sigma_{SV}^{\alpha^2}(\text{Penin}) - \sigma_{SV}^{\alpha^2}(\text{BabaYaga@NLO}) < 0.02\% \times \sigma_0$$

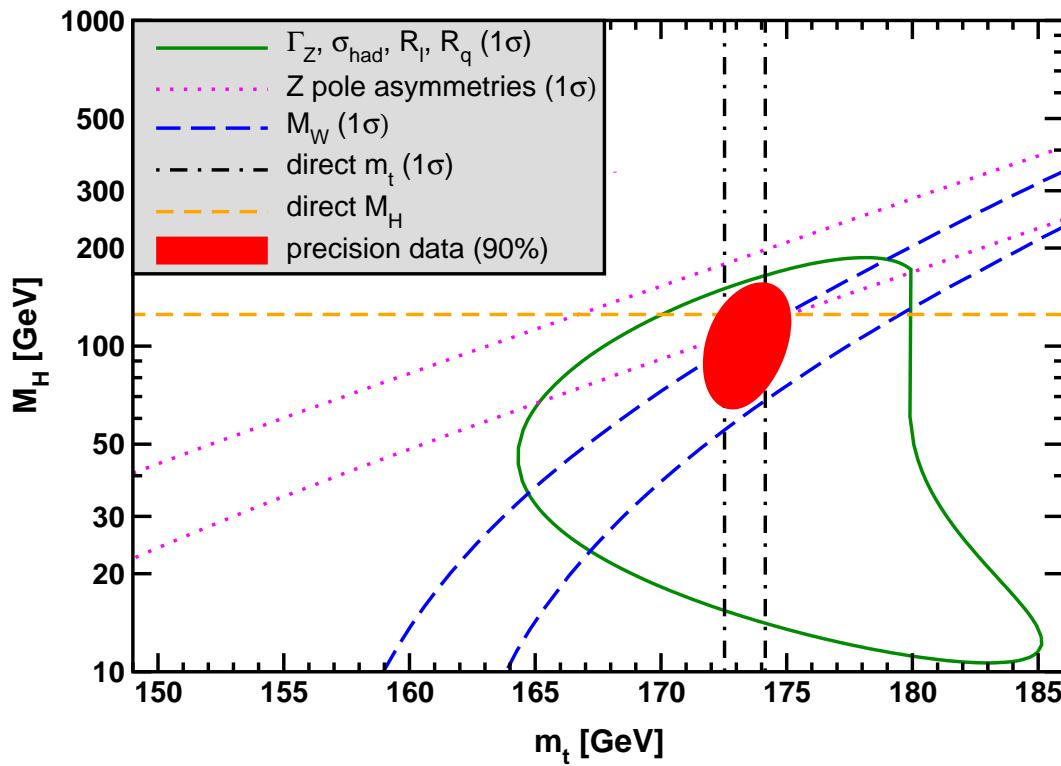
- State of the art: GAPP, Gfitter, Ciuchini et al.

Erler '00,11

Baak et al. '14

Ciuchini, Franco, Mishima, Silvestrini '13

- Based on pseudo-observables (Γ_Z , σ_{had}^0 , \mathcal{A}_f , etc.)
- Include all available NNLO and higher-order corrections
- Constraints on various types of new physics and/or of $d=6$ operators



Erler '16

EWPOs:

- FCC-ee will reduce exp. error by factor $\gtrsim 10$ compared to LEP/SLC
 - Current SM theory calculations not sufficient
 - 3-loop and partial 4-loop (5-loop?) corrections needed!
- Good control over input parameters m_t , M_W , α_s and $\Delta\alpha_{\text{had}}$ is crucial
 - New ideas may be helpful
 - Probably limited by theory uncertainties!
- Need for new/improved computer tools:
 - Monte-Carlo methods for multiple photon corrections
 - Consistent complex pole expansion for electroweak part