

EFT approach to top and W at FCC-ee

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FCC-ee workshop
Feb 3 2016 CERN

The SM EFT

$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{C_i O_i}{\Lambda^2}$$

- $tt\gamma/ttg$, EM/color dipole

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu} \quad O_{tG} = (\bar{Q}\sigma^{\mu\nu}T^A t)\tilde{\varphi}G_{\mu\nu}^A$$

- tbW

- ▶ V/A

$$O_{\varphi Q}^{(3)} = i(\varphi^\dagger D_\mu \tau^I \varphi)(\bar{Q}\tau^I \gamma^\mu Q) \quad O_{\varphi\varphi} = i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu b)$$

- ▶ Weak dipole

$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I \quad O_{bW} = (\bar{Q}\sigma^{\mu\nu}\tau^I b)\varphi W_{\mu\nu}^I$$

- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q) \quad O_{\varphi U} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

- ▶ Weak dipole O_{tW}

$(\bar{u}\gamma^\mu t)Z_\mu$

$$O_{\varphi Q}^{(3,1+3)} = i(\varphi^\dagger \tau^I D_\mu \varphi)(\bar{q}\gamma^\mu \tau^I Q)$$

$$O_{\varphi Q}^{(1,1+3)} = i(\varphi^\dagger D_\mu \varphi)(\bar{q}\gamma^\mu Q)$$

$$O_{\varphi U}^{(1+3)} = i(\varphi^\dagger D_\mu \varphi)(\bar{u}\gamma^\mu t)$$

$(\bar{u}\sigma^{\mu\nu}q_\nu t)V_\mu$, "weak dipole"

$$O_{uW}^{(13)} = (\bar{q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I$$

$$O_{uB}^{(13)} = (\bar{q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu}$$

$(\bar{u}\sigma^{\mu\nu}q_\nu t)G_\mu$, "color dipole"

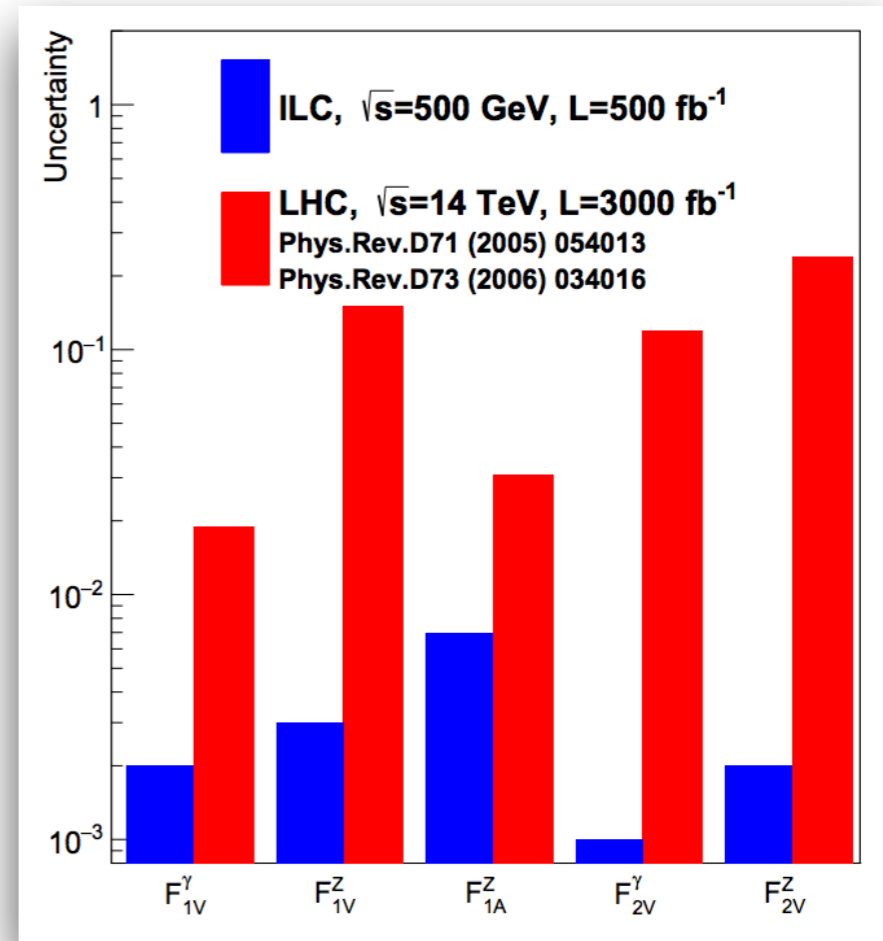
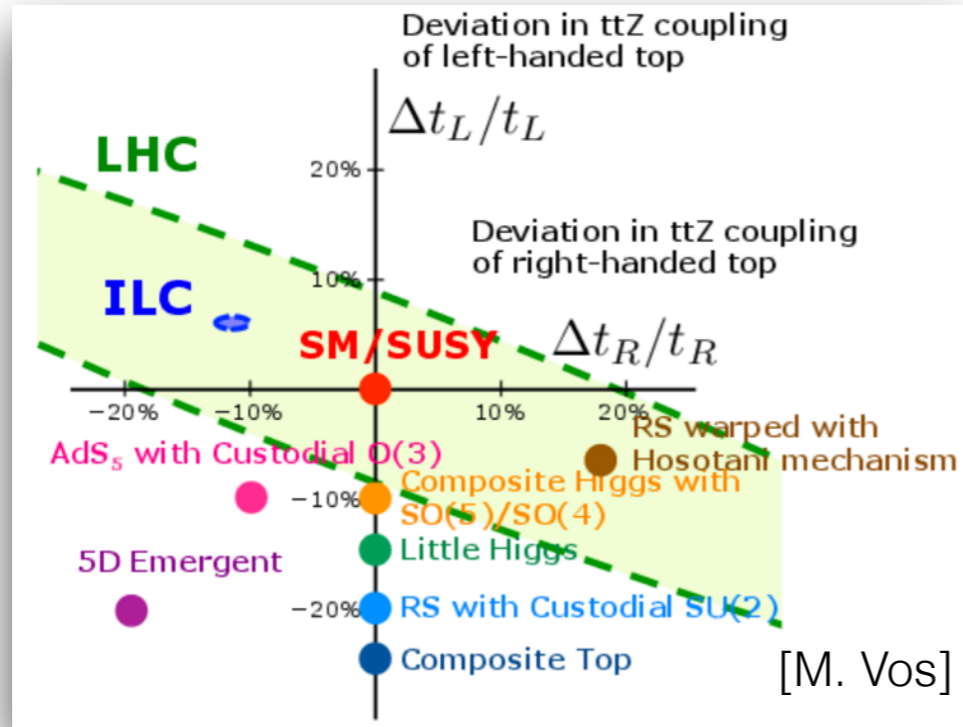
$$O_{uG}^{(13)} = (\bar{q}\sigma^{\mu\nu}T^A t)\tilde{\varphi}G_{\mu\nu}^A$$

$\bar{u}th$, "Yukawa"

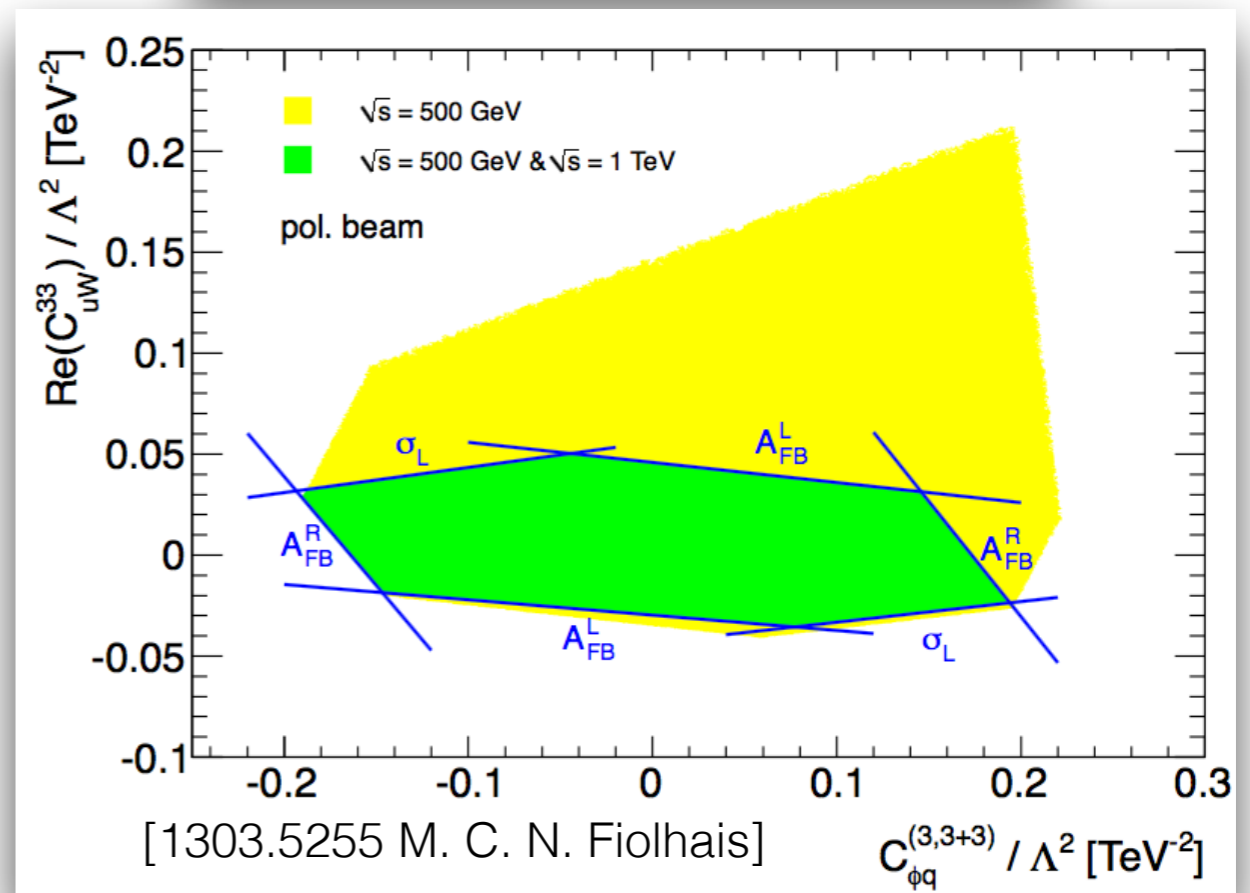
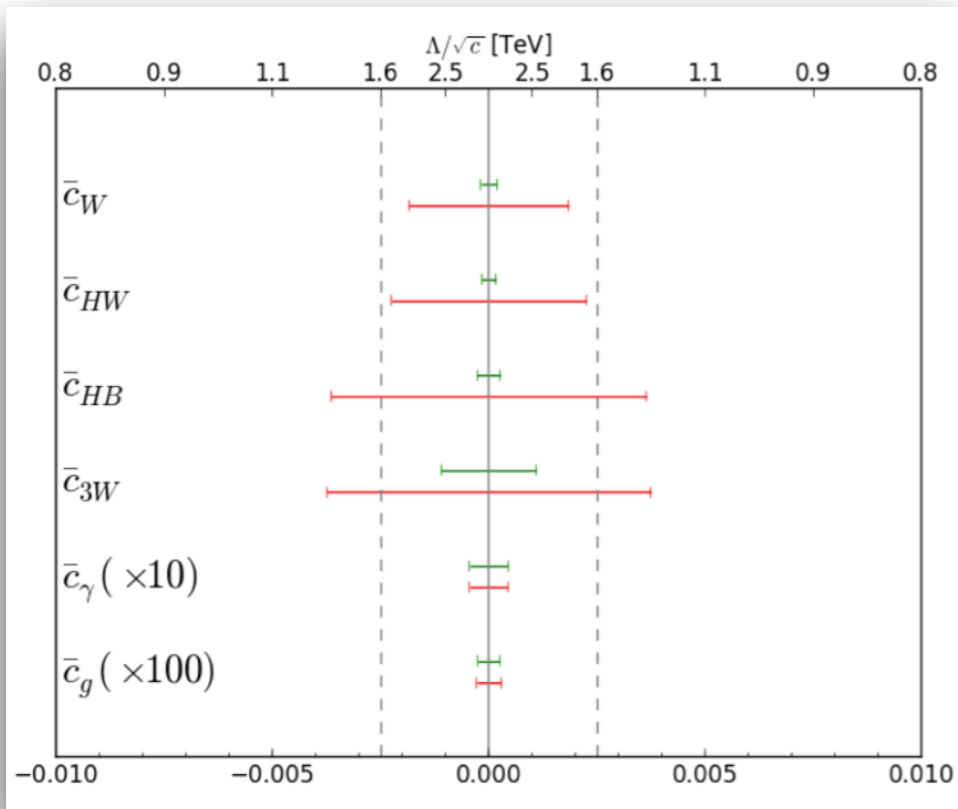
$$O_{u\varphi}^{(13)} = (\varphi^\dagger \varphi)(\bar{q}t)\tilde{\varphi}$$

Physics goal:
determine \mathcal{L}_{EFT} up to dim-6

Top & W couplings at LC



[1510.04561 J. Ellis, T. You]



Making the most use of your EFT

- A. Think globally. EFT combines all measurements, from different experiments, energies, processes, direct/indirect probes, which provide complementary information.
- B. More than just probing “Anomalous Couplings” of a few vertices. Consistent and complete SM deviations.
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Making the most use of your EFT

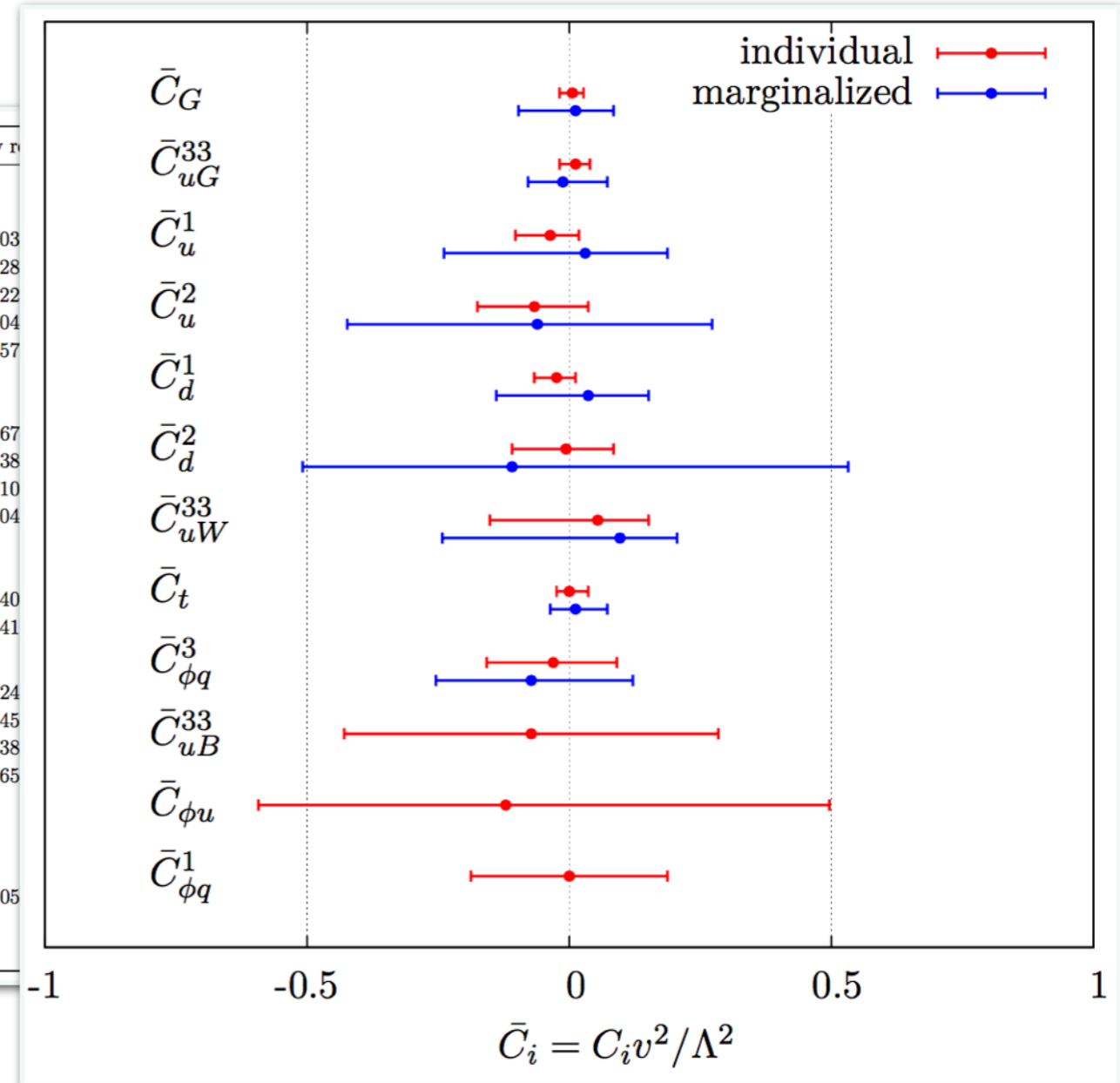
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A. A global approach

- On going fitting program at the LHC

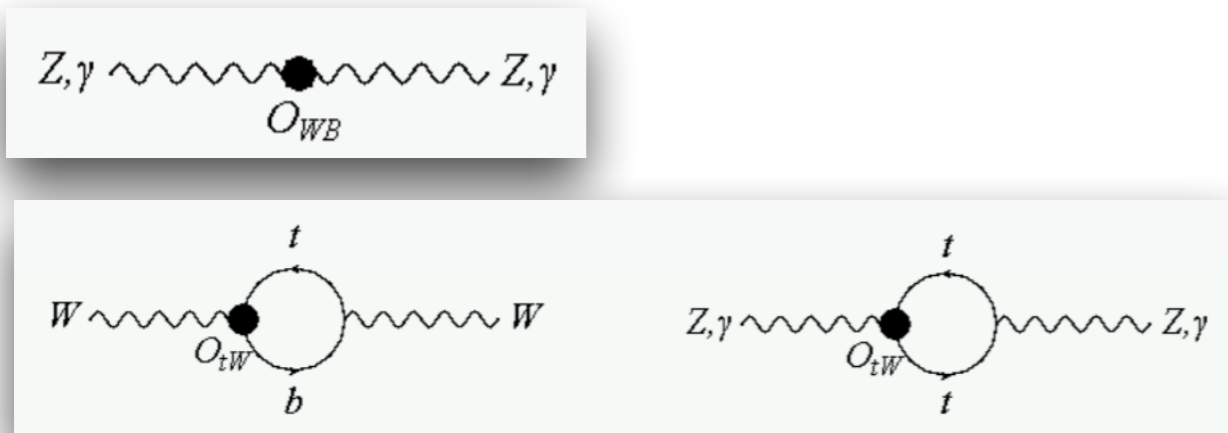
Dataset	\sqrt{s} (TeV)	Measurements	arXiv ref.	Dataset	\sqrt{s} (TeV)	Measurements	arXiv ref.
<i>Top pair production</i>				<i>Differential cross-sections:</i>			
Total cross-sections:				Charge asymmetries:			
ATLAS	7	lepton+jets	1406.5375	ATLAS	7	$p_T(t), M_{t\bar{t}}, y_{t\bar{t}} $	1407.03
ATLAS	7	dilepton	1202.4892	CDF	1.96	$M_{t\bar{t}}$	0903.28
ATLAS	7	lepton+tau	1205.3067	CMS	7	$p_T(t), M_{t\bar{t}}, y_t, y_{t\bar{t}}$	1211.22
ATLAS	7	lepton w/o b jets	1201.1889	CMS	8	$p_T(t), M_{t\bar{t}}, y_t, y_{t\bar{t}}$	1505.04
ATLAS	7	lepton w/ b jets	1406.5375	DØ	1.96	$M_{t\bar{t}}, p_T(t), y_t $	1401.57
ATLAS	7	tau+jets	1211.7205	<i>Top widths:</i>			
ATLAS	7	$t\bar{t}, Z\gamma, WW$	1407.0573	DØ	1.96	Γ_{top}	1308.40
ATLAS	8	dilepton	1202.4892	CDF	1.96	Γ_{top}	1201.41
CMS	7	all hadronic	1302.0508	<i>W-boson helicity fractions:</i>			
CMS	7	dilepton	1208.2761	ATLAS	7		1205.24
CMS	7	lepton+jets	1212.6682	CDF	1.96		1211.45
CMS	7	lepton+tau	1203.6810	CMS	7		1308.38
CMS	7	tau+jets	1301.5755	DØ	1.96		1011.65
CMS	8	dilepton	1312.7582	<i>Run II data</i>			
CDF + DØ	1.96	Combined world average	1309.7570	CMS	13	$t\bar{t}$ (dilepton)	1510.05
<i>Single top production</i>							
ATLAS	7	t -channel (differential)	1406.7844				
CDF	1.96	s -channel (total)	1402.0484				
CMS	7	t -channel (total)	1406.7844				
CMS	8	t -channel (total)	1406.7844				
DØ	1.96	s -channel (total)	0907.4259				
DØ	1.96	t -channel (total)	1105.2788				
<i>Associated production</i>							
ATLAS	7	$t\bar{t}\gamma$	1502.00586				
ATLAS	8	$t\bar{t}Z$	1509.05276				
CMS	8	$t\bar{t}Z$	1406.7830				

[1512.03360 A. Buckley et al.]



A. A global approach

- On going fitting program at the LHC
- More to add
- PEWM (improvement expected at future linear collider)



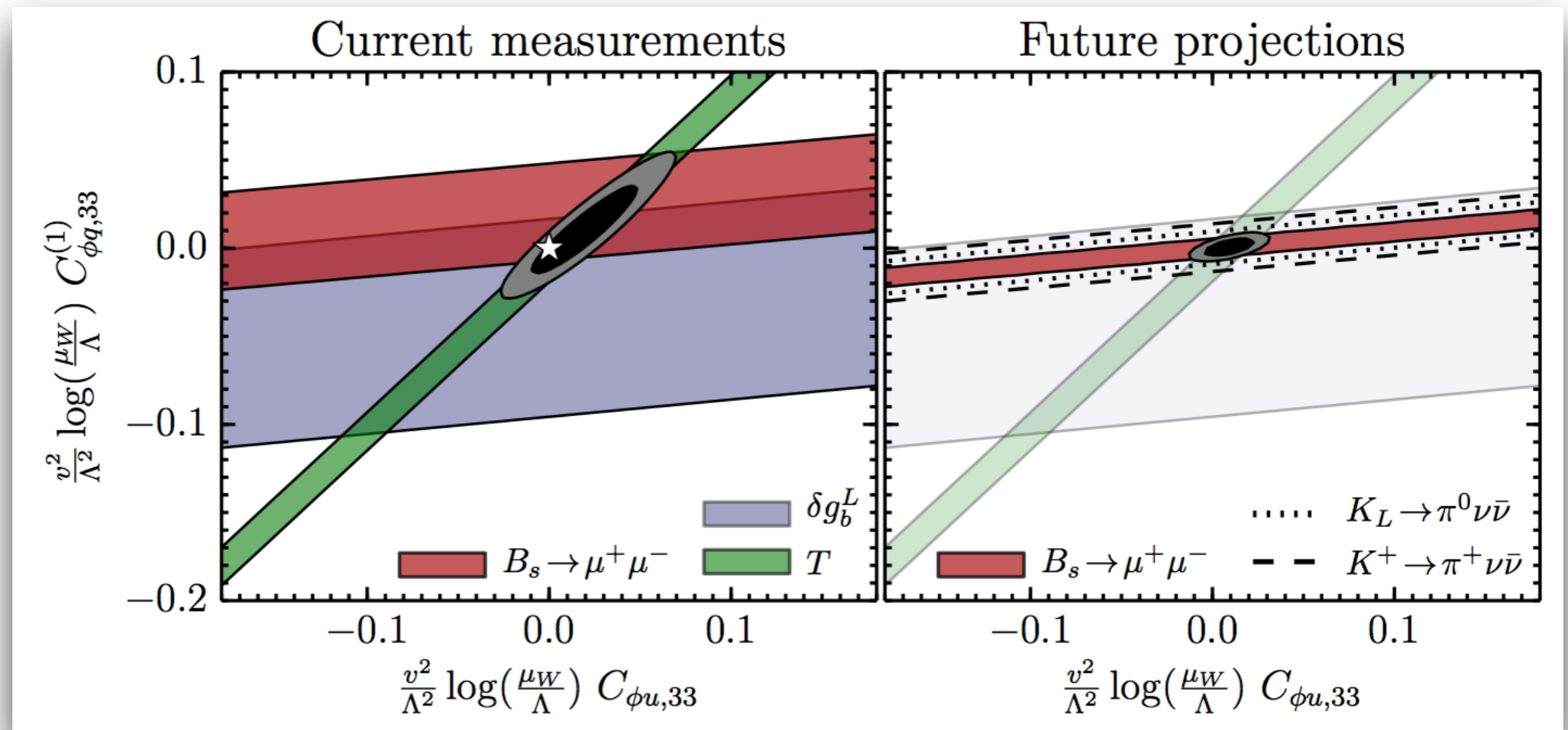
[1201.6670 CZ, Greiner, Willenbrock]

- Top loops mix into Dim-6 EW operators (S, T, ...). Need a theory to renormalize.
- Degeneracies need to be resolved.
- This is done by doing a global fit and making use of the form factors induced by the loops.

Coefficients	Electroweak data
$(C_{\phi q}^{(3)} + C_{\phi q}^{(1)}) / \Lambda^2$	0.016 ± 0.021
$(C_{\phi q}^{(3)} - C_{\phi q}^{(1)}) / \Lambda^2$	2.0 ± 2.7
$C_{\phi t} / \Lambda^2$	1.8 ± 1.9
$C_{\phi b} / \Lambda^2$	-0.16 ± 0.10
$C_{\phi\phi} / \Lambda^2$	
C_{tW} / Λ^2	-0.4 ± 1.2
C_{bW} / Λ^2	11 ± 13
C_{tB} / Λ^2	4.8 ± 5.3
C_{bB} / Λ^2	8 ± 19

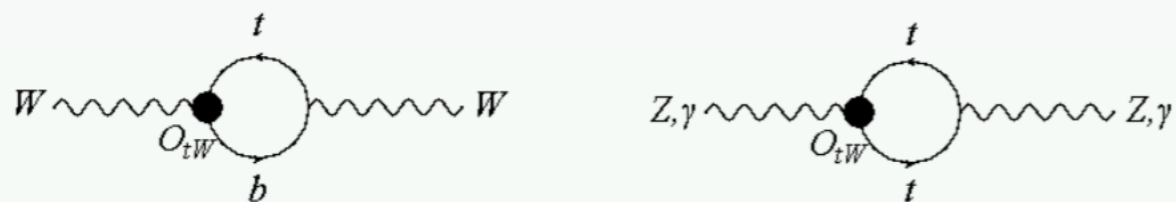
A. A global approach

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- Rare meson decays



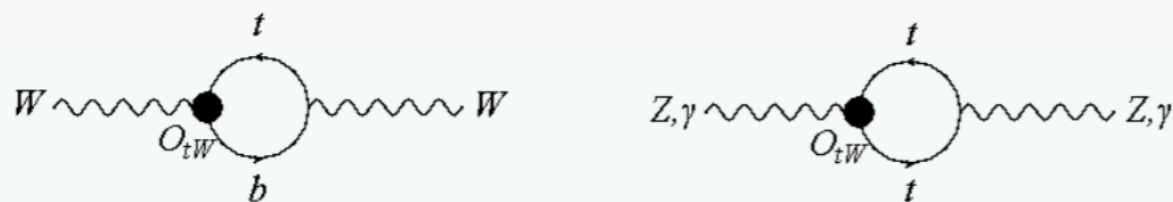
A. A global approach

- On going fitting program at the LHC
- More to add
- PEWM (improvement expected at future linear collider)
- Rare meson decays
- Future colliders: EFT allows for combining experiments that are
 - Performed **at different scales and need RG running** to be combined.
 - Probing **indirect effect through loops**, which can be consistently accommodated along with direct ones.



A. A global approach

- On going fitting program at the LHC
- More to add
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Use one framework, avoid specific underlying assumptions. Eventually everything can be combined.

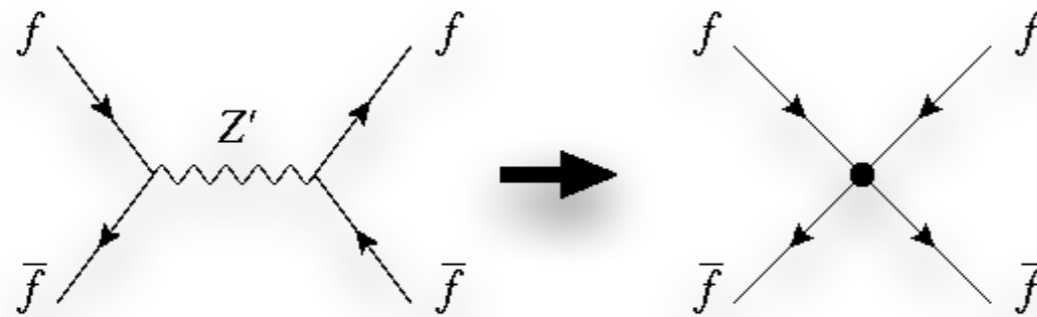
Making the most use of your EFT

- A. Think globally. EFT combines all measurements, from different experiments, energies, processes, direct/indirect probes, which provide complementary information.
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B. More than just “Anomalous Couplings”

- EFT may look less transparent BUT
 - Gauge invariance (requires $tbWZ/\gamma$, $ttZH$, etc., relevant in $pp \rightarrow tZ/\gamma$, $ee \rightarrow ttH, \dots$)
 - Consistently renormalizable, thus higher order possible
 - **More to explore.**
 - In particular, do not miss four-fermion operators in top-physics context!

B. More than just “Anomalous Couplings”



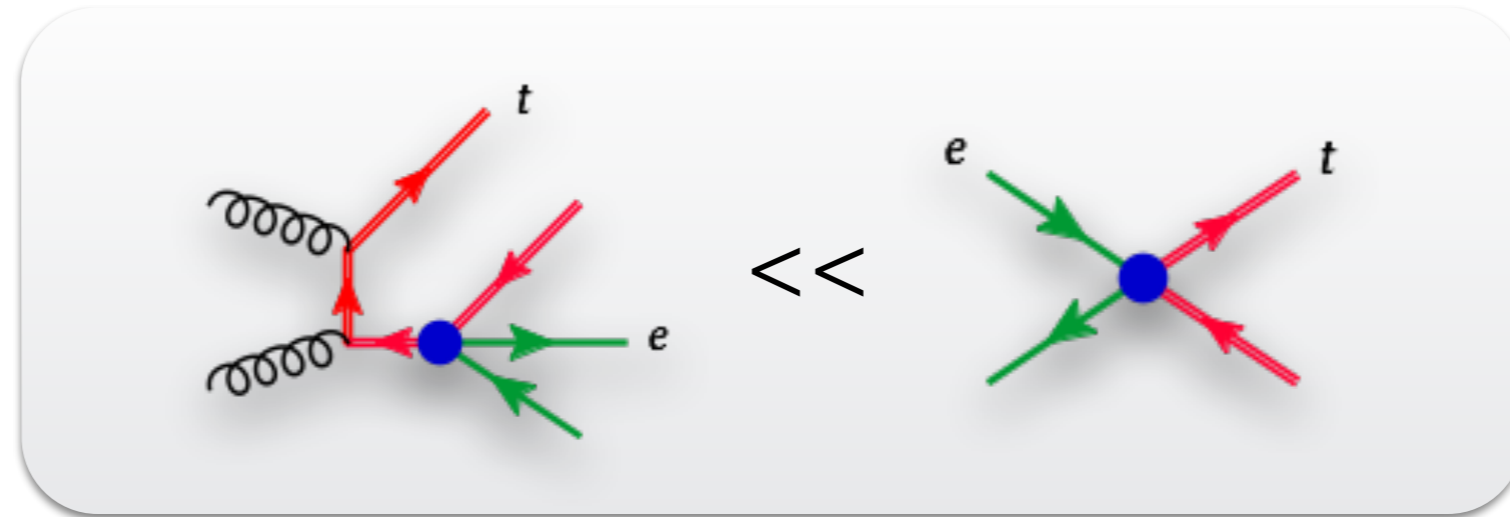
$$\begin{aligned} O_{lQ}^{(1)} &= (\bar{l}\gamma^\mu l) (\bar{Q}\gamma_\mu Q) \\ O_{eQ} &= (\bar{e}\gamma^\mu e) (\bar{Q}\gamma_\mu Q) \\ O_{lt} &= (\bar{l}\gamma^\mu l) (\bar{t}\gamma_\mu t) \\ O_{et} &= (\bar{e}\gamma^\mu e) (\bar{t}\gamma_\mu t) \end{aligned}$$

- The **four-fermion operators** are somewhat ignored in top physics. However they are well motivated:
 - **BSM with heavy mediators** generate 4-fs at tree level
 - **Consistency**: physically, there is not fundamental difference between 2-f vertices and 4-f contact interactions.
 - **A complete characterization** of BSM would always need all 4-f operators in any case.

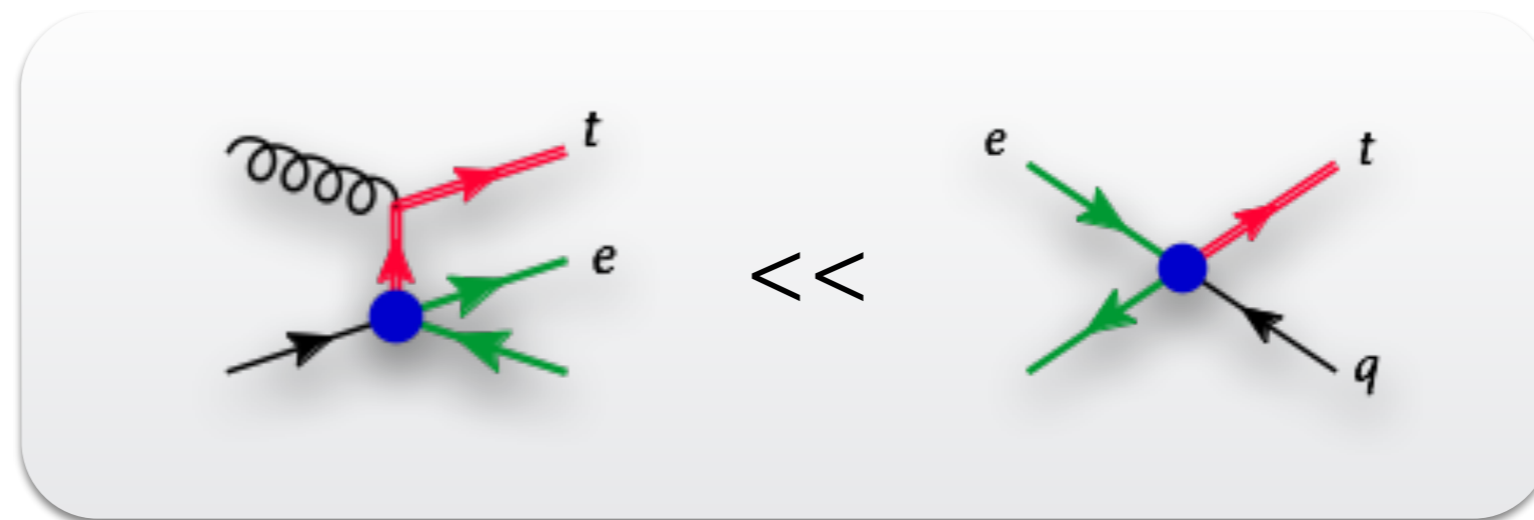
B. More than just “Anomalous Couplings”

LHC vs LC

Flavor-conserving case



Flavor-changing case



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- C. Higher-order corrections possible. Predictions can be systematically improved.

C. Radiative corrections

Measured data is only meaningful if it is matched with theoretical calculation at the same level of accuracy

- Predictions systematically improvable.

$$\mathcal{O}(1) + \mathcal{O}(\alpha_s) + \mathcal{O}(1/\Lambda^2) + \mathcal{O}(\alpha_s/\Lambda^2) + \dots$$

LO SM NLO SM LO EFT NLO EFT

- Precision efforts ongoing.

[G. Passarino 2012] [M. Ghezz et al. 2015] [R. Contino et al. 2014]
[C. Hartmann, M. Trott 2015] [C. Hartmann, M. Trott 2015]

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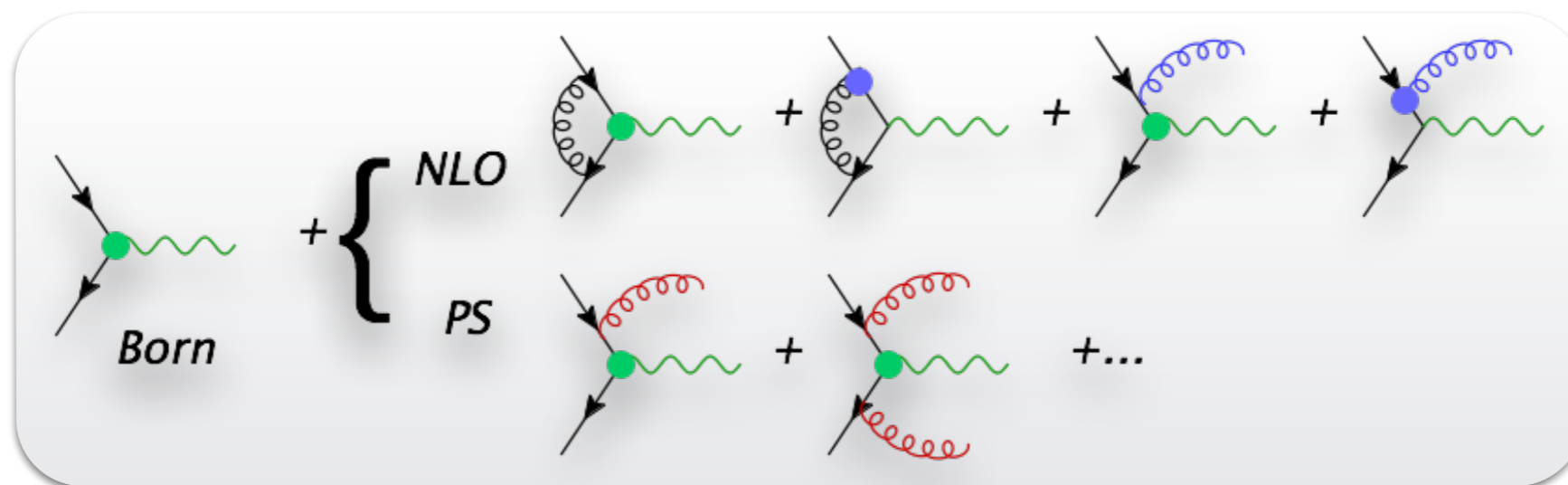
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$$\mathcal{O}(1) + \mathcal{O}(\alpha_s) + \mathcal{O}(1/\Lambda^2) + \mathcal{O}(\alpha_s/\Lambda^2) + \dots$$

LO SM
NLO SM
LO EFT
NLO EFT

- Precision efforts ongoing. *[G. Passarino 2012]* *[M. Ghezz et al. 2015]* *[R. Contino et al. 2014]*
[C. Hartmann, M. Trott 2015] *[C. Hartmann, M. Trott 2015]*

- Tools are becoming available.



being automated in MadGraph5_aMC@NLO

C. Radiative corrections

[1601.08193 Bylund, Maltoni, Tsinikos, Vryonidou, CZ]

QCD corrections for top
operators automated

Example #1:
tt production

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[1601.08193 Bylund, Maltoni, Tsirikos, Vryonidou, CZ]

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	virt	8	1	postscript	e- e+ > t t~ WEIGHTED=7 EFT=1 QED=2 [all = QCD]
	real	12	1	postscript	e- e+ > t t~ g WEIGHTED=8 EFT=1 QED=2 [all = QCD]

C. Radiative corrections

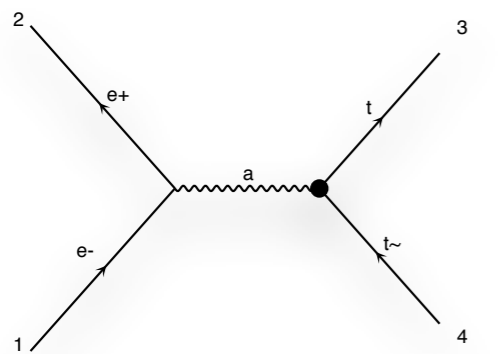
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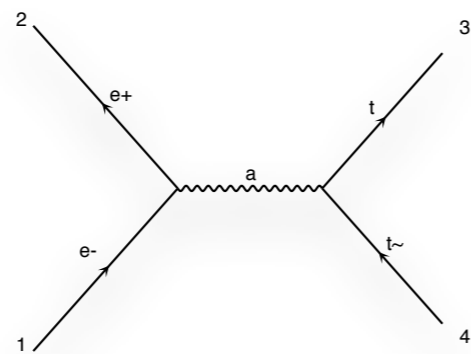
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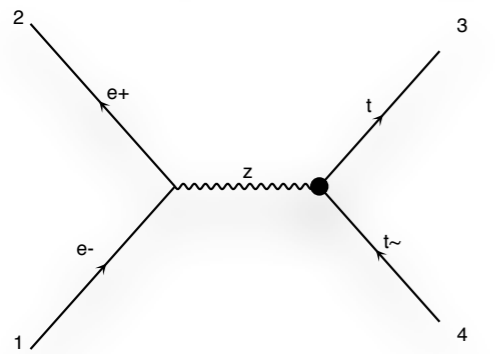
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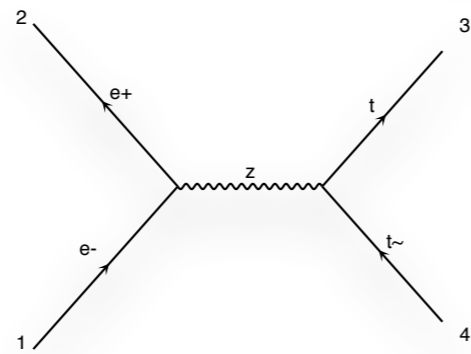
born diagram 1 EFT=1, QCD=0, QED=2



born diagram 2 EFT=0, QCD=0, QED=2



born diagram 3 EFT=1, QCD=0, QED=2



born diagram 4 EFT=0, QCD=0, QED=2

C. Radiative corrections

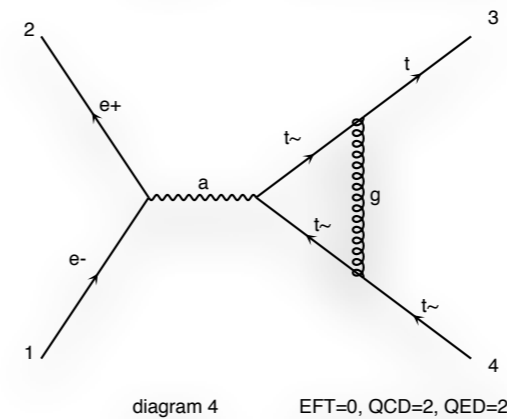
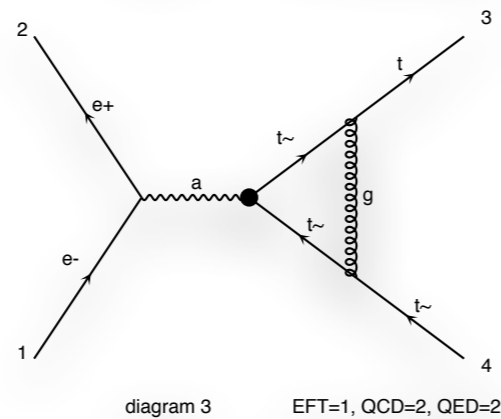
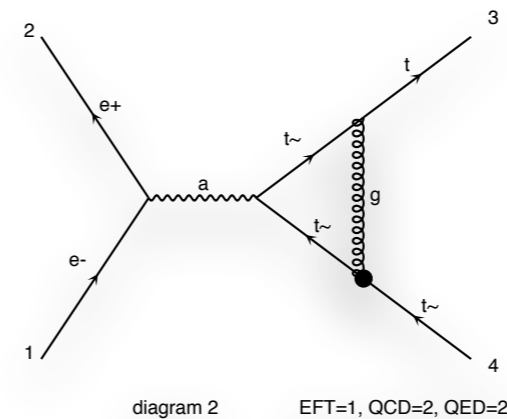
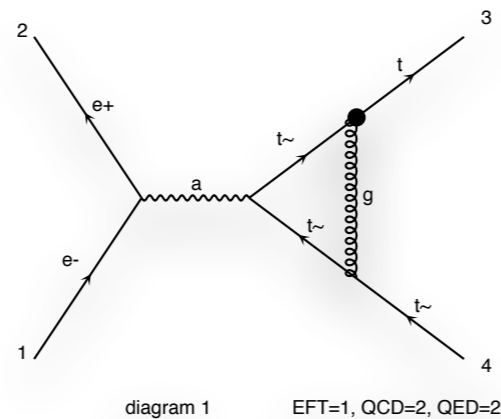
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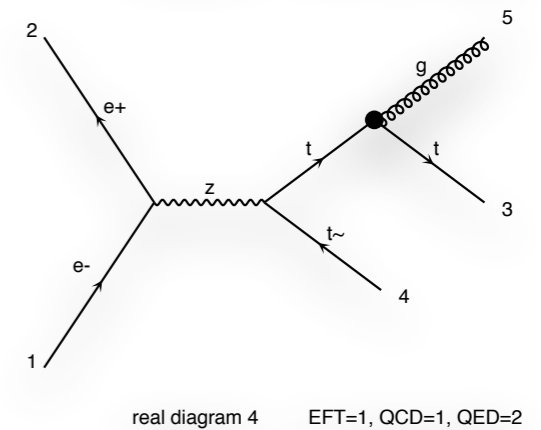
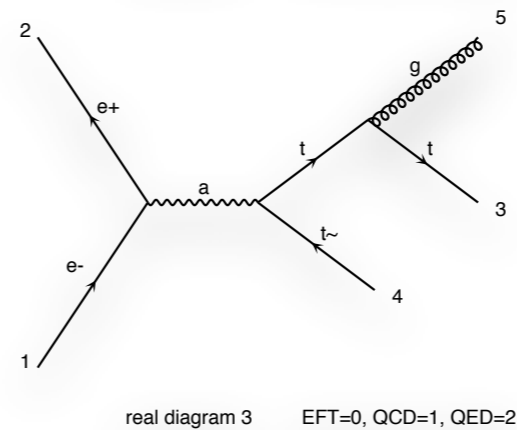
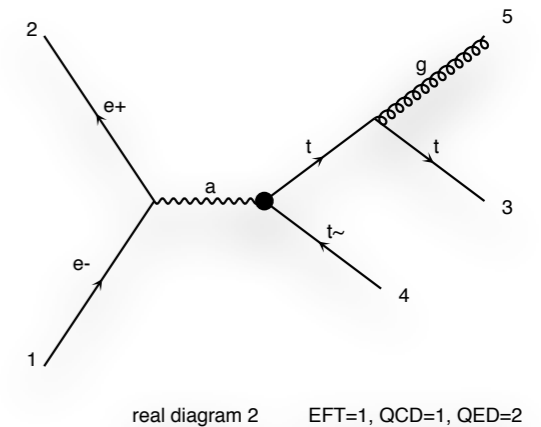
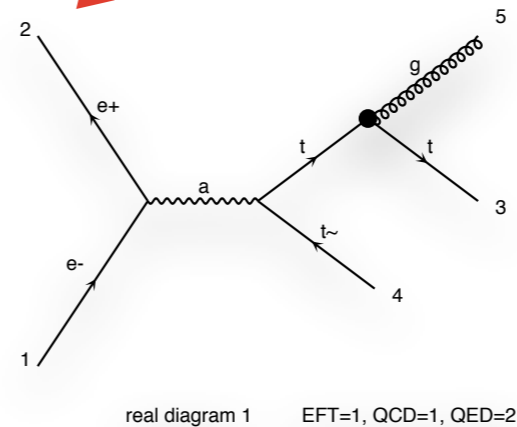
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C. Radiative corrections

$ee \rightarrow tt$ at NLO up to dim-6

$$\sigma = \sigma_{SM} + \sum_i \frac{C_i}{(\Lambda/1\text{TeV})^2} \sigma_i^{(1)} + \sum_{i,j}^{j \geq i} \frac{C_i C_j}{(\Lambda/1\text{TeV})^4} \sigma_{ij}^{(2)}$$

500GeV	SM	\mathcal{O}_{tG}	$\mathcal{O}_{\phi Q}^{(3)}$	$\mathcal{O}_{\phi Q}^{(1)}$	$\mathcal{O}_{\phi t}$	\mathcal{O}_{tW}	\mathcal{O}_{tB}
$\sigma_{i,LO}^{(1)}$	566	0	15.3	-15.0	-1.3	273	191
$\sigma_{i,NLO}^{(1)}$	647	-6.3	18.3	-18.5	-0.94	307	216
K-factor	1.14	N/A	1.20	1.23	0.7	1.12	1.13
$\sigma_{i,LO}^{(1)}/\sigma_{SM,LO}$		0	0.03	-0.03	-0.0022	0.48	0.34
$\sigma_{i,NLO}^{(1)}/\sigma_{SM,NLO}$		0.01	0.03	-0.03	-0.0014	0.47	0.33
$\sigma_{i,LO}^{(2)}/\sigma_{i,LO}^{(1)}$		N/A	0.06	-0.06	-0.56	0.22	0.14
$\sigma_{i,NLO}^{(2)}/\sigma_{i,NLO}^{(1)}$		\ll	0.04	-0.05	-0.9	0.22	0.14

Cross sections (in fb) for $t\bar{t}$ production at the ILC at $\sqrt{s} = 500$ GeV.

[1601.08193 Bylund, Maltoni, Tsirikos, Vryonidou ,CZ]

C. Radiative corrections

Process	O_{tG}	O_{tB}	O_{tW}	$O_{\varphi Q}^{(3)}$	$O_{\varphi Q}^{(1)}$	$O_{\varphi t}$	$O_{t\varphi}$	O_{4f}	$O_{\varphi G}$
$t \rightarrow bW \rightarrow bl^+\nu$	a		a	a				a	
$pp \rightarrow t\bar{q}$	A		A	A				A	
$pp \rightarrow tW$	A		A	A					
$pp \rightarrow t\bar{t}$	A							P	
$pp \rightarrow t\bar{t}\gamma$	A	A	A					P	
$pp \rightarrow t\gamma j$	A	A	A	A				P	
$pp \rightarrow t\bar{t}Z$	A	A	A	A	A	A		P	
$pp \rightarrow tZj$	A	A	A	A	A	A		P	
$pp \rightarrow t\bar{t}W$	A							P	
$e^+e^- \rightarrow t\bar{t}$	A	A	A	A	A	A		A	
$pp \rightarrow t\bar{t}H$	P						A	P	P
$pp \rightarrow tHj$	P		P	P			A	P	P
$gg \rightarrow H, Hj$	A						A		A
$gg \rightarrow HZ$	A			A	A	A	A		A

Table 1: Top-quark operators and key processes at the LHC, for the flavor-diagonal sector. Contributions which occur at the leading QCD order either at LO, or NLO, or through a top-quark loop, is marked by “a”, “A” or “P”. “a”: result is known analytically; “A”: NLO simulation within MG5_AMC is available; “P”: NLO simulation is not available currently, but is planned. O_{4f} denotes any four-fermion operator.

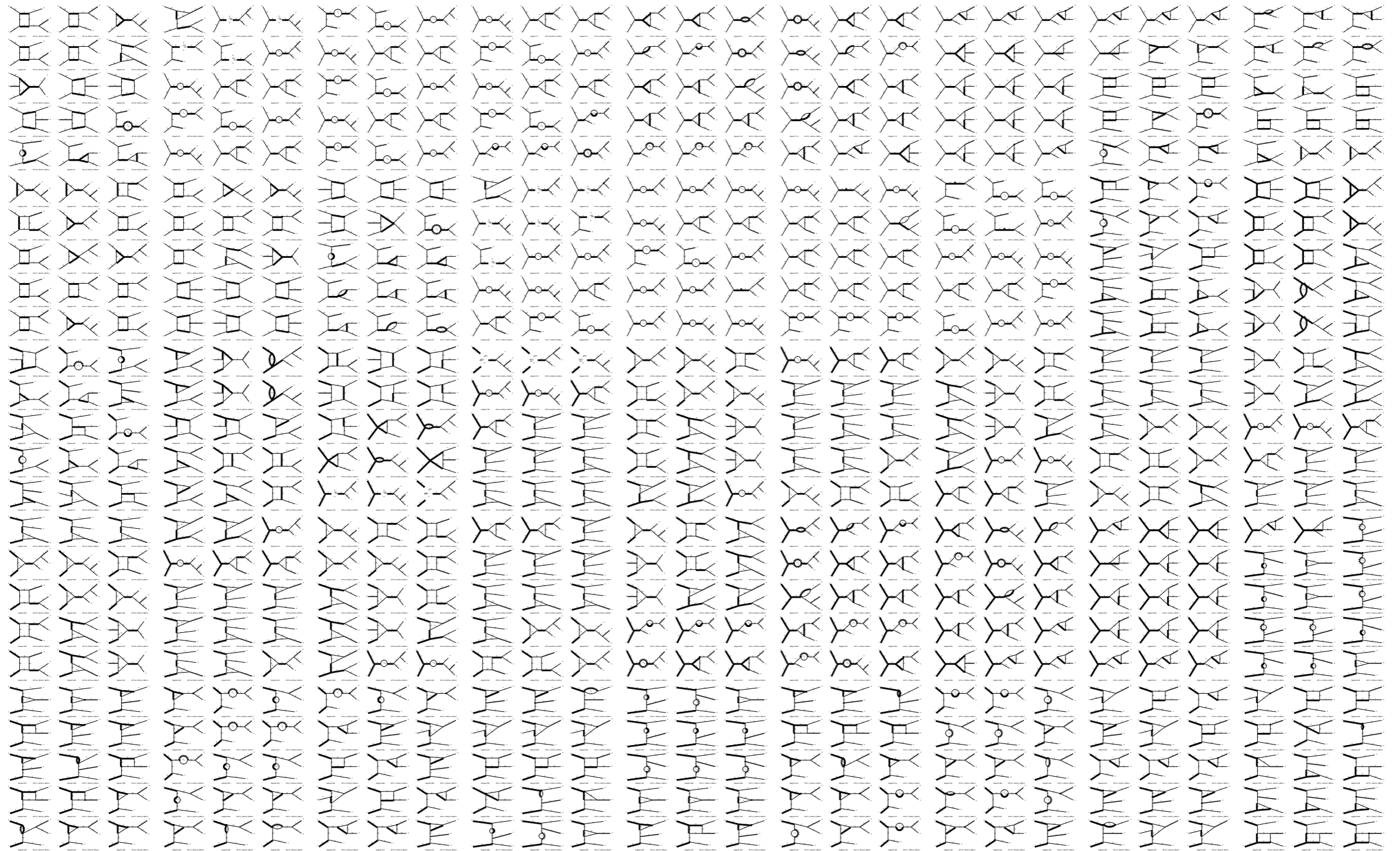
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ttZ production

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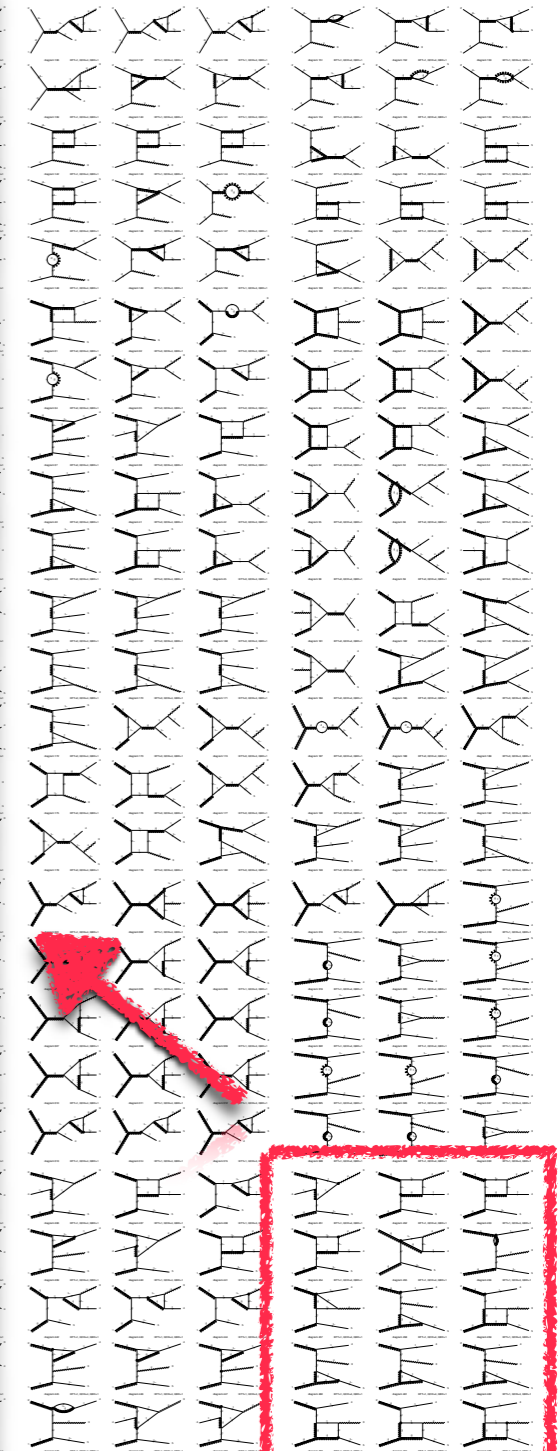
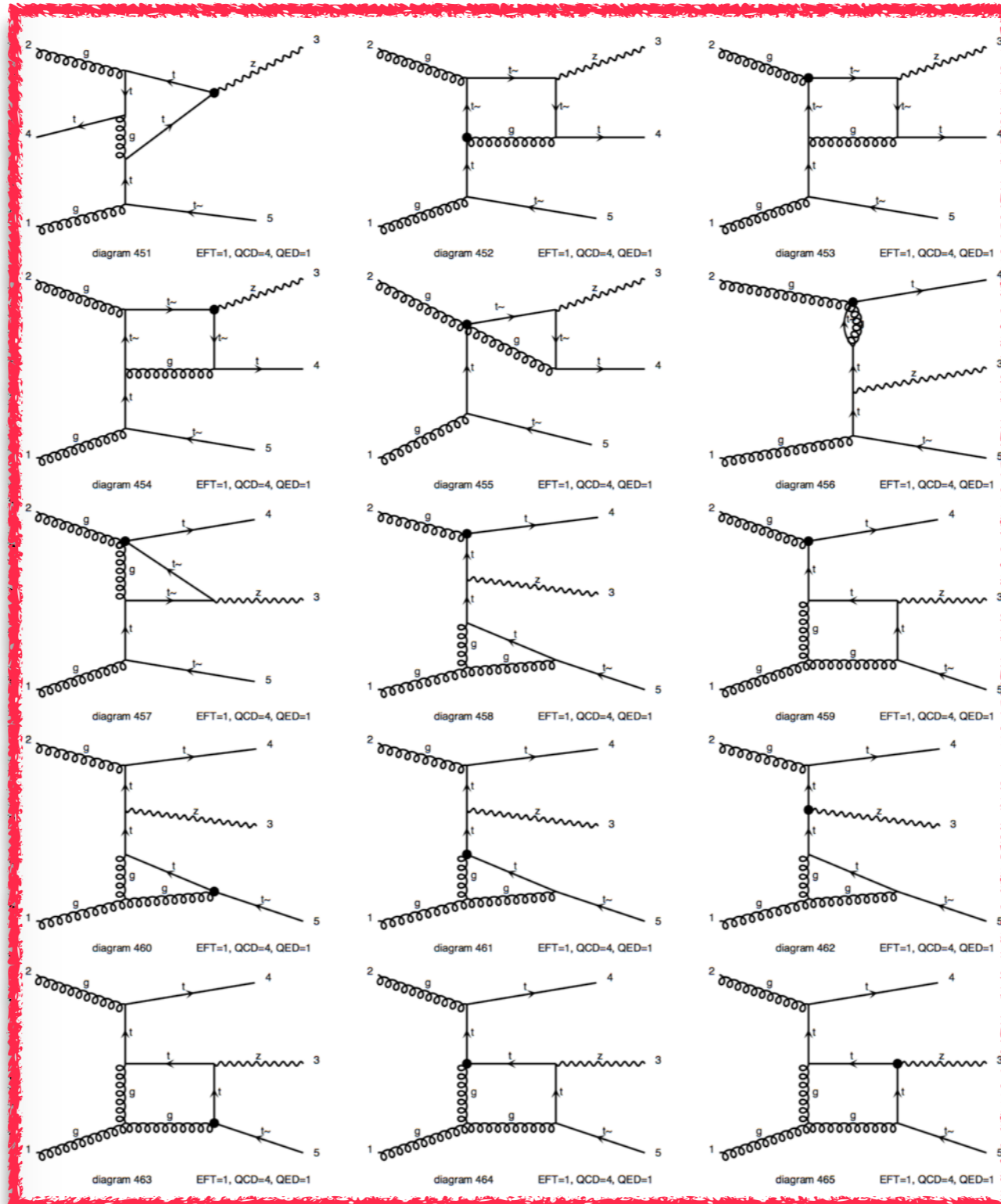
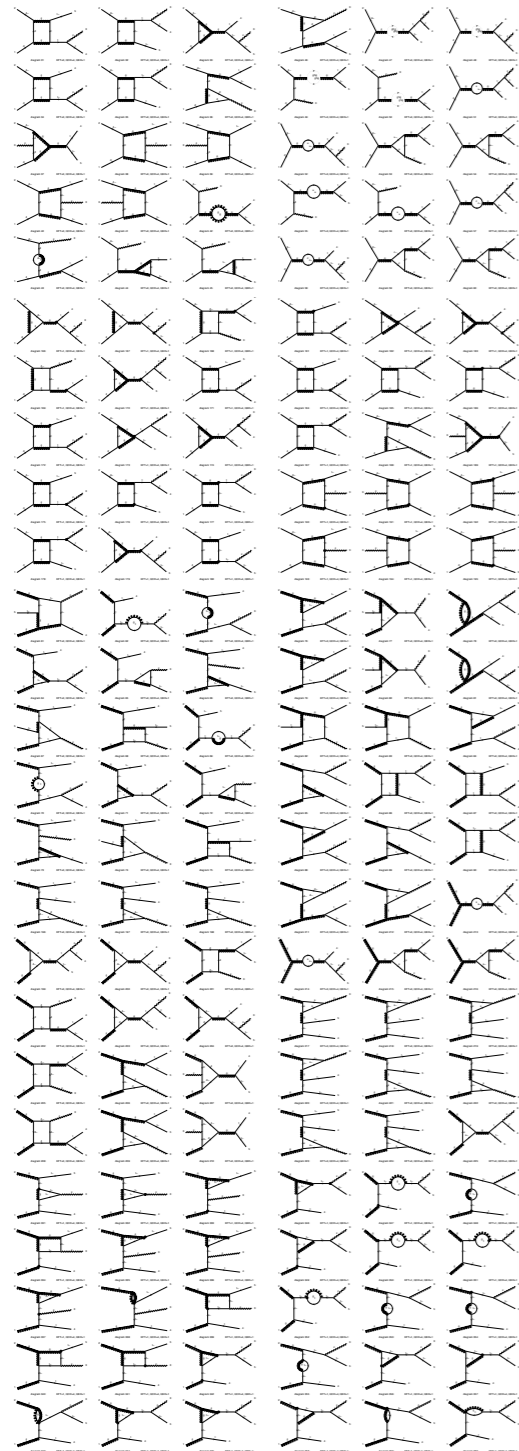
ttZ loops



Example #2: ttZ production

```
MG5_aMC>import model TEFT
MG5_aMC>generate p p > t t~ z EFT=1 [QCD]
MG5_aMC>output
MG5_aMC>launch
```

ttZ loops



C. Radiative corrections

EFT @ NLO:

we are not only advertising it as a possibility in principle; but we are making it possible in practice, in an automatic way, matched to MC generator, so that anyone can do NLO EFT in realistic analysis.

Top quark production

- $tt\gamma/tt g$, EM/color dipole

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu} \quad O_{tG} = (\bar{Q}\sigma^{\mu\nu}T^A t)\tilde{\varphi}G_{\mu\nu}^A$$

- tbW

- ▶ V/A

$$O_{\varphi Q}^{(3)} = i(\varphi^\dagger D_\mu \tau^I \varphi)(\bar{Q}\tau^I \gamma^\mu Q) \quad O_{\varphi\varphi} = i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu b)$$

- ▶ Weak dipole

$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I \quad O_{bW} = (\bar{Q}\sigma^{\mu\nu}\tau^I b)\varphi W_{\mu\nu}^I$$

- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q) \quad O_{\varphi u} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

- ▶ Weak dipole O_{tW}

$$O_{lQ}^{(1)} = (\bar{l}\gamma^\mu l)(\bar{Q}\gamma_\mu Q)$$

$$O_{eQ} = (\bar{e}\gamma^\mu e)(\bar{Q}\gamma_\mu Q)$$

$$O_{lt} = (\bar{l}\gamma^\mu l)(\bar{t}\gamma_\mu t)$$

$$O_{et} = (\bar{e}\gamma^\mu e)(\bar{t}\gamma_\mu t)$$

- $tt\gamma/ttg$, EM/color dipole

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$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I \quad O_{bW} = (\bar{Q}\sigma^{\mu\nu}\tau^I b)\varphi W_{\mu\nu}^I$$

- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q) \quad O_{\varphi t} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

- $tt\gamma/ttg$, EM/color dipole

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu} \quad O_{tG} = (\bar{Q}\sigma^{\mu\nu}T^A t)\tilde{\varphi}G_{\mu\nu}^A$$

- tbW

- ▶ V/A

$$O_{\varphi Q}^{(3)} = i(\varphi^\dagger D_\mu \tau^I \varphi)(\bar{Q}\tau^I \gamma^\mu Q)$$

$$O_{\varphi\varphi} = i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu b)$$

- ▶ Weak dipole

no interference

$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I$$

$$O_{bW} = (\bar{Q}\sigma^{\mu\nu}\tau^I b)\varphi W_{\mu\nu}^I$$

- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q)$$

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- tbW

- ▶ V/A

$$O_{\varphi Q}^{(3)} = i(\varphi^\dagger D_\mu \tau^I \varphi)(\bar{Q}\tau^I \gamma^\mu Q)$$

$$O_{\varphi\varphi} = i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu b)$$

- ▶ Weak dipole

no interference

$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I$$

$$O_{bW} = (\bar{Q}\sigma^{\mu\nu}\tau^I b)\varphi W_{\mu\nu}^I$$

- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q)$$

$$O_{\varphi t} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

Constraints from

- Pair production

- $tt\gamma/ttg$, EM/color dipole

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu}$$

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- ▶ Weak dipole

no interference

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$$O_{\varphi t} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

Constraints from

- Pair production
- Decay (W-hel)

- $tt\gamma/ttg$, EM/color dipole

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- ▶ V/A

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- ▶ Weak dipole

no interference

$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I$$

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- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q)$$

$$O_{\varphi t} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

Constraints from

- Pair production
- Decay (W-hel)
- Single top

- $tt\gamma/ttg$, EM/color dipole

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu}$$

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no interference

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- ttZ

- ▶ V/A

$$O_{\varphi Q}^{(1)} = i(\varphi^\dagger D_\mu \varphi)(\bar{Q}\gamma^\mu Q)$$

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Constraints from

- Pair production
- Decay (W-hel)
- Single top
- $Z \rightarrow bb$ (LEP)

will continue to be improved
at future LC

- $t\bar{t}\gamma/t\bar{t}g$, EM/color dipole

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- ▶ Weak dipole

no interference

$$O_{tW} = (\bar{Q}\sigma^{\mu\nu}\tau^I t)\tilde{\varphi}W_{\mu\nu}^I$$

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- $t\bar{t}Z$

- ▶ V/A

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$$O_{\varphi t} = i(\varphi^\dagger D_\mu \varphi)(\bar{t}\gamma^\mu t)$$

Constraints from

- Pair production
- Decay (W-hel)
- Single top
- $Z \rightarrow b\bar{b}$ (LEP)

will continue to be improved
at future LC

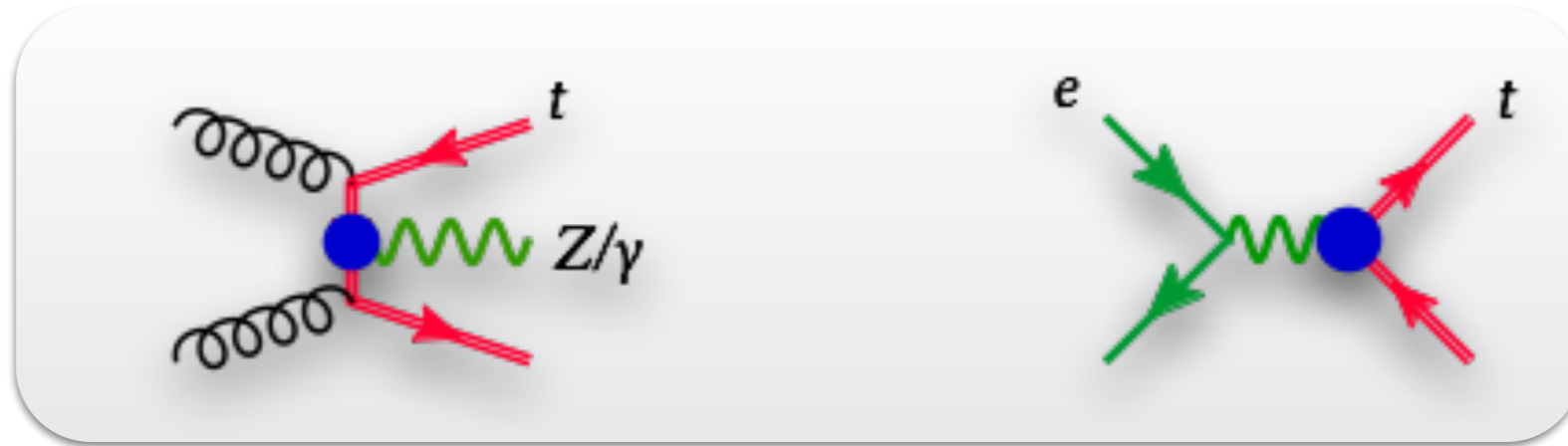
Neutral couplings can be
constrained from $t\bar{t}Z$ and $t\bar{t}\gamma$ at the
LHC, but

- Large uncertainties
- Much more sensitivity expected
at future lepton colliders

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu}$$

$$\frac{2m_t}{\Lambda^2} \frac{e}{s_W c_W} (C_{tB} s_W^2) \sigma^{\mu\nu} q_\nu$$

LHC vs LC



$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu}$$

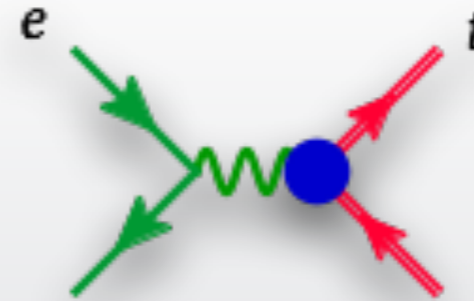
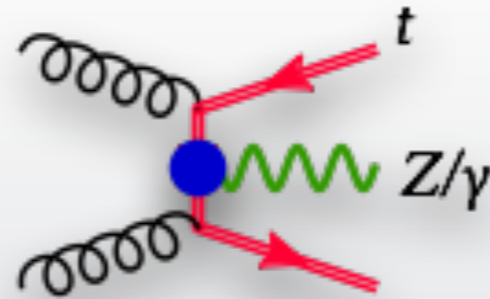
$$\frac{2m_t}{\Lambda^2} \frac{e}{s_W c_W} (C_{tB} s_W^2) \sigma^{\mu\nu} q_\nu$$

momentum plays an important role

LHC vs LC

Suppression:

SM favors soft radiation



Enhancement:

Large momentum guaranteed

$$O_{tB} = (\bar{Q}\sigma^{\mu\nu}t)\tilde{\varphi}B_{\mu\nu}$$

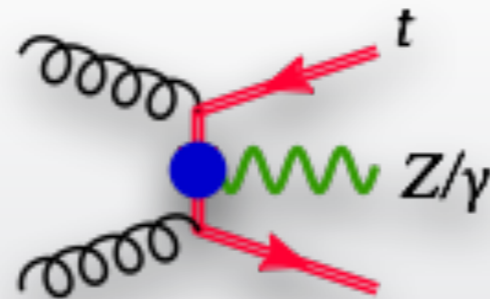
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momentum plays an important role

LHC vs LC

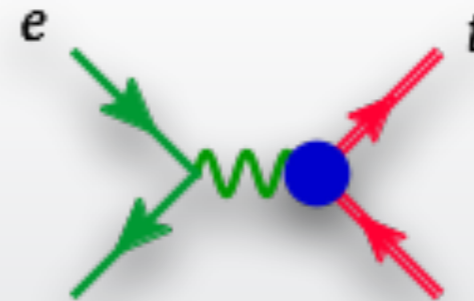
Suppression:

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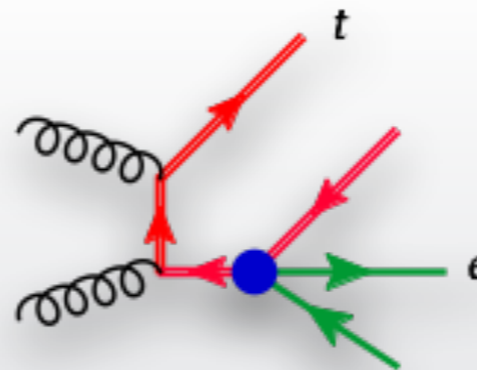
Enhancement:

Large momentum guaranteed



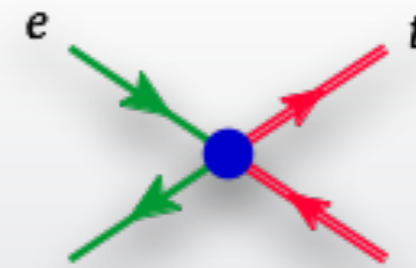
Four-fermion operators are similar

$$\begin{aligned} O_{lQ}^{(1)} &= (\bar{l}\gamma^\mu l) (\bar{Q}\gamma_\mu Q) \\ O_{eQ} &= (\bar{e}\gamma^\mu e) (\bar{Q}\gamma_\mu Q) \\ O_{lt} &= (\bar{l}\gamma^\mu l) (\bar{t}\gamma_\mu t) \\ O_{et} &= (\bar{e}\gamma^\mu e) (\bar{t}\gamma_\mu t) \end{aligned}$$



Suppression:

Four-body production



Enhancement:

No propagator (s/m²)

A comparison of HC and LC,
in terms of “**sensitivity**” = $\sigma_{\text{dim-6}}/\sigma_{\text{SM}}$
for $C/\Lambda^2 = 1 \text{ TeV}^{-2}$

Ops	PEWM bound	LHC		LC
		$t\bar{t}l^+l^-$ (13 TeV)	$t\bar{t}\gamma$ (13 TeV)	$ee \rightarrow t\bar{t}$ (500 GeV)
$O_{\phi t}$	[-3,7]	5%		-0.14%
O_{tB}	[-16,43]	0.1%(0.3%)	1%	33%
$O_{lQ}^{(1)}$		-1.7% (2.1%)		-468%
O_{eQ}		-0.4% (2.1%)		-190%
O_{lt}		-0.7% (2.1%)		-312%
O_{et}		-1.2% (2.1%)		-328%

proc.	TH	EXP	total
$pp \rightarrow t\bar{t}l^+l^-$	~ 10%	~ 15%	~ 20%
$pp \rightarrow t\bar{t}\gamma$	~ 15%	~ 25%	~ 30%
$e^-e^+ \rightarrow t\bar{t}$	< 1%	< 1%	~ 1%

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2. It also compensates the propagator suppression. Factor s/m^2 enhancement is expected.

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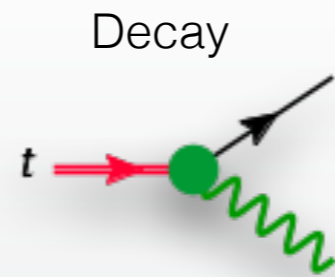
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There are operators for which the sensitivities of LHC and LC are very different. They are particularly interesting in the FCC-ee context and we shouldn't miss them.

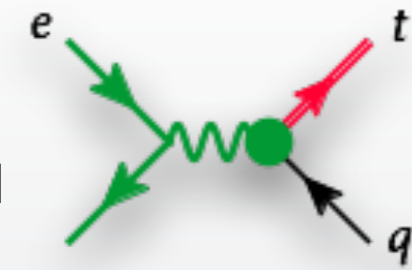
Top FCNC

LHC vs LC

2-fermion

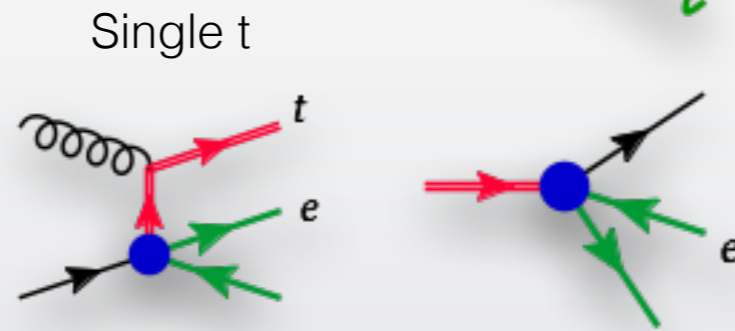


\approx comparable
[1311.2028 snowmass]

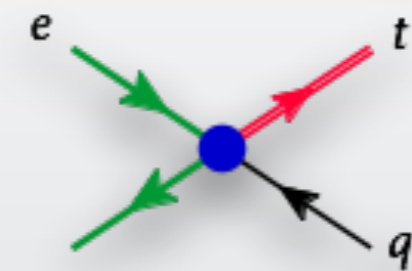


See also
Barbara Mele
this talk

4-fermion



\ll



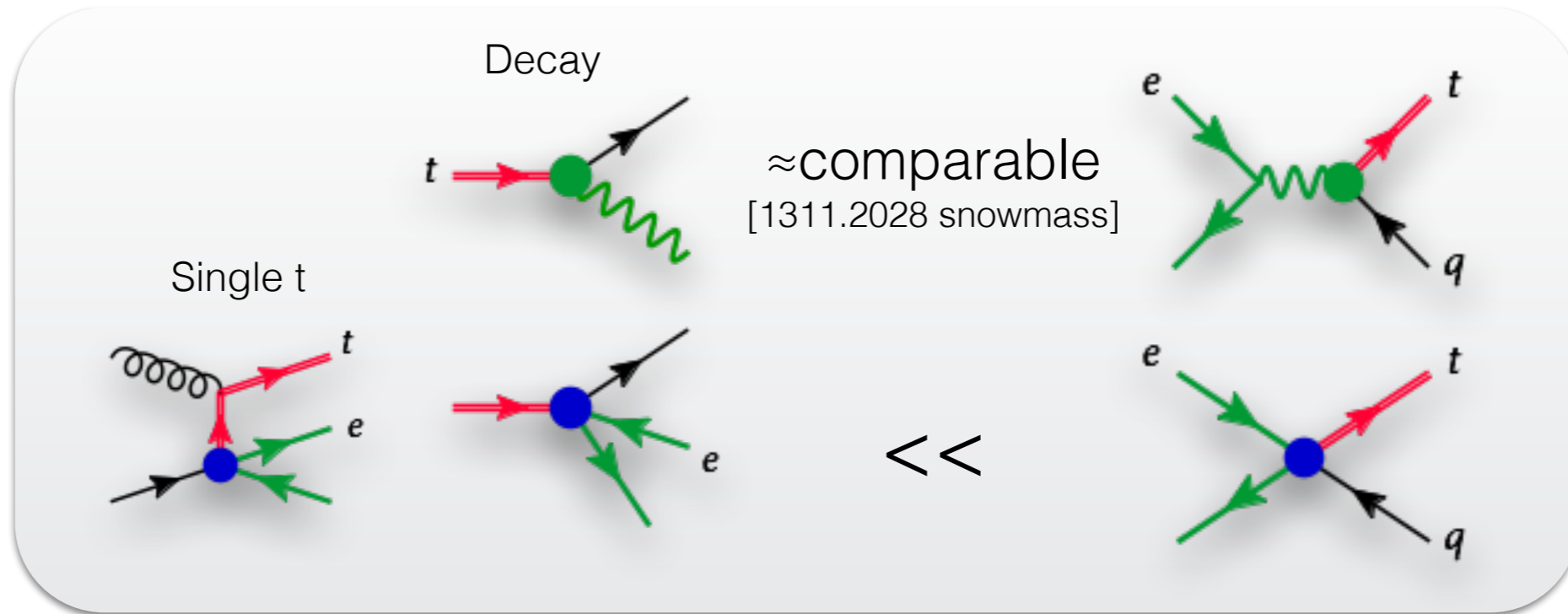
[1412.7166 Durieux, Maltoni, CZ]

Top FCNC

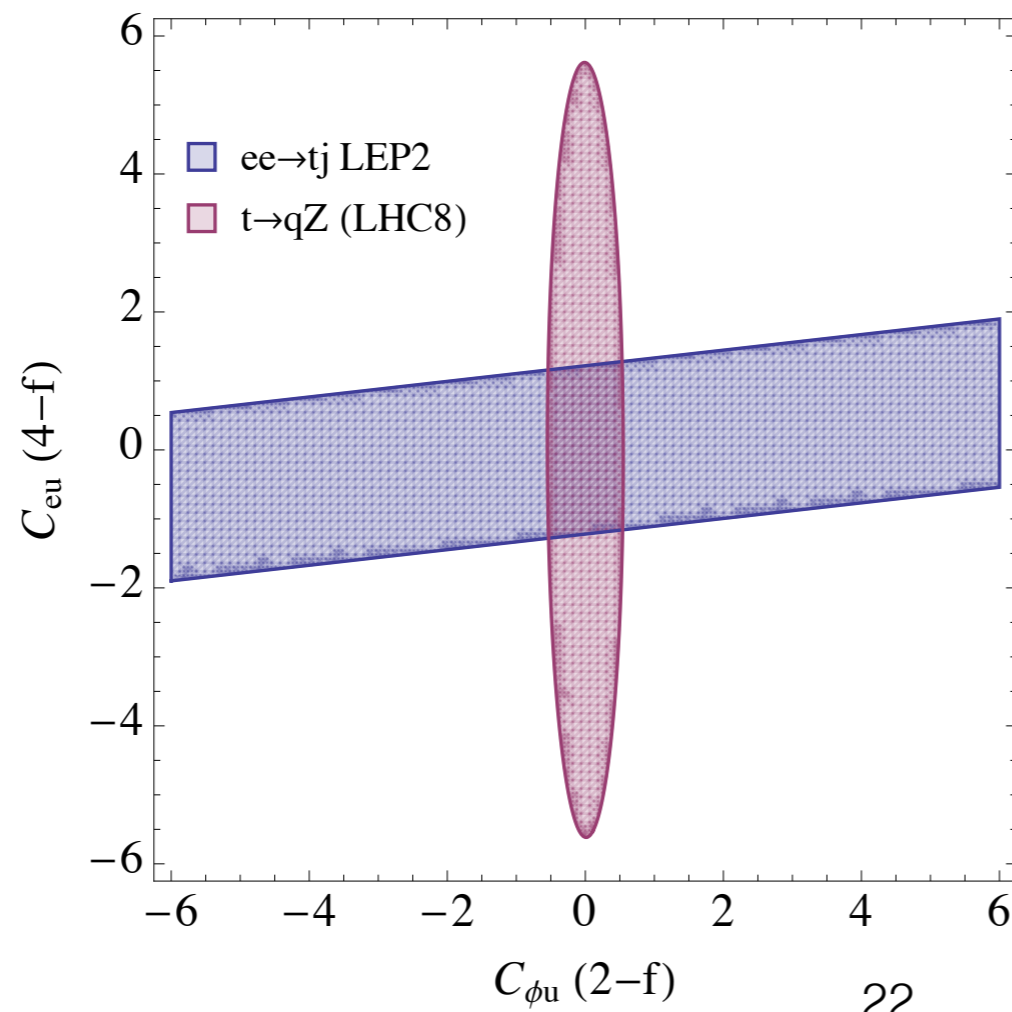
LHC vs LC

2-fermion

4-fermion



[1412.7166 Durieux, Maltoni, CZ]

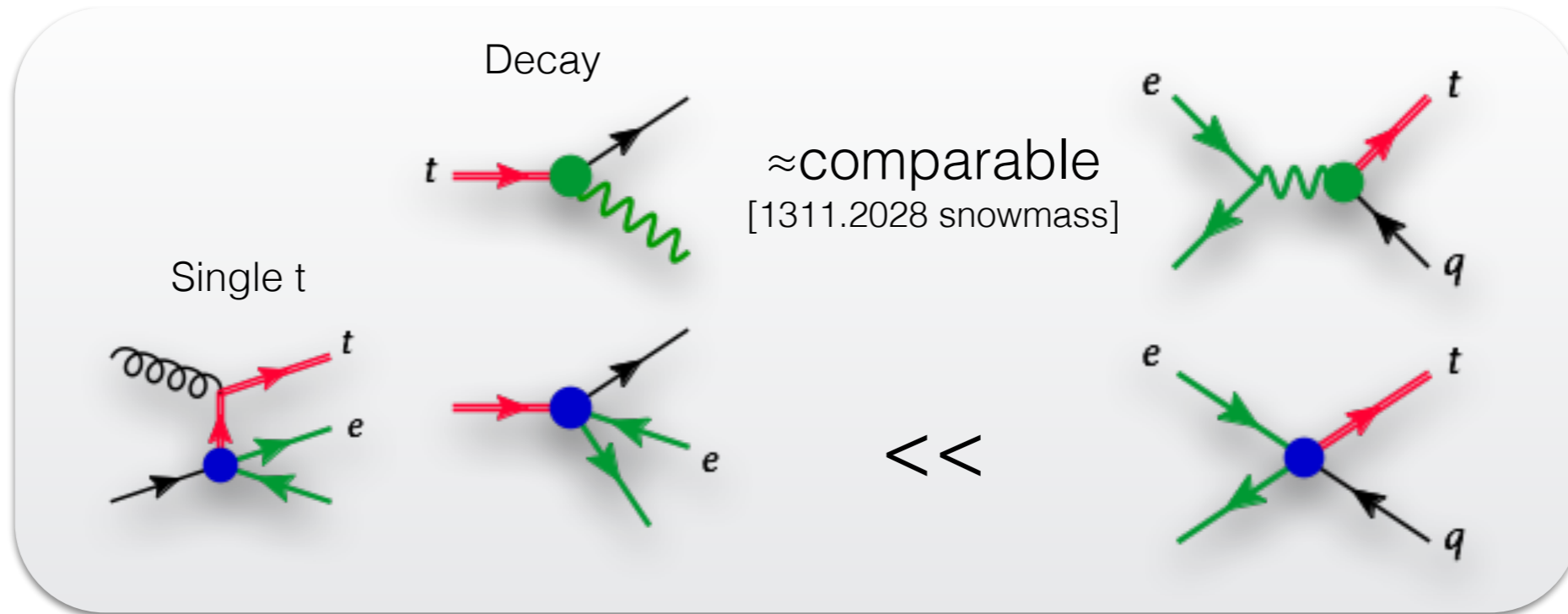


Top FCNC

LHC vs LC

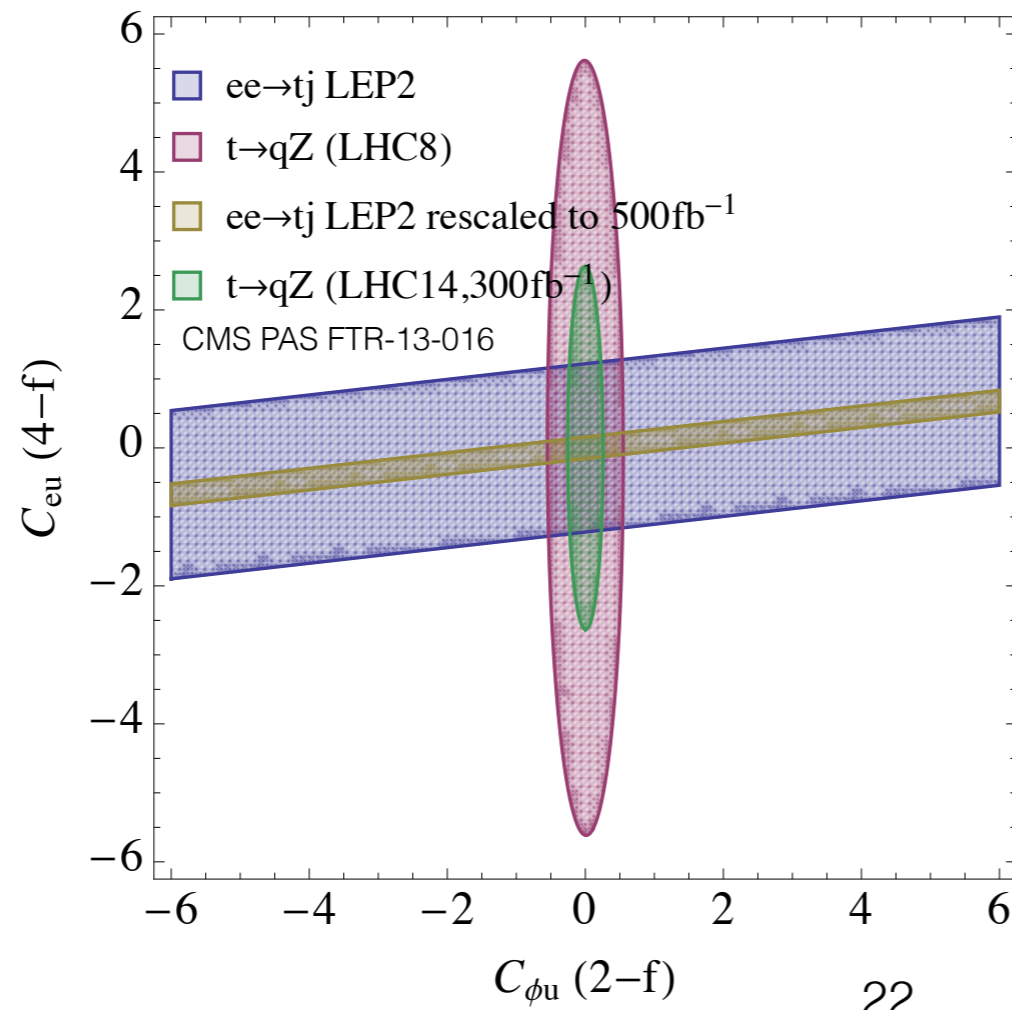
2-fermion

4-fermion



See also
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this talk

[1412.7166 Durieux, Maltoni, CZ]

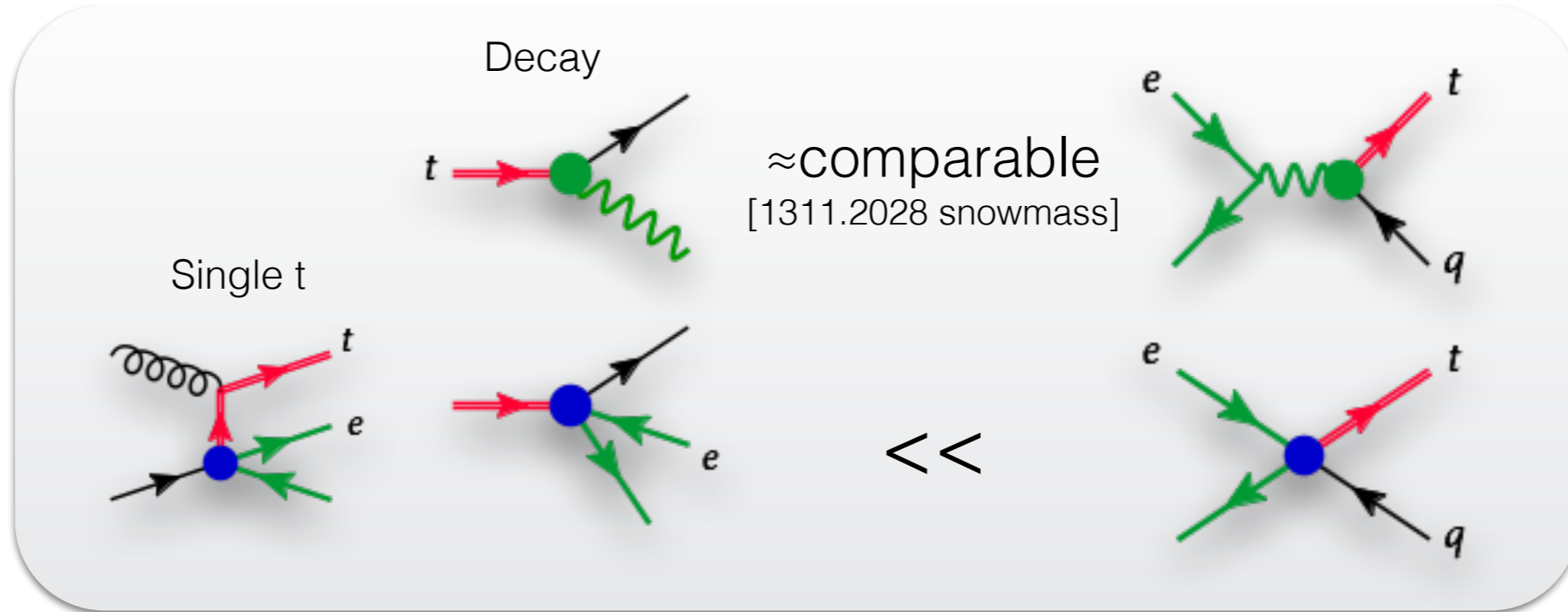


Top FCNC

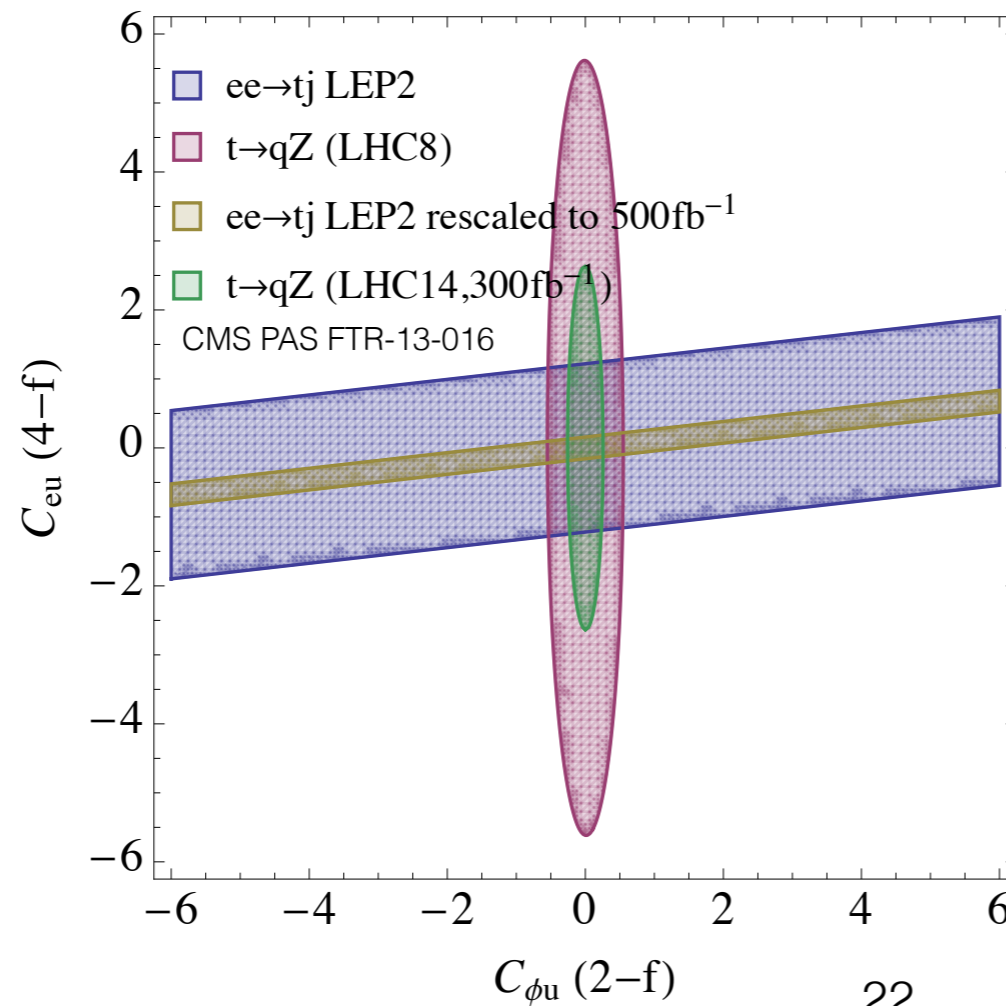
LHC vs LC

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4-fermion

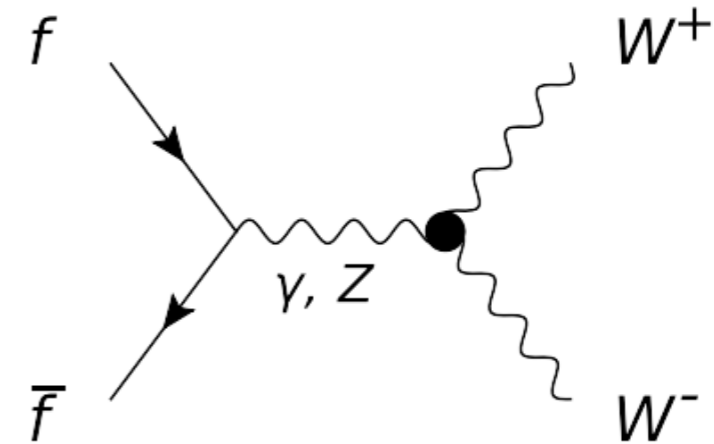


[1412.7166 Durieux, Maltoni, CZ]



- In terms of both 2-f and 4-f, LEP2 and LHC are complementary
- Same for HL-LHC and future ee collider

W-pair production



CP-even

$$\mathcal{O}_{WWW} = \text{Tr}[W_{\mu\nu}W^{\nu\rho}W_{\rho}^{\mu}]$$

$$\mathcal{O}_W = (D_{\mu}\Phi)^{\dagger}W^{\mu\nu}(D_{\nu}\Phi)$$

$$\mathcal{O}_B = (D_{\mu}\Phi)^{\dagger}B^{\mu\nu}(D_{\nu}\Phi),$$

CP-odd

$$\mathcal{O}_{\tilde{W}WW} = \text{Tr}[\tilde{W}_{\mu\nu}W^{\nu\rho}W_{\rho}^{\mu}]$$

$$\mathcal{O}_{\tilde{W}} = (D_{\mu}\Phi)^{\dagger}\tilde{W}^{\mu\nu}(D_{\nu}\Phi),$$

EW sector probed by PEWM+TGC+Higgs: 8+3+8

assuming flavor symmetry

1. PEWM (10^{-3})

$\mathcal{O}_W = \frac{ig}{2} \left(H^\dagger \overset{\leftrightarrow}{D}^\mu H \right) D^\nu W_{\mu\nu}^a$
(+) $\mathcal{O}_B = \frac{ig'}{2} \left(H^\dagger \overset{\leftrightarrow}{D}^\mu H \right) \partial^\nu B_{\mu\nu}$
$\mathcal{O}_T = \frac{1}{2} \left(H^\dagger \overset{\leftrightarrow}{D}_\mu H \right)^2$
$\mathcal{O}_{LL}^{(3)l} = (\bar{L}_L \sigma^a \gamma^\mu L_L) (\bar{L}_L \sigma^a \gamma_\mu L_L)$
$\mathcal{O}_R^e = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{e}_R \gamma^\mu e_R)$
$\mathcal{O}_R^u = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{u}_R \gamma^\mu u_R)$
$\mathcal{O}_R^d = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{d}_R \gamma^\mu d_R)$
$\mathcal{O}_L^{(3)q} = (iH^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \sigma^a \gamma^\mu Q_L)$
$\mathcal{O}_L^q = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \gamma^\mu Q_L)$

2. TGC (10^{-2})

$\mathcal{O}_W = \frac{ig}{2} \left(H^\dagger \sigma^a \overset{\leftrightarrow}{D}^\mu H \right) D^\nu W_{\mu\nu}^a$
(-) $\mathcal{O}_B = \frac{ig'}{2} \left(H^\dagger \overset{\leftrightarrow}{D}^\mu H \right) \partial^\nu B_{\mu\nu}$
$\mathcal{O}_{HB} = ig' (D^\mu H)^\dagger (D^\nu H) B_{\mu\nu}$
$\mathcal{O}_{3W} = \frac{1}{3!} g \epsilon_{abc} W_\mu^{a\nu} W_{\nu\rho}^b W^{c\rho\mu}$

3. New from Higgs

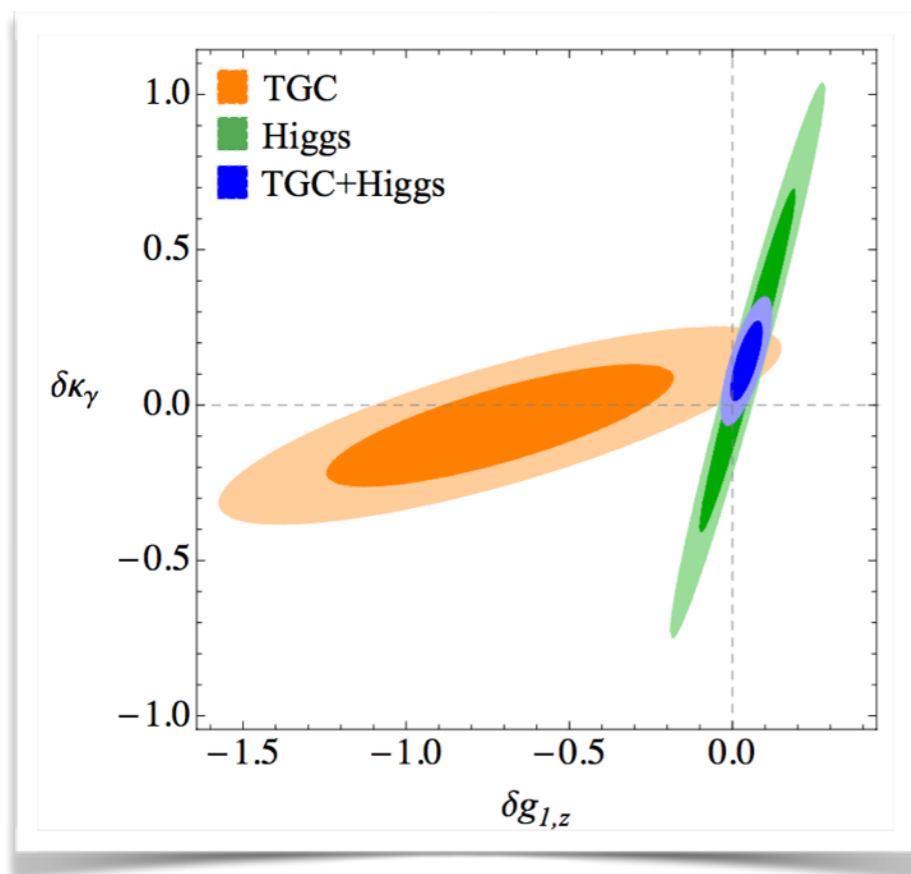
$\mathcal{O}_H = [\partial_\mu (H^\dagger H)]^2$
$\mathcal{O}_{BB} = \frac{g'^2}{4} H ^2 B_{\mu\nu} B^{\mu\nu}$
$\mathcal{O}_{WW} = \frac{g^2}{4} H ^2 W_{\mu\nu}^a W^{a\mu\nu}$
$\mathcal{O}_{GG} = \frac{g_s^2}{4} H ^2 G_{\mu\nu}^A G^{A\mu\nu}$
$\mathcal{O}_{y_u} = y_u H ^2 \bar{Q}_L \tilde{H} u_R$
$\mathcal{O}_{y_d} = y_d H ^2 \bar{Q}_L H d_R$
$\mathcal{O}_{y_e} = y_e H ^2 \bar{L}_L H e_R$
$\mathcal{O}_6 = \lambda H ^6$

W-pair production: $g_1^Z, \kappa_\gamma, \lambda_Z$

- In principle, WW also test 4 more parameters, from the PEWM class.
- They decouple because at LEP, they are measured with $\sim O(10)$ better accuracy.
- However this may change in future.

Mainly from Z-pole: only 8 d.o.f
(10 would appear in BW/Warsaw)

[A. Falkowski and F. Riva 2014]



“Accidental flat direction”

$$\delta g_{1,z} = -0.83 \pm 0.34, \quad \delta \kappa_\gamma = 0.14 \pm 0.05, \quad \lambda_Z = 0.86 \pm 0.38,$$

$$\rho = \begin{pmatrix} 1 & -0.71 & -0.997 \\ \cdot & 1 & 0.69 \\ \cdot & \cdot & 1 \end{pmatrix} \quad [1411.0669 \text{ Falkowski and Riva}]$$

To lift the degeneracy,

- Add Higgs data
- Use polarized beam
- Go to higher energy

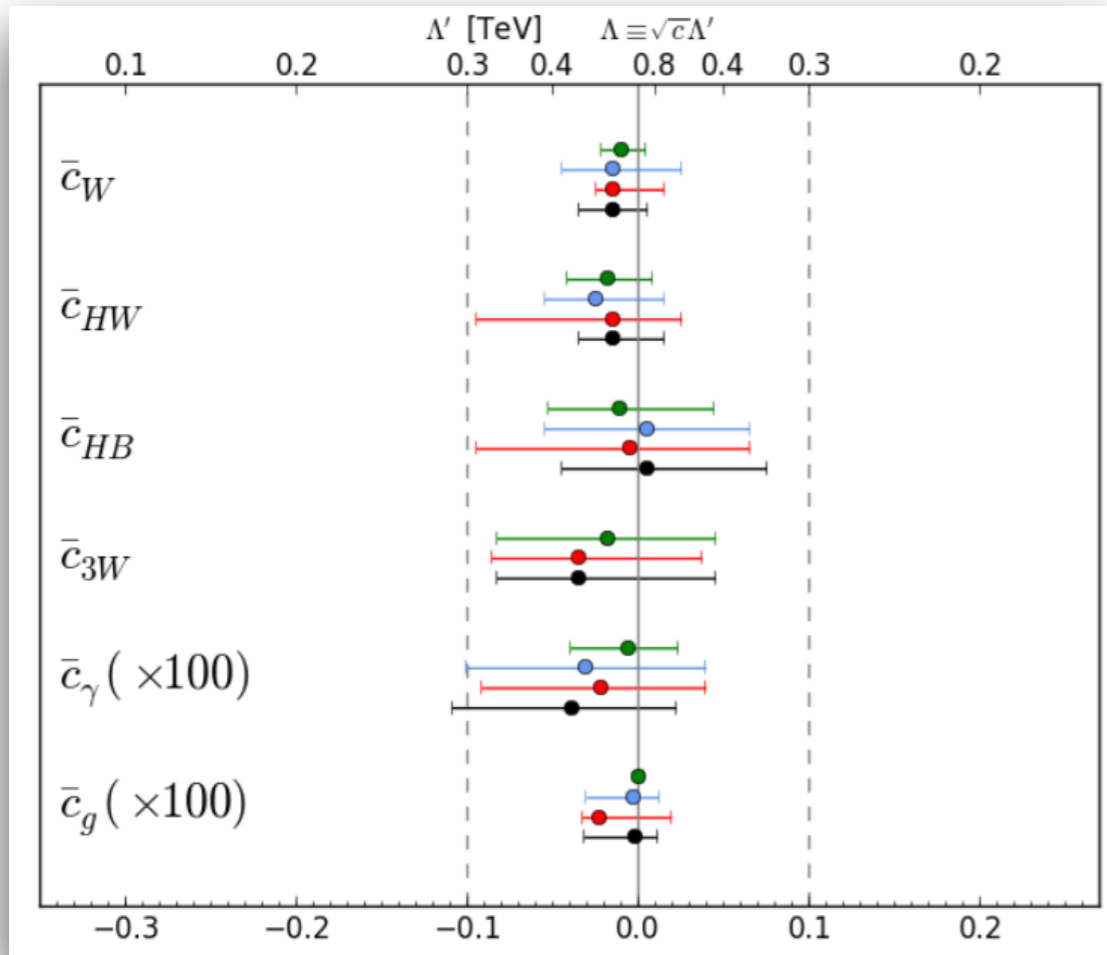
$$\begin{aligned} 200 \text{ GeV} \quad \bar{\delta}^{\text{NP}} \sigma_L &= 0.0192(c_{HW} + c_W) + 0.00345(c_{HW} + c_{HB} - 4c_{WB}) + 0.00667c_{3W} \\ \bar{\delta}^{\text{NP}} \sigma_R &= -1.32(c_{HW} + c_W) + 0.640(c_{HW} + c_{HB} - 4c_{WB}) - 0.0898c_{3W} \end{aligned}$$

$$\begin{aligned} 500 \text{ GeV} \quad \bar{\delta}^{\text{NP}} \sigma_L &= 0.0835(c_{HW} + c_W) + 0.0277(c_{HW} + c_{HB} - 4c_{WB}) + 0.0191c_{3W} \\ \bar{\delta}^{\text{NP}} \sigma_R &= -8.25(c_{HW} + c_W) + 6.64(c_{HW} + c_{HB} - 4c_{WB}) - 0.0426c_{3W} \end{aligned}$$

[1507,01594 D. Wells and Z. Zhang]

LHC current status

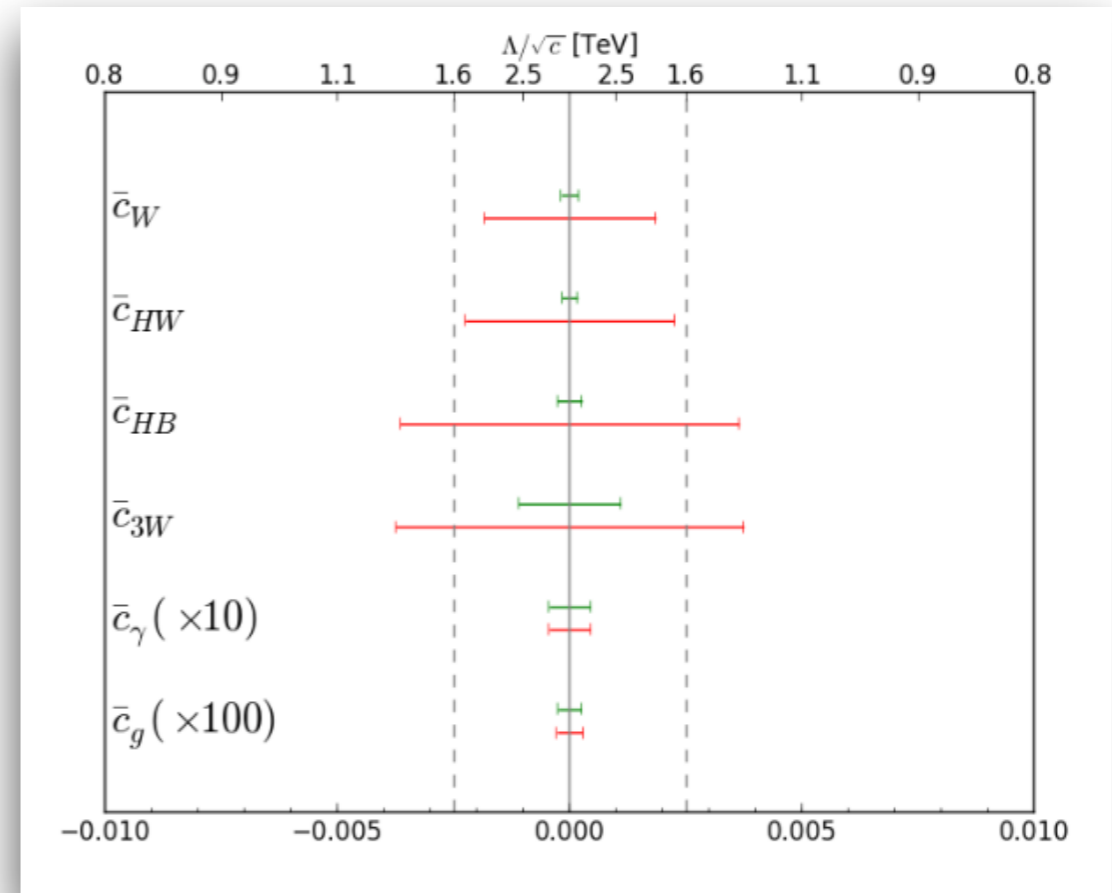
[1404.3667, 1410.7703 Ellis, Sanz, You]



Higgs, TGC, combination

FCC-ee projection

[1510.04561 Ellis, You]



single, marginalized

$O(10^2)$ improvement

ILC projection [1506.07830 Barklow] [J. List]	g_1^Z	8.5×10^{-4}
	κ_γ	9.2×10^{-4}
	λ_γ	7×10^{-4}

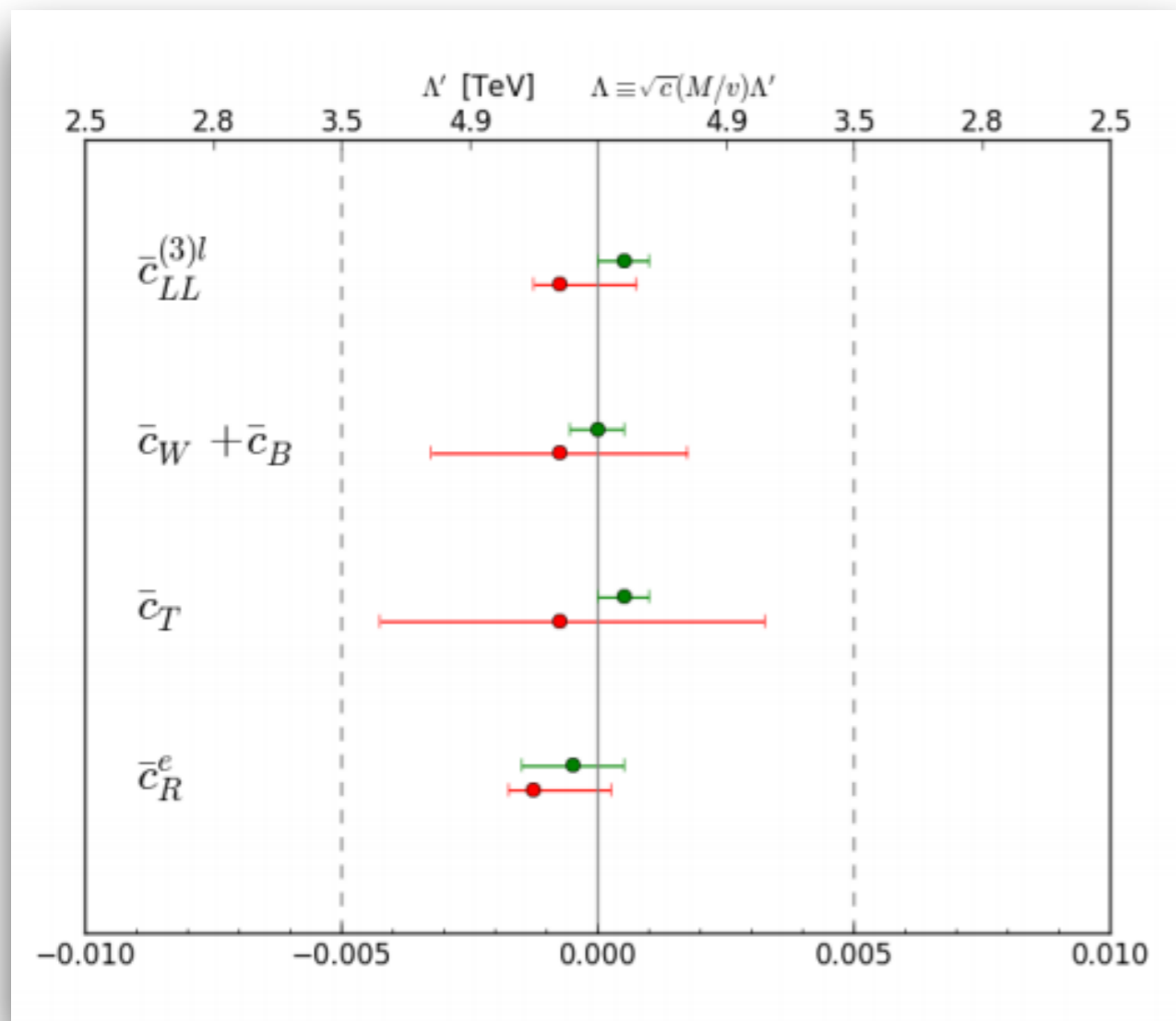
Summary

- EFT approach to top/W physics at future linear colliders.
- In EFTtheorests' mind, one should make the most use of SM EFT, by
 - **Taking it as a global approach**, include current and future measurements in one framework.
 - Apart from anomalous couplings, investigating all possible SM deviations, **in particular 4-fermions** are very interesting in this context.
 - Including **higher-order corrections in EFT are becoming available**. More than just a “possibility in principle”, but will become possible for anybody in any analysis.

Thank you for your attention

LHC current status

[1404.3667, 1410.7703 Ellis, Sanz, You]



FCC-ee projection

[1510.04561 Ellis, You]

