

S($\gamma\gamma$) and FCCee

Fitting the $\gamma\gamma$ peak:

- 1) Widths
- 2) Models
- 3) Theories
- 4) What next?

Alessandro Strumia
talk at the 10th FCC-ee physics workshop
CERN, February 5, 2016



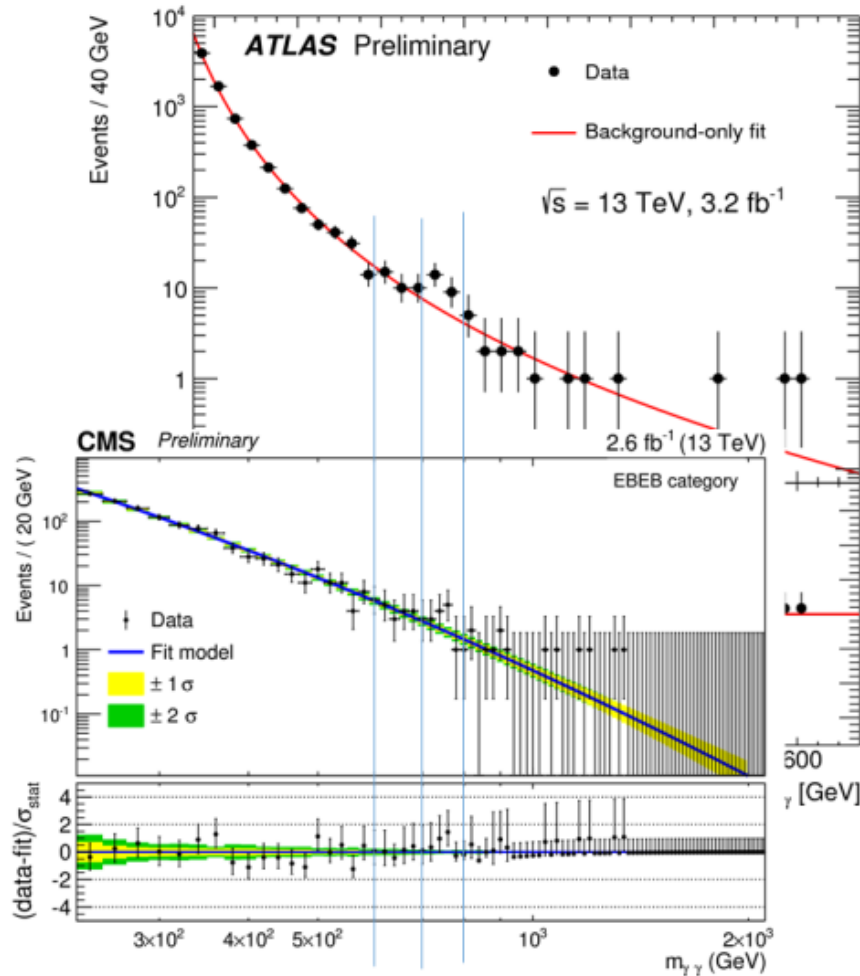
European
Commission

Horizon 2020
European Union funding
for Research & Innovation



European
Research
Council

The data



Flatland plus a $\gamma\gamma$ peak around 750 GeV

$\sigma(pp \rightarrow \gamma\gamma)$	8 TeV	13 TeV
CMS	$(0.5 \pm 0.6) \text{ fb}$	$(6 \pm 3) \text{ fb}$
ATLAS	$(0.4 \pm 0.8) \text{ fb}$	$(10 \pm 3) \text{ fb}$

Theoretically clean

Experimentally simple

ATLAS prefers large width $\Gamma/M \sim 0.06$.

CMS prefers narrow width.

$\gamma\gamma$ not accompanied by hard extras.

Full energy distribution? Angular distribution? Full events?

Needless to say

Maybe a fluke.

Gold does not come to you spontaneously.

The Gold Rush

INSPIRES list

Date	papers
16 Dec	10
19 Dec	46
25 Dec	101
1 Jan	137
1 Feb	212
1 Apr	?

All that is gold does not glitter

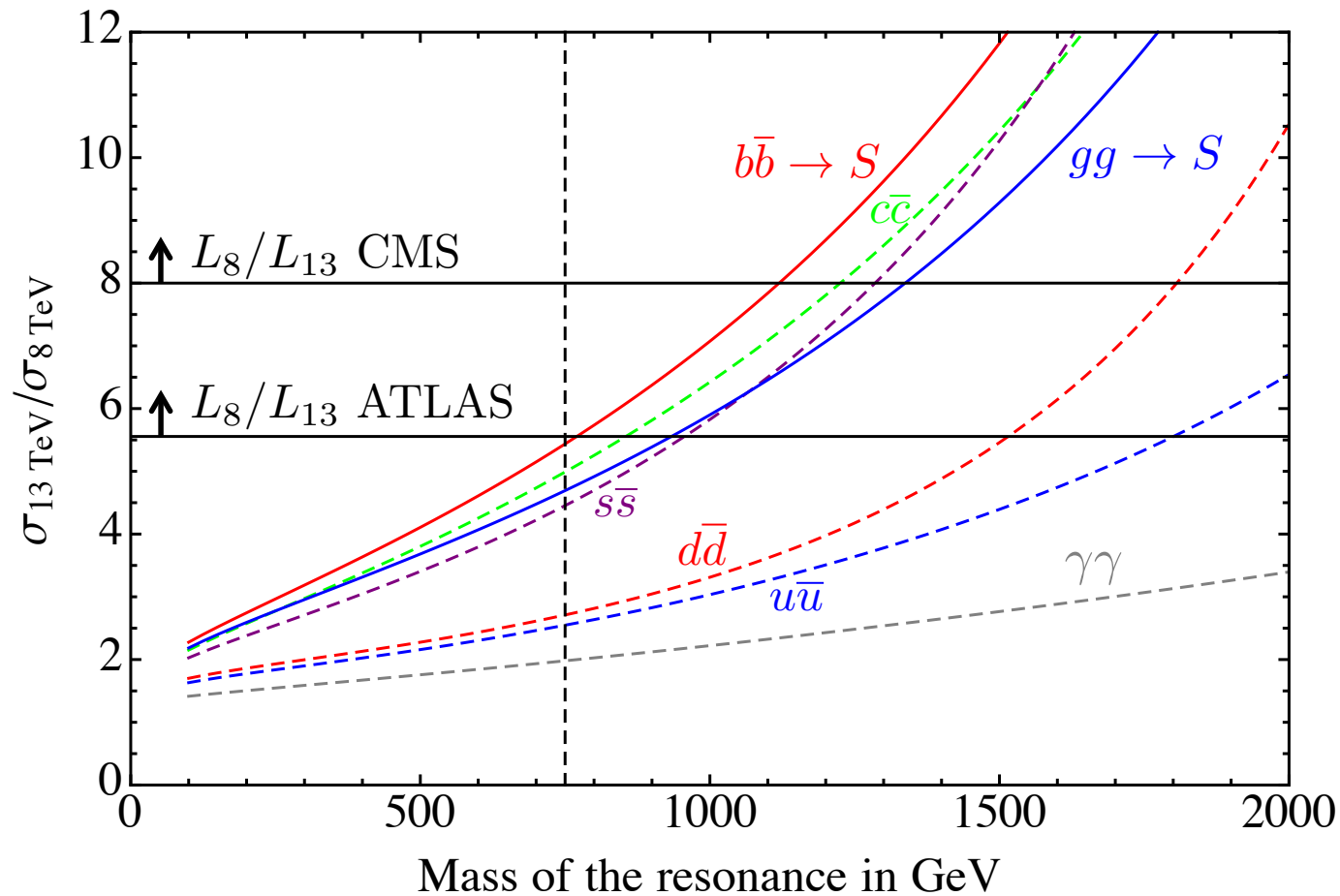
physics = experiment + *i* theory

It's time to present a
review of the new boson.



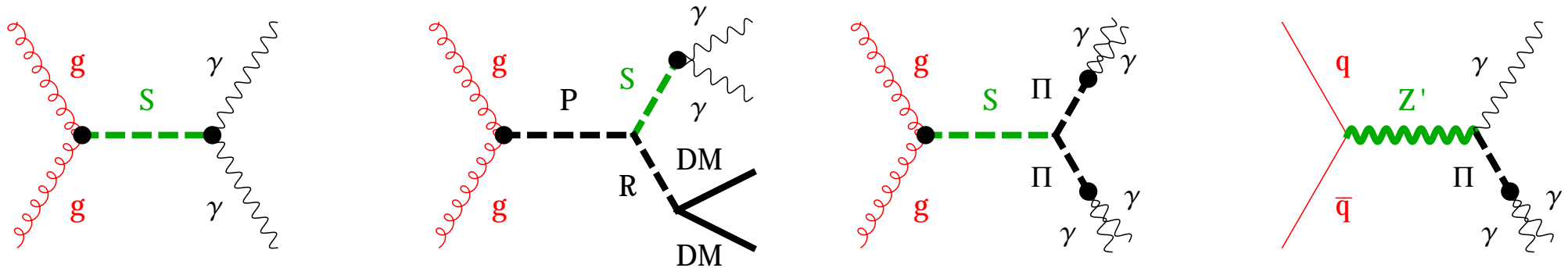
A new boson at 750 GeV?

Run 1 compatible with Run 2 if S is produced as gg , $b\bar{b}$, $c\bar{c}$, $s\bar{s}$.
The SM background $q\bar{q} \rightarrow \gamma\gamma$ at 750 GeV grows only by 2.3



A more complicated kinematics?

Compatibility between runs 1, 2 improved if S decays from a heavier particle.



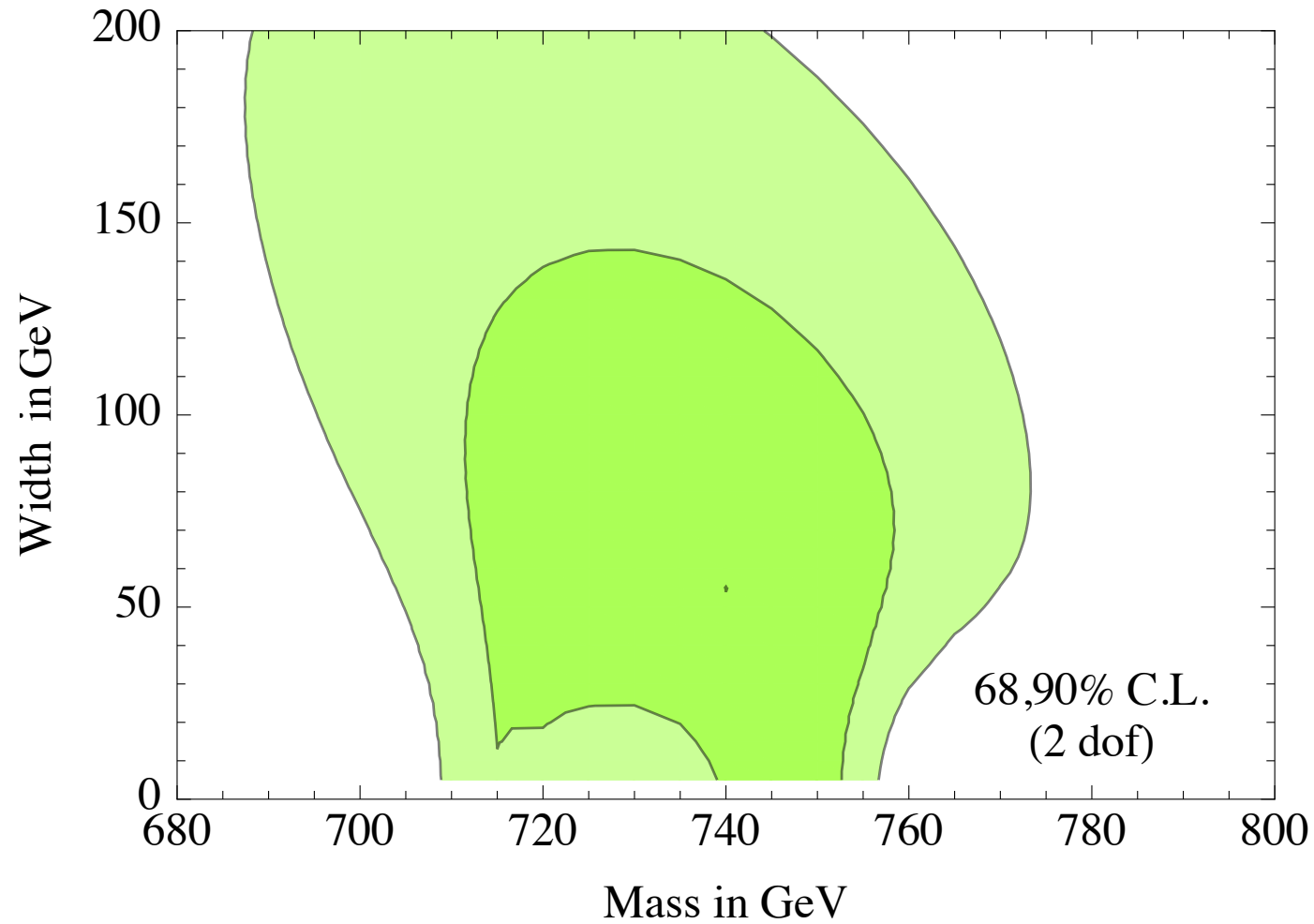
Tuning $M_P \approx M_S + M_R$ needed to avoid p_T . S virtuality can fake S width.

Or large $S \rightarrow \Pi\Pi$ with $\Pi \rightarrow \gamma\gamma$, collimated and seen as a single γ if $M_\Pi \ll M_S$.

Or many collimated γ . Or not a peak. Or two nearby narrow resonances.

Widths

M and Γ from data



Apology: public ATLAS data only

Cross section

It can be computed in terms of (narrow) widths:

$$\sigma(pp \rightarrow S \rightarrow \gamma\gamma) = \frac{2J_S + 1}{M\Gamma_S} \left[\sum_{\wp} C_{\wp\bar{\wp}} \Gamma(S \rightarrow \wp\bar{\wp}) \right] \Gamma(S \rightarrow \gamma\gamma)$$

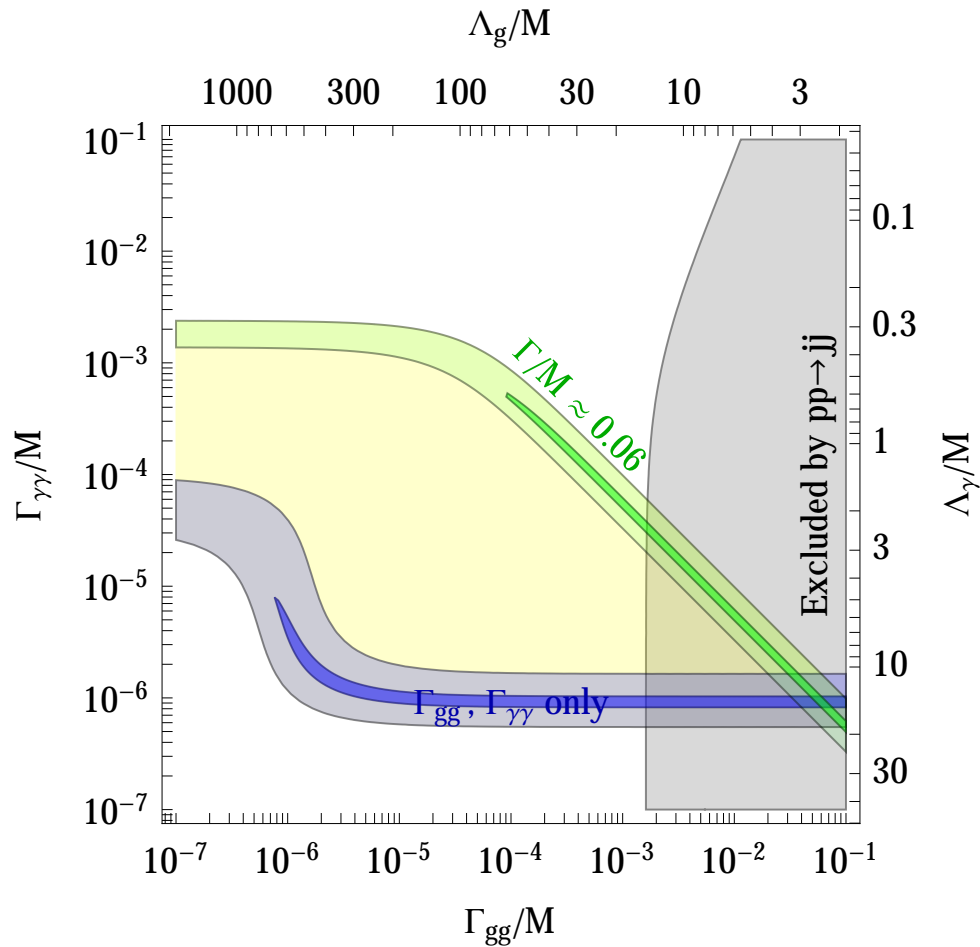
Parton \wp luminosities:

\sqrt{s}	$C_{b\bar{b}}$	$C_{c\bar{c}}$	$C_{s\bar{s}}$	$C_{d\bar{d}}$	$C_{u\bar{u}}$	C_{gg}	$C_{\gamma\gamma}$
8 TeV	1.07	2.7	7.2	89	158	174	54
13 TeV	15.3	36	83	627	1054	2137	11

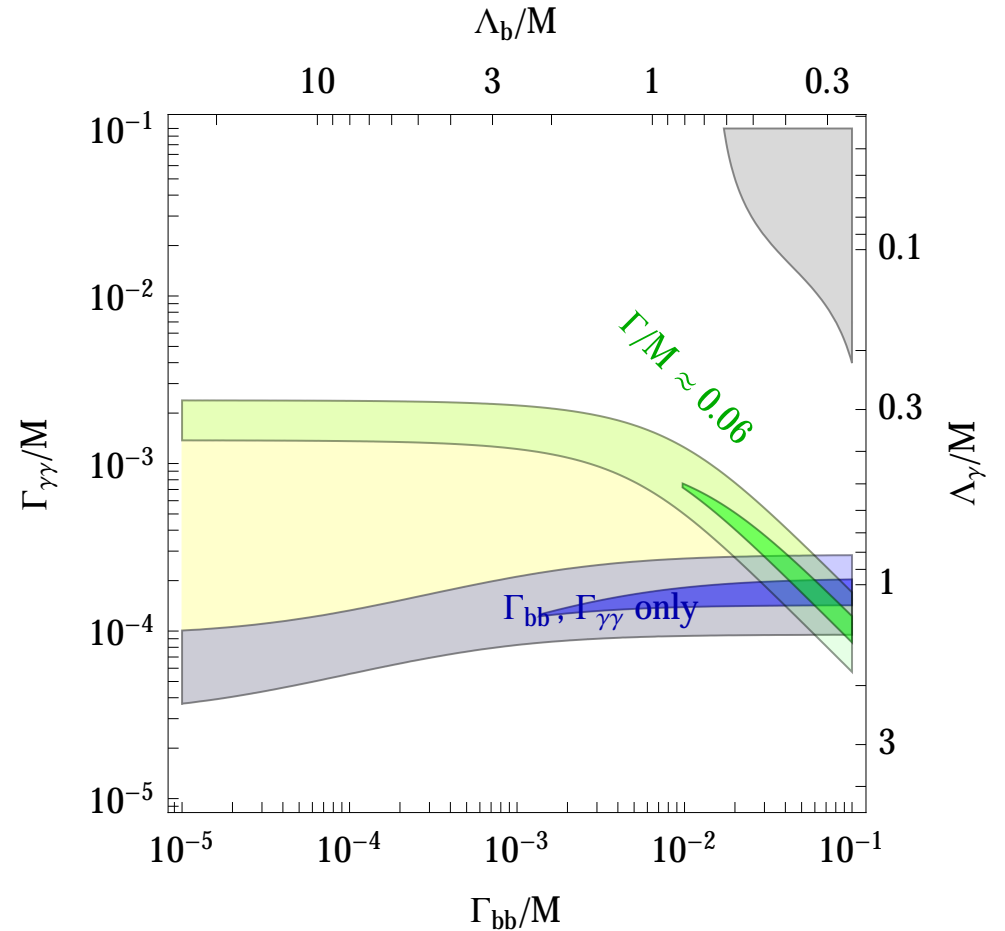
(Partonic $\gamma\gamma \rightarrow S \rightarrow \gamma\gamma$ would be minimal but run1/run2 compatibility is poor).

Extreme cases: gg and $b\bar{b}$

$$\mathcal{L}_{\text{scalar}} = S \left[g_3^2 \frac{G_{\mu\nu}^2}{2\Lambda_g} + e^2 \frac{F_{\mu\nu}^2}{2\Lambda_\gamma} + \frac{HQ_3 D_3}{\Lambda_b} \right] \quad \text{or} \quad \mathcal{L}_{\text{pseudo scalar}} = S \left[g_3^2 \frac{G_{\mu\nu} \tilde{G}_{\mu\nu}}{2\tilde{\Lambda}_g} + e^2 \frac{F_{\mu\nu} \tilde{F}_{\mu\nu}}{2\tilde{\Lambda}_\gamma} + \frac{HQ_3 i\gamma_5 D_3}{\tilde{\Lambda}_b} \right]$$



$S \leftrightarrow \gamma\gamma, gg, ?$



$S \leftrightarrow \gamma\gamma, b\bar{b}$

Bounds on other decay modes

final state f	σ at $\sqrt{s} = 8$ TeV		implied bound on $\Gamma(S \rightarrow f)/\Gamma(S \rightarrow \gamma\gamma)_{\text{obs}}$
	observed	expected	
$\gamma\gamma$	< 1.5 fb	< 1.1 fb	$< 0.8 (r/5)$
$e^+e^-, \mu^+\mu^-$	< 1.2 fb	< 1.2 fb	$< 0.6 (r/5)$
$\tau^+\tau^-$	< 12 fb	< 15 fb	$< 6 (r/5)$
$Z\gamma$	< 11 fb	< 12 fb	$< 6 (r/5)$
ZZ	< 12 fb	< 20 fb	$< 6 (r/5)$
Zh	< 19 fb	< 28 fb	$< 10 (r/5)$
hh	< 39 fb	< 42 fb	$< 20 (r/5)$
W^+W^-	< 40 fb	< 70 fb	$< 20 (r/5)$
$t\bar{t}$	< 450 fb	< 600 fb	$< 300 (r/5)$
invisible	< 0.8 pb	-	$< 400 (r/5)$
$b\bar{b}$	$\lesssim 1$ pb	$\lesssim 1$ pb	$< 500 (r/5)$
jj	$\lesssim 2.5$ pb	-	$< 1300 (r/5)$

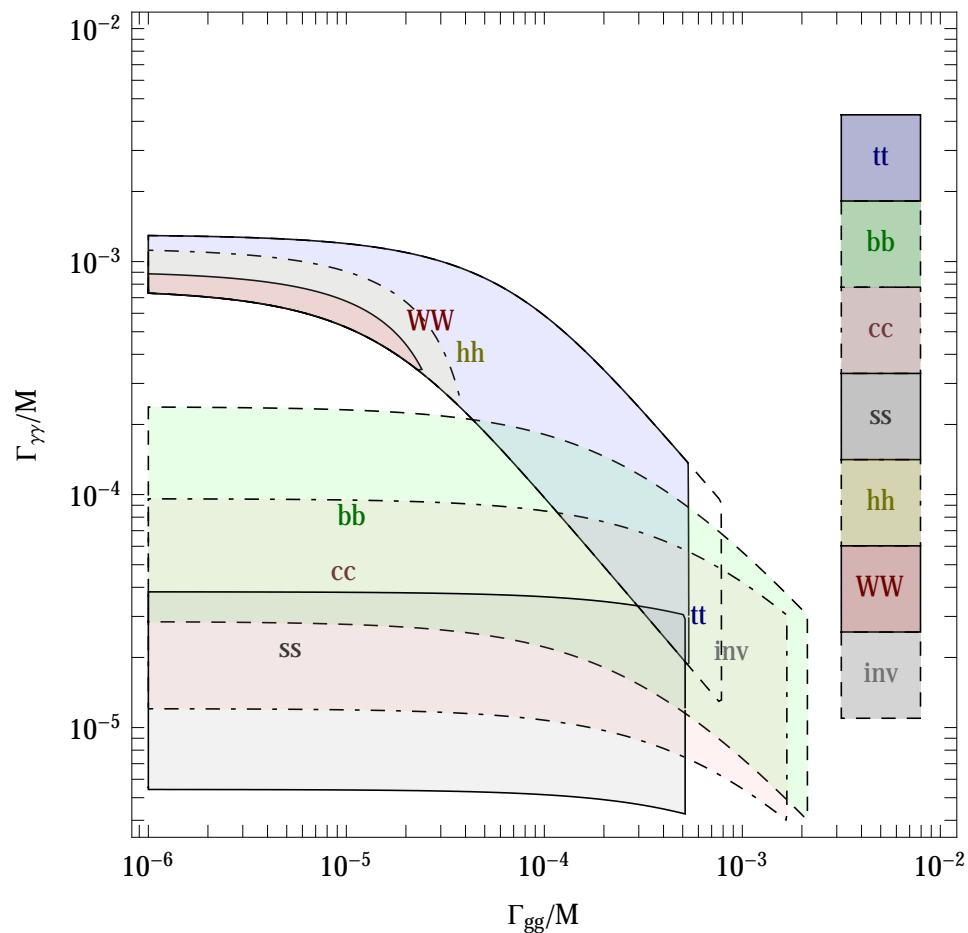
Here $r = \sigma_{13\text{ TeV}}/\sigma_{8\text{ TeV}}$. Using run 2 data only would be safer. Run 2 jj ?

Even invisible modes are constrained

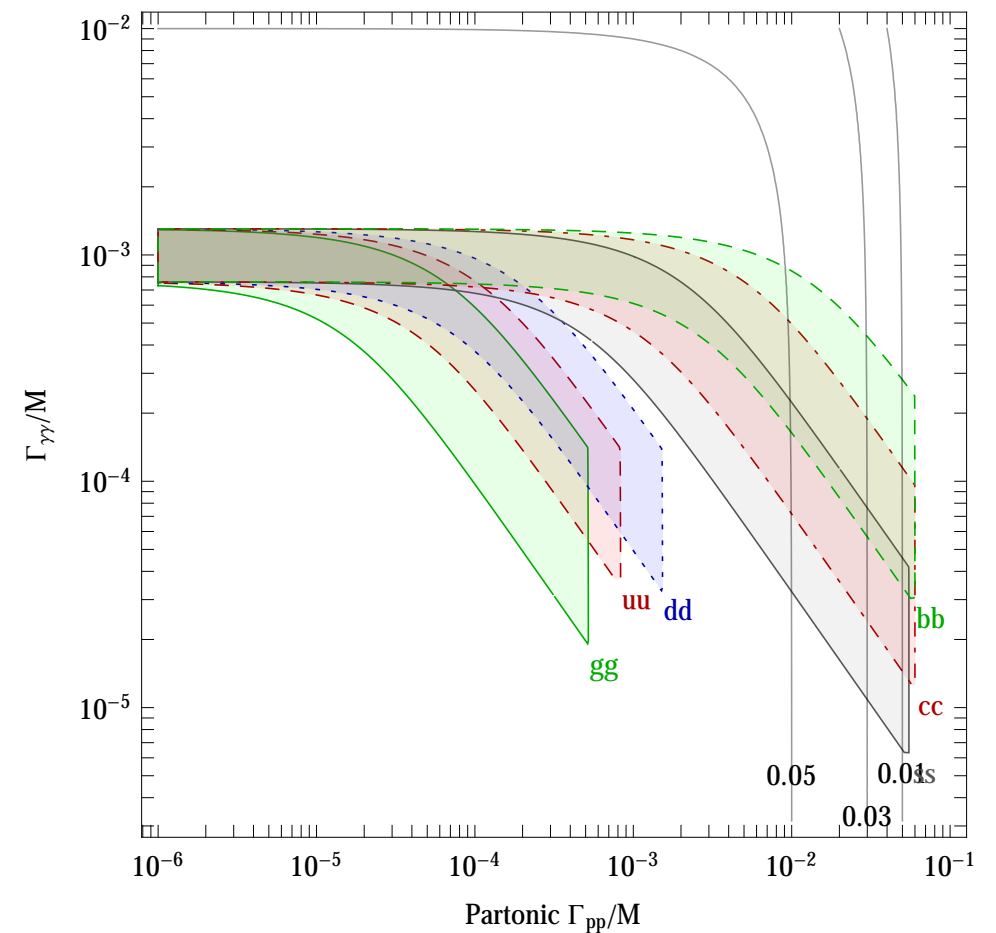
Global fits

Regions that fit the rate, the width $\Gamma/M \sim 0.06$ and that satisfy all bounds:

$S \leftrightarrow gg, \gamma\gamma, X$



$S \leftrightarrow \rho\rho, \gamma\gamma, \text{invisible}$



$\Gamma(S \rightarrow \gamma\gamma)/M \gtrsim \text{few} \times 10^{-5}$ needed: it's big!

Spin 2?

$S_{\mu\nu}$ coupled to the various components $T_{\mu\nu}^{(p)}$ of the energy-momentum tensor:

$$\mathcal{L}_{\text{eff}} = S^{\mu\nu} \sum_{p=\{\gamma,g,\ell,\dots\}} \frac{T_{\mu\nu}^{(p)}}{\Lambda_p}$$

The width into photons can be large. Angular distributions: not flat.

Randall-Sundrum graviton disfavored: it predicts a common Λ such that

$$\sigma(pp \rightarrow e^+e^- + \mu^+\mu^-) = \sigma(pp \rightarrow \gamma\gamma)$$

But no peaks seen in leptons, $\sigma(pp \rightarrow \ell^+\ell^-) < 5 \text{ fb}$ (ATLAS) and $\lesssim 3 \text{ fb}$ (CMS).

Spin 0: $SU(2)_L$ singlet or doublet?

S as a doublet: coupling to $\gamma\gamma$ and gg suppressed by $v/M \sim 0.2$.

S as a singlet: coupling to SM fermions suppressed by $v/M \sim 0.2$:

$$\mathcal{L}_{\text{eff}} = S \left[g_3^2 \frac{G_{\mu\nu}^2}{2\Lambda_g} + g_2^2 \frac{W_{\mu\nu}^2}{2\Lambda_W} + g_1^2 \frac{B_{\mu\nu}^2}{2\Lambda_B} + \left(\frac{H\bar{\psi}_L\psi_R}{\Lambda_\psi} + \text{h.c.} \right) + \frac{|D_\mu H|^2}{\Lambda_H} \right]$$

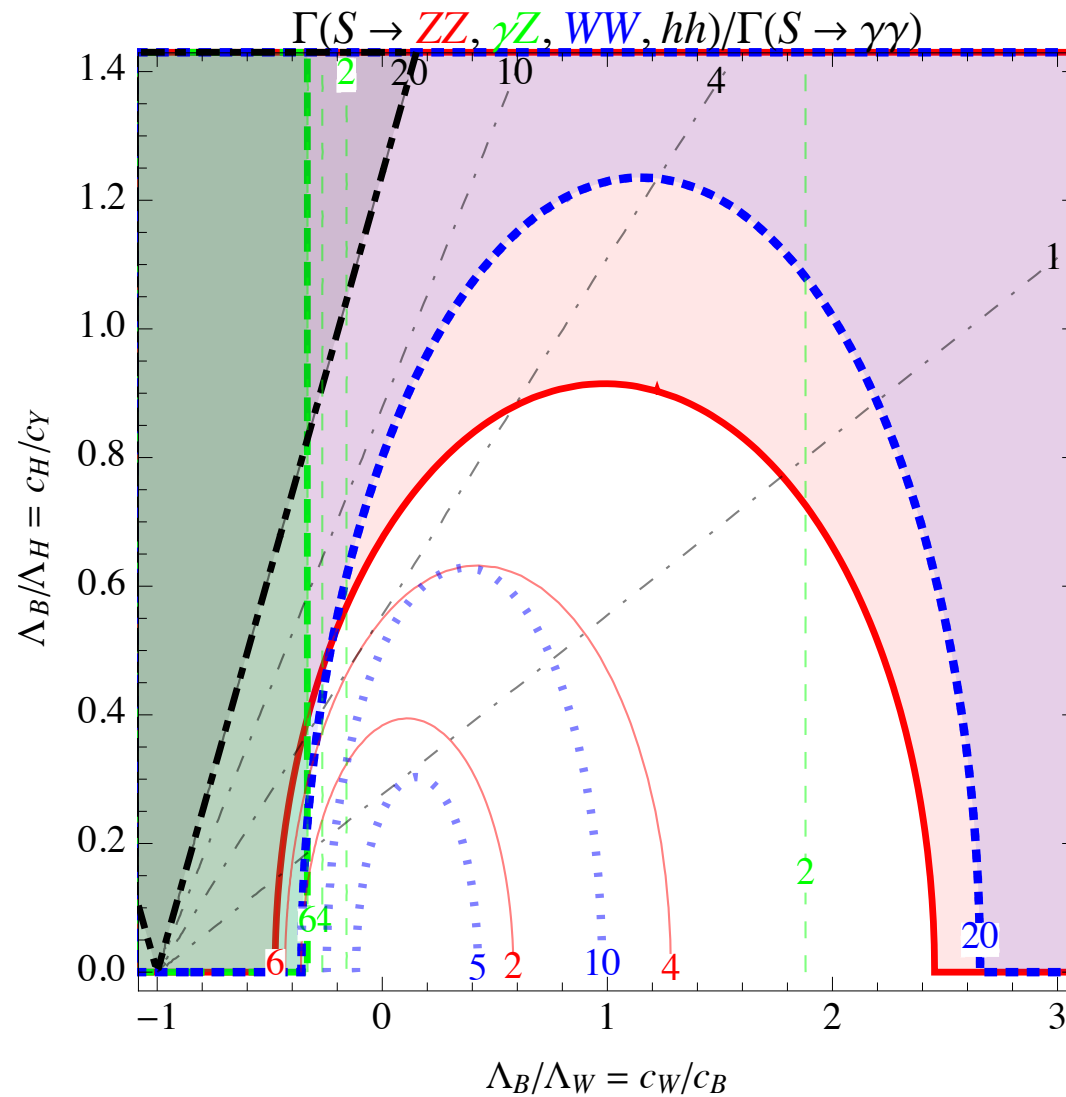
$SU(2)$ -invariance implies $S \rightarrow Z\gamma, ZZ, WW$ nearby

operator	$\frac{\Gamma(S \rightarrow Z\gamma)}{\Gamma(S \rightarrow \gamma\gamma)}$	$\frac{\Gamma(S \rightarrow ZZ)}{\Gamma(S \rightarrow \gamma\gamma)}$	$\frac{\Gamma(S \rightarrow WW)}{\Gamma(S \rightarrow \gamma\gamma)}$
WW only	$2/\tan^2 \theta_W \approx 7$	$1/\tan^4 \theta_W \approx 12$	$2/\sin^4 \theta_W \approx 40$
BB only	$2 \tan^2 \theta_W \approx 0.6$	$\tan^4 \theta_W \approx 0.08$	0

Bounds satisfied for $-0.2 < \Lambda_B/\Lambda_W < 2.5$

Future

Measuring the EW widths will over-constrain the operators

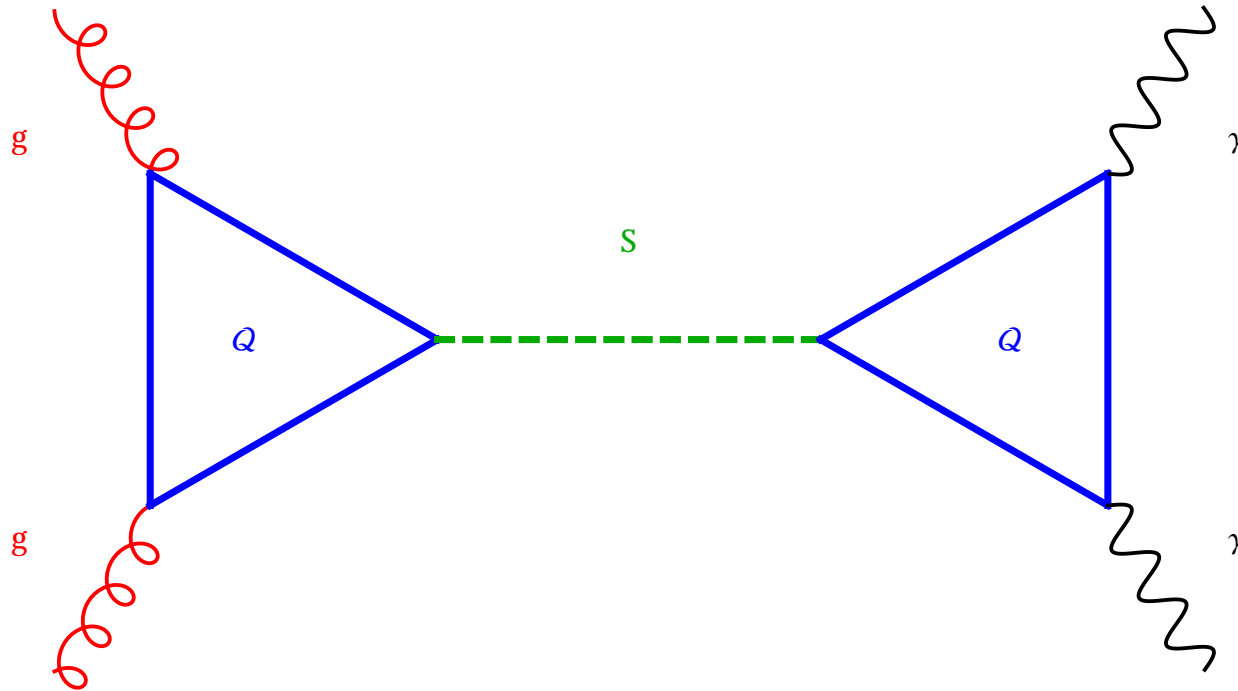


and maybe discover coupling to h

Models



Volkmodell (the everybody's model)



Extra fermions Q or scalars \tilde{Q} needed

SM loop excluded: the tree level decay would be too large e.g. $\frac{\Gamma_{t\bar{t}}}{\Gamma_{\gamma\gamma}} \approx 10^5$.

More particles needed

The Sgg and $S\gamma\gamma$ operators can be generated if S couples to charged particles

$$S\bar{Q}_f(y_f + i y_{5f}\gamma_5)Q_f + SA_s\bar{Q}_s^*Q_s$$

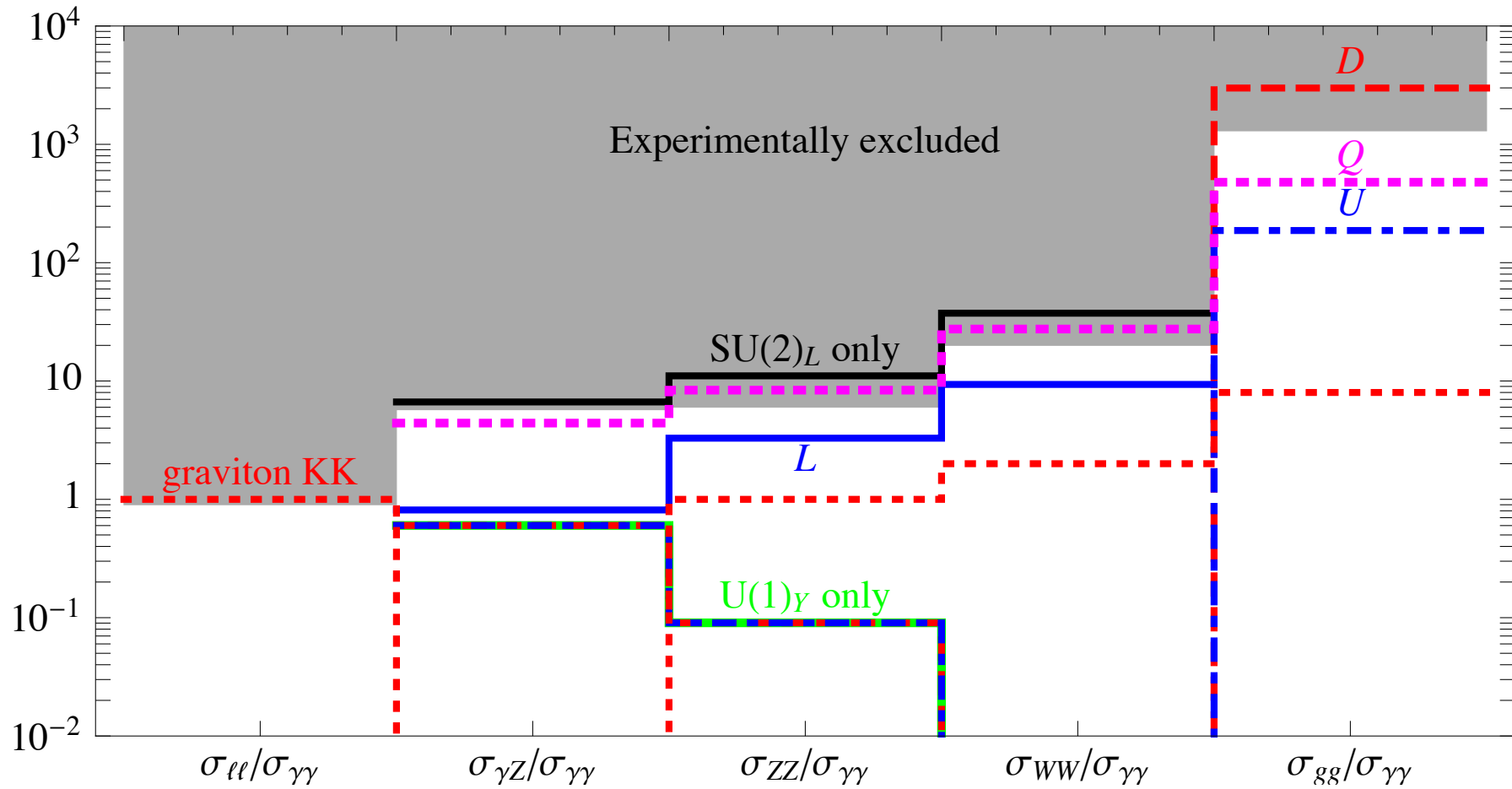
At one loop

$$\frac{\Gamma(S \rightarrow gg)}{M} \approx 7.2 \times 10^{-5} \left| \sum_f I_{r_f} y_f \frac{M}{2M_f} + \sum_s I_{r_s} \frac{A_s M}{16M_s^2} \right|^2,$$

$$\frac{\Gamma(S \rightarrow \gamma\gamma)}{M} \approx 5.4 \times 10^{-8} \left| \sum_f d_{r_f} Q_f^2 y_f \frac{M}{2M_f} + \sum_s d_{r_s} Q_s^2 \frac{A_s M}{16M_s^2} \right|^2,$$

- Such loop processes cannot make $\Gamma/M \sim 0.06$.
The large width is typical of a $1 \rightarrow 2$ tree level decay with coupling $y \sim 1$.
- If Γ is small, data want $\Gamma(S \rightarrow \gamma\gamma) \gtrsim 10^{-6}M$, which can be done. E.g. a H' , with S and P splitted by $\Delta M = \lambda v^2/M = \lambda \times 40 \text{ GeV}$ ($< 6 \text{ GeV}$ in MSSM)
- If Γ is large, data want $\Gamma(S \rightarrow \gamma\gamma) \gtrsim 10^{-4}M$, which seems too large?

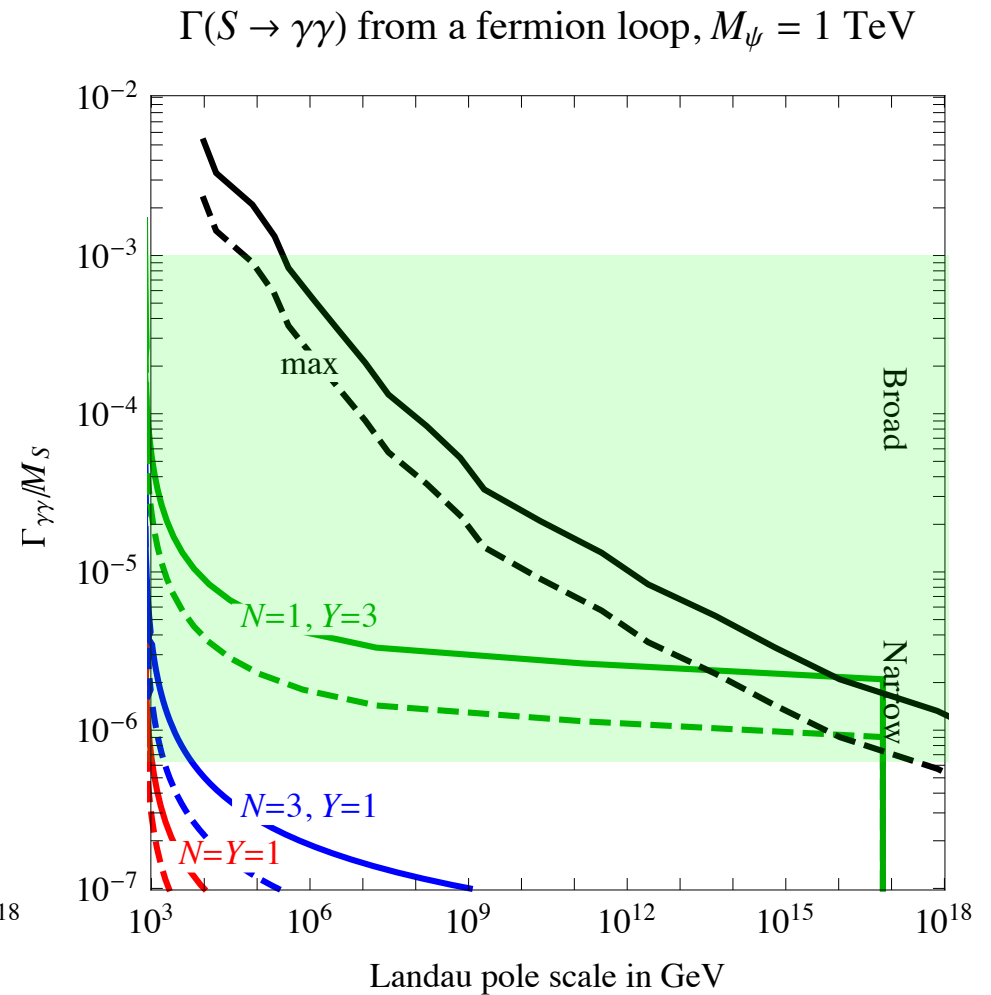
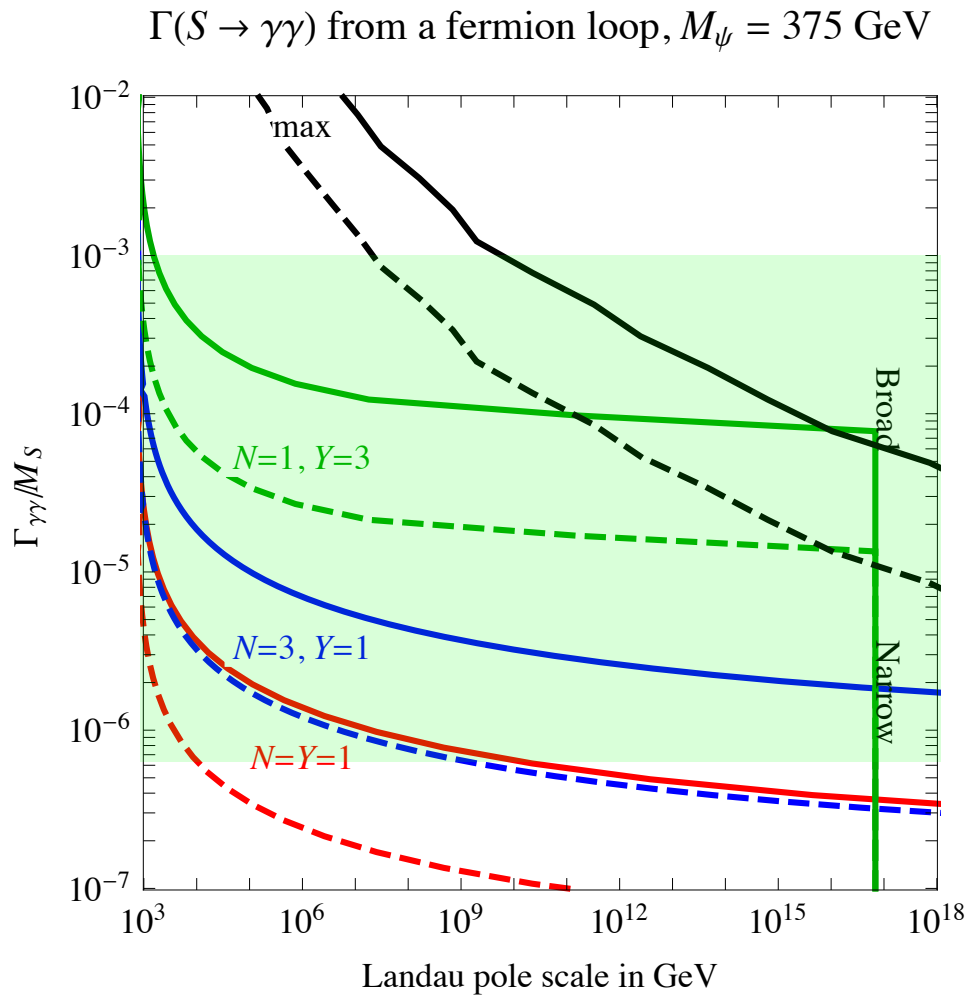
Good particles in the loop: L, E, U



Large width \Rightarrow non-perturbativity

Enhance $\Gamma(S \rightarrow \gamma\gamma)$ with: a) many states; b) big Yukawa y ; c) big charge.

In any case: nearby Landau poles for g_3 or e or y :

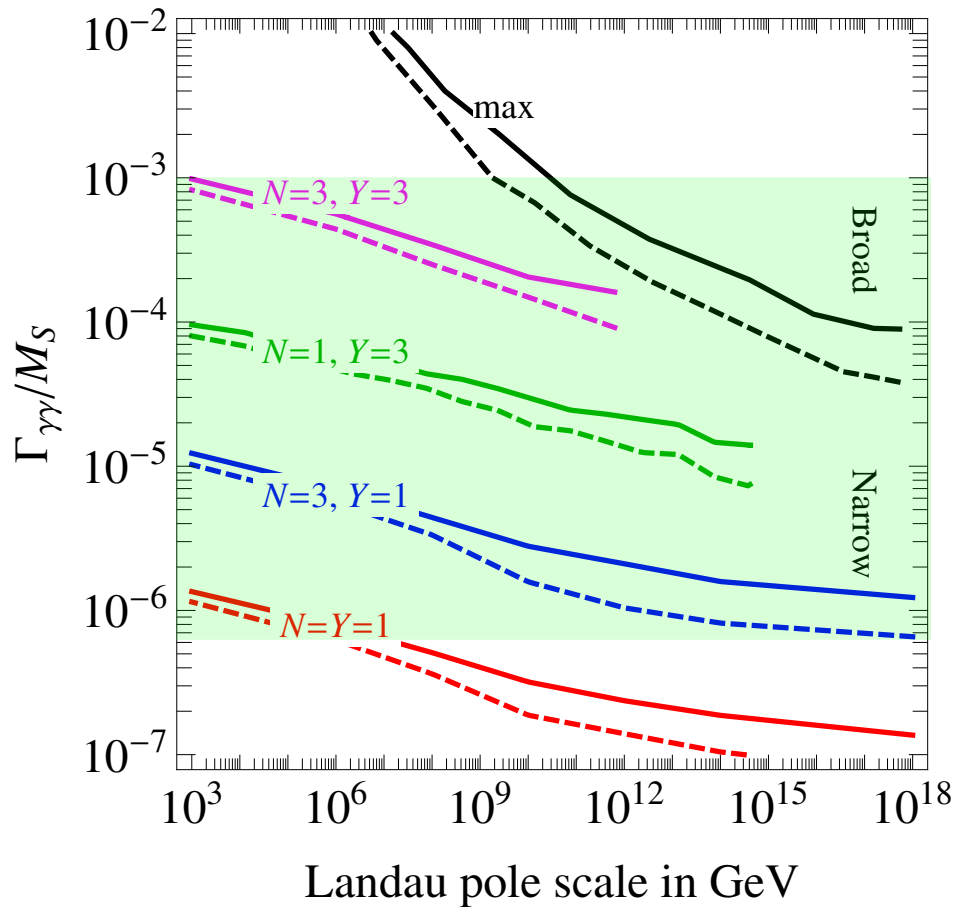


Much larger y and $\Gamma_{\gamma\gamma}$ if gauged $SU(N)$ with IR fixed point. Then $pp \rightarrow SS$.

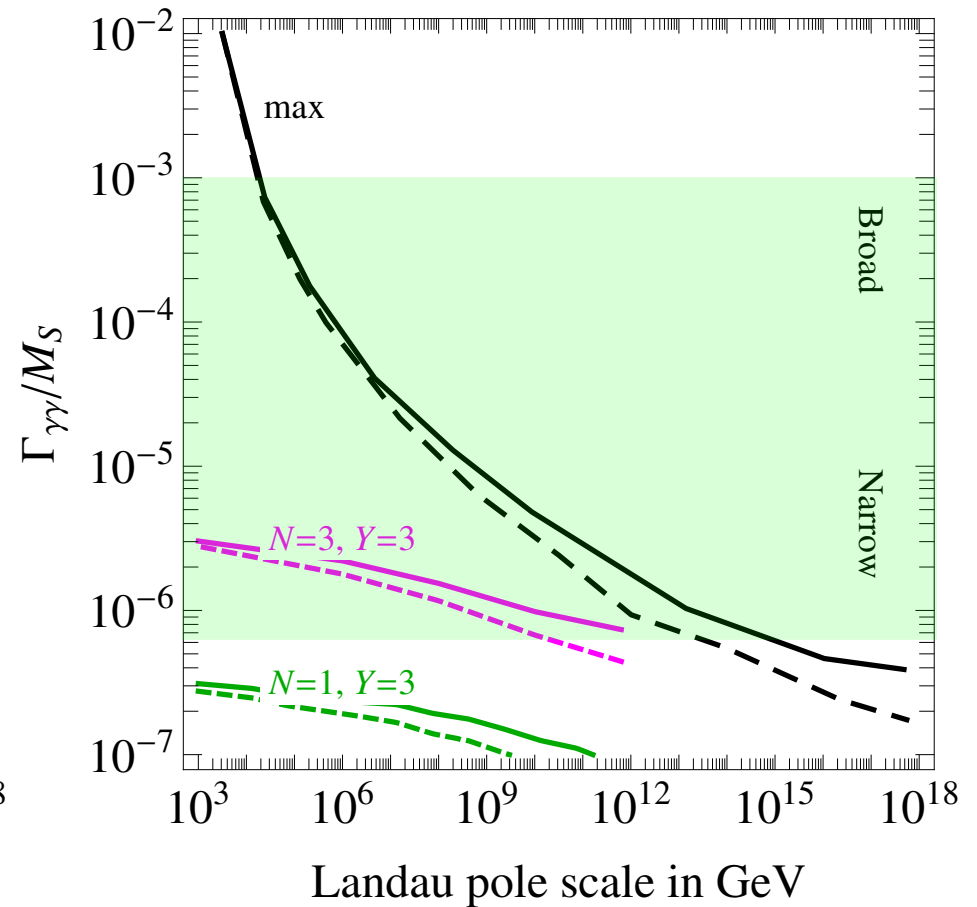
Similar results with extra scalars

A large cubic does not give Landau poles, but it is limited by vacuum decay.

$\Gamma(S \rightarrow \gamma\gamma)$ from a scalar loop, $M_X = 375$ GeV



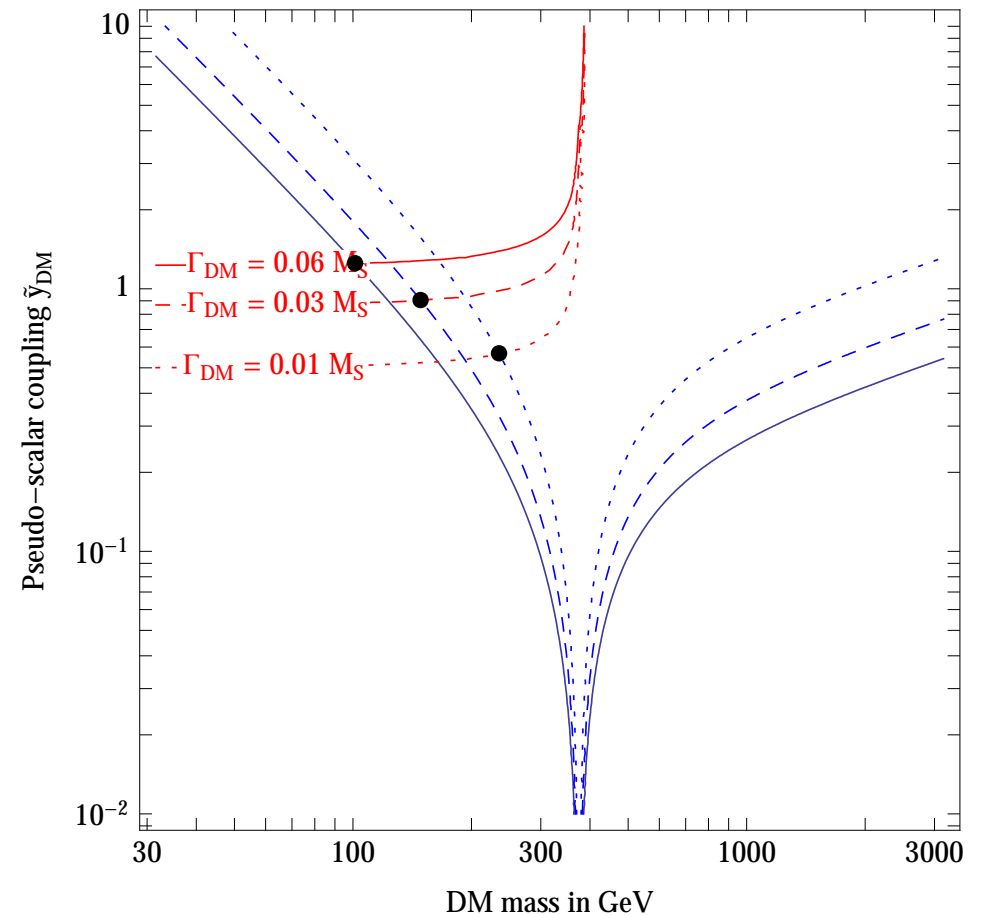
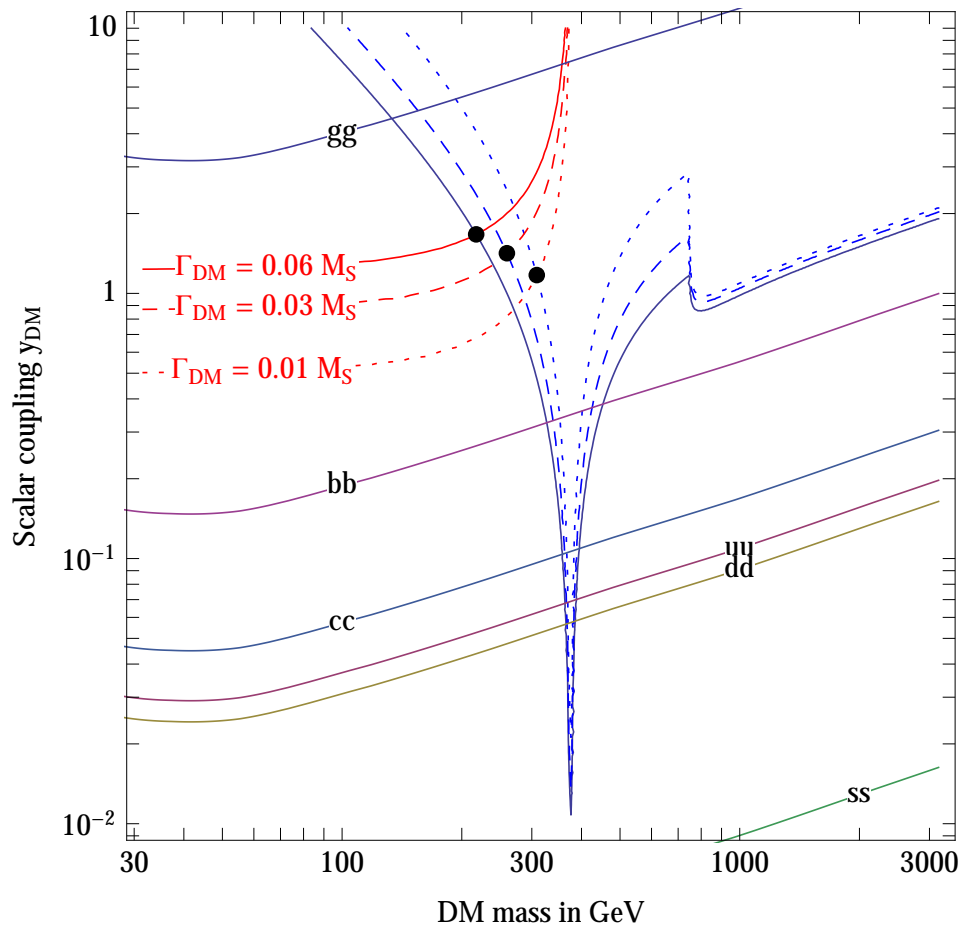
$\Gamma(S \rightarrow \gamma\gamma)$ from a scalar loop, $M_X = 1$ TeV



$\Gamma_{\gamma\gamma}$ can be much larger if gauged $SU(N)$ with IR fixed point

Extra $\mathcal{Q} = \text{Dark Matter?}$

- 1) The connection with Ω_{DM} is interesting on its own;
- 2) as a way to hide many particles, if they are needed to enhance $S \rightarrow \gamma\gamma$;
- 3) as a way to get tree level $S \rightarrow \text{DM DM}$ decays, if $\Gamma/M \sim 0.06$ is true.



Direct detection bounds are (weak) irrelevant if S is a scalar (pseudo-scalar).

Strongly coupled models

Large width natural. Main options:

$SU(2)_L$ broken by strong dynamics aka technicolor. Bonus/malus:

- + Simple UV-complete fundamental theories. E.g. extra fermions Q chiral under $SU(2)_L$ and charged under extra $SU(N)$ strong at v .
- Dead. RIP.

Composite H and S . Bonus/malus:

- Flavor problems avoided imagining partial compositeness \mathcal{L}_{eff} , no theory.
- + Allows large width trough $S \rightarrow t\bar{t}$.
- + 750 GeV compatible with usual (fine-tuned) naturalness.

Composite S , elementary H and SM. Bonus/malus:

- + No flavor problems, simple UV-complete fundamental theories. E.g. extra vectorlike particles Q charged under SM and under extra strong $SU(N)$.
- + Dark Matter could be a G -stable $TC\pi$. Large width as $S \rightarrow DM DM$?
- + 750 GeV ok with modified naturalness, dynamical generation of v .

S could be:

- 1) a pseudo-scalar $TC\eta$ with anomalous coupling to vectors enhanced by N ;
- 2) a scalar, maybe a dirty TC-dilaton;
- 3) a TC-charmonium.

Theories

Ferrari 125



Ferrari 750
MONZA SCAGLIETTI 1955

MILLE MIGLIA 2014

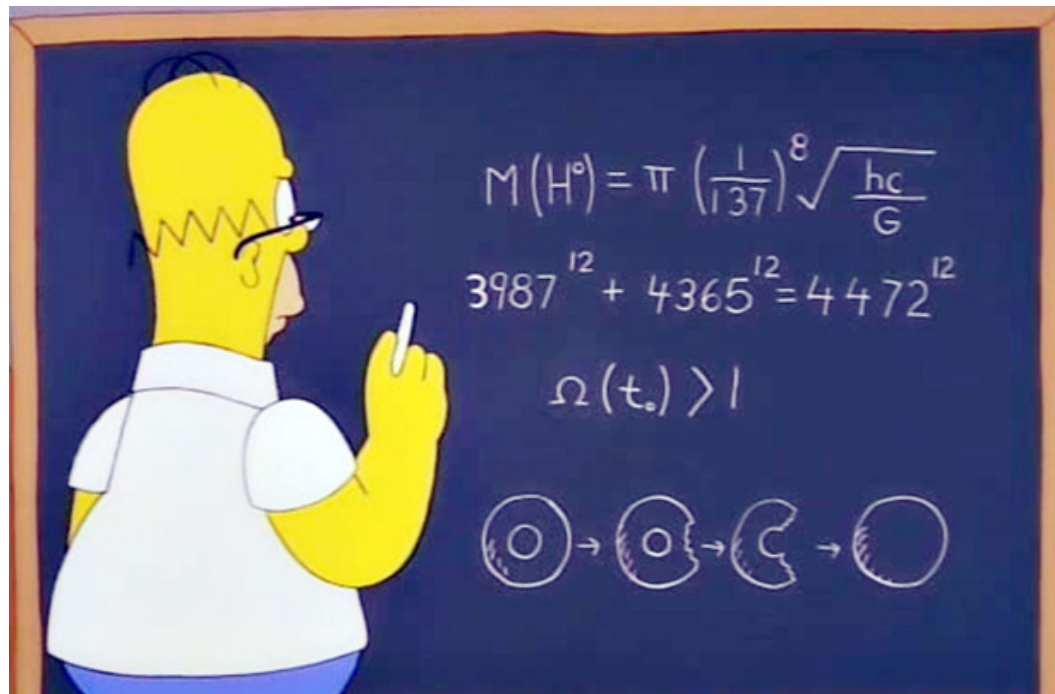


The Big Picture

Too early. If narrow, everybody who has a singlet can postdict.

E.g. S could be some SUSY singlet ($H, A, \tilde{\nu}$, NMSSM, sgoldstinos...).

Trinification $SU(3)_L \otimes SU(3)_R \otimes SU(3)_c$ predicts extra L, D ($\subset 27$ of E_6) and can fit diphoton + diboson ($g_R \approx 0.444$). String models often have extra states. Homer Simpson predicted $M = 774 \text{ GeV}$ in 1998



What next?

LHC

Sgg gives more jets than $Sq\bar{q}$. Measure the transverse momentum of S :

$$\frac{\sigma(20 \text{ GeV} < p_T^S < 40 \text{ GeV})}{\sigma(p_T^S < 20 \text{ GeV})} = \begin{cases} 1.4 & gg \\ 0.6 & q\bar{q} \\ \sim 1.1 & b\bar{b} \end{cases}$$

$Sb\bar{b}$ gives extra b jets [arXiv:1512.08478].

Scalar/pseudoscalar can be discriminated by

- 1) observing $S \rightarrow hh$,
- 2) studying $S \rightarrow ZZ \rightarrow 4\ell$,
- 3) studying Sjj .

$pp \rightarrow SS$ if big couplings. [...]

Implications for FCCee

1) WOW

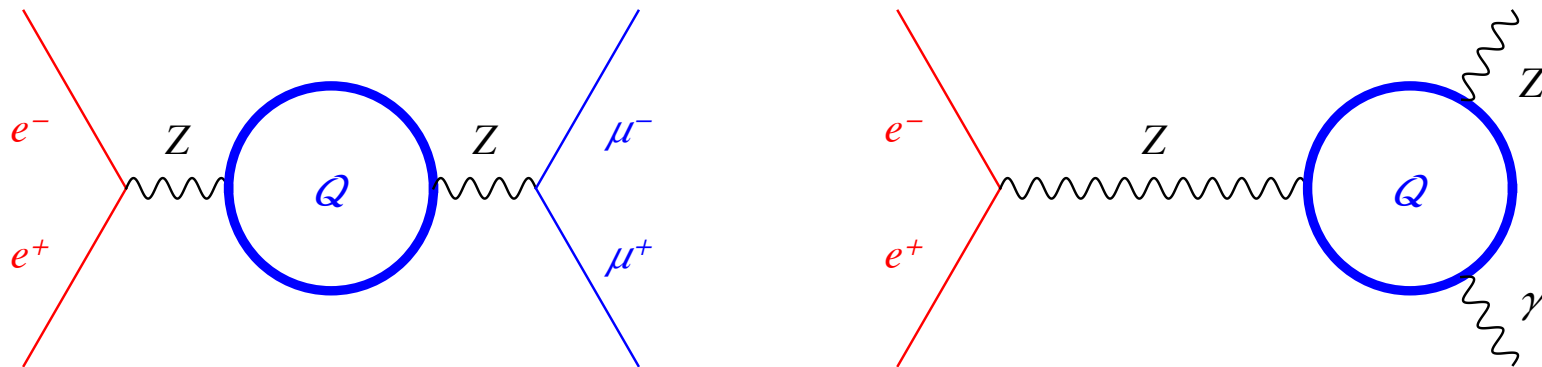
2) New physics at 750 GeV! Is FCCee dis-f**ed?

3) Maybe at $\lesssim \frac{1}{2} 750$ GeV to give invisible S decays + loop width into $\gamma\gamma$?

LHC can miss DM multiplets, especially if quasi degenerate (soft tag). Then:

A) High-energy tails of $pp \rightarrow \ell^- \ell^+$, sensitive to Δb (BSM running of g_Y, g_2).

B) e^+e^- collider: even if Q is too heavy, it could be probed indirectly



Universal new physics

Heavy new physics that affects the SM vector propagators is condensed in:

S, T: affected by new particles coupled to *H*, FCC sensitivity already studied

W, Y: affected by any new particle with gauge interactions

$W, Y \geq 0$ are cumulative and related to RGE coefficients by

$$Y = \frac{\alpha_Y}{20\pi} \left[\sum_s \frac{\Delta b_Y^{(s)} M_W^2}{2 M_s^2} + \sum_f \Delta b_Y^{(f)} \frac{M_W^2}{M_f^2} \right] = 1.1 \cdot 10^{-5} \left[\frac{\Delta b_Y^{(s)}}{2} + \Delta b_Y^{(f)} \right] \left(\frac{300 \text{ GeV}}{M} \right)^2$$

and the same for $Y \rightarrow W$. For example heavy SUSY particles give:

$$Y = \frac{\alpha_Y}{40\pi} M_W^2 \left(\frac{1}{m_{\tilde{E}}^2} + \frac{1}{2m_{\tilde{L}}^2} + \frac{1}{3m_{\tilde{D}}^2} + \frac{4}{3m_{\tilde{U}}^2} + \frac{1}{6m_{\tilde{Q}}^2} + \frac{1}{6m_A^2} \right) + \frac{\alpha_Y}{30\pi} \frac{M_W^2}{\mu^2} > 0$$

$$W = \frac{\alpha_2}{80\pi} M_W^2 \left(\frac{1}{m_{\tilde{L}}^2} + \frac{3}{m_{\tilde{Q}}^2} + \frac{1}{3m_A^2} \right) + \frac{\alpha_2}{30\pi} \left(\frac{M_W^2}{\mu^2} + \frac{2M_W^2}{M_2^2} \right) > 0$$

FCC sensitivity to W, Y

Mostly through the weak angle from Z couplings, $g_V/g_A \equiv 1 - 4 \sin^2 \theta_W$:

$$\delta \sin^2 \theta_W = -0.33Y - 0.09W = \pm 10^{-6}|_{\text{exp}} \pm 1.8 \cdot 10^{-5}|_{\alpha_{\text{em}}}$$

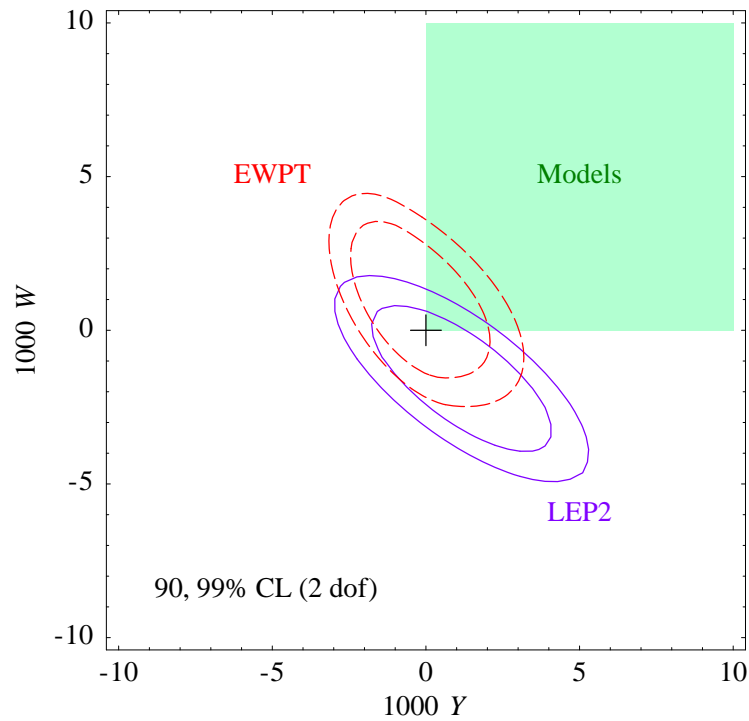
The present α_{em} error might be irreducible. At least Y to $\pm \text{few} \times 10^{-5}$.

Band becomes a long ellipse assuming standard estimates for other observables.

$W < 0.6 \cdot 10^{-4}$ means $m_{\tilde{L}} > 120 \text{ GeV}$, $\mu > 200 \text{ GeV}$, $M_2 > 270 \text{ GeV}$.

LEP1 + LEP2

LEP2 measured W, Y as well as LEP1. Indeed their effects grow with energy. At $s \gg M_W^2$



	obs/obs _{SM} - 1
$\sigma(e^+e^- \rightarrow \mu^+\mu^-)$	$-(0.67W + 1.33Y)s/M_W^2$
$\sigma(e^+e^- \rightarrow q\bar{q})$	$-(1.35W + 0.65Y)s/M_W^2$
A_{FB}^μ	$-0.40(W + Y)s/M_W^2$

Can FCCee do the same?
Expected accuracy:

\sqrt{s}/GeV	160	240	350
$\sigma(e^+e^- \rightarrow \mu^+\mu^-)$	0.015%	0.025%	0.075%
W, Y	$0.4 \cdot 10^{-4}$	$0.3 \cdot 10^{-4}$	$0.4 \cdot 10^{-4}$

Conclusions

We will know much more, soon.