## Overview: <br> from (many) qubits to space-time

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c.f. talks by Latorre, Molina-Vilaplana, Pastawski, Wen, Maldacena

## Motivation

Condensed Matter
Quantum Information

(Objects that look like) space-times seem to emerge from the entanglement structure of quantum many-body states
(and we were not thinking about gravity at all...)

This talk: overview of some ideas along these lines

## Outline

1) Review of TNs
2) PEPS and emergent Hamiltonians
3) Symmetric TNs and emergent spin networks
4) MERA and emergent AdS/CFT
5) Summary \& open questions

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## Entanglement obeys area-law

## Entanglement

## key resource in quantum information

teleportation, quantum algorithms, quantum error correction, quantum cryptography...

## Entanglement

## 2d system


key resource in quantum information teleportation, quantum algorithms, quantum error correction, quantum cryptography...

$$
\rho_{A}=\operatorname{tr}_{E}(|\Psi\rangle\langle\Psi|)
$$

$S(A)=-\operatorname{tr}\left(\rho_{A} \log \rho_{A}\right)$

Reduced density matrix of subsystem A

Entanglement entropy (von Neumann entropy)

For many ground states


$$
(L>\xi)
$$

## Entanglement

## 2d system

$$
\begin{aligned}
\rho_{A} & =\operatorname{tr}_{E}(|\Psi\rangle\langle\Psi|) \\
S(A) & =-\operatorname{tr}\left(\rho_{A} \log \rho_{A}\right)
\end{aligned}
$$



Reduced density matrix of subsystem A

For many ground states

## In d dimensions

key resource in quantum information teleportation, quantum algorithms, quantum error correction, quantum cryptography...

## Entanglement entropy (von Neumann entropy)



$$
(L>\xi)
$$

Generic $\quad S(A) \sim L^{d}$ state (volume)
Ground states of (most) local Hamiltonians
$S(A) \sim L^{d-1}$

Srednicki, Plenio, Eisert, Dreißig, Cramer, Wolf..

## Many-body Hilbert space is far too large

## Hilbert space is a convenient illusion

Hilbert space of a N -body many-body system

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Hilbert space of a N -body many-body system

Set of area-law states
Y. Ge, J. Eisert, arXiv:1411.2995

Set of TN states (low-energy eigenstates of local Hamiltonians)

Set of product states (mean field)

## Hilbert space is a convenient illusion

Hilbert space of a N -body many-body system
"Exploration" time $\sim O\left(10^{10^{23}}\right)$ sec.
Compare to...
Most states here are not Age of the universe $\sim O\left(10^{17}\right)$ sec. even reachable by a time evolution with a local Hamiltonian in polynomial time

Poulin, Qarry, Somma, Verstraete, PRL 106170501 (2011)

Set of area-law states
Y. Ge, J. Eisert, arXiv:1411.2995

Set of TN states (low-energy eigenstates of local Hamiltonians)

Set of product states (mean field)

We need a language to target the relevant corner of quantum states directly

Tensor Networks

## A new language

$$
\left.|\Psi\rangle=\sum_{i, s} \Psi_{i, i_{2}, i_{\nu}}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle
$$

## A new language

$$
\left.|\Psi\rangle=\sum_{i^{\prime} s} \Psi_{i i_{1}, \ldots, i_{V}}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle
$$



## A new language

$$
\left.|\Psi\rangle=\sum_{i^{\prime} s} \Psi_{i i_{1}, \ldots, i_{\nu}}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle
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$$



## Tensor network diagrams

vector $\vec{v}$
matrix $A$

matrix product $A B$

trace of matrix product $\operatorname{tr}(A B C)$
tensor contraction $\quad f(A, B, C, D)$


# Tensor Networks <br> $A \cdot B \rightarrow$<div class="inline-tabular"><table id="tabular" data-type="subtable">
<tbody>
<tr style="border-top: none !important; border-bottom: none !important;">
<td style="text-align: right; border-left: none !important; border-right: none !important; border-bottom-style: solid !important; border-bottom-width: 1px !important; border-top: none !important; width: auto; vertical-align: middle; ">$A$</td>
<td style="text-align: right; border-bottom-style: solid !important; border-bottom-width: 1px !important; border-top: none !important; width: auto; vertical-align: middle; ">$B$</td>
</tr>
<tr style="border-top: none !important; border-bottom: none !important;">
<td style="text-align: right; border-left: none !important; border-right: none !important; border-bottom-style: solid !important; border-bottom-width: 1px !important; border-top: none !important; width: auto; vertical-align: middle; " class="_empty"></td>
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</tr>
</tbody>
</table>
<table-markdown style="display: none">| $A$ | $B$ |
| ---: | ---: |
|  | - |</table-markdown></div> <br> e.g. RO, Annals of Physics 349 (2014) 117-158 <br> $$
|\Psi\rangle=\sum_{i^{\prime} s} \Psi_{i_{1} i_{2} \ldots i_{N}}\left|i_{1}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle \quad \begin{aligned} & \text { p-level } \\ & \text { systems } \end{aligned}
$$ 

## Tensor Networks

## $A \cdot B \rightarrow \stackrel{A}{-}-\stackrel{B}{0}$

e.g. RO, Annals of Physics 349 (2014) 117-158

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|\Psi\rangle=\sum_{i^{\prime} s} \Psi_{i_{1} i_{2} \ldots i_{N}}\left|i_{1}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle \begin{aligned}
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\text { p-level } \\
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\end{array}
\end{aligned}
$$

# Tensor Networks 

e.g. RO, Annals of Physics 349 (2014) 117-158


Matrix Product States (MPS)
 DMRG, PWFRG, TEBD...

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## Tensor Networks

e.g. RO, Annals of Physics 349 (2014) 117-158

\section*{$A \cdot B \rightarrow$| $A$ | $B$ |
| ---: | ---: |
|  | - |}


$|\Psi\rangle=\sum_{i^{\prime} s} \Psi_{i_{1} i_{2} \ldots i_{N}}\left|i_{1}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle$| p-level |
| :--- |
| systems |



Matrix Product States (MPS)

DMRG, PWFRG, TEBD...
physical 1...p bond 1..D (entanglement)

Projected Entangled Pair States (PEPS), Tensor Product States (TPS)


Tensor Product Variational Approach, PEPS \& iPEPS algorithms, Tensor-Entanglement Renormalization, TRG/SRG/HOTRG/HOSRG...

## Tensor Networks

e.g. RO, Annals of Physics 349 (2014) 117-158

$\downarrow$ Scale-invariant
Multiscale Entanglement Renormalization Ansatz (MERA)


AdS/CFT, Entanglement Renormalization

$$
|\Psi\rangle=\sum_{i^{\prime} s} \Psi_{i_{1} i_{2} \ldots i_{N}}\left|i_{1}\right\rangle \otimes\left|i_{2}\right\rangle \otimes \cdots \otimes\left|i_{N}\right\rangle \begin{aligned}
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physical 1...p bond 1..D (entanglement)

## Projected Entangled Pair States (PEPS), Tensor Product States (TPS)



Tensor Product Variational Approach, PEPS \& iPEPS algorithms, Tensor-Entanglement Renormalization, TRG/SRG/HOTRG/HOSRG...

Efficient $\mathrm{O}($ poly $(\mathrm{N})$ ), satisfy area-law, low-energy eigenstates of local Hamiltonians

## Comparison

| MPS in 1d |  |  |  |
| :---: | :---: | :---: | :---: |
| Ent. entropy | $S(L)=O(1)$ | $S(L)=O(L)$ | $S(L)=O(\log L)$ |
| contraction | efficient | inefficient | efficient |
| Corr. length | finite | finite \& infinite | finite \& infinite |
| To/from | 1d Ham. | 2d Ham. | 1d Ham. |
| Tensors | arbitrary | arbitrary | constrained |



Exact in many cases
Variational ansatz for numerical simulations (e.g. DMRG)

## Exact example 1: Kitaev’s Toric Code

$A_{s}=\prod_{i \in s} \sigma_{i}^{*} \quad$ star operator
$B_{p}=\prod_{i \in p} \sigma_{i}^{z} \quad$ plaquette operator


Simplest known model with "topological order"
Ground state (and in fact all eigenstates) are PEPS with $D=2$


And another tensor rotated $90^{\circ}$


## Exact example 2: Kitaev’s honeycomb model jg|u <br> Bogoliubov modes Bogoliubov transformation $\{$ Dirac momentum modes <br> Dirac real-space modes <br> Majorana braidings + Jordan-Wigner <br> P. Schmoll, RO, arXiv:1605.04315 <br> $$
\begin{aligned} H= & -J_{x} \sum_{x-\text { links }} \sigma_{j}^{x} \sigma_{k}^{x}-J_{y} \sum_{y-\text { links }} \sigma_{j}^{y} \sigma_{k}^{y} \\ & -J_{z} \sum_{z-\text { links }} \sigma_{j}^{z} \sigma_{k}^{z} \end{aligned}
$$ <br>  <br> Abelian and non-abelian, chiral and non-chiral topological phases <br> Vortex modes <br> Spins on the honeycomb

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## Projected Entangled Pair States (PEPS)



2d systems


3d systems

## PEPS obey 2d area-law







$$
\begin{aligned}
& \left.\left.\rho_{\text {in }}=\operatorname{tr}_{\text {out }}(|\Psi\rangle\langle\Psi|)=\sum_{\bar{\alpha}, \overline{\alpha^{\prime}}} X_{\bar{\alpha}, \bar{\alpha}^{\prime}} \mid \text { in }(\bar{\alpha})\right\rangle\left\langle\text { in }\left(\bar{\alpha}^{\prime}\right)\right| \quad X_{\bar{\alpha}, \bar{\alpha}^{\prime}}=\left\langle\text { out }\left(\bar{\alpha}^{\prime}\right)\right| \text { out }(\bar{\alpha})\right\rangle \\
& \operatorname{rank}\left(\rho_{\text {in }}\right) \leq D^{4 L} \quad S(L)=-\operatorname{tr}\left(\rho_{\text {in }} \log \rho_{\text {in }}\right) \leq \log (D) 4 L
\end{aligned}
$$



$$
\begin{gathered}
\left.\rho_{\text {in }}=\operatorname{tr}_{\text {out }}(|\Psi\rangle\langle\Psi|)=\sum_{\bar{\alpha}, \bar{\alpha}^{\prime}} X_{\bar{\alpha}, \bar{\alpha}^{\prime}}|\operatorname{in}(\bar{\alpha})\rangle\left\langle\text { in }\left(\bar{\alpha}^{\prime}\right)\right| \quad X_{\bar{\alpha}, \bar{\alpha}^{\prime}}=\left\langle\text { out }\left(\bar{\alpha}^{\prime}\right)\right| \text { out }(\bar{\alpha})\right\rangle \\
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\text { prefactor size of the boundary }
\end{gathered}
$$

## PEPS \& Entanglement Hamiltonians <br> e.g. I. Cirac et al, PRB 83, 245134 (2011), N. Schuch et al, PRL 111, 090501 (2013)




$\langle\Psi \mid \Psi\rangle$




Boundary
How is physics described here?

Boundary
How is physics described here?

1-dim transfer matrix: dominant eigenvector?


Can be approximated using infinite MPS

iTEBD, iDMRG, PWFRG, etc

## Emergent Hamiltonians

$\cdots p-p-p-p-\bar{p}$

Remember it has
double indices...

## Emergent Hamiltonians

Virtual indices of bra Boundary virtual index $1 \ldots \chi$
$\begin{aligned} & 1 \ldots D \\ & \text { Virtual indices of ket } \\ & 1 \ldots D\end{aligned}$
It is also hermitian and
(up to finite- $\chi$ effects)

## Emergent Hamiltonians

Virtual indices of bra Boundary virtual index 1... $\chi$
1...D

Virtual indices of ket 1...D
 It is also hermitian and positive by construction (up to finite- $\chi$ effects)

1d Entanglement Hamiltonian
$\rho=\exp \left(-\hat{H}_{E}^{\prime}\right) \quad$ moss $H_{s}$ m

## Emergent Hamiltonians

Virtual indices of bra
1...D

Virtual indices of ket 1...D

Boundary virtual index 1... $\chi$
It is also hermitian and positive by construction (up to finite- $\chi$ effects)

1d Entanglement Hamiltonian


Gapped 2d systems, trivial phase
Critical 2d systems
Gapped 2d systems, topological order
Chiral topological order, gapless
RO, M. Mambrini, D. Poilblanc, work in progress

1d Hamiltonian, short-range 1d Hamiltonian, long-range

Completely non-local (projector)
(1+1)d Conformal field theory

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## Symmetric tensors and Schur's lemma

e.g., S. Singh, R. N. C. Pfeifer, G. Vidal, PRA 82, 050301 (2010)


## Symmetric tensors and Schur's lemma <br> e.g., S. Singh, R. N. C. Pfeifer, G. Vidal, PRA 82, 050301 (2010)



Structural part depends only on the group properties (intertwiners)

## Emergent spin networks

e.g., S. Singh, R. N. C. Pfeifer, G. Vidal, PRA 82, 050301 (2010)

Symmetric TN


Spin network


## Emergent spin networks

e.g., S. Singh, R. N. C. Pfeifer, G. Vidal, PRA 82, 050301 (2010)

Symmetric TN


Coefficient


Spin network


States of quantum geometry in loop quantum gravity...

## Emergent spin networks

e.g., S. Singh, R. N. C. Pfeifer, G. Vidal, PRA 82, 050301 (2010)


Global and gauge symmetries are handled naturally

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## Multiscale Entanglement Renormalization Ansatz (MERA)




2d systems

## 1d MERA



## 1d MERA


spatial dimension

## Tensors obey constraints





## Reason:

## entanglement is built locally at all length scales

L


日月


月
$\sharp$ entangle locally

L/2


L/2






Extra dimension defines an RG flow: Entanglement Renormalization


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## Entropy of 1d MERA






Entanglement as boundary in holographic geometry: $S(L) \leq \log (\chi)\left|\partial \Omega_{L}\right|$


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Constant contribution at every layer 1d MERA can produce logarithmic violations to the area-law: $S(L) \approx \log L$
(like 1d critical systems!)

## MERA \& AdS/CFT

e.g. B. Swingle, PRD 86, 065007 (2012), G. Evenbly, G. Vidal, JSTAT 145:891-918 (2011)

## Emergent space-time

## MERA



$$
S_{A} \propto \operatorname{Min}\left[\# \operatorname{Bonds}\left(\gamma_{A}\right)\right] \quad S_{A} \propto \operatorname{Min}[\text { Area }]
$$

AdS/CFT



Picture from M. Nozaki, S. Ryu, T. Takayanagi, JHEP10(2012)193

MERA entropy ~ Ryu-Takayanagi prescription

Picture from G. Evenbly, G. Vidal, (2011) JSTAT 145:891-918

Picture from G. Evenbly, G. Vidal, (2011) JSTAT 145:891-918

Bulk is a discretized AdS space

Picture from G. Evenbly, G. Vidal, (2011) JSTAT 145:891-918
(and we were not thinking about gravity at all...)


## $\mathrm{CFT}_{1+1}$

For a scale-invariant MERA, the tensors of a critical model with a CFT limit correspond to a "gravitational" description in a discretized AdS space: „lattice" realization of AdS/CFT correspondence

Let's now play some jazz...


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Finite correlation length (gapped systems) $=$ finite number of layers


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Product state = trivial fixed point

(time slice)
If arbitrary, then we can have nontrivial thermal states.

If isommetry, then all information is encoded in the network of correlations and

$$
\rho_{i n}=I
$$

Finite correlation length (gapped systems) $=$ finite number of layers

$$
\left.\begin{array}{l}
\rho_{\text {in }}=\operatorname{tr}_{\text {out }}(|\Psi\rangle\langle\Psi|) \\
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\end{array}\right\} \quad \begin{gathered}
\text { Same thermal spectrum (entanglement Hamiltonian) } \\
\text { finite temperature, scale invariance broken }
\end{gathered}
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Thermofield double state
Eternal AdS black-hole
$|T F D\rangle=\frac{1}{\sqrt{Z(\beta)}} \sum_{n} e^{-\beta E_{n} / 2}|n\rangle_{1}|n\rangle_{2}$

MPO
wormhole

CFT1

MERA

cMERA
(continuum)

$$
|\psi\rangle=P e^{-i \int_{u 2}^{u 1}(K(u)+L) d u}|\Omega\rangle
$$

J. Haegeman et al, Phys. Rev. Lett. 110, 100402 (2013)
$K(u)$ Disentangler generator
$L$ Isommetry generator

## cMERA

(continuum)

$$
-i \int^{u 1}(K(u)+L) d u
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$|\psi\rangle=P e$
J. Haegeman et al,

Phys. Rev. Lett. 110, 100402 (2013)
$K(u)$ Disentangler generator
$L$ Isommetry generator

$$
\left.g_{u u}(u) d u^{2}=\left.\mathcal{N}^{-1}\left(1-\left|\langle\Psi(u)| e^{i L \cdot d u}\right| \Psi(u+d u)\right\rangle\right|^{2}\right)
$$

Measures the density of strength of disentanglers.
Compatible with AdS metric
curvature ~ change of entanglement at every length scale
M. Nozaki, S. Ryu, T. Takayanagi, JHEP10(2012)193

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Many-body entanglement (QM, non-relativistic)

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Entanglement spectrum (energy = eigenvalues)

Quasiparticles (eigenvectors)

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Mass, statistics...

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Many-body entanglement (QM, non-relativistic)


## Other developments

- Exact holographic mapping
X.-Liang Qi, arXiv:1309.6282
- AdS/CFT as Quantum Error-Correcting Code A. Almheiri, X. Dong, D. Harlow, JHEP 1504:163 (2015)
- Holographic Quantum Error Correcting Codes
F. Pastawski, B. Yoshida, D. Harlow, J. Preskill, JHEP 06149 (2015)
J. I. Latorre, G. Sierra, arXiv:1502.06618
- Einstein's equations from Entanglement Entropy T. Faulkner, M. Guica, T. Hartman, R. C. Myers, M van Raamsdonk, JHEP 03051 (2013); B. Swingle, M. van Raamsdonk, arXiv:1405.2933


## Some cross-over open questions

- cMERA, wavelets, and AdS/CFT?
- Superpositions of TTNs? Linear optics?
C. Papadopoulos, RO work in progress
- Consistent AdS/TN? What about other correspondences?
- TN structure of, e.g., $\mathrm{N}=4 \mathrm{SYM}$ ? (Type-IIB on $\mathrm{AdS}_{5} \times \mathrm{S}^{5}$ ) Can one derive string theory from entanglement?
J. Molina-Vilaplana
- Non-classical gravity from „exotic" TNs? (topological order \& D-branes, TNs with symmetries...)
- Holographic multipartite entanglement? Holographic mixed-state entanglement? D. Pang, RO, work in progress
- "Gravitational" interpretation of branching MERA?
- Numerical simulations of gravity with TN methods?
- „entanglement renormalization" $\longleftrightarrow$ „,holographic renormalization"?
- Lorentz invariance from Lieb-Robinson bounds?


