

# Conformal field theory and energy-momentum tensor on the lattice

Etsuko Ito (KEK)

## Conformal theory:

E.I. , PTEP (2013)8, 083B01

E.I. , PTEP (2015)4, 043B08

E.I. and A.Tomiya, PoS LATTICE2014 (2014) 252

## Energy-momentum tensor:

Asakawa, Hatsuda, E.I., Kitazawa, Suzuki, Phys.Rev. D90 (2014) 1, 011501

E.I.,Suzuki, Taniguchi, Umeda PoS(LATTICE 2015)303

## Entanglement entropy:

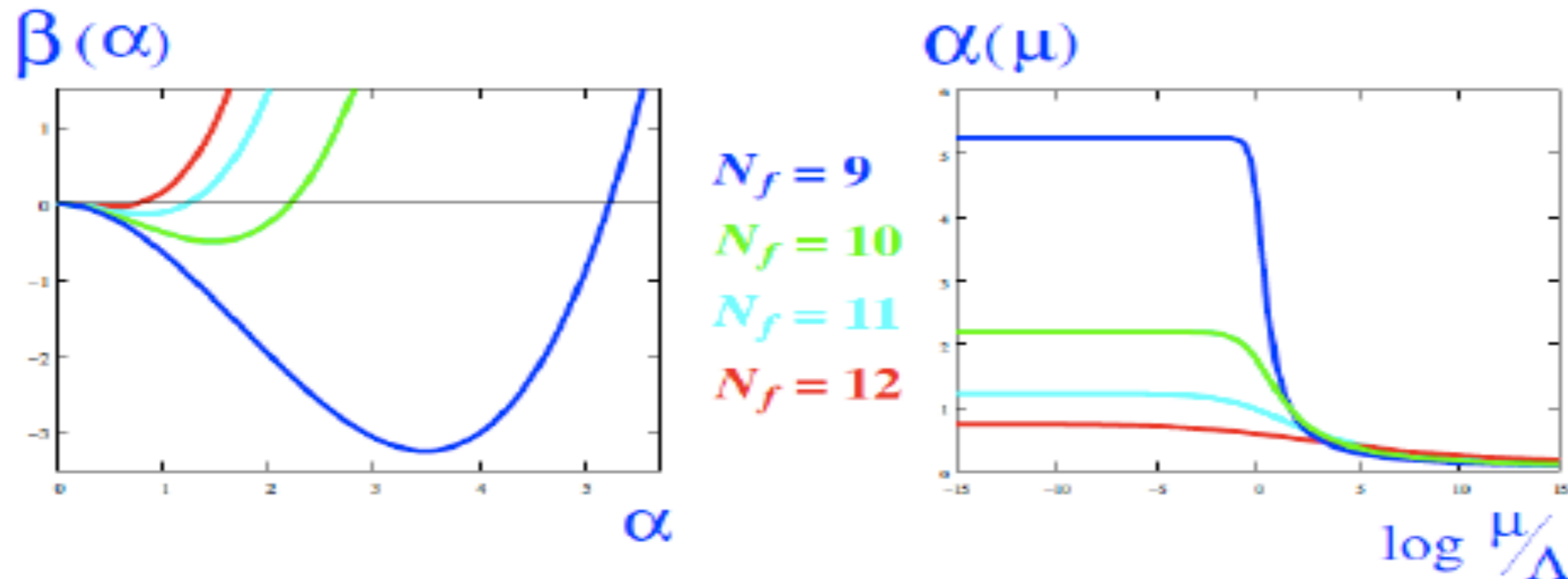
E.I., K. Nagata, Y. Nakagawa, A. Nakamura and V.I.Zakharov

PTEP 2016 (2016) no.6, 061B01, arXiv: 1512.01334 [hep-th]

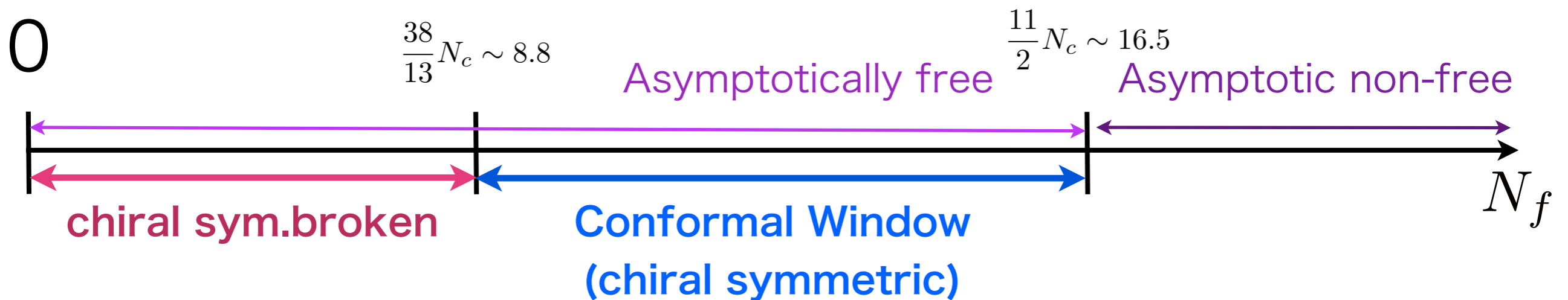
conformal window

# SU(3) gauge theory with massless $N_f$ fermions

Two loop analysis  $\beta(\alpha) = -b\alpha^2 - c\alpha^3$

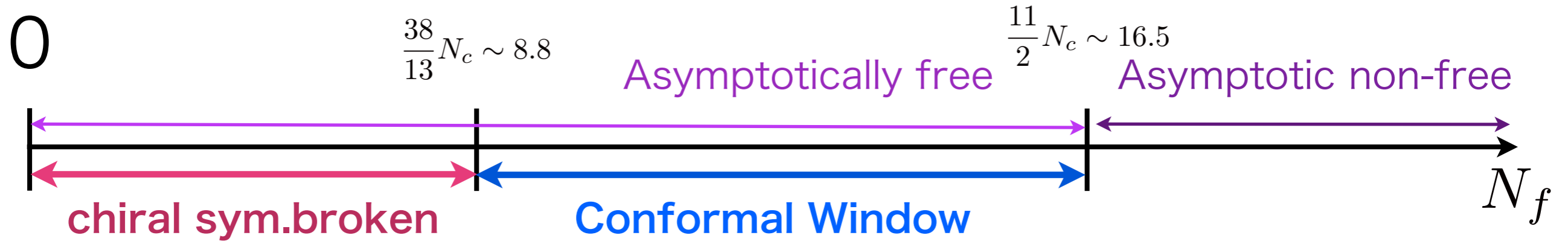


Phase structure based on two loop



# Conformal Window of SU(3) gauge

## Perturbation



## Schwinger-Dyson

$$N_f^* = 11.9$$



## Exact RG

$$N_f^* = 10.0^{+1.6}_{-0.7}$$



## Lattice simulation

only known nonperturbative and gauge invariant regularization

# Methods for IRFP study using lattice simulations

- \* Step scaling method (discretized renormalization group)  
measurement of the growth ratio of the Z-factor when the IR scale is changed

Luescher, Weisz and Wolff, NPB 359 (1991) 221

$$g_R^2(1/sL)/g_R^2(1/L)$$

- \* Hyperscaling for mass of composite state in mass deformed theory

Miransky, PRD59(1999)105003

Luty, JHEP 0904(2009)050

Del Debbio and Zwicky, PRD82(2010)014502

$$M_X \sim m_q^{1+\gamma_m^*}$$

- \* Hyperscaling for the Dirac eigenmodes for massless fermions

Patella, PRD86(2012)025006

Cheng, Hasenfratz, Retropulos and Schaich, JHEP1307(2013)061

$$\nu(\lambda) = V \int_{-\lambda}^{\lambda} \rho(\omega) d\omega \propto V \lambda^{\frac{4}{1+\gamma_m(\lambda)}}$$

- \* Shape of the correlation fn. for the composite op.

Ishikawa, Iwasaki, Nakayama and Yoshie, PRD87(2013)071503

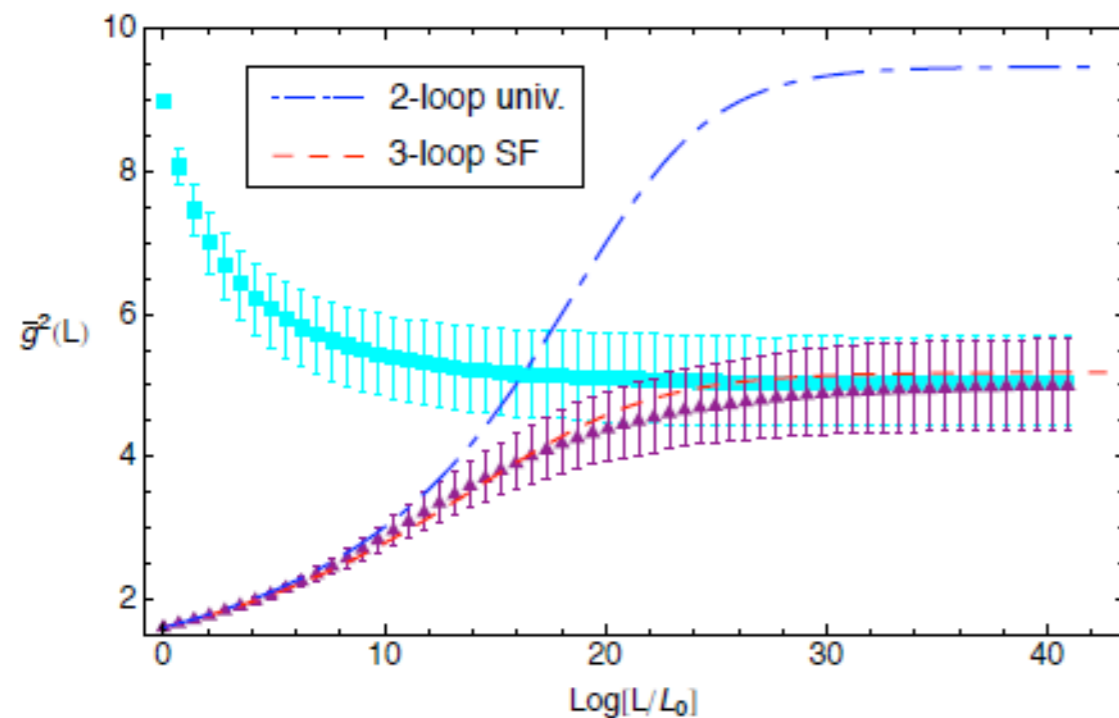
$$G_H(t) = \tilde{c}_H \frac{\exp(-\tilde{m}_H t)}{t^{\alpha_H}}$$

# Step scaling method

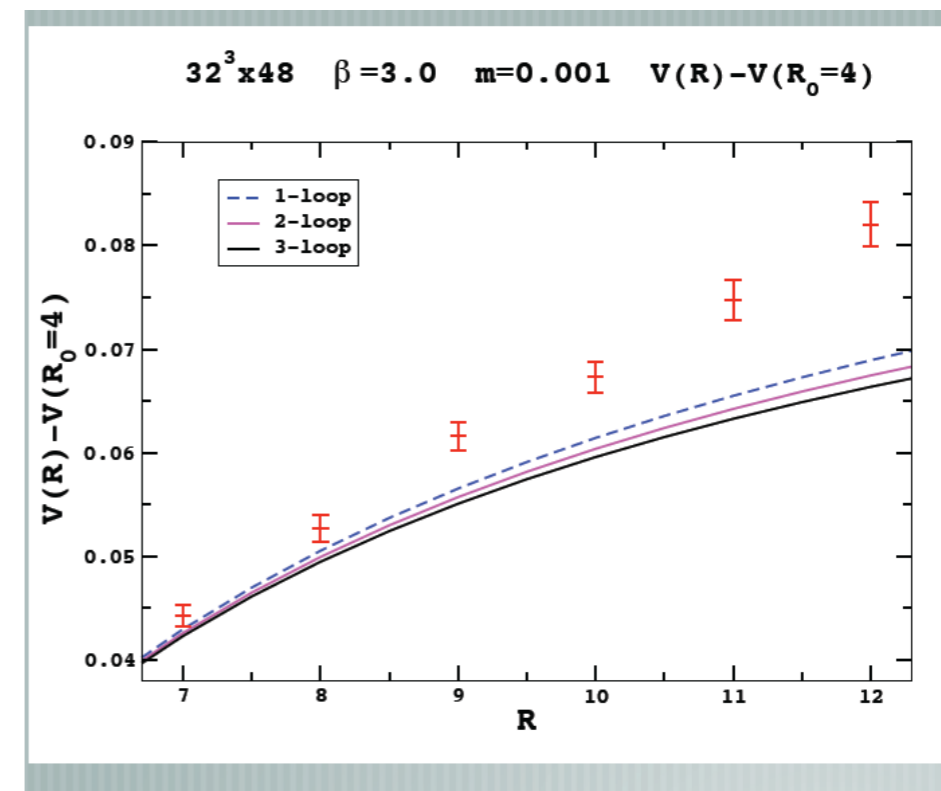
Running coupling constant/ renormalized coupling SU(3) Nf=12 theory

Appelquist et al. (SF scheme)  
Phys.Rev.D79:076010,2009

Taking constant extrapolation in (a/L)



Fodor et al. (potential scheme)  
PoS LAT2009:055,2009, talk at Lattice2010



Plot: Slide of K.Holland's talk at Lattice2010

The continuum extrapolation was not considered.  
( $O(a)$  effects depends on the renormalization scheme)

# Renormalization scheme and universality

scheme transformation

$$g_1 \rightarrow g_2 = f(g_1)$$

$f(g_1)$  is an analytic fn. of  $g_1$

$$\text{beta fn. } \beta(g_2) = \frac{\partial f(g_1)}{\partial g_1} \beta(g_1)$$

The existence of the fixed point is scheme independent.

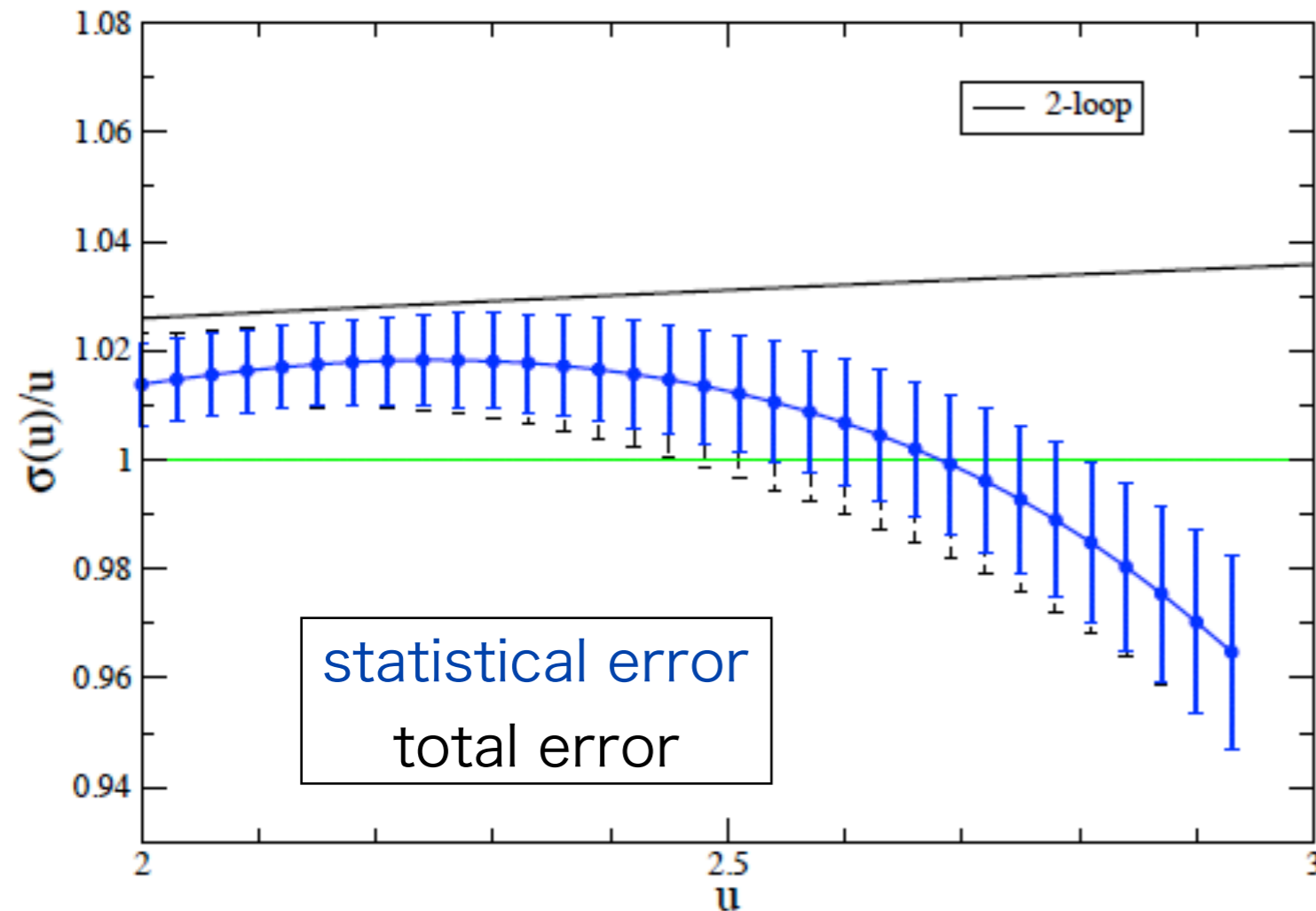
## Our work

- ◆ independent renormalization scheme
- ◆ taking continuum limit carefully
- ◆ exact massless simulation using twisted b.c.

# Growth ratio of TPL coupling (SU(3) Nf=12)

E.I. PTEP (2013) 083B01

$$\sigma(u)/u = g_R^2(1/sL)/g_R^2(1/L)$$



$$g_{\text{TPL}}^{*2} = 2.69 \pm 0.14(\text{stat.})_{-0.16}^0(\text{syst.})$$

Critical exponent of the beta fn.

$$\beta(g^2) \sim \gamma_g^*(g^{2*} - g^2)$$

Our result

$$\gamma_g^* = 0.57_{-0.31}^{+0.35}(\text{stat.})_{-0.16}^0(\text{syst.})$$

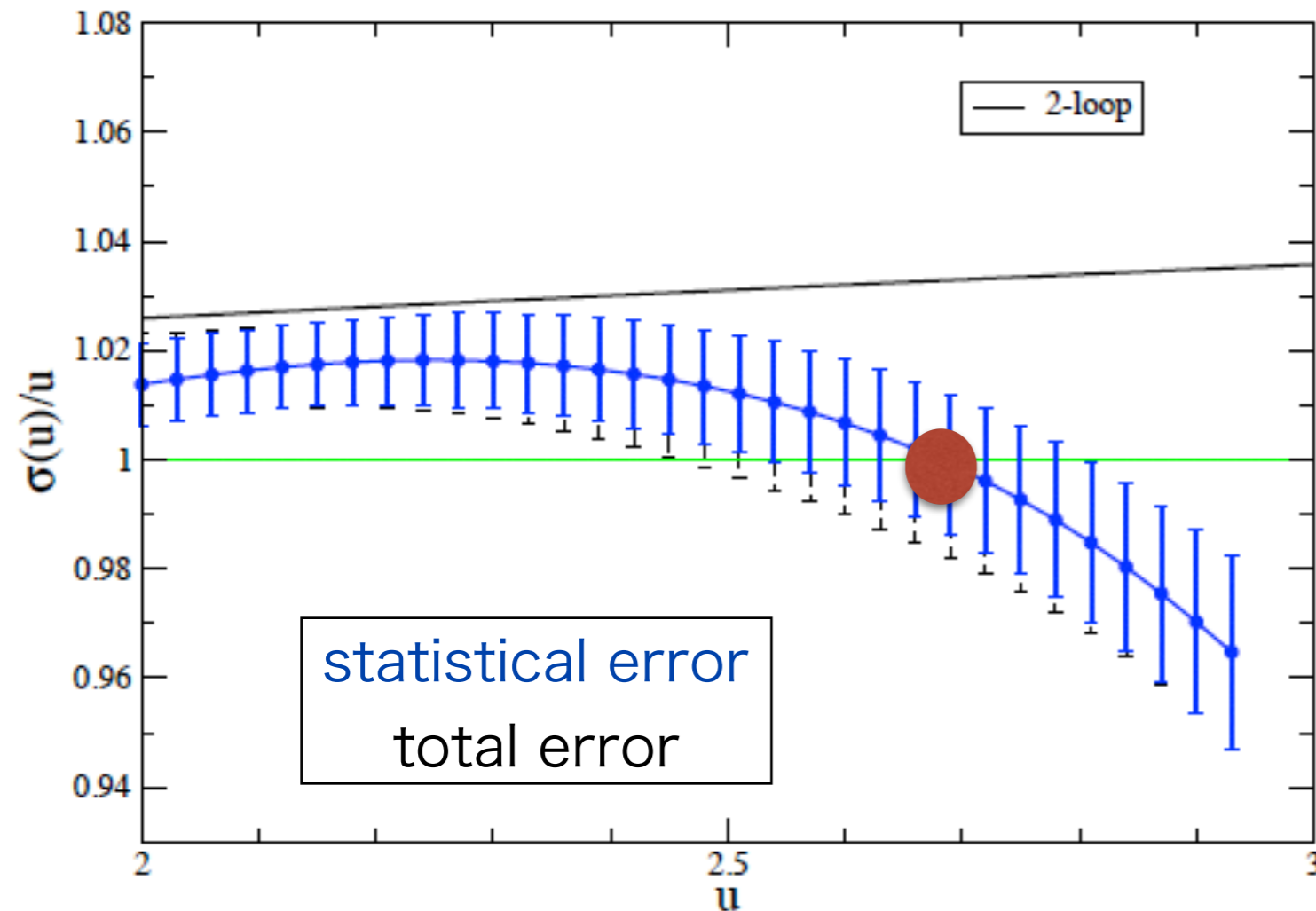
SF scheme	2 loop at $g^{2*} = 9.4$	4 loop (MS bar)
$\gamma_g^* = 0.13 \pm 0.03$	$\gamma_g^* = 0.36$	$\gamma_g^* = 0.28$



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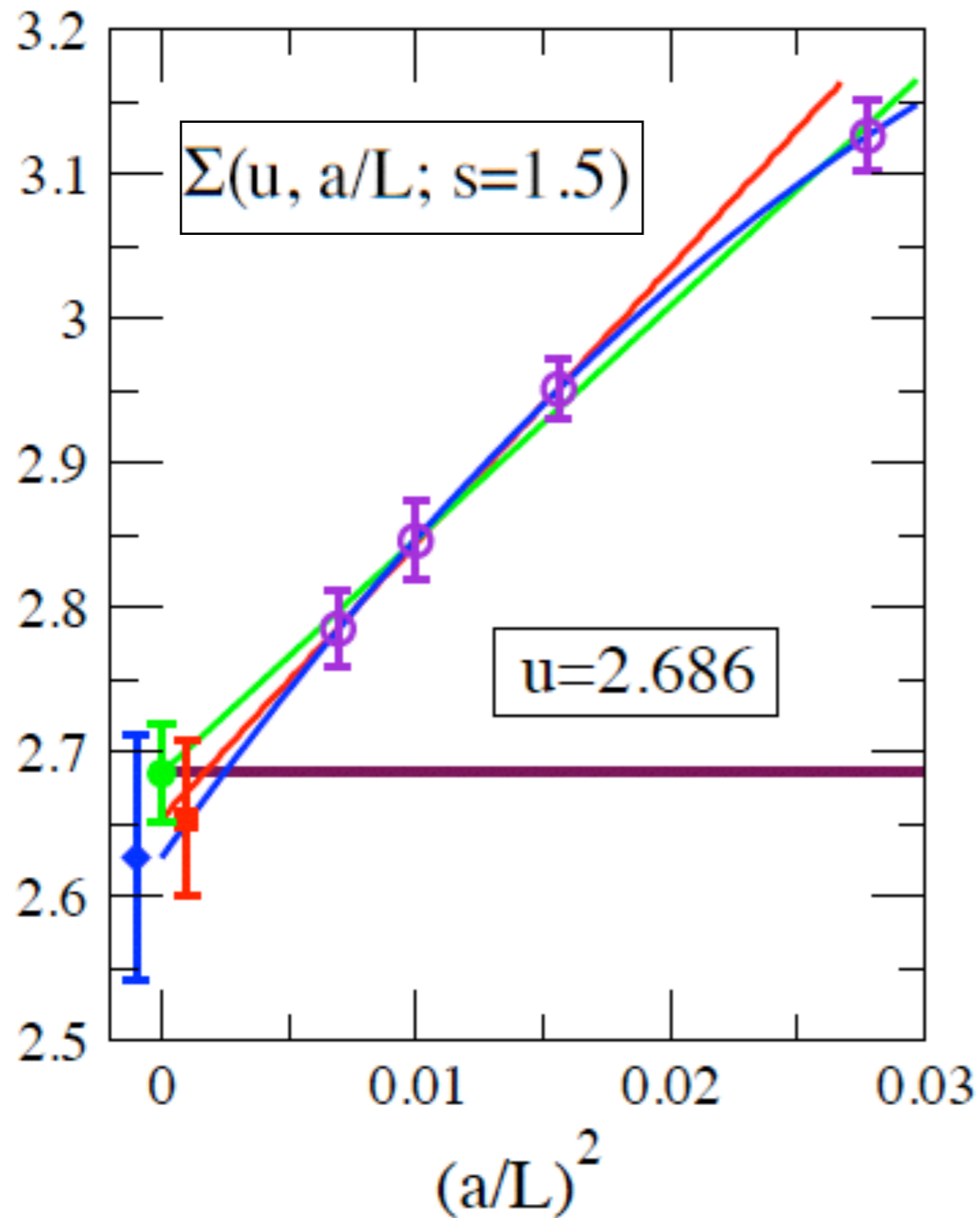
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# Continuum extrapolation



**s=1.5 step scaling**

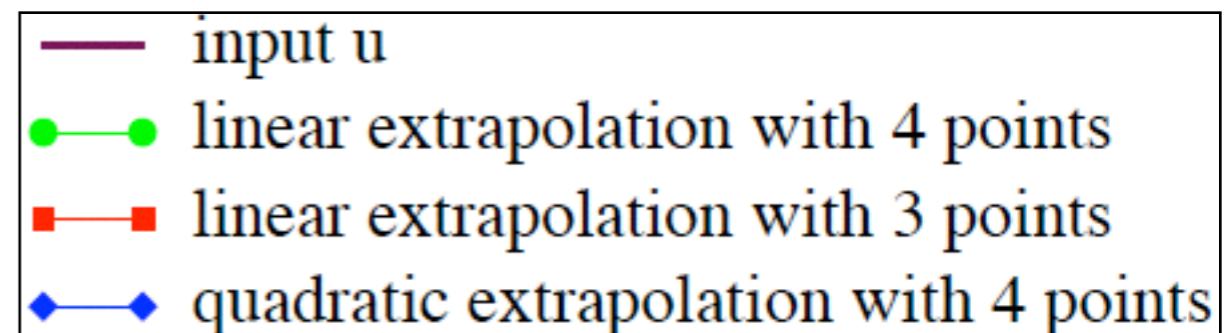
L/a=6 → L/a=9

L/a=8 → L/a=12

L/a=10 → L/a=15

L/a=12 → L/a=18

2 loop prediction  
in this region is  
 $\sigma(u = 2.69) \sim 2.78$



The systematic error is small in the strong coupling region in this scheme.

(Fit range dependence and “s” (step scaling parameter) dependence are also small.)

# Methods to obtain $\gamma_m^*$ using lattice simulations

- \* Step scaling method (discretized renormalization group)

Luescher, Weisz and Wolff, NPB 359 (1991) 221

mass anomalous dim  $\Leftrightarrow$  anomalous dim. of pseudo-scalar op.

- \* Hyperscaling for mass of composite state in mass deformed theory

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Luty, JHEP 0904(2009)050

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Patella, PRD86(2012)025006

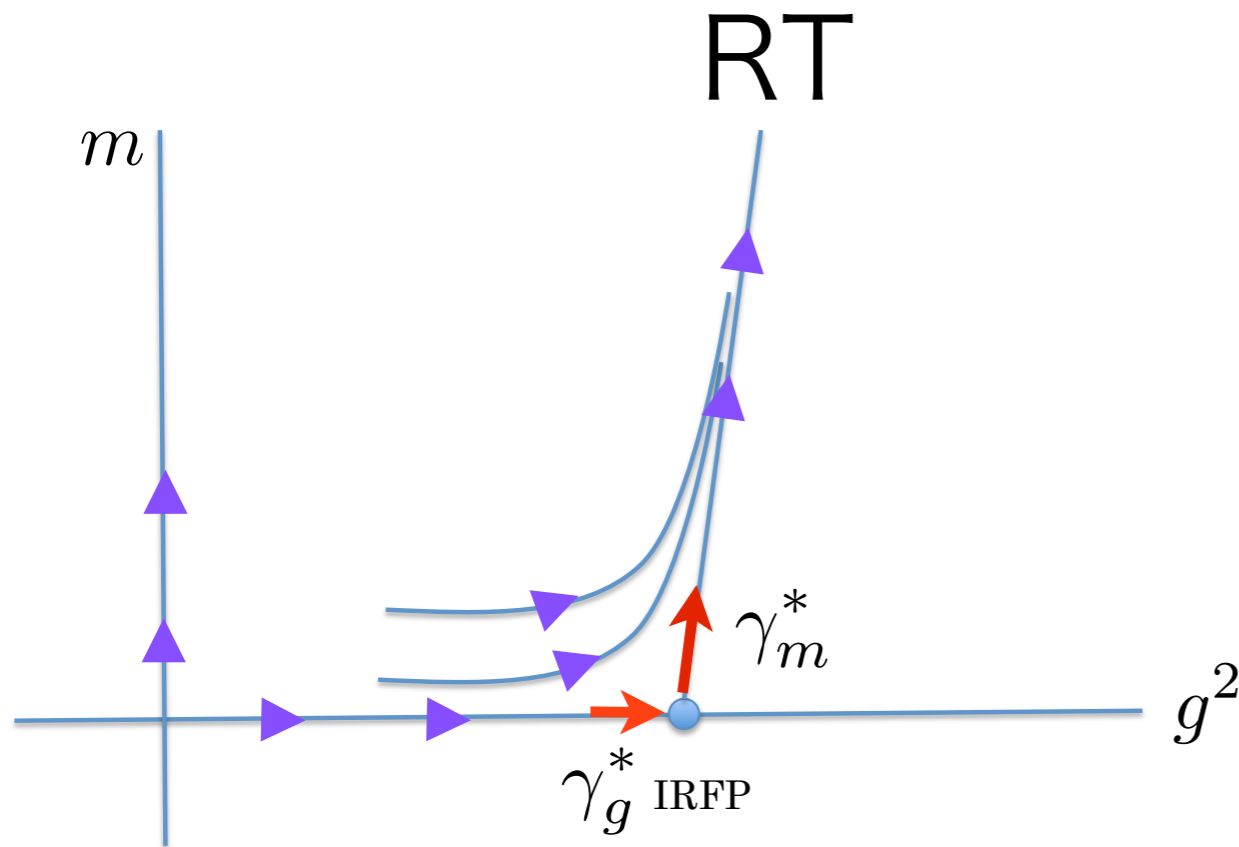
Cheng, Hasenfratz, et.al., JHEP1307(2013)061

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# Summary of critical exponents



Further studies are necessary to find the universal quantities.

**SF scheme** : PRD79 (2009) 076010  
**Fodor's data**: PLB703 (2011) 348-358  
**Fit(I)**: PRD84(2011) 054501  
**Fit(II)**: PRD84 (2011) 116901  
**LatKMI** : PRD86 (2012) 054506  
**Cheng et.al** : JHEP1307 (2013) 061  
**Ours** : PTEP (2013)083B01  
 arXiv: 1307.6645

	$\gamma_g^*$	$\gamma_m^*$
2loop	0.36	0.77
4loop (MS bar)	0.28	0.25
S.-D.		0.80
SF scheme	0.13(3)	
Fodor's data		0.403(13) 0.35(23)
LatKMI		0.4-0.5
Cheng et. al.		0.32(3)
Ours	0.57(35)	$0.081^{+0.03}_{-0.02}$

hyperscaling (mass)  $M_X \sim m_q^{\frac{1}{1+\gamma_m^*}}$

Del Debbio and Zwicky: Phys.Rev.D89(2014)014503

$\delta g \equiv g - g^*$  corrections (cont. limit)

$$m(b) = mb^{\gamma_m^*} \exp \left[ -\frac{\gamma_m^{(1)}}{\beta_1} \delta g f(b) \right]$$

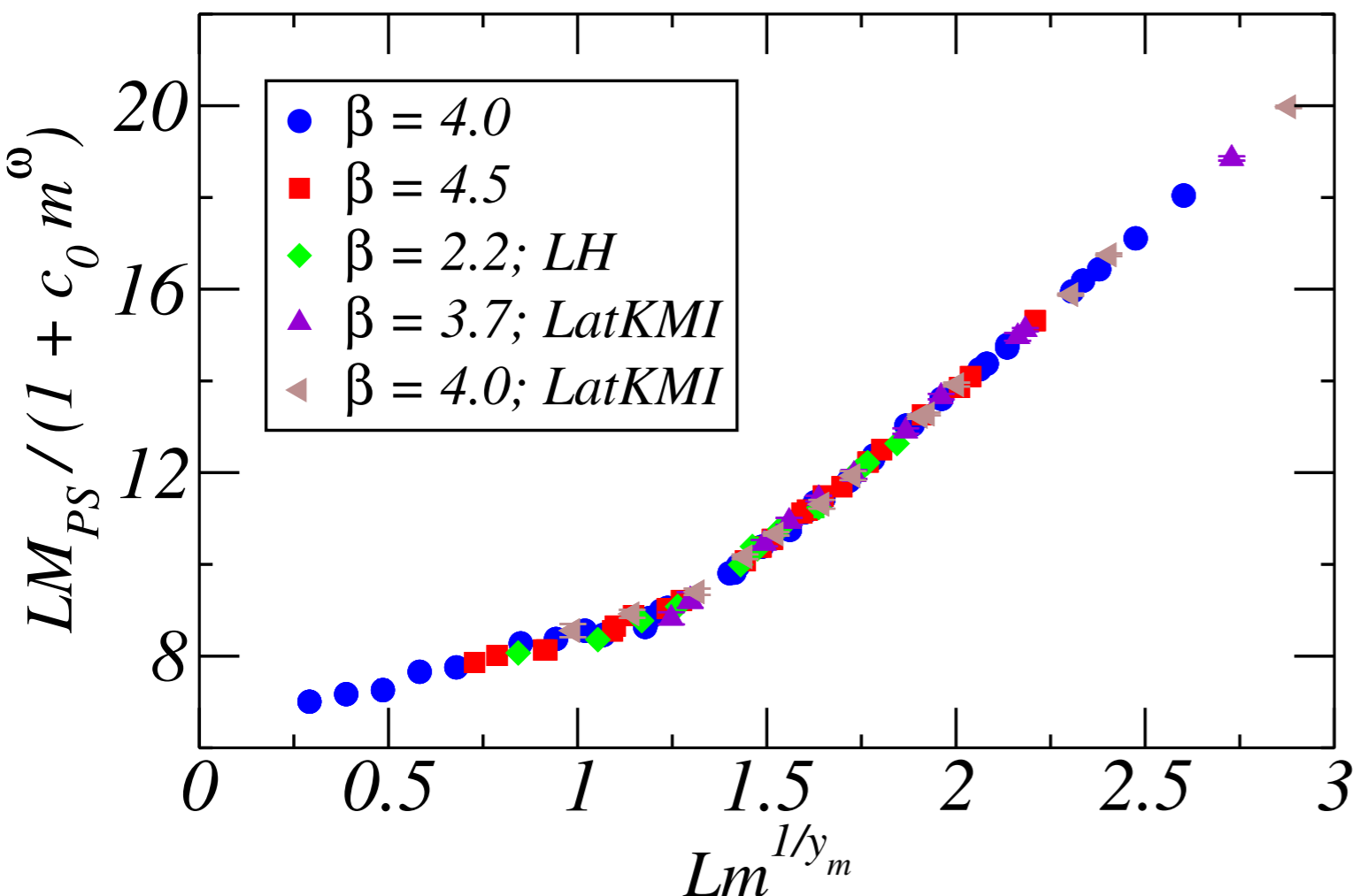
A.Hasenfratz

$\delta g \equiv g - g^*$  corrections (Finite vol. effect)

$$M_H = L^{-1} f(x, g_0 m^\omega), \quad \omega = -y_0 / y_m$$

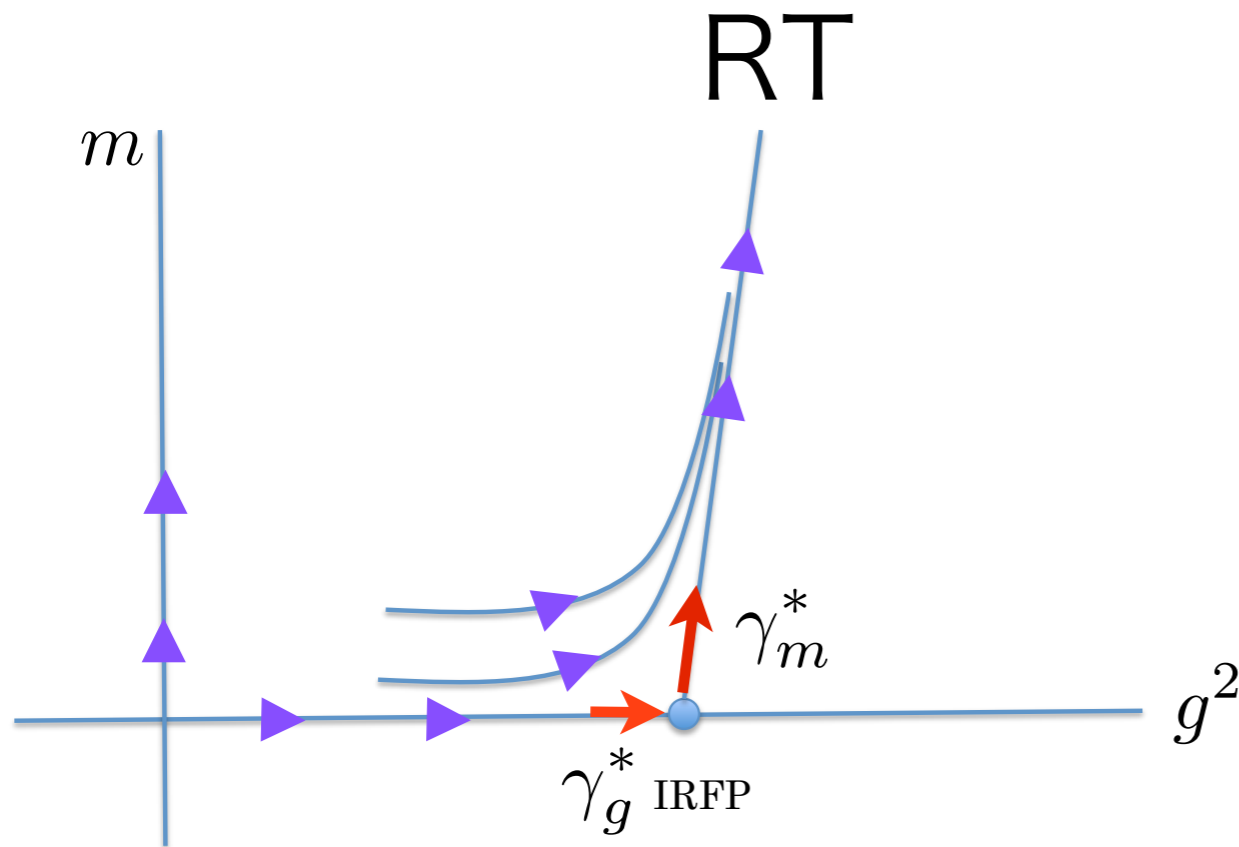
$$LM_H = F(x)(1 + g_0 m^\omega G(x))$$

pion at  $\beta=4.0, 4.5$ , LH, KMI :  
 $y_m=1.24[1]$ ,  $y_0=-0.51[5]$  ;  
 $\chi^2 / \text{dof} = 1.4 [95]$



Slide of seminar @KEK  
 by A.Hasenfratz

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LatKMI		
Cheng et. al.		
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# Conformal bootstrap

Assume that the absence of a relevant operator around the IRFP  
for the any reps.,  
the upper bound of the anomalous dimension is determined.

Iha, Makino, Suzuki arXiv:1603.01995

$$SU(N_f)_L \times SU(N_f)_R$$

staggered fermion Nf=12  $\gamma^* \leq 1.29$

staggered fermion Nf=8  $\gamma^* \leq 1.33$

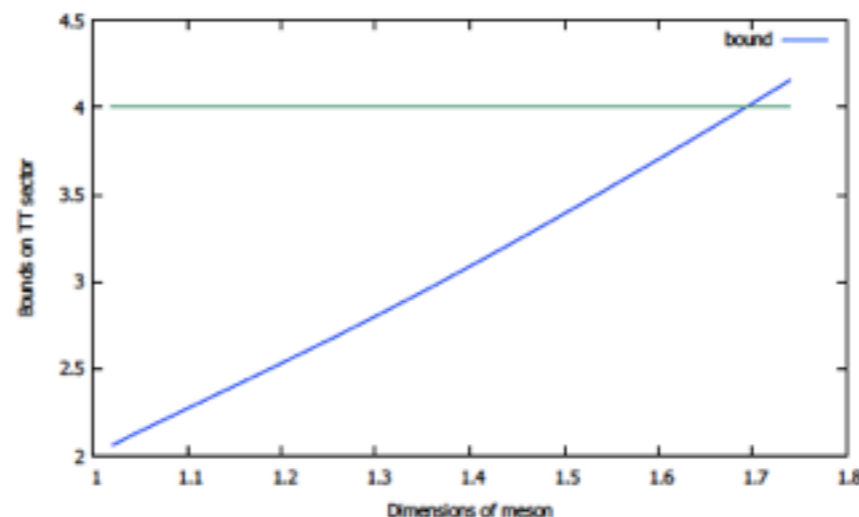
Y.Nakayama arXiv:1605.04052

$$SU(N_f)_V$$

$$SU(N_f)_L \times SU(N_f)_R$$

Wilson (or Domain-Wall) fermion Nf=8

$$\gamma^* \leq 1.31$$





# Conclusion and Directions (CFT part)

- ◆ Methods to find the IRFP in QCD-like theory
- ◆  $SU(3)$   $N_f=12$  has an IRFP
- ◆ Two critical exponents for beta-fn. and anomalous dim. are calculated
- ◆ The anomalous dim. at IRFP is quite smaller than the prediction by 2-loop or S-D eq.
- ◆ Critical  $N_f^*$  almost determine  $N_f^* \sim 8$
- ◆ Universality class as a CFT (central charge?)
  - ◆ Energy-momentum tensor
  - ◆ Entanglement entropy

# Plenary talks of Lattice conference

- ◆ G. T. Fleming, PoS (Lattice 2008) 021
- ◆ E. Pallante, PoS (Lattice 2009) 015
- ◆ L. Del Debbio, PoS (Lattice 2010) 004
- ◆ E. Neil, PoS (Lattice 2011) 009
- ◆ D. Negradi, PoS (Lattice 2011) 010
- ◆ J. Giedt, PoS (Lattice 2012) 006
- ◆ J. Kuti, PoS (Lattice 2013) 004
- ◆ E. Itou, PoS (Lattice 2013) 005
- ◆ Y. Aoki, PoS (LATTICE 2014) 011
- ◆ A. Hasenfratz, 2015

EMT on the lattice

# EMT on Lattice

## Energy-momentum tensor (EMT)

- ♦ generator of general coord. transformation
- ♦ conserved quantity (energy density, momentum, pressure)
- ♦ universal quantity (central charge in conformal theory)

## Lattice regularization

- ♦ nonperturbative regularization
- ♦ gauge invariant
- ♦ discretized space-time coordinate

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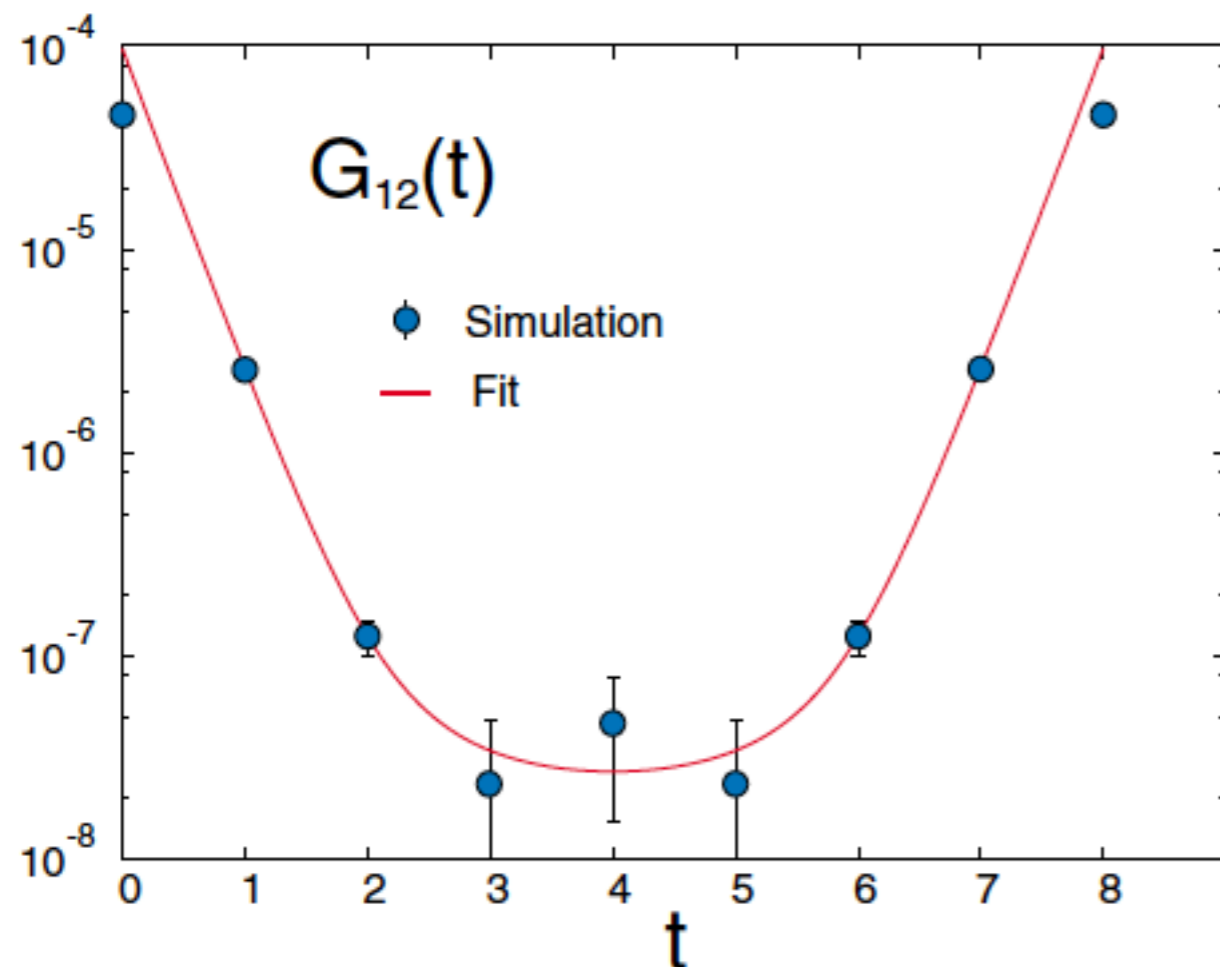
## Lattice regularization

- ♦ nonperturbative regularization
- ♦ gauge invariant
- ♦ discretized space-time coordinate

Same quantum number with the vac. (signal is noisy)

How to define the renormalized EMT

# Shear viscosity in QGP phase



Matsubara-Green's function  $G_{12}(t)$ ,  
Nakamura-Sakai(2005)  
800,000 conf.

shear viscosity: retarded Green's fn.

$$\eta = - \int \langle T_{12}(\vec{x}, \tau) T_{12}(\vec{x}', 0) \rangle_{\text{ret.}}$$

obtained by the analytic continuation  
of Matsubara Green's fn.

$$G_{\beta}(\vec{p}, t) = \sum_n e^{i\omega_n t} \int d\omega \frac{\rho(\vec{p}, \omega)}{i\omega_n - \omega}$$

quenched QCD  
(pure Yang-Mills theory)

Before going to the conformal theory...

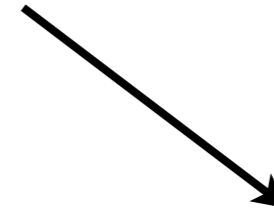
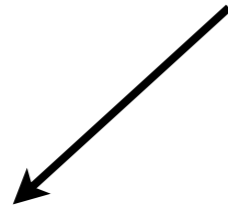
- ◆ give a better def. of EMT on the lattice
- ◆ show the result for the finite-T (quenched) QCD



# Basic Idea

Quantum field theory

(UV divergence)



perturbation with  
dim. reg.

+ **YM gradient flow**

(general covariance OK!)



lattice reg.

+ **Wilson flow**

(with  $a \rightarrow 0$  limit)

At finite flow time, **UV finite!**

Firstly, we obtain the relation between them perturbatively.

Assume that it applies to the nonperturbative regime.

# YM gradient flow

## Flow equation

Luescher, JHEP 1008, 071 (2010)

Yang–Mills gradient flow (continuum theory)

$$\partial_t B_\mu(t, x) = D_\nu G_{\nu\mu}(t, x) = \Delta B_\mu(t, x) + \dots, \quad B_\mu(t=0, x) = A_\mu(x)$$

Wilson flow (lattice theory)

$$\partial_t V(t, x, \mu) V(t, x, \mu)^{-1} = -g_0^2 \partial \mathcal{S}_{\text{Wilson}}, \quad V(t=0, x, \mu) = U(x, \mu)$$

link variable:  $U_\mu(x) = e^{ig_0 A_\mu(x)}$

t: fictitious time direction (flow-time)

$$x = (\vec{x}, \tau)$$

# UV finiteness of the gradient flow

## Flow equation (continuum)

$$\partial_t B_\mu(t, x) = D_\nu G_{\nu\mu}(t, x) \quad \text{initial condition: } B_\mu(t = 0, x) = A_\mu(x)$$

## perturbative solution in the leading order

$$B_\mu(t, x) = \int d^D y K_t(x - y) A_\mu(y)$$

$$K_t(z) = \int \frac{d^D p}{(2\pi)^D} e^{ipz} e^{-tp^2}$$

signal becomes clear?

$p^2 > 1/t$  modes are suppressed (a smooth UV cutoff)

Smearred in the range  $|x| < \sqrt{8t}$

Finiteness is shown perturbatively in all order

Luescher and Weisz, JHEP 1102, 051 (2011)

# Energy-momentum tensor

$$T_{\mu\nu}(x) = \frac{1}{g_0^2} \left[ F_{\mu\rho}^a(x) F_{\nu\rho}^a(x) - \frac{1}{4} \delta_{\mu\nu} F_{\rho\sigma}^a(x) F_{\rho\sigma}^a(x) \right]$$

Renormalized EMT within dim. reg.

$$\{T_{\mu\nu}\}_R(x) = T_{\mu\nu}(x) - \langle T_{\mu\nu}(x) \rangle$$

Dim=4 rank2 symmetric operator

$$U_{\mu\nu}(t, x) \equiv G_{\mu\rho}^a(t, x) G_{\nu\rho}^a(t, x) - \frac{1}{4} \delta_{\mu\nu} G_{\rho\sigma}^a(t, x) G_{\rho\sigma}^a(t, x)$$

$$E(t, x) \equiv \frac{1}{4} G_{\mu\nu}^a(t, x) G_{\mu\nu}^a(t, x)$$

Here, ops. are constructed by flowed field.

# “Suzuki method”

- small flow-time expansion -

Suzuki, PTEP 2013, no8, 083B03, [Erratum: PTEP2015,079201 (2015)],

relation... dim.=4 op on the lattice vs. renormalized EMT at small flow-time

$$U_{\mu\nu}(t, x) = \alpha_U(t) \left[ \{T_{\mu\nu}\}_R(x) - \frac{1}{4} \delta_{\mu\nu} \{T_{\rho\rho}\}_R(x) \right] + O(t),$$

$$E(t, x) = \langle E(t, x) \rangle + \alpha_E(t) \{T_{\rho\rho}\}_R(x) + O(t),$$

coefficients... given by renormalized coupling and coeff. of beta fn.

$$\alpha_U(t)(g; \mu) = g^2 \left\{ 1 + 2b_0 \left[ \ln(\sqrt{8t}\mu) + s_1 \right] g^2 + O(g^4) \right\},$$

$$\alpha_E(t)(g; \mu) = \frac{1}{2b_0} \left\{ 1 + 2b_0 s_2 g^2 + O(g^4) \right\},$$

$b_0$  1-loop coeff. of beta fn.

MSbar scheme

$$s_1 = -0.0863575$$

$$s_2 = 0.05578512$$

cf.) Nonperturbative method:

L.DelDebbio, A.Patella, A.Rogo, JHEP 1311,212(2013)

# How to get EMT

## Step 1 for quenched QCD

Generate gauge configuration at  $t=0$  (usual process)

## Step 2

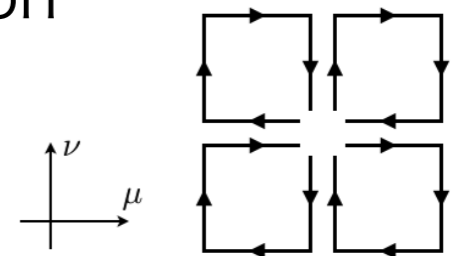
Solve the Wilson flow eq. and generate the gauge configuration at flow time ( $t$ )

$$a \ll \sqrt{8t} \ll \Lambda_{QCD}^{-1} \text{ or } T^{-1}$$

## Step 3

Measure two dim=4 ops. using flowed gauge configuration

$$U_{\mu\nu}(t, x), E(t, x)$$



## Step 4

Take the continuum limit. Then take  $t \rightarrow 0$  limit.

(Take care the feasible window of flow time)

$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left\{ \frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} [E(t, x) - \langle E(t, x) \rangle_0] \right\}$$

# One-point fn. of EMT in finite temperature quenched QCD

Asakawa, Hatsuda, E.I., Kitazawa, Suzuki (FlowQCD coll.)  
Phys.Rev. D90 (2014) 1, 011501

# Simulation setup

- Wilson plaquette gauge action
- lattice size ( $N_s=32$ ,  $N_t=6,8,10,32$ )
- # of confs. is 100 - 300
- simulation parameters

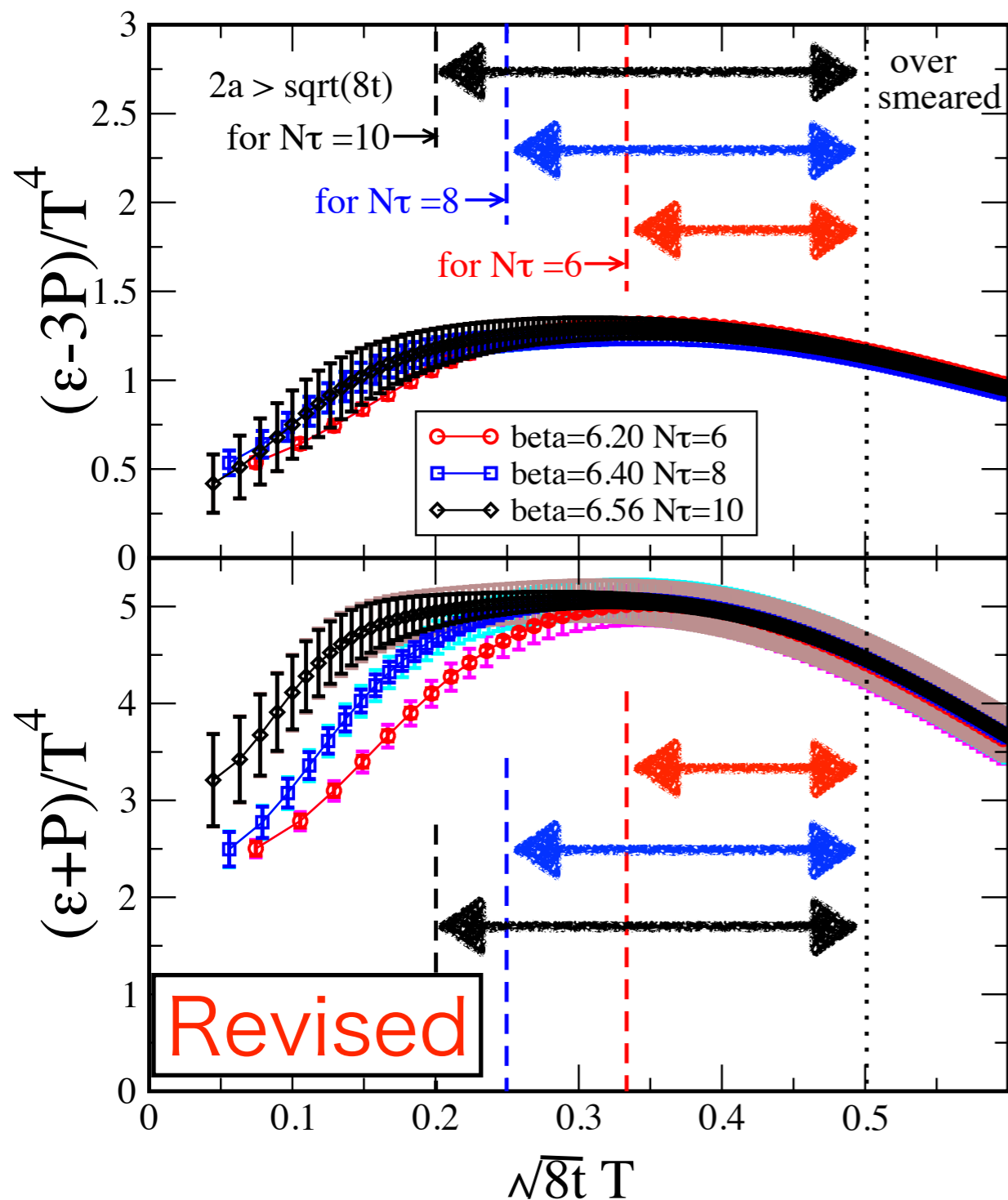
$N_\tau$	6	8	10	$T/T_c$
	6.20	6.40	6.56	1.65
$\beta$	6.02	6.20	6.36	1.24
	5.89	6.06	6.20	0.99

Temperature is determined by  
[Boyd et. al. NPB469,419 \(1996\)](#)

Parametrization is given by  
[alpha collaboration NPB538,669 \(1999\)](#)



# flow time dependence ( $T=1.65T_c$ )



trace anomaly  $\sum_{i=1}^4 T_{ii} = \frac{\epsilon - 3P}{T^4}$

entropy density  $T_{44} - T_{11} = \frac{\epsilon + P}{T^4}$

## feasible flow time

longer than lattice cutoff

avoid an over-smeared regime

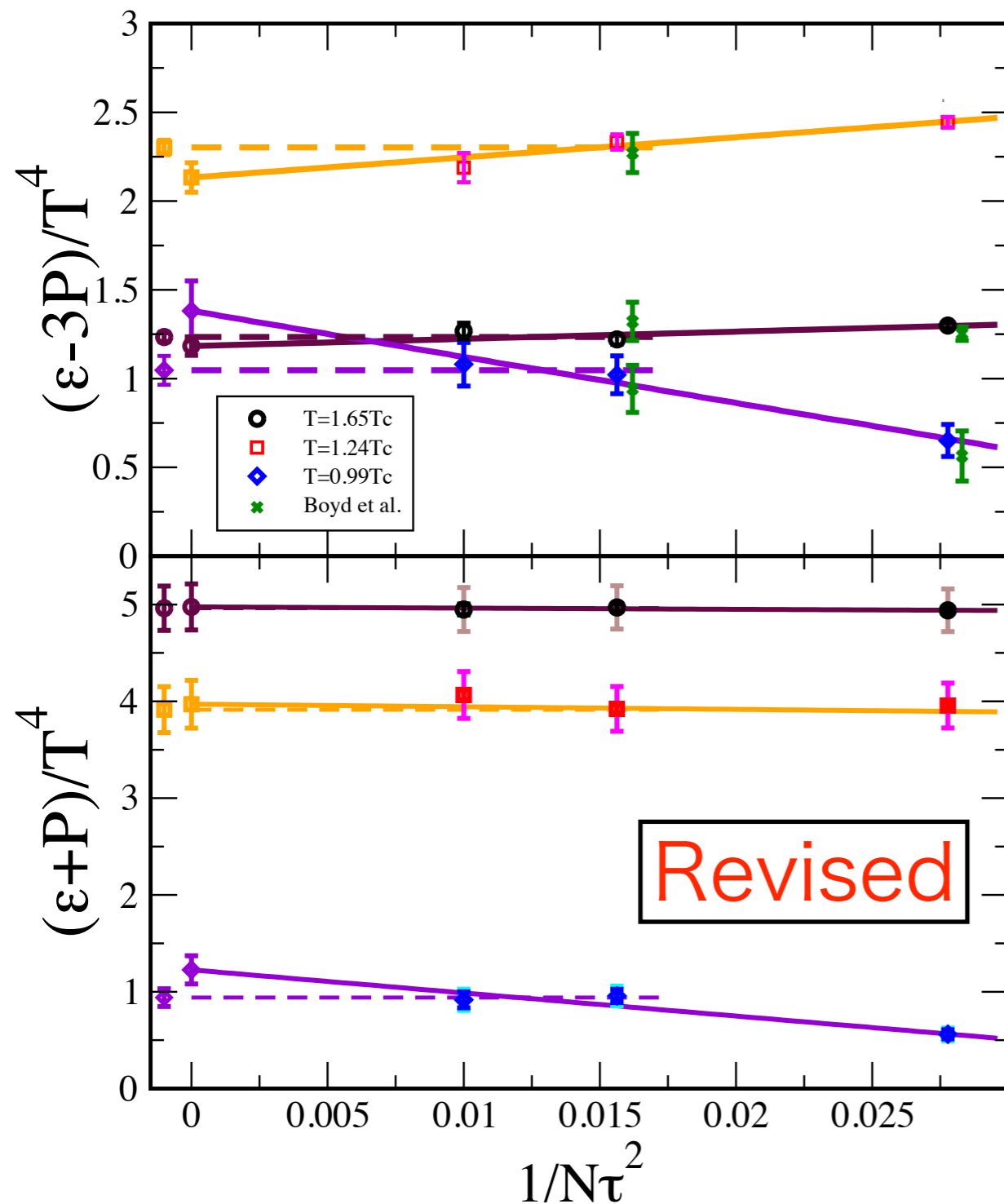
$$2a < \sqrt{8t} < N_\tau a/2$$

- show a plateau  
(small higher dimensional op.)  
Practically, no need  $t \rightarrow 0$  limit  
**\*\*finer lattice simulation shows a slope**

- systematic error coming from scale setting is dominated in entropy density

each dark color shows statistical error  
each light color includes systematic error

# Continuum extrapolation



$$\sqrt{8tT} = 0.40$$

3point linear extrap.  
(2pt. const. extrap.)

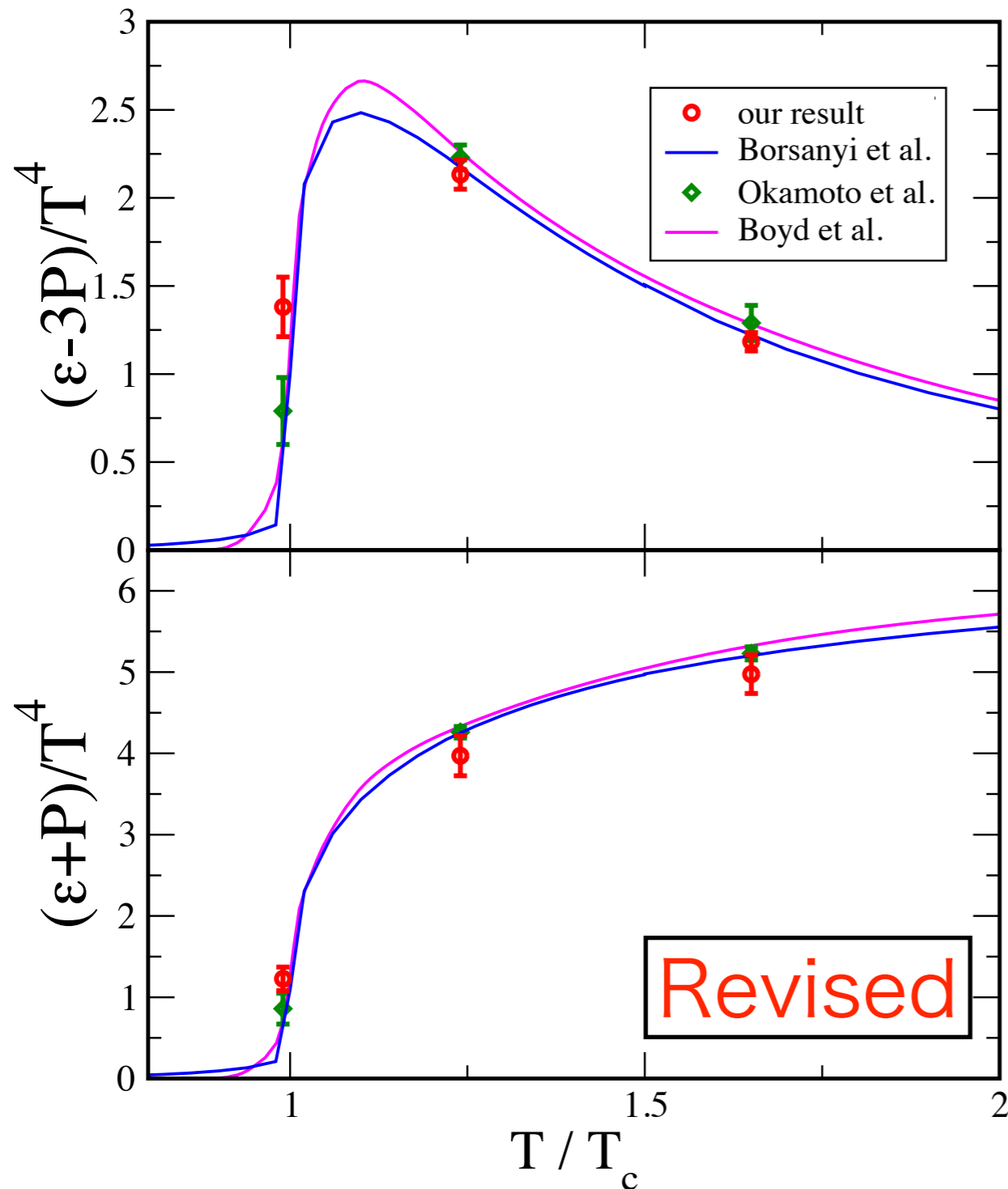
We also see the data at  $\sqrt{8tT} = 0.35$   
In cont.lim. the result is consistent.



$t \rightarrow 0$  limit is not needed  
in this case

# Comparison with the results given by integration method

Phys.Rev. D90 (2014) 1, 011501, arXiv:1312.7492v3[hep-lat]



Boyd et. al. NPB469,419 (1996)

Okamoto et. al. (CP-PACS) PRD60, 094510 (1999)

Borsanyi et. al. JHEP 1207, 056 (2012)

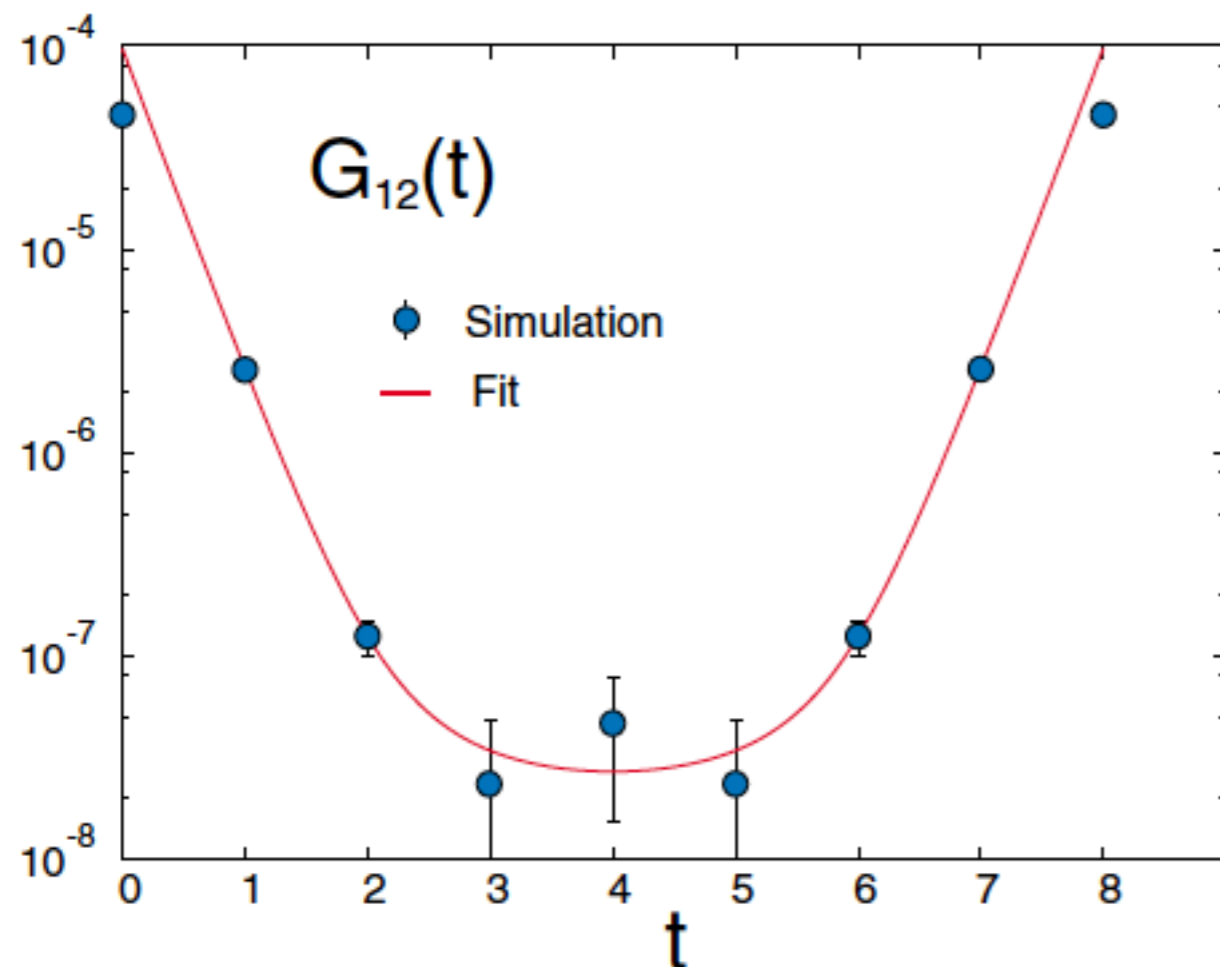
Integration method is based on (macroscopic) thermodynamics.

Our method is based on the (microscopic) quantum field theory.

two-point fn. of EMT

preliminary

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quenched QCD  
 (pure Yang-Mills theory)

# Renormalization

$$T_{\mu\nu}^{(R)}(g_0) = Z(g_0)T_{\mu\nu}^{(bare)}$$

Meyer (2007) ... 1 loop approximation

Fodor et al. (2013) ... calculate Z-factor from entropy density

This work ... Not necessary

(usage of Suzuki coefficient and MSbar coupling)

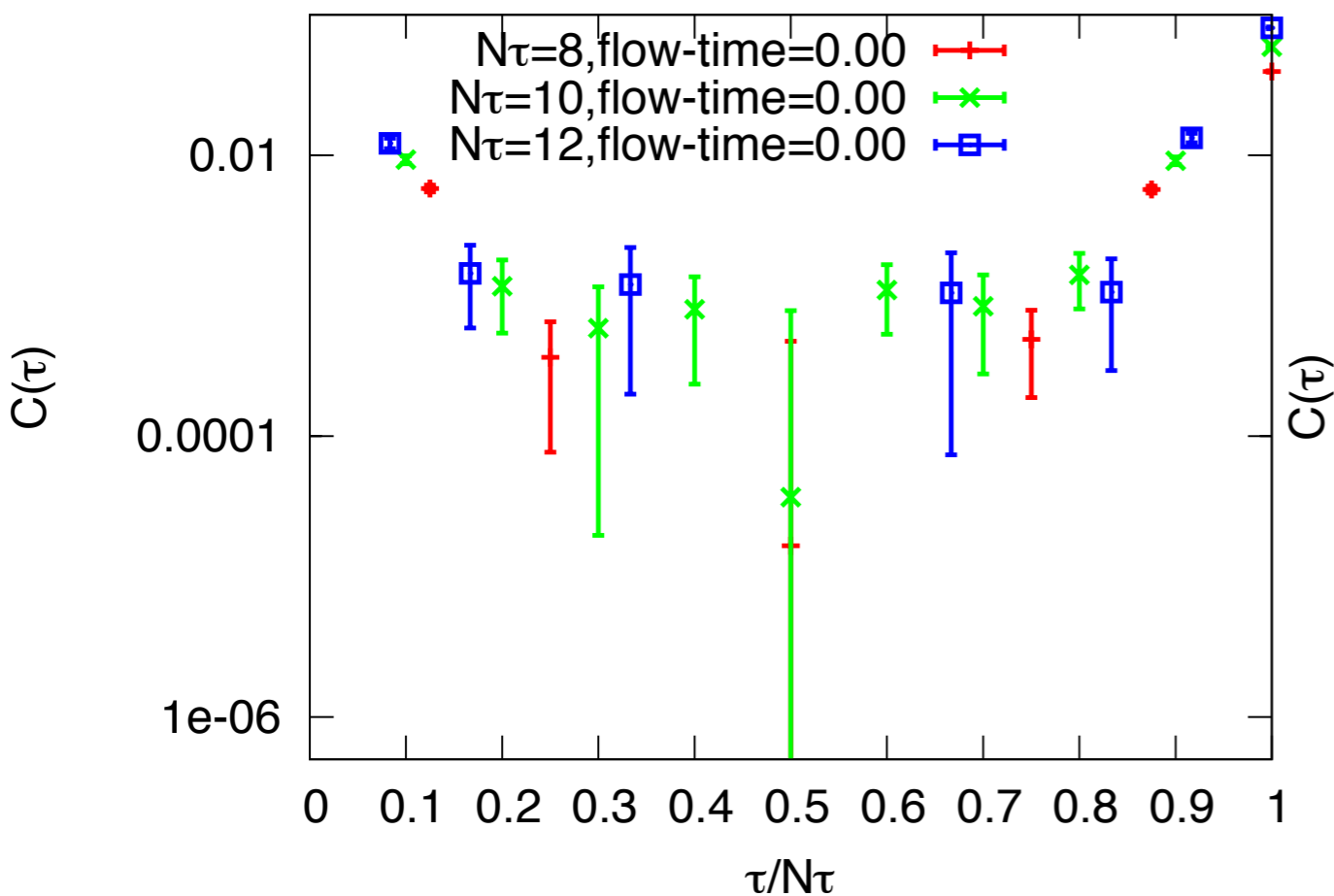
cf.)  $\frac{sT}{4} = \langle T_{11}^{(R)} \rangle$

$$\langle T_{12}T_{12} \rangle = \frac{1}{4} \langle (T_{11} - T_{22})(T_{11} - T_{22}) \rangle$$

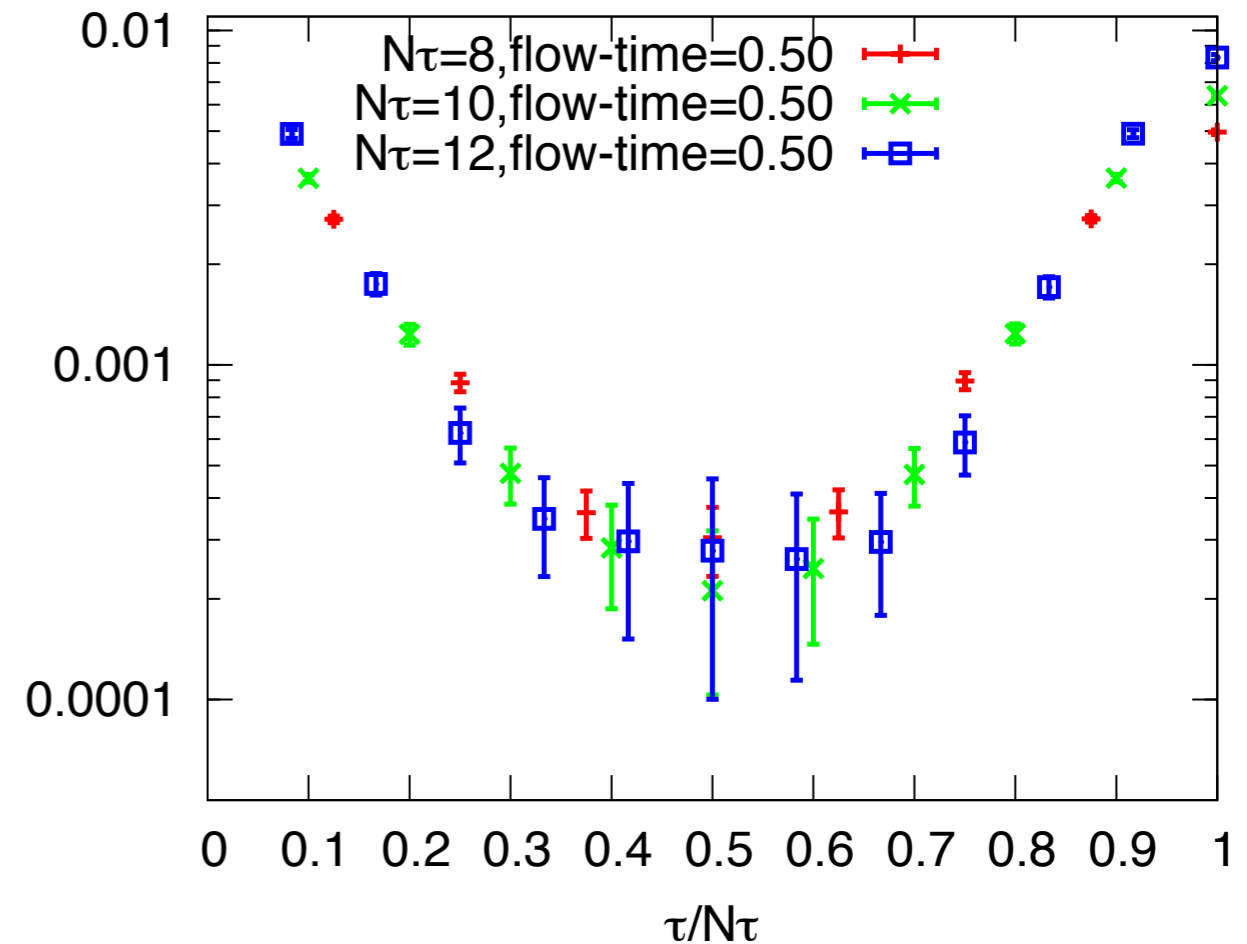
# lattice raw data

$$C(\tau) = \left\langle \frac{1}{N_s^3} \sum_{\vec{x}} U_{12}(\vec{x}, \tau) \frac{1}{N_s^3} \sum_{\vec{y}} U_{12}(\vec{y}, 0) \right\rangle$$

flow-time=0



flow-time  $t/a^2=0.50$



fixed smeared length in lattice unit

beta=6.40,  $N_t=8$ , 2,000 conf.

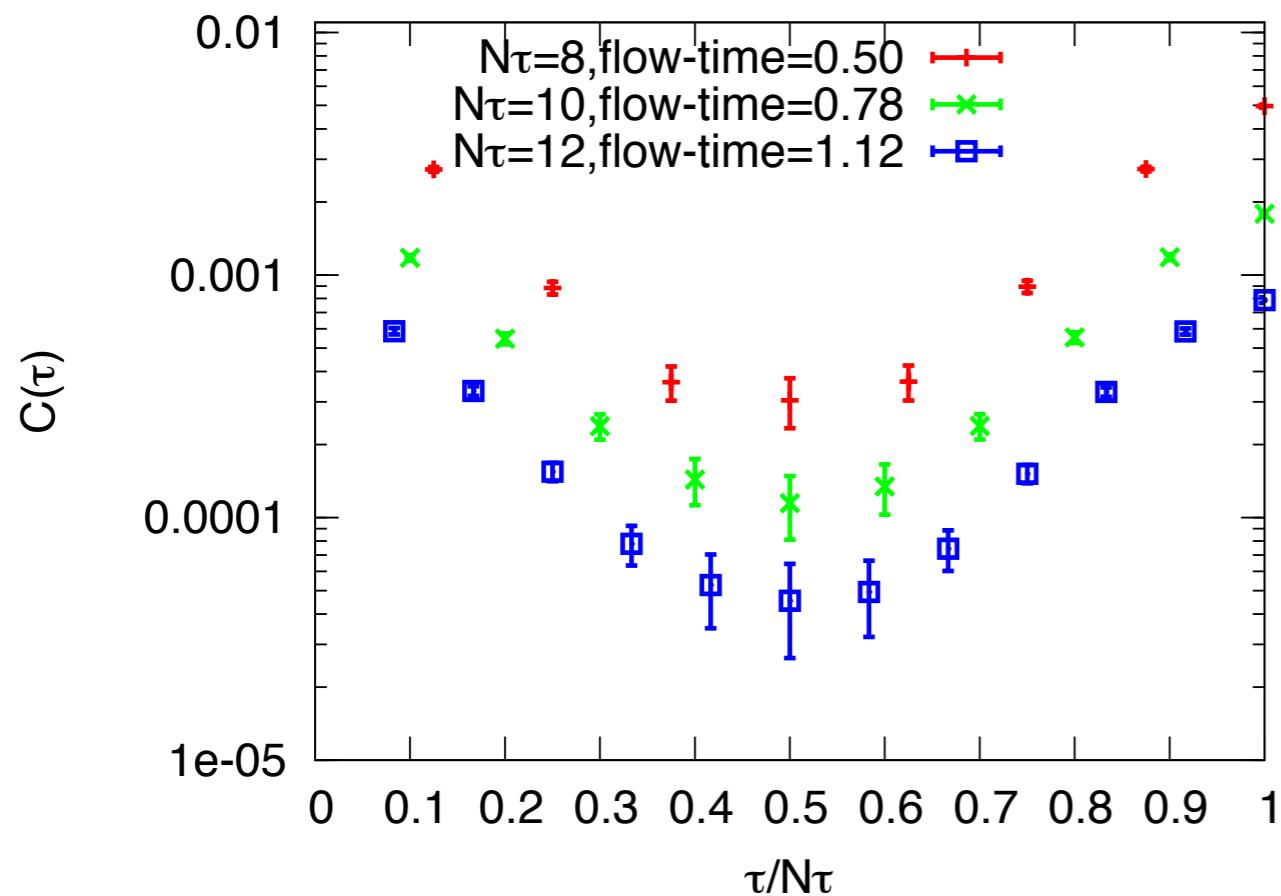
beta=6.57,  $N_t=10$ , 1,100 conf.

beta=6.72,  $N_t=12$ , 650 conf.

# EMT correlator

$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left\{ \frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} [E(t, x) - \langle E(t, x) \rangle_0] \right\}$$

$$C(\tau) = \left\langle \frac{1}{N_s^3} \sum_{\vec{x}} U_{12}(\vec{x}, \tau) \frac{1}{N_s^3} \sum_{\vec{y}} U_{12}(\vec{y}, 0) \right\rangle$$



fixed smeared length in physical unit  $\sqrt{8tT} = 0.25$

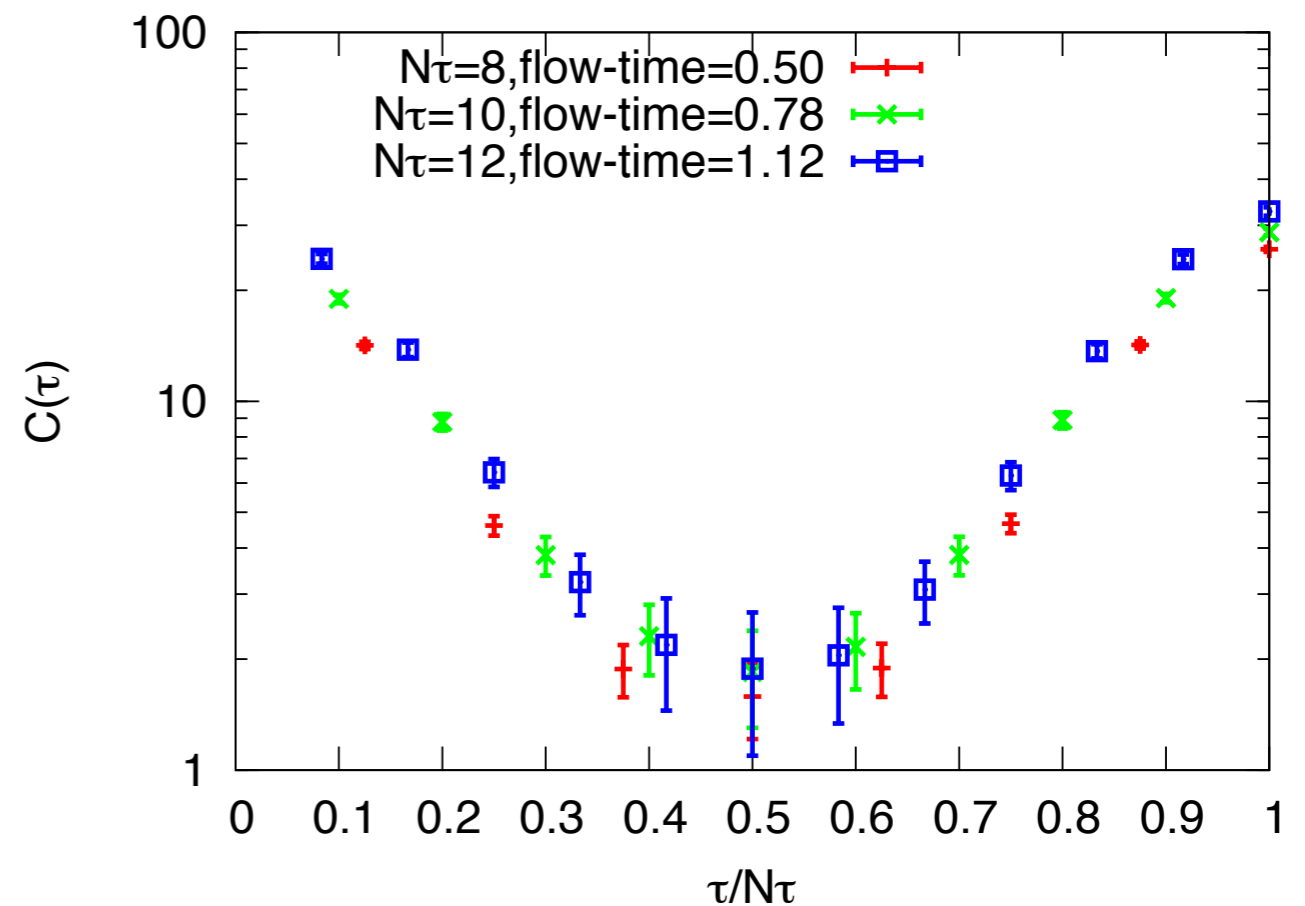
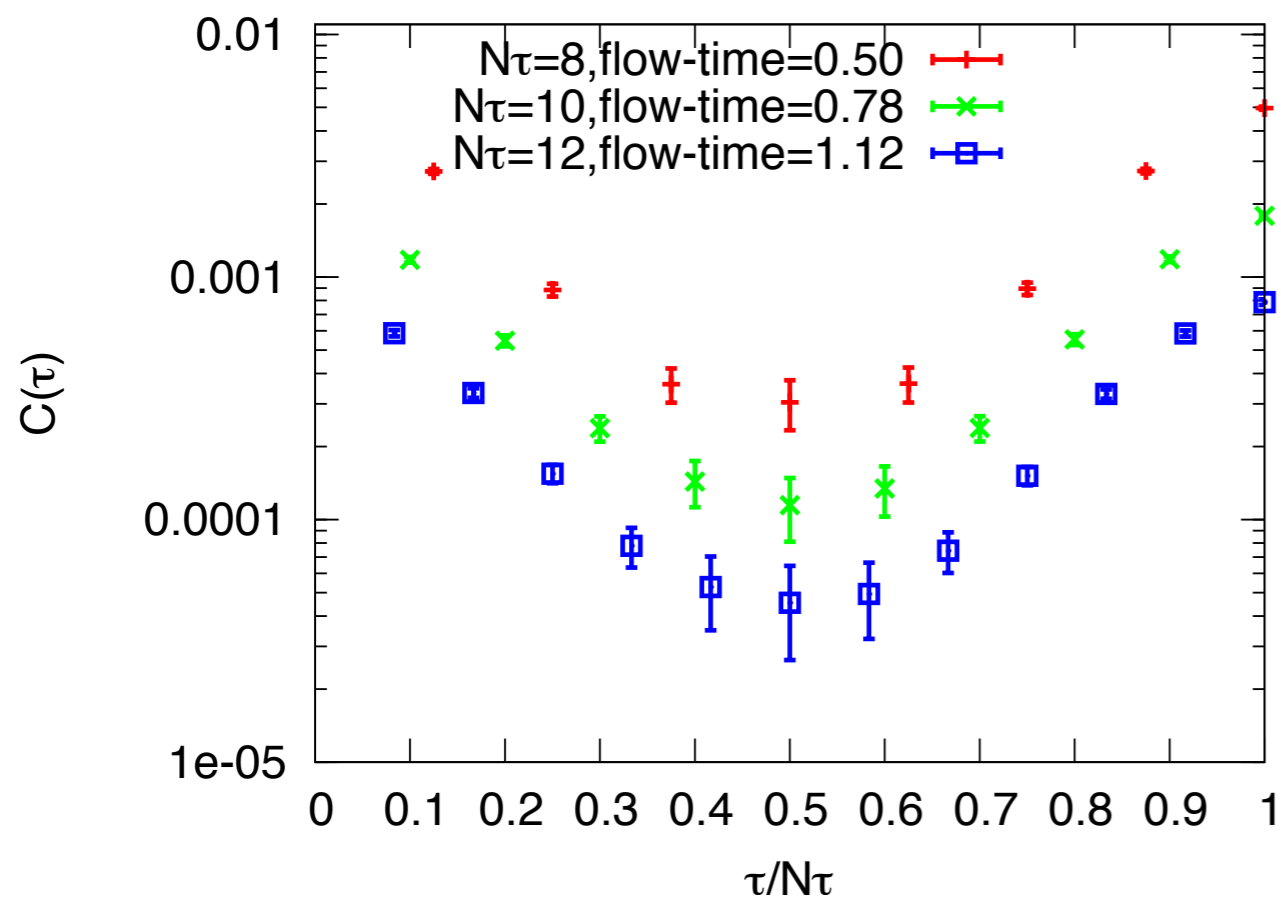


# EMT correlator

$$T_{\mu\nu}^R(x) = \lim_{t \rightarrow 0} \left\{ \frac{1}{\alpha_U(t)} U_{\mu\nu}(t, x) + \frac{\delta_{\mu\nu}}{4\alpha_E(t)} [E(t, x) - \langle E(t, x) \rangle_0] \right\}$$

$$C(\tau) = \left\langle \frac{1}{N_s^3} \sum_{\vec{x}} U_{12}(\vec{x}, \tau) \frac{1}{N_s^3} \sum_{\vec{y}} U_{12}(\vec{y}, 0) \right\rangle$$

$$C(\tau) = \frac{1}{T^5} \left\langle \sum_{\vec{x}} T_{12}(\vec{x}, \tau) \sum_{\vec{y}} T_{12}(\vec{y}, 0) \right\rangle$$



fixed smeared length in physical unit  $\sqrt{8tT} = 0.25$

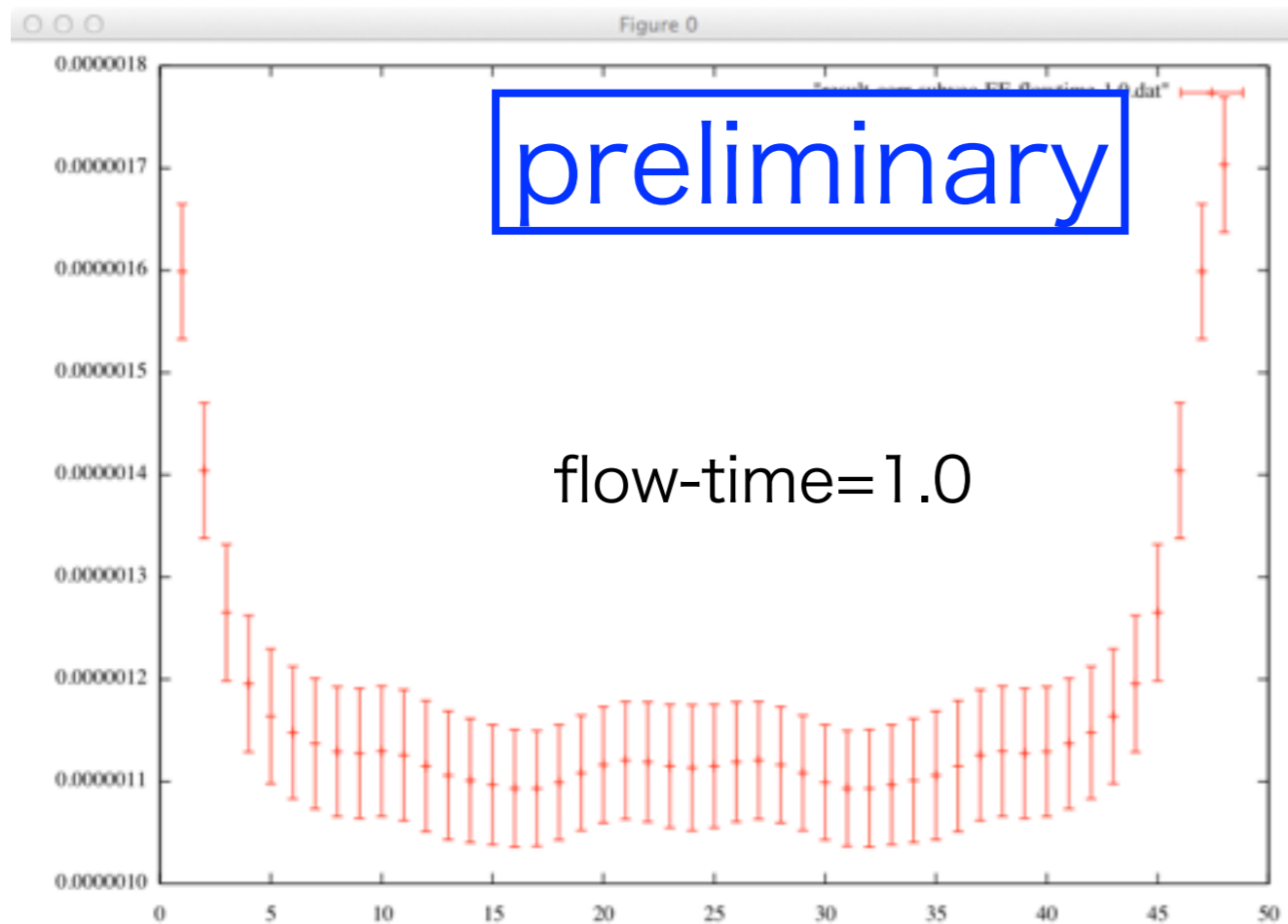
Nf=12 case  
(conformal theory)

very preliminary

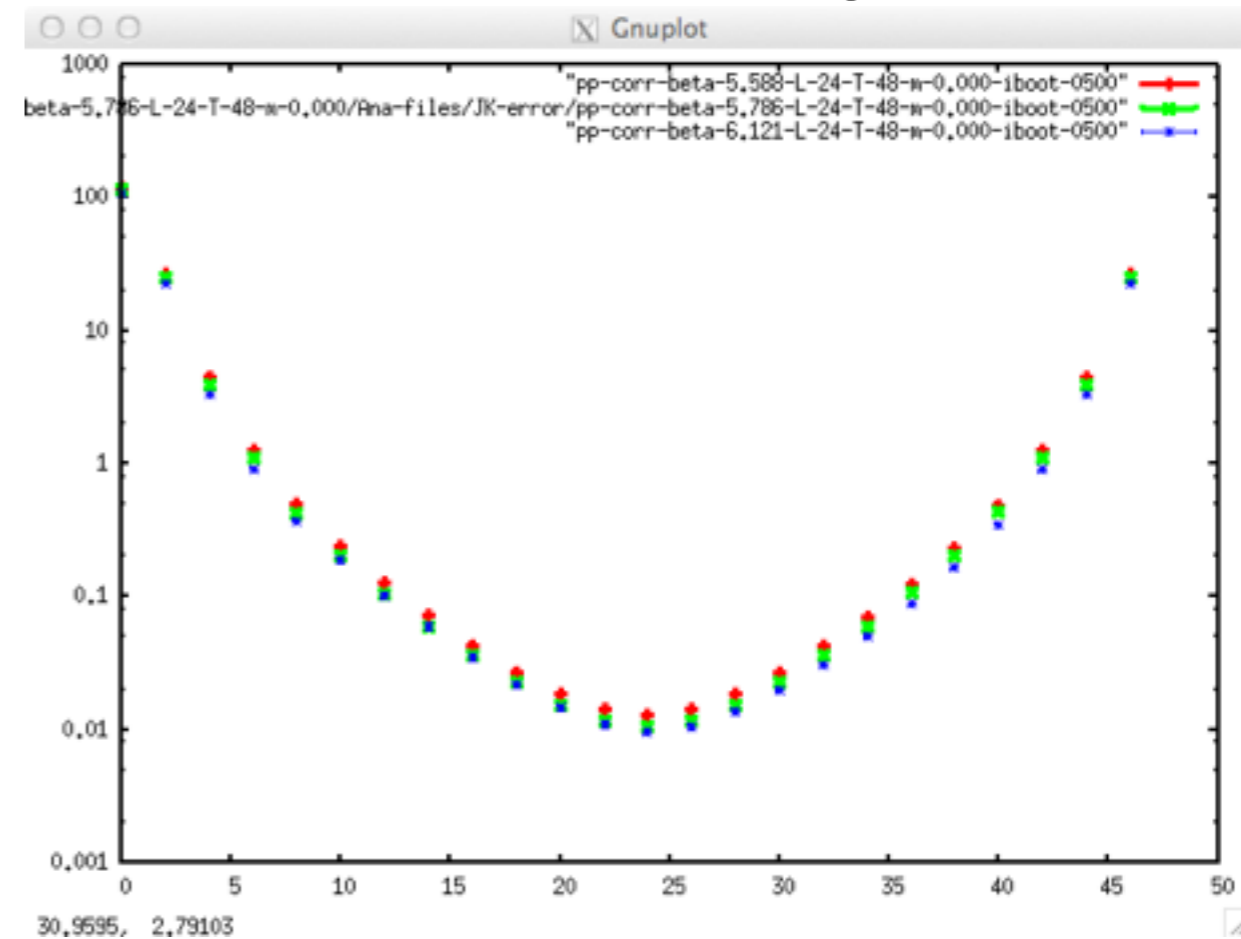
# Correlation fn. of the fluctuation

massless Nf=12 at IRFP

Trace part of the EMT  
vertical axis: linear scale



pseudo scalar op.  
vertical axis: log scale



Conservation law of energy-momentum tensor?

$$\partial_\tau T_\mu^\mu = 0$$

# Conclusion

## (EMT part)

- ◆ Novel method to obtain EMT using the lattice simulation
- ◆ quenched results (1pt.fn) show that the small flow time expansion is promising
- ◆ clear statistical signal, small systematic error
- ◆ 2pt. fn. calculation is also doable!!

## Future directions

- ◆ two-point function of EMT (shear and bulk viscosity, heat capacity)
- ◆ include dynamical fermion fields (need the calculation of Z-factor)
- ◆ application to conformal field theory (central charge, dilation physics)
- ◆ supersymmetry on the lattice