

What aspects of initial stages will NOT be tractable

(or will, but without sufficient precision)

with current or foreseeable theoretical tools?

Derek Teaney  
Stony Brook University



Stony Brook University

Thanks to: A. Kurkela, U. Heinz, A. Mazeliauskas, J. Jia, P. Steinberg, J. F. Paquet, and S. Schlichting

What aspects of initial stages are ALMOST  
tractable with current theoretical tools?

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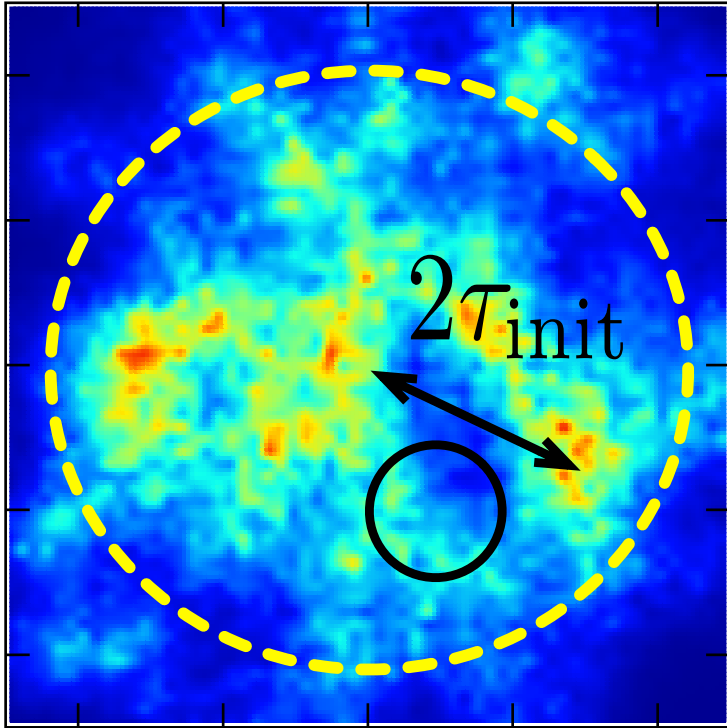


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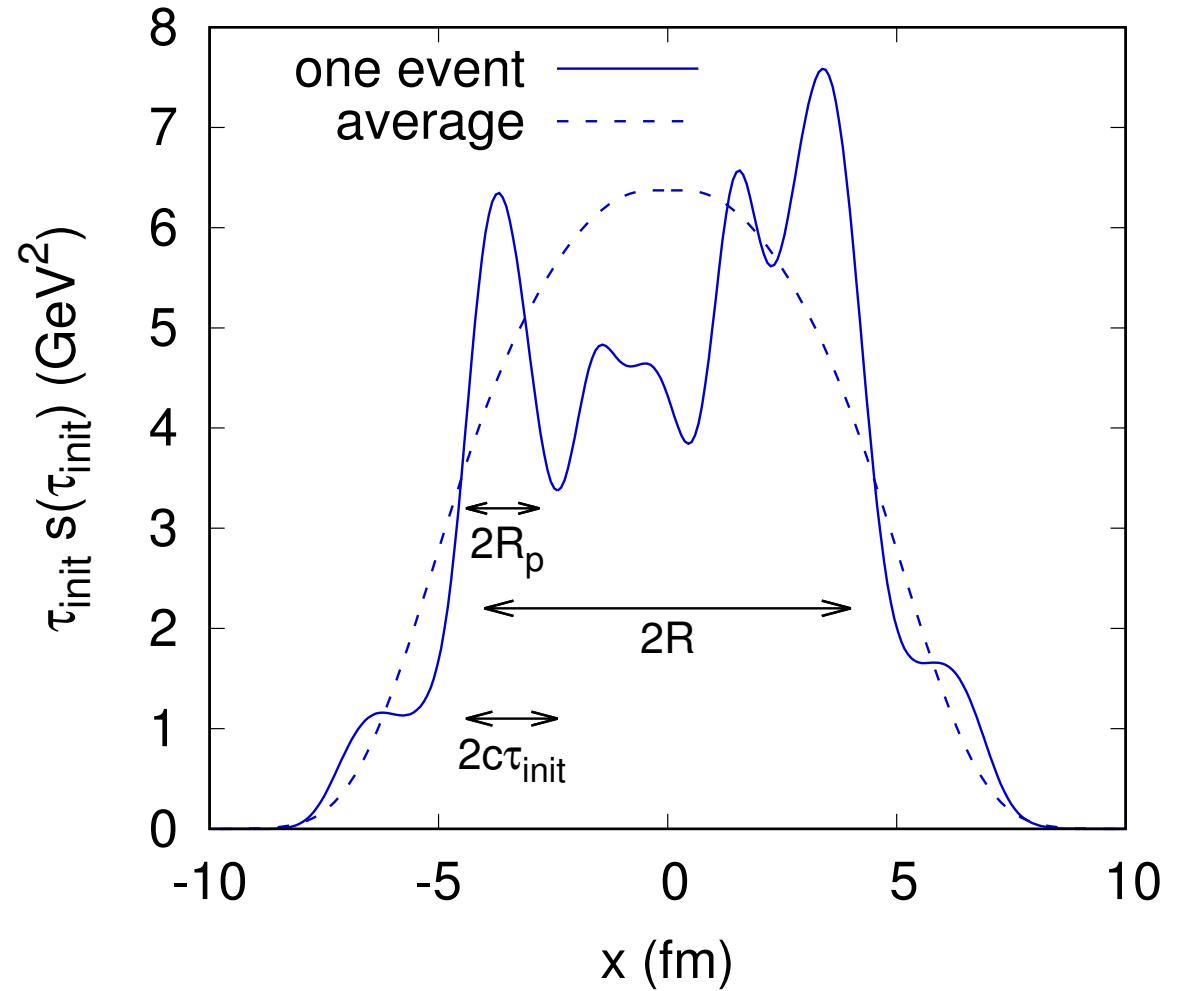
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## Spatial scales for initial stages in A+A

### IP Glasma A+A



### Glauber

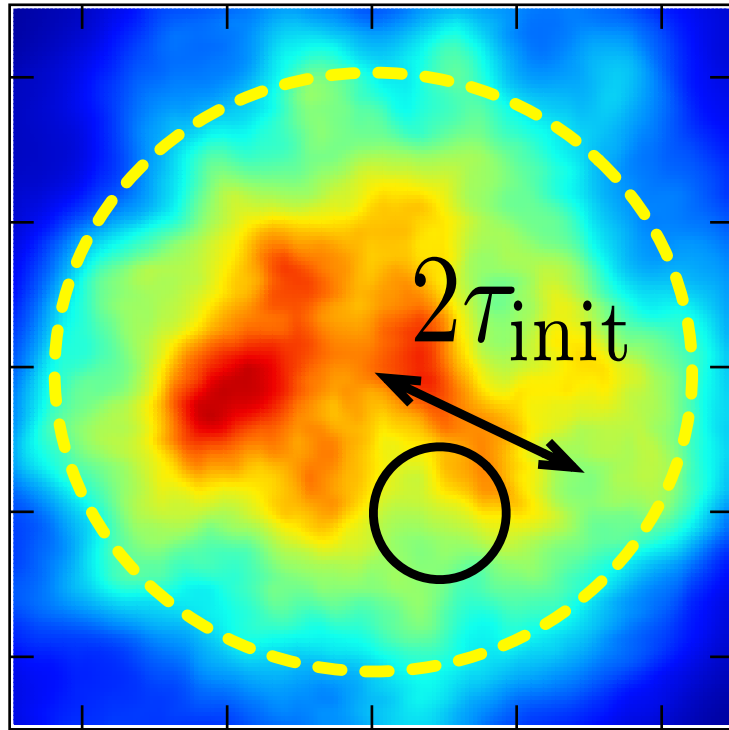


Many scales are the same:

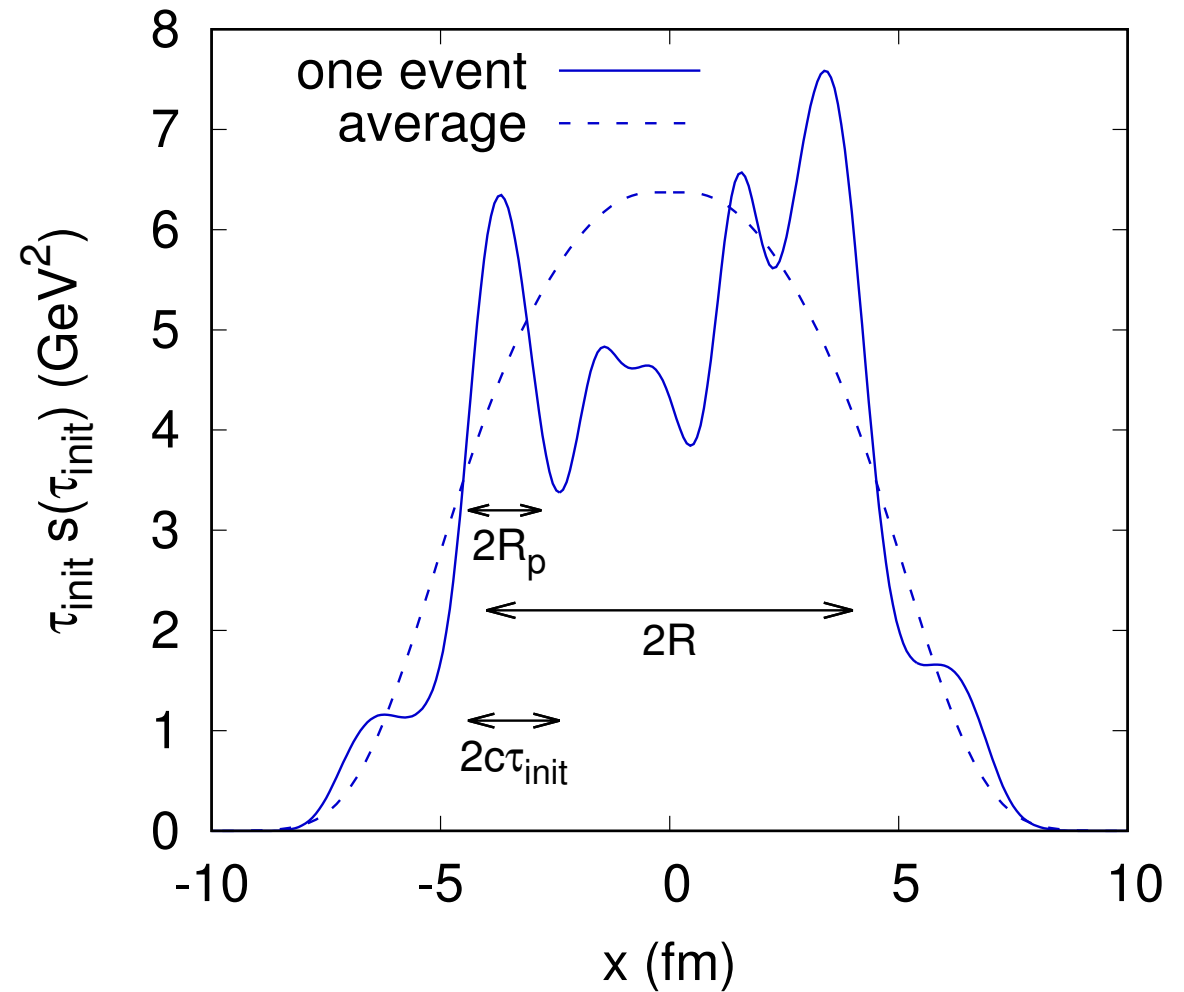
$$R \gg c\tau_{\text{init}} \sim R_p \sim \ell_{\text{mfp}} \gg \frac{1}{Q_s}$$

Need kinetics !

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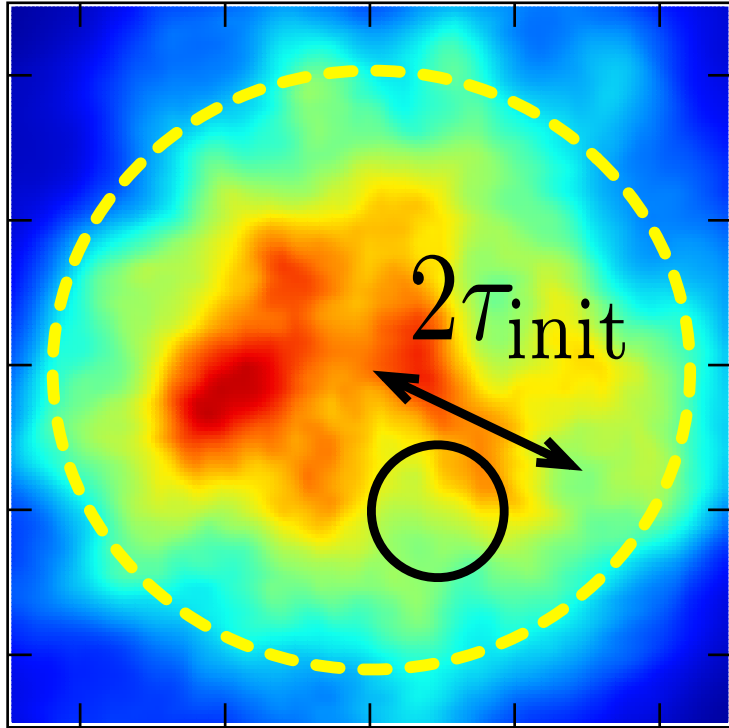
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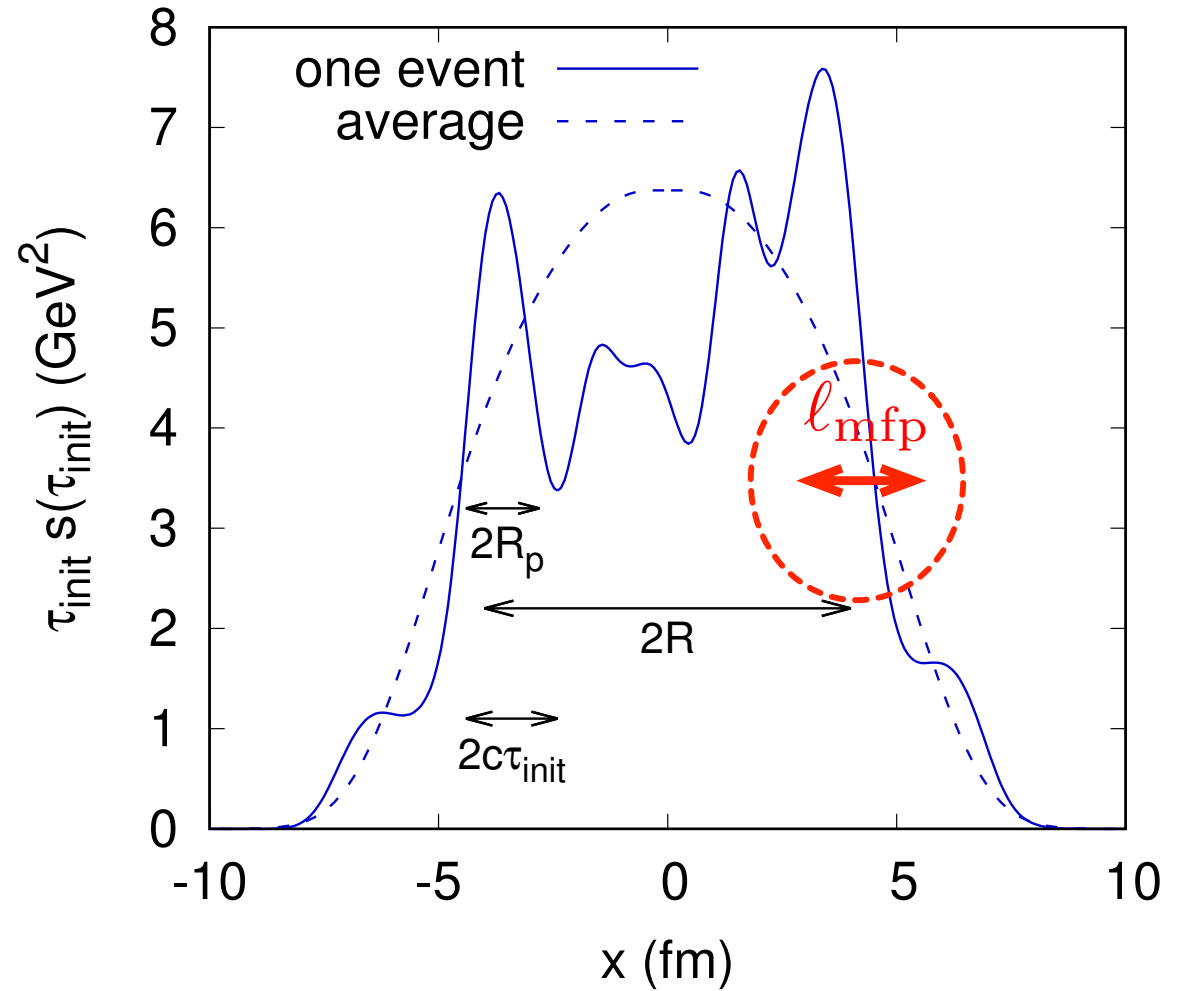


## Spatial scales for initial stages in A+A

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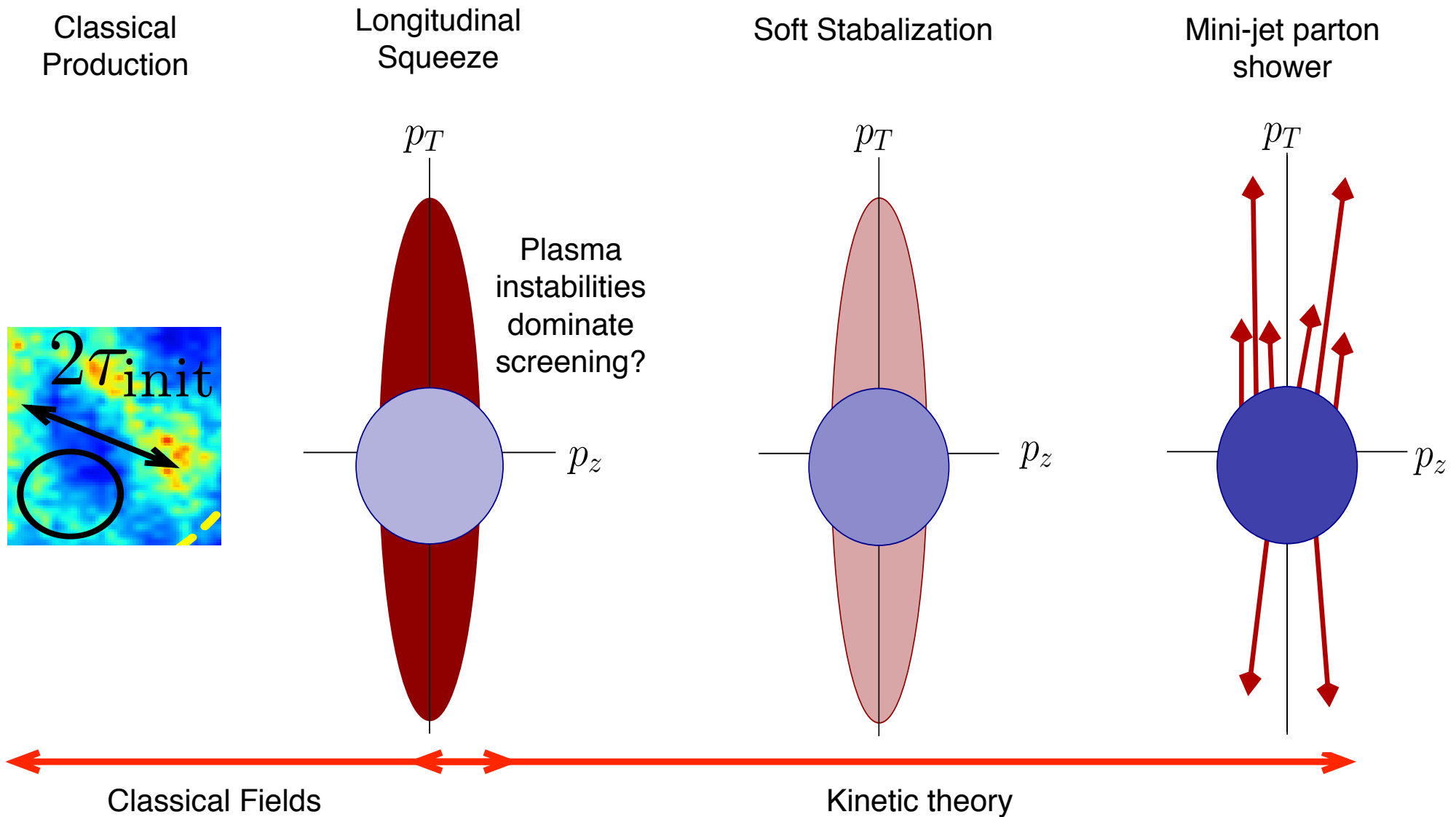
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# Momentum scales and “bottom-up” thermalization

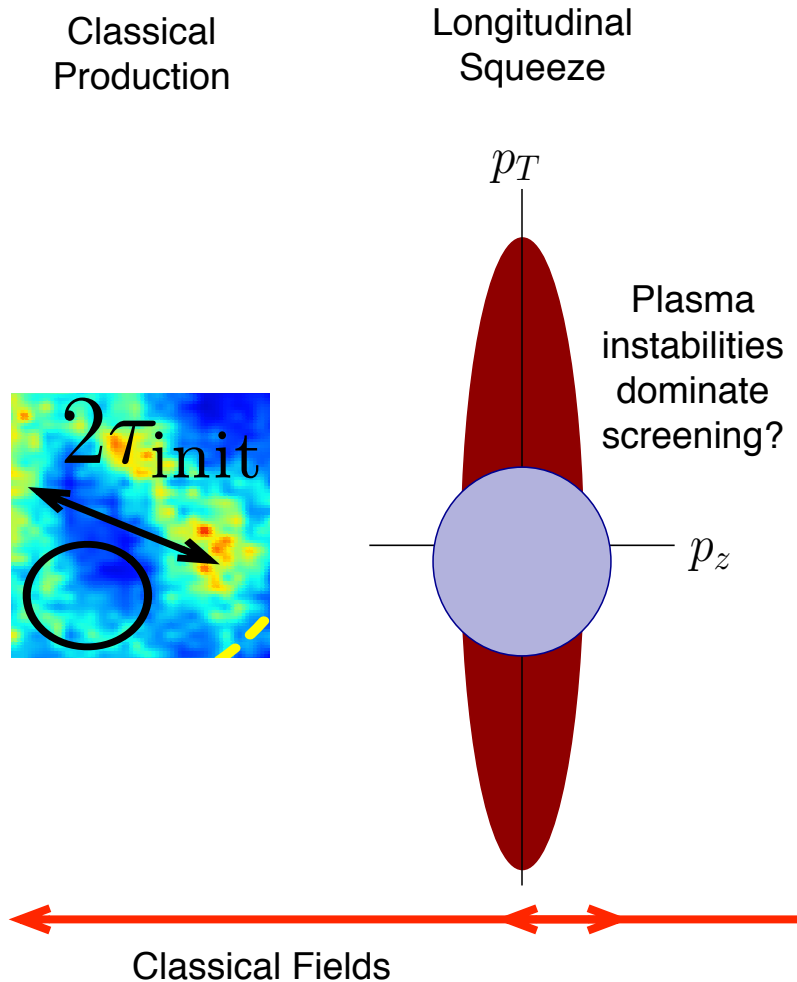
Baier, Mueller, Schiff, Son



Reach a thermal state in  $\tau_{\text{init}} \sim 1/(\alpha_s^{2.6} Q_s)$

## The classical phase of “bottom-Up” in 2014

Berges, Boguslavski, Schlichting, Venugopalan



### Recent progress on the first phase of “bottom-up”

- Scaling solution for phase space distribution

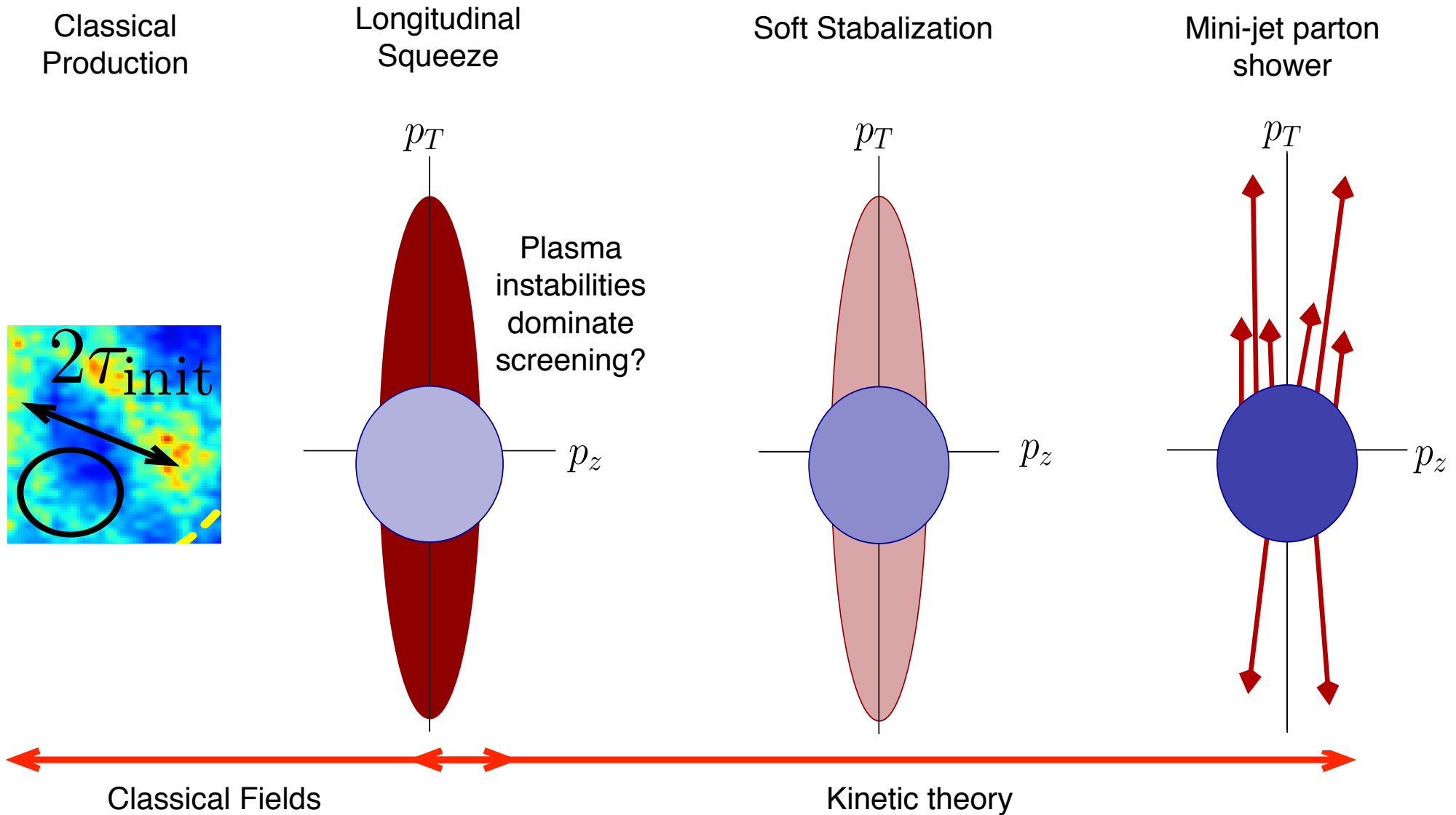
$$f(\tau, p_z, p_T) = \frac{(Q\tau)^{1/3}}{\tau} f_o(p_z(Q\tau)^{1/3}, p_T)$$

- Need to compare classical and kinetics in detail
- Instabilities do not seem to play a significant role

Reach the end of the first phase at  $\tau \sim 1/(\alpha_s^{3/2} Q_s)$

# Momentum scales and “bottom-up” thermalization

Baier, Mueller, Schiff, Son

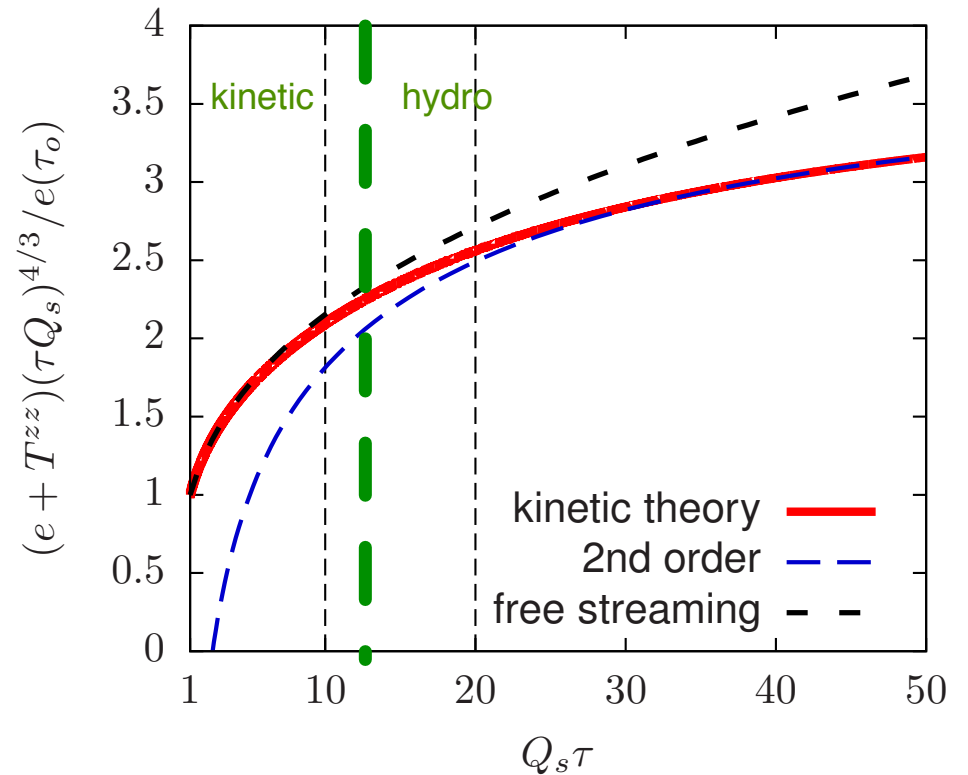
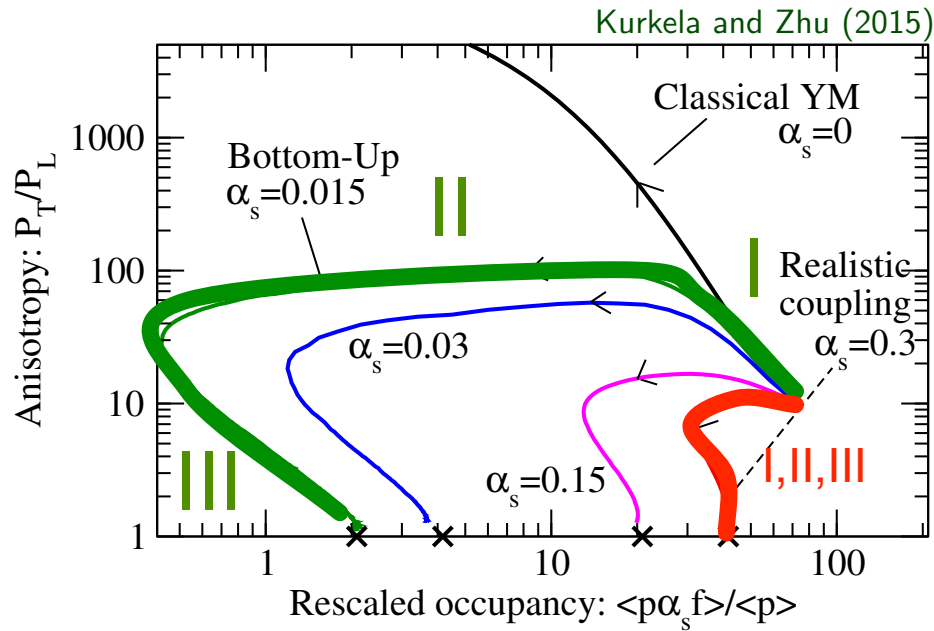


Reach a thermal state in  $\tau_{\text{init}} \sim 1/(\alpha_s^{2.6} Q_s)$

## The kinetic phase of “bottom-Up” in 2015

Kurkela, Zhu

- Compare the longitudinal pressure  $T^{zz}$  with hydro prediction to calibrate equilibration

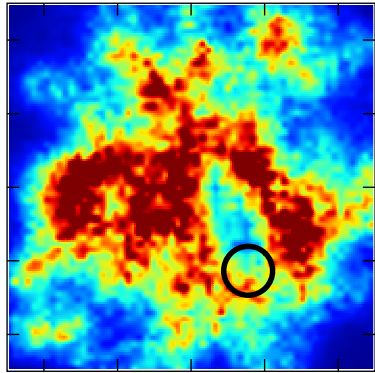


For realistic coupling three stages still “exists” and

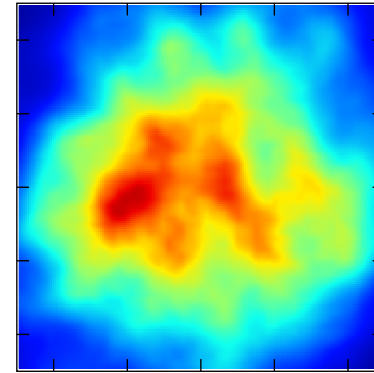
$$\tau_{\text{init}} \simeq \frac{10}{Q_s} \left( \frac{0.3}{\alpha_s} \right)^{2.6}$$

constitutive relations satisfied to 10%

$$\int d^2\mathbf{x}' \underbrace{\frac{\delta e(\tau_0, \mathbf{x}')}{e(\tau_0)}, \frac{\vec{\nabla} e(\tau_0, \mathbf{x}')}{e(\tau_0)}}_{\text{ip-glasma}} \times \underbrace{E(|\mathbf{x} - \mathbf{x}'|; \tau, \tau_0)}_{\text{Green fcn}} = \underbrace{\frac{\delta e(\tau, \mathbf{x})}{e(\tau)}, \frac{\vec{g}(\tau, \mathbf{x})}{e(\tau)}}_{\text{hydro}}$$



Green function



1. Assumes that  $\tau_{\text{init}} \ll R$
2. Matches the early glasma kinetics, to an effective theory of hydro initial states Blaizot, Ollitrault

$$\langle e(\mathbf{x}) \rangle \quad \langle e(\mathbf{x})e(\mathbf{y}) \rangle = A \langle e(\mathbf{x}) \rangle \delta^2(\mathbf{x} - \mathbf{y})$$

3. Include momentum into the hydro initial state effective theory at first order in  $\tau_{\text{init}}/R$

When the pre-equilibrium response is in, the hydro results will (hopefully) not depend at  $\tau_{\text{init}}$

Can we measure this ?

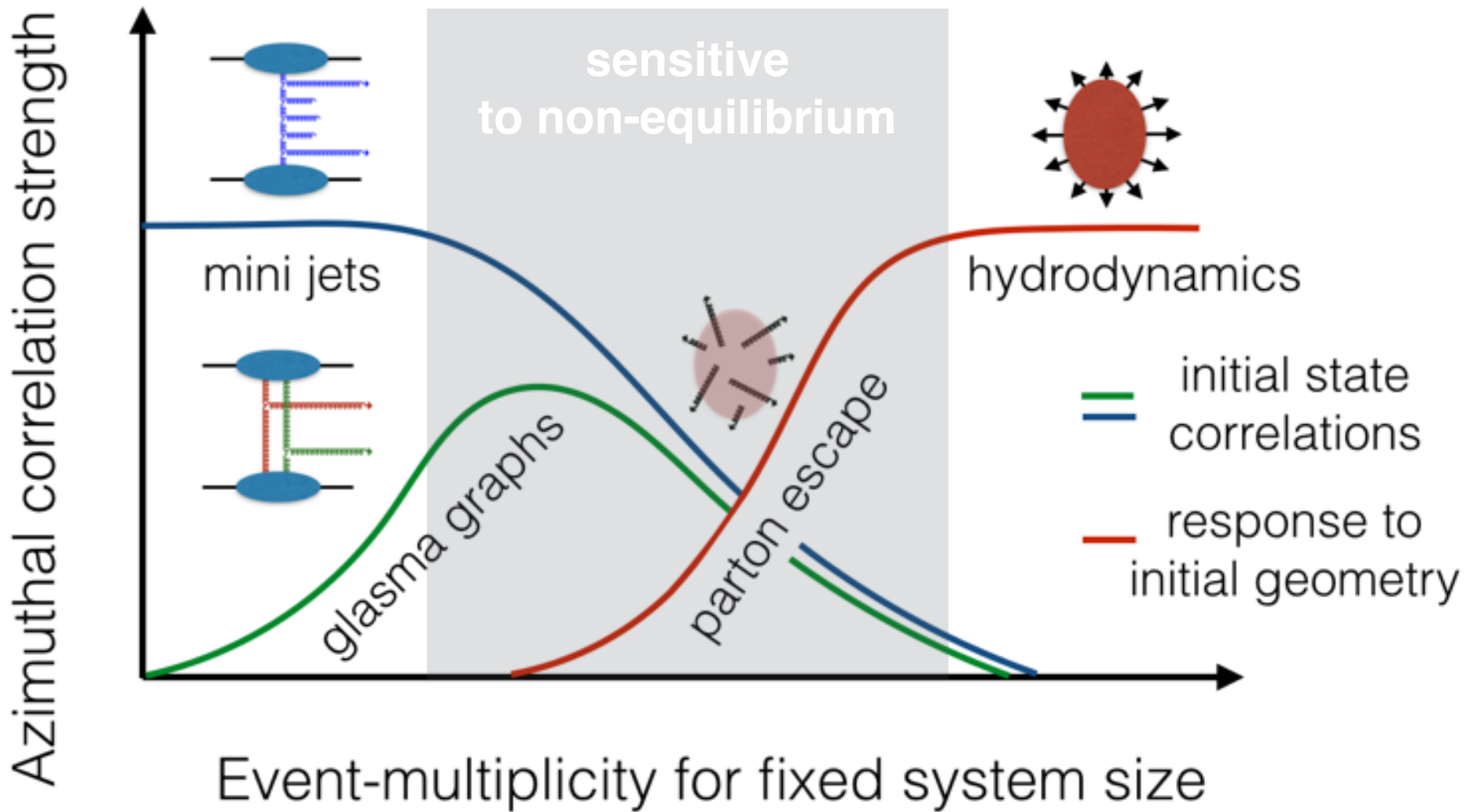
Need control parameters – system size,  $p_T$ , rapidity, . . .

Can we measure this ?

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## Sören Schlichting's dream plot



Let's put some numbers on the  $x$  axis . . .

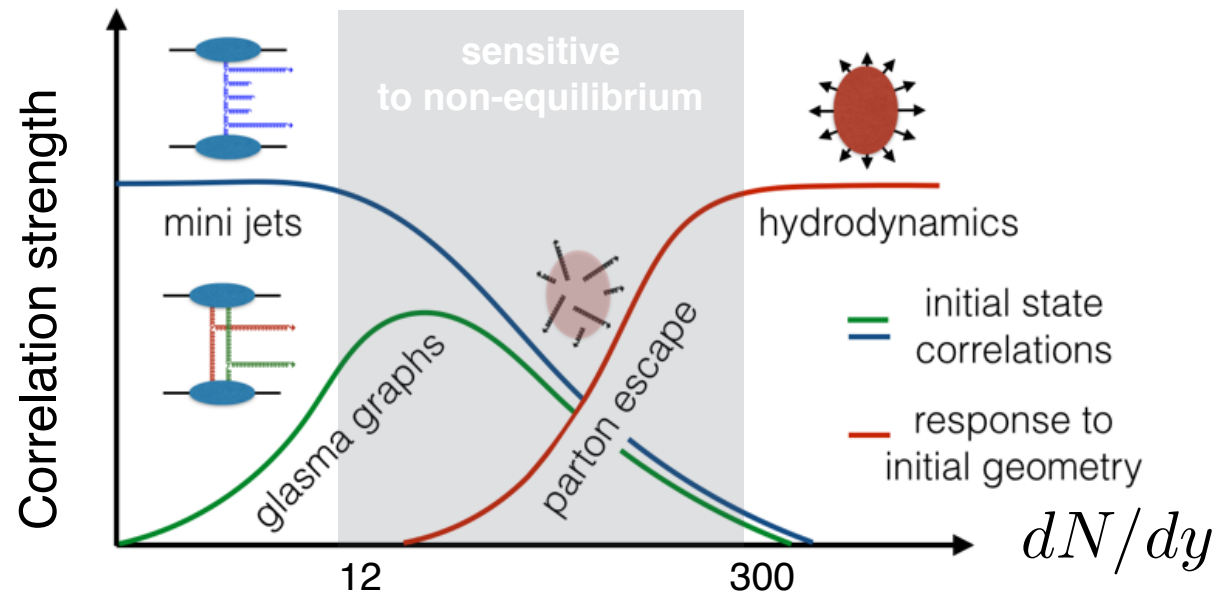
## When do you see initial state correlations ?

1. To see the initial state correlations we require

$$R \ll \tau_{\text{init}} \quad \text{or} \quad R/\tau_{\text{init}} \ll 1$$

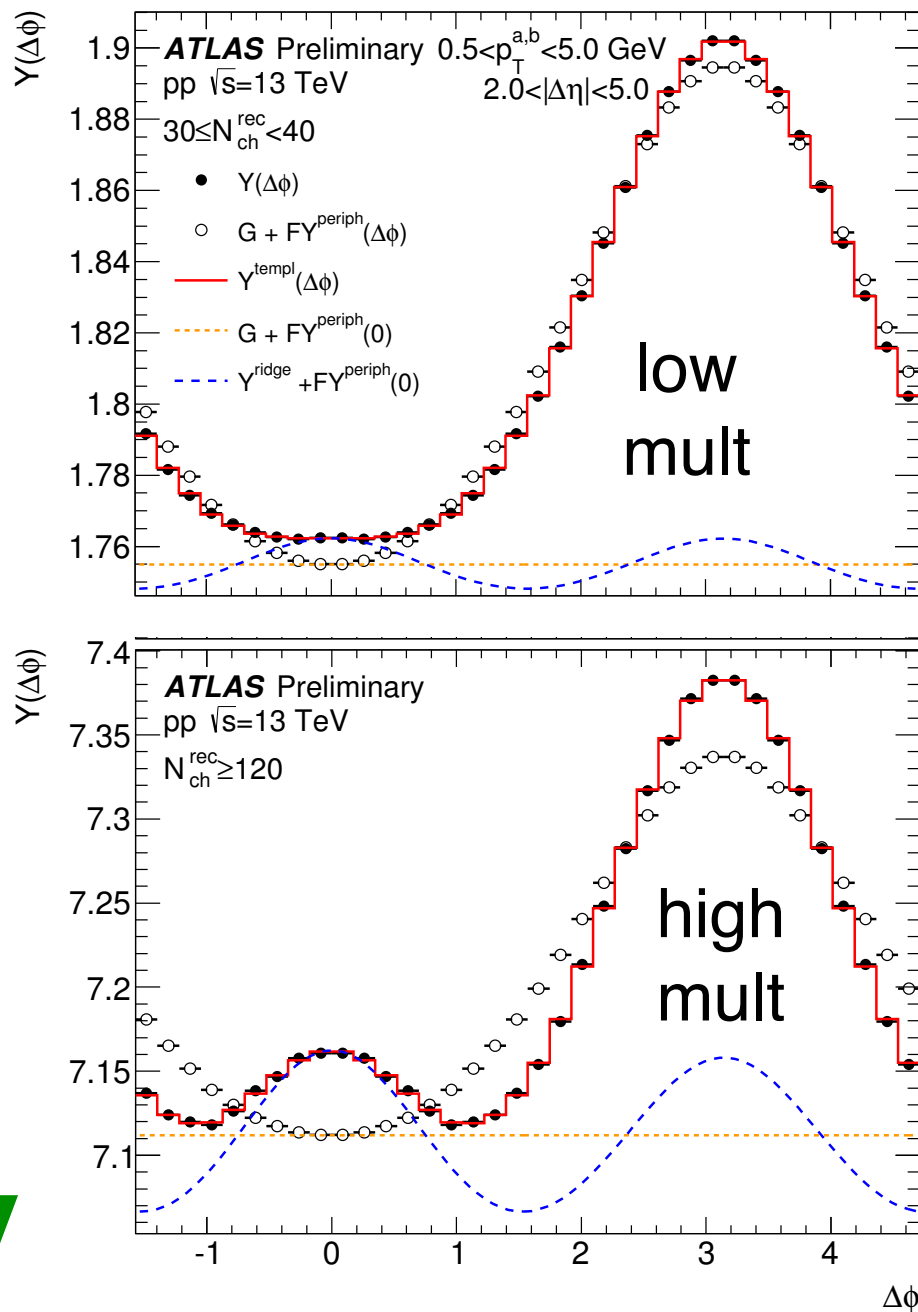
2. Now we reach a thermal state in:

$$\tau_{\text{init}} \simeq \frac{10}{Q_s} \left( \frac{0.3}{\alpha_s} \right)^{2.6} \quad \text{and} \quad Q_s^2 = \alpha_s \frac{dN/dy}{\pi R^2}$$



$$\frac{R}{\tau_{\text{init}}} = 0.2 \left( \frac{\alpha_s}{0.3} \right)^{3.1} \left( \frac{dN/dy}{12} \right)^{1/2}$$

# 13 TeV pp



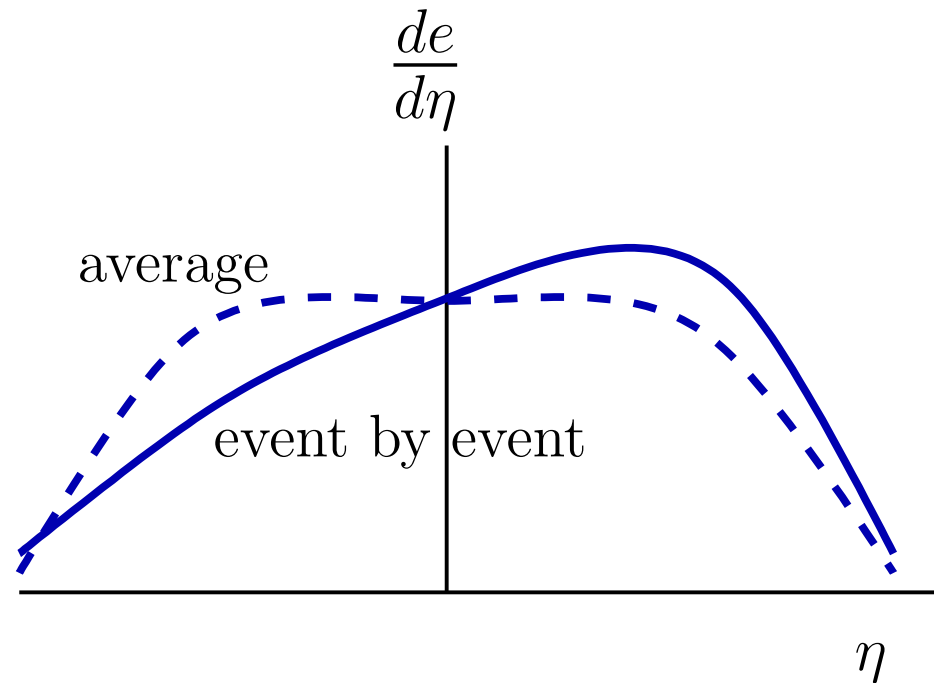
## Remarks on $pp$ :

1. Gradual transition from low to high mult.
2. Is  $v_2$  at small multiplicity is important?  
 - See CMS data on  $v_2\{4\}$ ,  $v_2\{6\}$ , ...
3. The  $v_2$  correlation does not decrease with multiplicity, and factorizes
4. How are minijets modified in  $p_T$  ?

The last stage of thermalization is  
 minijet-quenching

Can we measure this ?

Need control parameters – system size,  $p_T$ , rapidity, . . .

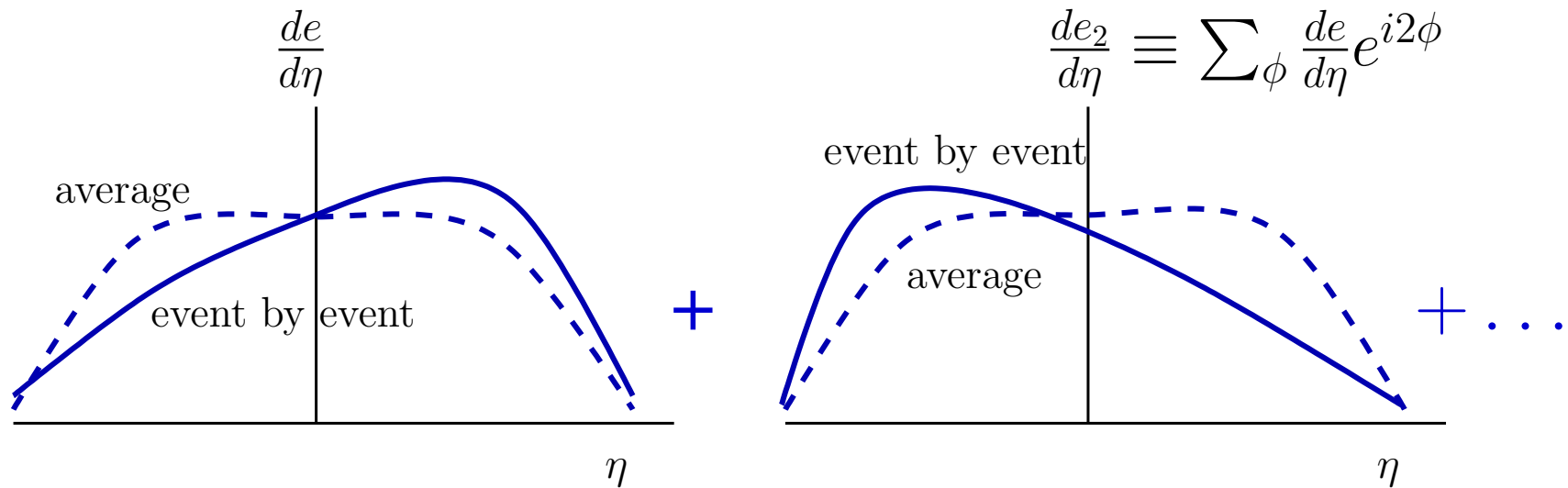


## 1. An effective theory for hydro initial states must specify

Bozek, Broniowski arxiv:1512.01945

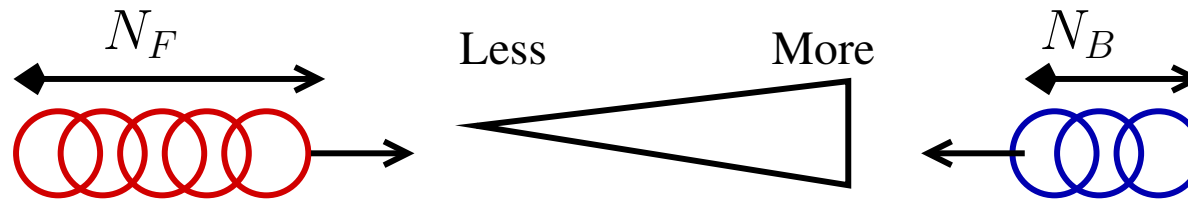
$$\left\langle \frac{de(\mathbf{x})}{d\eta} \right\rangle \quad \text{and} \quad \left\langle \frac{de(\mathbf{x})}{d\eta_1} \frac{de(\mathbf{y})}{d\eta_2} \right\rangle_{\text{conn}} = \underbrace{A(\eta_1, \eta_2)}_{\text{unknown}} \langle e(\mathbf{x}) \rangle \delta(\mathbf{x} - \mathbf{y})$$

## Rapidity correlations



A model for rapidity fluctuations that doesn't work:

Teaney&Bzdak



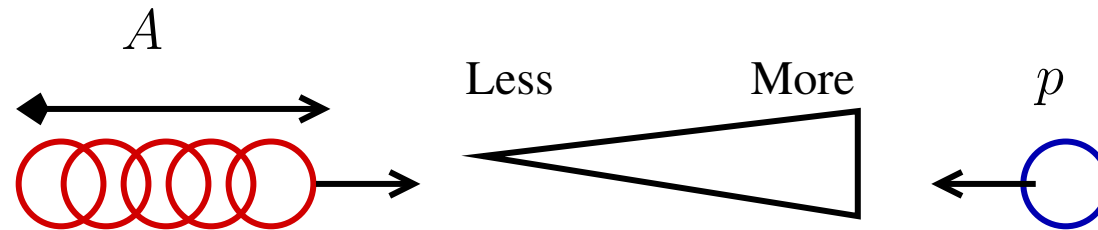
then

$$\frac{de}{d\eta} = C_0 (N_F + N_B) + C_1 \eta (N_F - N_B)$$

and

$$\left\langle \frac{de}{d\eta_1} \frac{de}{d\eta_2} \right\rangle = C_0^2 \langle (N_F + N_B)^2 \rangle + \underbrace{C_1^2 \eta_1 \eta_2 \langle (N_F - N_B)^2 \rangle}_{\text{measurable rapidity odd correlations}}$$

1. In  $pA$  events the number of forward backward participants doesn't fluctuate much



Fluctuations reflect the correlations  
in the tube, not the participants

2. But, significant (rapidity-odd) longitudinal fluctuations of  $v_2$  are seen

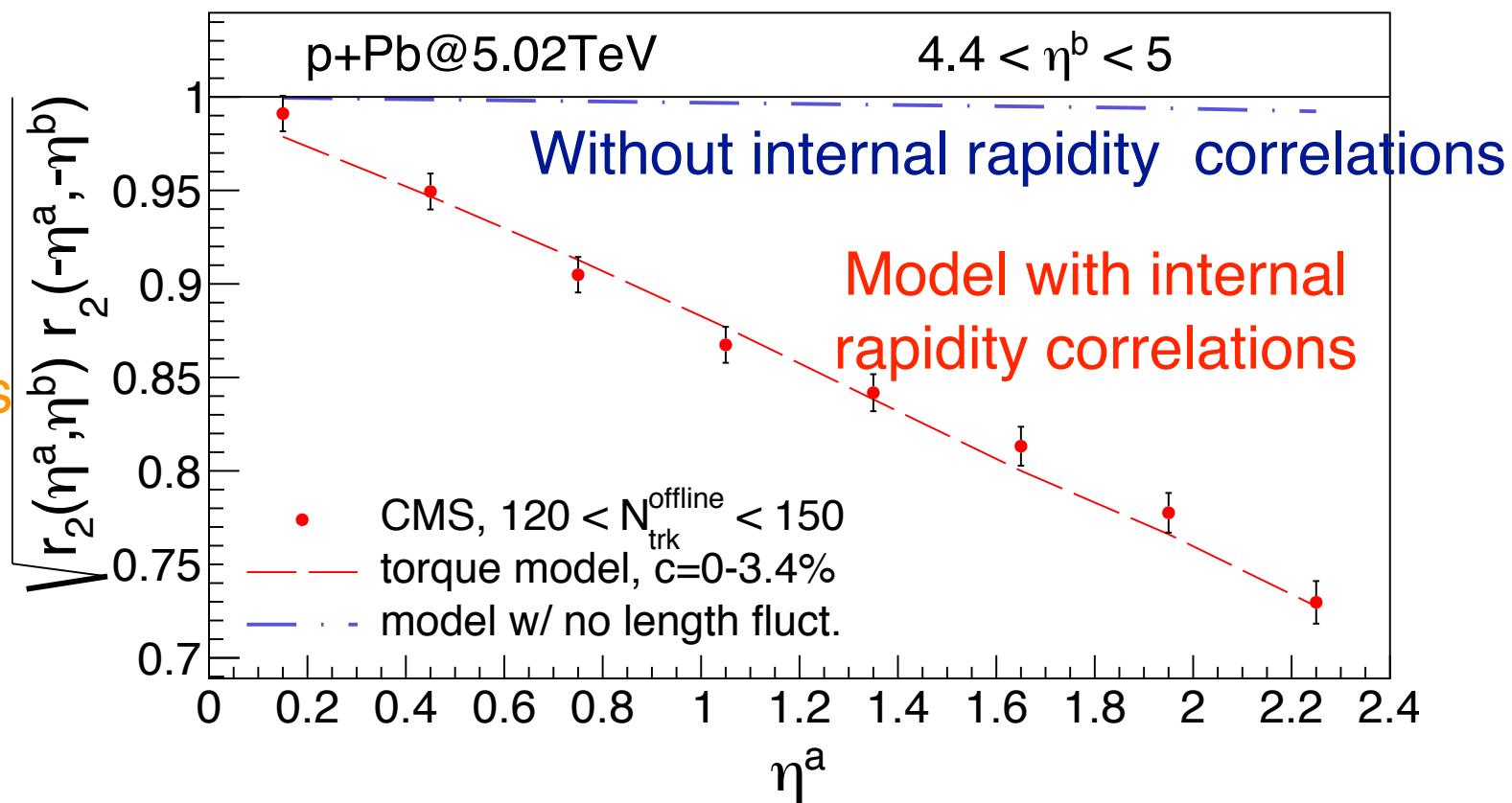
Successful string-like models incorporate the  $dx/x$  parts of the splitting function  
into the rapidity correlation function  $\langle e(\eta_1)e(\eta_2) \rangle$

3. Opportunity to predict the functional form in CGC or EKRT model of

Schenke, Schlichting

$$\left\langle \frac{de}{d\eta_1} \frac{de}{d\eta_2} \right\rangle_{\text{conn}}$$

Rapidity  
odd  
v2 fluctutations





From the almost to truly intractable . . .

1. A path to document theoretically and experimentally the path to hydro
  - System size,  $p_T$  mini-jet quenching, rapidity
2. Going beyond weak coupling:
  - Non-perturbatively calculate screening?
  - Is the temperature scale perturbative?

Thank you!