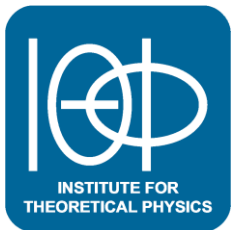

Colored particle-in-cell simulations for heavy-ion collisions

Initial Stages 2016, Lisboa, Portugal
25.05.2016

David Müller

with Andreas Ipp and Daniil Gelfand

Institute for Theoretical Physics, Vienna University of Technology, Austria



Introduction

Goal: Simulate heavy-ion collisions in the color-glass-condensate (CGC) framework with finite nucleus thickness. Possible with colored particle-in-cell (CPIC).

Access to lower collision energies, break boost-invariance!

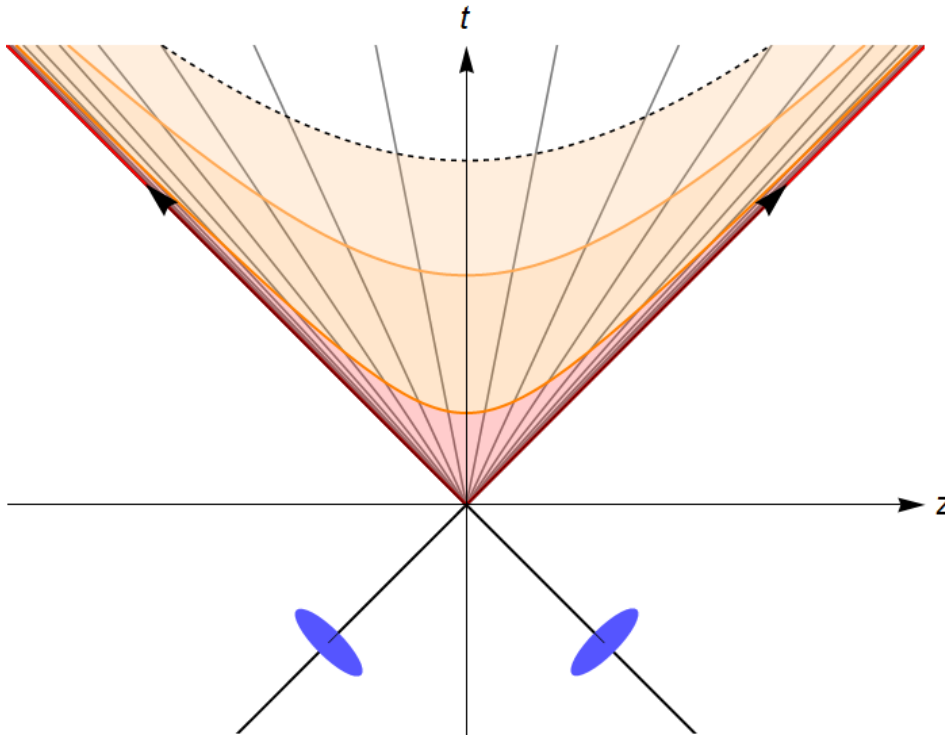
- Formation and evolution of the quark-gluon-plasma (QGP)?
- How does the QGP become isotropic and thermalized?
- **What is the role of boost-invariance?**

Various experiments with wide range of the gamma factor γ :

- LHC (ALICE) @ CERN: Pb+Pb with **~ 2.76 TeV** per nucleon pair. ($\gamma \approx 2700$)
- RHIC @ BNL: Au+Au with **~ 200 GeV** ($\gamma \approx 100$)
- RHIC beam energy scan: **$\sim 7.7 - 62.4$ GeV** ($\gamma \approx 4 - 30$)

Need to go beyond boost-invariant approximation. → Simulations with “thick” nuclei!

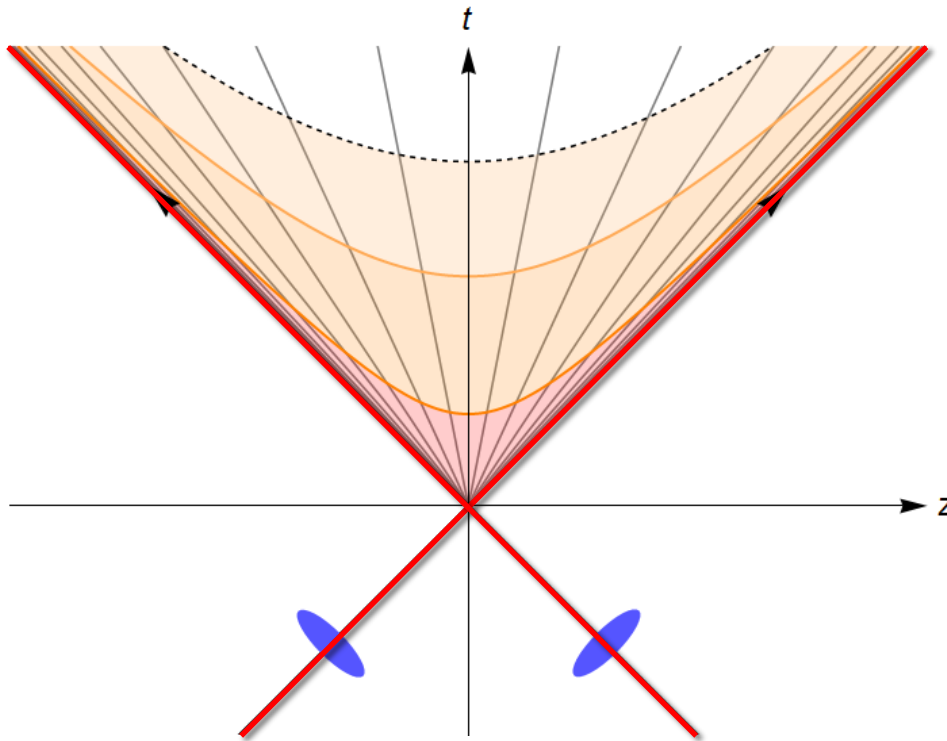
Boost-invariant CGC collision



- CGC: Separation of hard and soft degrees of freedom, weak coupling
- Color currents of the nuclei restricted to the light cone and infinitely thin
- Analytical solutions exist for everything except the forward light cone
- Fields in the forward light cone are independent of rapidity. Reduction from 3D+1 to 2D+1
- Need to solve 2D+1 source-free Yang-Mills equations in the forward light cone with special initial conditions on the light cone

$$D_\mu F^{\mu\nu}(\tau, x_T) = 0$$

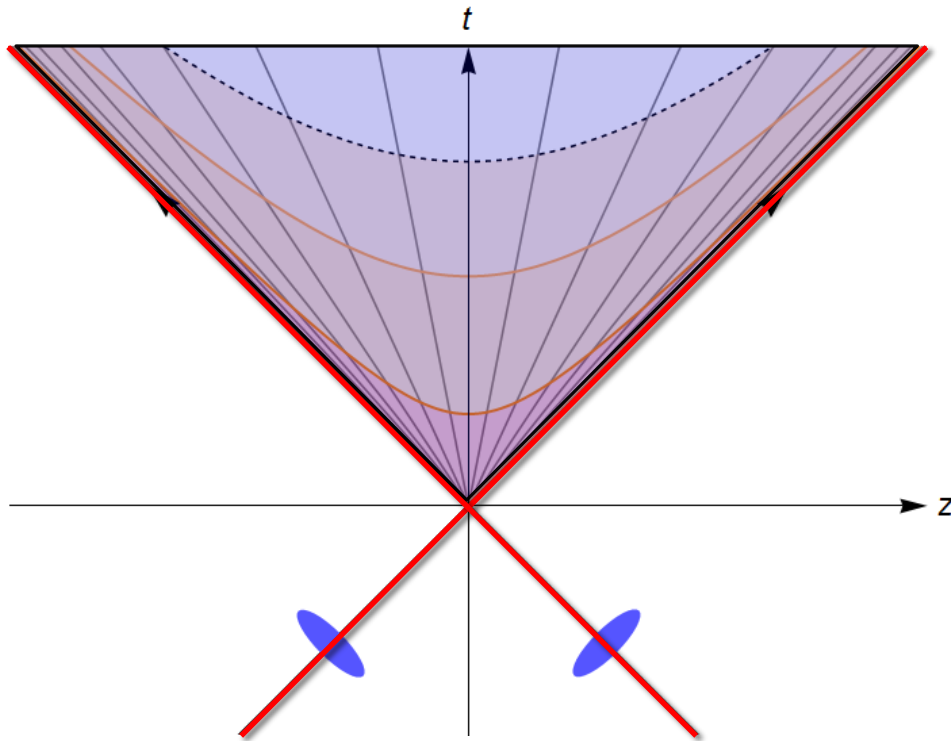
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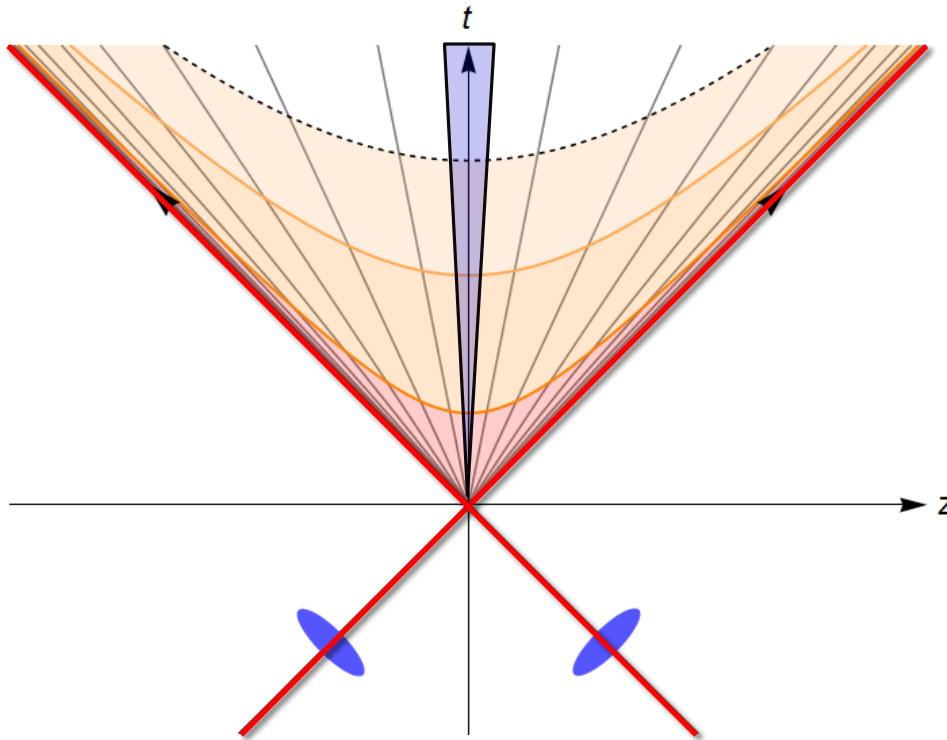
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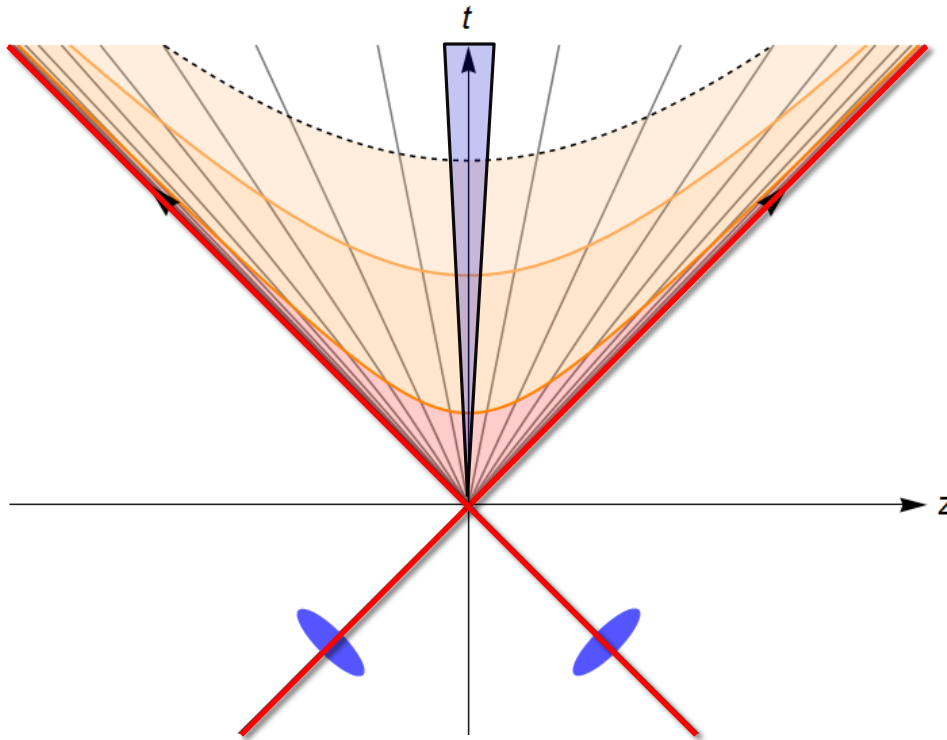
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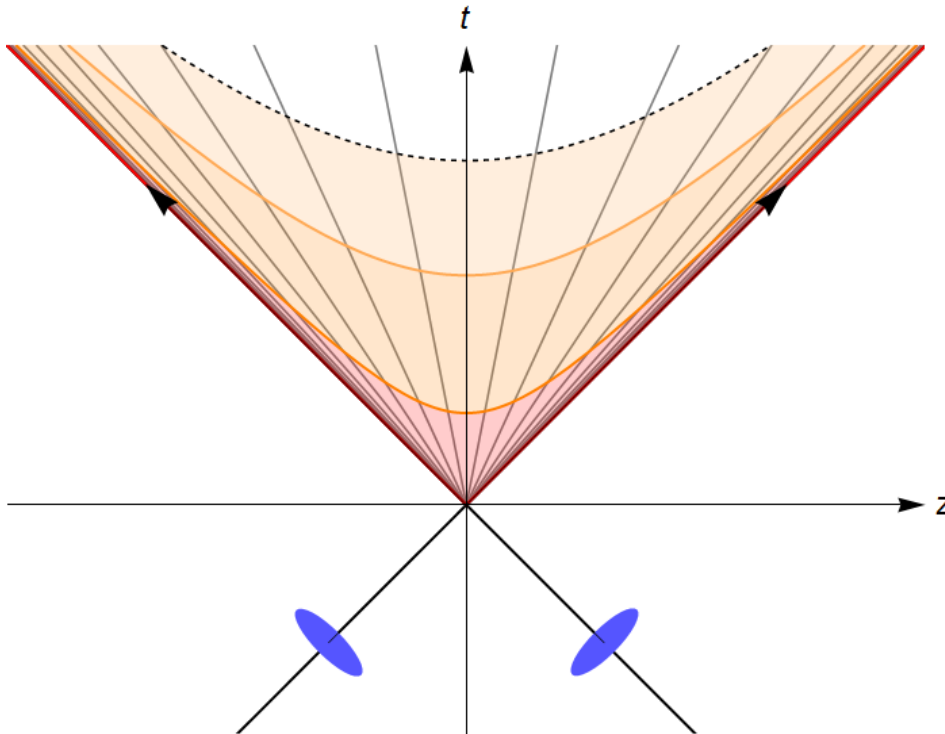
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Finite nucleus thickness

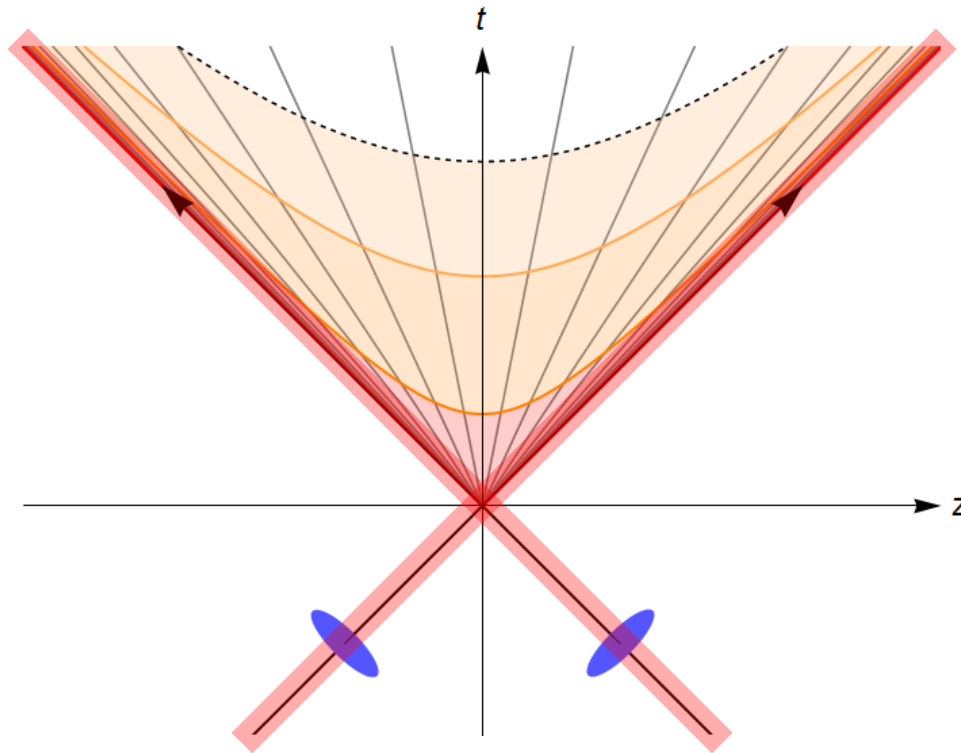


- Extended color currents need to be taken into account.
- Fields depend on rapidity.
- Need to solve full 3D+1 Yang-Mills equation with currents.

$$\begin{aligned} D_\mu F^{\mu\nu}(t, z, x_T) &= J^\nu \\ D_\mu J^\mu(t, z, x_T) &= 0 \end{aligned}$$

Colored particle-in-cell (CPIC) provides a framework to numerically solve the field and current equations on a lattice.

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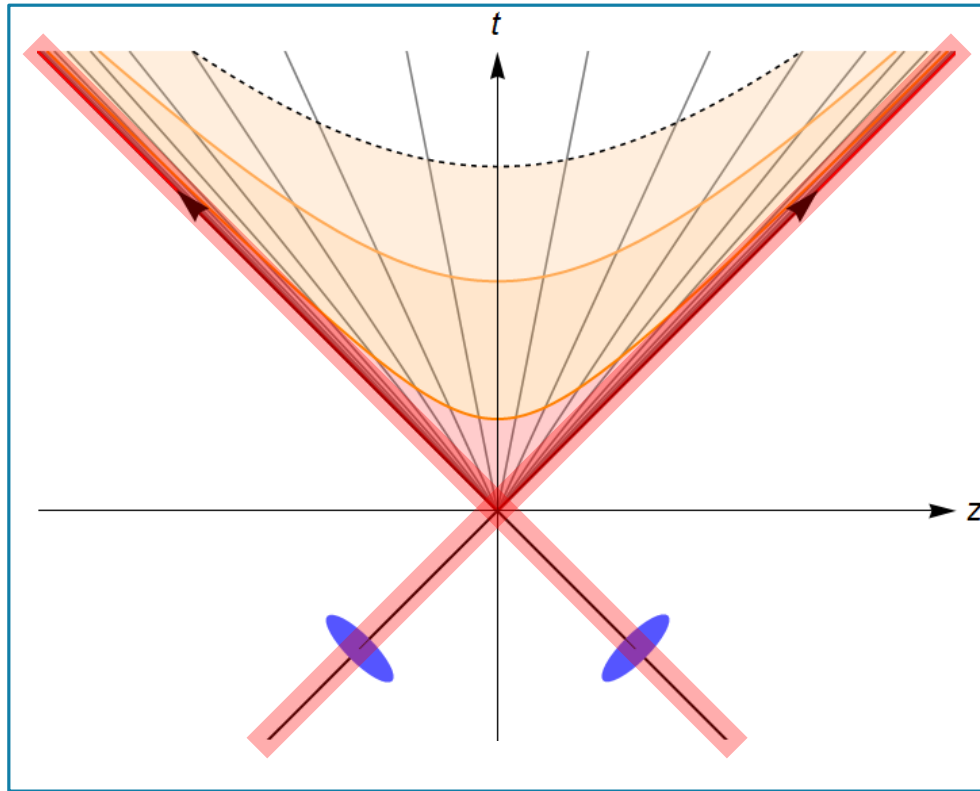


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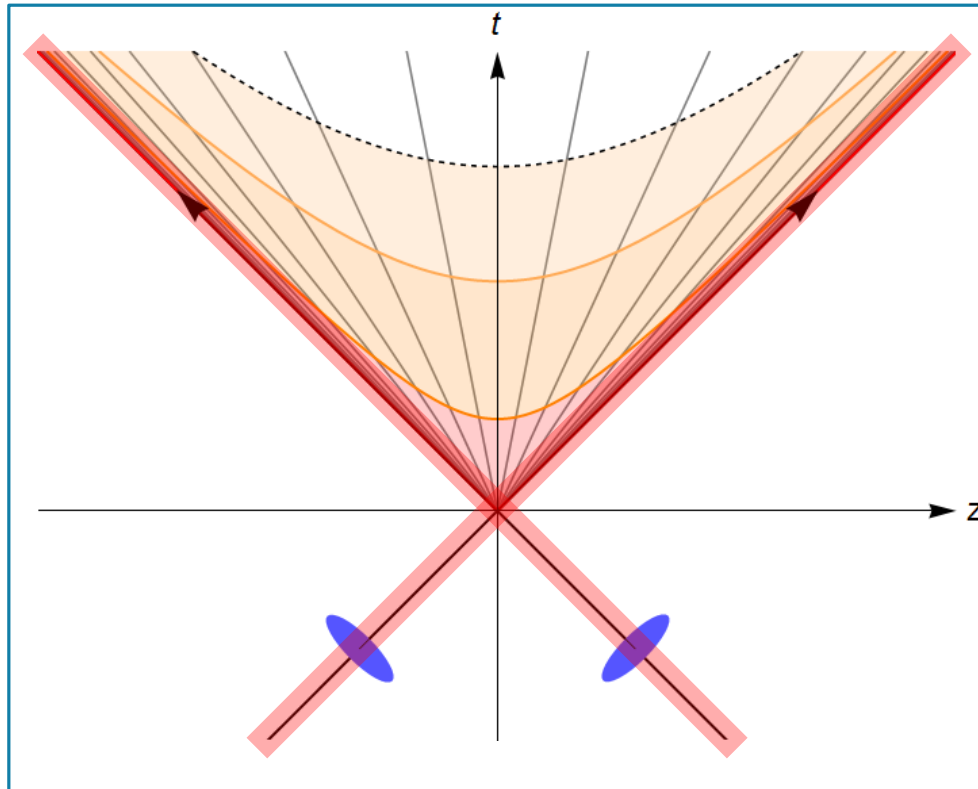


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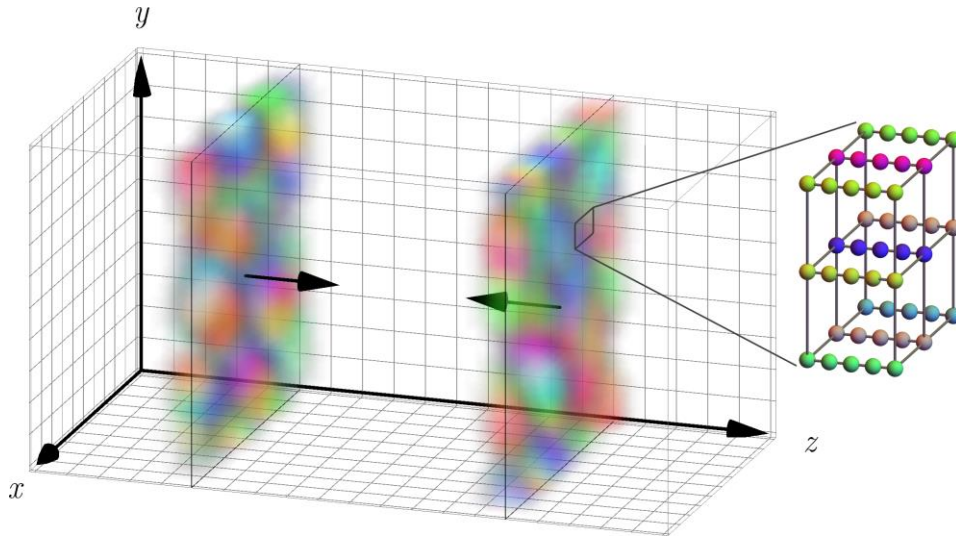


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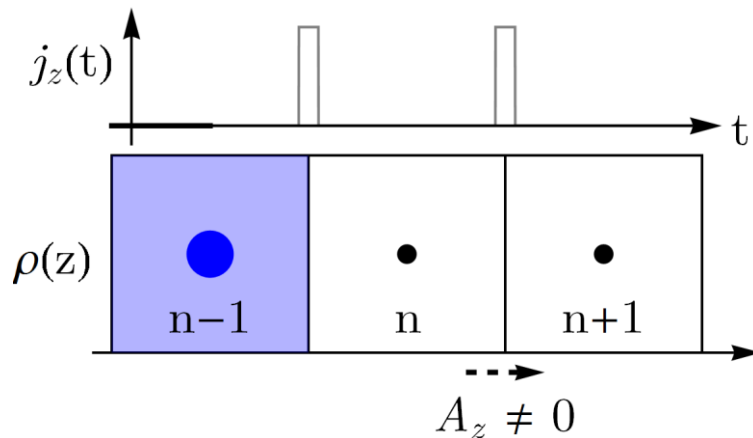
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[A. Dumitru, Y. Nara, M. Strickland: PRD75:025016 (2007)]

Collision of two nuclei in the lab frame:

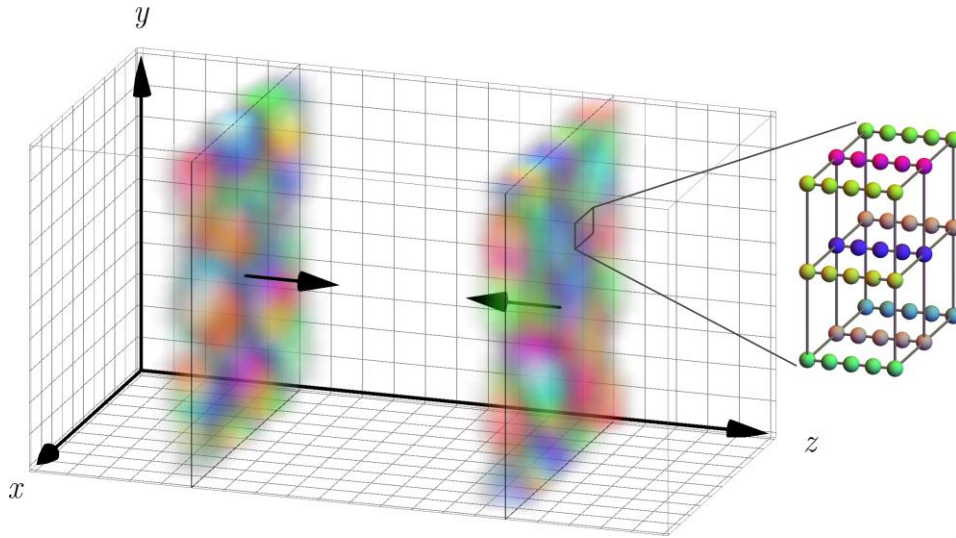
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Nearest-grid-point method (NGP)



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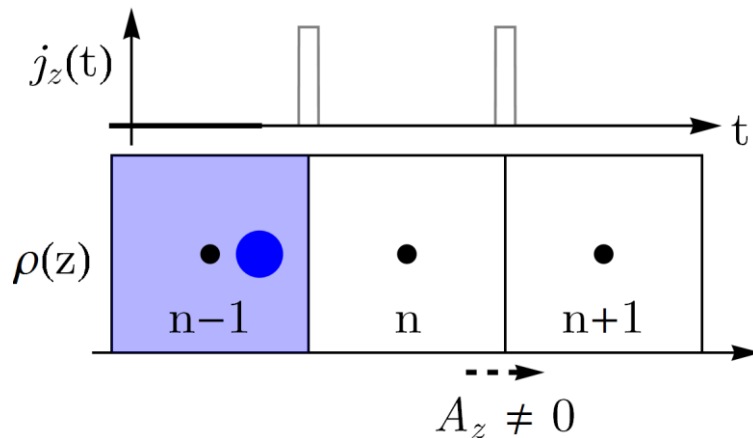
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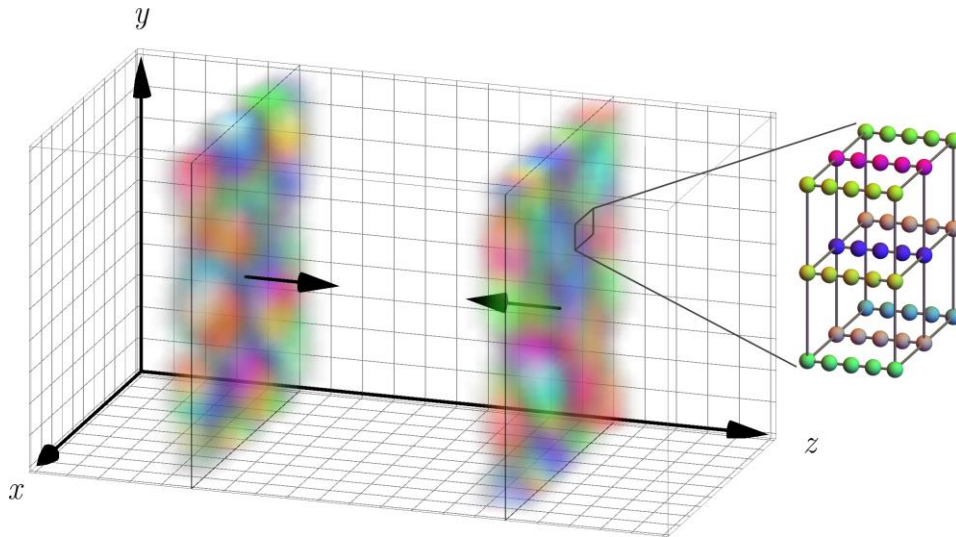
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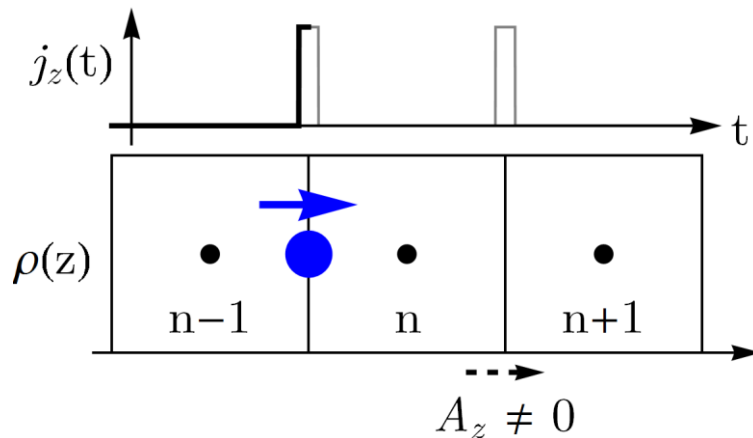
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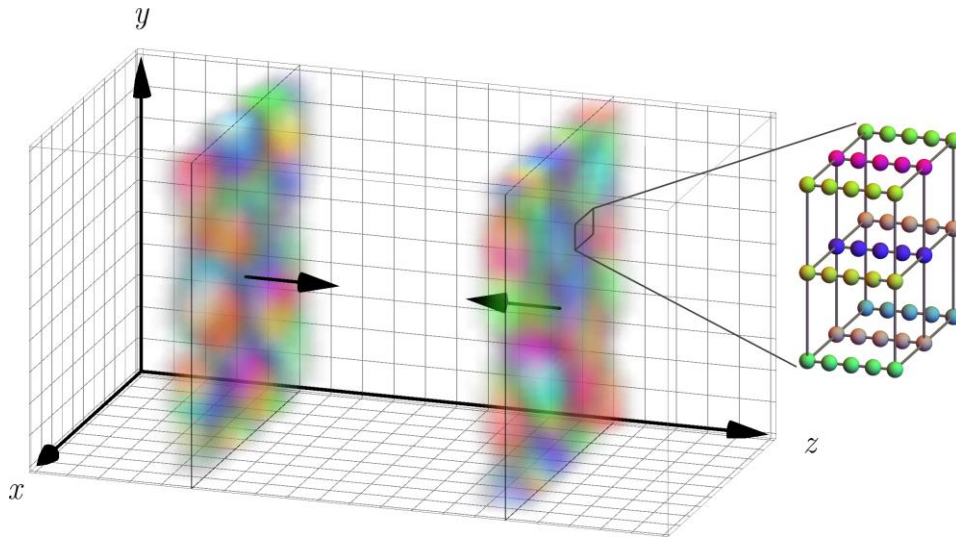
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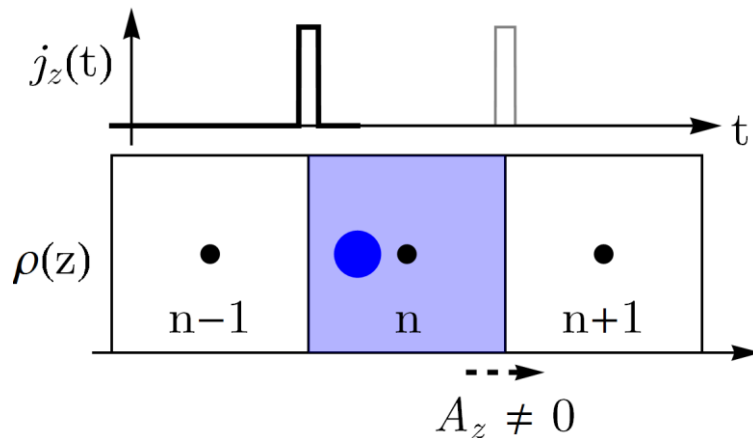
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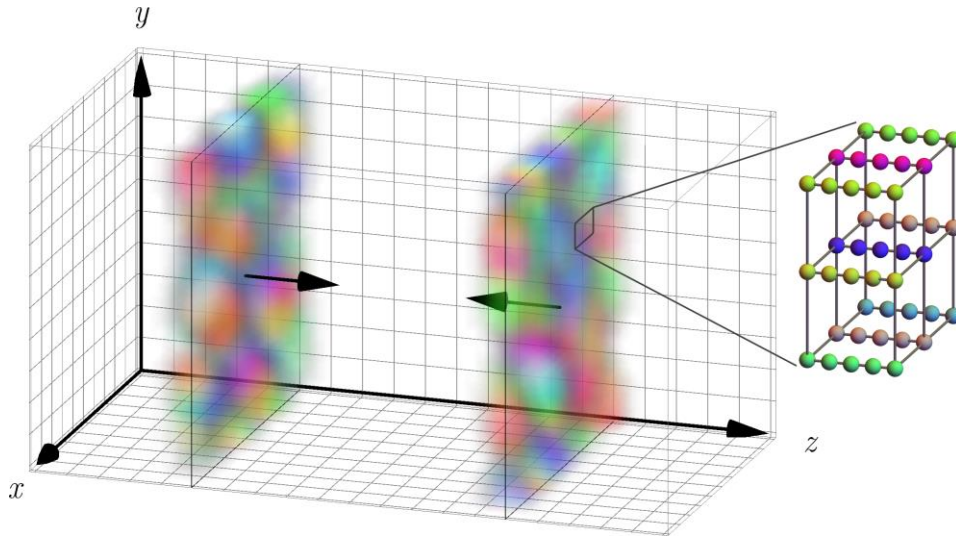
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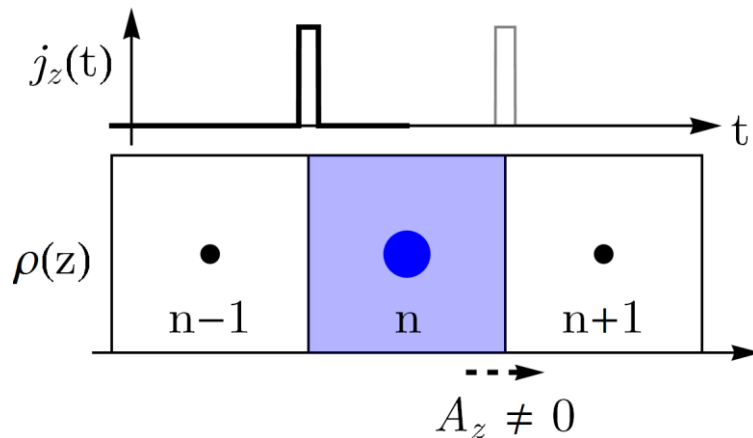
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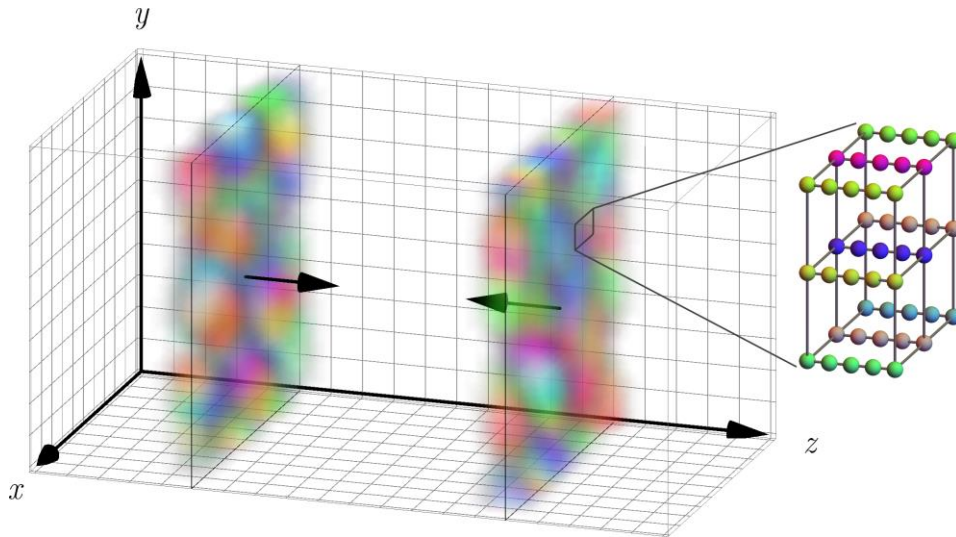
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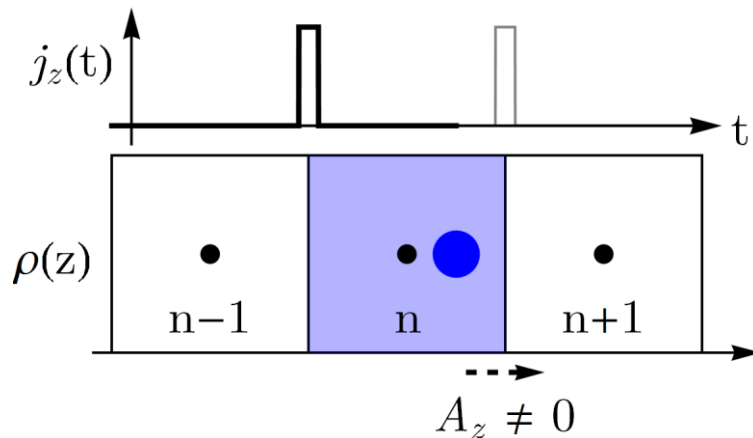
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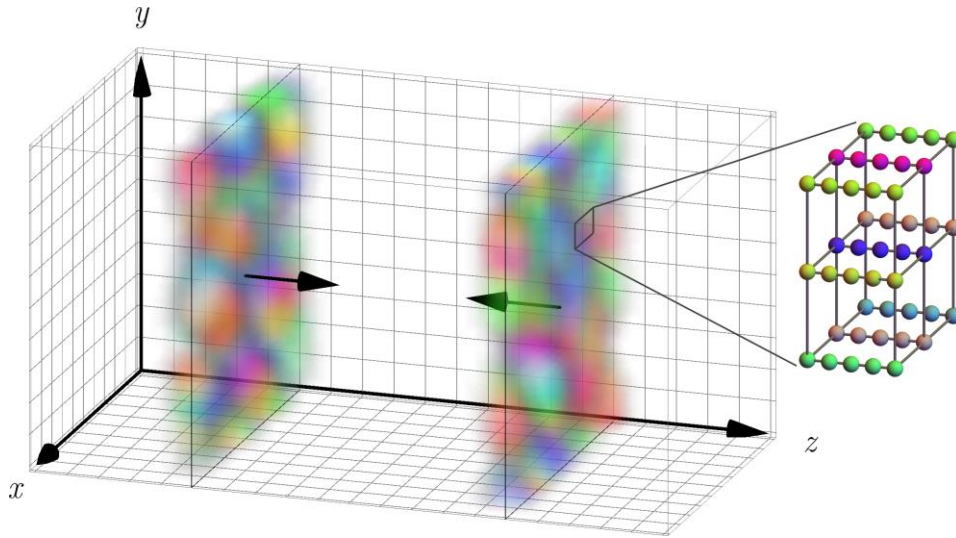
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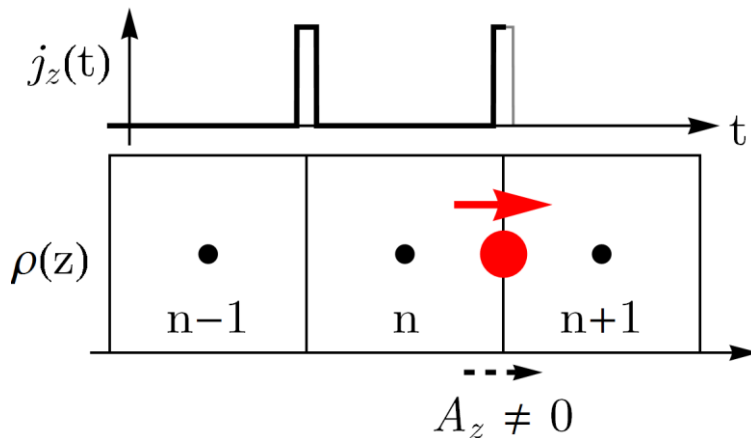
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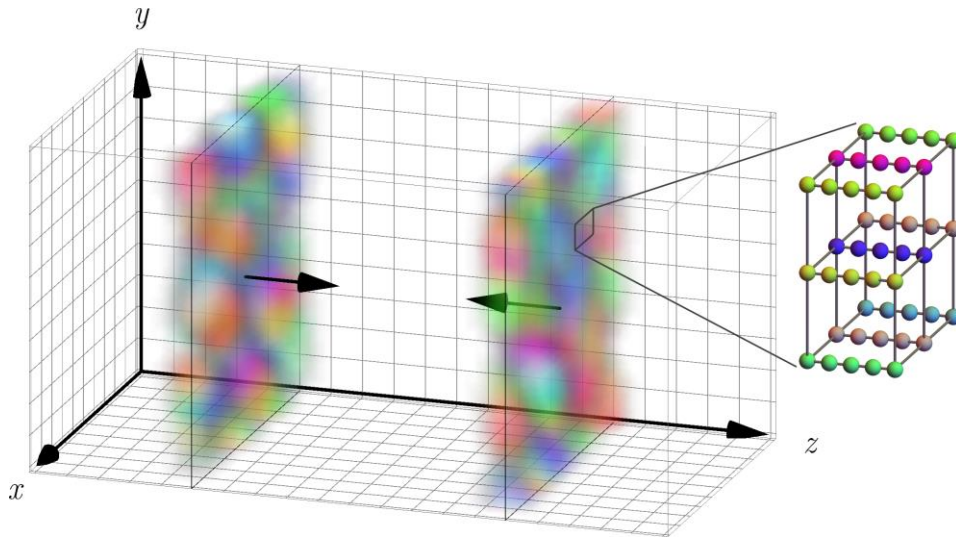
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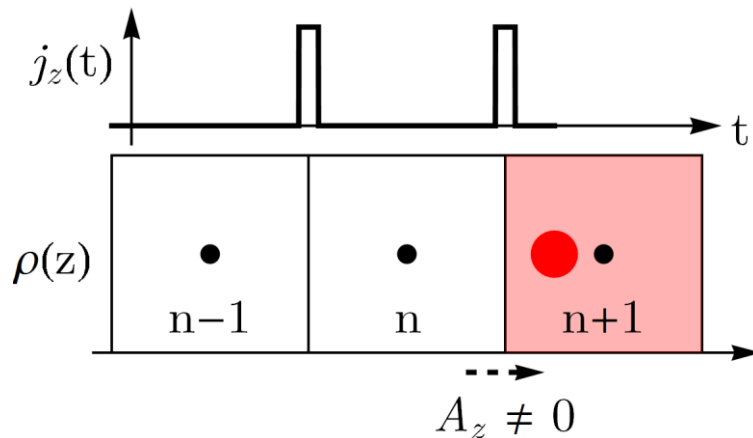
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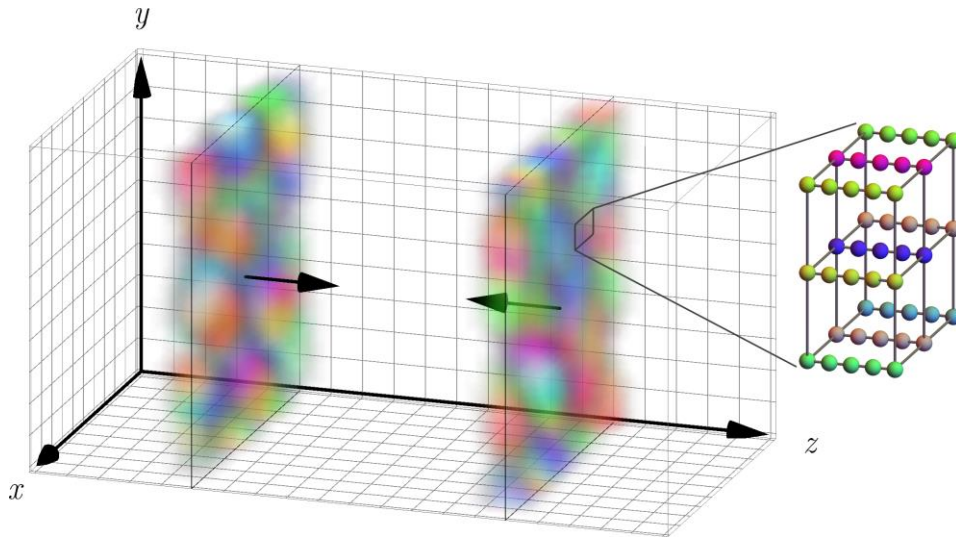
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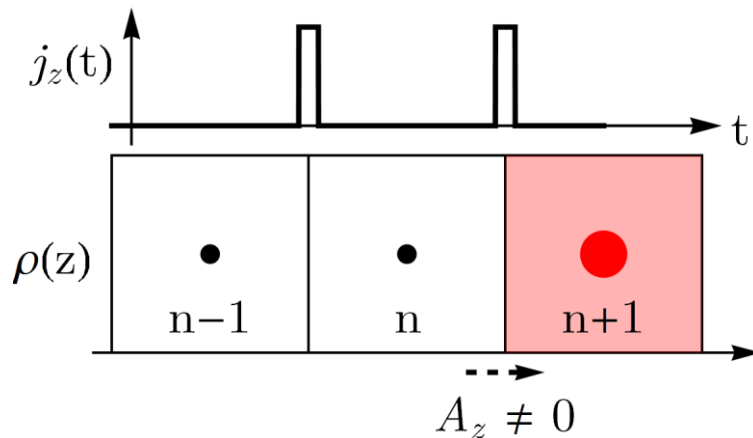
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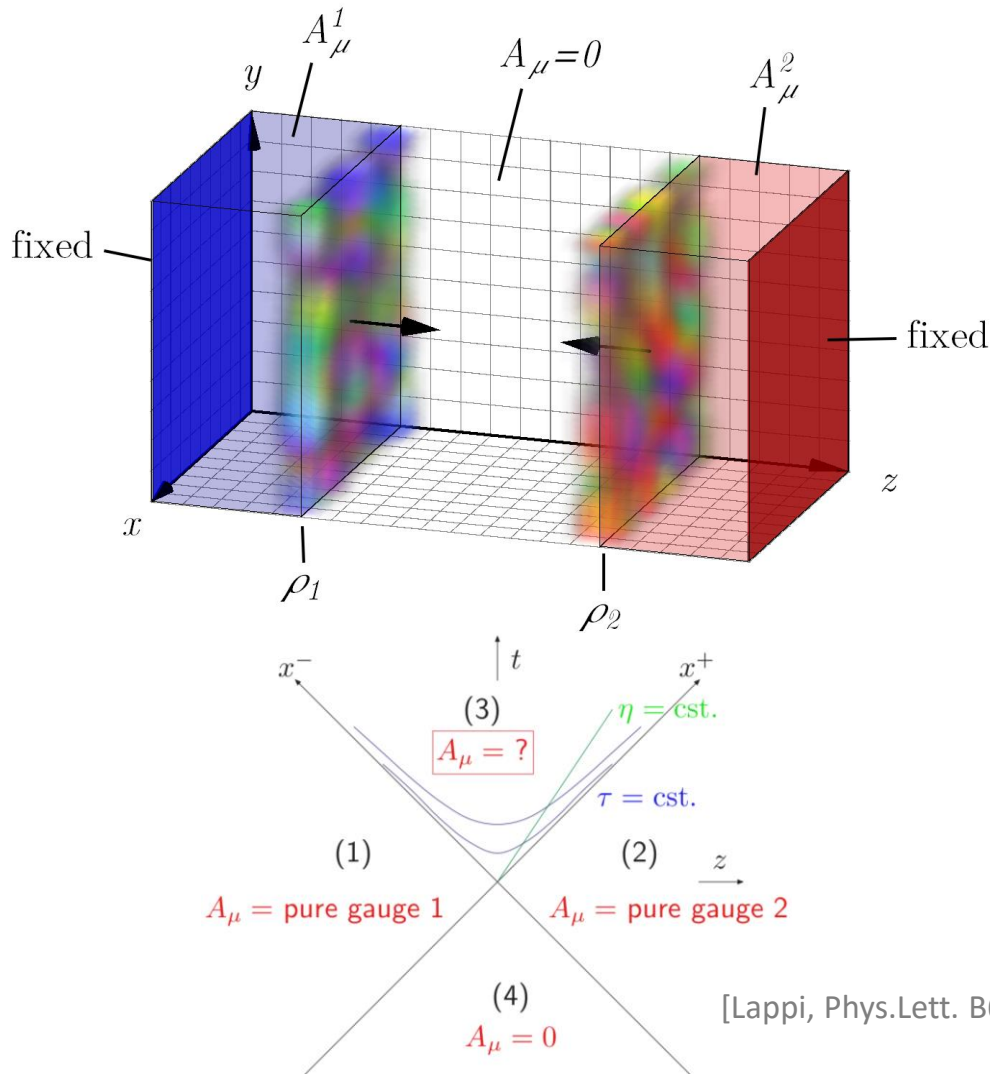
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Initial conditions

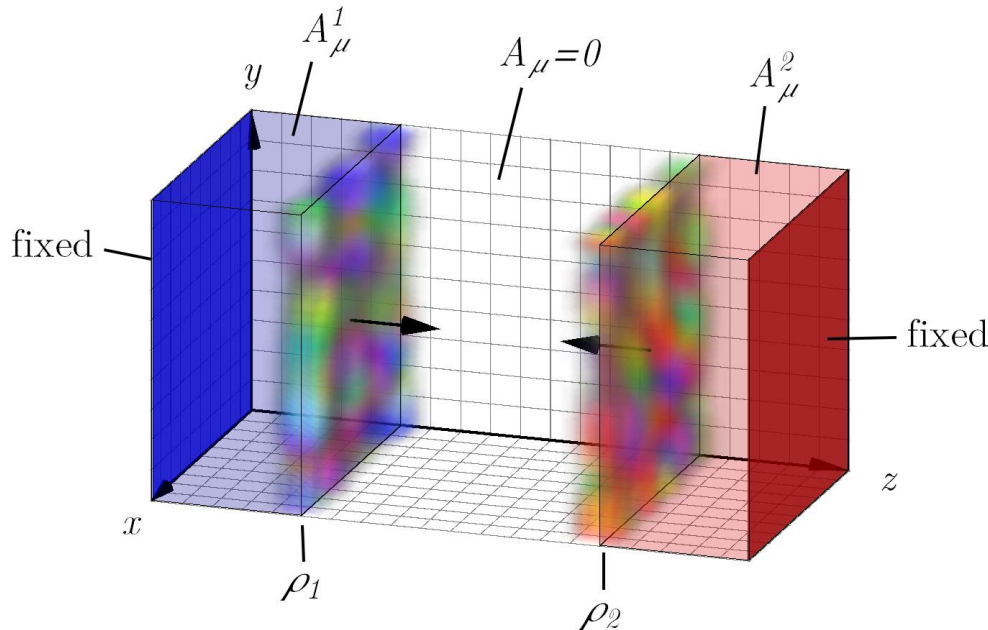


- Temporal gauge ($A_0 = 0$) suitable for numerical time evolution.
- Asymptotically pure gauge “trails” behind nuclei.
- Fixed boundary conditions on the longitudinal boundaries are required.
- Random charge densities $\rho_{(1,2)}$ are sampled from **McLerran-Venugopalan (MV)** model.

[McLerran, Venugopalan: PRDD49 (1994) 3352-3355]

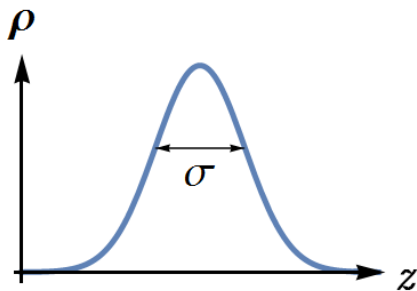
[Lappi, Phys.Lett. B643 (2006) 11-16]

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finite longitudinal thickness

$$\langle \rho^a(x_T) \rho^b(y_T) \rangle = g^2 \mu^2 \delta^{(2)}(x_T - y_T)$$

MV parameter $\mu \approx 0.5 \text{ GeV (Au)}$

Simulation overview

1. Initialize random charges and fields of two colliding nuclei.

2. Simulation cycle:

a. Move particles and apply parallel transport.

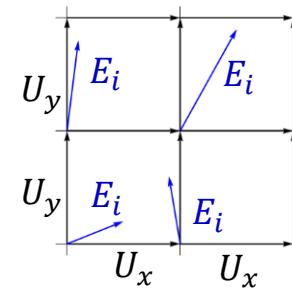
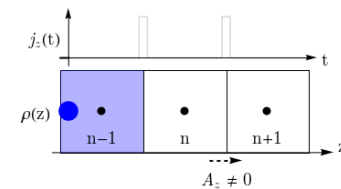
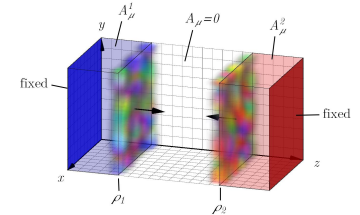
b. Generate currents from particle movement.

c. Evolve fields in time with currents as input.

d. Compute observables ($T_{\mu\nu}$, ε , p_L , p_T , ...).

3. Average over many random events.

➡ $\langle T_{\mu\nu} \rangle, \langle \varepsilon \rangle, \langle p_L \rangle, \langle p_T \rangle, \dots$



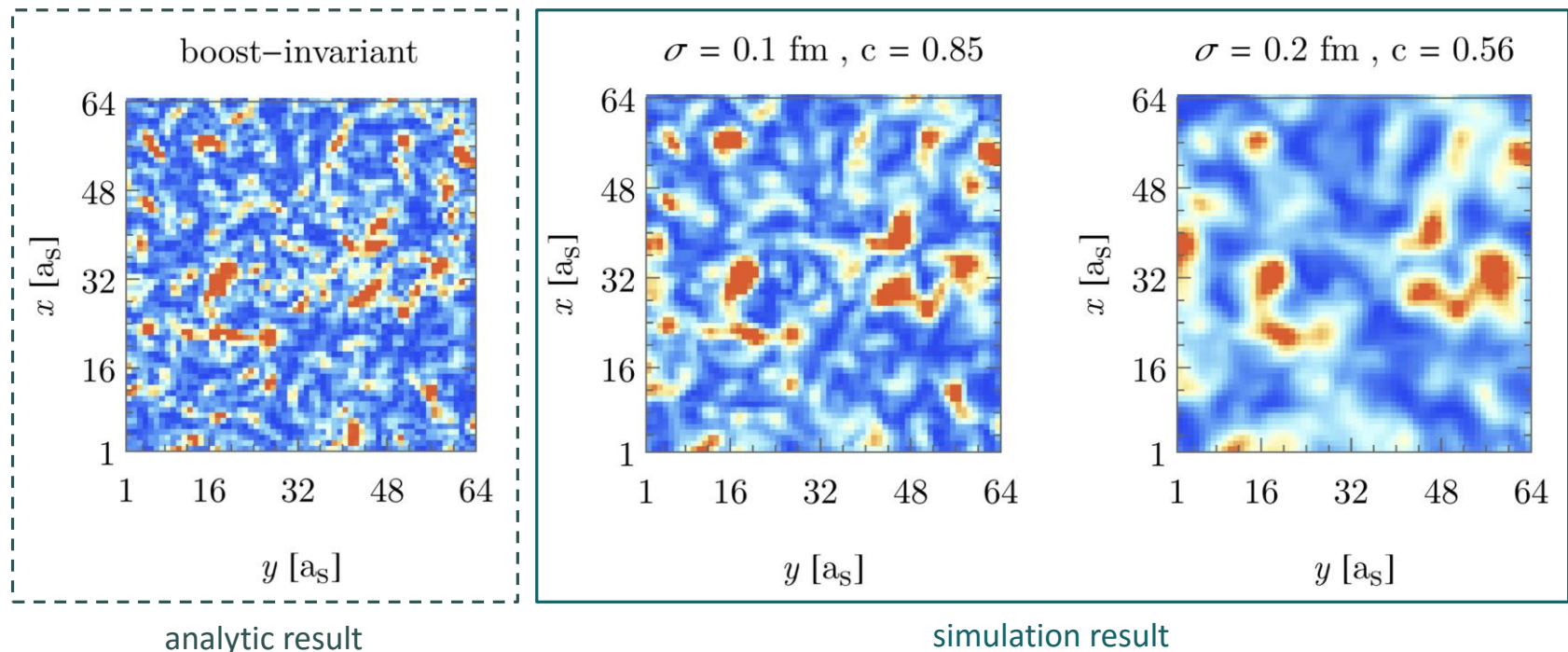
Numerical results

Au-Au collision in the MV model, SU(2)

Comparison to boost-invariant results

- Check validity of simulation results with finite nucleus thickness by comparing to analytical boost-invariant results.
- Compare boost-invariant Glasma initial conditions to simulated fields and vary thickness parameter σ .

Energy density component $\text{tr}E_L^2(x_T)$ in the transverse plane at $\eta = 0$.

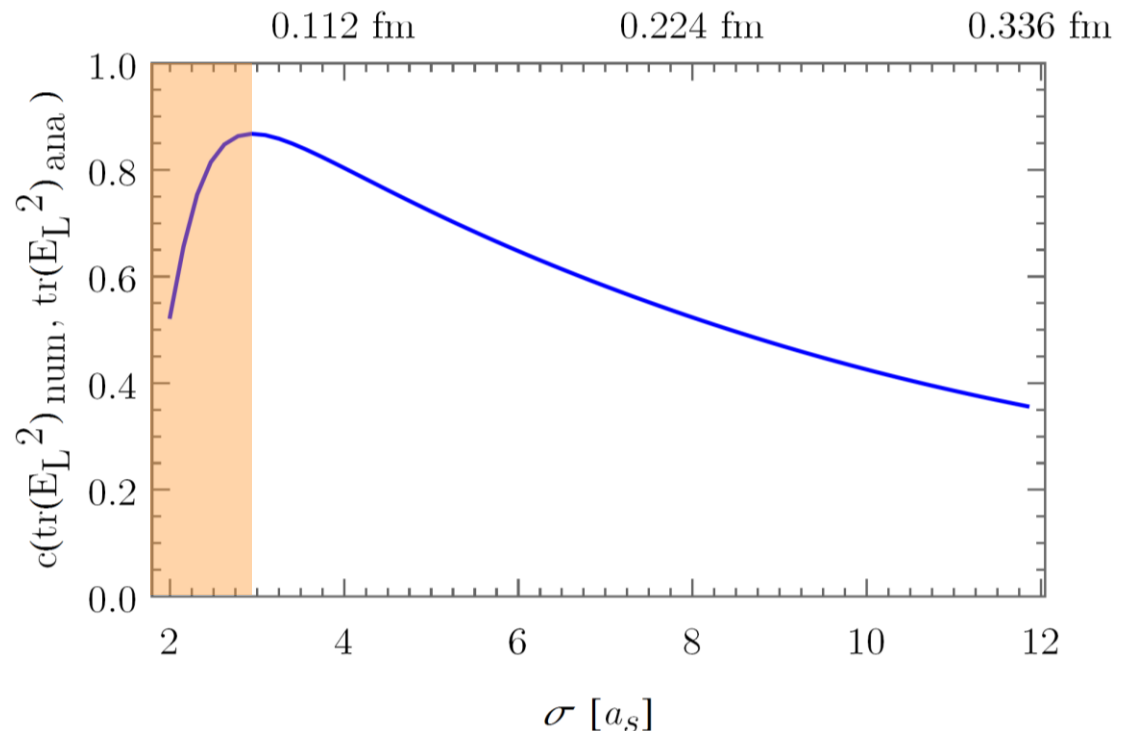


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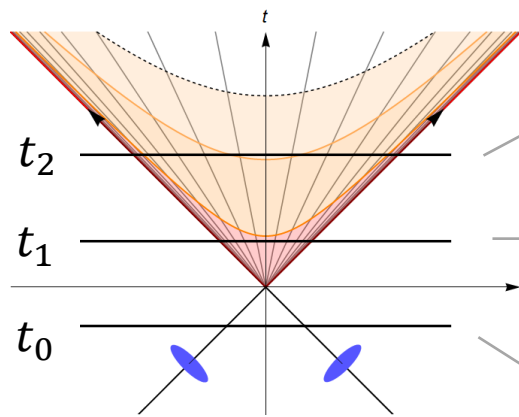
- Compute correlation between analytic and numerical results as a function of σ .
- Thick nuclei: low correlation
- Thin nuclei: high correlation
- Numerical instabilities prohibit very thin nuclei.

(but it's just a question of lattice sizes)

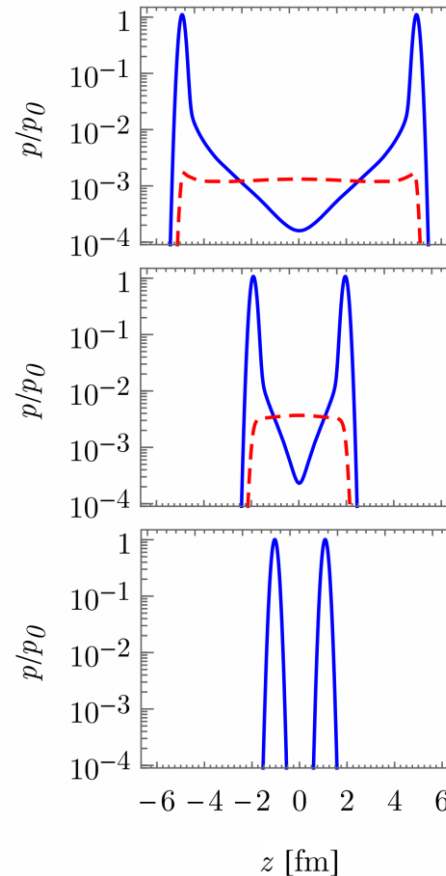


Pressure anisotropy (1)

- Compute longitudinal and transverse pressure $p_L(z)$ and $p_T(z)$ as a function of the longitudinal coordinate z .



→ Pronounced pressure anisotropy



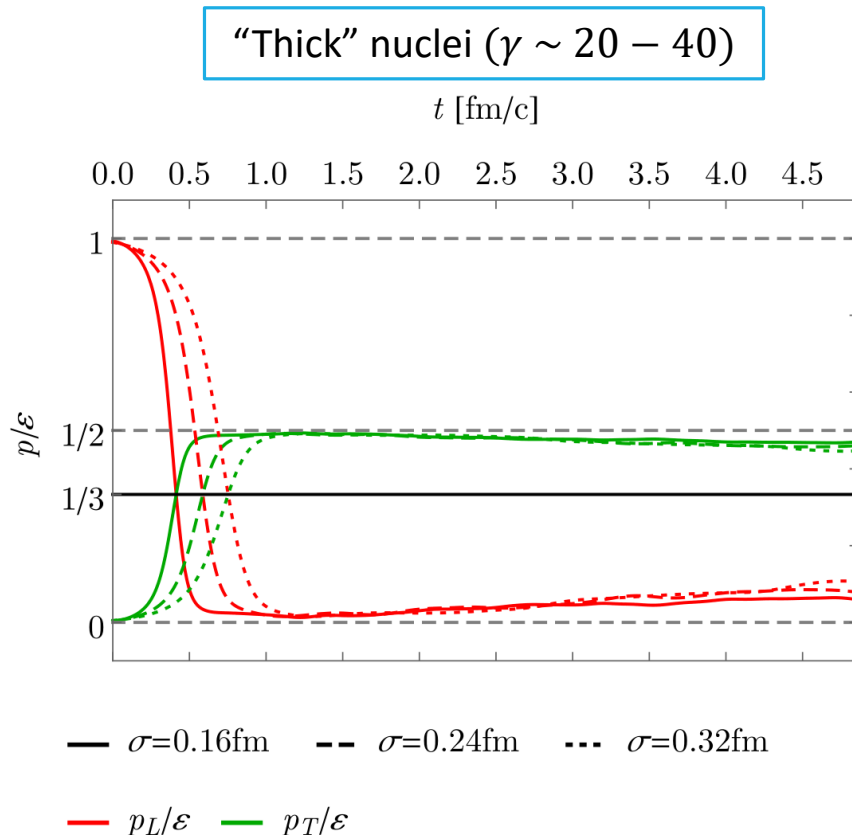
$t_2 = +5 \text{ fm}/c$
(late times)

$t_1 = +2 \text{ fm}/c$
(after collision)

$t_0 = -1 \text{ fm}/c$
(before collision)

Pressure anisotropy (2)

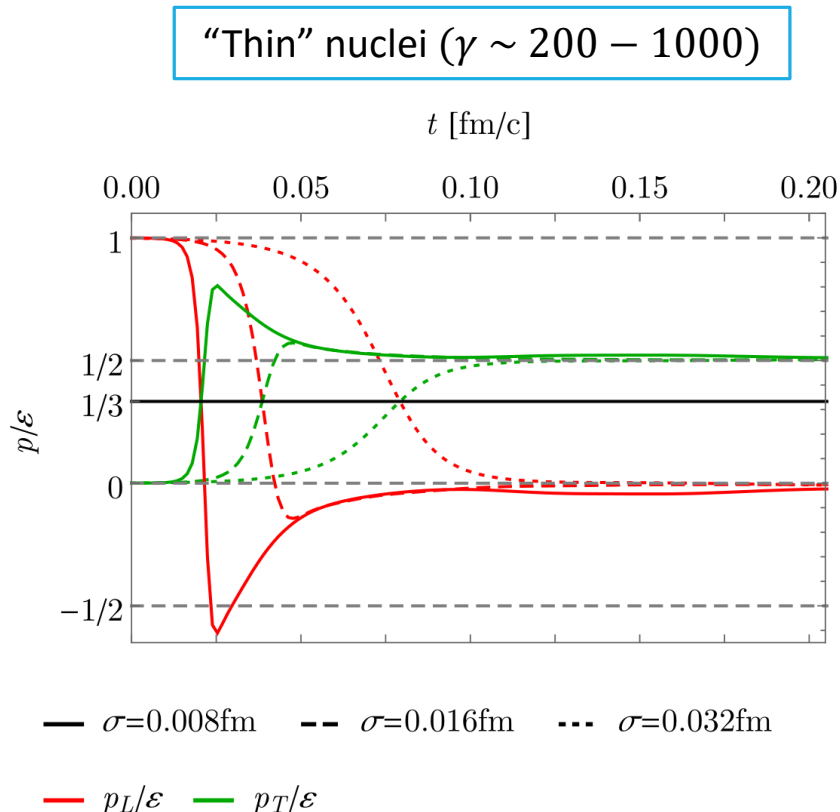
- **Isotropization:** initial pressure anisotropy should vanish after ~ 0.1 fm/c to a few fm/c.
- Boost-invariance breaking perturbations drive system towards isotropization. [Epelbaum, Gelis, PRL 111 (2013) 232301]. Finite thickness breaks boost-invariance.



- Analyze pressure to energy density ratio in the central region at $\eta = 0$.
- Thick nuclei: pronounced pressure anisotropy (free-steaming).
- Slight movement towards isotropization visible, but it is too slow.
- Negative longitudinal pressures?

Pressure anisotropy (2)

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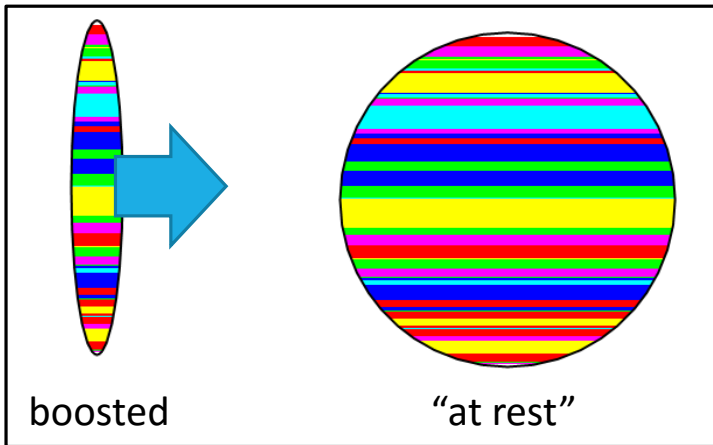


- Analyze pressure to energy density ratio in the central region at $\eta = 0$.
- Thin nuclei: negative longitudinal pressures
- Observables always influenced by presence of the nuclei at early times.

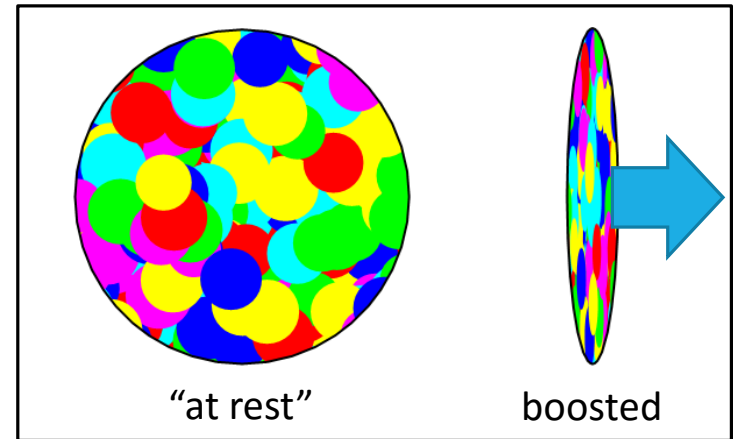
Longitudinal structure

- Initial conditions are still missing **random longitudinal structure**.
- Longitudinal randomness...
 - leads to higher energy density in the glasma.
[Fukushima, PRD 77 (2008) 074005]
 - could provide boost-invariance breaking perturbations.
- Possible consequence: **faster isotropization times? → future work!**

Current implementation



Longitudinal randomness



Conclusions

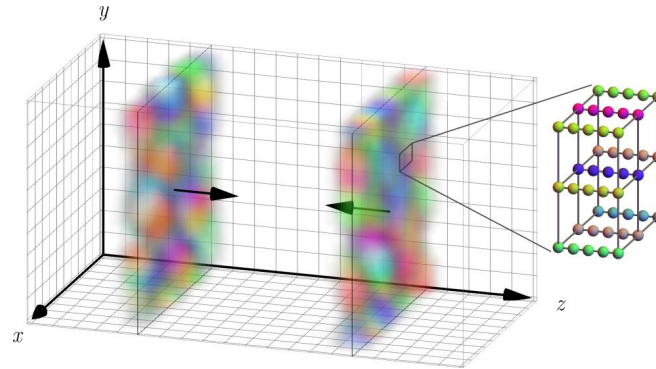
- Simulating CGC collisions in 3D+1 with finite nucleus thickness in the laboratory frame using CPIC is viable.
- Boost-invariant results reproduced in the limit of thin nuclei.
- We observe a pronounced pressure anisotropy after the collision.
- Observed isotropization too slow.
- **Future:** Study effects of initial conditions with random longitudinal structure on isotropization.

Open



arXiv:1605.07184 [hep-ph]
open source: <https://github.com/openpixi>

Thank you for your attention!



Open



arXiv:1605.07184 [hep-ph]
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Backup slides

Pressure in comoving frame

- Pressure components $p_L(t, z)$ and $p_T(t, z)$ computed in laboratory frame.
- Rapidity dependence?
- Transform $p_L(t, z)$ to $\bar{p}_L(\tau, \eta)$.

$$\bar{p}_L(\tau, \eta) = p_L(\tau, \eta) \cosh^2 \eta + \varepsilon(\tau, \eta) \sinh^2 \eta - 2S_L(\tau, \eta) \cosh \eta \sinh \eta$$

$$S_L \equiv 2\text{tr} \left(\vec{E} \times \vec{B} \right)_z$$

“Thick” nucleus ($\gamma \sim 20$)

