

# Possible 750 GeV diphoton signal via light pseudoscalars

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with C. Hugonie, arXiv:1602.03344,

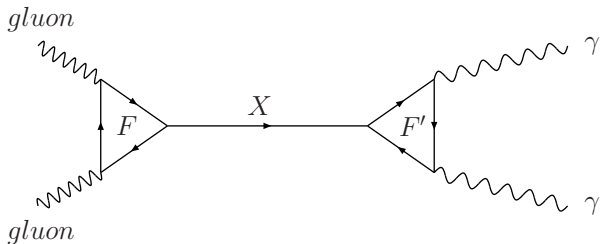
see also F. Domingo et al., arXiv:1602.07691

Data ( $\lesssim$  Moriond 2016):

- ATLAS at 13 TeV,  $710 \text{ GeV} < M_{\gamma\gamma} < 790 \text{ GeV}$  (two bins):  
21 events vs. 11.3 expected; local excess  $3.9 \sigma$  ( $2.0 \sigma$  incl. LLE);  
compatible with 8 TeV at the  $1.2 \sigma$  level (assuming  $ggF$ )
- CMS at 13 TeV,  $750 \text{ GeV} < M_{\gamma\gamma} < 770 \text{ GeV}$  (one bin):  
11 events vs. 5.4 expected; local excess  $2.8 \sigma$  ( $\sim 1 \sigma$  incl. LLE);  
combined with 8 TeV: local excess  $3.4 \sigma$  ( $1.6 \sigma$  incl. LLE)
- Signal cross sections of  $\sim 3 - 8 \text{ fb}$  would explain the excesses



"Standard" interpretation:



- $X$ : Scalar or pseudoscalar (possibly composite) with  $M_X \sim 750$  GeV
- Coupling to gluons through loops of coloured fermions  $F$
- Coupling to photons through loops of charged fermions  $F'$  ( $\sim F$ ?)
- Possibly a large width ( $\gtrsim$  a few GeV) in order to explain the ATLAS data



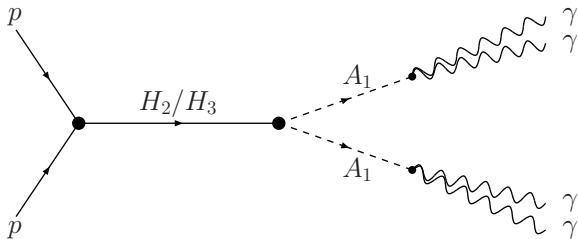
## Challenges:

- Need large (loop induced) production cross section  
→ need large ( $\sim$  non-perturbative)  $XFF$  Yukawa coupling
- Need large (loop induced) width into  $\gamma\gamma$   
→ need large ( $\sim$  non-perturbative)  $XF'F'$  Yukawa coupling
- Tree level decays of  $X$  must be (practically) forbidden, otherwise the loop induced decay into  $\gamma\gamma$  would have a too small branching fraction  
→  $X$  must not couple to Standard Model fermions (or Higgs), the new fermions  $F$  ( $F'$ ) must be heavier than  $M_X/2 \sim 375$  GeV
- A large width into  $\gamma\gamma$  is tough to get...
- $\gtrsim 200$  BSM scenarios of this type... (more than events)



## Alternative scenario with light pseudoscalars $A_1$ :

(S. Knapen et al., P. Agrawal et al., J. Chang et al., ...)



Viable if  $M_{A_1} \lesssim 800$  MeV; then the photons from  $A_1$  decays are sufficiently collimated such that they appear (mostly) as a single photon in the electromagnetic calorimeters

## Constraints on resonance(s) $H_{(i)}$ at $\approx 750$ GeV:

- Sufficient production cross section in  $ggF$  or *ass. prod. with  $b$ -quarks*
- Large branching fraction into  $A_1A_1$

## Constraints on a light pseudoscalar $A_1$ below $\approx 800$ MeV:

- Not ruled out by low energy experiments
- Large branching fraction into  $\gamma\gamma$
- Decay length  $\lesssim 1$  m, preferably shorter



## A concrete scenario: the NMSSM

featuring 3 scalars  $H_{1,2,3}$  and two pseudoscalars  $A_{1,2}$

Two candidates for scalar(s)  $H_2/H_3$  at  $\approx 750$  GeV:  
(with  $H_1 = \text{SM-Higgs}$  at 125 GeV)

- the “MSSM-like” scalar  $H$  with potentially large production cross section via  $bbH$  if  $\tan\beta \gtrsim 10$
- the singlet-like scalar  $H_S$  with potentially large branching fraction into singlet-like  $A_1A_1$  ( $A_2$  is the MSSM-like pseudoscalar with  $M_{A_2} \sim M_H$ )
- > Best solution: both scalars have masses of  $\approx 750$  GeV,  $H$  and  $H_S$  mix strongly and form  $H_2/H_3$ ; two nearby narrow states can imitate a large width as seen by ATLAS



A light pseudoscalar  $A_1$  can be a (pseudo-) Goldstone boson of an  $R$ -symmetry ( $\leftrightarrow$  small trilinear couplings  $A_\lambda, A_\kappa$  in the scalar potential);

Impossible in the MSSM where the  $\mu$ -term breaks  $R$ -symmetry;  
in the NMSSM,  $\mu$  is replaced by the vev of a singlet field  $S$   
 $\rightarrow$  a (weakly broken)  $R$ -symmetry is possible

**But:** Broken by radiative corrections  $\sim A_{top}$ , gaugino masses  
 $\rightarrow$  Tuning is still required for  $M_{A_1} \lesssim 800$  MeV



## Possible $A_1$ masses satisfying the above constraints:

(1)  $M_{A_1} \sim M_{\pi^0} \sim 135 \text{ MeV}$  (Domingo et al., arXiv:1602.07691):

- $A_1$  mixes with  $\pi^0$ , hence  $A_1$  decays with a similar width (short decay length) into  $\gamma\gamma$ ; calculable using PCAC

Heavier  $A_1$ :  $135 \text{ MeV} < M_{A_1} < 2m_\mu$ :

- Susy loops generate flavour changing couplings of the extra (MSSM-like) Higgs bosons, hence also for  $A_1$  (through mixing with the MSSM-like  $A_2$ )
- dangerous rare decays  $K^\pm \rightarrow \pi^\pm e^+ e^-$  (less constraining:  $B^\pm \rightarrow K^\pm e^+ e^-$ ) unless the soft Susy breaking terms are chosen such that contributions to flavour changing couplings cancel, which is possible (see arXiv:1602.07691)
- $A_1$  decays dominantly into  $e^+ e^-$  with a decay length  $\gtrsim 40 \text{ m}$  → useless





(2)  $M_{A_1} \lesssim 2m_\mu \sim 211 \text{ MeV}$  (U.E., C. Hugonie, arXiv:1602.03344):

- The muon loop induced BR into  $\gamma\gamma$  is enhanced up to  $\sim 75\%$  if  $M_{A_1}$  is just below the threshold (see A. Bharucha et al., arXiv:1603.04464)
- The decay length is reduced to 2–5 m, but the production cross section can be large enough such that enough  $A_1 \rightarrow \gamma\gamma$  decays take place before the EM calorimeter
- Soft Susy breaking terms have to be chosen such that flavour changing couplings are cancelled

$M_{A_1} \gtrsim 500 \text{ MeV}$ : Constraints from rare  $K$  decays disappear



### (3) $M_{A_1} \sim 510$ MeV (U.E., C. Hugonie, arXiv:1602.03344):

- At the parton level, the dominant decays of  $A_1$  are into  $s\bar{s}$  and gluons
- But: one is still in the nonperturbative regime of QCD  
Best guess:  $s\bar{s}$  and  $F\tilde{F}_{(QCD)}$  act as interpolating fields;  
these are part of the  $\eta$  wave function in Fock space ( $M_\eta \sim 548$  MeV),  
hence  $A_1$  decays like the  $\eta$  meson:  
 $BR(\eta \rightarrow \gamma\gamma) \sim 39\%$  ,  $BR(\eta \rightarrow 3\pi^0) \sim 33\%$  ,  $BR(\eta \rightarrow \pi^+\pi^-\pi^0) \sim 23\%$
- $BR(A_1 \rightarrow \gamma\gamma) \sim 39\%$  ,  $BR(A_1 \rightarrow 3\pi^0 \rightarrow 6\gamma) \sim 33\%$   
with a decay length below 1 mm (?to be confirmed?)
- Dominant constraint: Now from searches for  $\Upsilon(1S) \rightarrow \gamma\eta$  decays by CLEO where no events were seen  
(but 2 events for  $M_{\pi^+\pi^-\pi^0} \sim 510$  MeV in the  $\eta \rightarrow \pi^+\pi^-\pi^0$  search channel)  
→ constraints on the coupling  $A_1 b\bar{b}$ ;  
if too large, CLEO would have observed  $\Upsilon(1S) \rightarrow \gamma A_1 \rightarrow 3\pi^0$  decays
- These constrain the  $BR(H_{2,3} \rightarrow A_1 A_1)$ , still:  
a signal cross section up to 6.7 fb is possible



(4)  $M_{A_1} \sim M_\eta \sim 550$  MeV (U.E., C. Hugonie, arXiv:1602.03344):

- $A_1$  mixes strongly with the  $\eta$  meson, its corresponding branching fractions are no longer educated guesses (calculable using PCAC)
- But: Constraints from CLEO from unseen  $\Upsilon(1S) \rightarrow \gamma A_1$  decays are somewhat stronger, still:
  - a signal cross section up to 3.4 fb is possible

These are the only known scenarios for the 750 GeV diphoton excess without extra “ad hoc” fermions, but based on an old ( $\sim 35$  years) Susy extension of the SM



If the excess of events persists, these scenarios can be distinguished (or ruled out) experimentally:

- If  $M_{A_1} \sim 211$  MeV: The  $A_1$  decay length is macroscopic, and  $A_1$  may decay inside the EM calorimeters (before the EM calorimeters, the  $A_1 \rightarrow \gamma\gamma$  vertex is invisible)
- The photons can convert in the material before the EM calorimeter leading to electrons which are visible, but usually added to the photon signal in the EM (20% for rapidity  $\eta \sim 0$  to 45% for  $\eta \sim 1.6$ )
  - photon-jets lead to more converted photons than a single photon
  - one can potentially distinguish single photons from collinear diphotons or, in the case  $A_1 \rightarrow 3\pi^0 \rightarrow 6\gamma$ , from collinear 6 photons (B. Dasgupta et al., arXiv:1602.04692) iff the  $A_1$  decays occur inside the material
- If the signal originates from two nearby states  $H_2/H_3$ , their masses can potentially be separated (depending on the actual  $H_2/H_3$  mass splitting)

Exciting times may lie ahead of us!

