

Advances in Particle Astrophysics

Session IV: IceCube's neutrinos

Kfir Blum

Weizmann Institute

CERN academic training 11-15/04/2016

Gton of instrumented ice at the South Pole.

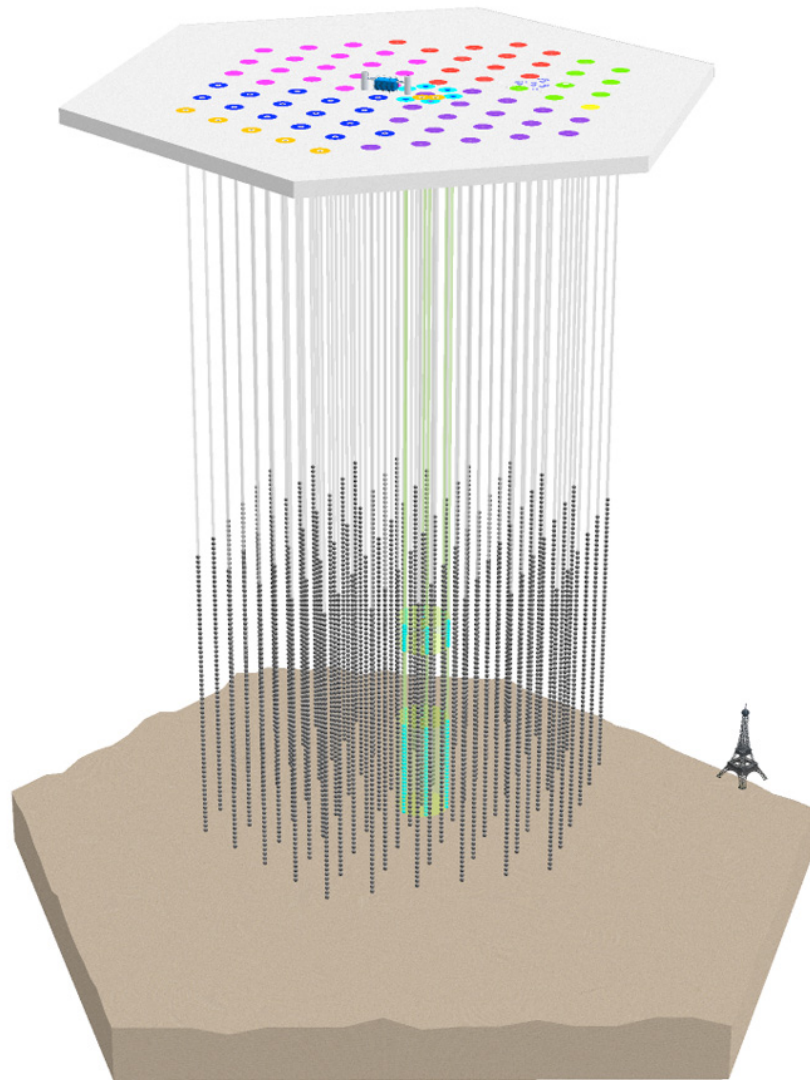
We will focus on theory implications from discoveries made w/ this Gton of ice.

Some basics you have to know:

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$10^{14} \times$



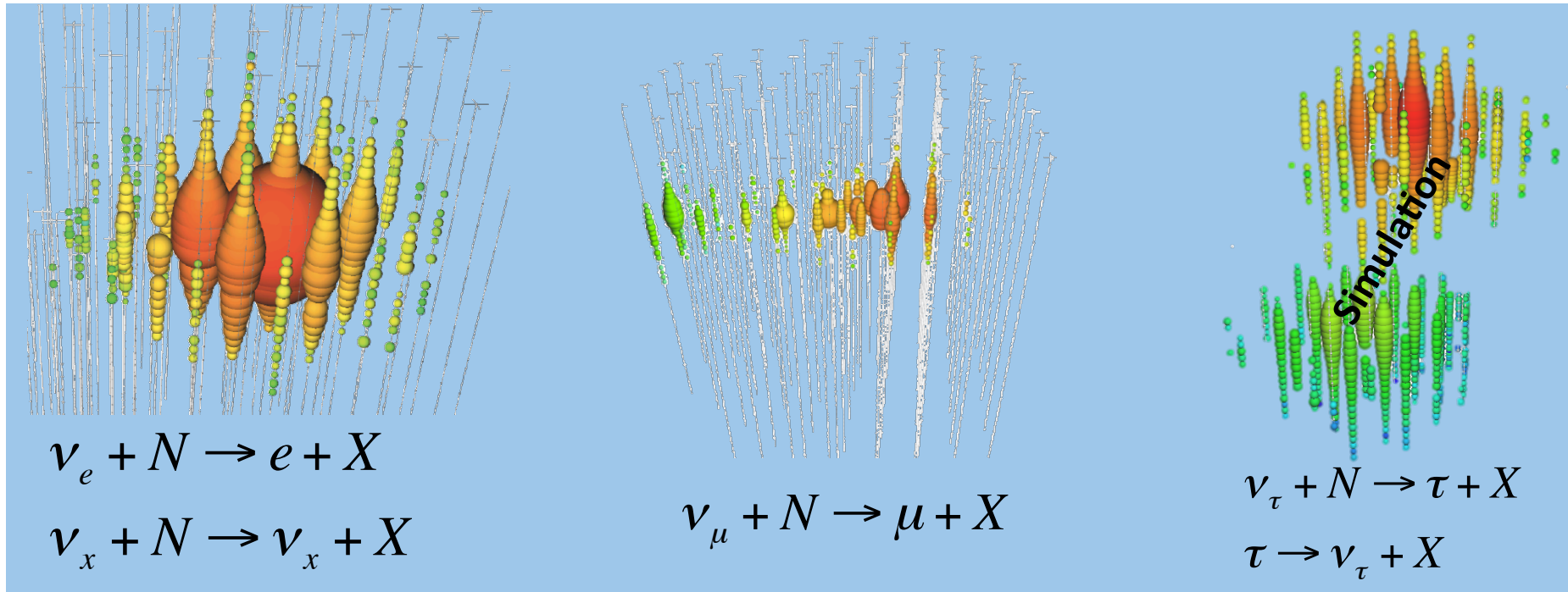
5000 x



Gton of instrumented ice at the South Pole.

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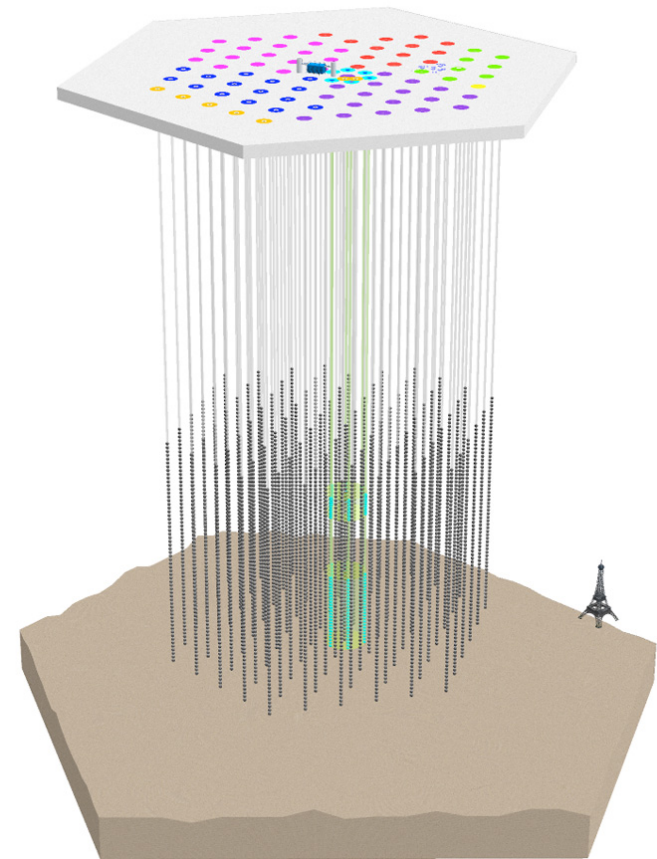
Jan Auffenberg, Moriond 2016

Outline:

...in 2013, IC **announced a discovery** of high-energy astrophysical neutrinos.

What's making it?

What will it teach us?



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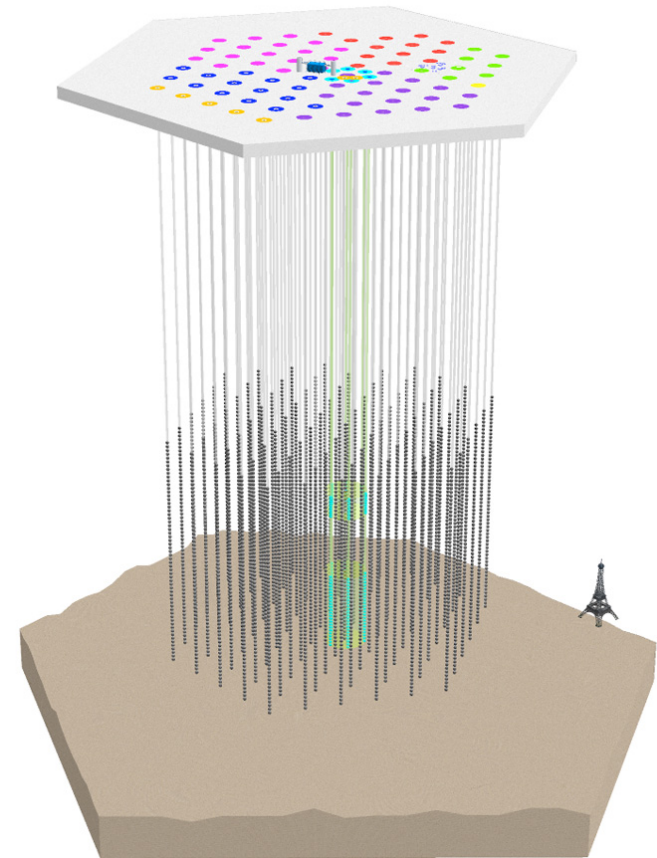
What's making it?

We will put it in context: (U)HECRs, gamma rays

WB going *backwards*

Some predictions.

What will it teach us?



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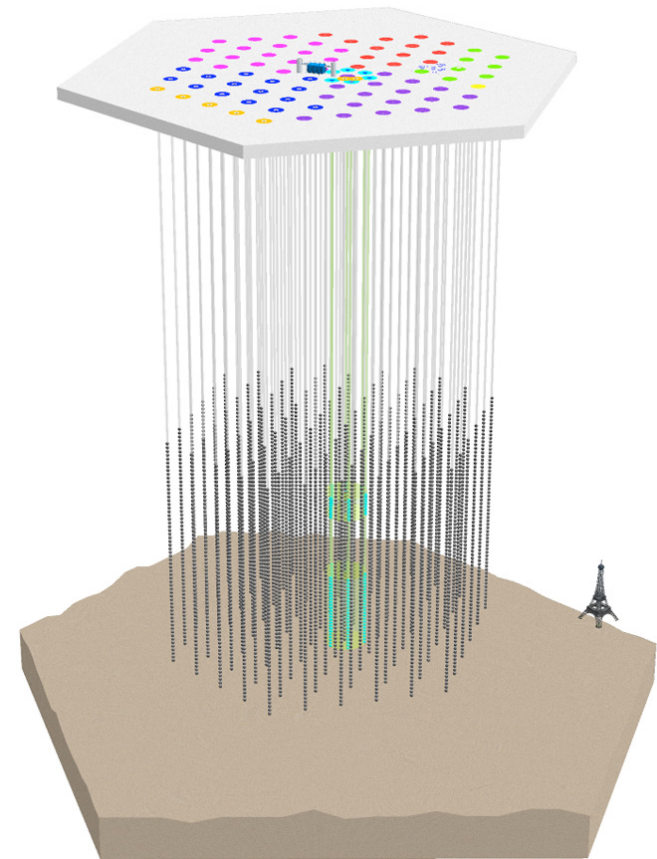
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Likely: sources of (U)HECRs

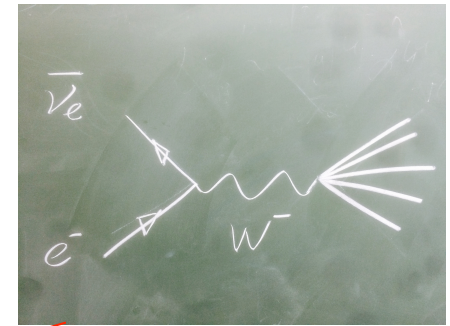
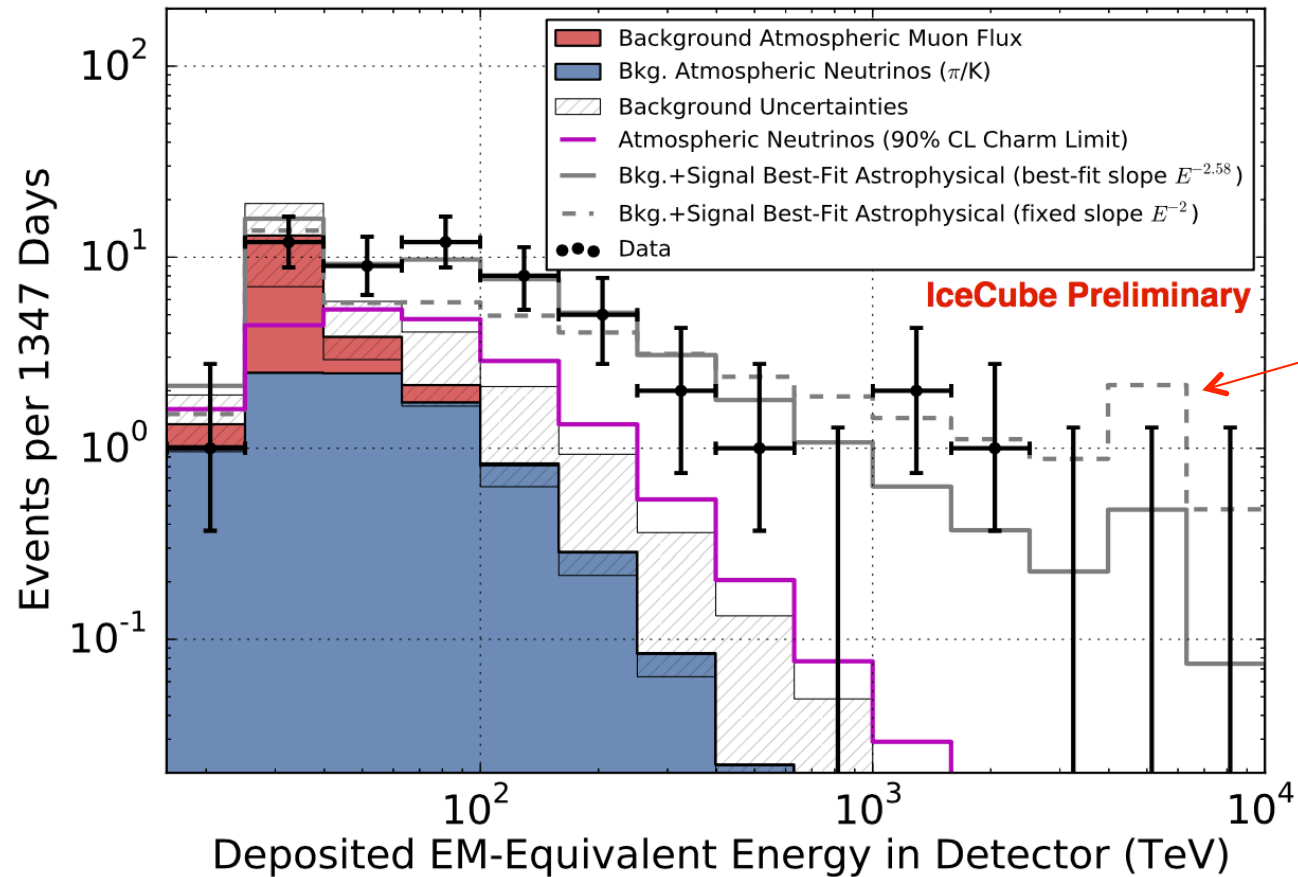
Possibly: neutrino surprises?

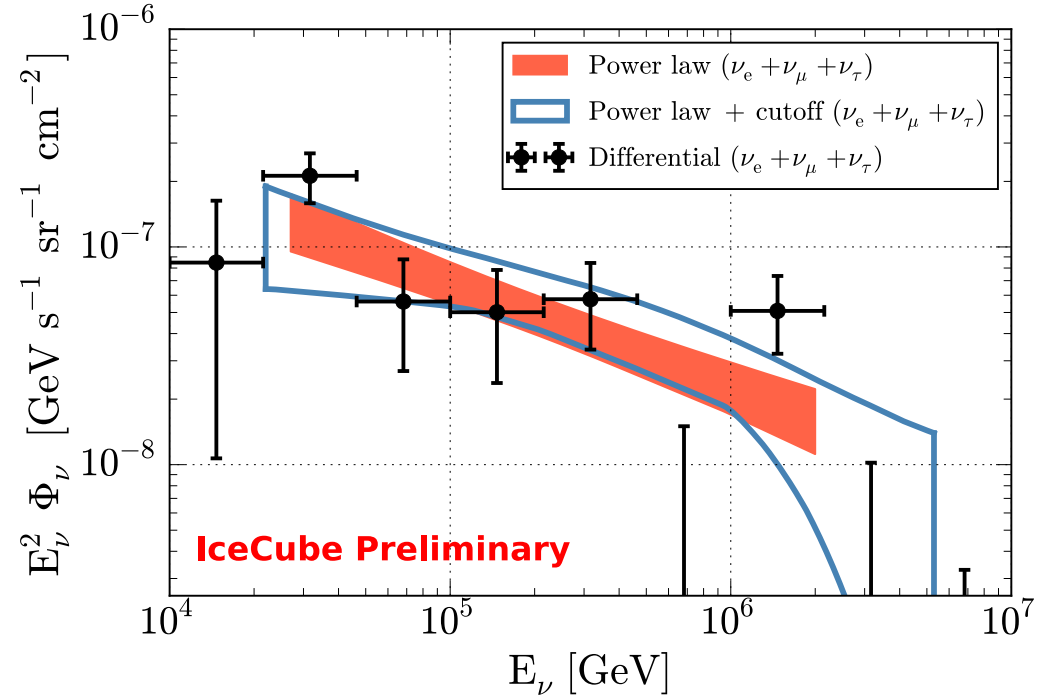
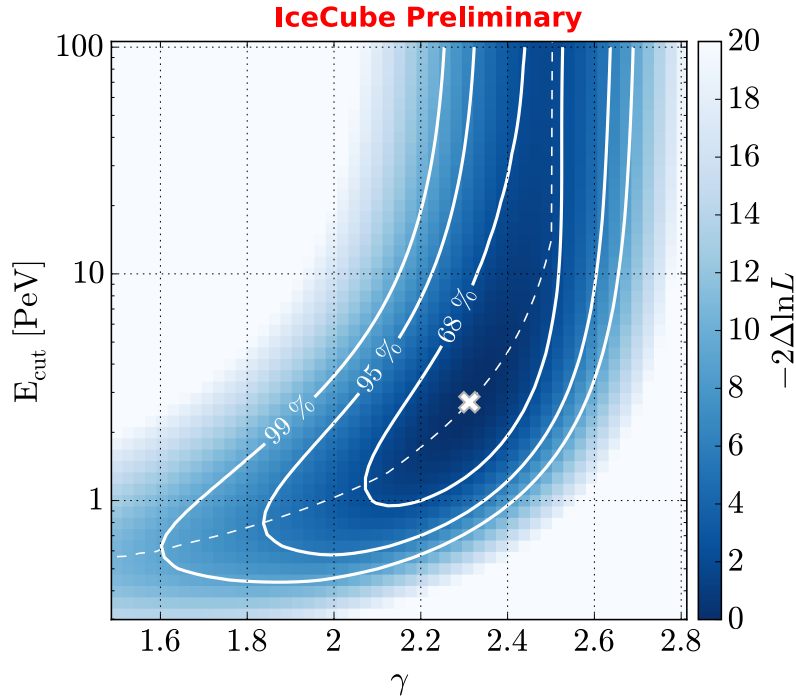
Time permits, we will toy with
exotic possibilities



Moriond 2016 (Jan Auffenberg)
also 1510.05223

Consistent w/ power law, $dn/dE \sim E^{-\gamma}$, with $\gamma \sim 2$





Break around PeV: reasonable.

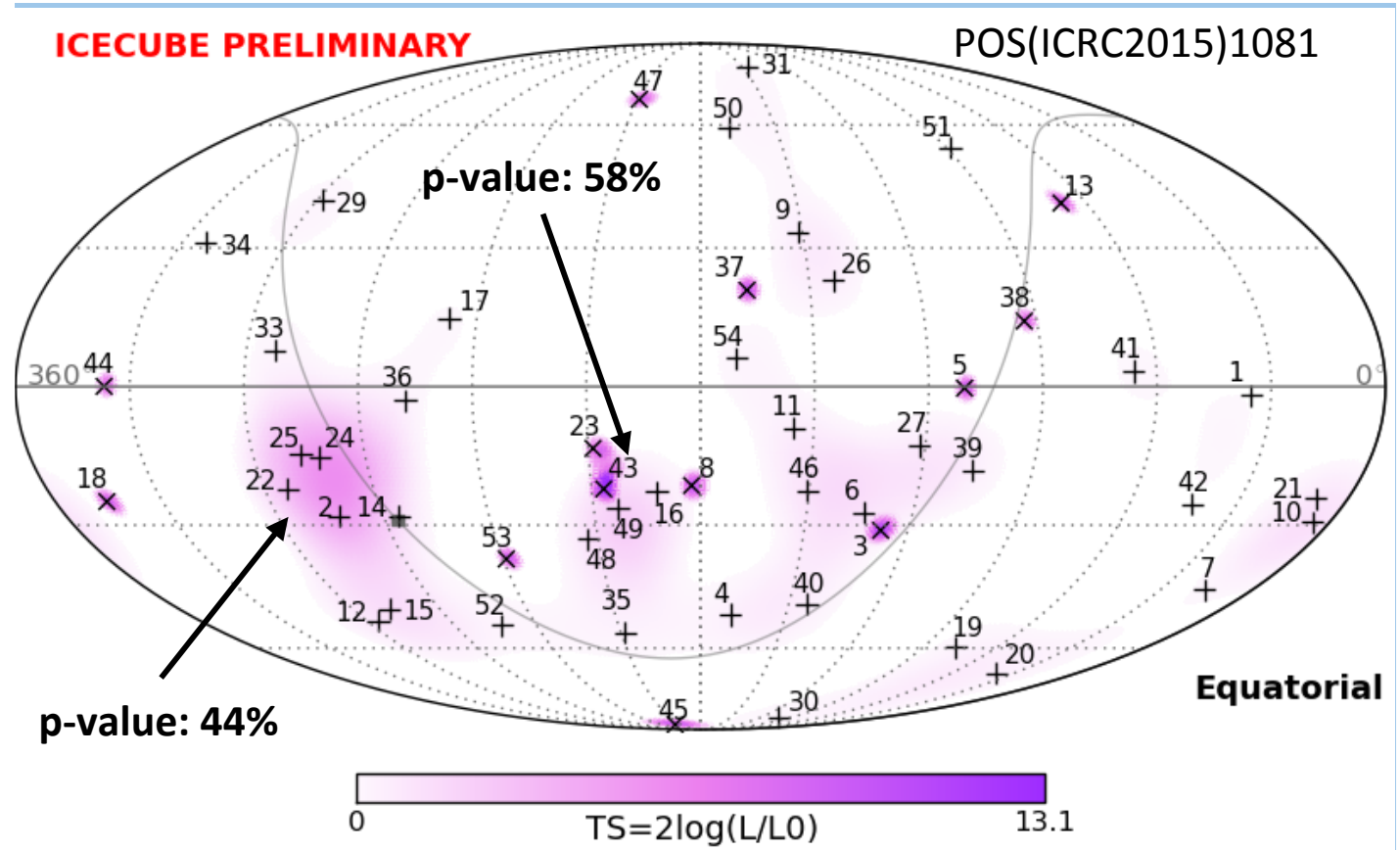
Suggests γ close to 2.

| Param. | Unit | Hyp. A | Hyp. B |
|------------------------|--|------------------------|------------------------|
| ϕ_{conv} | HKKMS | $1.10^{+0.20}_{-0.15}$ | $1.11^{+0.20}_{-0.15}$ |
| ϕ_{prompt} | ERS | $0.0^{+0.7}_{-0.0}$ | $0.0^{+0.8}_{-0.0}$ |
| ϕ | $10^{-18} \text{ GeV}^{-1} \text{ s}^{-1} \text{ sr}^{-1} \text{ cm}^{-2}$ | $7.0^{+1.0}_{-1.0}$ | $8.0^{+1.3}_{-1.2}$ |
| γ | — | $2.49^{+0.08}_{-0.08}$ | $2.31^{+0.14}_{-0.15}$ |
| E_{cut} | PeV | — | $2.7^{+7.7}_{-1.4}$ |
| $-2\Delta\ln L$ | | +1.94 | 0 |

Table 2: Best-fit results for the energy spectrum. The quoted uncertainties are at 1 σ confidence level.

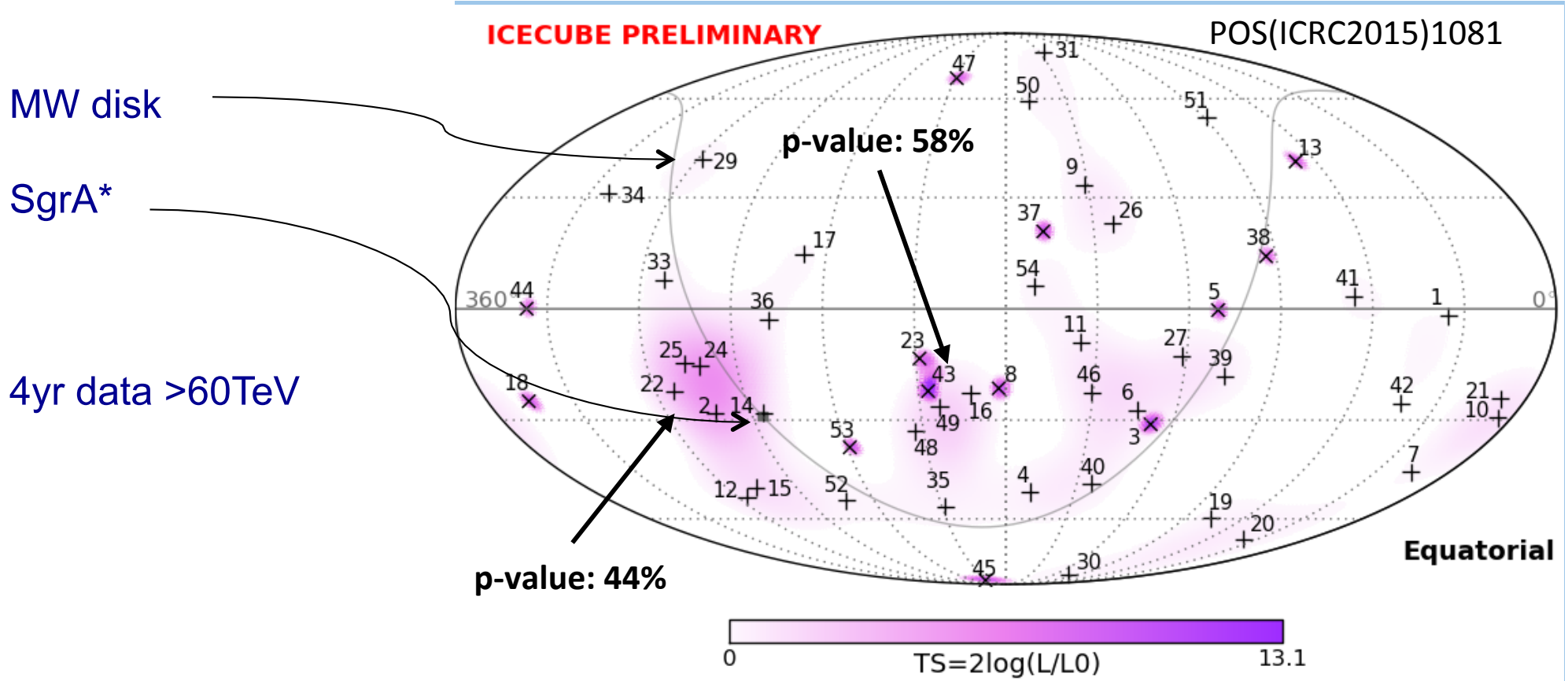
Moriond 2016
1510.05223

4yr data >60TeV



Many events point away from MW disk/bulge.
Consistent w/ isotropy.

Moriond 2016
1510.05223

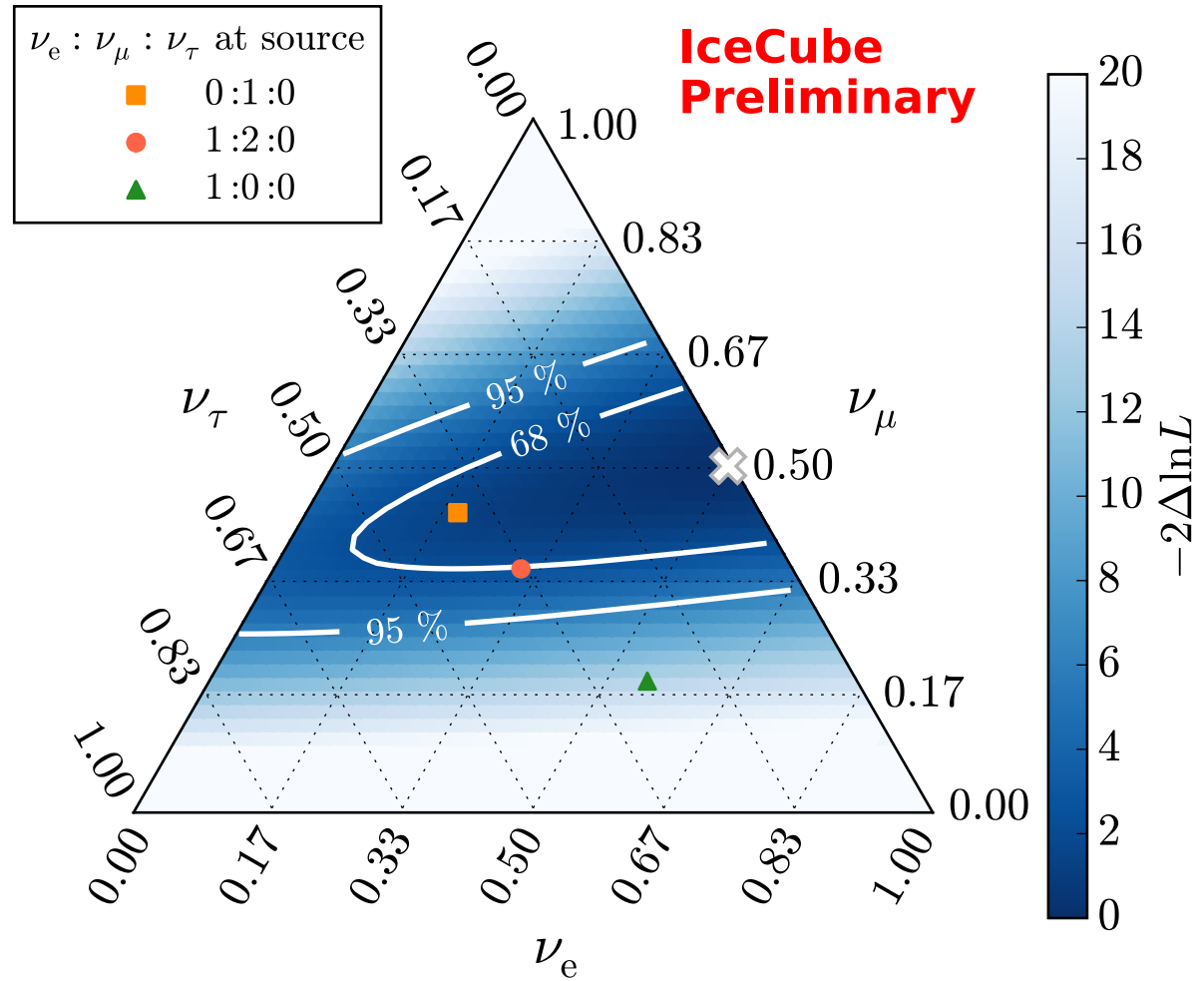


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Moriond 2016
1510.05223

Pion source
1:2:0 → 1:1:1

Muon-damped
0:1:0 → 1:1.8:1.8



IceCube's neutrinos in context:

Flux normalization and rough spectral shape

Angular distribution

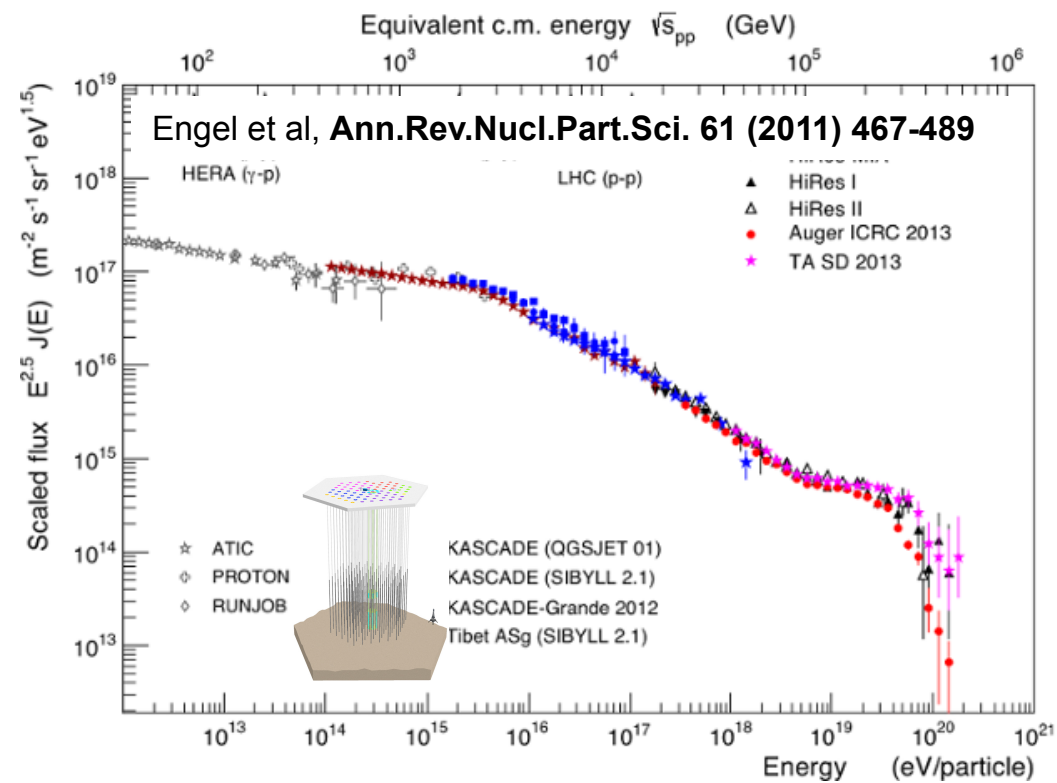
Flavor content

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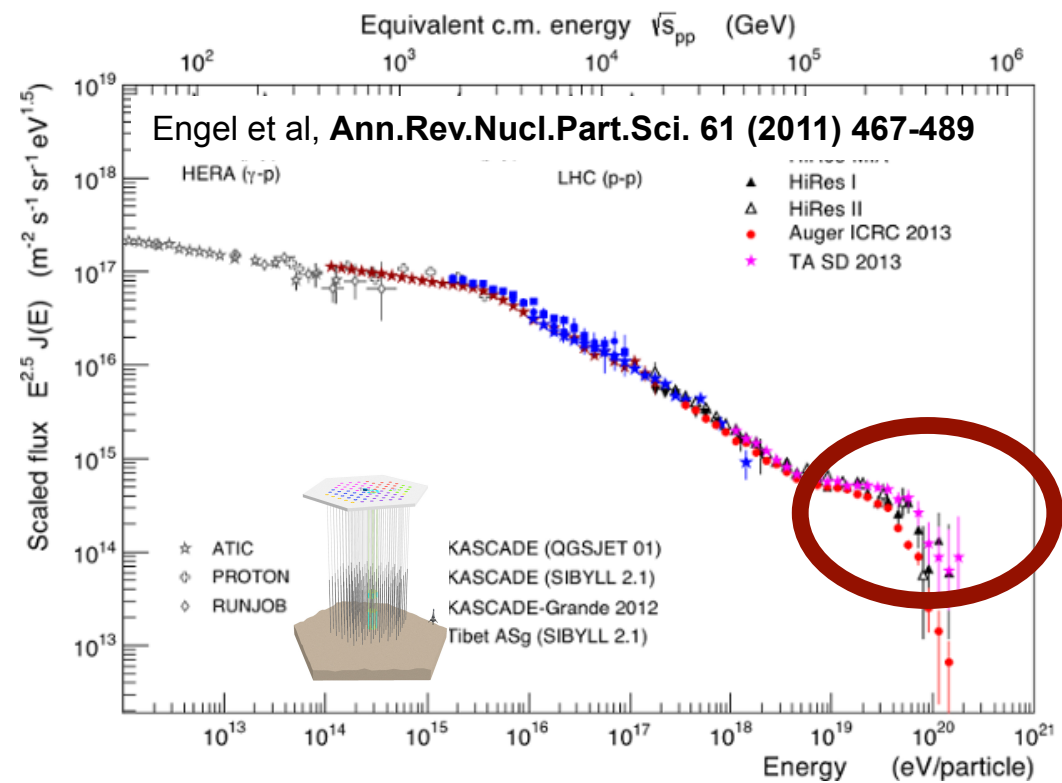


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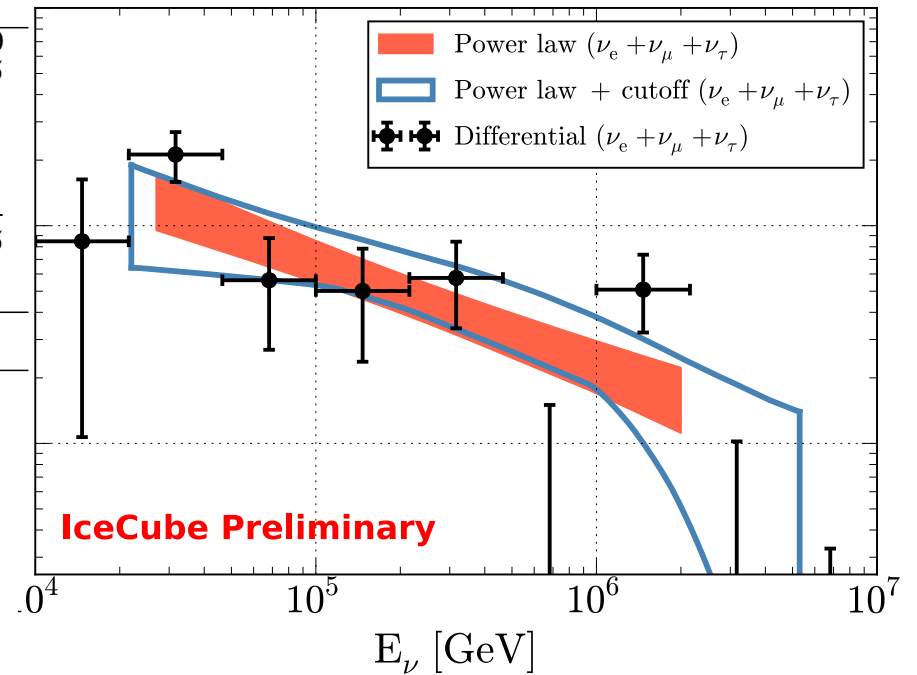


IceCube's neutrinos and UHECRs

Waxman & Bahcall, PRD59, 023 (1998)

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Waxman & Bahcall, PRD59, 023 (1998)

Advances in particle astrophysics:

use IceCube data, run WB argument *backwards*.

see recently, e.g. Yoshida & Takami, PRD90 (2014) no.12, 123012

UHECR and the Waxman-Bahcall bound

$$\begin{aligned}n_{\nu}(>E) &= \int_0^{t_0} dt \dot{n}_{\nu} [>(1+z)E; z] \\ &= \int_0^{\infty} \frac{dz}{(1+z)H(z)} \dot{n}_{\nu} [>(1+z)E; z] \\ &= \frac{\dot{n}_{\nu}(>E, z=0)}{H_0} \int_0^{\infty} \frac{dz f(z)}{(1+z)^{\gamma} \sqrt{0.7 + 0.3(1+z)^3}} \\ \Rightarrow \left. \frac{dn_{\nu}}{dE_{\nu}} \right|_{z=0} &= \left. \frac{d\dot{n}_{\nu}}{dE_{\nu}} \right|_{z=0} \times \frac{1}{H_0} \int_0^{\infty} \frac{dz f(z)}{(1+z)^{\gamma} \sqrt{0.7 + 0.3(1+z)^3}}\end{aligned}$$

UHECR and the Waxman-Bahcall bound

If each CR proton loses fraction η of its energy at the CR source to py collisions before escaping:

$$\underbrace{E_\nu^2 \frac{d\dot{N}_\nu}{dE_\nu}}_{\text{all flavor summed}} = \eta \times \frac{1}{2} \times \frac{1}{2} \times \frac{3}{2} \times \left[E_{CR}^2 \frac{d\dot{N}_{CR}}{dE_{CR}} \right]_{E_{CR} \approx 20 E_\nu}$$

(result the same to factor of 2 if losses are hadronic, pp instead of py)

UHECR and the Waxman-Bahcall bound

$$E_\nu^2 \frac{d\dot{N}_\nu}{dE_\nu} = \eta \times \frac{1}{2} \times \frac{1}{2} \times \frac{3}{2} \times \left[E_{CR}^2 \frac{d\dot{N}_{CR}}{dE_{CR}} \right]_{E_{CR} \approx 20 E_\nu}$$

$$E_{\nu_\mu} \approx E_{\bar{\nu}_\mu} \approx E_{\nu_e}$$

$$\underbrace{E_\nu^2 \frac{d\dot{N}_\nu}{dE_\nu}}_{\text{all flavor summed}} = \eta \times \frac{1}{2} \times \frac{1}{2} \times \frac{3}{2} \times \left[E_{CR}^2 \frac{d\dot{N}_{CR}}{dE_{CR}} \right]_{E_{CR} \approx 20 E_\nu}$$

$$\frac{p\gamma \rightarrow \pi^+ n}{p\gamma \rightarrow \pi^+ n, \pi^0 p}$$

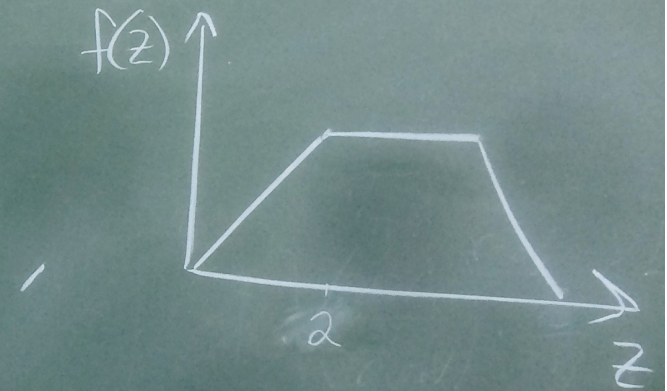
$$\frac{E_{\nu_\mu} + E_{\bar{\nu}_\mu}}{E_\pi}$$

$$\pi^+ \rightarrow \mu^+ \nu_\mu \rightarrow e^+ \nu_e \bar{\nu}_\mu \nu_\mu$$

UHECR and the Waxman-Bahcall bound

$$\eta \times \left[\frac{2}{E_{CR}} \frac{d\dot{N}_{CR}}{dE_{CR}} \right]_{10^{11} \text{ GeV}}^{z=0} = \frac{8}{3} \left(\frac{20 E_\nu}{10^{11} \text{ GeV}} \right)^{\gamma-2} \frac{4\pi H_0}{c} \times \left[E_\nu^2 \Phi_\nu(E_\nu) \right] \times \mathcal{F}_z(\gamma)$$

$$\mathcal{F}_z(\gamma) = \left[\int_0^\infty \frac{dz f(z)}{(1+z)^\gamma \sqrt{0.7 + 0.3(1+z)^3}} \right]^{-1}$$

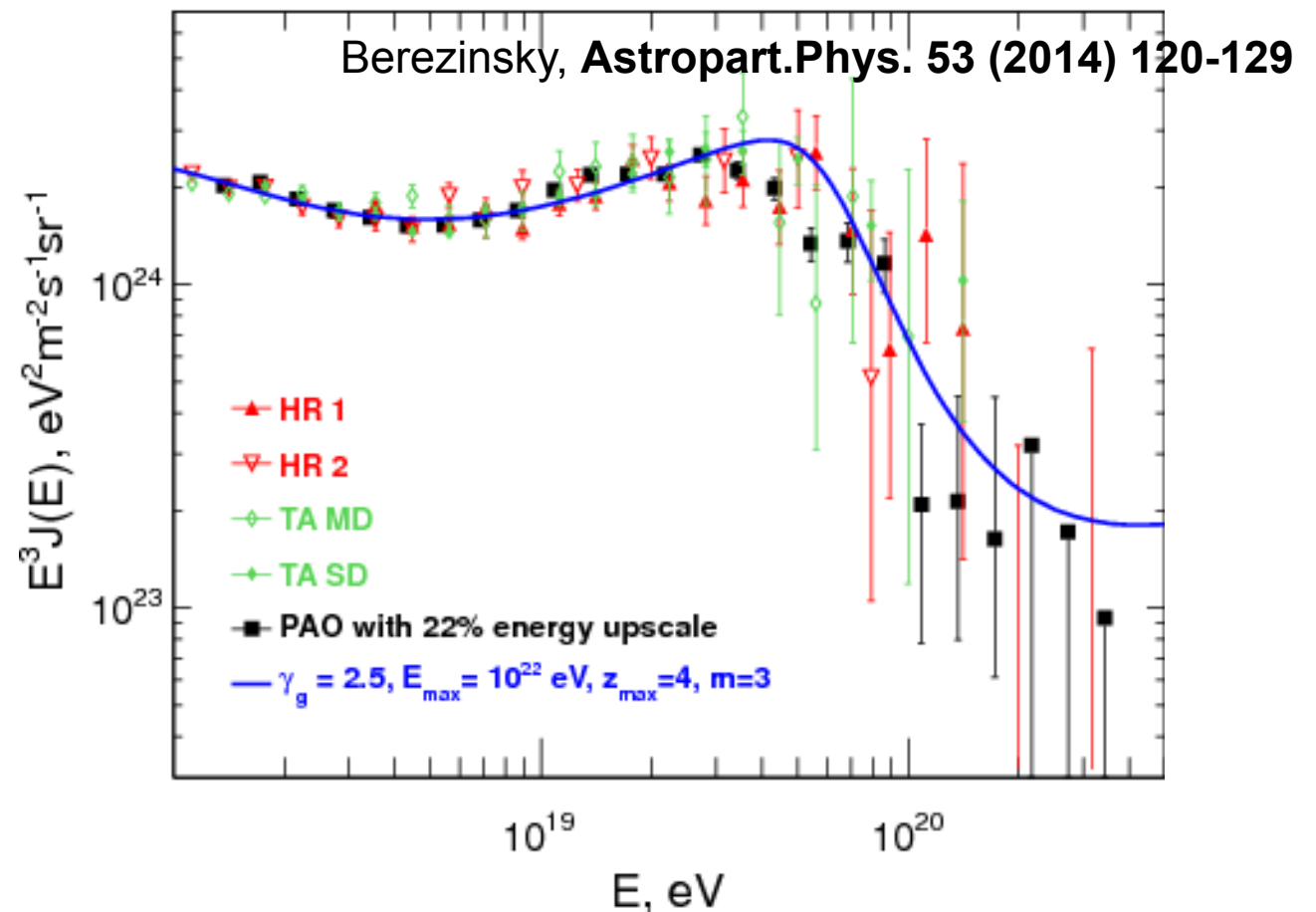


UHECR and the Waxman-Bahcall bound

UHECR source power at $z=0$ constrained from observation of GZK cutoff.

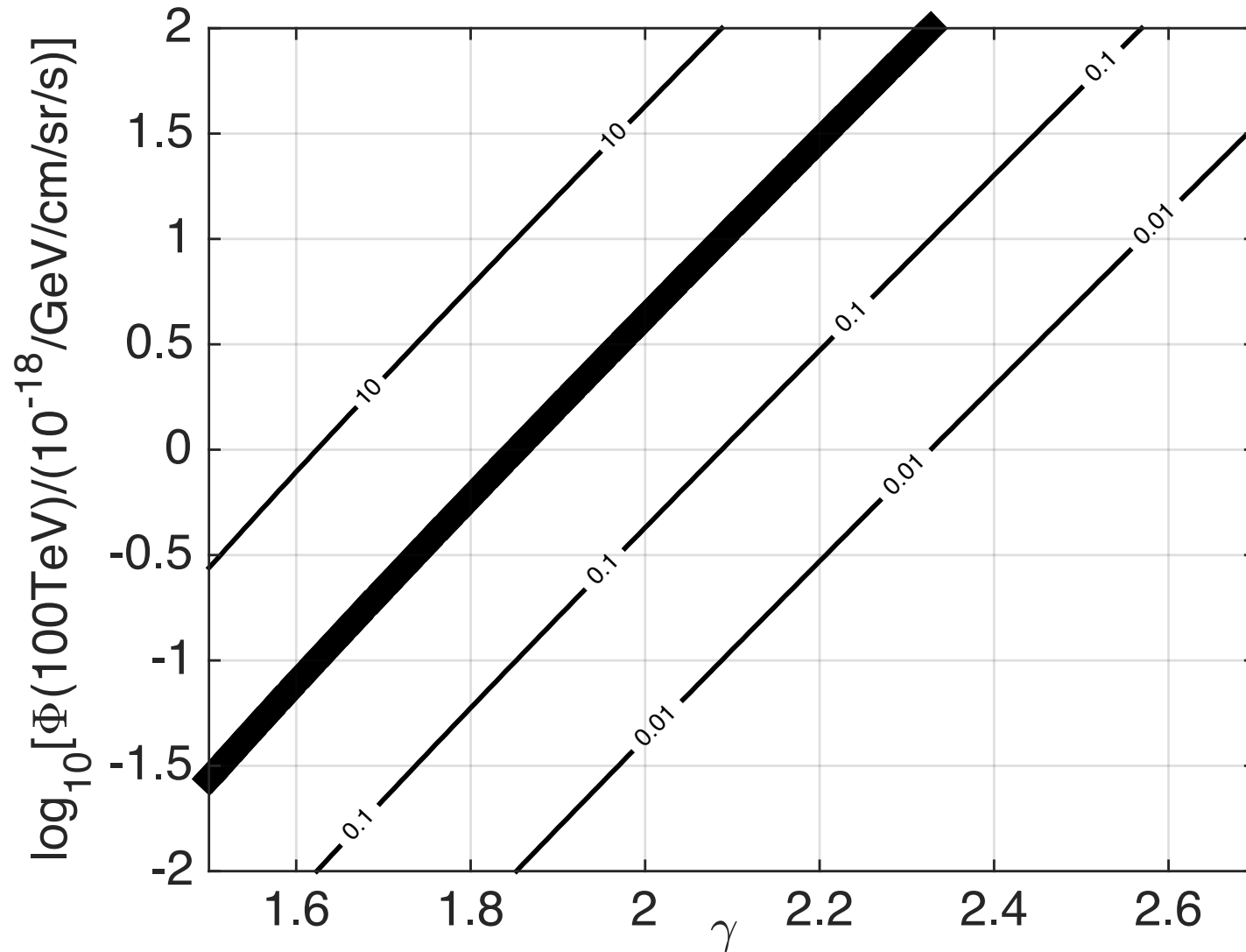
Katz, Budnik, Waxman, **JCAP 0903 (2009) 020**

$$\varepsilon^2 d\dot{n}/d\varepsilon(z=0) = (0.45 \pm 0.15)(\alpha - 1) \times 10^{44} \text{ erg Mpc}^{-3} \text{ yr}^{-1}$$



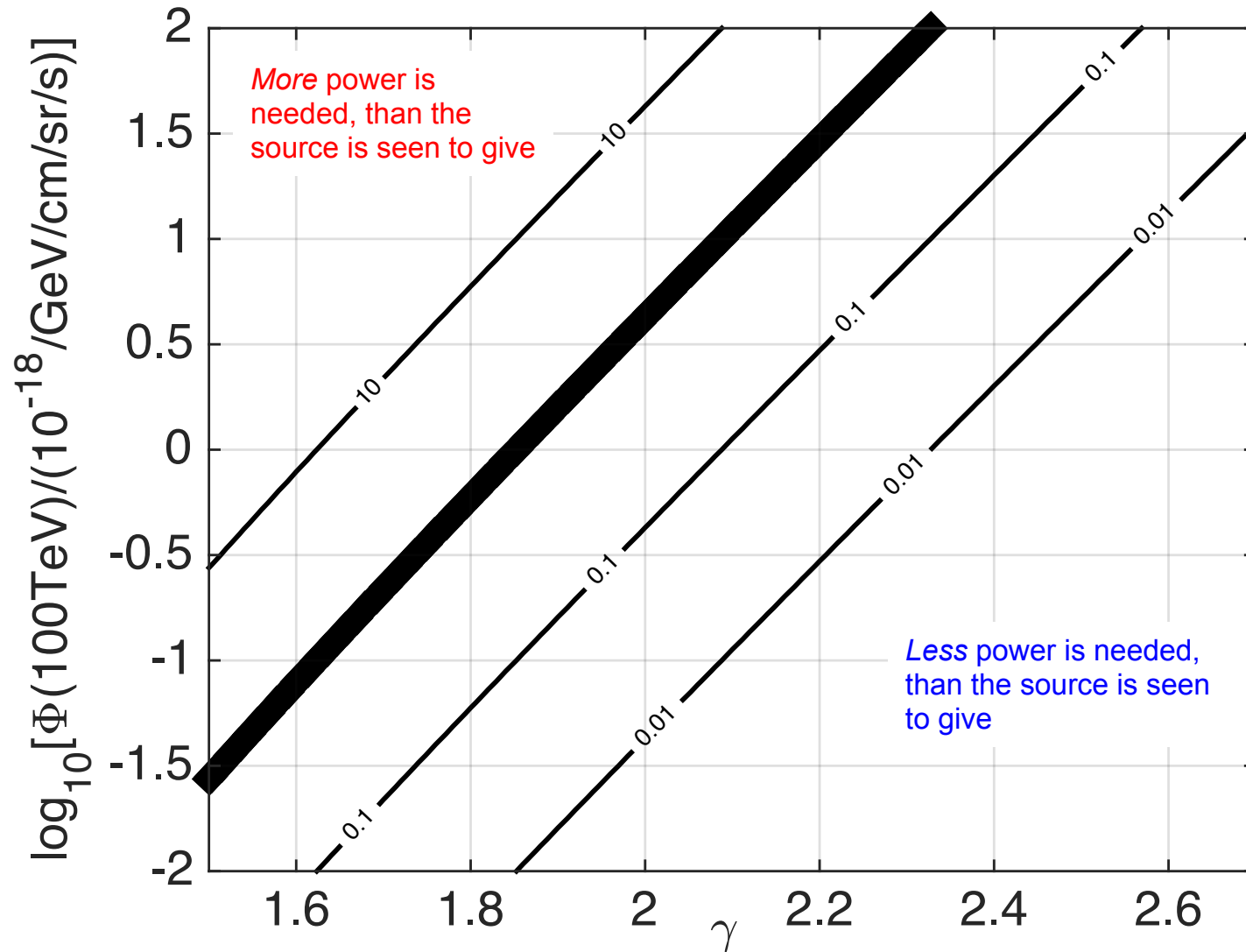
UHECR and the Waxman-Bahcall bound

Contours: required UHECR η x power @ 10^{11} GeV normalized to observed $(0.45 \pm 0.15)(\gamma - 1)10^{44}$ erg/Mpc³/yr



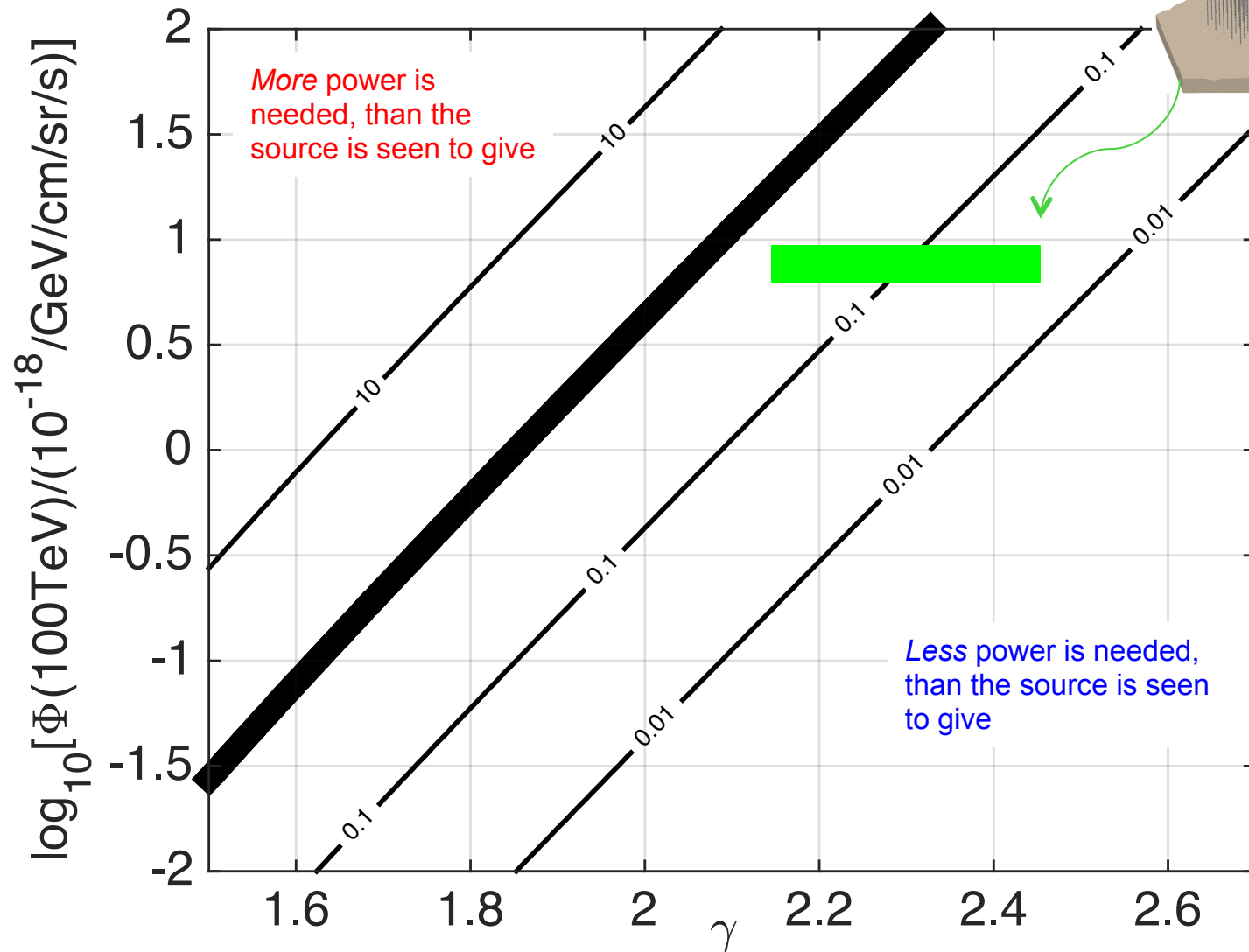
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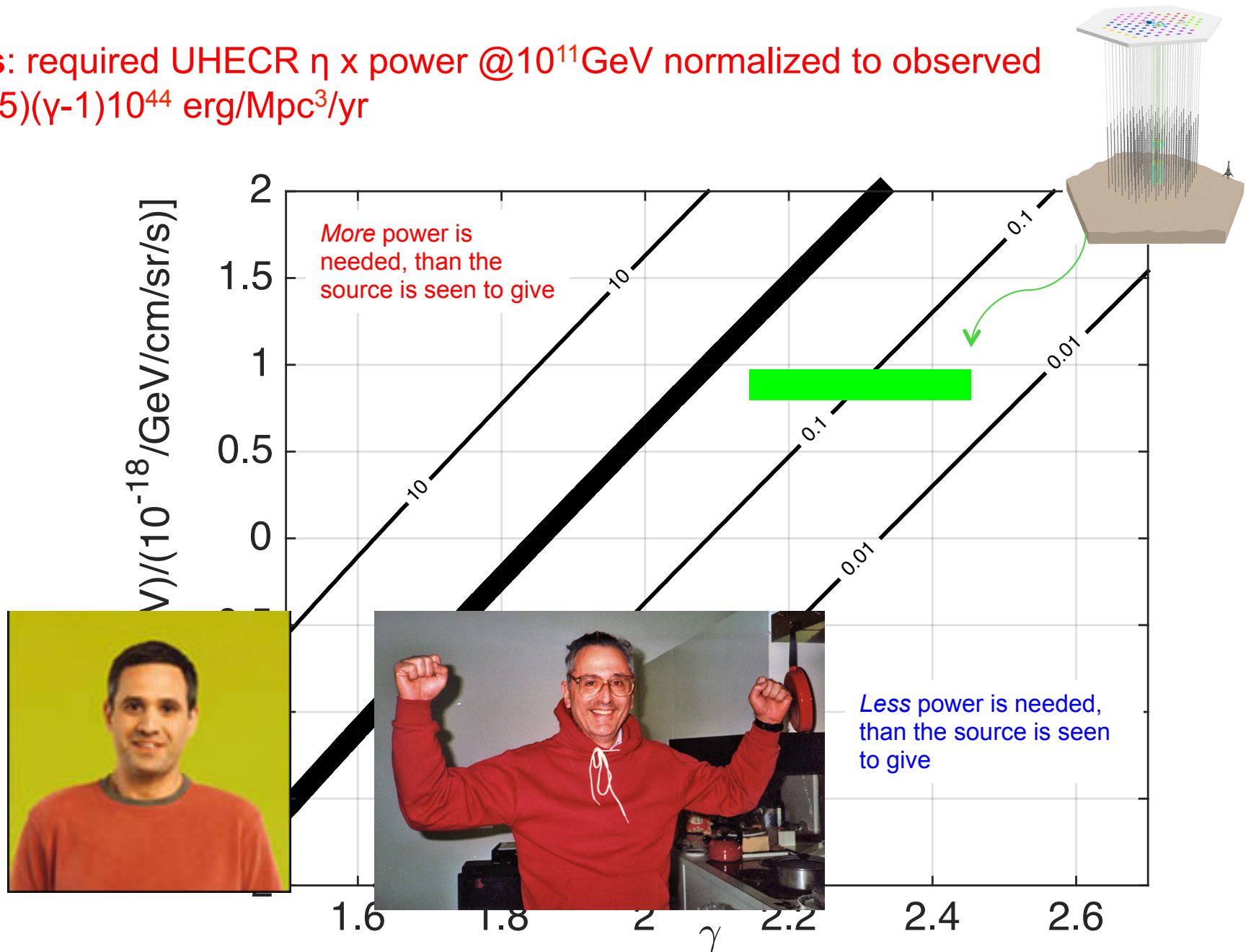
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IceCube's neutrinos and diffuse gamma rays

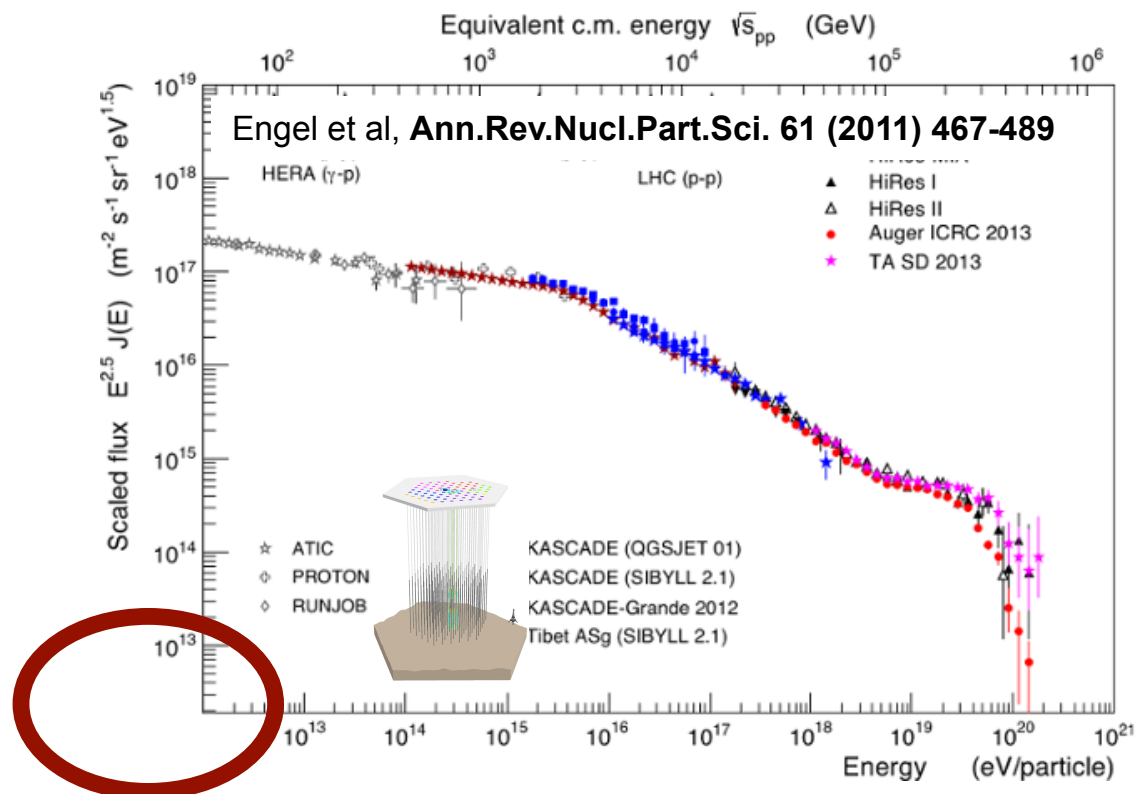
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$$pp \rightarrow pn\pi^+$$

$$pp \rightarrow pp\pi^0$$

$$\pi^+ \rightarrow \mu^+ \nu_\mu \rightarrow e^+ \nu_e \bar{\nu}_\mu \nu_\mu$$

$$\pi^0 \rightarrow \gamma\gamma$$



IceCube's neutrinos and diffuse gamma rays

$$pp \rightarrow pn\pi^+$$

$$pp \rightarrow pp\pi^0$$

$$\pi^+ \rightarrow \mu^+ \nu_\mu \rightarrow e^+ \nu_e \bar{\nu}_\mu \nu_\mu$$

$$\pi^0 \rightarrow \gamma\gamma$$

$\gamma < 2.2$

to not overshoot
diffuse gamma rays

Murase, Ahlers, Lacki, **PRD88 (2013) no.12, 121301**

