

# Energy density of light quark jet using AdS/CFT

Razieh Morad

In Collaboration with Dr. W. A. Horowitz

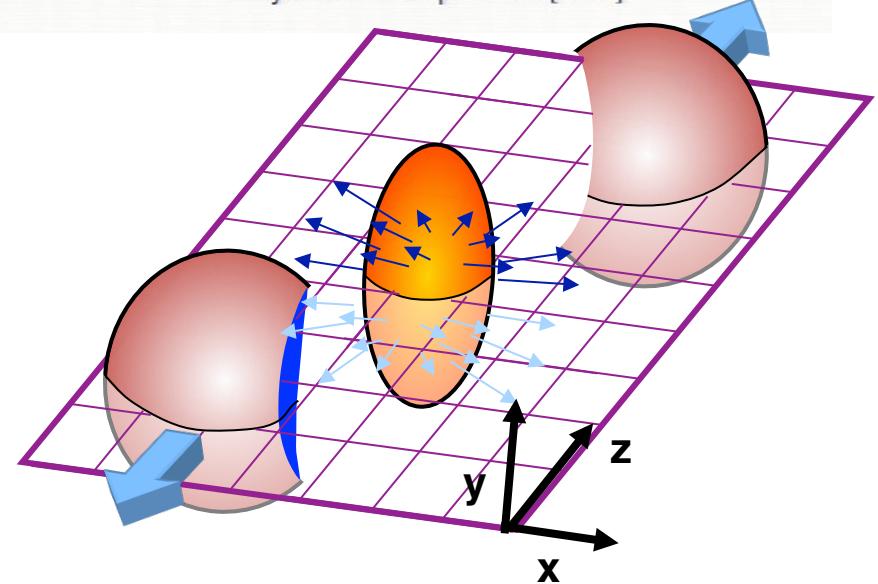
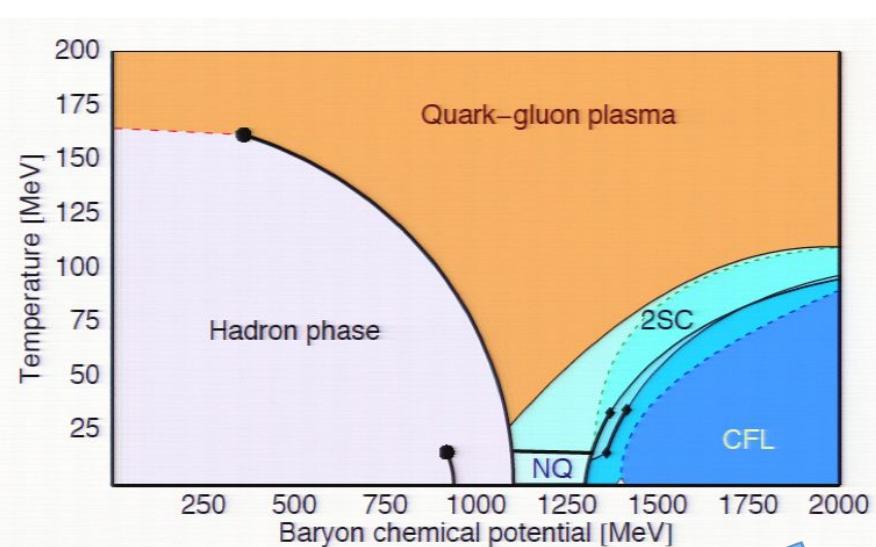
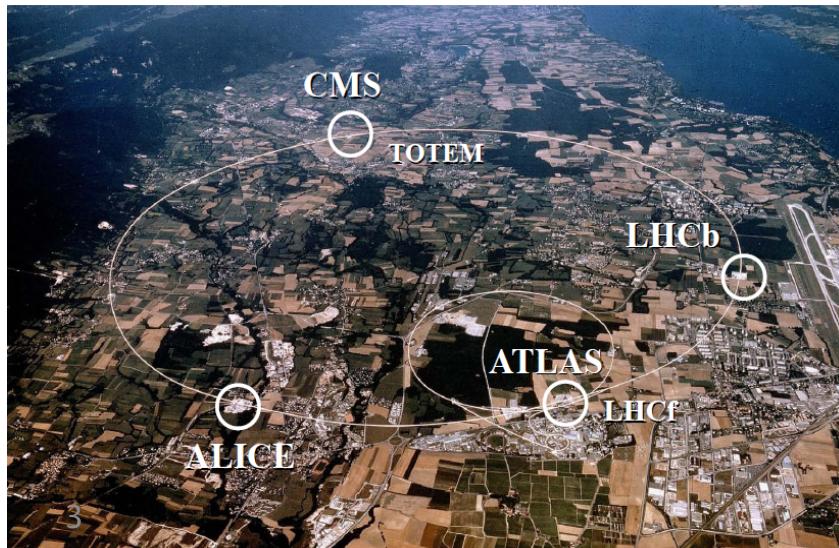
University of Cape Town  
Workshop on High Energy Particle Physics  
iThemba LABS, Gauteng, Johannesburg  
February 2016

# Outline

- Quark-Gluon Plasma
- AdS/CFT Correspondence
- Jets
- Light Quark Jet in String Setup
- N=4 SYM Stress Tensor

# Quark-Gluon Plasma

Quark-Gluon Plasma is formed in Heavy Ion Collision at RHIC and LHC.



# Quark-Gluon Plasma

## Shear viscosity

Hydrodynamics prediction:

$$\frac{\eta}{s} < 0.1 - 0.2$$

Teaney (2003)

Lattice:

$$\frac{\eta}{s} = 0.13 \pm 0.03, \text{ at } T=1.65 T_c$$

Meyer (2007)

Naive pQCD:

$$\frac{\eta}{s} \sim 1$$

N=4 SYM:

$$\frac{\eta}{s} = \frac{1}{4\pi} \approx 0.08$$

Policastro, Son, and Starinets (2001)

AdS/CFT predicts a universal lower bound for the ratio of shear viscosity to entropy.

Kovton, Son and Starinets (2003)

## Rapid thermalization

Chesler and Yaffe (2010)  
Janik et all (2012),(2014)

# AdS/CFT Correspondence

## Maldacena *Conjecture*

Classical gravity on  $\text{AdS}_{d+1}$



Strongly coupled  $d$  - dimensional CFT which lives on  
boundary of  $\text{AdS}_{d+1}$

Maldacena 98

Duality unproven, but many consistency checks performed.

# AdS/CFT Correspondence

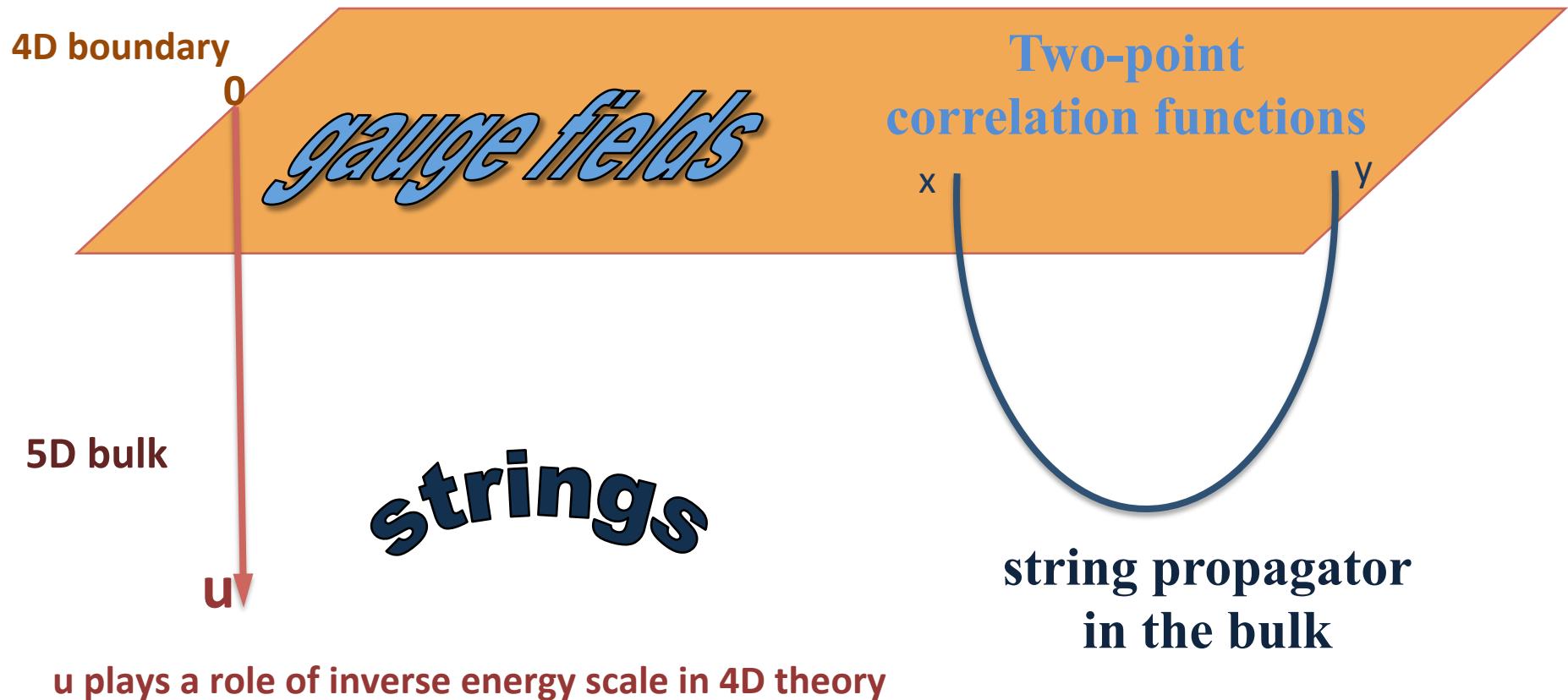
## *AdS/CFT Dictionary*

CFT <sub>d</sub>	<i>AdS</i> <sub>d+1</sub>
Conformal symmetry SO(2,d)	Isometry SO(2,d)
Charges	Charges
Global Symmetry “G”	Gauge Symmetry “G”
Local Operators	Quantum Fields
$\left\langle e^{\int d^4x \phi_0(\vec{x}) O(\vec{x})} \right\rangle_{\text{CFT}}$	$Z[\phi(\vec{x}, z=0) = \phi_0(\vec{x})]$
Partition Function of operator	Classical Action Gubser, Klebanov, Polyakov'98, Witten'98

# AdS/CFT Correspondence

Anti-de-Sitter space (AdS<sub>5</sub>)

$$ds^2 = \frac{dx^\mu dx_\mu + du^2}{u^2}$$



# Probing the hot matter

QGP exists for a few fm, making it impossible to study it using any external probes.

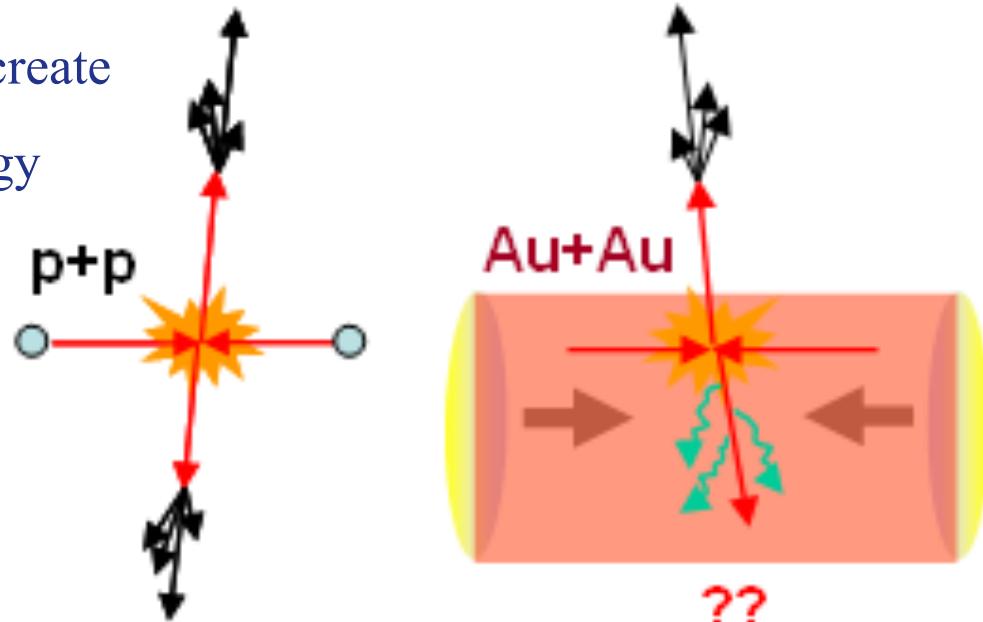
Use self-generated quarks/gluons/photons as probes of the medium

## Hard Probes:

Jets are produced within the expanding fireball and probe the QGP.

Before they become hadronized and create jets, the scattered quarks radiate energy ( $\sim \text{GeV/fm}$ ) in the colored medium.

The presence of hot matter modifies the properties of jets.



# Jet suppression

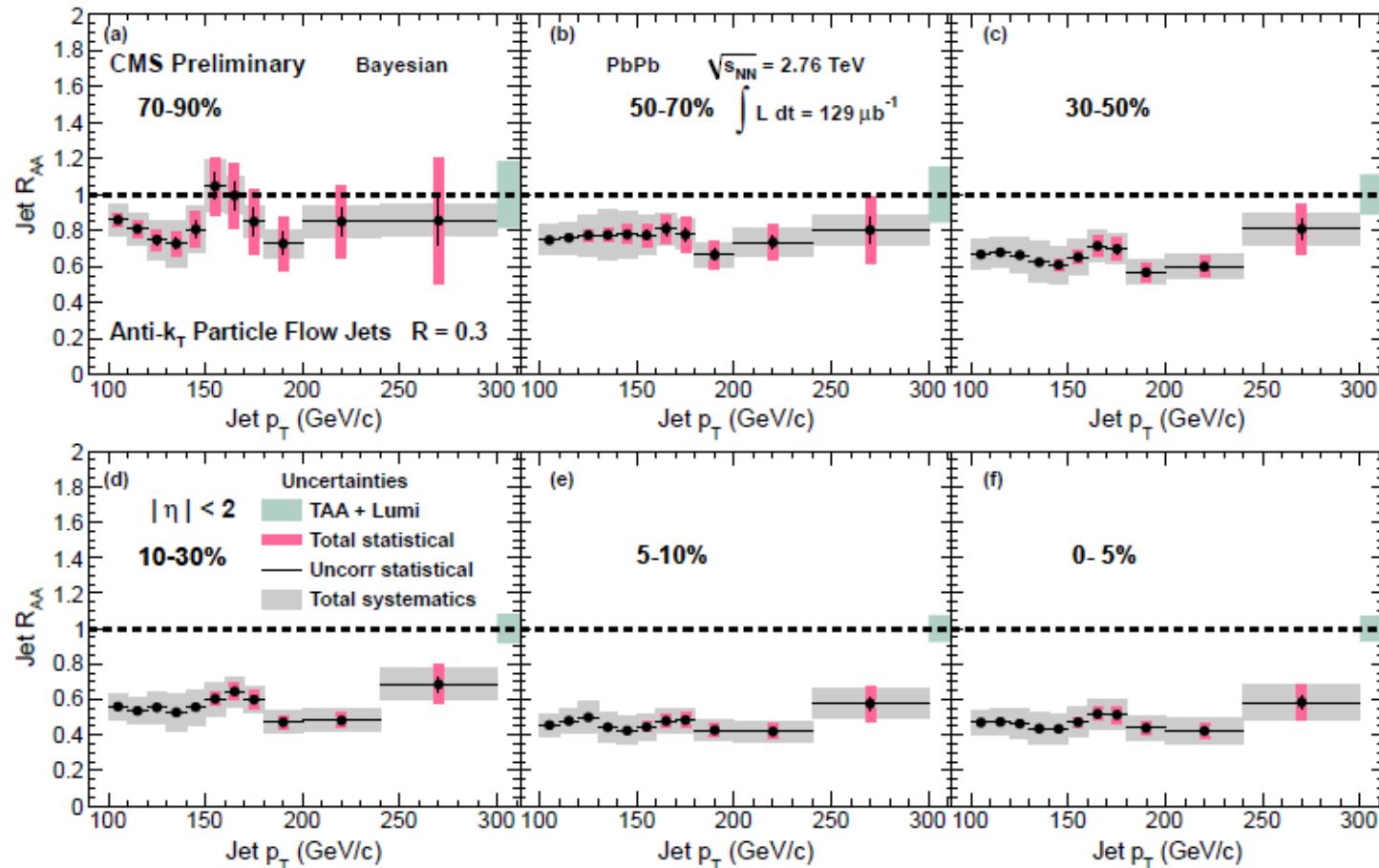
## Nuclear Modification Factor:

$$R_{AA}(p_T) = \frac{\# \text{ of particles observed at } p_T}{\# \text{ of particles expected from pp collisions}}$$

CMS Preliminary data  
(2012)

Naively, if medium has no effect, then  $R_{AA} = 1$ .

$R_{AA} < 1$  means jet quenching



# Light-Quark in string Setup

$N = 4$  Super-Yang-Mills theory in 4d in large  $N_C$  and strong coupling limit  $\lambda$



A Classical supergravity on the 10d  $AdS_5 \times S^5$

Studying the theory at finite temperature



Adding black hole to the geometry:  
AdS-schwarzchild metric

Fundamental quarks in theory



Open strings moving in the 10d geometry

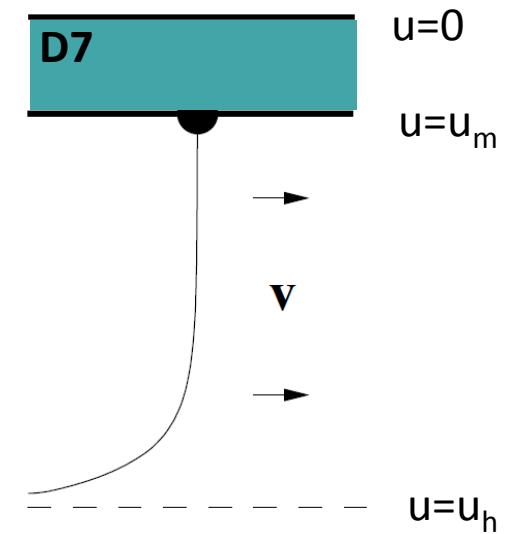
Fundamental quark is dual to a string in the bulk with an endpoint attached to a D7-brane ending at  $u_m$ .

$$\text{For a massive quark at rest: } m_Q = T_0 L^2 \left( \frac{1}{u_h} - \frac{1}{u_m} \right)$$

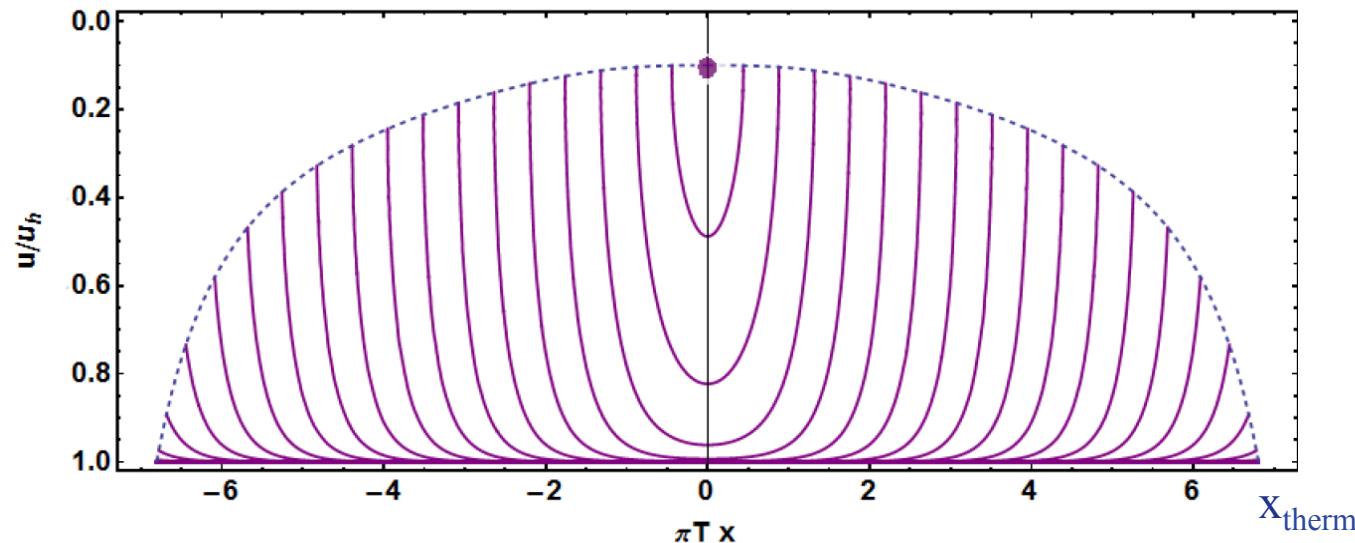
**Light Quarks**



**Falling Strings**



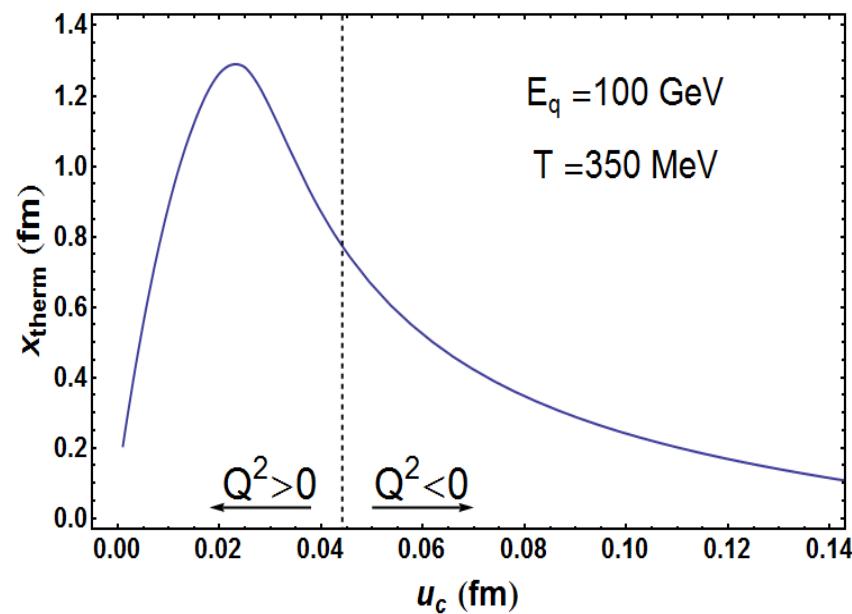
# Thermalization distance



**Light quark dynamics highly depends on the initial conditions of the string:**

There is no known map between the string initial profiles and states in dual field theory.

The only way, is calculating the **energy-momentum tensor of the string** on the boundary and compare with the QCD results.



# SYM Stress-Tensor

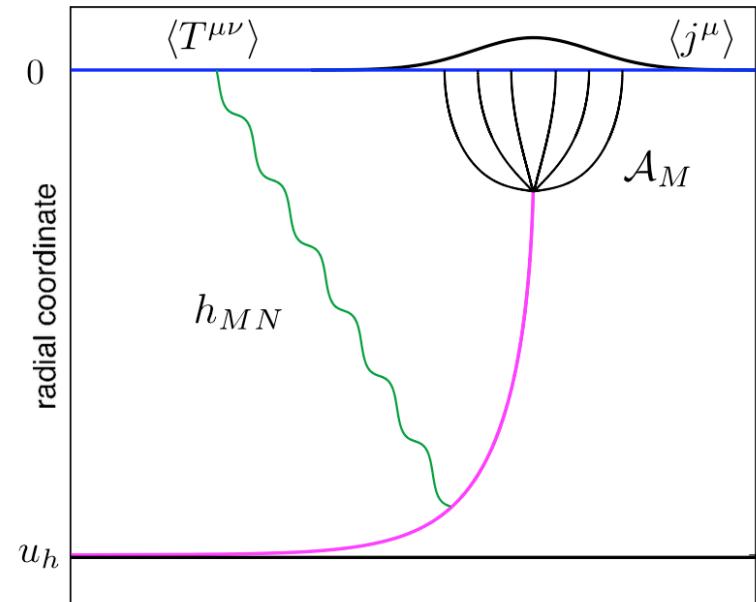
- Presence of string source with the following energy-momentum profile in the bulk perturb the metric:

$$t^{MN} = -\frac{T_0}{\sqrt{-G}} \sqrt{-g} g^{ab} \partial_a X^M \partial_b X^N \delta^3(\mathbf{r} - \mathbf{r}_s)$$

- Metric perturbation  $h_{MN}$ :  $G_{MN} = G_{MN}^{(0)} + h_{MN}$
- Linearized Einstein equation for  $h_{MN}$ :

$$\begin{aligned} & -D^2 h_{MN} + 2D^P D_{(M} h_{N)P} - D_M D_N h + \frac{8}{L^2} h_{MN} \\ & + (D^2 h - D^P D^Q h_{PQ} - \frac{4}{L^2} h) G_{MN}^{(0)} = 2\kappa_5^2 t_{MN}, \end{aligned}$$

- On-shell gravitational action:  $S_G = \frac{1}{2\kappa_5^2} \int d^5x \sqrt{-G} \left( \mathcal{R} + \frac{12}{L^2} \right) + S_{GH}$
- SYM energy-momentum tensor:  $T^{\mu\nu}(x) = \frac{2}{\sqrt{-g}} \frac{\delta S_G}{\delta g_{\mu\nu}(x)}$



**h<sub>MN</sub> has 15 degrees of freedom**  $\longleftrightarrow$  **T<sub>\mu\nu</sub> has 5 degrees of freedom**

?

# Gauge-Invariants

It is possible to construct gauge invariant quantities out of linear combinations of  $h_{MN}$  and its derivatives:

- ✓ There are just 5 of them.
- ✓ Their equation of motions are completely decoupled.

The gauge invariant which is transformed as scalar under rotation  $Z$ , can give us the energy density of the SYM stress tensor on the boundary:

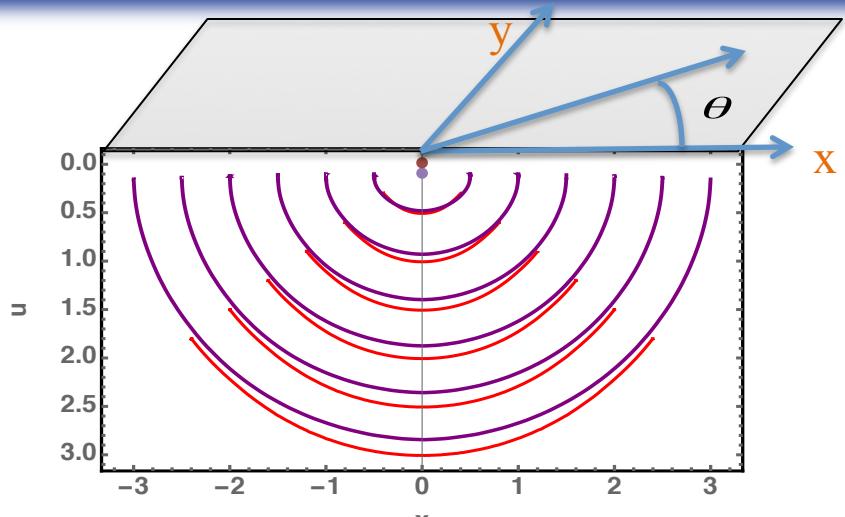
$$Z'' + AZ' + BZ = S$$

$$\begin{aligned} A &\equiv -\frac{24 + 4q^2u^2 + 6f + q^2u^2f - 30f^2}{uf(u^2q^2 + 6 - 6f)}, \quad B \equiv \frac{\omega^2}{f^2} + \frac{q^2u^2(14 - 5f - q^2u^2) + 18(4 - f - 3f^2)}{u^2f(q^2u^2 + 6 - 6f)} \\ \frac{S}{\kappa_5^2} &\equiv \frac{8}{f}t'_{00} + \frac{4(q^2u^2 + 6 - 6f)}{3uq^2f}(q^2\delta^{ij} - 3q^iq^j)t_{ij} + \frac{8i\omega}{f}t_{05} + \frac{8u[q^2(q^2u^2 + 6) - f(12q^2 - 9f'')]}{3f^2(q^2u^2 - 6f + 6)}t_{00} - \frac{8q^2u}{3}t_{55} - 8iq^it_{i5} \end{aligned}$$

Asymptotic behavior of  $Z$ :  $Z(u) = Z_{(2)}u^2 + Z_{(3)}u^3 + \dots$

Energy density:  $\mathcal{E} = -\frac{L^3}{8\kappa_5^2}Z_{(3)}$

# Boundary Jet Energy Density in AdS<sub>5</sub>

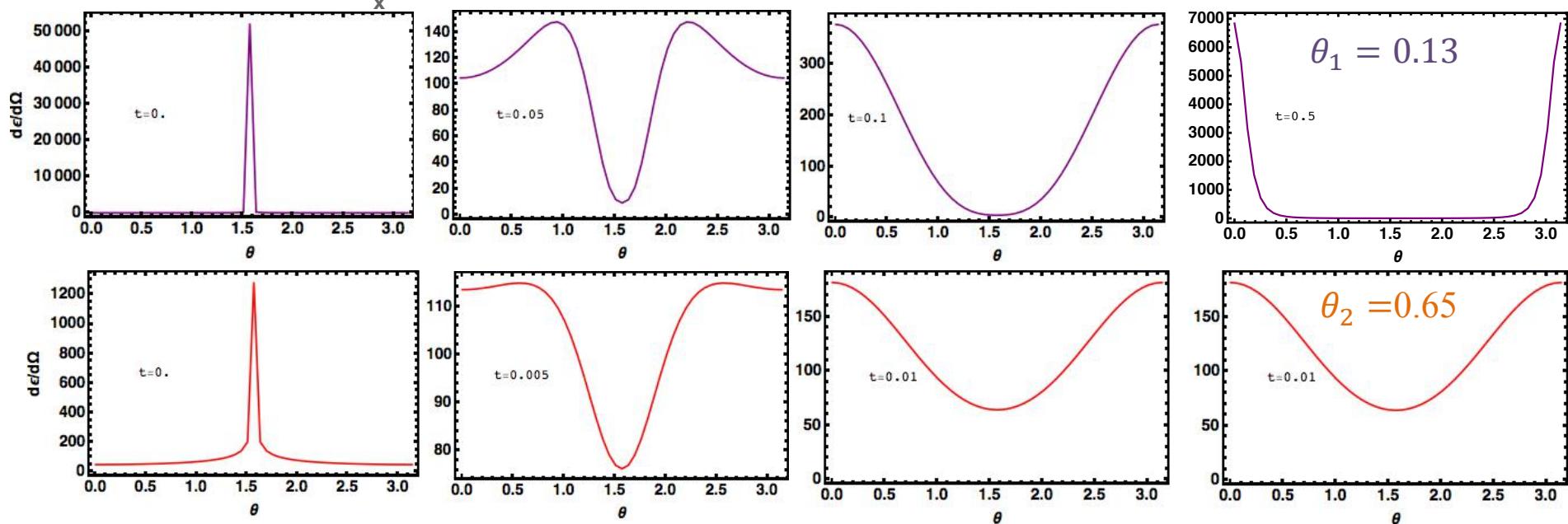


String 1 created at  $u_c=0.1$  :

$$E_q = 100 \text{ GeV}, Q^2 = 176 \text{ GeV}^2$$

String 2 created at  $u_c=0.01$  :

$$E_q = 100 \text{ GeV}, Q^2 = 6000 \text{ GeV}^2$$



One can define the opening angle of jet:  $\theta_j = \text{ArcTan}[\frac{\sqrt{Q^2}}{E_q}]$

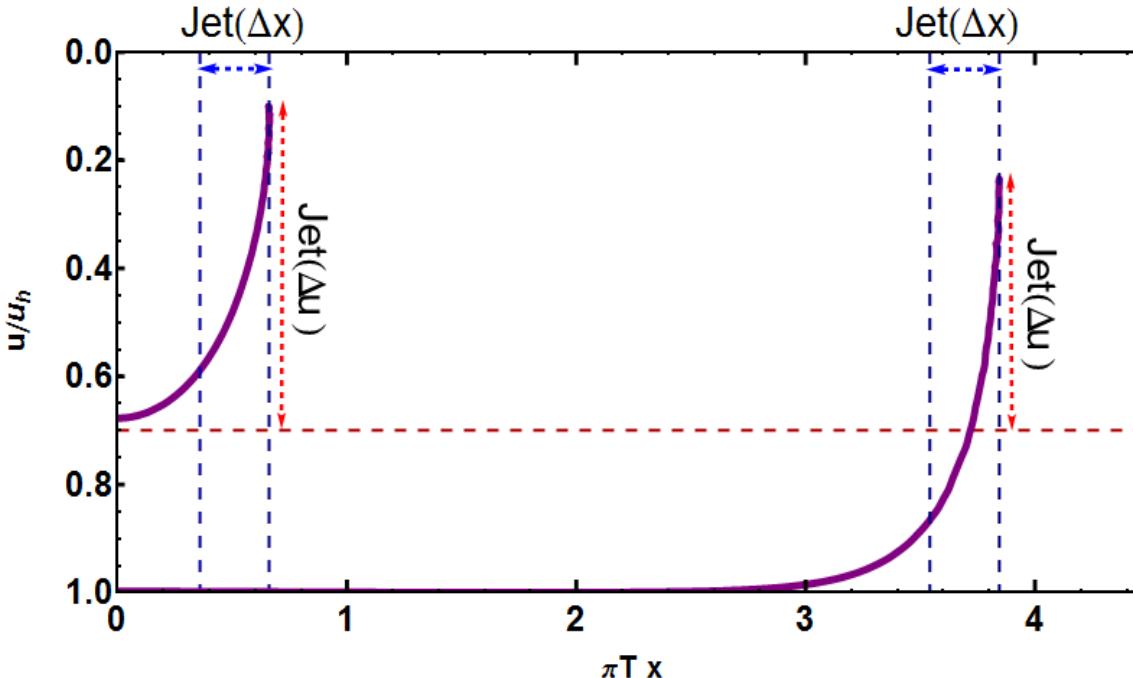
# Thank you



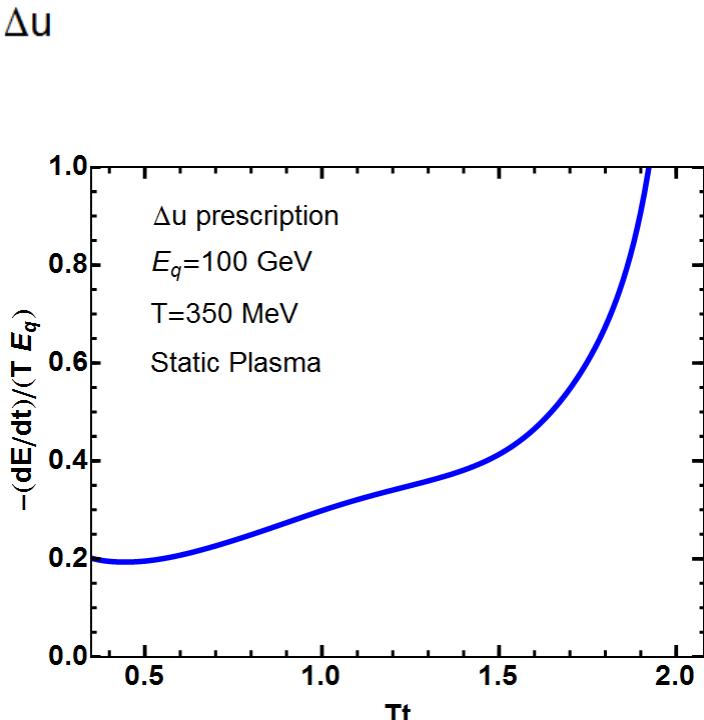
# Back up slides

# Jet Prescription in AdS/CFT

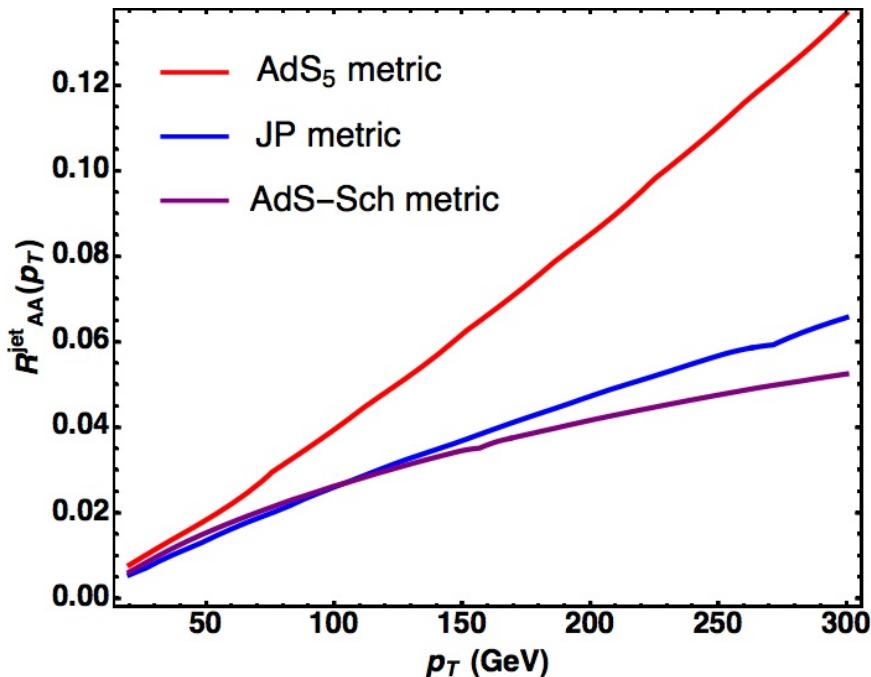
Our prescription of jet in AdS/CFT: based on separation of hard and soft sectors



Jet energy lose rate:



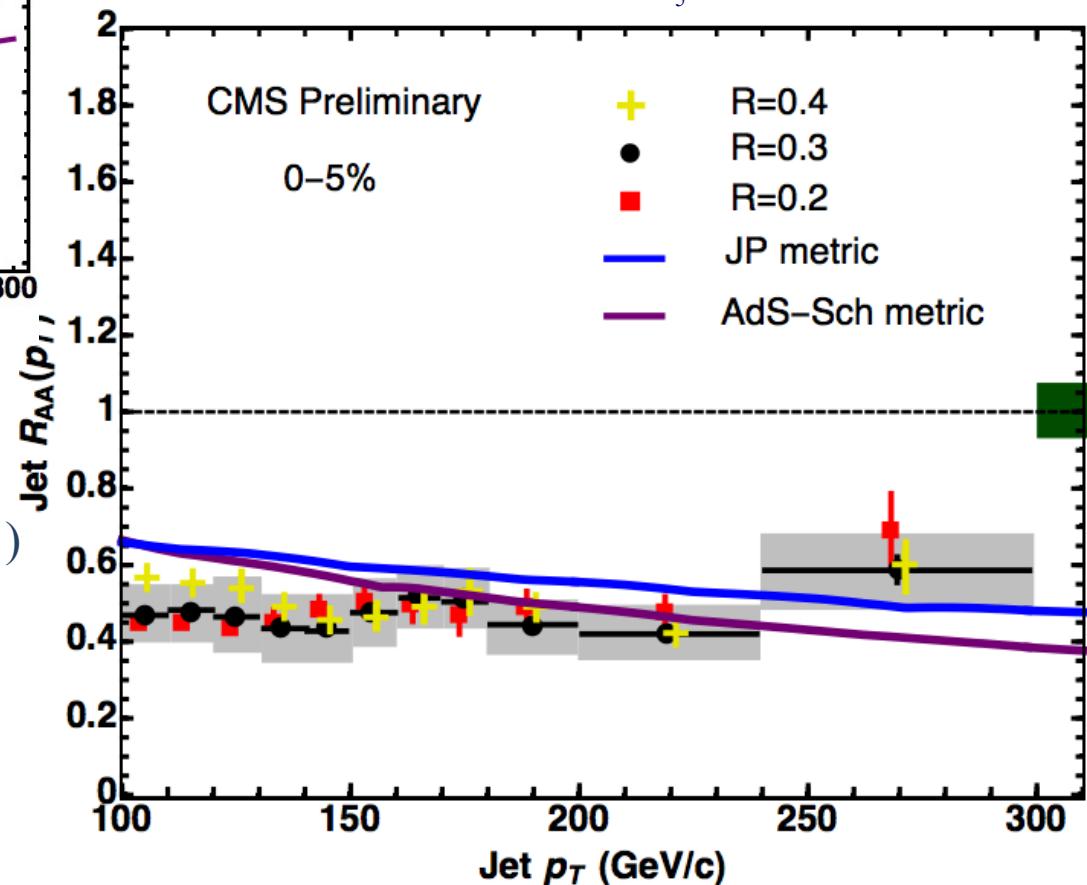
# Jet Nuclear Modification Factor



$$\Delta E^{\text{sub, ren}}(\pi, L, T) \equiv \Delta E^{\text{medium}}(\pi, L, T) - \Delta E^{\text{vacuum}}(\pi, L, T)$$

We define a renormalized  $R_{AA}$  in AdS/CFT:

$$R_{AA}^{jet}(p_T)_{AdS/CFT} \equiv \frac{R_{\text{medium}}^{jet}(p_T)}{R_{AdS_5}^{jet}(p_T)}$$



R. Morad and W. A. Horowitz,  
JHEP 11 (2014) 017