# SEARCH for INTRINSIC CHARM in Z/y+c PRODUCTION in p-p at ATLAS



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in collaboration with
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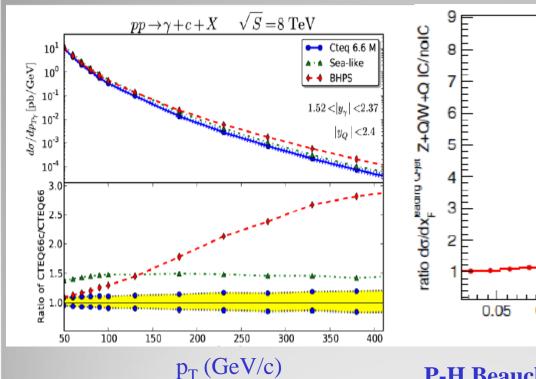
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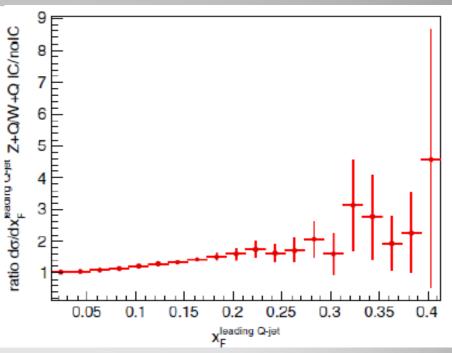
### **OUTLINE**

- 1. Motivation of this study.
- 2.Intrinsic charm (*IC*) and ratio  $\sigma_{pp}(\gamma+c)/\sigma_{pp}(\gamma+b)$  as a function of  $p_{T}$ .
- 3. Ratio  $\sigma_{pp}(Z+c)/\sigma_{pp}(Z+b)$  as a function of the IC probability.
- 4. Ratio  $\sigma_{pp}(Z+c)/\sigma_{pp}(Z+b)$  as a function of  $p_T$  and the highest possible *IC* signal .
- 5. Summary.

## $PP \rightarrow \gamma + c + X$ , $s^{1/2} = 8 \ TeV \ PP \rightarrow Z/W + Q + X$ , $s^{1/2} = 8 \ TeV$

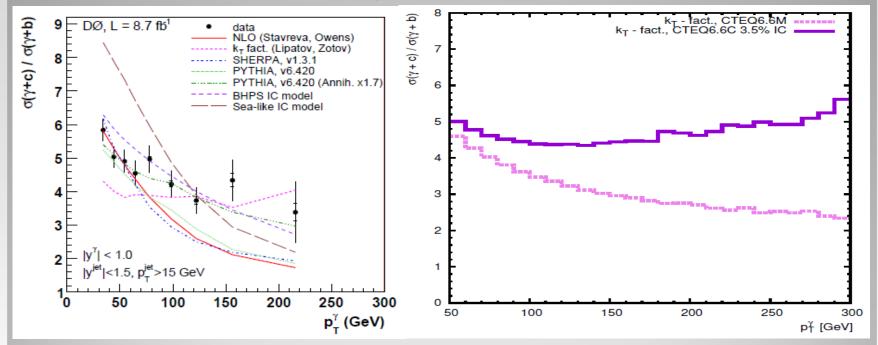


V.A.Bednyakov, M.A.Demichev, G.L., T.Stavreva, M.Stockton, Phys.Lett. B728, 602(2014).



P-H Beauchemin, V.A.Bednyakov, G.L., Yu.Yu. Stepanenko, Phys.Rev. D92 034014 (2015). Q = c or b.

## New results on $\gamma$ + c production in p-p

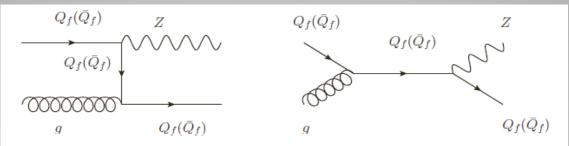


Ratio of differential cross sections of  $\gamma$  +c and  $\gamma$  +b productions as function of photon transsverse momentum Left: D0 data from TEVATRON and theoretical calculations

D0 Collaboration, Phys. Lett. B 719, 354 (2013); arXiv:1210.5033 [hep-ex].

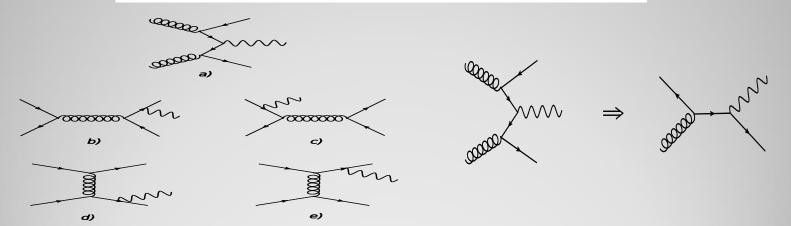
Right: our latest predictions with the IC (top line) and without it (bottom line)

## $pp \rightarrow Z + heavy flavour jets$



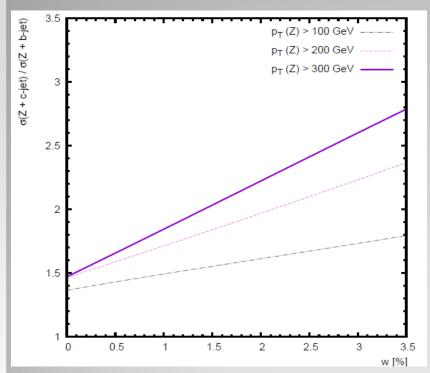
 $Q_f = c.b$  and  $Q'_f = b, c$  respectively.

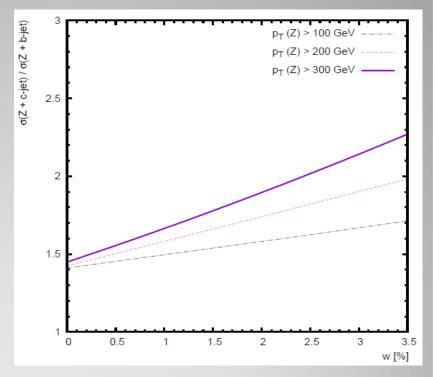
Feynman diagram for the process  $Q_f(\bar{Q}_f)g \to ZQ_f(\bar{Q}_f)$ 



Feynman diagrams, which contribute to the Z+Q production in p-p collision

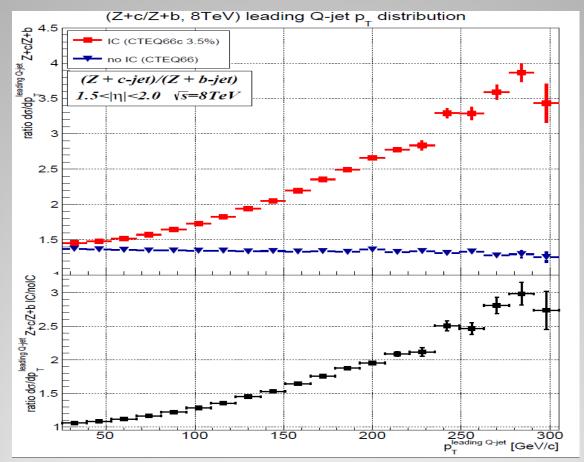
### $PP \rightarrow Z + c(b) + X$





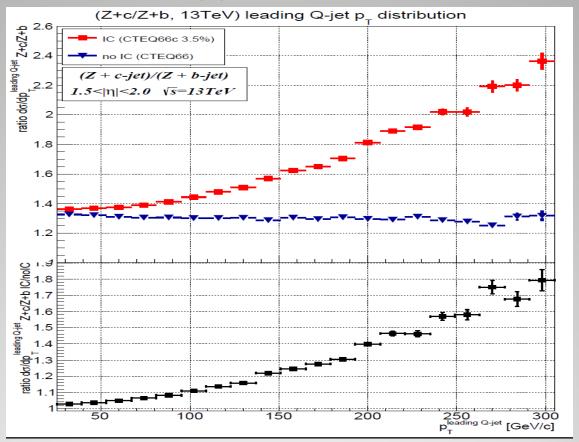
Ratio  $\sigma_{pp}(Z+c)/\sigma_{pp}(Z+b)$  as a function of the *IC* probability w at  $p_T^Z > 100$  GeV/c (bottom dash-dotted line),  $p_T^Z > 200$  GeV/c middle dotted line) and  $p_T^Z > 300$  GeV/c (top solid line). Left:  $s^{1/2} = 8$  TeV; right:  $s^{1/2} = 13$  TeV.

#### $PP -> 8 \ TeV \ at \ s^{1/2} = 8 \ TeV$



Ratio  $\sigma_{pp}(Z+c)/\sigma_{pp}(Z+b)$  as a function of  $p_T$  of leading jet. Top: with the *IC* probability w=3.5% (red points) and without *IC* (blue points). Bottom: ratio of red points to blue ones.

## $PP -> Z + c(b) + X at s^{1/2} = 13 TeV$



Ratio  $\sigma_{pp}(Z+c)/\sigma_{pp}(Z+b)$  as a function of  $p_T$  of leading jet. Top: with the *IC* probability w=3.5% (red points) and without *IC* (blue points). Bottom: ratio of red points to blue ones.

### **SUMMARY**

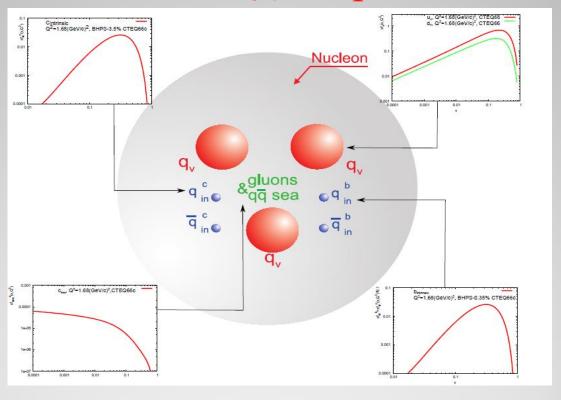
- 1. We suggest to measure the ratio of cross sections  $(\mathbf{Z}/\gamma + \mathbf{c})/(\mathbf{Z}/\gamma + \mathbf{b})$  as a function of  $\mathbf{p_T}^{\mathbf{Z}}$  or  $\mathbf{p_T}^{\text{leadingQ}}$
- 2. At  $p_T^Z > 300$  GeV/c the *IC* signal is maximum and this ratio becomes about 2.5-2.8 at the *IC* probability w = 3.5 %.
- 3. We calculate this ratio within the  $k_T$  –factorization of QCD as a function of w at different intervals of  $p_T^Z$ , when 0 < w < 3.5%.
- 4. We estimate the uncertainty of our predictions, it is less than 10%.
- 5. The similar predictions are made within the PYTHIA

# THANK YOU VERY MUCH FOR YOUR ATTENTION!

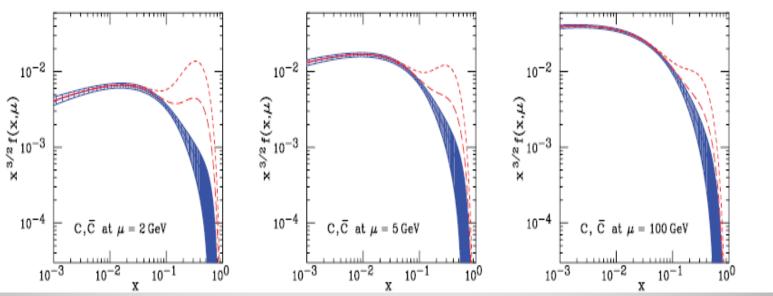


BHPS model: S.J. Brodsky,P. Hoyer, C. Peterson and N.Sakai, Phys.Lett.B9(1980) 451; S.J. Brodsky, S.J. Peterson and N. Sakai, Phys.Rev. D23 (1981) 2745.

Intrinsic  $Q\bar{Q}$  in proton



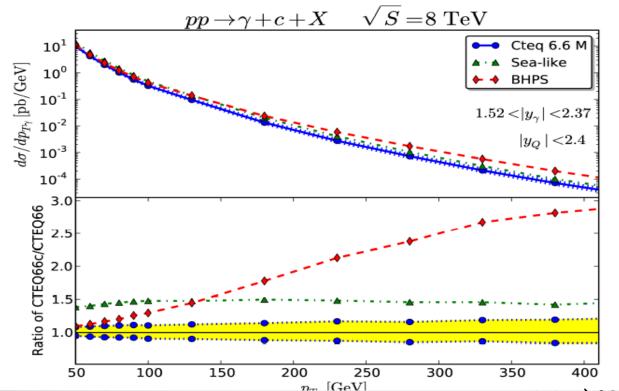
#### **CHARM QUARK DISTRIBUTIONS IN PROTON**



Charm quark distributions within the BHPS model. The three panels correspond to the renormalization scales  $\mu$ =2,5,101 GeV respectively. The long-dashed and the short-dashed curves correspond to  $\langle x_{c\bar{c}} \rangle$ =0.57%2.% respectively using the PDF CTEQ66c. The solid curve and shaded region show the central value and uncertainty from CTEQ6.5, which contains no *IC*.

There is an enhancement at x>0.1 due to the IC contribution

## IC signal in $pp \rightarrow \gamma + c(jet) + X$



V.A.Bednyakov, M.A.Demichev, G.L., T.Stavreva, M.Stockton, hep-ph/1305.3548

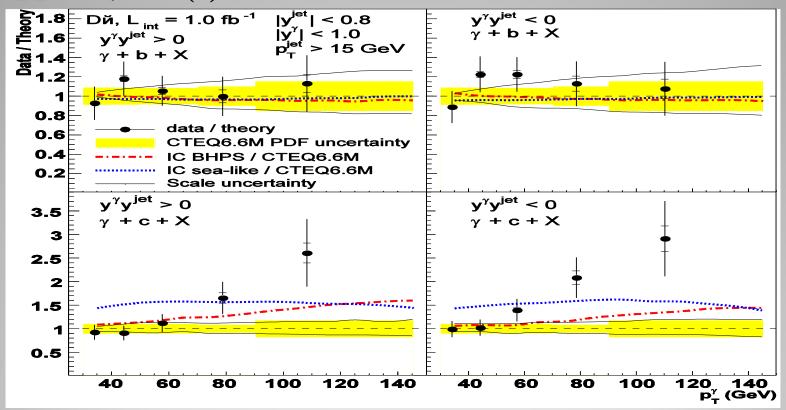
Phys.Lett. B728 (2014) 602.

 $p_T$ —distribution of photons produced in  $pp \rightarrow \gamma + c(jet) + X$ The blue line is calculation without the IC. The red curve includes the IC, its **probability is about 3.5** % (top). The ratio of spectra with and without the IC

probability is about 3.5 % (top). The ratio of spectra with and without the IC The IC signal is about 200%-250% at  $p_T \sim 150-200$  where the cross section

is about 20-80 fb (400-3200 events) and can be measuered

# $pp \longrightarrow \gamma + c(b) + X$ D0 experiment at Tevatron $S^{1/2} = 1.9$ TeV



The data-to-theory ratio of cross sections as a function of  $P_T^{\gamma}$  for  $pp \rightarrow \gamma + c(b) + X$ . There is the three time excess of the data above the theory for  $\gamma + c$  at  $p_T > 15$  GeV c. It stimulates us to study  $pp \rightarrow \gamma + c(b) + X$ 

# INTRINSIC HEAVY QUARK STATES

Two types of parton contributions

The extrinsic quarks and gluons are generated on a short time scale in association with a large transverse-momentum reaction.

The intrinsic quarks and gluons exist over a time scale independent of any probe momentum, they are associated with the bound state hadron dynamics.

$$P(x_1,...,x_5)=N_5\delta\left(1-\sum_{i=1}^5 x_i\right)M_p^2-\sum_{i=1}^5 \frac{m_i^2}{x_i}$$

# INTRINSIC HEAVY QUARK DISTRIBUTION IN PROTON

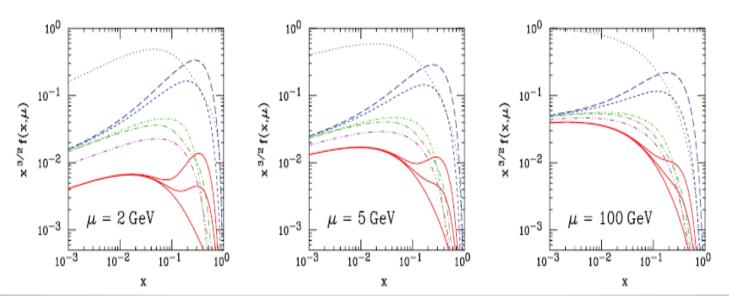
Integrating  $P(x_1,...x_5)$  over  $dx_1...dx_5$  and neglecting of all quark masses except the charm quark mass we get

$$P(x_{5}) = \frac{1}{2}N_{5}x_{5}^{2} \left[ \frac{1}{3}(1-x_{5})(1+10x_{5}+x_{5}^{2}) + 2x_{5}(1+x_{5})\ln\left(\frac{1}{x_{5}}\right) \right]$$

Where  $N_5 = N_5/m_{4,5}^4$  normalization constant. Here  $m_4 = m_5 = m_c = m_c$  is the bar mass of the charmed quark.  $N_5$  determines some probability  $W_{IQ}$  to find the Fock state |uudQ| in the proton.

One can see qualitatively that  $P(x_5)$  vanishes at  $x_5 \rightarrow 0$  and  $x_5 \rightarrow 1$  and has an enhancement at  $0 < x_5 < 1$ 

# COMPARISON OF LIGHT AND HEAVY QUARK DISTRIBUTIONS IN PROTON



The dotted line is the gluon distribution, the blue long-dashed curve is the valence u-distribution, the blue short-dashed line is the valence d-distribution, the green long-dashed-dotted line is the intrinsic  $\bar{u}$ , the short dashed-dotted line is the intrinsic  $\bar{d}$  distribution, the dashed-dot-dotted is the intrinsic  $S=\bar{S}$  and the solid curves are  $C=\bar{C}$  with **no IC** (lowest) and with  $IC\langle x_{c\bar{c}}\rangle=0.5$  % at x>0.2 respectively. It is shown that IC contribution is larger than  $u,\bar{d},\bar{s}$  at x>0.2

#### PRODUCTION OF HEAVY FLAVOURS IN HARD P-P COLLISIONS

$$E\frac{d\sigma}{d^3p} = \sum_{i,i} \int d^2k_{iT} \int d^2k_{jT} \int_{x_i^{\min}}^1 dx_i \int_{x_j^{\min}}^1 dx_j f_i(x_i, k_{iT}) f_j(x_j, k_{jT}) \frac{d\sigma_{ij}(\hat{s}, \hat{t})}{d\hat{t}} \frac{D_{i,j}^h(z_h)}{\pi z_h}$$

$$x_i^{\min} = \frac{x_T \cot(\frac{\theta}{2})}{2 - x_T \tan(\frac{\theta}{2})} \qquad x_F \equiv \frac{2p_z}{\sqrt{s}} = \frac{2p_T}{\sqrt{s}} \frac{1}{\tan \theta} = \frac{2p_T}{\sqrt{s}} \sinh(\eta)$$

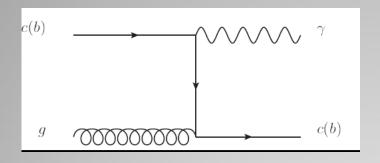
$$x_i^{\min} = \frac{x_R + x_F}{2 - (x_R - x_F)}$$

$$x_R = 2p/\sqrt{s}$$

One can see that  $x_i \ge x_F$  If  $x_F > 0.1$  then,  $x_i > 0.1$  and the conventional sea heavy quark (extrinsic) contributions are suppressed in comparison to the intrinsic ones.  $x_F$  is related to  $p_T$  and  $p_T$ . So, at certain values of these

 $x_F$  is related to  $P_T$  and  $rac{1}{1}$ . So, at certain values of these variables, in fact, there is no conventional sea heavy quark (extrinsic) contribution. And we can study the IQ contributions in hard processes at the certain kinematical region.

#### PHOTON (DI-LEPTON) AND c(b)-JETS PRODUCTION IN P-P



for the process  $c(b)+g\rightarrow \gamma+c(b)$  the process  $c(b)+g\rightarrow \gamma/Z^0+c(b)$ 

Fig.a. Feynman diagram

Fig.b. Feynman graph for

$$x_{F} = \frac{2p_{T}}{S^{1/2}} sh(\eta); p_{T\gamma} = -p_{Tc}. \quad x_{c(b)} = \frac{m_{Tc}^{2}}{x_{o}S} + x_{c(b)}^{f}$$

$$oxed{x_{c(b)}} = rac{oldsymbol{m_{l^+l^-}}}{oldsymbol{\mathcal{X}_g} oldsymbol{S}} + oldsymbol{\mathcal{X}_{c(b)}^f}$$

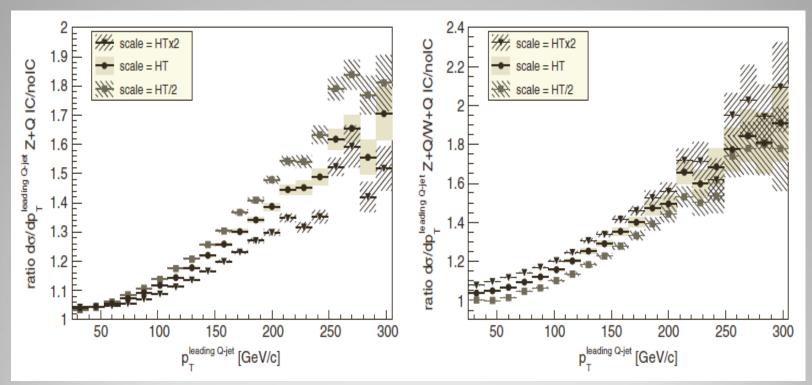
To observe the IC

for Fig.a

$$x_c \ge x_F > 0.1$$

for Fig.b
$$x_{c(b)} = \frac{m_{f+f-}^2}{x_g S} + x_{c(b)}^f > 0.$$

### Scale uncertainty for Z+Q and Z+Q/W+Q

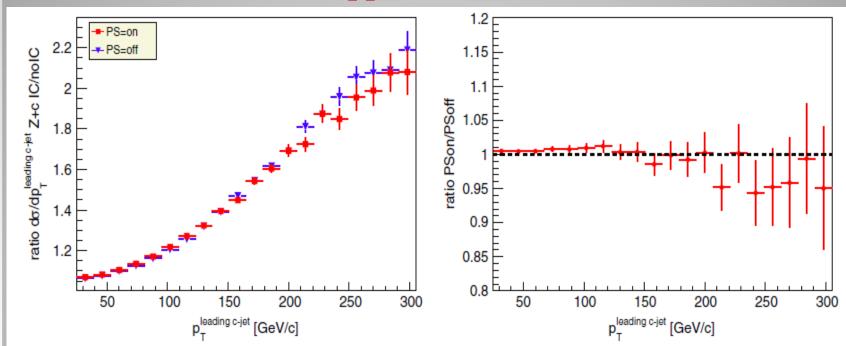


**Left: Z+Q with IC and without IC at different scales** 

Right: Z+Q/W+Q with IC and without IC at different

scales

# Inclusion of parton shower by Z+Q production in pp at $s^{1/2} = 8 \text{ TeV}$

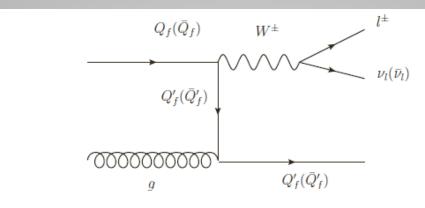


**Left:**  $p_T$  – spectra with parton shower (red points) and without

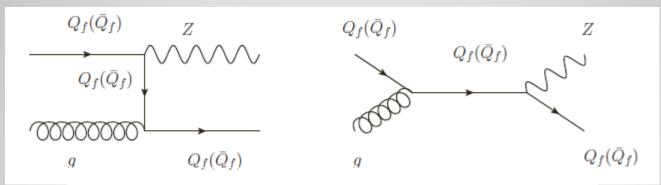
it (blue points) using the PDF of type CTEQ66c (3.5% of IC)

**Right:** ratio of red points to blue points

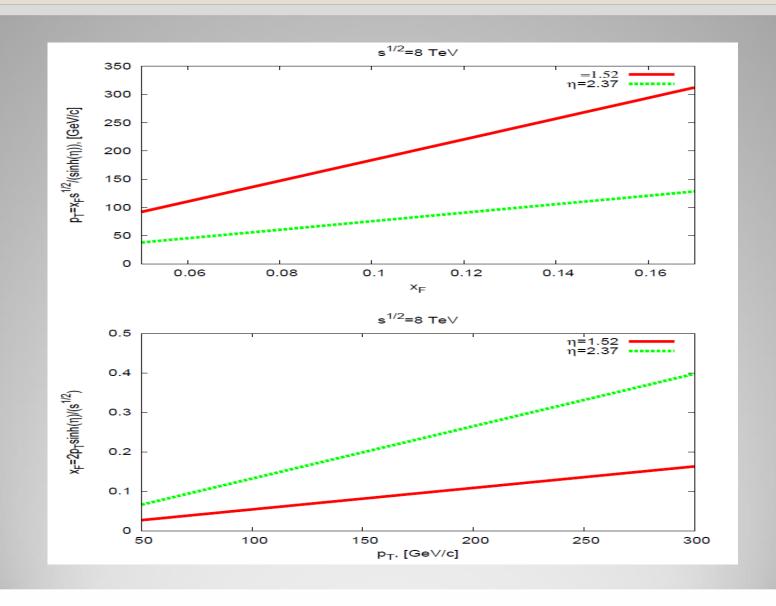
### $pp \rightarrow W/Z + heavy flavour jets$

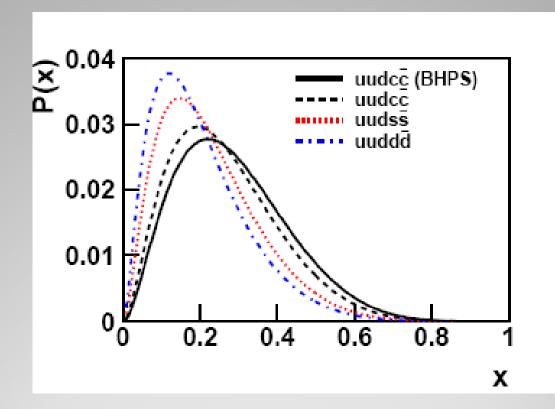


The LO Feynman diagrams for the process  $Q_f(\bar{Q}_f)g \to W^{\pm}Q'_f(\bar{Q}'_f)$ , where  $Q_f = c.b$  and  $Q'_f = b, c$  respectively.

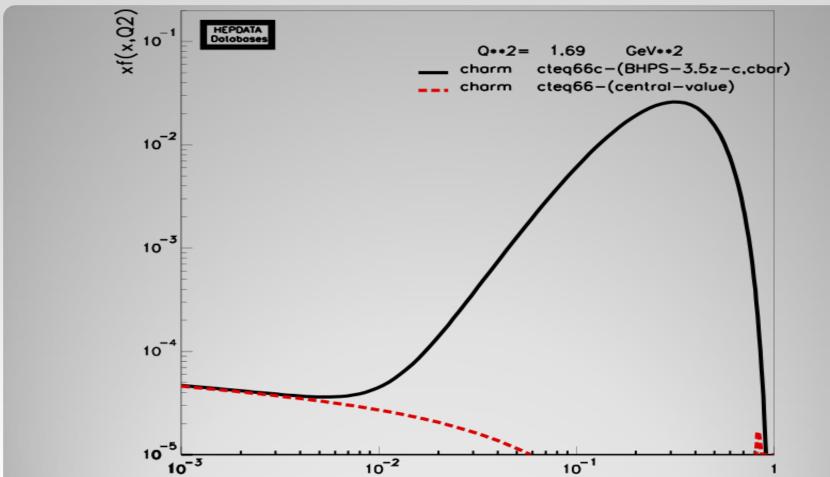


Feynman diagram for the process  $Q_f(\bar{Q}_f)g \to ZQ_f(\bar{Q}_f)$ 

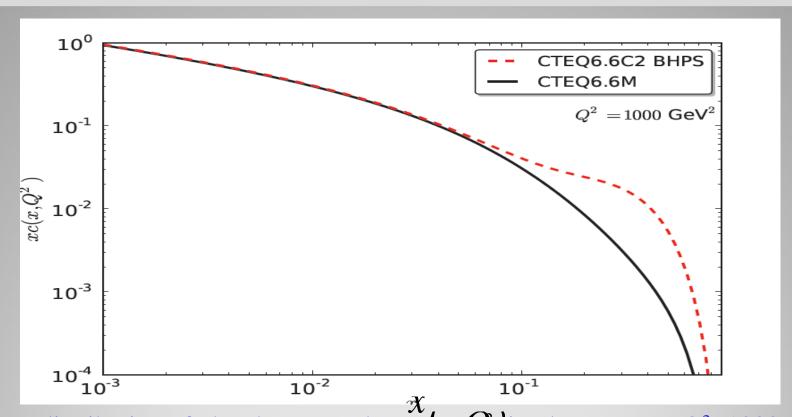




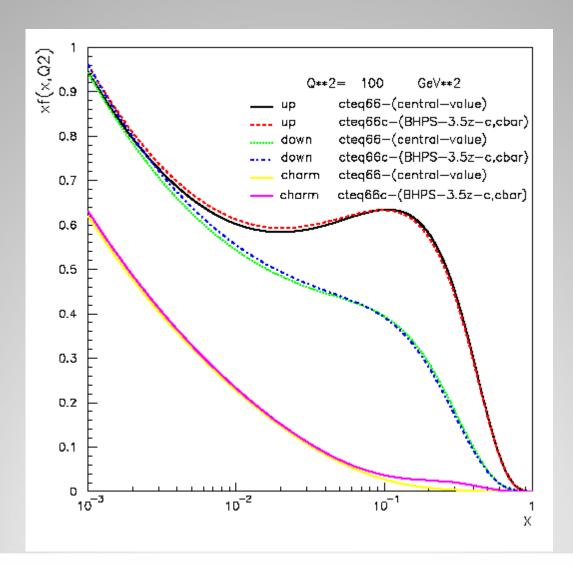
The x-distribution of the intrinsic Q calculated within the BHPS model. There is an enhancement at x > 0.1 Jen-Chieh Peng & We-Chen Chang, hep-ph/1207.2193.



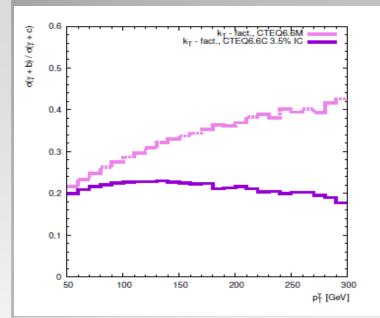
The x-distribution of the charm quarks  $xc(x,Q^2)$  in proton; the solid black line is the IC contribution with its probability about 3.5 %, the dash green curve is the see charm quark contribution  $xc_{sea}(x,Q^2)$  at  $Q^2=1.69$  GeV<sup>2</sup>. There is enhancement at x>0.1.

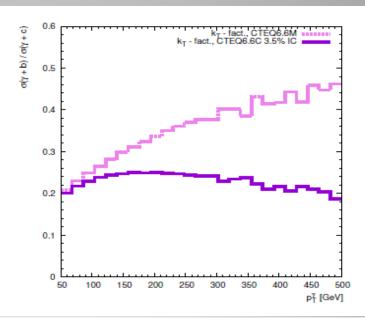


The x-distribution of the charm quarks  $\hat{x}(x,Q)$  in the proton at  $Q^2=1000$  GeV<sup>2</sup>; the solid black line is the radiatively generated charm density  $XC_{rg}(x,Q)$  distribution only, whereas the dashed curve is the sum of  $XC_{rg}(x,Q)$  and the intrinsic charm density  $XC_{rg}(x,Q)$  with its probability about 3.5 %. There is the sizable enhancement at x > 0.1.



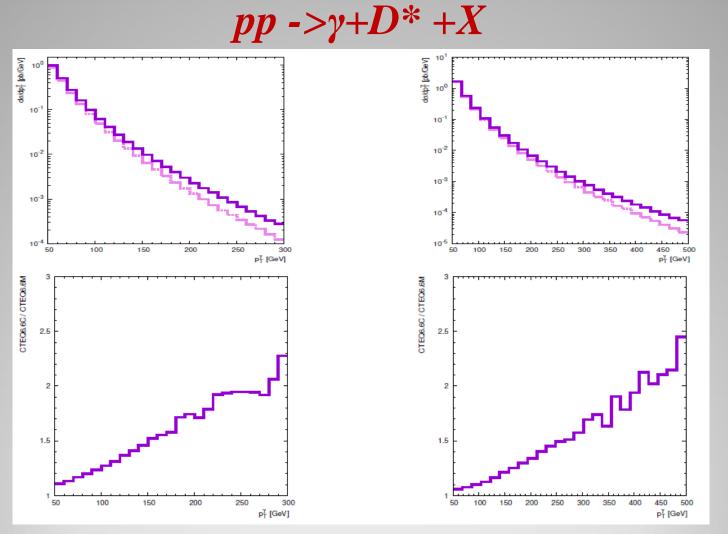
# $pp \rightarrow \gamma + c(b) + X$





The ratio  $\sigma(\gamma + b)/\sigma(\gamma + c)$  as a function of the photon transverse momentum at  $1.5 < |y^{\gamma}| < 2.4$ ,  $|y^{\text{jet}}| < 2.4$ ,  $p_T^{\text{jet}} > 20$  GeV,  $\sqrt{s} = 8$  TeV (left panel) and  $\sqrt{s} = 13$  TeV (right panel)

One can see a grow of this ratio by increasing of  $p_T^{\gamma}$ , when the intrinsic charm contribution (*IC*) in PDF is ignored. This ratio is approximately flat or slowly decreasing when the *IC* about 3.5 % is included. That is an additional *IC* signal, which could be measered at ATLAS.



Top:  $p_T$ —spectra of a photon at  $s^{1/2}$  =8 Tev (left) and 13 TeV(right) Bottom: ratio of  $p_T$ —spectra with IC (3.5%) and without it at 8 TeV (left) and 13 TeV (right).