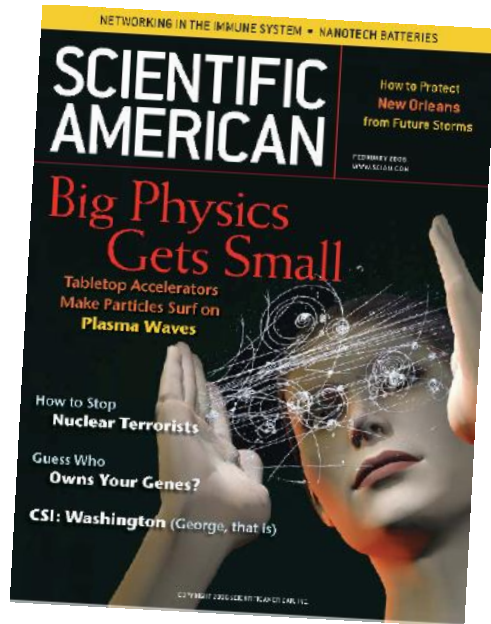


Laser Plasma Acceleration



R.W. Aßmann

Leading Scientist DESY

JUAS Seminar

Archamps, 3.2.2016



1. Accelerators – From Conventional Techniques to Plasmas
2. The Linear Regime
3. The Non-Linear Regime
4. Tolerances
5. Outlook for Europe



- 1. Accelerators – From Conventional Techniques to Plasmas**
2. The Linear Regime
3. The Non-Linear Regime
4. Tolerances
5. Outlook for Europe



First RF Linac Setup: PhD Wideröe 1927 in Aachen



Über ein neues Prinzip zur Herstellung hoher Spannungen

Von der Fakultät für Maschinenwirtschaft der Technischen Hochschule zu Aachen

zur Erlangung der Würde eines Doktor-Ingenieurs

genehmigte

Dissertation

vorgelegt von

Rolf Wideröe, Oslo

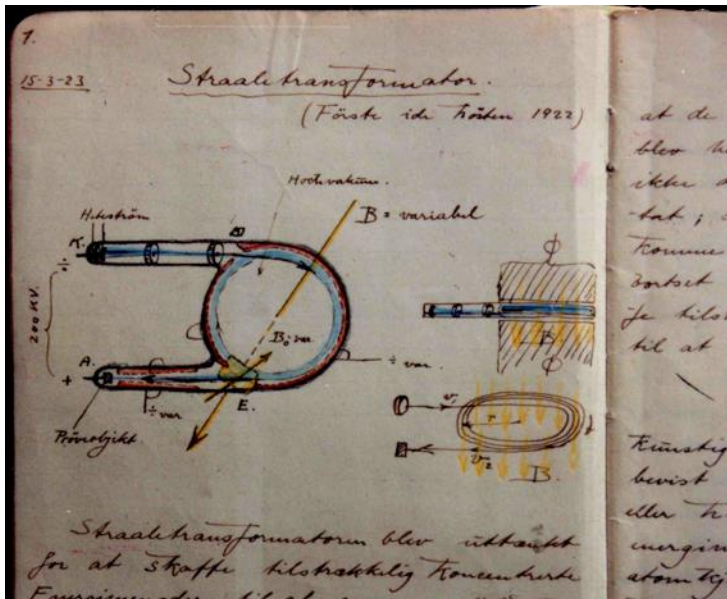
Referent: Professor Dr.-Ing. W. Rogowski

Korreferent: Professor Dr. L. Finzi

Tag der mündlichen Prüfung: 28. November 1927

27 pages

Sonderdruck aus Archiv für Elektrotechnik 1928, Bd. XXI, Heft 4
(Verlag von Julius Springer, Berlin W 9)



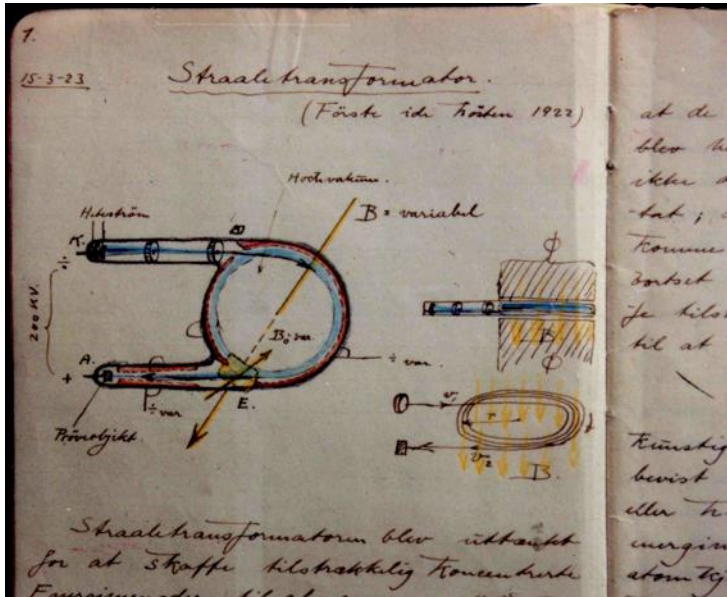
First RF Linac Setup: PhD Wideröe 1927 in Aachen



Über ein neues Prinzip zur Herstellung hoher Spannungen

First refused at University Karlsruhe as not feasible!

Wideröe had to go to Technical University Aachen



Tag der mündlichen Prüfung: 28. November 1927

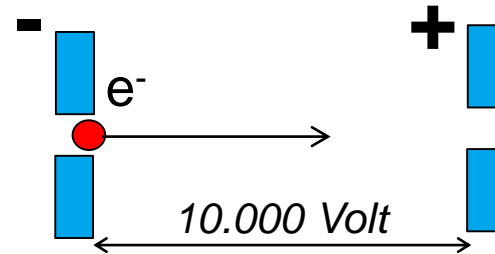
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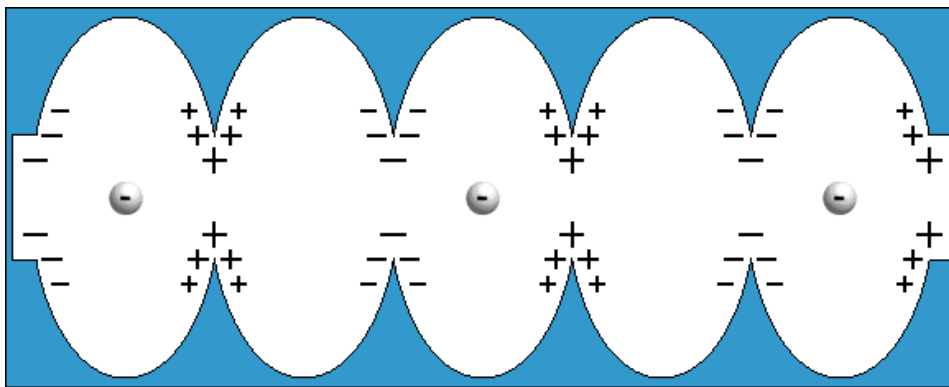


Electron Acceleration: The Success of RF

- > Areas with positive and negative charge; free electrons in between.
- > Free electron (e^-) is accelerated towards the positive charge (anode) („**Gegensätze ziehen sich an**“).



- > For a voltage of 10.000 Volt the electron gains 10.000 electron-Volt („eV“).
- > Higher energies with **alternating voltage („RF“)**:



Sketch Padamse, Tigner

“Runzelröhre”

20.000.000 Volt per Meter



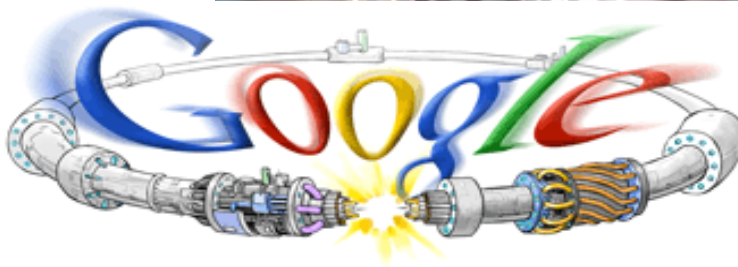
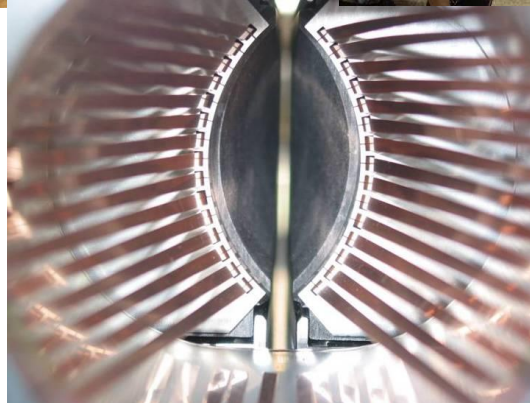
80 Years (and many inventions) later: LHC as a Masterpiece of Accelerator Science



First beam
10.9. 2008



Higgs
Sem.
4.7.
2012

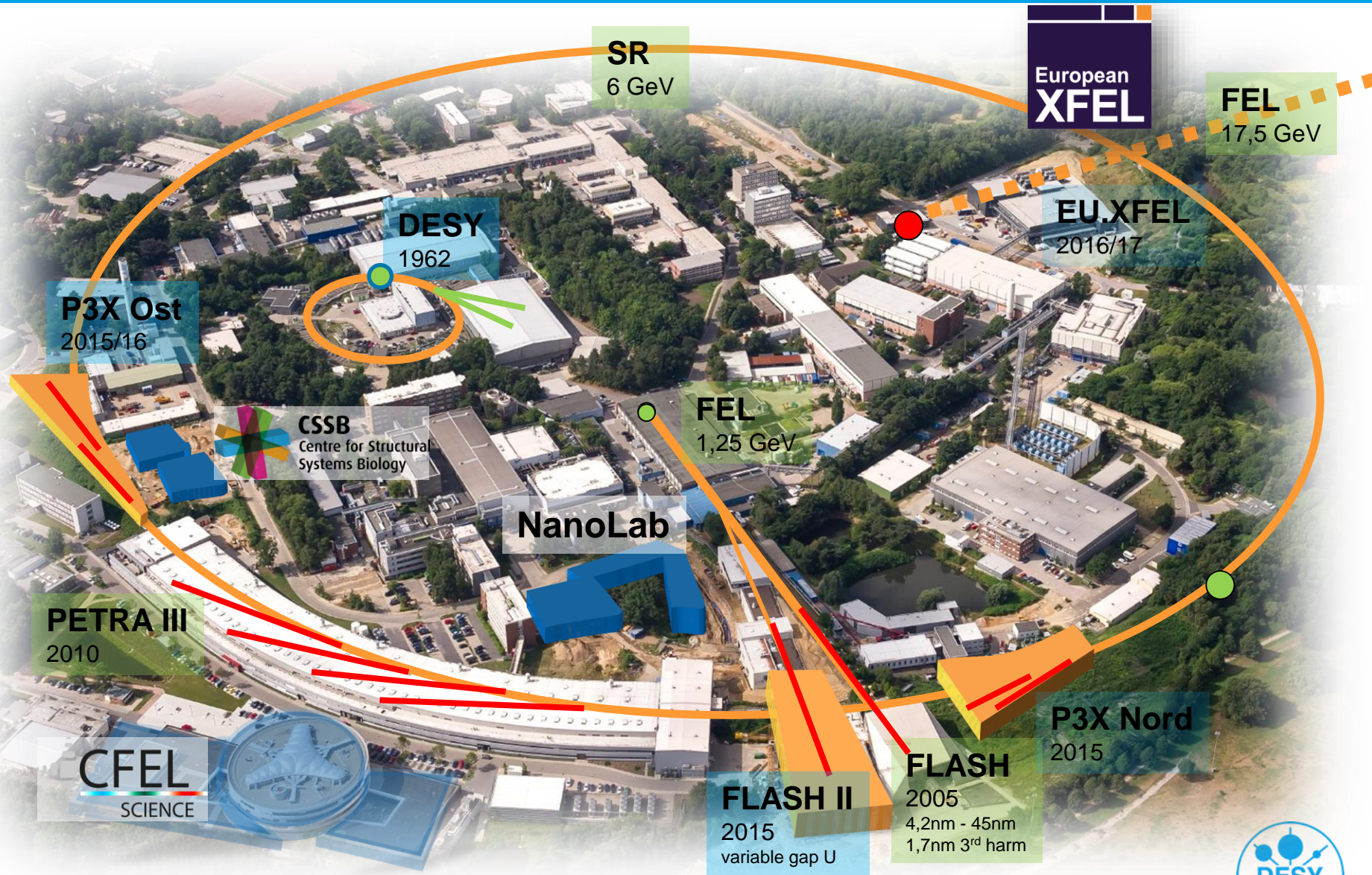


DESY 50 Years ago...



9284

Today: X-Ray Facilities at DESY. *Masterpieces for photon science*



Acceleration: Conventional and Advanced

Surfer gain velocity and energy by riding the water wave!

Charged particles gain energy by riding the electromagnetic wave!



Modern lasers generate light pulses with very large transverse fields:

Many 1.000 billion volt per meter

Plasma or metallic structures couple fields to our particles!

ANGUS Laser Lab (200 TW, DESY & University Hamburg)



The Laser Promise: Transverse Electrical Field

$$E_0 = \sqrt{2 \cdot \frac{I_0}{c \epsilon_0}}$$

ϵ_0 = Dielectric constant

c = Light velocity

$$P = 100 \text{ TW}$$

$$r_0 = 10 \mu\text{m}$$

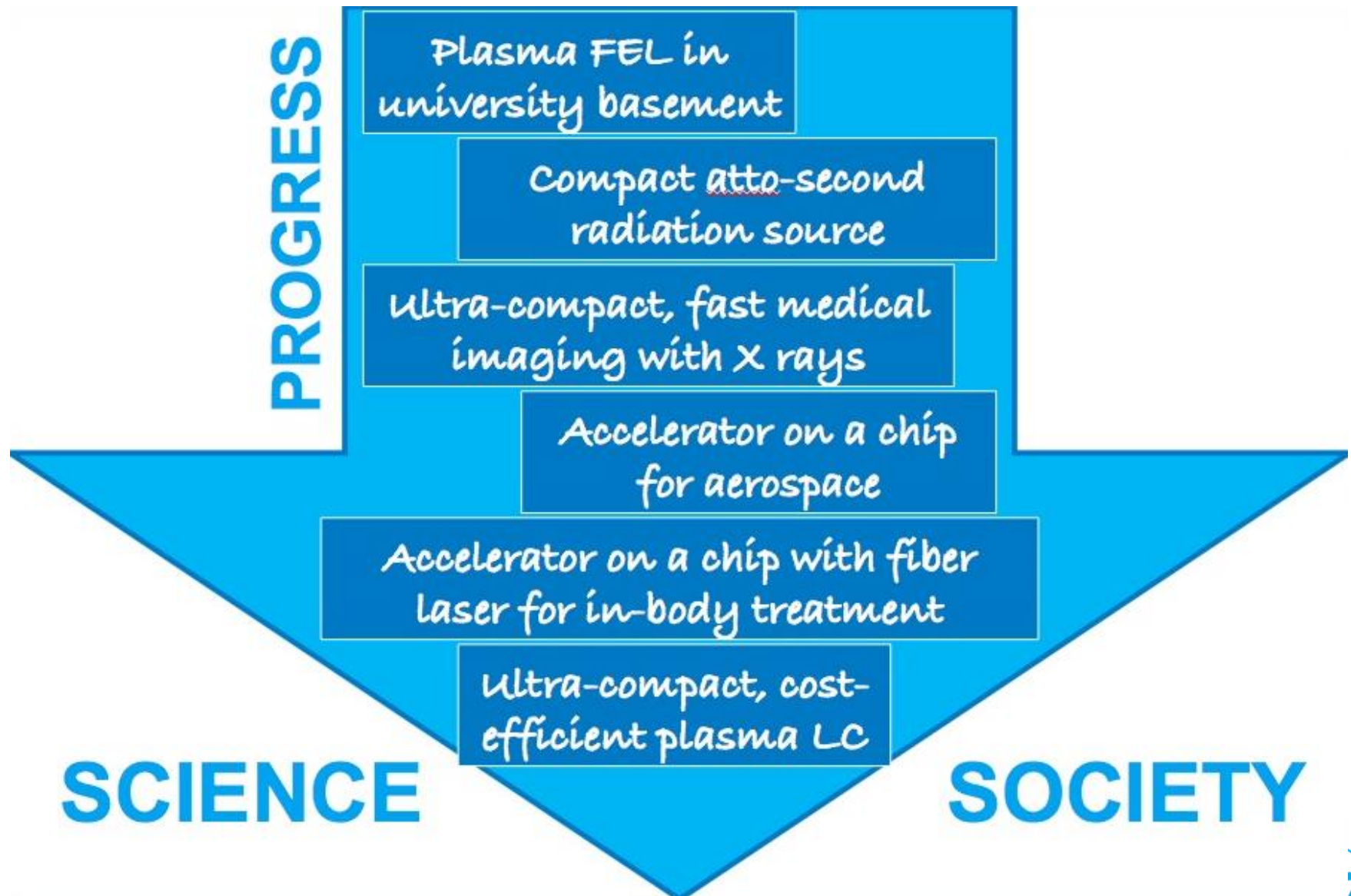
$$I_0 = 6.4 \cdot 10^{19} \text{ W/cm}^2$$


$$E_0 = 22 \text{ TV/m}$$

This is
what we
need!

Scientists wonder: Can we use
the strong transverse electrical
fields to accelerate our beam?

High fields trigger imagination of scientists and public...



High Gradient Accelerators

High
Gradients
(1 – 100 GV/m)



High
Frequencies
(> 100 GHz)



Small
Dimensions
(< 1 mm)

- > No **klystrons** for high frequencies!
- > Use particle bunches or laser pulses as drivers.
- > Material limitations solved through “new cavities”: dielectric materials, plasma cavities, ...
- > Two main directions:

1 Microstructure Accelerator

Laser- or beam driven
Vacuum accelerators
Conventional field design

2 Plasma Accelerator

Laser- or beam driven
Dynamic Plasma Structure
Plasma field calculations



Laser-Driven Micro Structures (Vacuum) – 1

- > 1 GeV/m possible but low absolute energies achieved so far
- > **AXSIS project (ERC synergy grant)** at DESY/ Uni Hamburg: THz laser-driven accelerator with atto-second science → *Kärtner/Fromme/Chapman/Assmann*

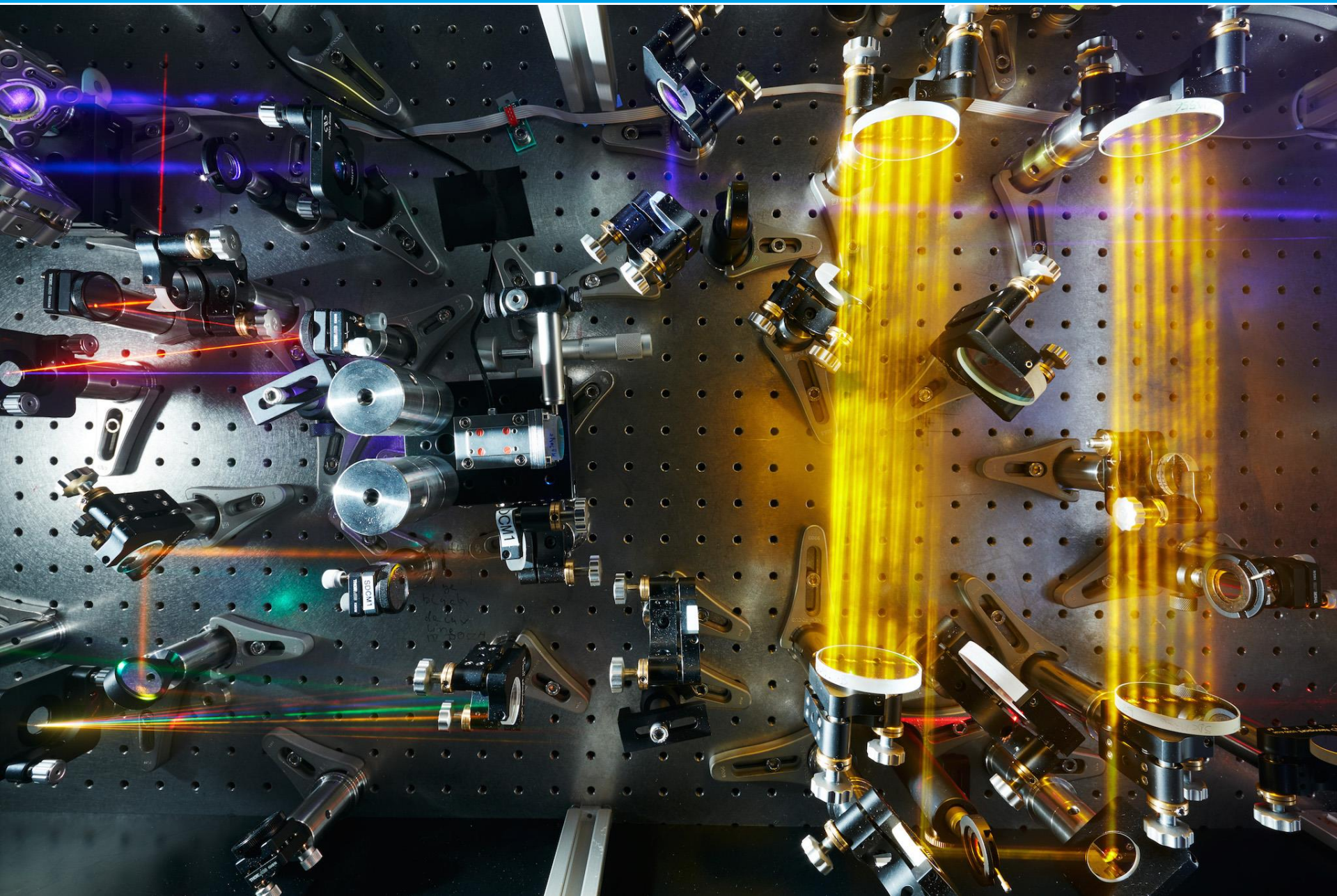


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from anywhere in the world

European Research Council
Established by the European Commission

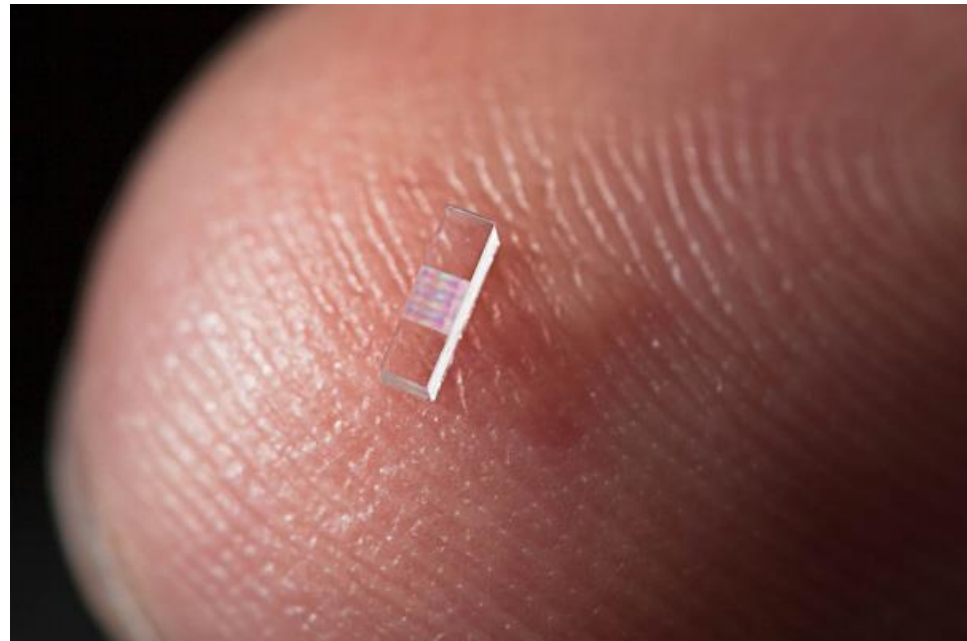


THz Laser Lab (DESY, CFEL, University Hamburg)



Laser-Driven Micro Structures (Vacuum) – 2

- > **“Accelerator on a Chip”** grant from Moore foundation for work by/at Stanford, SLAC, University Erlangen, DESY, University Hamburg, PSI, EPFL, University Darmstadt, CST, UCLA
- > Lasers drive **structures that are engraved on microchips** (e.g. Silicium)
- > Major breakthroughs can be envisaged:
 - **Mass production**
 - **Implantable accelerators** for in-body irradiation of tumors
 - Accelerators for **outer space**



Financed by Silicon Valley billionaire...



ABOUT

Intel co-founder Gordon and his wife Betty established the foundation to create positive change for future generations.



Courtesy of Hawley Peterson Snyder

Our Science Program invests in the development of new technologies, supports the world's top research scientists and brings together new—often groundbreaking—scientific partnerships. Our passion for discovery reflects that of our founders, Gordon and Betty Moore.

We believe in the inherent value of science and the sense of awe that discovery inspires. Scientific advancement and societal benefits will occur if we find ways to unleash the potential of inquiry and exploration. So we take risks, we incubate change, and we foster the kind of excitement that inspires third-graders to become scientists. We look for opportunities to transform, or even create, entire fields.

Our Founders

We're inspired by the innovation, compassion and focus of our founders, Gordon and Betty Moore. Gordon's thinking was part of the birth of Silicon Valley in the late 1950s. Betty's commitment to improving the lives of patients resulted in a regional collaborative that is making a difference in the care that Californians receive. Together, they've identified places where they, and the foundation, could create positive change for future generations.

[Read More](#)



Courtesy of Susanna Frohman, San Jose Mercury News



The Laser Promise: Transverse Electrical Field

$$E_0 = \sqrt{2 \cdot \frac{I_0}{c \epsilon_0}}$$

ϵ_0 = Dielectric constant

c = Light velocity

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$$E_0 = 22 \text{ TV/m}$$

This is
what we
need!

Scientists wonder: Can we use
the strong transverse electrical
fields to accelerate our beam?

Lorentz Force F

$$\mathbf{F} = q (\mathbf{E} + \mathbf{v} \times \mathbf{B})$$

q = Charge

\mathbf{v} = Velocity

Longitudinal
electrical field to
accelerate a particle

Transverse
magnetic field to
guide a particle

Laser Plasma Accelerator: Transverse to Longitudinal

- > Idea: Use a plasma to convert the transverse space charge force of a beam driver (or the electrical field of the laser) into a longitudinal electrical field in the plasma!

VOLUME 43, NUMBER 4

PHYSICAL REVIEW LETTERS

23 JULY 1979

Laser Electron Accelerator

T. Tajima and J. M. Dawson

Department of Physics, University of California, Los Angeles, California 90024

(Received 9 March 1979)

An intense electromagnetic pulse can create a wave of plasma oscillations through the action of the nonlinear ponderomotive force. Electrons trapped in the wake can be accelerated to high energy. Existing glass lasers of power density 10^{18} W/cm² shone on plasmas of densities 10^{18} cm⁻³ can yield giga-electronvolts of electron energy per centimeter of acceleration distance. This acceleration mechanism is demonstrated through computer simulation. Applications to accelerators and pulsers are examined.

Collective plasma accelerators have recently received considerable theoretical and experimental investigation. Earlier Fermi¹ and McMillan² considered cosmic-ray particle acceleration by moving magnetic fields¹ or electromagnetic waves.² In terms of the realizable laboratory technology for collective accelerators,

the wavelength of the plasma waves in the wake:

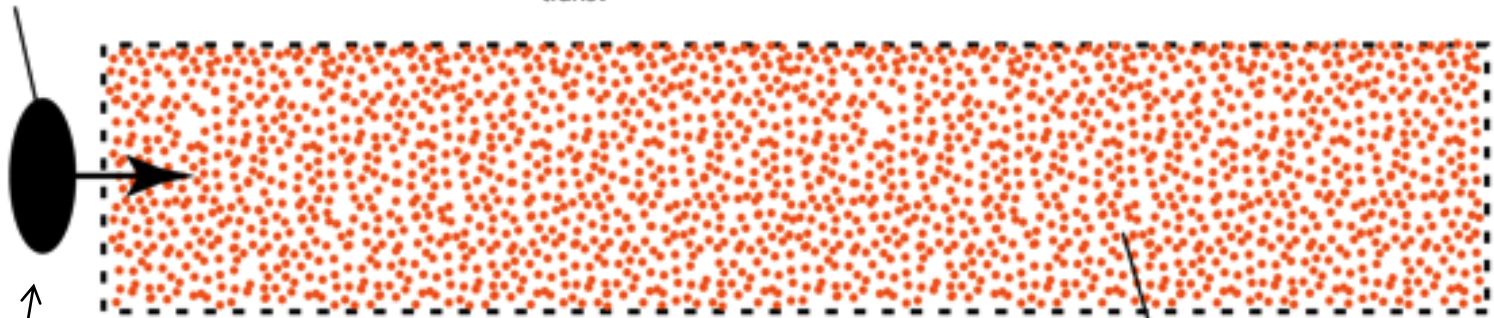
$$L_1 = \lambda_w/2 = \pi c/\omega_p. \quad (2)$$

An alternative way of exciting the plasmon is to inject two laser beams with slightly different frequencies (with frequency difference $\Delta\omega \sim \omega_p$) so that the beat distance of the packet becomes



Laser Plasma-Acceleration (Internal Injection)

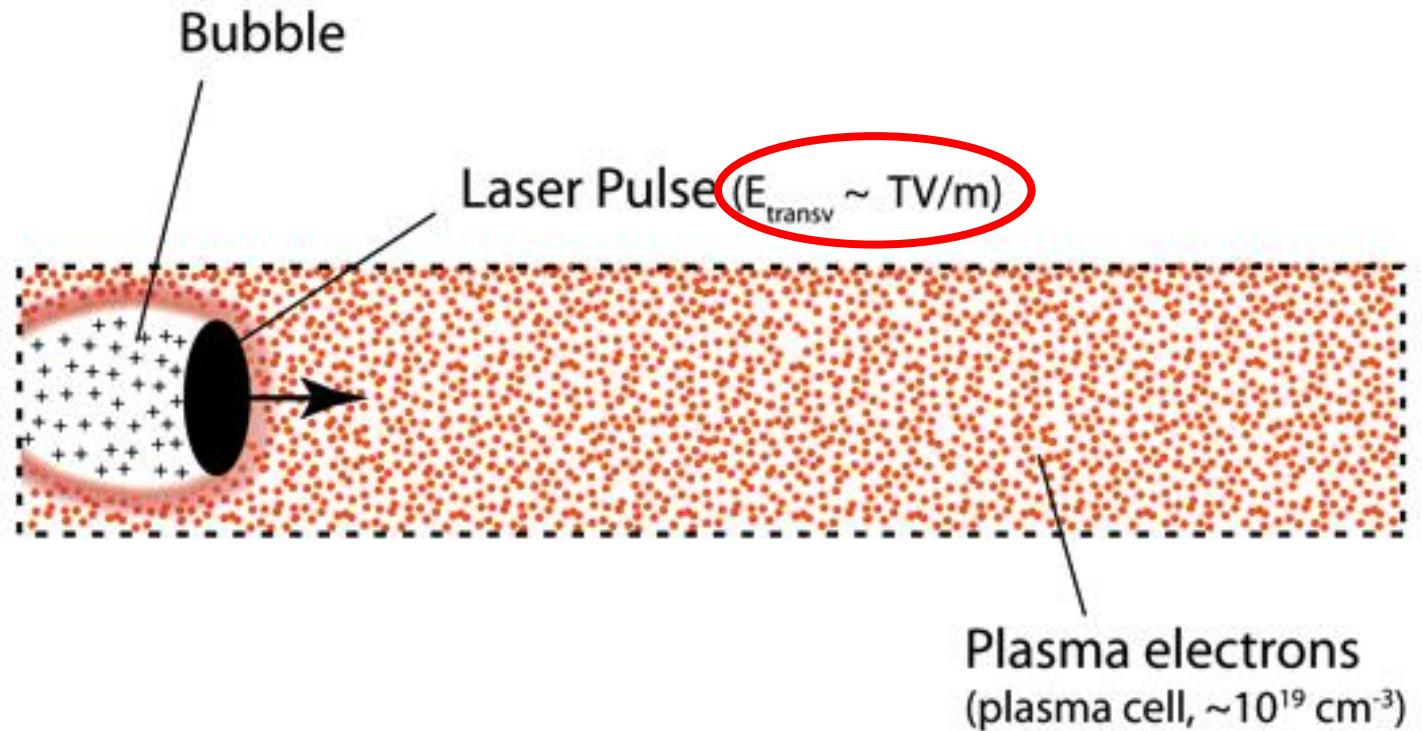
Laser Pulse (200 TW, ~ 30 fs, $E_{\text{transv}} \sim \text{TV/m}$)



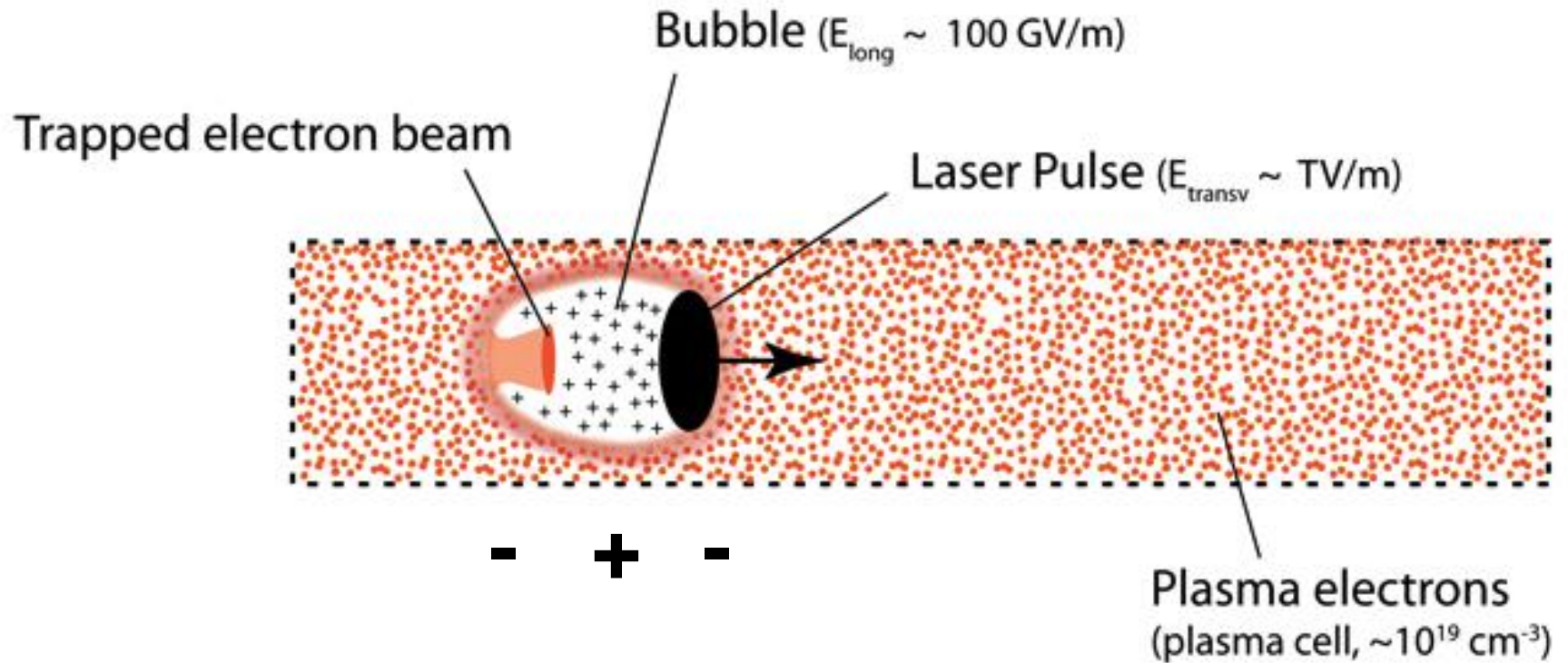
Plasma electrons
(plasma cell, $\sim 10^{19} \text{ cm}^{-3}$)

Works the same way with an **electron beam as wakefield driver**. But then usually lower plasma density. Ponderomotive force of laser is then replaced with space charge force of electrons on plasma electrons (repelling).

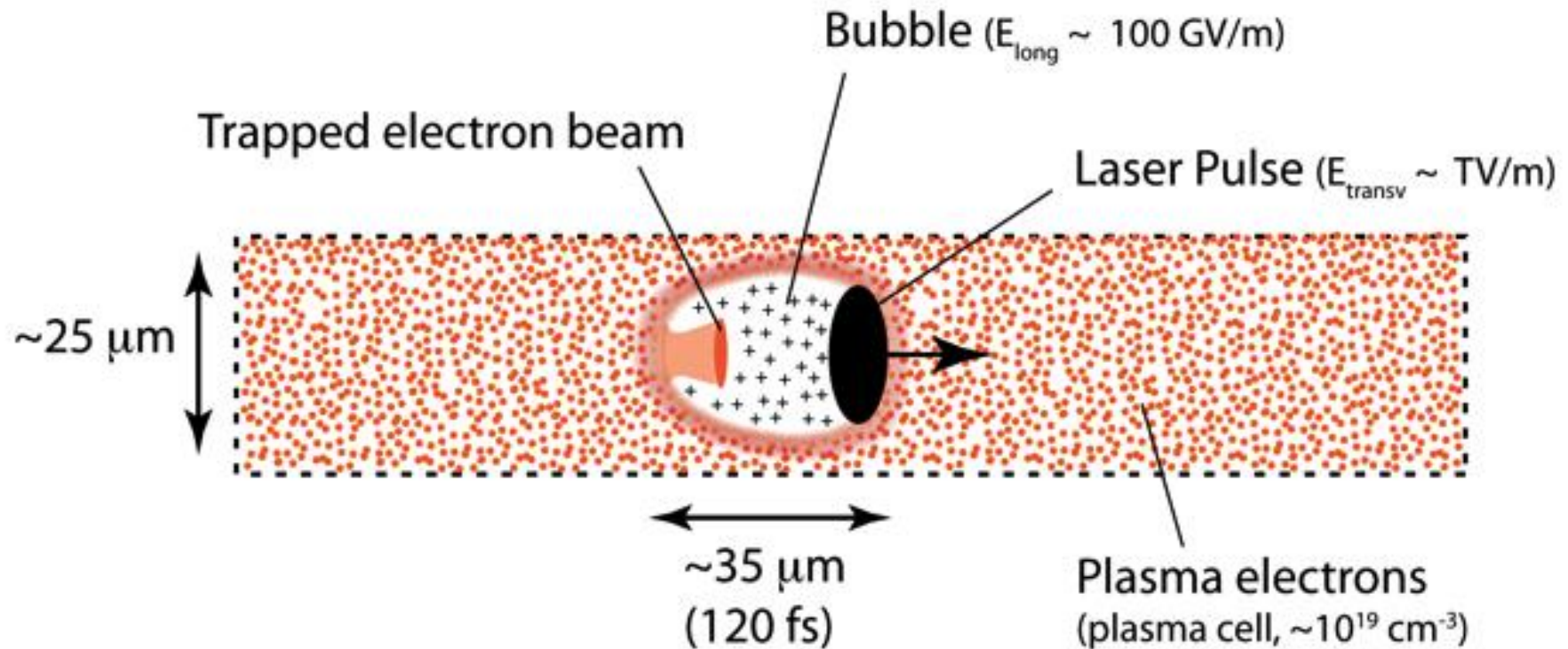
Laser Plasma-Acceleration (Internal Injection)



Laser Plasma-Acceleration (Internal Injection)

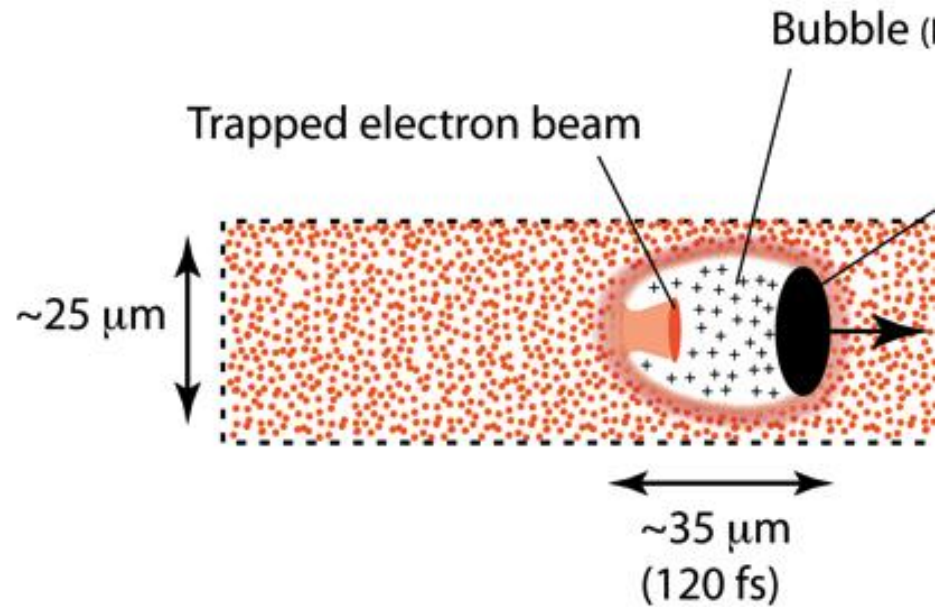


Laser Plasma-Acceleration (Internal Injection)



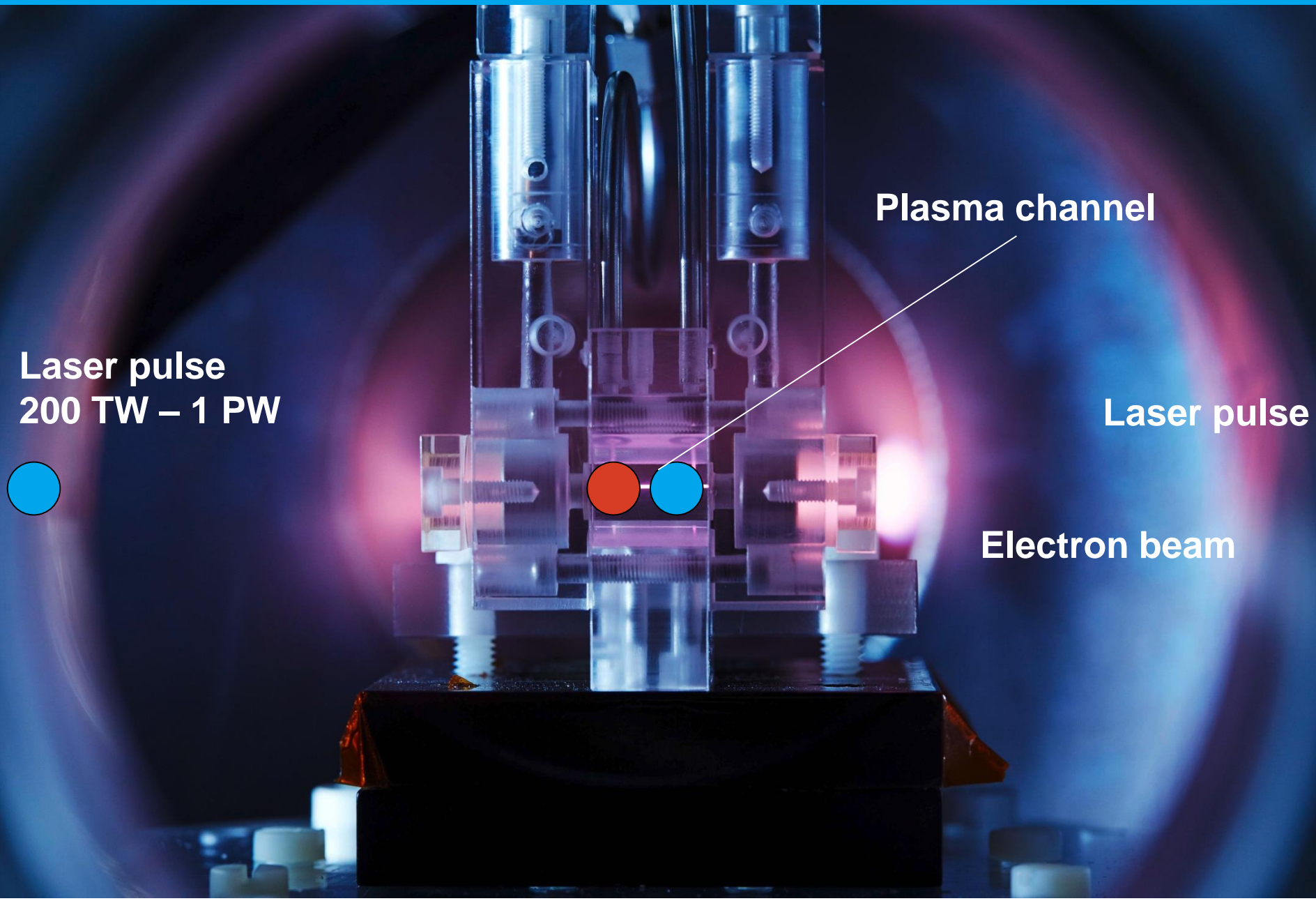
This accelerator fits into a human hair!

Laser Plasma-Acceleration (Internal Injection)



- This proved highly successful with electron bunches of **up to 4.25 GeV produced over a few cm.**
- Small dimensions involved → few **micron tolerances!**
- Highly compact but also **highly complex** accelerator: generation, bunching, focusing, acceleration, (wiggling) all in one small volume.
- Energy spread and stability at the few % level.

Laser Plasma Accelerators for Electron Beams



And the Plasma Accelerator is Compact...

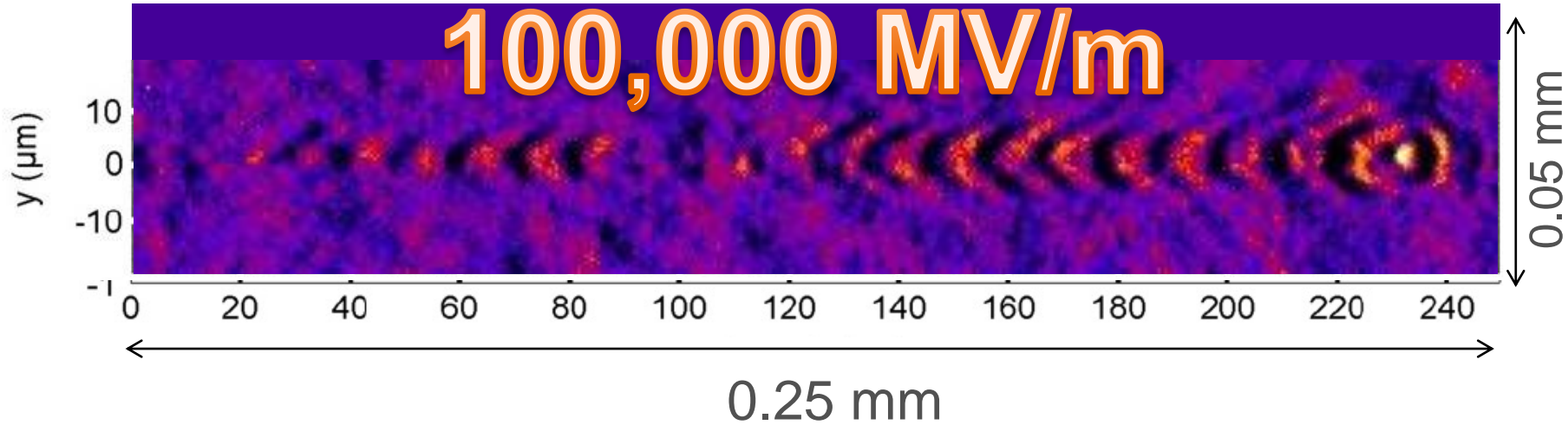


Foto Laser-Plasmabeschleuniger

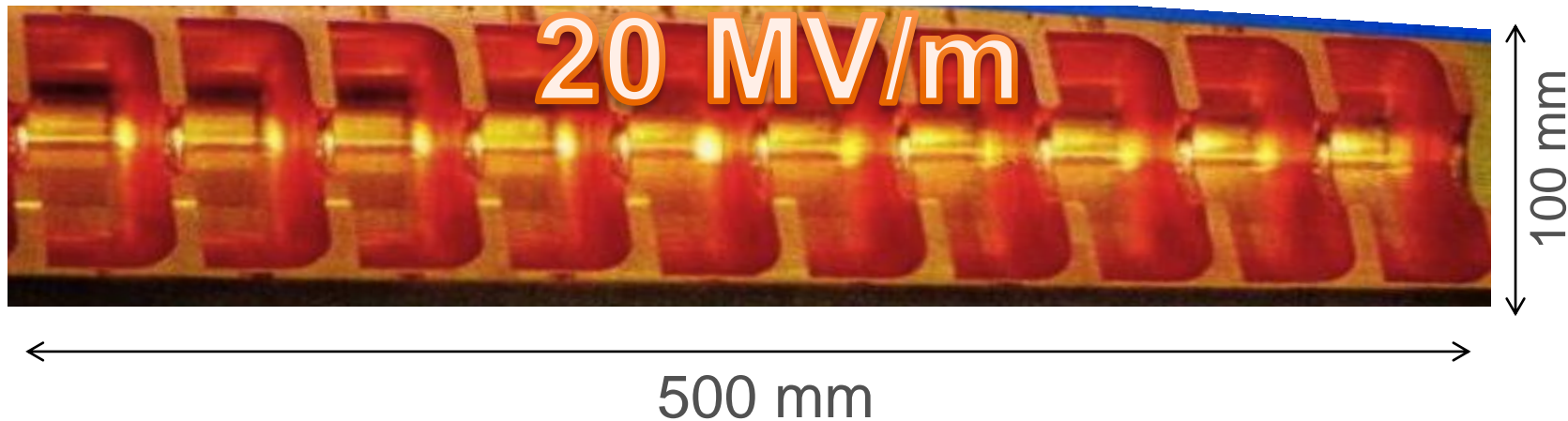
Few-cycle optical probe-pulse for investigation of relativistic laser-plasma interactions

M. B. Schwab,^{1,a)} A. Sävert,¹ O. Jäckel,^{1,2} J. Polz,¹ M. Schnell,¹ T. Rinck,¹ L. Veisz,³
 M. Möller,¹ P. Hansinger,¹ G. G. Paulus,^{1,2} and M. C. Kaluza^{1,2}
¹Institut für Optik und Quantenelektronik, Max-Wien-Platz 1, 07743 Jena, Germany
²Helmholtz-Institut Jena, Fröbelstieg 3, 07743 Jena, Germany
³Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Straße 1, 85748 Garching, Germany

2013



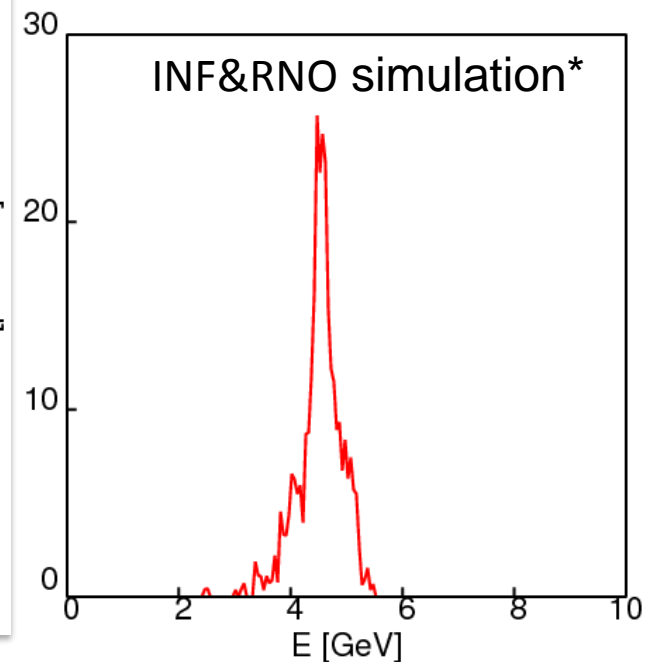
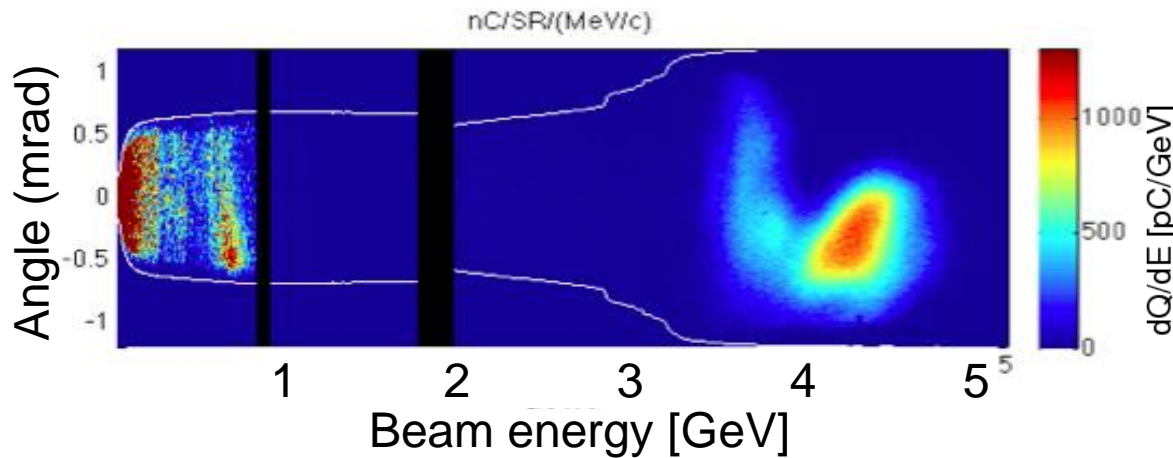
Metall
(Kupfer)
S band
Linac
Struktur
Mikro-
Wellen
zur
Wellener-
zeugung



4.25 GeV beams have been obtained from 9 cm plasma channel powered by 310 TW laser pulses (15 J)

*C. Benedetti et al., proceedings of AAC2010, proceedings of ICAP2012

Electron beam spectrum



Slide: W. Leemans

- **Laser (E=15 J):**
 - Measured longitudinal profile ($T_0 = 40$ fs)
 - Measured far field mode ($w_0 = 53$ μm)
- **Plasma:** parabolic plasma channel (**length 9 cm**, $n_0 \sim 6 \times 10^{17} \text{ cm}^{-3}$)

W.P. Leemans et al., PRL 2014, in print

	Exp.	Sim.
Energy	4.25 GeV	4.5 GeV
$\Delta E/E$	5%	3.2%
Charge	~ 20 pC	23 pC
Divergence	0.3 mrad	0.6 mrad

Wait one moment... Compact and Cost-Effective?

- > Consider laser-driven plasma: Presently one can buy **1 Peta-Watt Ti:Sa** lasers from industry for a low double digit million € cost.
- > The most compact 1 PW laser is installed in HZDR, Dresden, Germany (part of ARD):

Required space: 120 m²
(can be visited)

- > The laser size drives the size of such an accelerator facility. With such a 1 PW laser electrons of **4.25 GeV have been produced within 9 cm** (see LBNL result).
- > The 1 PW laser should be sufficient for a **10 GeV accelerator within about 20 cm**. Total footprint: about 200-300 m² (incl. all infrastructure).
- > Now do this conventionally and compare size and cost!
(e.g. **10 GeV = 500 m of conventional acceleration with 20 MV/m**)
- > Need to bring up quality, efficiency and repetition rate.



1. Accelerators – From Conventional Techniques to Plasmas
2. **The Linear Regime**
3. The Non-Linear Regime
4. Tolerances
5. Outlook for Europe



The Linear Regime

- Analytical treatment
- Placement of beams in the plasma accelerating structure
- Maximum acceleration (transformer ratio)
- Optimizations: Energy spread, phase slippage, stability, reproducibility



Linear Wakefields (R. Ruth / P. Chen 1986)

$$\mathcal{E}_z \simeq -A \left(1 - \frac{r^2}{a^2}\right) \cos(k_p z - \omega_p t)$$

$$r \ll a$$

$$\mathcal{E}_r \simeq 2A \frac{r}{k_p a^2} \sin(k_p z - \omega_p t)$$

\mathcal{E} = electrical field

z = long. coord.

r = radial coord.

a = driver radius

ω_p = plasma frequency

k_p = plasma wave number

t = time variable

e = electron charge

N = number e- drive bunch

$$A = \begin{cases} \frac{\omega_p \tau k_p e E_0^2}{8 \omega^2 m} & PBWA \\ \frac{8eN}{a^2} & PWFA \end{cases}$$

ω = laser frequency

τ = laser pulse length

E_0 = laser electrical field

m = mass of electron

Can be analytically solved and treated. Here comparison beam-driven and laser-driven (beat wave).



Linear Wakefields (R. Ruth / P. Chen 1986)

Accelerating field

Depends on radial position r

Changes between accelerating and decelerating as function of longitudinal position z

$$\mathcal{E}_z \simeq -A \left(1 - \frac{r^2}{a^2}\right) \cos(k_p z - \omega_p t)$$

$\pi/2$ out of phase

$$\mathcal{E}_r \simeq 2A \frac{r}{k_p a^2} \sin(k_p z - \omega_p t)$$

Transverse field

Depends on radial position r

Changes between focusing and defocusing as function of longitudinal position z



The Useful Regime of Plasma Accelerators

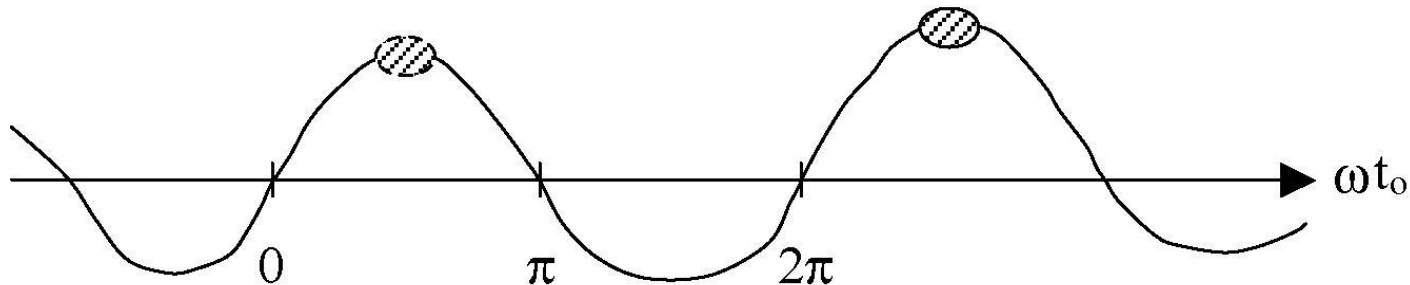
Two conditions for an accelerator:

1. **Accelerated bunch must be in accelerating regime.**
2. **Accelerated bunch must be in focusing regime.**

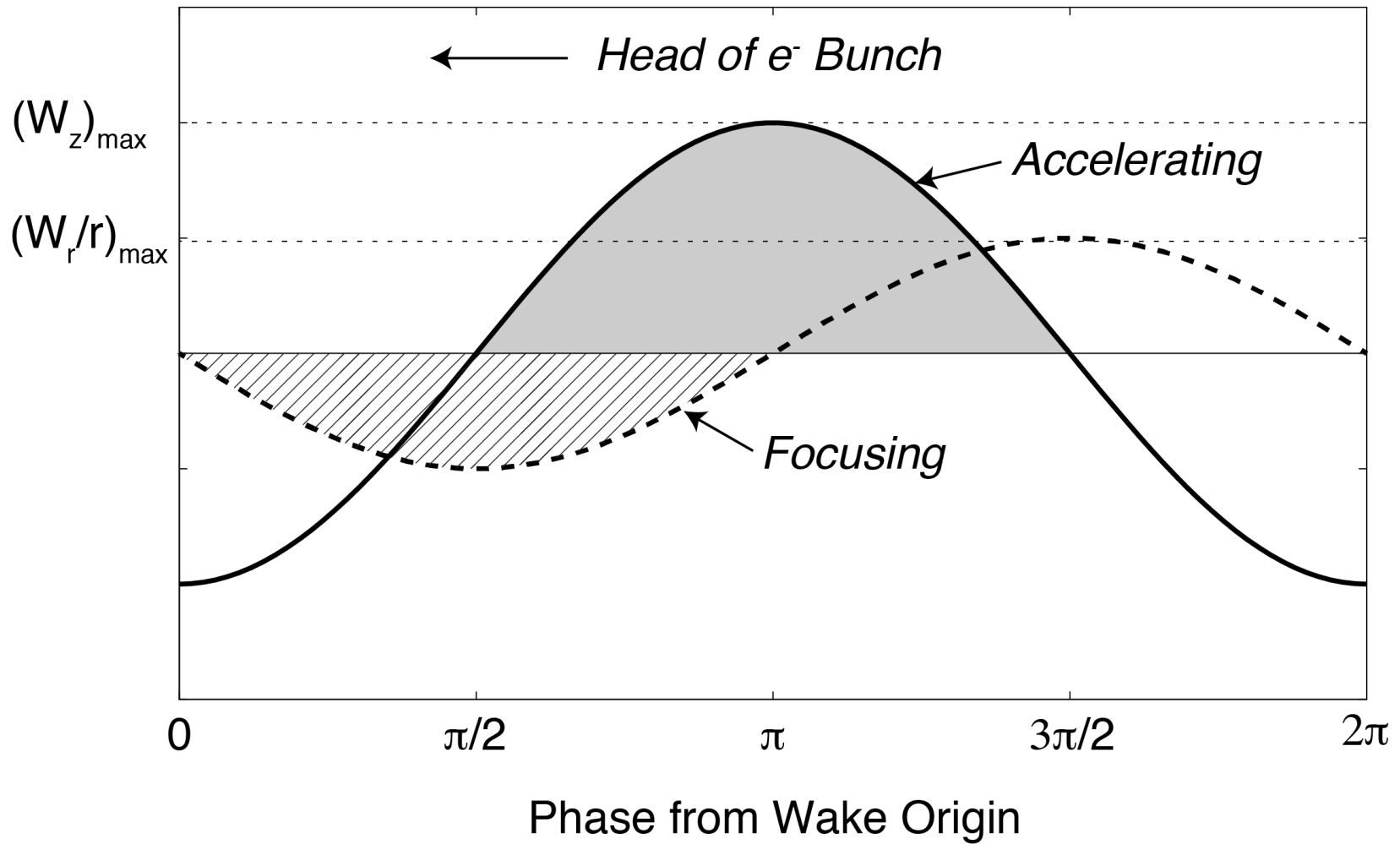
These two conditions define a useful range of acceleration!

Reminder metallic RF accelerator structures:

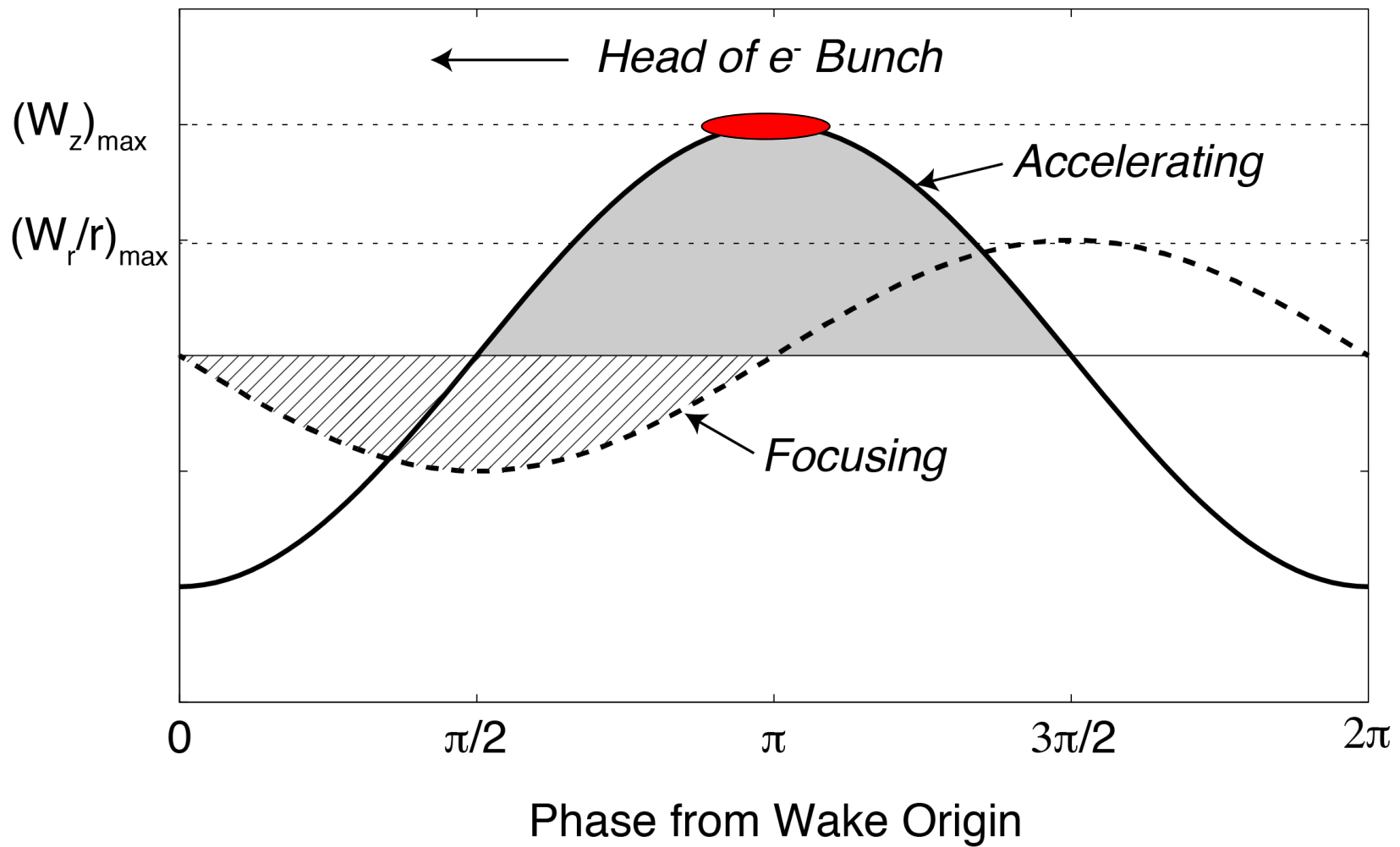
no net transverse fields for beam particles \rightarrow full accelerating range is available for beam \rightarrow usually place the beam on the crest of the accelerating voltage



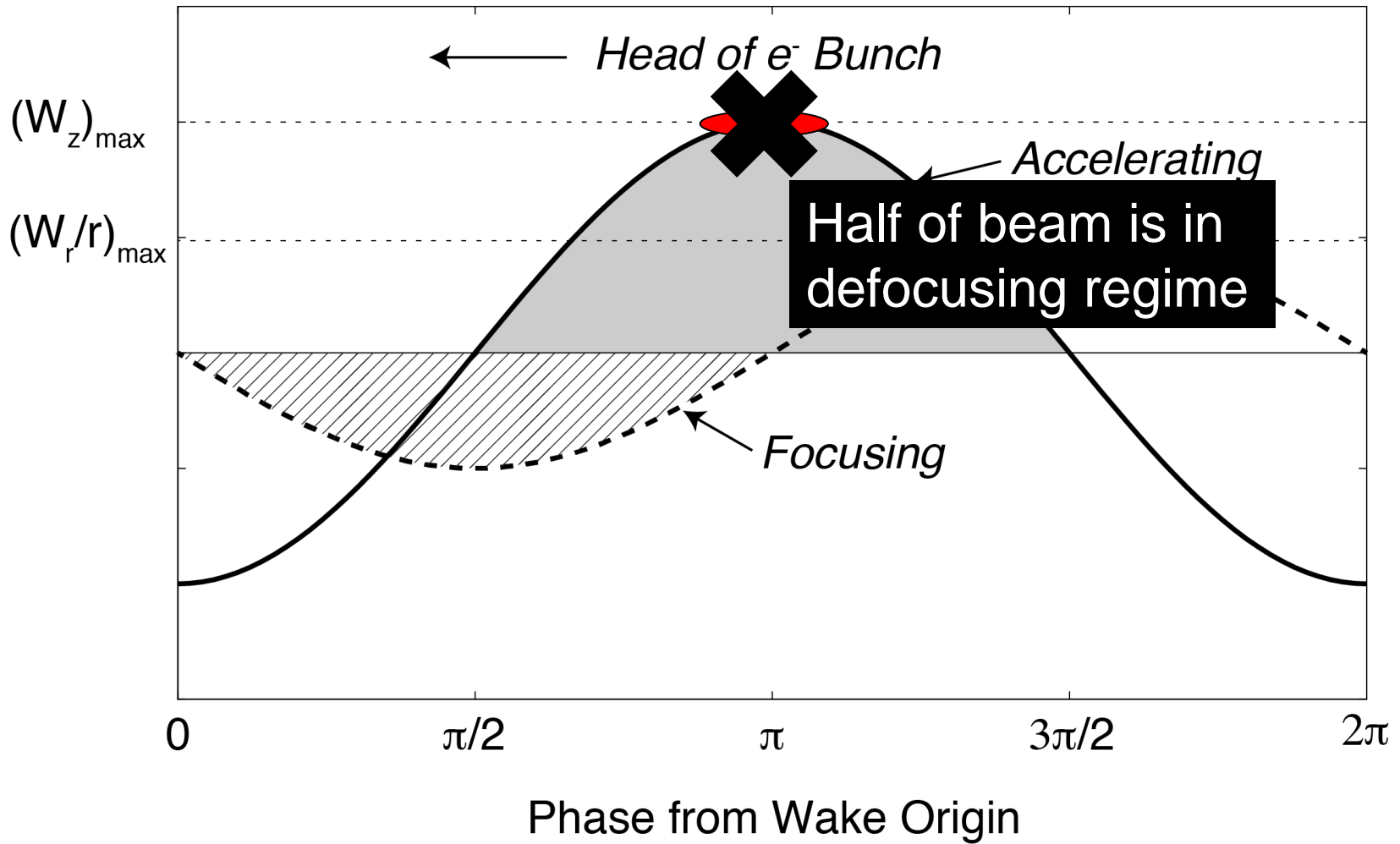
Plasma Accelerator



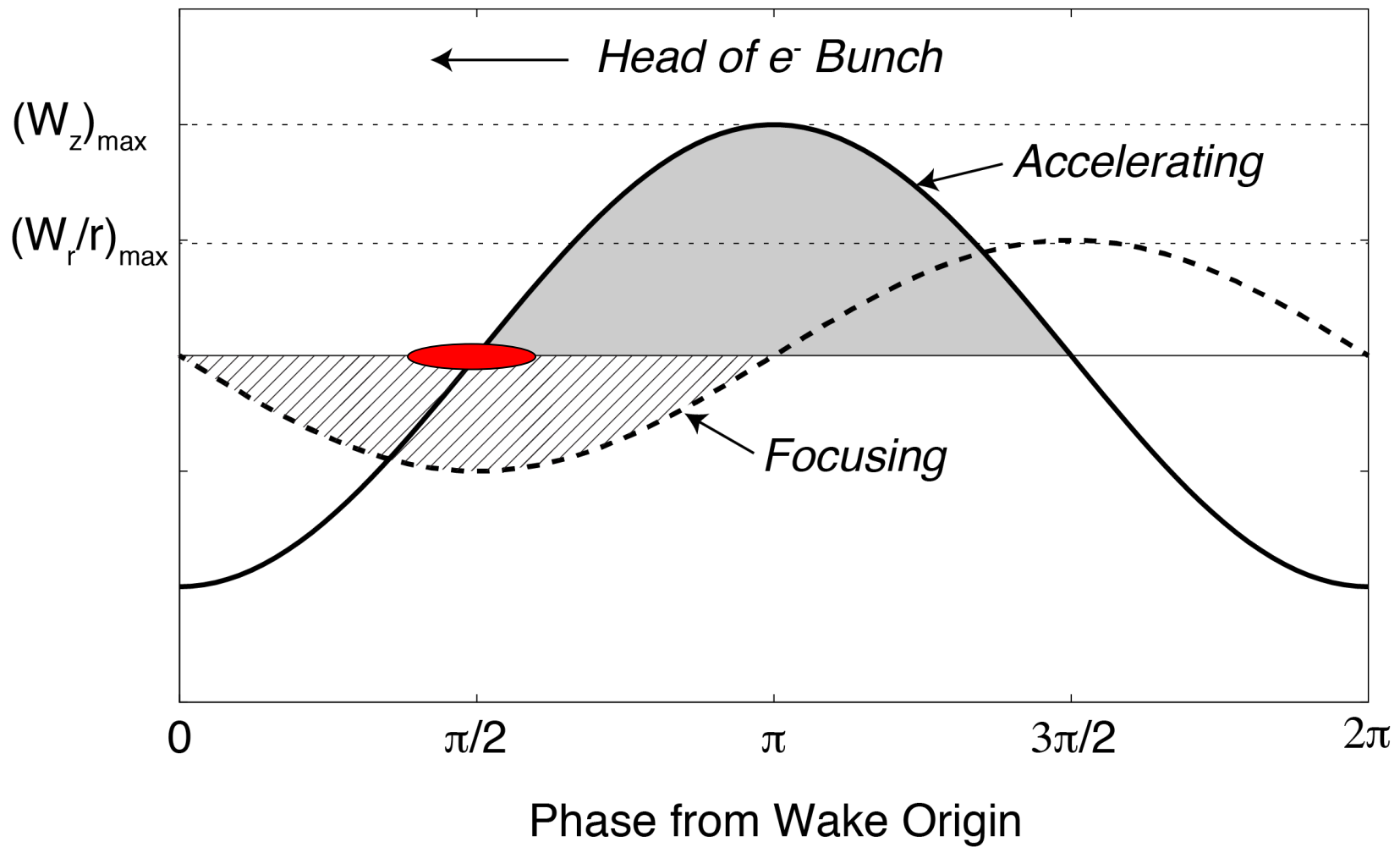
Plasma Accelerator



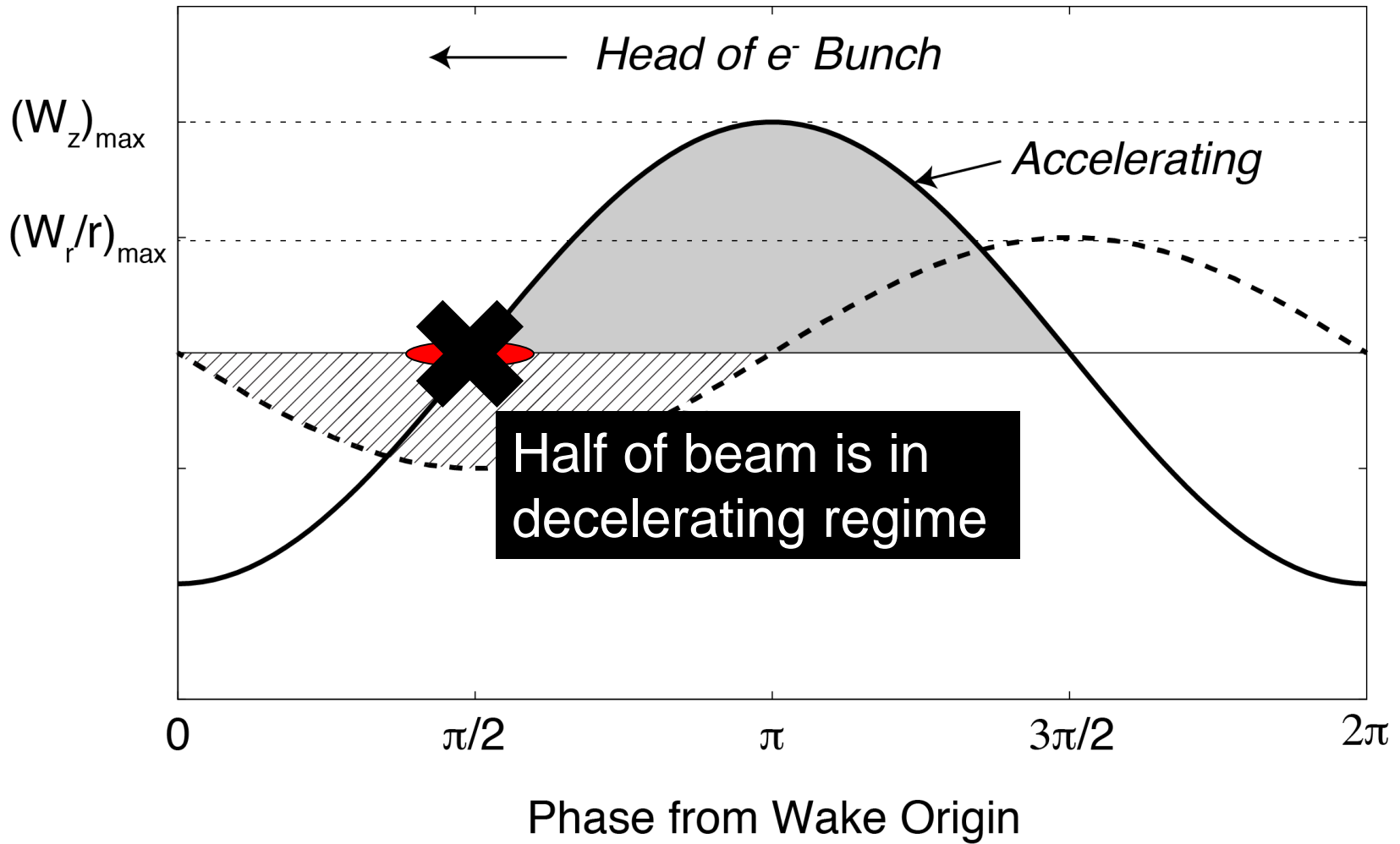
Plasma Accelerator



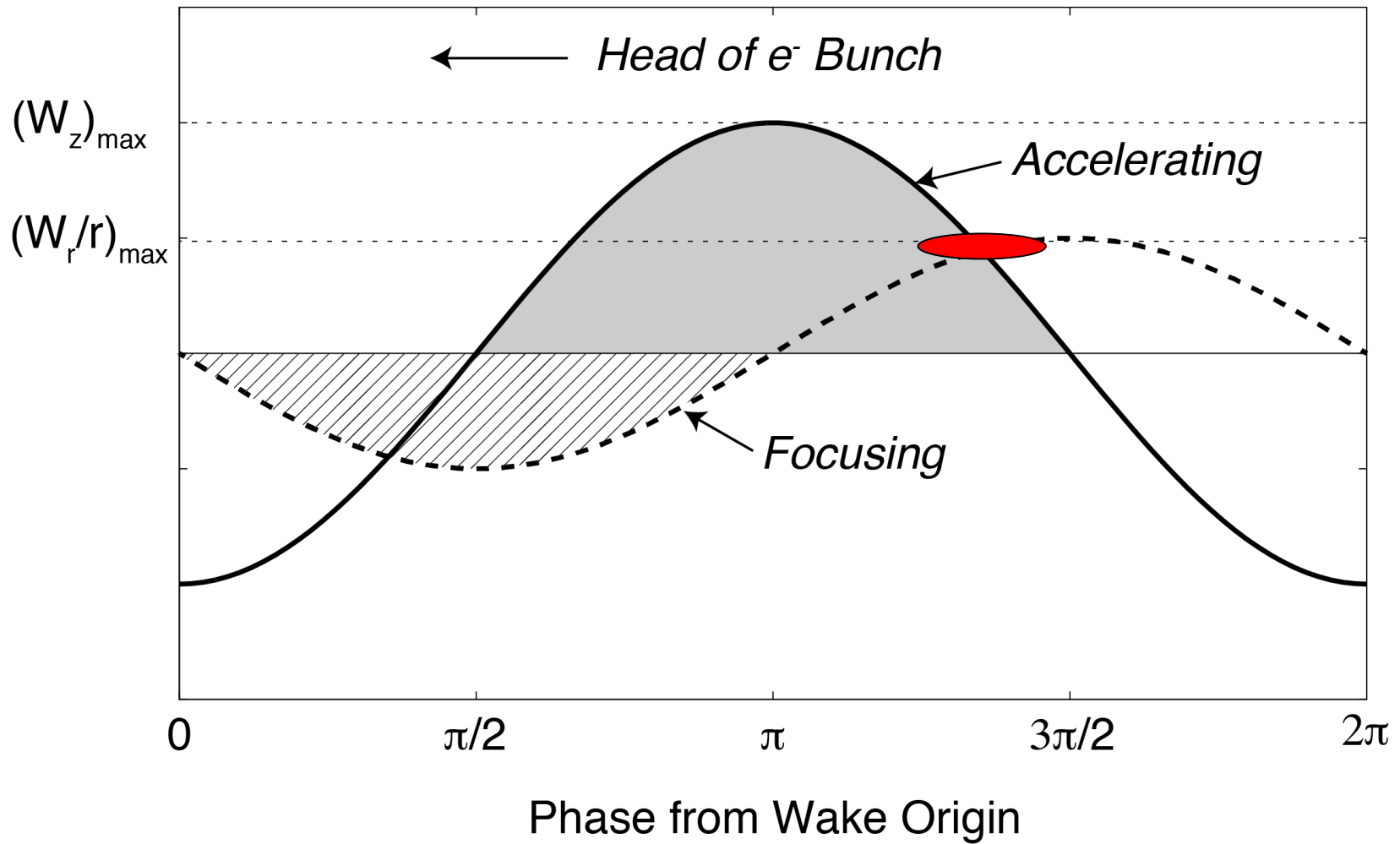
Plasma Accelerator



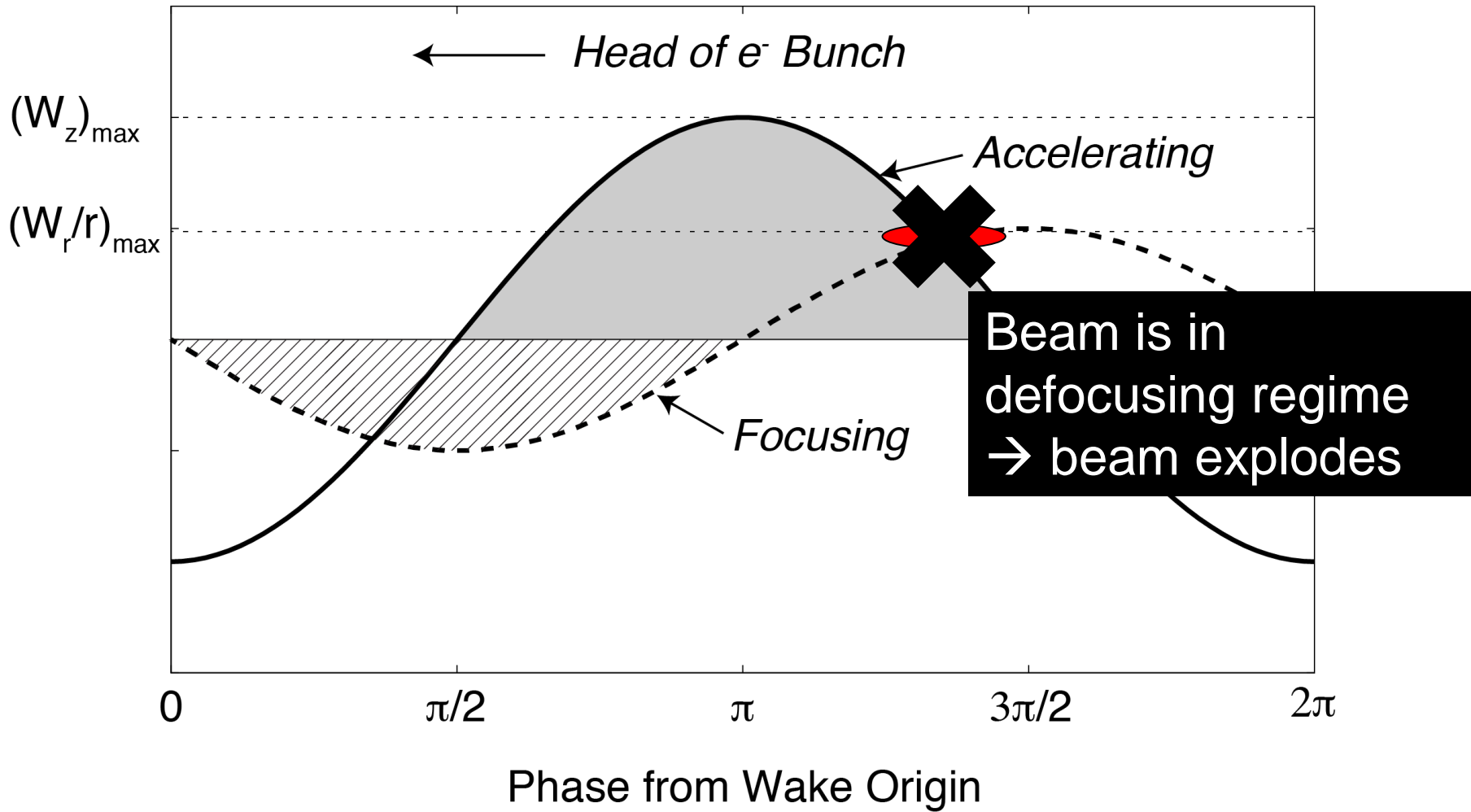
Plasma Accelerator



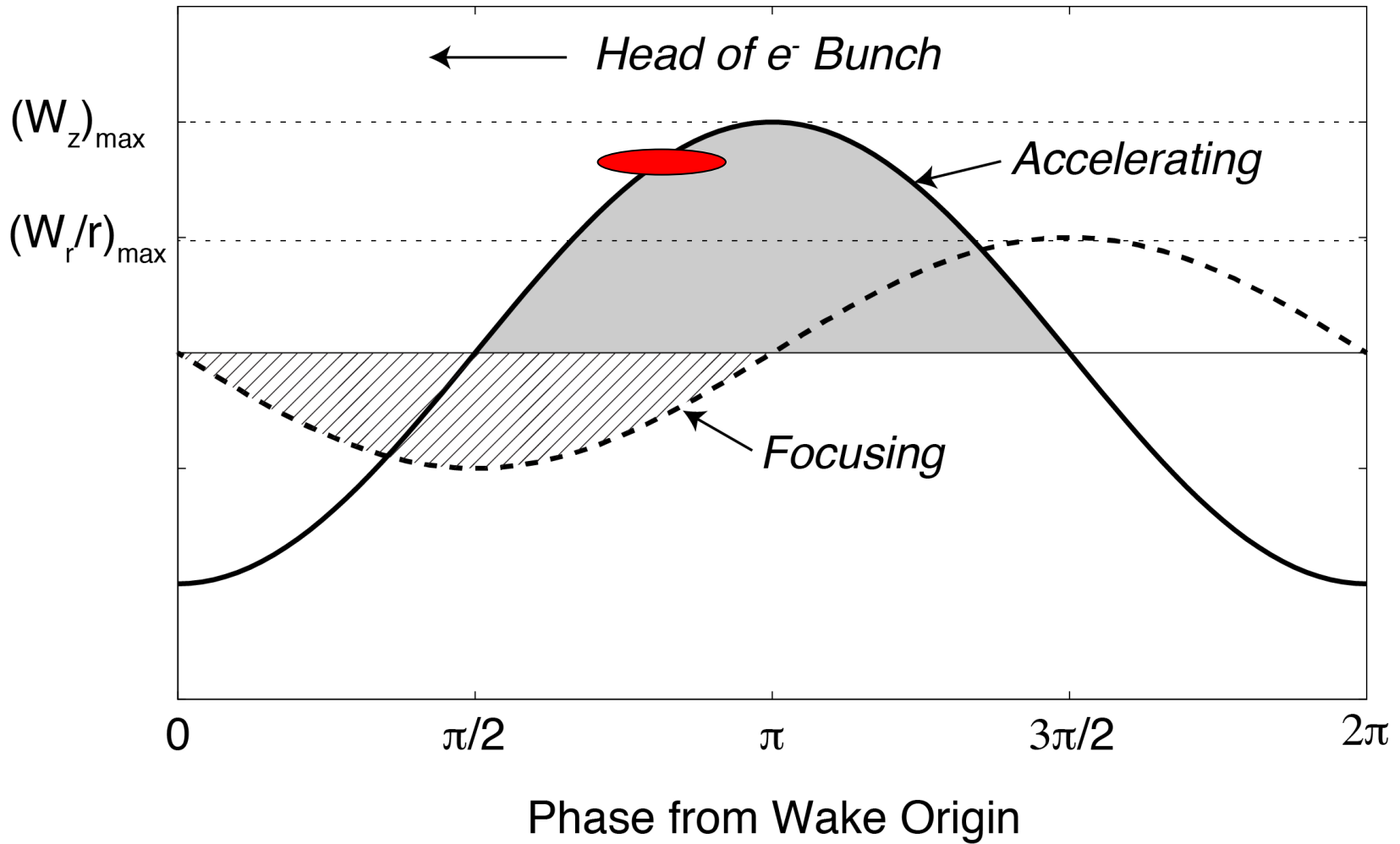
Plasma Accelerator



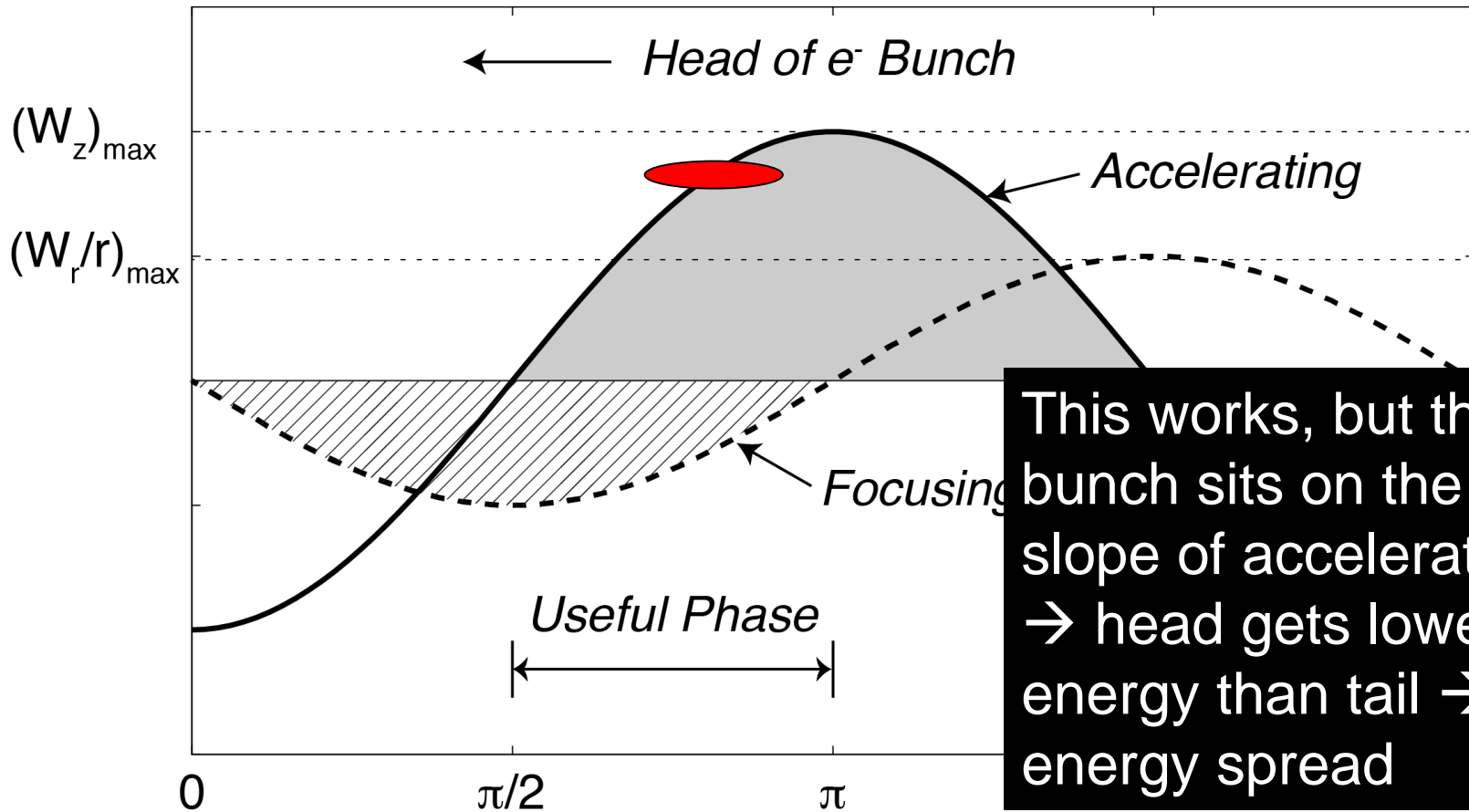
Plasma Accelerator



Plasma Accelerator

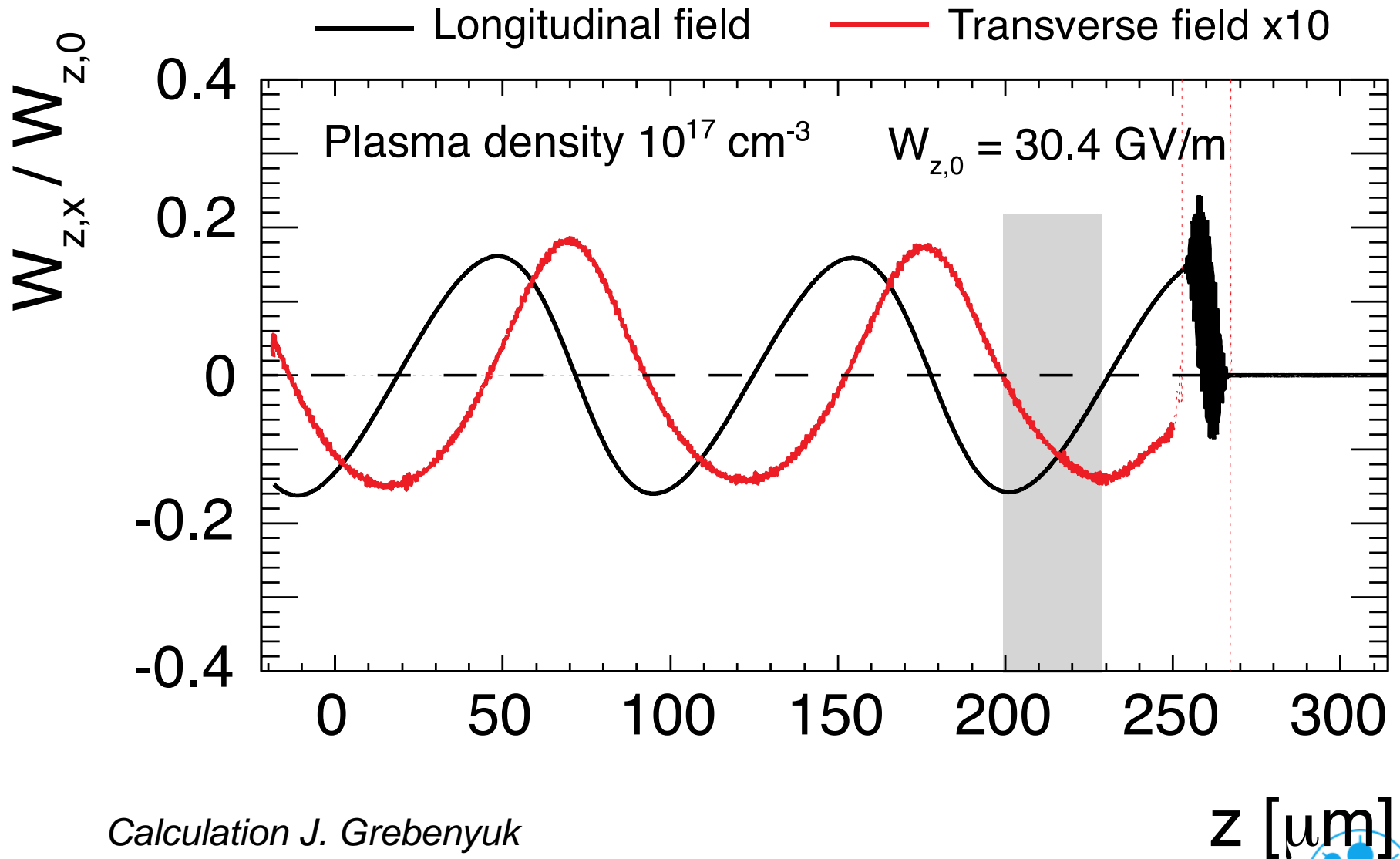


Plasma Accelerator



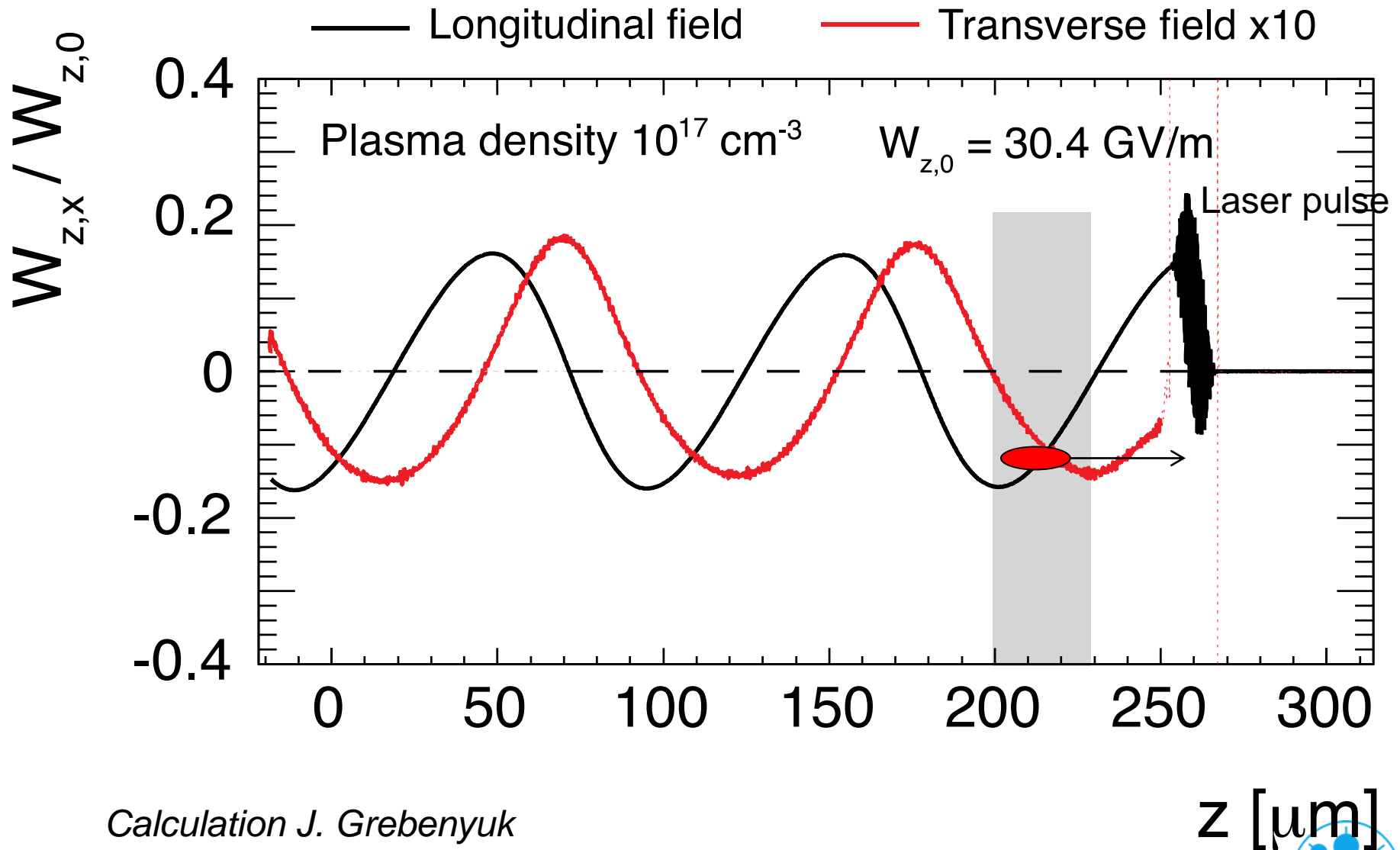
Phase from Wake Origin

Comparison with OSIRIS simulation

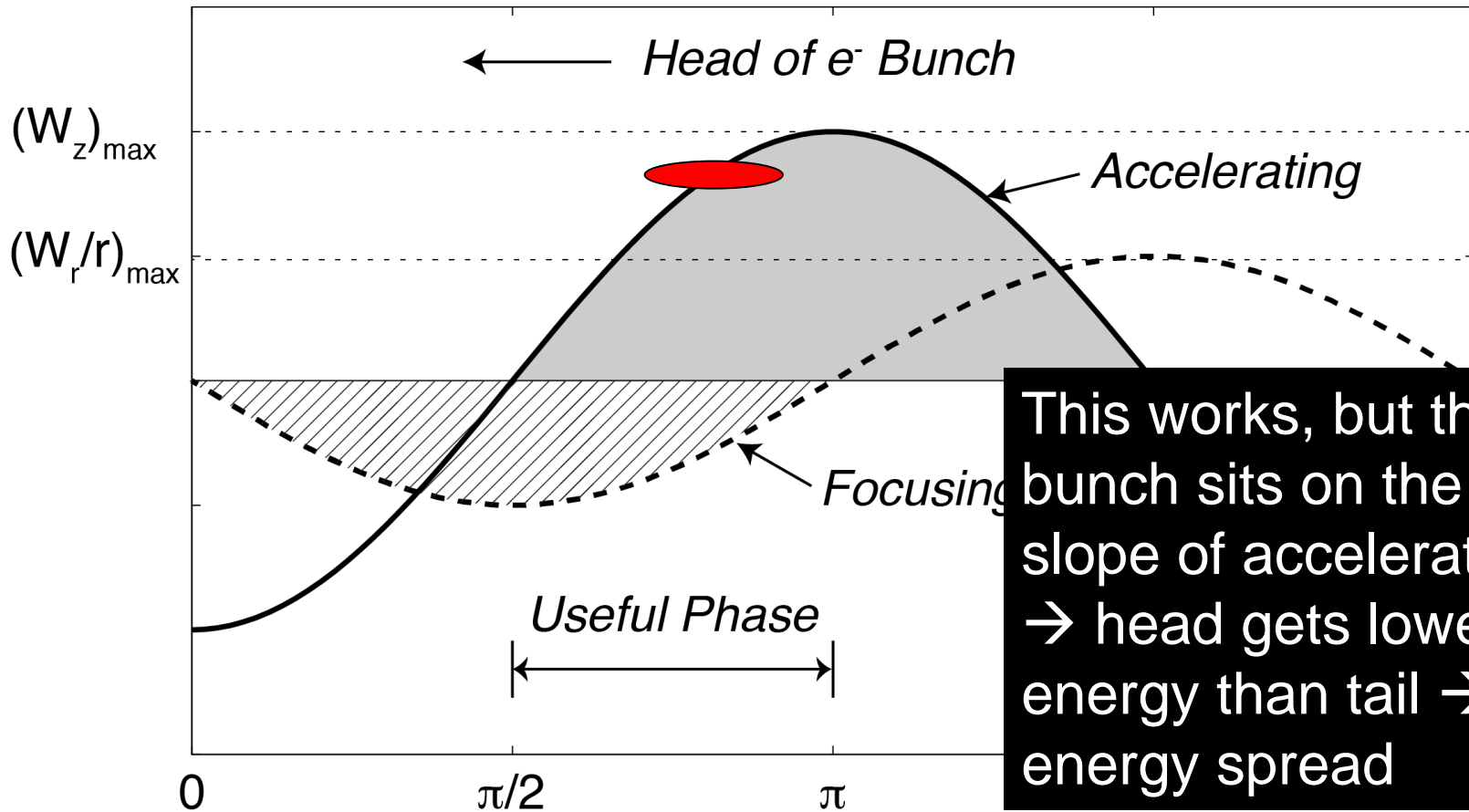


Calculation J. Grebenyuk

Comparison with OSIRIS simulation



Plasma Accelerator



This works, but the bunch sits on the slope of acceleration → head gets lower energy than tail → energy spread

Phase from Wake Origin



Optimization 1: Energy Spread

Reduce energy spread (head to tail \rightarrow correlated with z)

Minimize: Ratio of accelerated bunch length over $\frac{1}{4}$ plasma wavelength!

Minimize length accelerated bunch

Ultra-short bunches (fs, as)

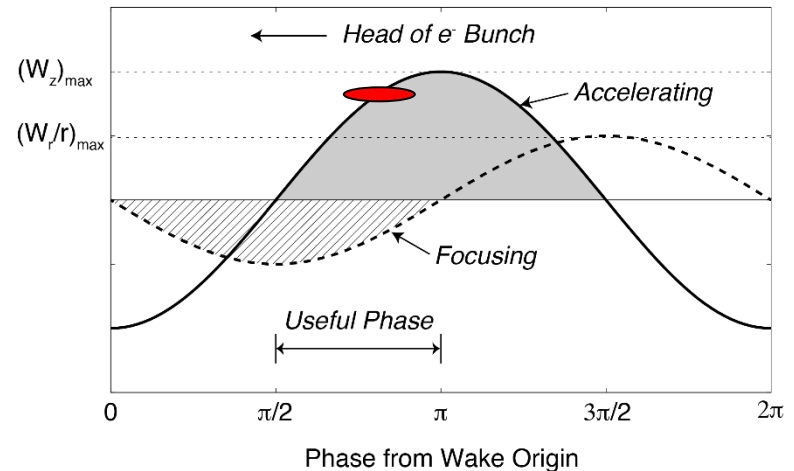
Ultra-fast science

and/or

Increase plasma wavelength

Lower plasma density

Lower accelerating gradient

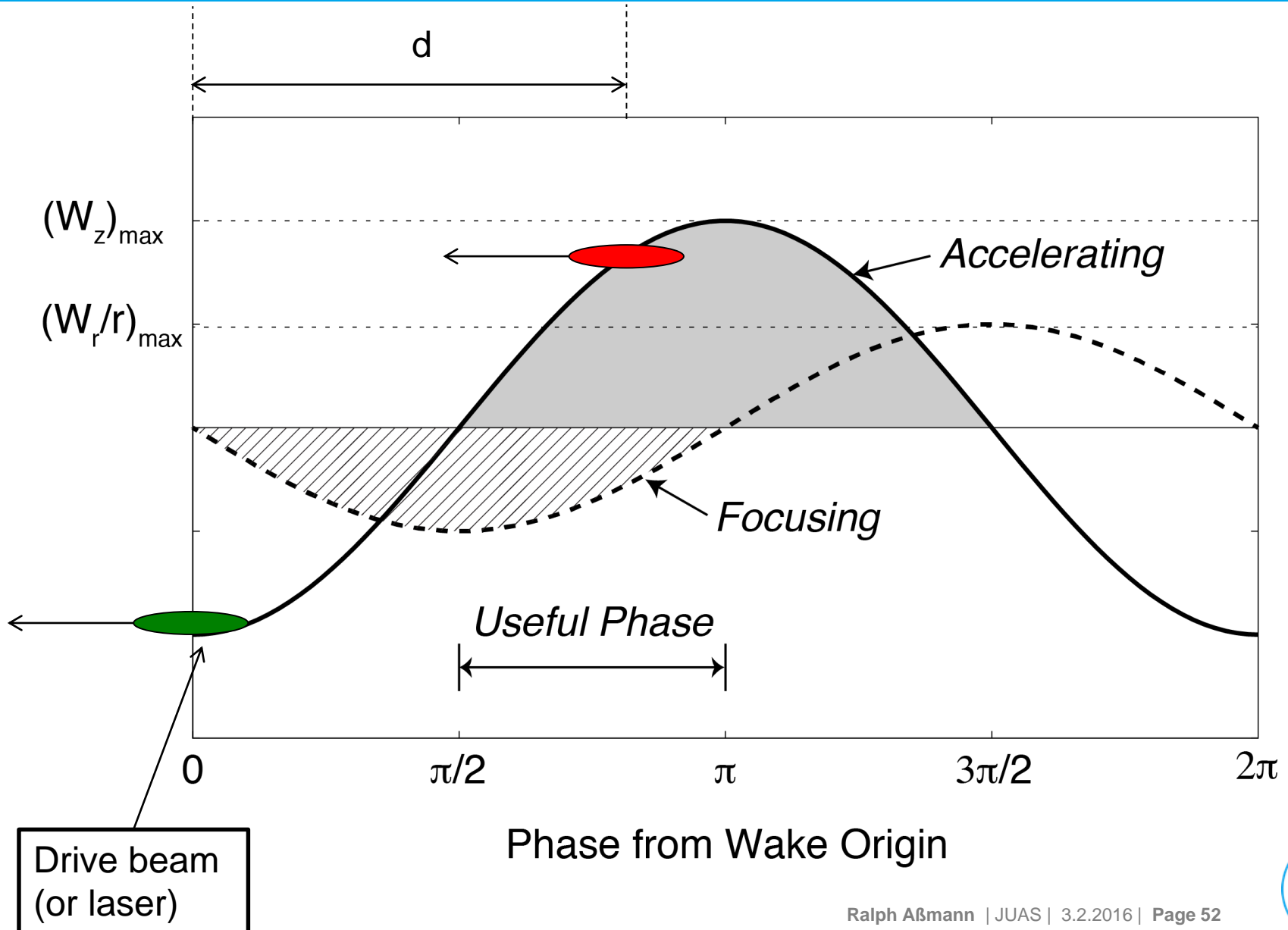


1 fs = 0.3 μm

when travelling with light velocity c

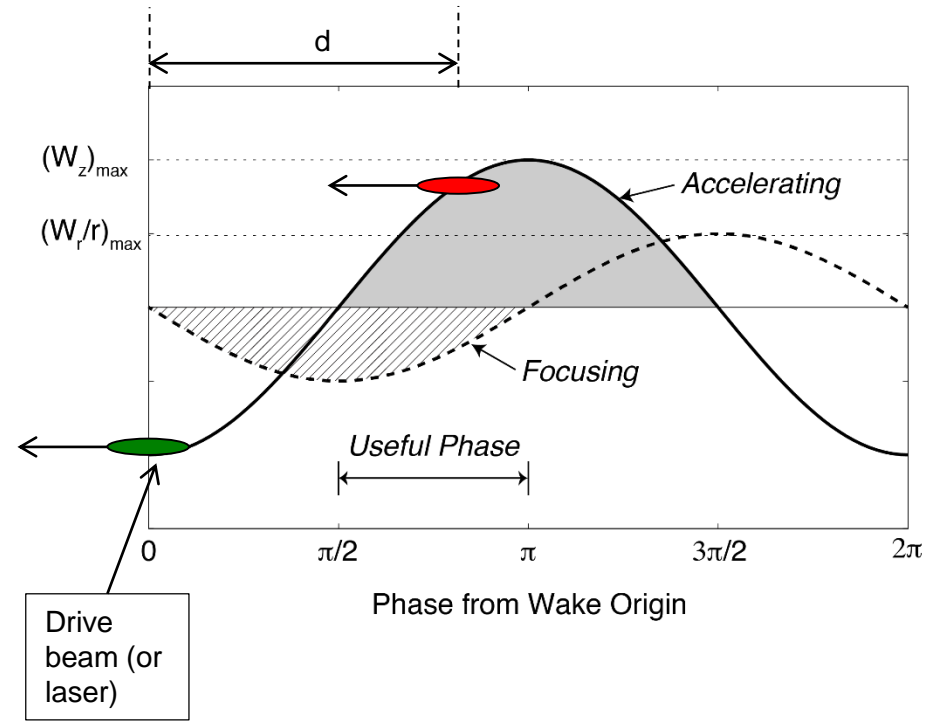


Phase Slippage

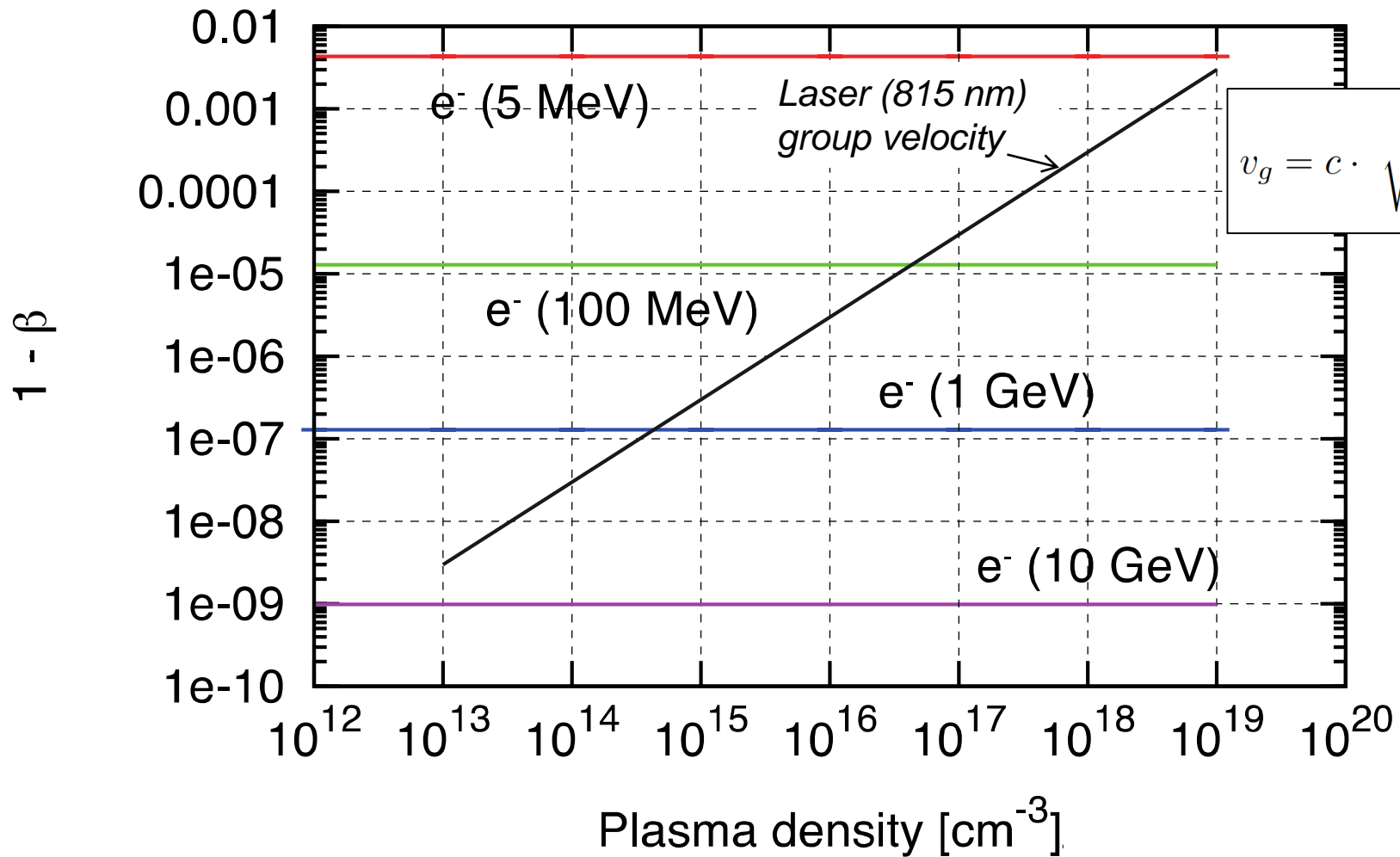


Phase Slippage

- > Keep distance d constant for maximum acceleration and minimum energy spread.
- > **Problem 1:** Drive beam loses energy and (slightly) slows down.
- > **Problem 2:** Accelerated beam starts at low energy, gains energy and (slightly) speeds up.
- > **Problem 3 (for lasers):** Laser group velocity depends on plasma density and is slower than light velocity c .



Dephasing ($\beta = v/c$, here consider relativistic beams)



Dephasing ($\beta = v/c$, here consider relativistic beams)

$$\Delta L = \frac{\Delta v}{c} \cdot L = \Delta\beta \cdot L$$

➤ Imagine **10 GeV beam driver**.

➤ Imagine **initial energy of accelerated electrons to be 100 MeV**.

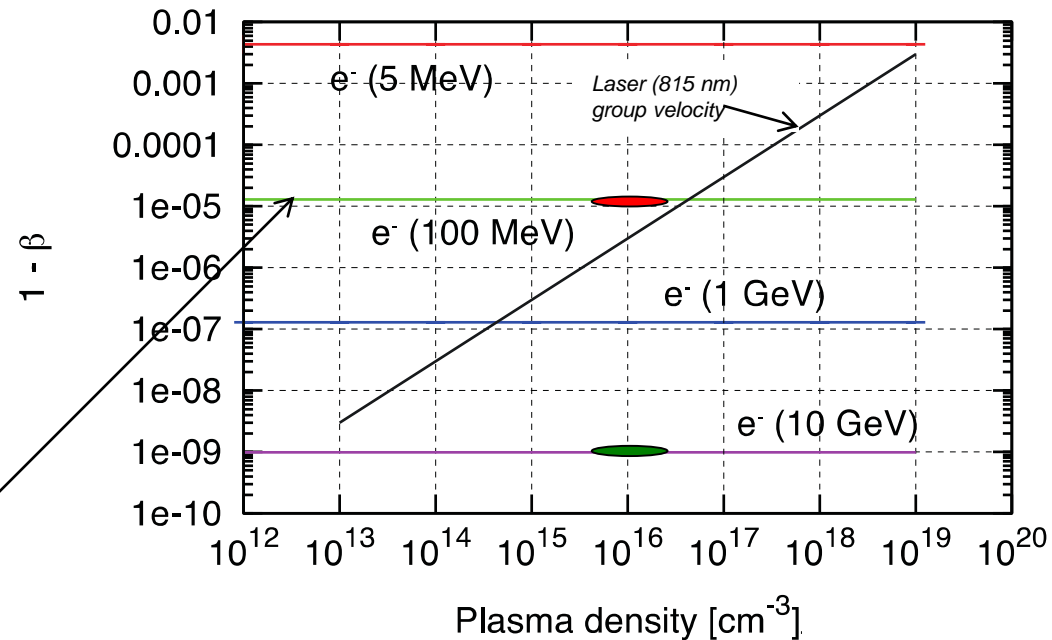
➤ After 1 m slippage by $\approx 10^{-5}$ m = **10 μm** .

➤ Plasma wavelength: **10 μm ($n_0=1e19$) – 1 mm ($n_0=1e15$)**

➤ However:

- Driver electrons are decelerated and slow down.
- Accelerated electrons speed up.

➤ Big advantage of beam-driven...



$$\lambda_p \approx 1\text{mm} \cdot \sqrt{\frac{10^{15}\text{cm}^{-3}}{n_0}}$$



Dephasing ($\beta = v/c$, here consider relativistic beams)

$$\Delta L = \frac{\Delta v}{c} \cdot L = \Delta\beta \cdot L$$

➤ Imagine **10 GeV beam driver**.

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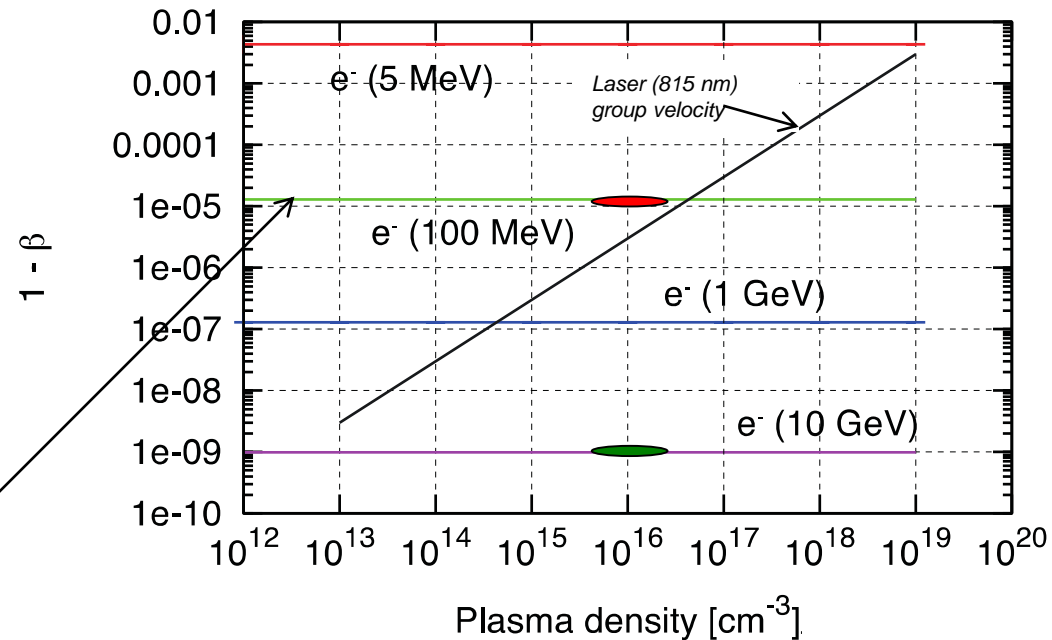
➤ However

- Driver e
- Accele

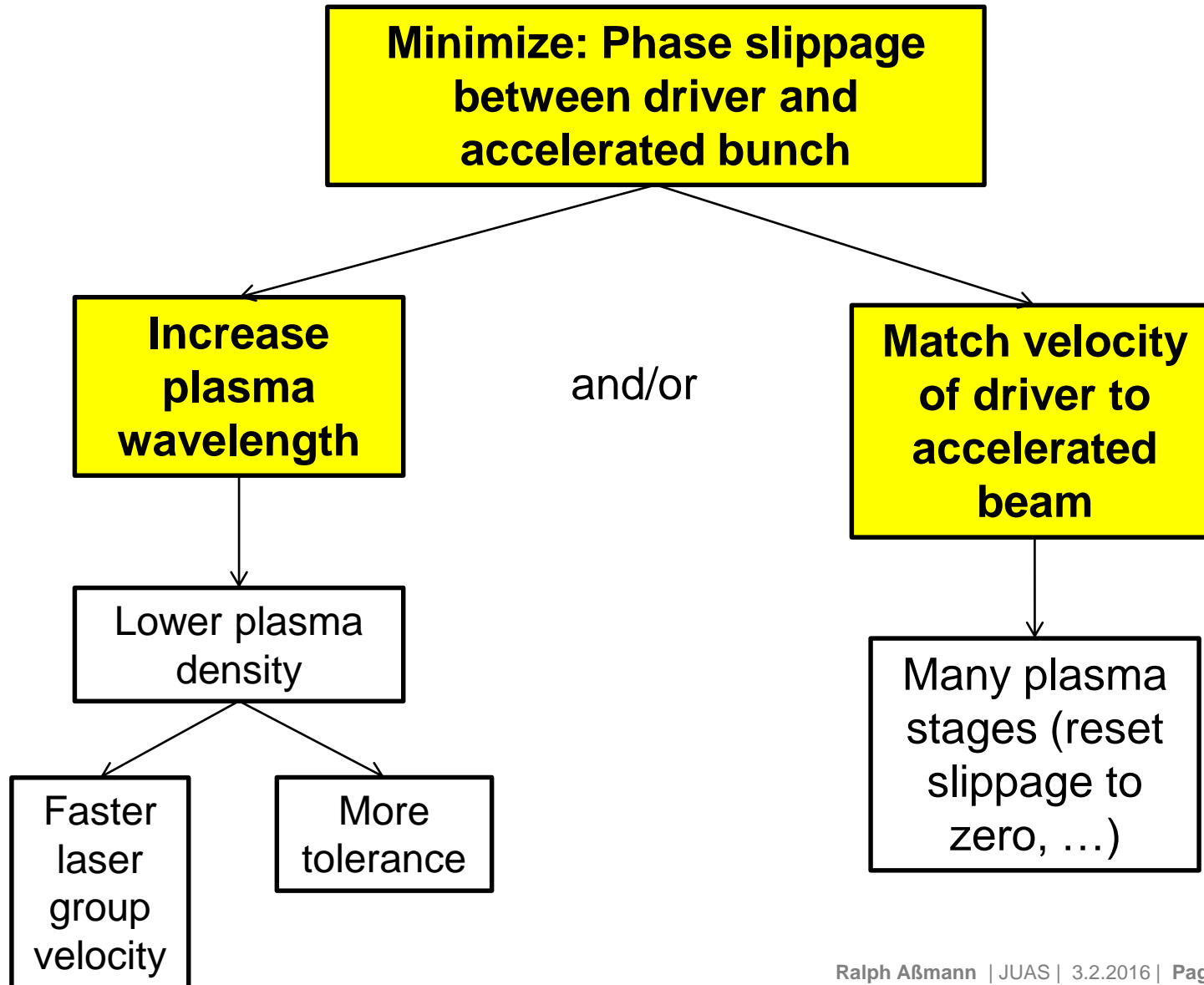
3.6° - 360°

$$\lambda_p \approx 1\text{mm} \cdot \sqrt{\frac{10^{15}\text{cm}^{-3}}{n_0}}$$

➤ Big advantage of beam-driven...



Optimization 2: Phase Slippage



Optimization 3: Stability / Reproducibility

Stabilize: Distance between driver and accelerated bunch.

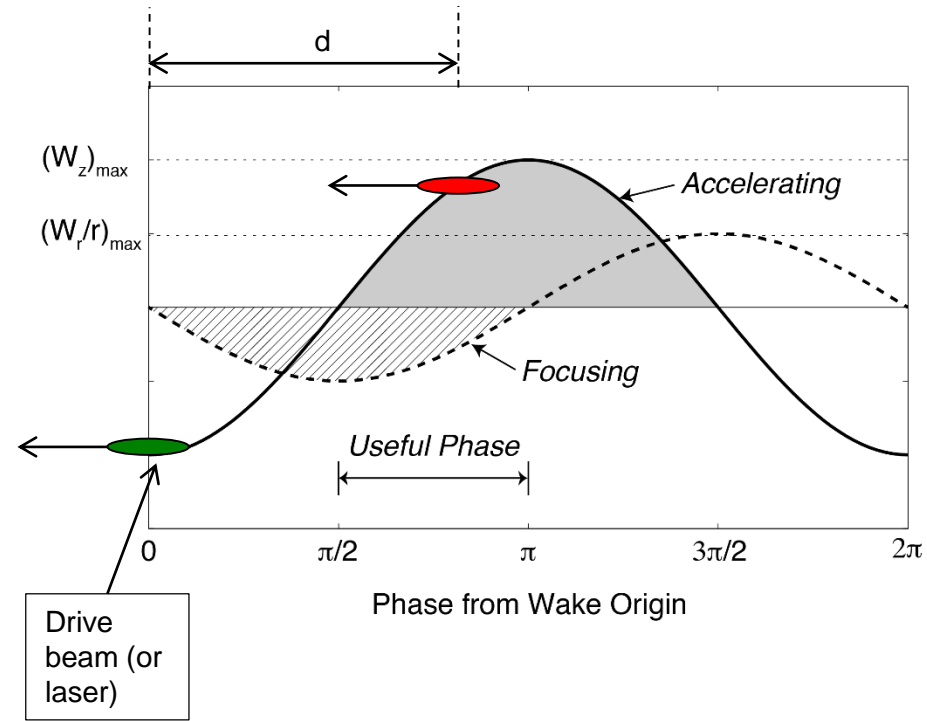
Internally generate accelerated electron bunch

High accelerating fields

High plasma density

Synchronize externally injected beam to driver

Synchronize with **30 fs to 0.3 fs accuracy** for few degree phase stability



Maximum Acceleration

> (Beam 1) **drive beam of N_1 electrons at E_1 (GeV)** that pumps its energy into the plasma wakefield.

(Beam 2) **acc. electron beam of N_2 electrons** gets at **maximum ΔE**

> Energy conservation must be fulfilled:

$$E_{stored,1} = N_1 \cdot \frac{E_1}{(\text{GeV})} \cdot 1.6022 \times 10^{-10} \text{ J}$$

$$E_{stored,2} = N_2 \cdot \frac{\Delta E}{(\text{GeV})} \cdot 1.6022 \times 10^{-10} \text{ J}$$

> From $E_{stored,1} \geq E_{stored,2}$ we find:

$$\Delta E \leq E_1 \cdot \frac{N_1}{N_2}$$



Maximum Acceleration

$$\Delta E \leq E_1 \cdot \frac{N_1}{N_2}$$

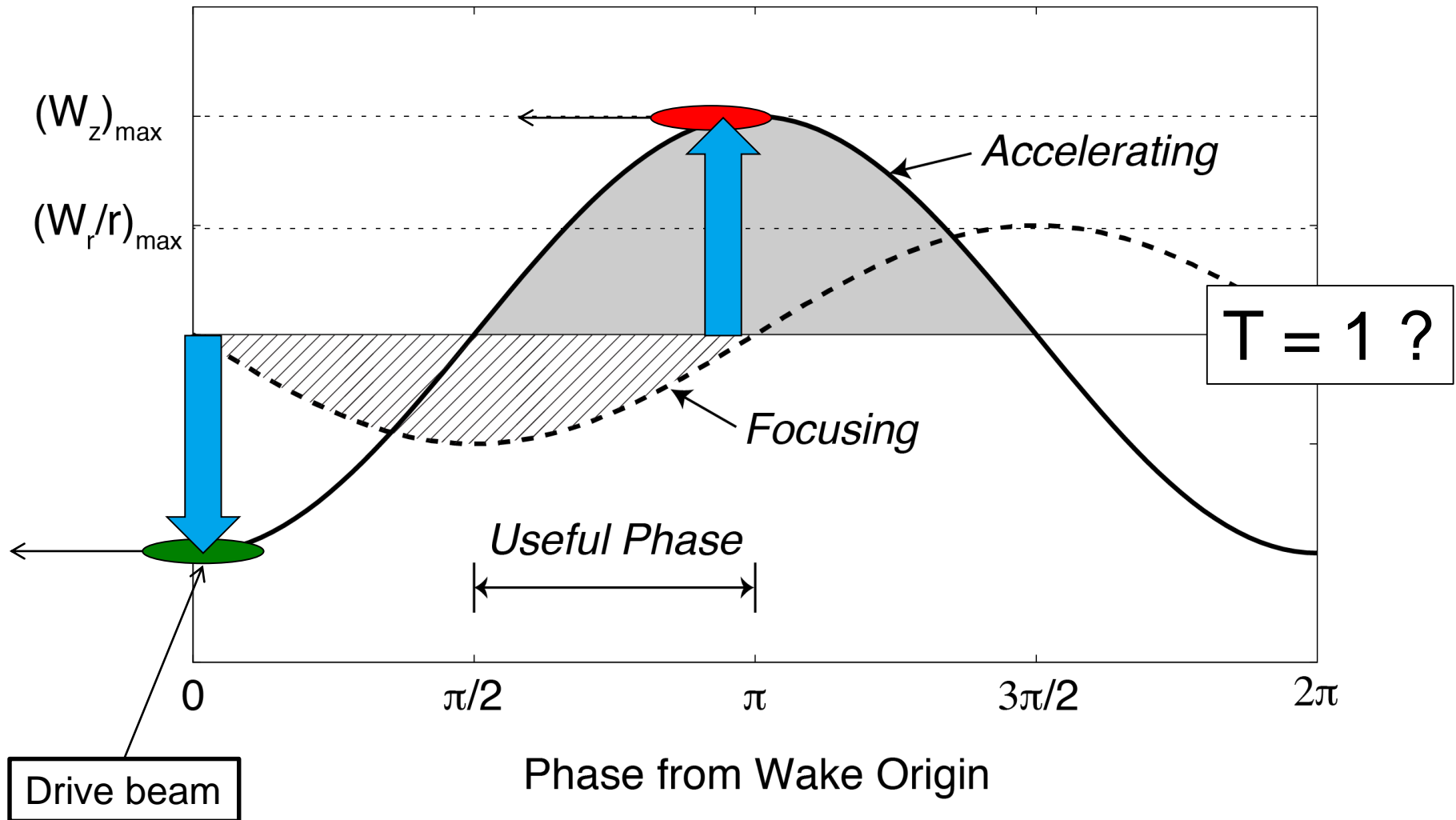
- > Would be great. E.g. take a 1 GeV electron drive beam with 10^{11} electrons to accelerate 10^9 electrons by 100 GeV!
- > This is, however, not possible in reality!
- > Limited by **transformer ratio T** (short, symmetric bunches):

$$T = \left| \frac{\Delta E}{E_1} \right| = \left| \frac{\Delta E_{acc}}{\Delta E_{drive}} \right| \leq 2$$

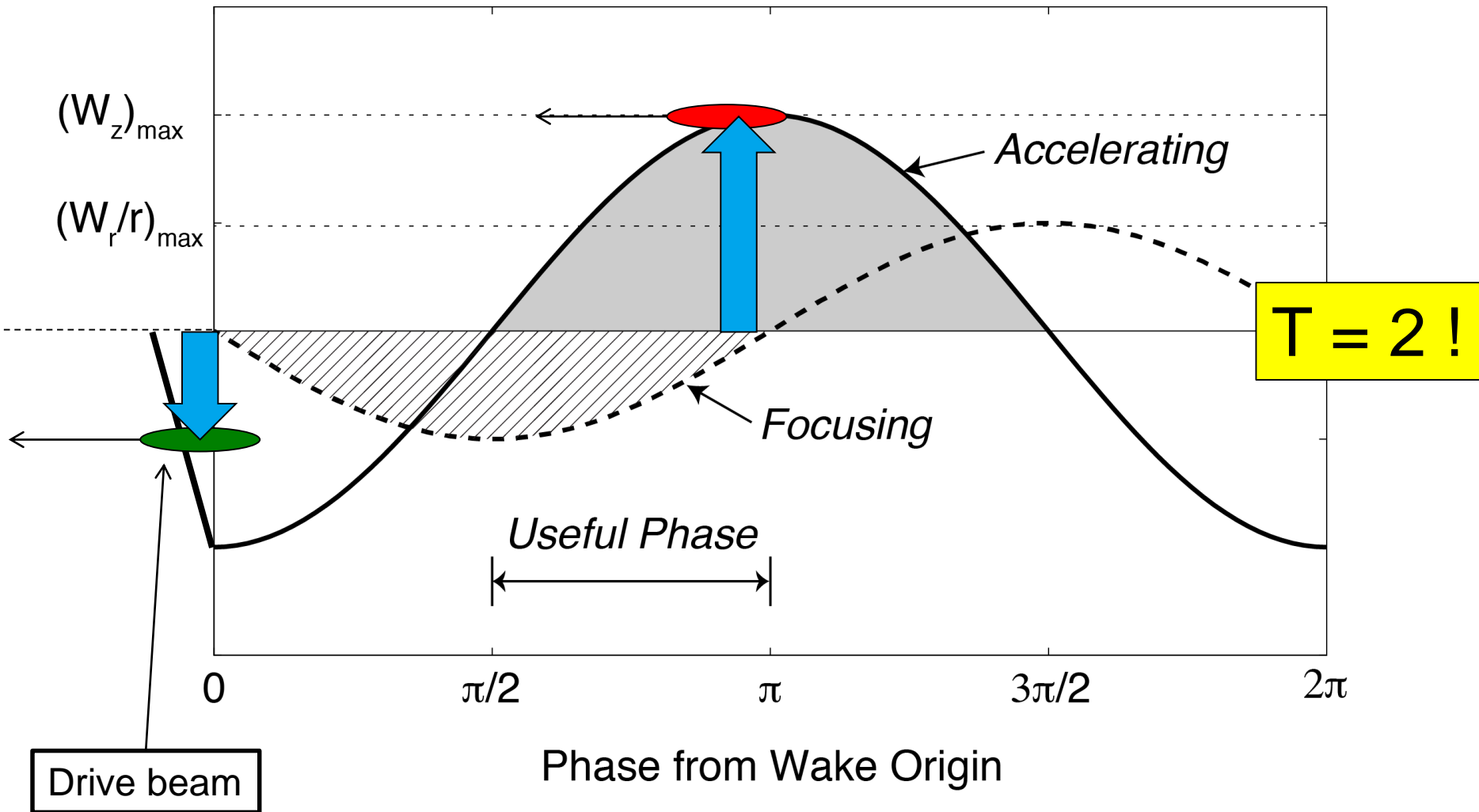
Here it is assumed, drive beam loses all its energy



Transformer Ratio (Short Symmetric Bunches)



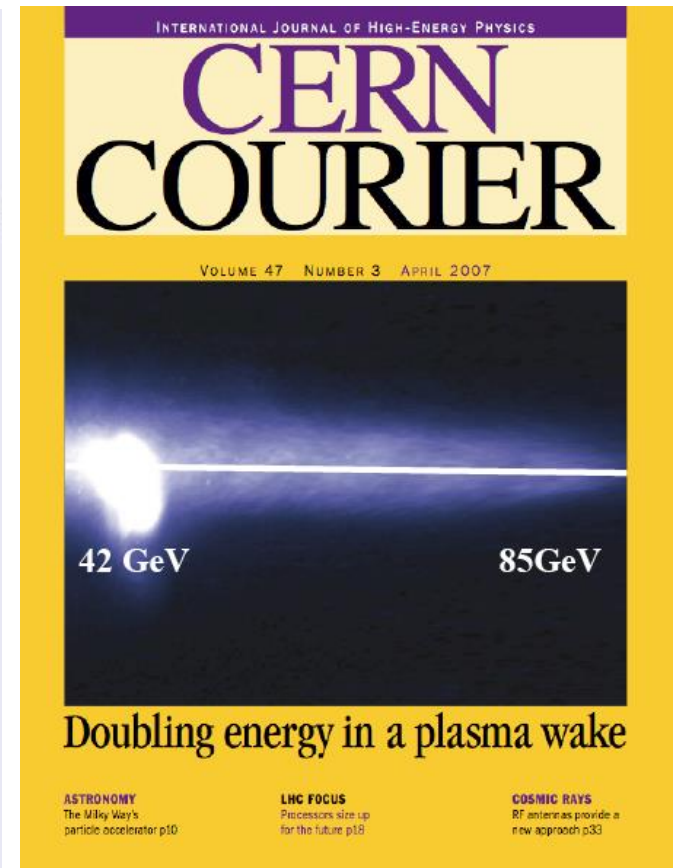
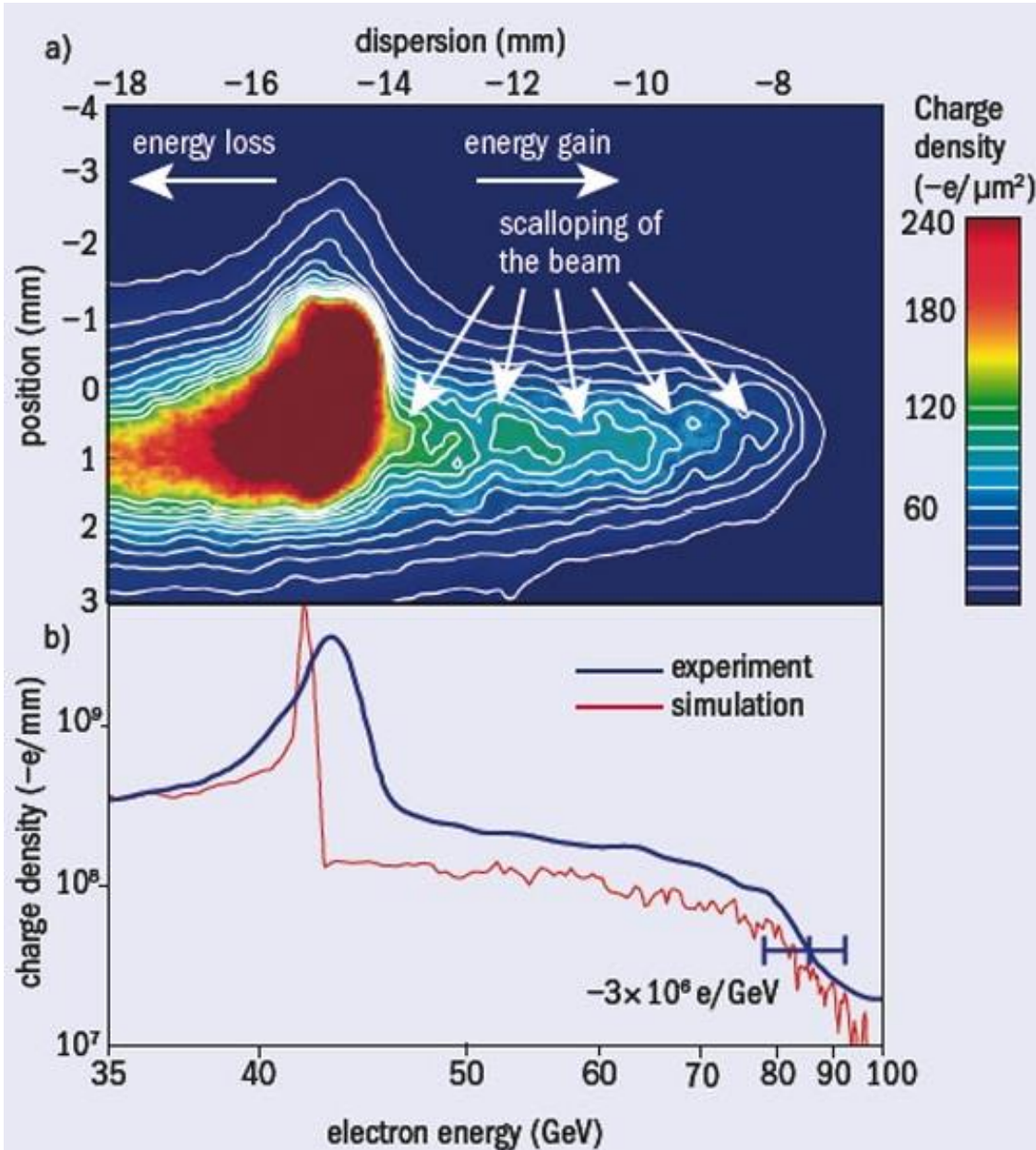
Transformer Ratio (Short Symmetric Bunches)



Head of driver sees no decelerating field



Record Acceleration: 42 GeV (beam-driven)



E167 collaboration
SLAC, UCLA, USC

I. Blumenfeld et al, Nature 445,
p. 741 (2007)



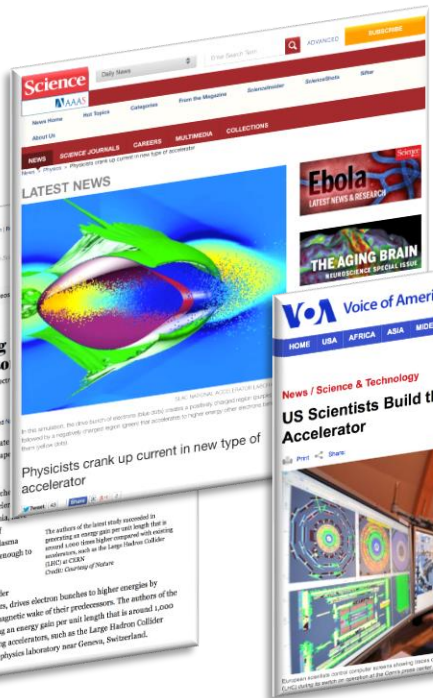
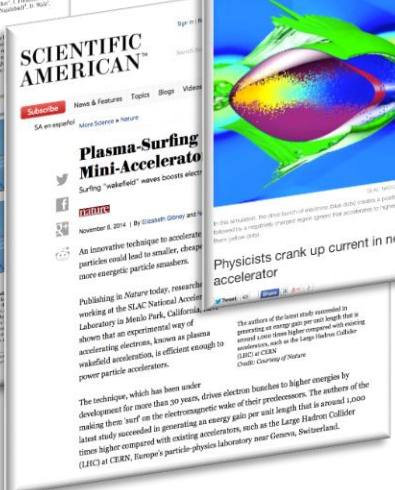
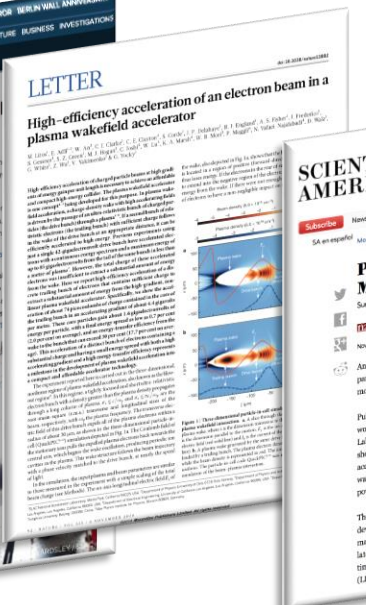
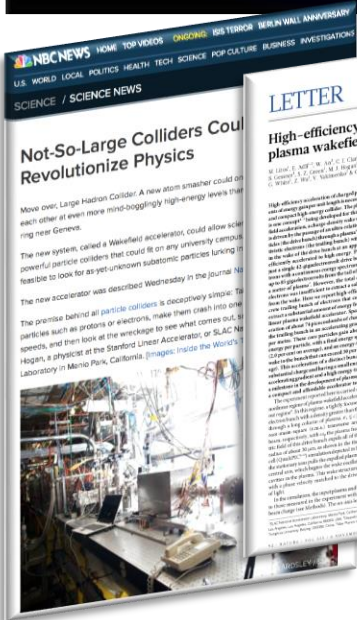


FACET: A National User Facility based on high-energy beams and their interaction with plasmas and lasers

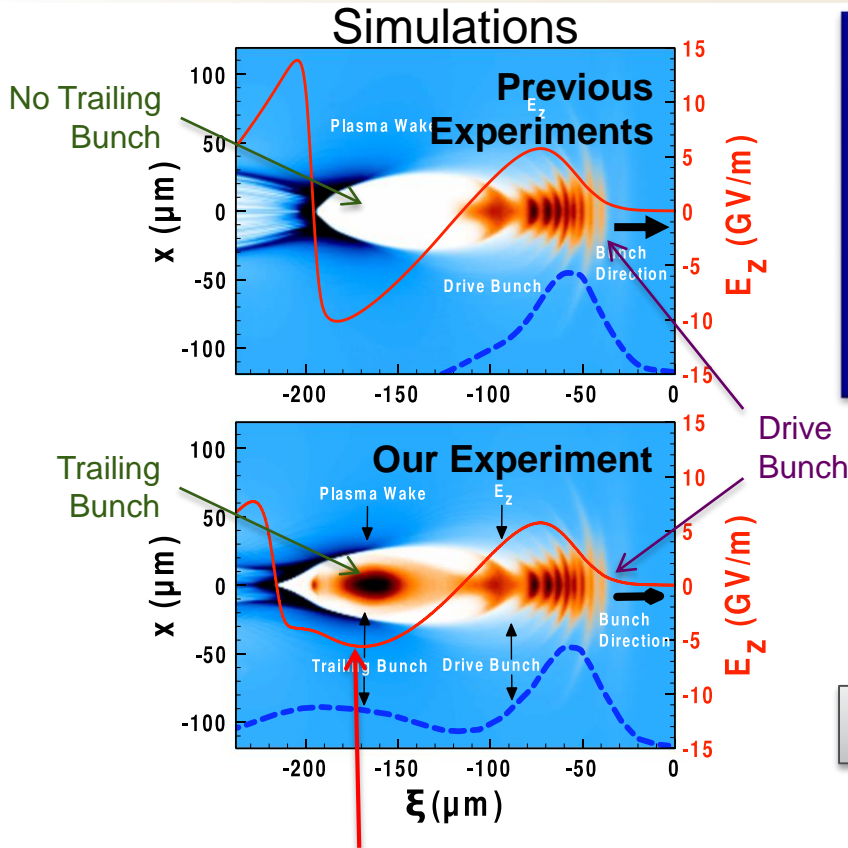
UCLA SLAC

- Facility hosts more than 150 users, 25 experiments
- One high profile result a year
- Priorities balanced between focused plasma wakefield acceleration research and diverse user programs with ultra-high fields

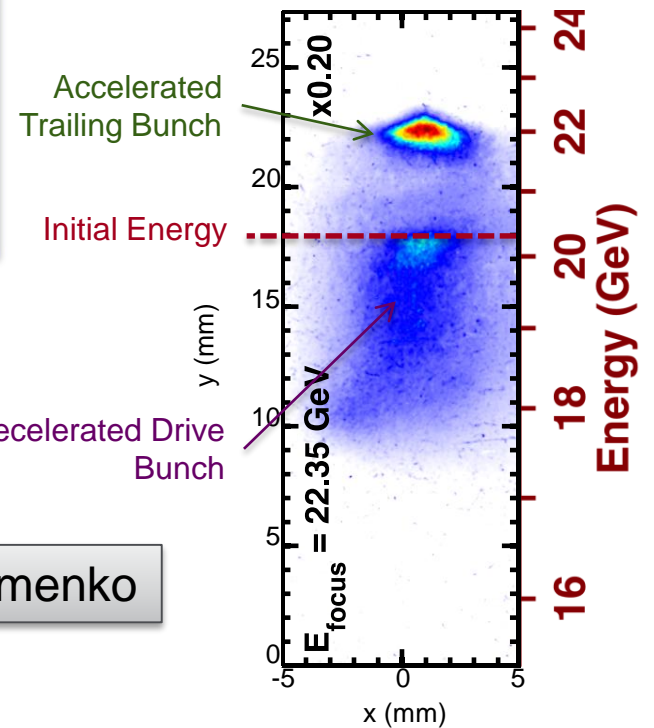
Slide: V. Yakimenko



High-Efficiency Acceleration of an Electron Bunch in a Plasma Wakefield Accelerator



Energetically Dispersed Beam After Plasma (Data)



Slide: V. Yakimenko

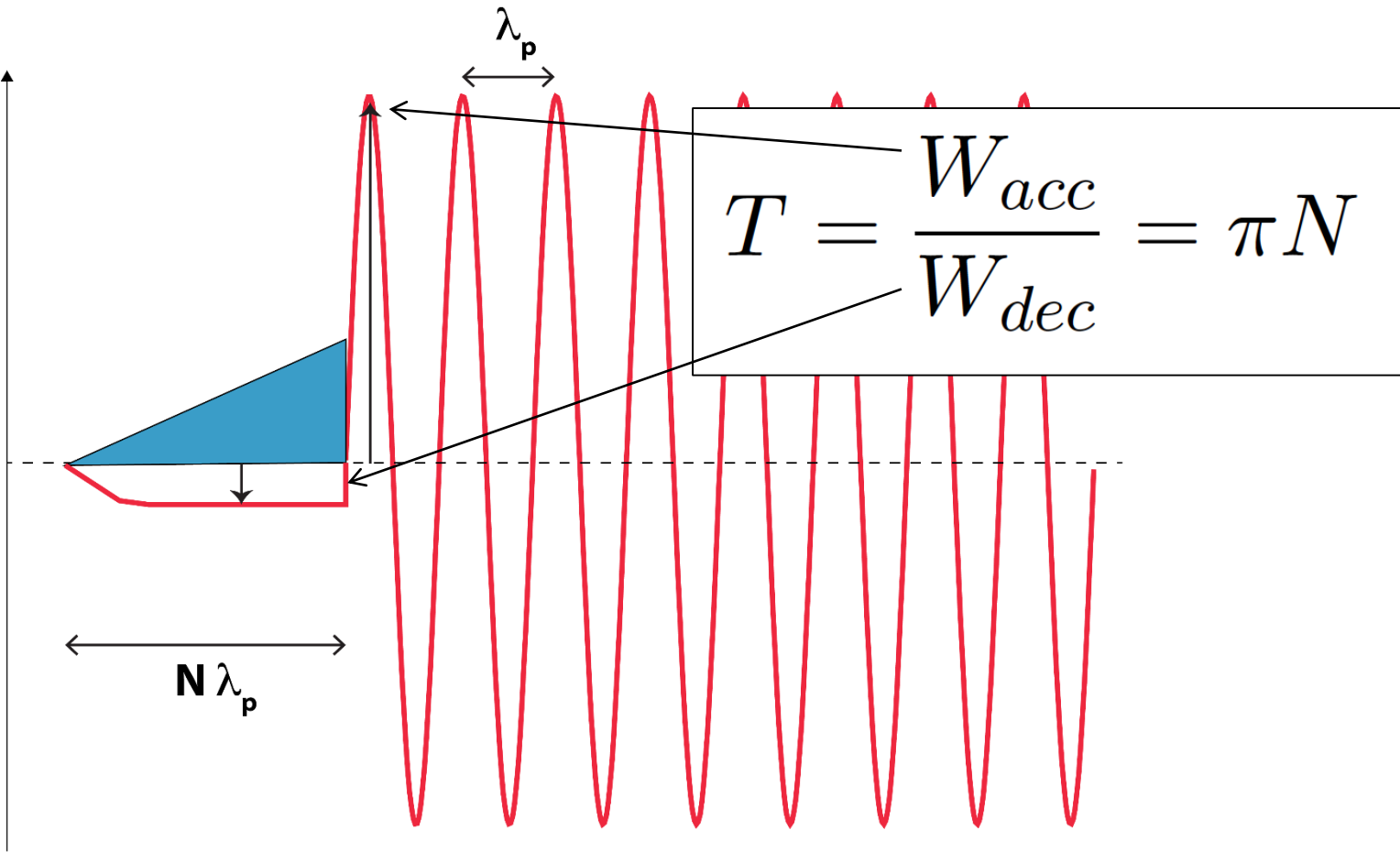
- Electric field in plasma wake is loaded by presence of trailing bunch
- Allows efficient energy extraction from the plasma wake

This result is important for High Energy Physics applications that require very efficient high-gradient acceleration

Increasing the Transformer Ratio

Longitudinal
Field W_z

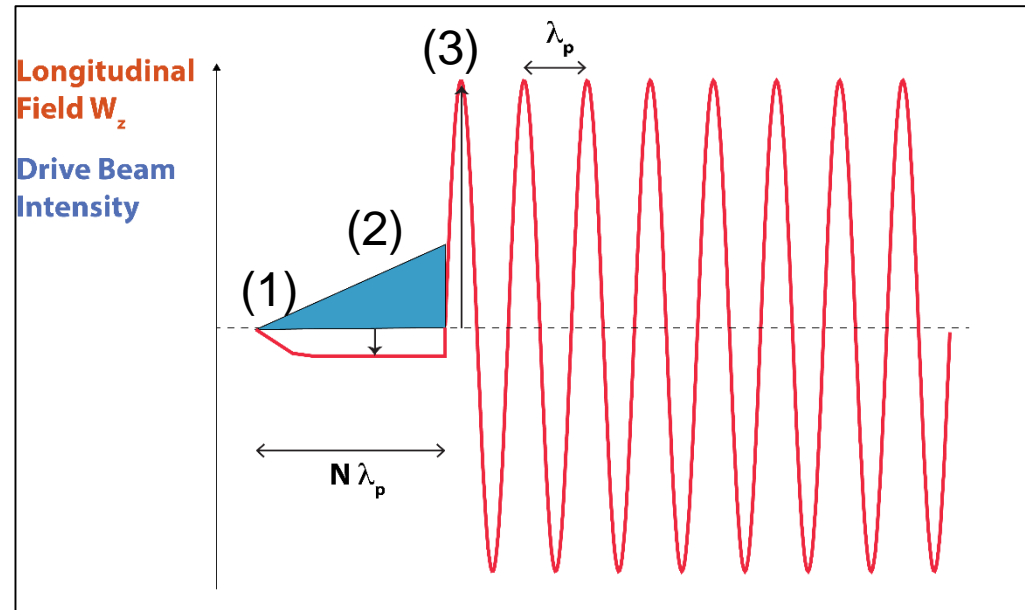
Drive Beam
Intensity



Physics of the Triangular Bunch Driver

(1) **Leading component** acts as precursor:

- Gives plasma electrons an impulse so that they flow out of the beam driver region with an increasing flux.
- End of precursor: depletion rate of plasma electrons is balanced by replacement rate of electrons in the drive bunch.



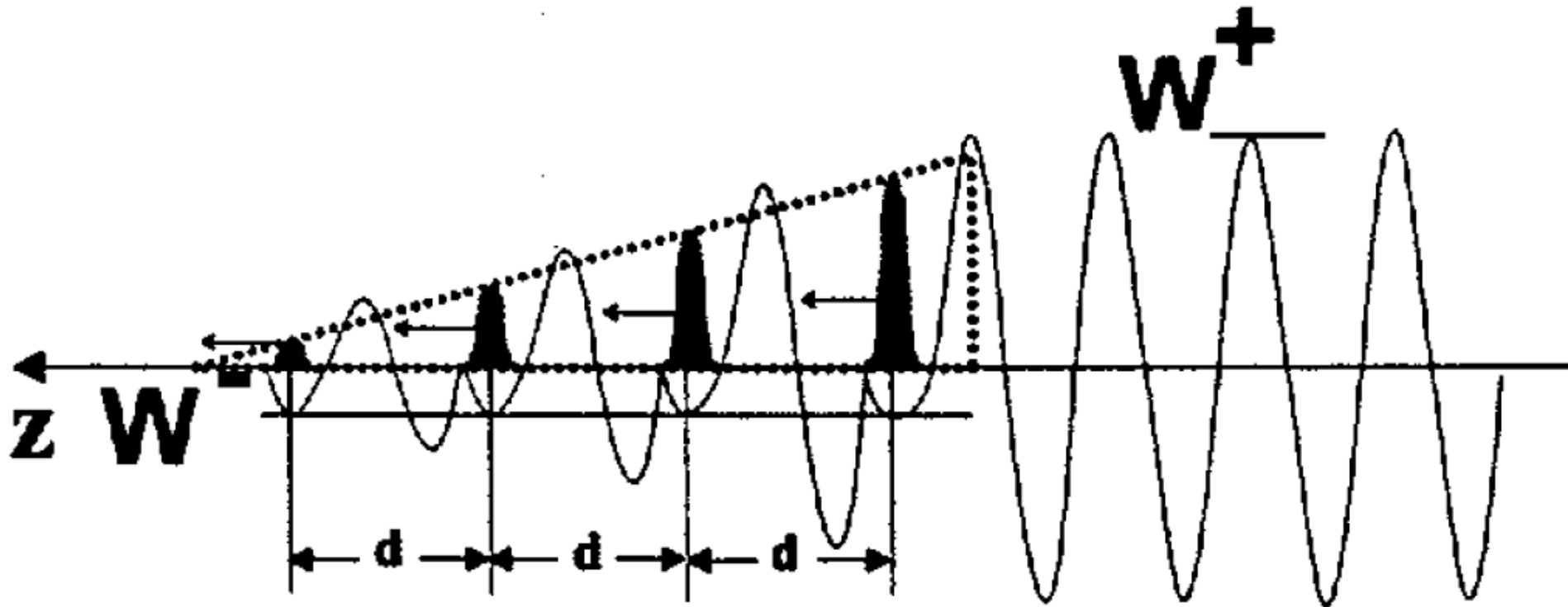
(2) **Long ramp component:**

- Charge neutrality is maintained. Same decelerating field maintained.

(3) **After driving bunch** (sharp edge):

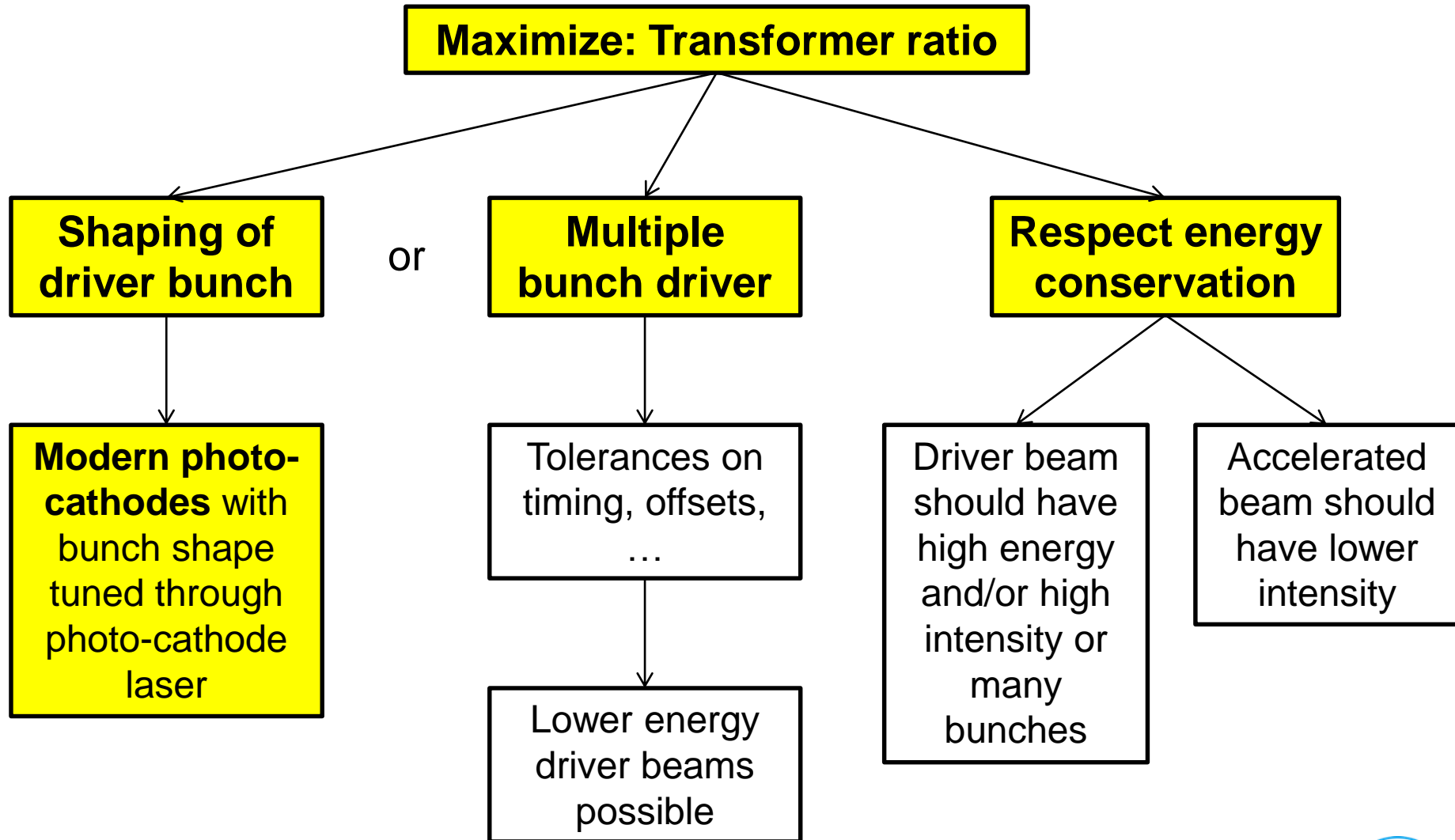
- Plasma channel becomes non-neutral.
- Plasma electrons are strongly attracted back to the ions and large scale plasma oscillations begin.

Alternative: Multi-bunch driven PWFA



Power et al, 2000

Optimization 4: Maximum Energy Gain



1. Accelerators – From Conventional Techniques to Plasmas
2. The Linear Regime
3. **The Non-Linear Regime**
4. Tolerances
5. Outlook for Europe

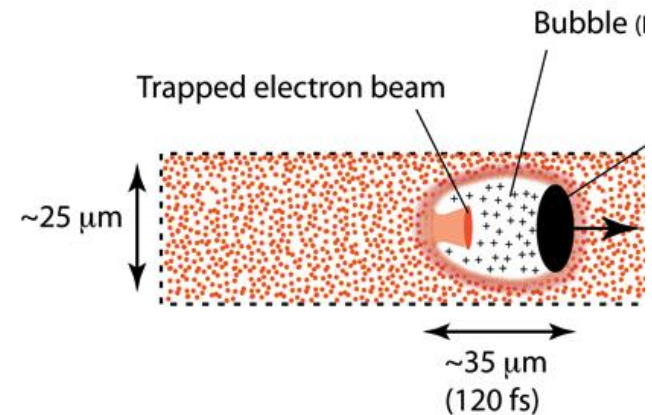
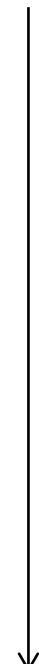
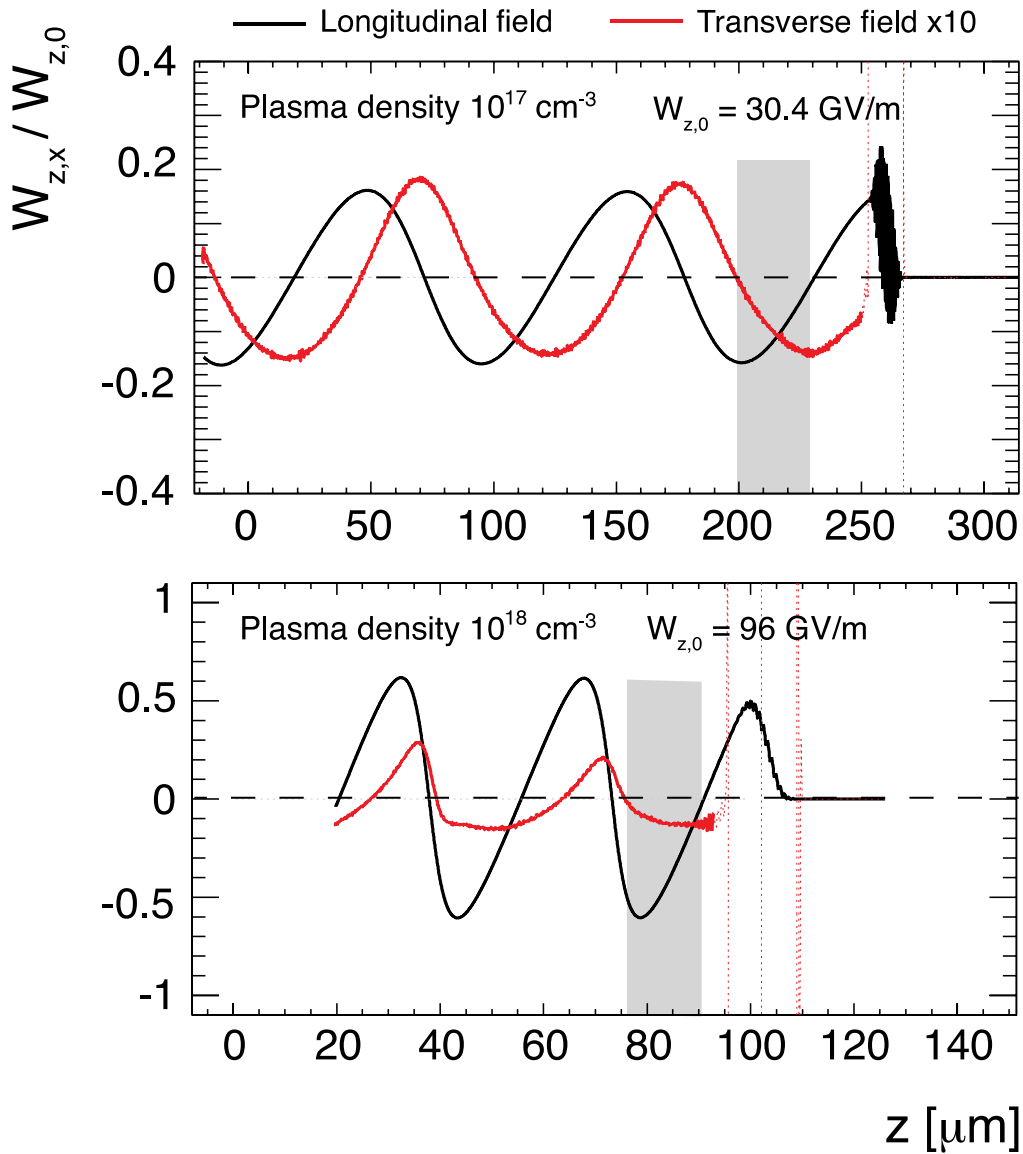


The Non-Linear Regime

- Blow-out and Non-Linear Regime
- Wave-Breaking as Limit to Maximum Energy Gain
- Self-injection in wave-breaking regime
- Hybrid Schemes → Trojan Horse



Fields Calculated with OSIRIS Code: Non-Linear Regime



Plasma Accelerator Physics I

- > A plasma of density n_0 (same density electrons - ions) is characterized by the **plasma frequency**:

$$\omega_p = \sqrt{\frac{n_0 e^2}{\epsilon_0 m_e}}$$

- > This translates into a **wavelength** of the plasma oscillation:

$$\lambda_p \approx 1\text{mm} \cdot \sqrt{\frac{10^{15}\text{cm}^{-3}}{n_0}}.$$

0.3 mm for $n_0 = 10^{16}\text{cm}^{-3}$

- > The wavelength gives the longitudinal size of the plasma cavity...
Lower plasma density is good: larger dimensions.



Plasma Accelerator Physics II

- > The plasma oscillation leads to **longitudinal accelerating fields** with a gradient of (higher plasma densities are better):

$$W_z = 96 \frac{V}{m} \cdot \sqrt{\frac{n_0}{\text{cm}^{-3}}} \propto N_b / \sigma_z^2 \quad \text{9.6 GV/m for } 10^{16} \text{ cm}^{-3}$$

- > The **group velocity of the laser in a plasma** is as follows for $\omega_p \ll \omega_l$: (note ω_l is laser frequency)

$$v_g = c \cdot \sqrt{1 - \frac{\omega_p^2}{\omega_l^2}}$$

- > The laser-driven wakefield has a lower velocity than a fully relativistic electron \rightarrow slippage and dephasing. Lower densities are better.



Plasma Accelerator Physics III

- > The ion channel left on axis, where the beam passes, induces an **ultra-strong focusing field**. In the simplest case:

$$g = 960 \pi \cdot \left(\frac{n_0}{10^{14} \text{ cm}^{-3}} \right) \text{ T/m} \quad \text{300 kT/m for } 10^{16} \text{ cm}^{-3}$$

- > This can be converted into a **optical beta function** (lower density is better, as beta function is larger)::

$$k_\beta^2 = 0.2998 \frac{g}{E} \quad \beta = \frac{1}{k_\beta} \quad \beta = 1.1 \text{ mm for } 100 \text{ MeV}$$

- > The **phase advance** in the plasma channel is rapid:

$$\psi(s) = \int k_\beta s \, ds \propto \sqrt{E}$$



Plasma Accelerator Physics IV

- > The **matched beam size** in the ion channel is small:

$$\sigma_0 = \sqrt{\beta \varepsilon}$$

$$\sigma_0 = 1.3 \mu\text{m for } \gamma \varepsilon = 0.3 \mu\text{m}$$

- > Offsets between laser and beam centres will induce betatron oscillations. Assume: full dilution into emittance growth (energy spread and high phase advance).
- > Tolerances for **emittance growth** due to offsets $\Delta x = \sigma_x$:

$$\frac{\Delta \varepsilon}{\varepsilon_0} = \left(\frac{\sigma_x}{\sigma_0} \right)^2$$

$$100\% \text{ for } 1.3 \mu\text{m offset}$$

- > Lower plasma density better: larger matched beam size, bigger tolerances.



Makes Things Difficult...

> Conventional acceleration structures:

- Optimized to provide longitudinal acceleration and no transverse forces on the beam.
- Due to imperfections, transverse forces can be induced. These “wakefields” caused major trouble to the first and only linear collider at SLAC.

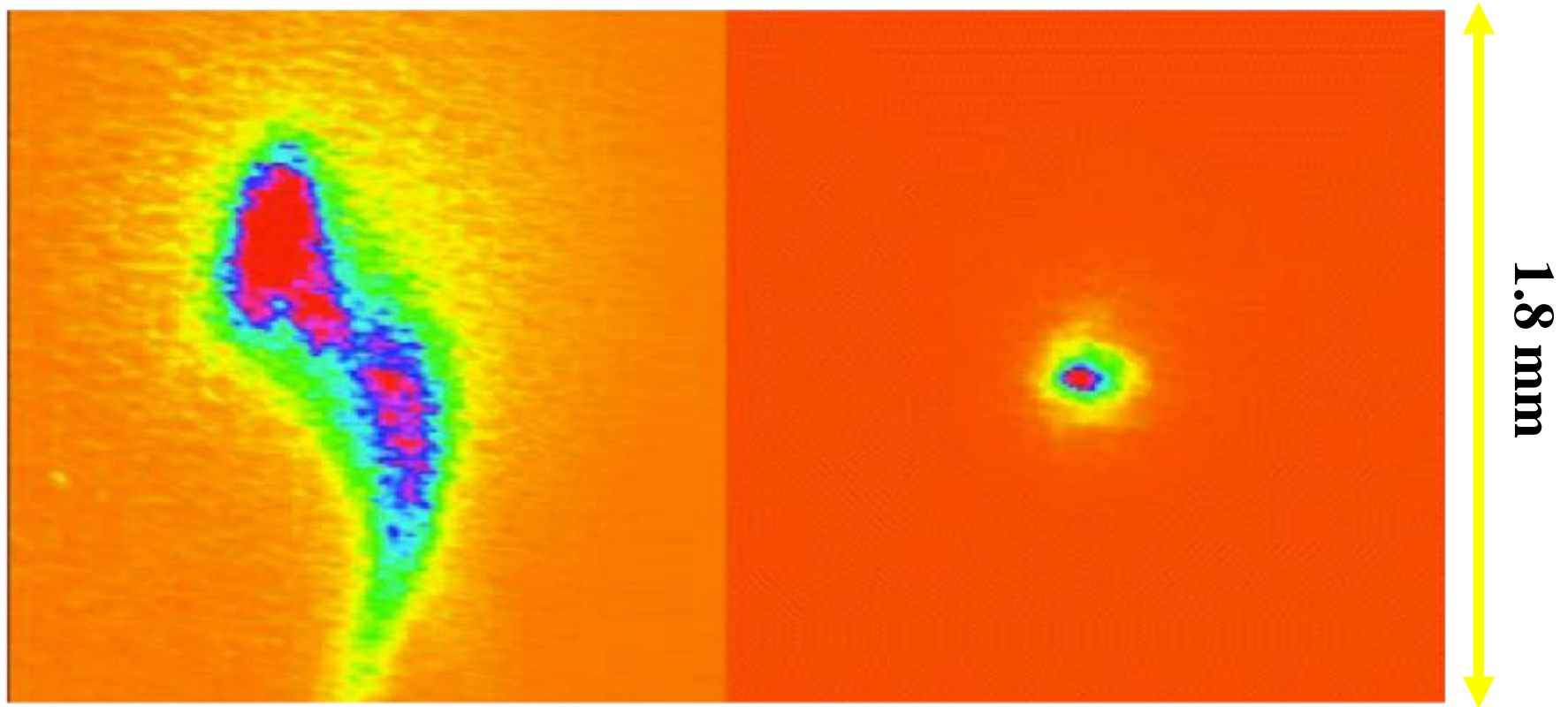
> Plasma acceleration:

- Ultra-strong longitudinal fields → high accelerating gradient.
- Ultra-strong transverse fields → transverse forces cannot be avoided and must be controlled.

> For fun: A look at the SLAC linac beam before entering the plasma!



Seeing Electron Beam...



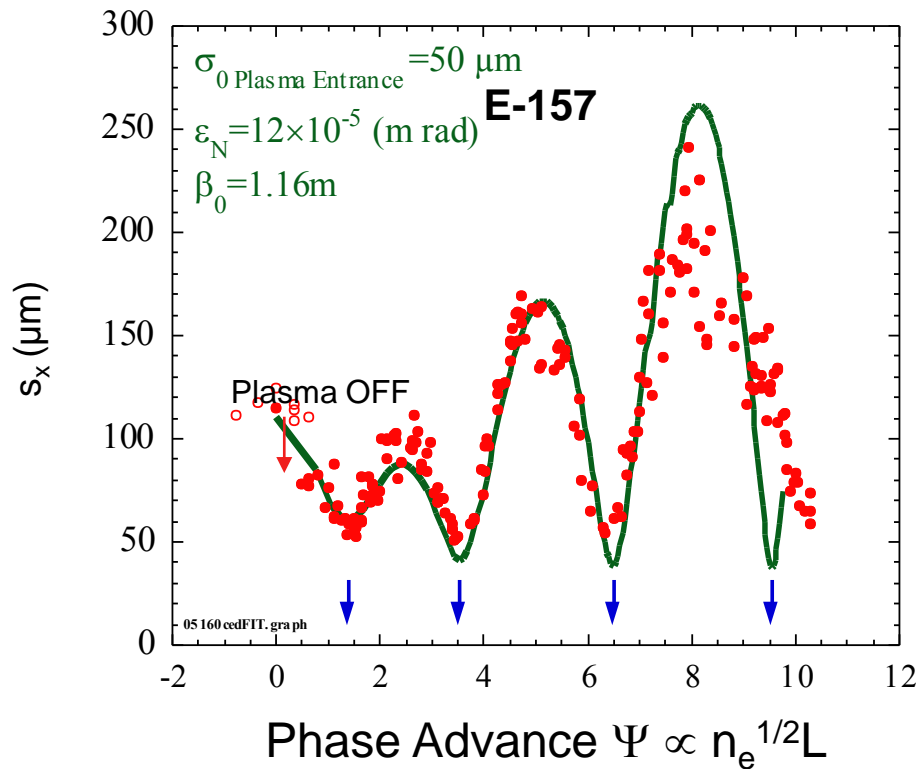
The transverse and longitudinal fields of the accelerator are set up to achieved small transverse beam sizes (right).

~ 2×10^{10} electrons, 30 GeV

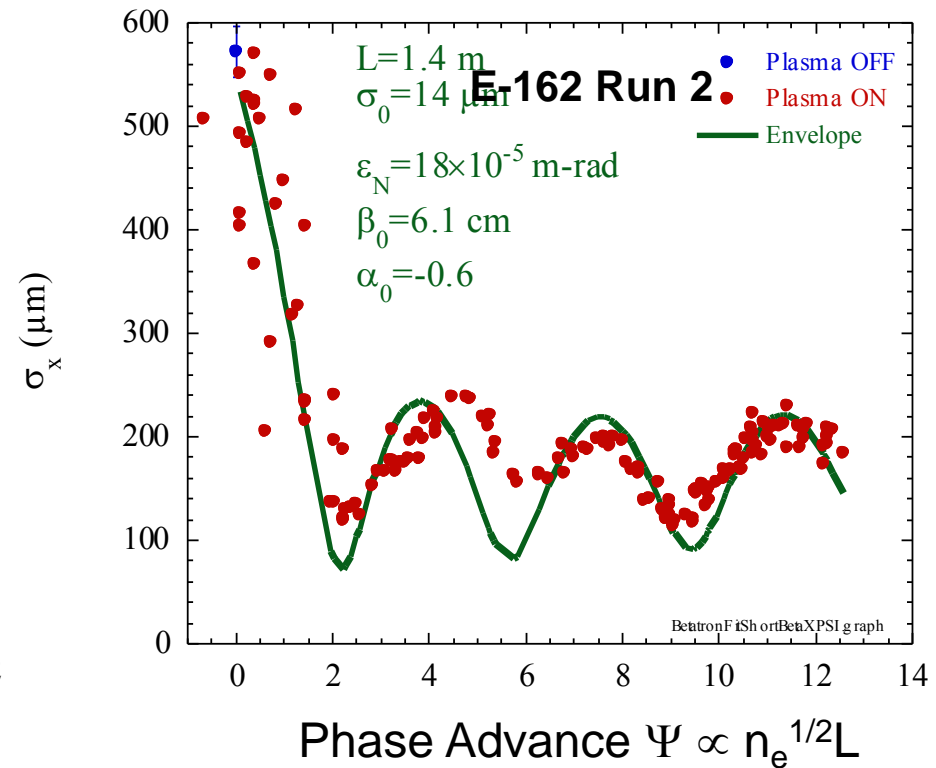
Beam Propagation Through A Long Plasma



- Smaller “matched” beam size at the plasma entrance reduces amplitude of the betatron oscillations measured at the OTR downstream of the plasma
- Allows stable propagation through long plasmas (> 1 meter)



C. E. Clayton et al., PRL 1/2002



E-157/E-162 collaboration

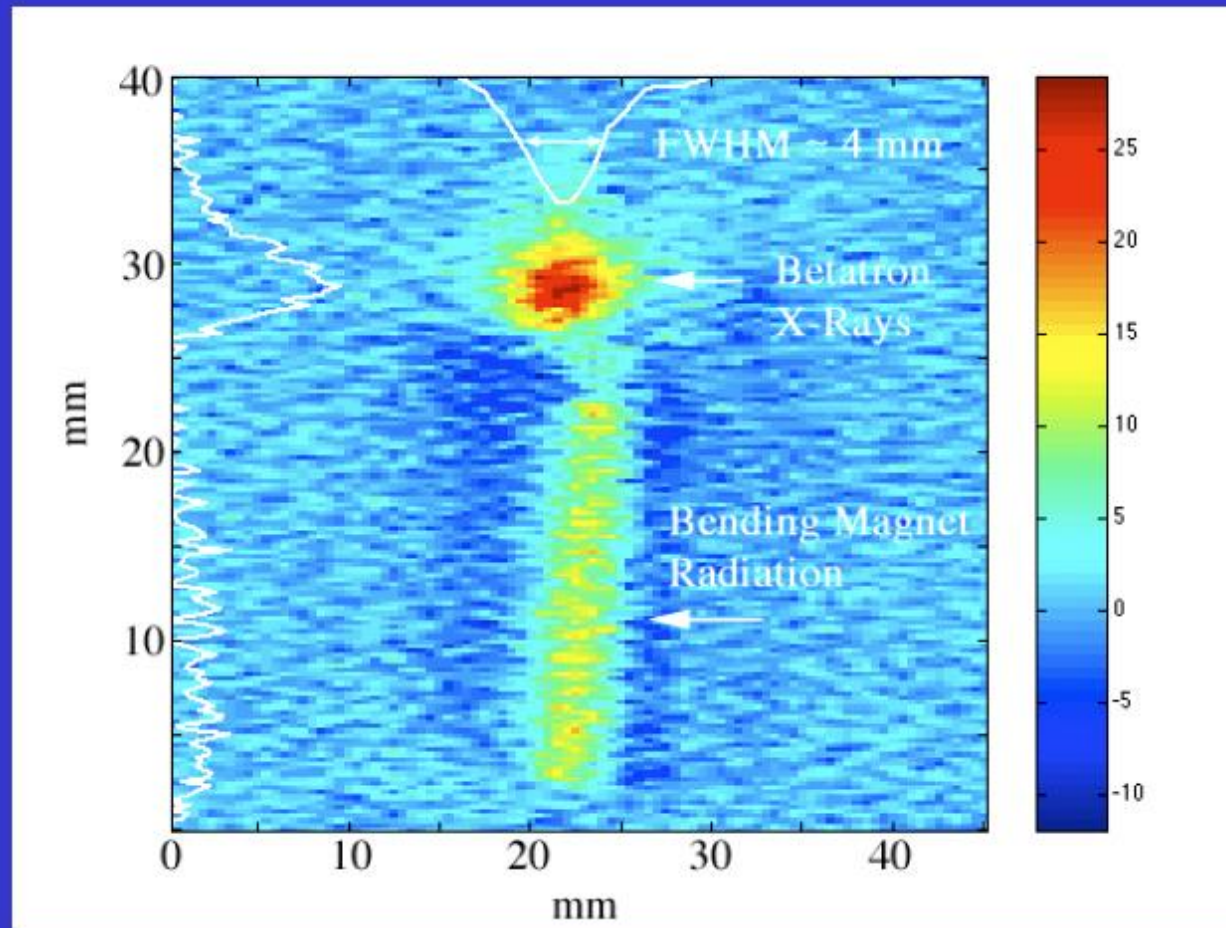
Betatron Radiation of X-rays



UCLA



Plasma focusing strength of 6000T/m acts as a strong undulator



Peak brightness $\sim 10^{19}$ photons/sec-mm²-mrad²-.1%bw!

Wave Breaking



Water velocity becomes larger than phase velocity of the wave

→ **Wave is breaking...**

Dawson 1959: if plasma modeled with one-dimensional sheets, then wave breaking equivalent to crossing of neighboring sheets.

Wavebreaking Limit

- > Dawson 1959: if plasma modeled with one-dimensional sheets, then **wave breaking equivalent to crossing of neighboring sheets**.
- > Non-relativistic wavebreaking field E_0 :

$$W_z = 96 \frac{V}{m} \cdot \sqrt{\frac{n_0}{\text{cm}^{-3}}}$$



Wave Breaking in Plasma Wakefields

- > Relativistic wavebreaking: **capturing of electrons with velocity close to phase velocity of plasma wave → absorption of energy in plasma wave.**
- > Relativistic wavebreaking field (**higher fields than E_0 are possible**):

$$E_{\text{WB}} = \sqrt{2}(\gamma_p - 1)^{1/2} E_0$$

$$\gamma_p = (1 - v_p^2/c^2)^{-1/2}$$

- > Thermal electron effects lead to reduction in wavebreaking field.
Physics: A large fraction of the electron distribution will become trapped in the plasma wave → wave breaks.



Trapped in the Breaking Wave



Appl. Phys. B 74, 355–361 (2002)

DOI: 10.1007/s003400200795

Applied Physics B

Lasers and Optics

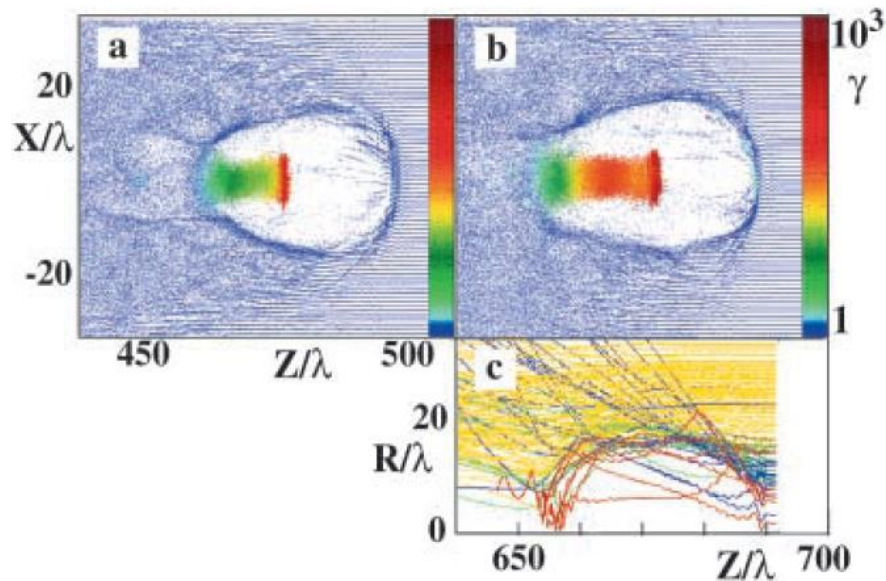
A. PUKHOV^{1,✉}

J. MEYER-TER-VEHN²

Laser wake field acceleration: the highly non-linear broken-wave regime

¹ Institut für Theoretische Physik I, Heinrich-Heine-Universität Düsseldorf, 40225 Düsseldorf, Germany

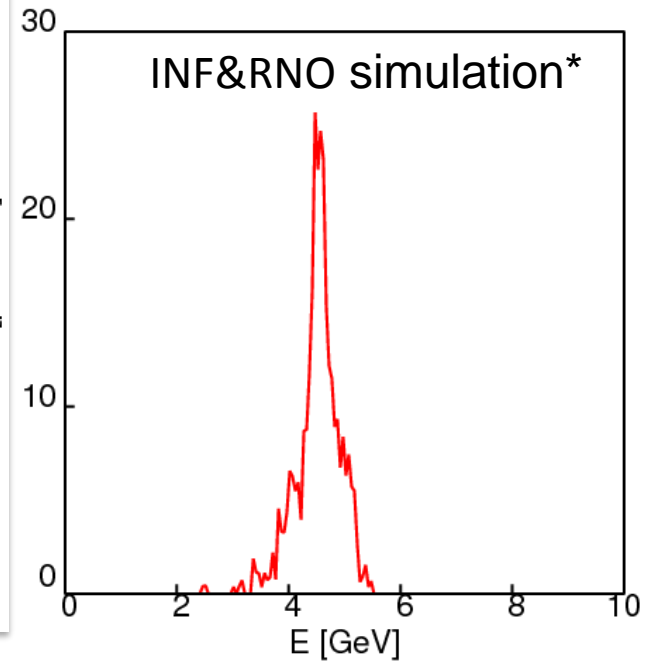
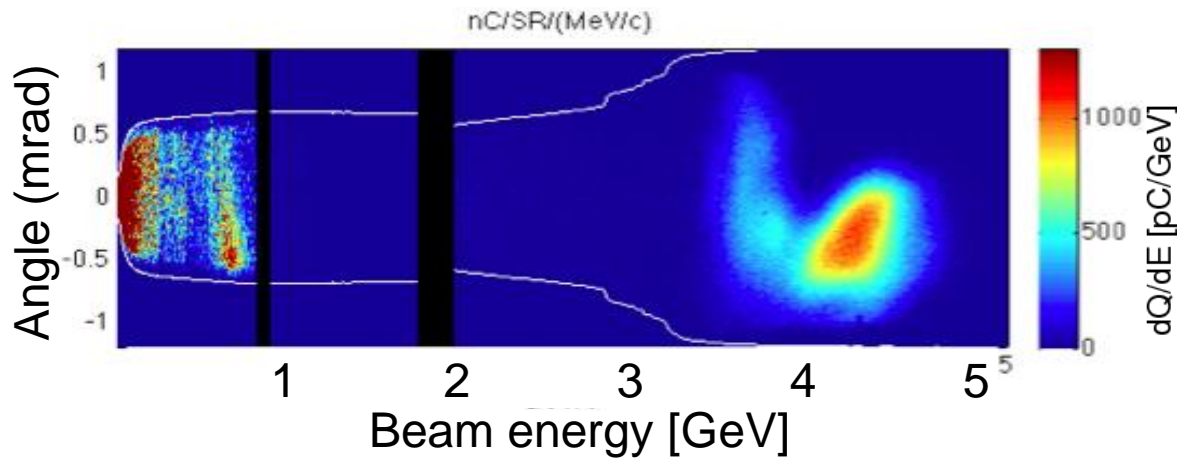
² Max-Planck-Institut für Quantenoptik, Hans-Kopfermann-Str. 1, 85748 Garching, Germany



4.25 GeV beams have been obtained from 9 cm plasma channel powered by 310 TW laser pulses (15 J)

*C. Benedetti et al., proceedings of AAC2010, proceedings of ICAP2012

Electron beam spectrum



Slide: W. Leemans

- **Laser** (E=15 J):
 - Measured longitudinal profile ($T_0 = 40$ fs)
 - Measured far field mode ($w_0 = 53 \mu\text{m}$)
- **Plasma:** parabolic plasma channel (length 9 cm, $n_0 \sim 6 \times 10^{17} \text{ cm}^{-3}$)

	Exp.	Sim.
Energy	4.25 GeV	4.5 GeV
$\Delta E/E$	5%	3.2%
Charge	~ 20 pC	23 pC
Divergence	0.3 mrad	0.6 mrad

W.P. Leemans et al., PRL 2014, in print

Bringing in the Trojan Horse



The Hybrid Scheme (Trojan Horse)

PRL 108, 035001 (2012)

PHYSICAL REVIEW LETTERS

week ending
20 JANUARY 2012

Ultracold Electron Bunch Generation via Plasma Photocathode Emission and Acceleration in a Beam-Driven Plasma Blowout

B. Hidding,^{1,2} G. Pretzler,² J.B. Rosenzweig,¹ T. Königstein,² D. Schiller,¹ and D.L. Bruhwiler³

¹*Department of Physics and Astronomy, University of California Los Angeles, Los Angeles, California 90095, USA*

²*Institut für Laser- und Plasmaphysik, Heinrich-Heine-Universität Düsseldorf, 40225 Düsseldorf, Germany*

³*Tech-X Corporation, Boulder, Colorado 80303, USA*

(Received 30 March 2011; published 17 January 2012)

Beam-driven plasma wakefield acceleration using low-ionization-threshold gas such as Li

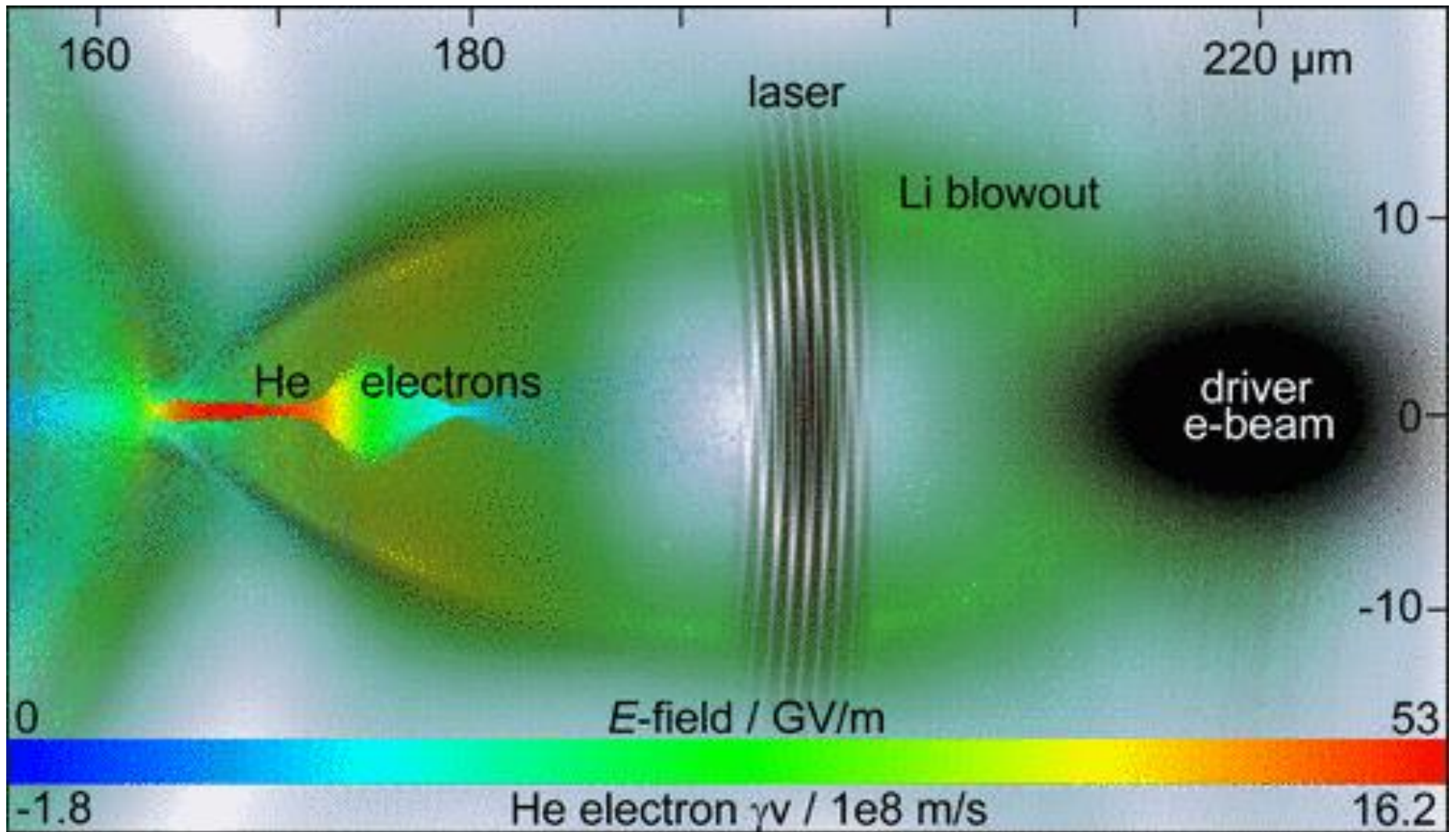


Laser-controlled electron injection via ionization of high-ionization-threshold gas such as He.

He electrons are released with low transverse momentum in the focus of the copropagating, non-relativistic intensity laser pulse directly inside the accelerating or focusing phase of the Li blowout → **generation of sub- μm -size, ultralow-emittance, highly tunable electron bunches.**



Sketch (Hidding et al, 2012)



Other approaches being studied, e.g. injection on the plasma density ramp

1. Accelerators – From Conventional Techniques to Plasmas
2. The Linear Regime
3. The Non-Linear Regime
4. **Tolerances**
5. Outlook for Europe



Accelerator Builder's Challenge *(simplified to typical values)*

- > Match into/out of plasma with **beam size $\approx 1 \mu\text{m}$** (about 1 mm beta function). Adiabatic matching (Whittum, 1989).
- > Control **offsets** between the wakefield driver (laser or beam) and the accelerated electron bunch at **$1 \mu\text{m}$ level**.
- > Use **short bunches (few fs)** to minimize energy spread.
- > Achieve **synchronization stability of few fs** from injected electron bunch to wakefield (energy stability and spread).
- > Control the **charge and beam loading** to compensate energy spread (idea Simon van der Meer).
- > Develop and demonstrate **user readiness of a 5 GeV plasma accelerated beam**.



Relax conditions...

- > As as low as possible plasma densities to start in most simple conditions. Larger matched beam size, relaxed tolerances, ...
- > The success will be all in accuracy, tolerances, precision! We mastered this in conventional accelerators.
- > Do the same for plasma accelerators!



1. Accelerators – From Conventional Techniques to Plasmas
2. The Linear Regime
3. The Non-Linear Regime
4. Tolerances
5. **Outlook for Europe**



Intensive work in Europe...



- 258 registered participants + about 50 accompanying persons.
- 45 sponsored students.
- Participants from 23 countries in 4 continents (11 EU member states).
- 16 % female participation.



SCAPA

LC Lund Laser Centre

STFC
ASTeC

LAOLA

Laboratory for Laser- and beam-driven plasma Acceleration

ILPP

JuSPARCO

ELBE

HELMHOLTZ
ZENTRUM DRESDEN
ROSSENDORF

STFC
Central Laser Facility

PHELIX

eli | beamlines

Cilex

Centre Interdisciplinaire Lumière Extrême



CALA

Laboratoire d'optique appliquée
UMR 7077 - Palaiseau / France

AIVAKE

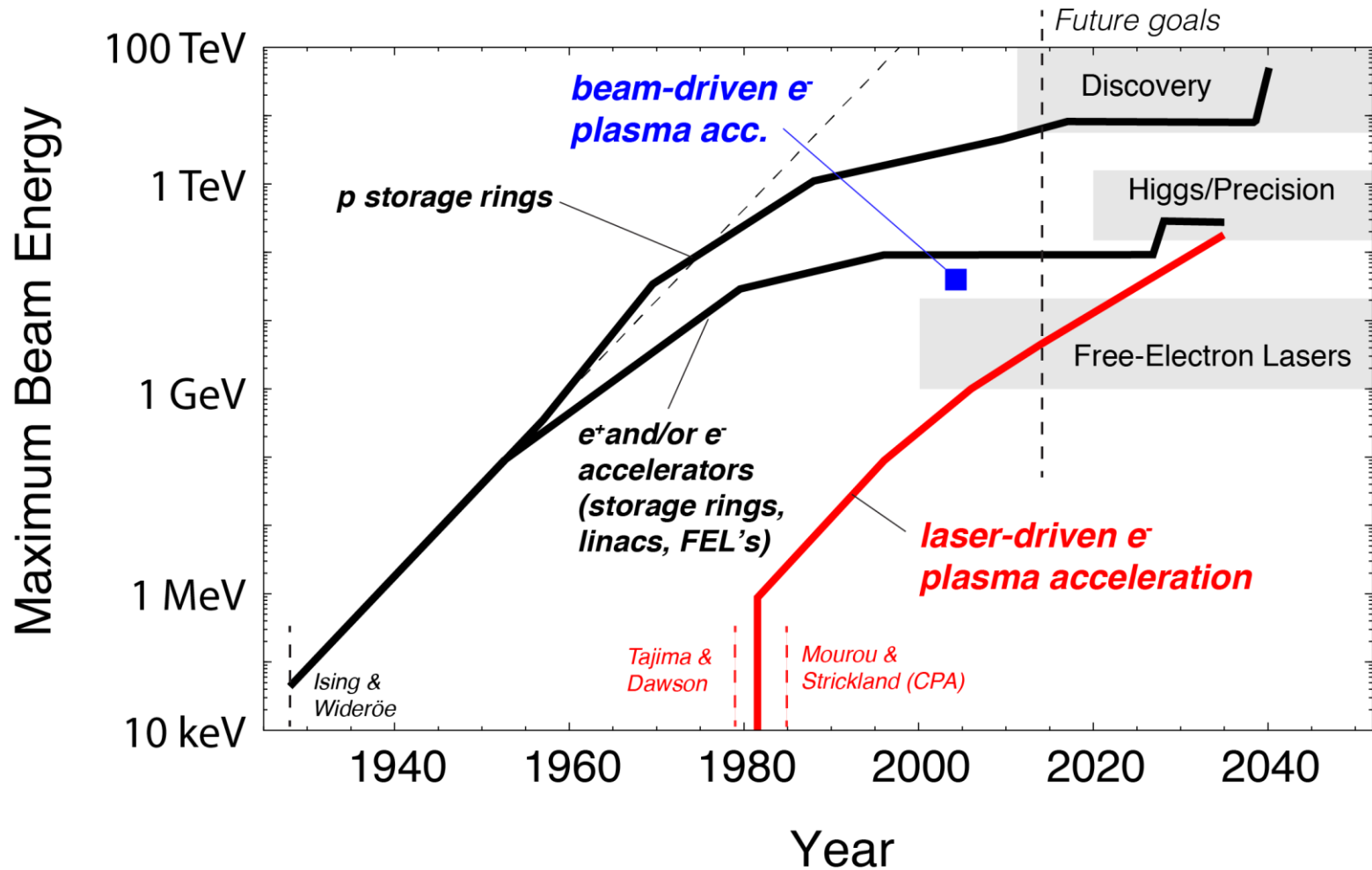
Consiglio Nazionale delle Ricerche

SPARC

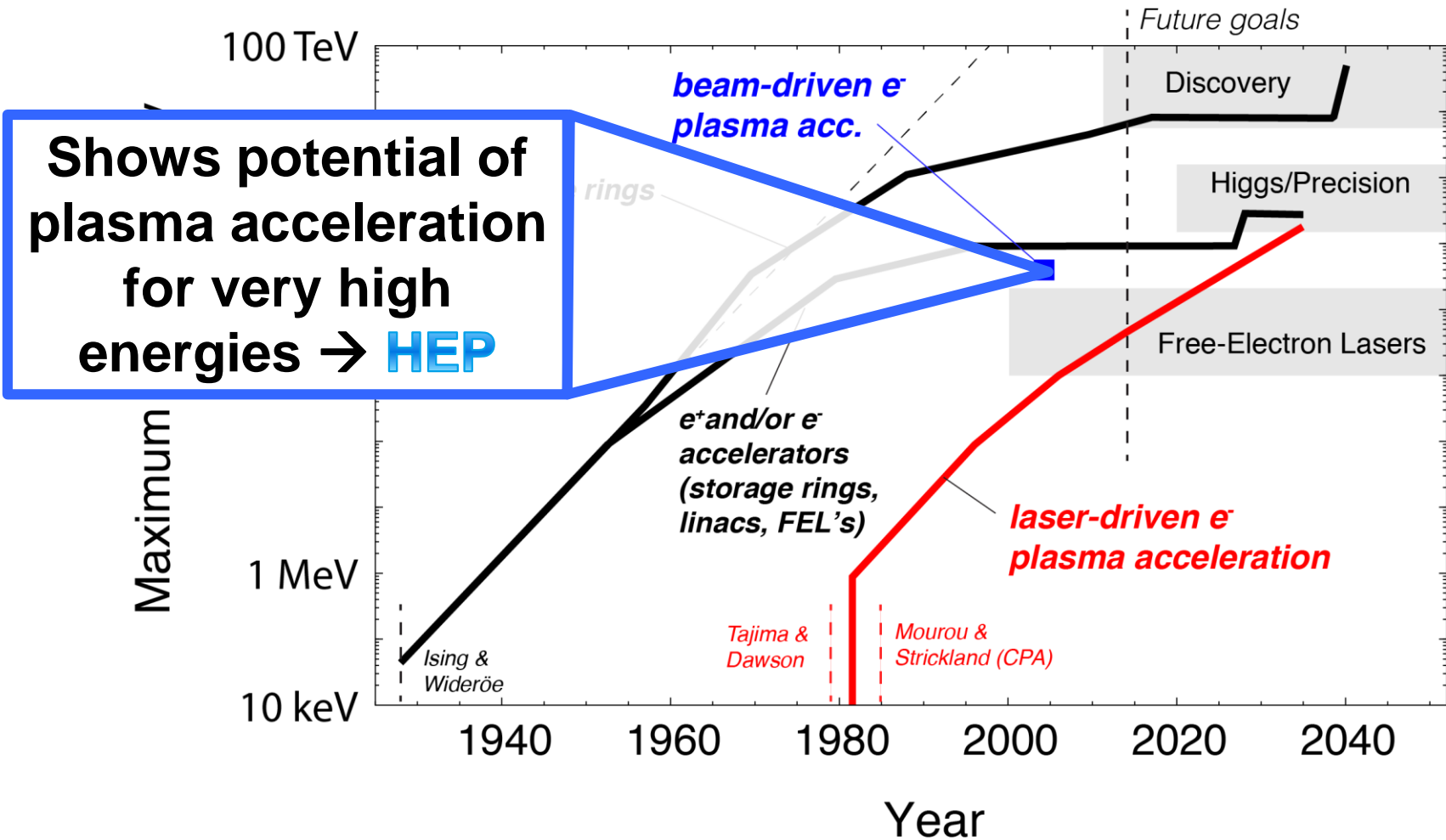
INFN
Sezione di Roma



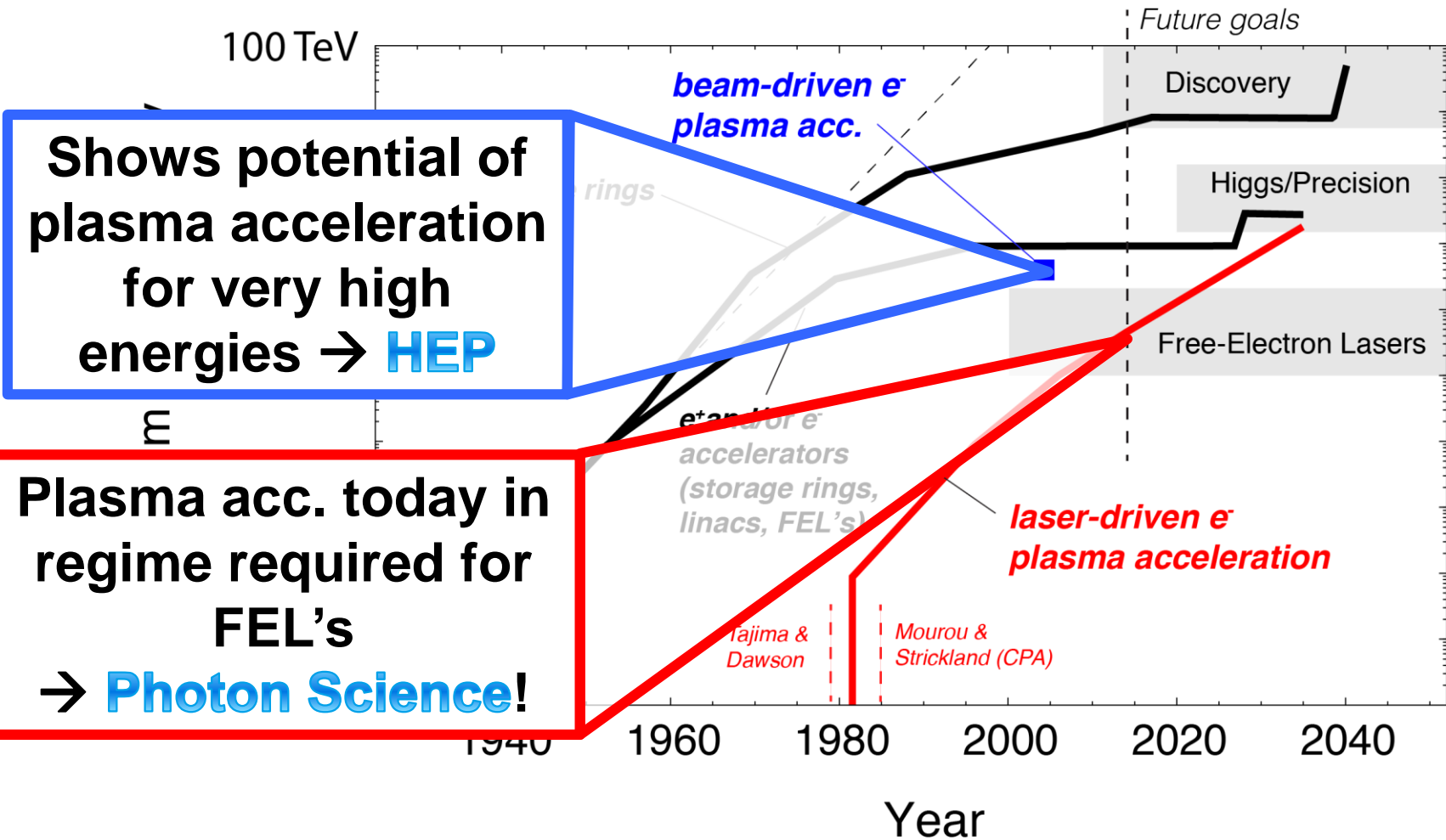
Livingston and Accelerators at the Energy Frontier



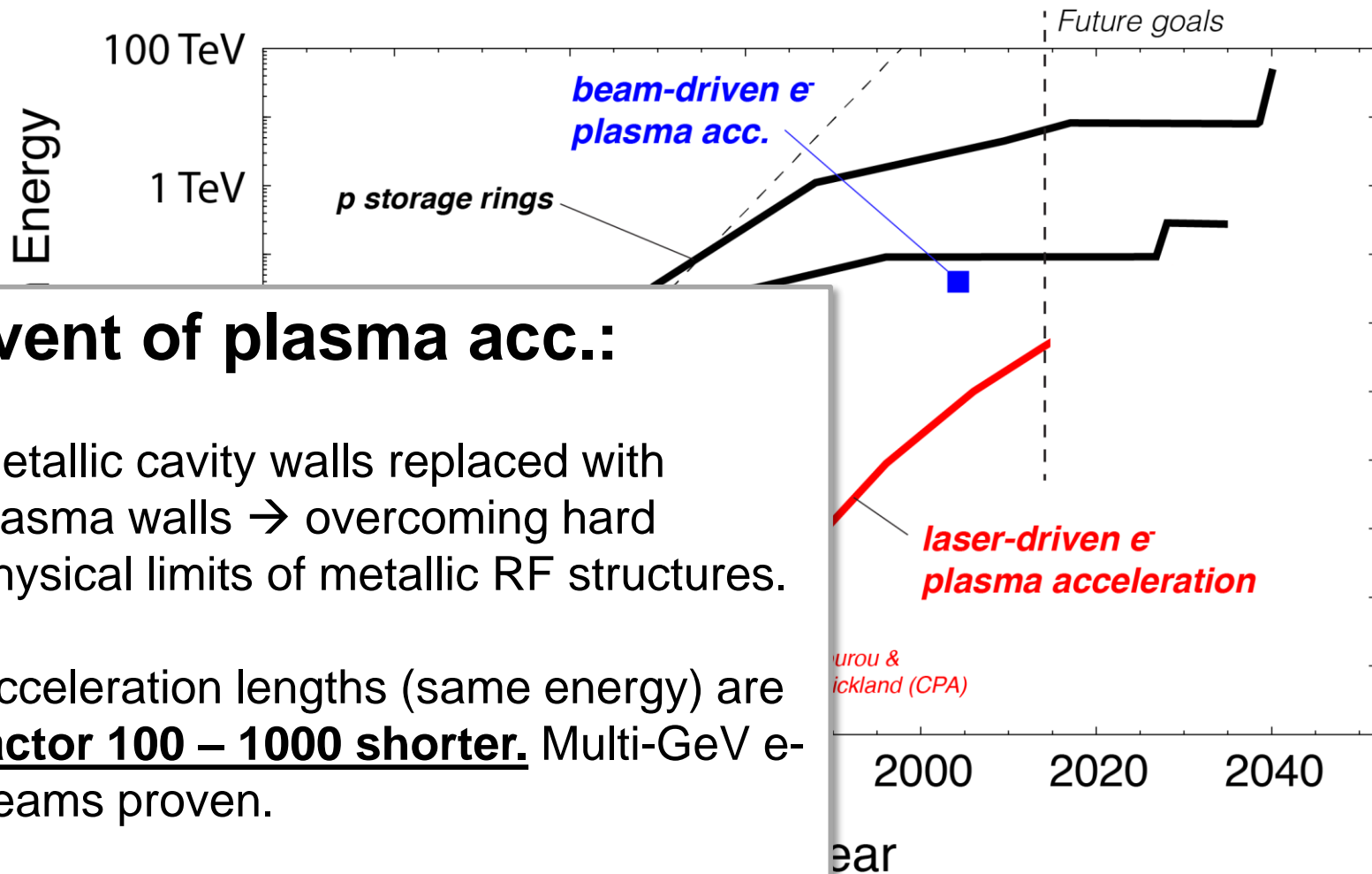
Livingston and Accelerators at the Energy Frontier



Livingston and Accelerators at the Energy Frontier



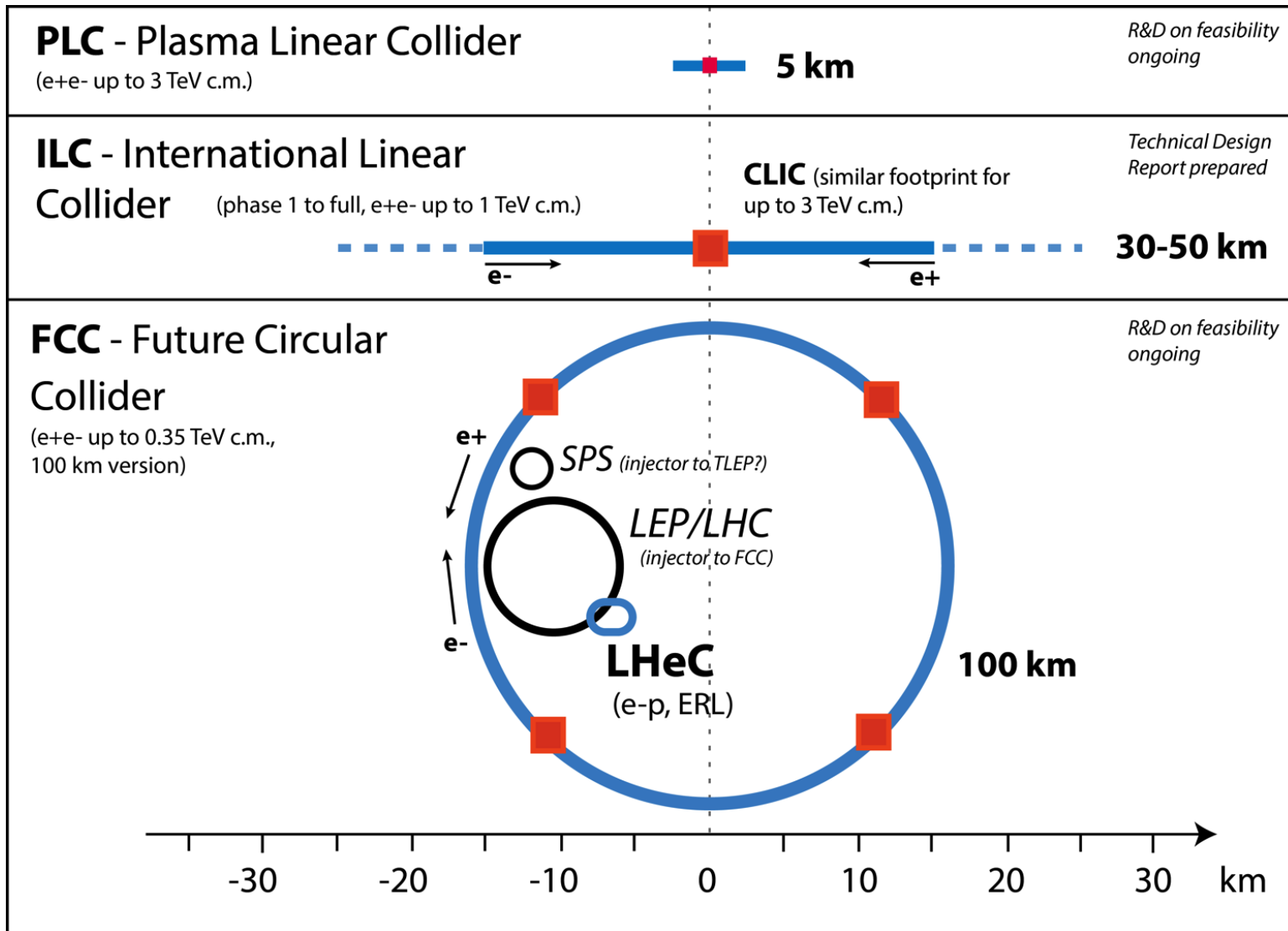
Livingston and Accelerators at the Energy Frontier



Advent of plasma acc.:

1. Metallic cavity walls replaced with plasma walls → overcoming hard physical limits of metallic RF structures.
2. Acceleration lengths (same energy) are **factor 100 – 1000 shorter.** Multi-GeV e-beams proven.
3. Still short-comings but **no fundamental limit.**

Small is Beautiful!? Is it?





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DISSEMINATION

EVENTS

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<http://eupraxia-project.eu>

NOVEL FUNDAMENTAL RESEARCH COMPACT EUROPEAN PLASMA ACCELERATOR WITH SUPERIOR BEAM QUALITY

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OUR TECHNOLOGY

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PARTICIPANTS

A consortium of 16 laboratories and universities from 5 EU member states has formed to produce a conceptual design report.

[LEARN MORE](#)



WORK PACKAGES

The project is structured into 14 work packages of which 8 are included into the EU design study.

[LEARN MORE](#)



MANAGEMENT

The management bodies will organise, lead and control the project's activities and make sure that objectives are met

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OPENING NEW HORIZONS

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COMPACT EUROPEAN PLASMA ACCELERATOR WITH SUPERIOR BEAM QUALITY

Design report for a 5 GeV facility by end of 2019, including science case for pilot users, cost and site study. Second design study (“plan B”) after FCC/EuroCirCol.

Kick-off meeting at DESY on Nov 26th – 27th



Wideröe 1992 at age 90

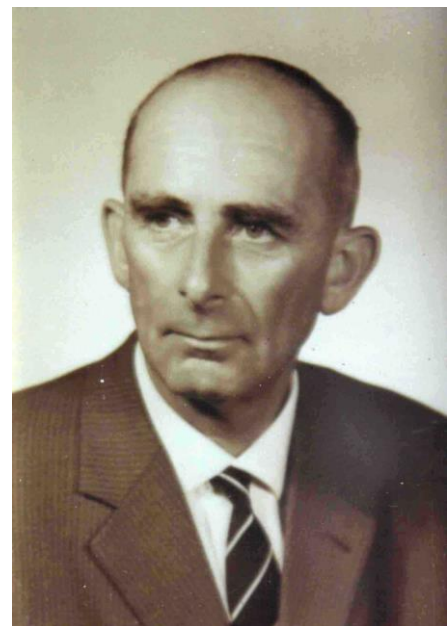


After all, **plans can only be made for those accelerators which can realistically be built with the means available, and obviously, these means are limited.**

Ideas are not subject to any such considerations. The **limitations are set only by the intellect of human beings themselves.**

The **theoretical possibilities** with regard to accelerating particles by electromagnetic means (i.e. within the scope of the Maxwell equations which have been known since the 19th century), **are nowhere near being exhausted**, and technology surprises us almost daily with innovations which in turn allow us to broach new trains of thought.

...there are yet **more fundamental breakthroughs** to be made. They could allow us to advance to **energies unimaginable today.**



Thank you for your attention...

