

FCC-hh Detectors and Experiments

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FCC Academic Training Lectures
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Material from Discussions on FCC-hh Detector Meetings:

FCC-hh Detector Magnets

<https://indico.cern.ch/category/6244/>

FCC-hh Detectors

<https://indico.cern.ch/category/6069/>

e-mail-list: fcc-experiments-hadron@cern.ch

FCC-hh machine detector interface

<https://indico.cern.ch/category/5901/>

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FCC-hh Detectors and Experiments

Abstract

This lecture will outline detector concepts for a 100 TeV proton collider that is being investigated in the context of the Future Circular Collider (FCC) study. An ultimate instantaneous luminosity of about $3 \times 10^{35} \text{ cm}^{-2} \text{ s}^{-1}$ and a total integrated luminosity in the range of 20 ab^{-1} pose significant challenges on detector technology beyond the LHC Phase II upgrades. The baseline specifications for detectors at such a machine are presented and first implementation ideas are being shown. Key areas of necessary detector R&D are discussed as well.

Baseline Parameters for the FCC-hh Machine

The present working hypothesis is:

- peak luminosity baseline: $5 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$
- peak luminosity ultimate: $\leq 30 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$

- integrated luminosity baseline $\sim 250 \text{ fb}^{-1}$ (average per year)
- integrated luminosity ultimate $\sim 1000 \text{ fb}^{-1}$ (average per year)

An operation scenario with:

- 10 years baseline, leading to 2.5 ab^{-1}
- 15 years ultimate, leading to 15 ab^{-1}

would result in a total of $O(20) \text{ ab}^{-1}$ over 25 years of operation.

Parameters Assumed for the Detector Design

$$L_{\text{peak}} [5 \times 10^{34}, 30 \times 10^{34}] \text{ cm}^{-2}\text{s}^{-1}$$

→ Average N_{pileup} [170, 1020] at 25ns

→ Average N_{pileup} [34, 204] at 5ns

$$L_{\text{int}} [3, 30] \text{ ab}^{-1}$$

These upper limits of L_{peak} and L_{int} should be read as Phase II goals that we use for detector studies and not as numbers promised by the machine!

The 5ns vs. 25ns bunch crossing time will stay an open parameter for some time.



FCC Motivation: Pushing the Energy Frontier

The name of the game of a hadron collider is **energy reach**

$$E \propto B_{\text{dipole}} \times R_{\text{bending}}$$

Cf. LHC: factor ~ 4 in radius, factor ~ 2 in field \rightarrow **O(10) in E_{cms}**

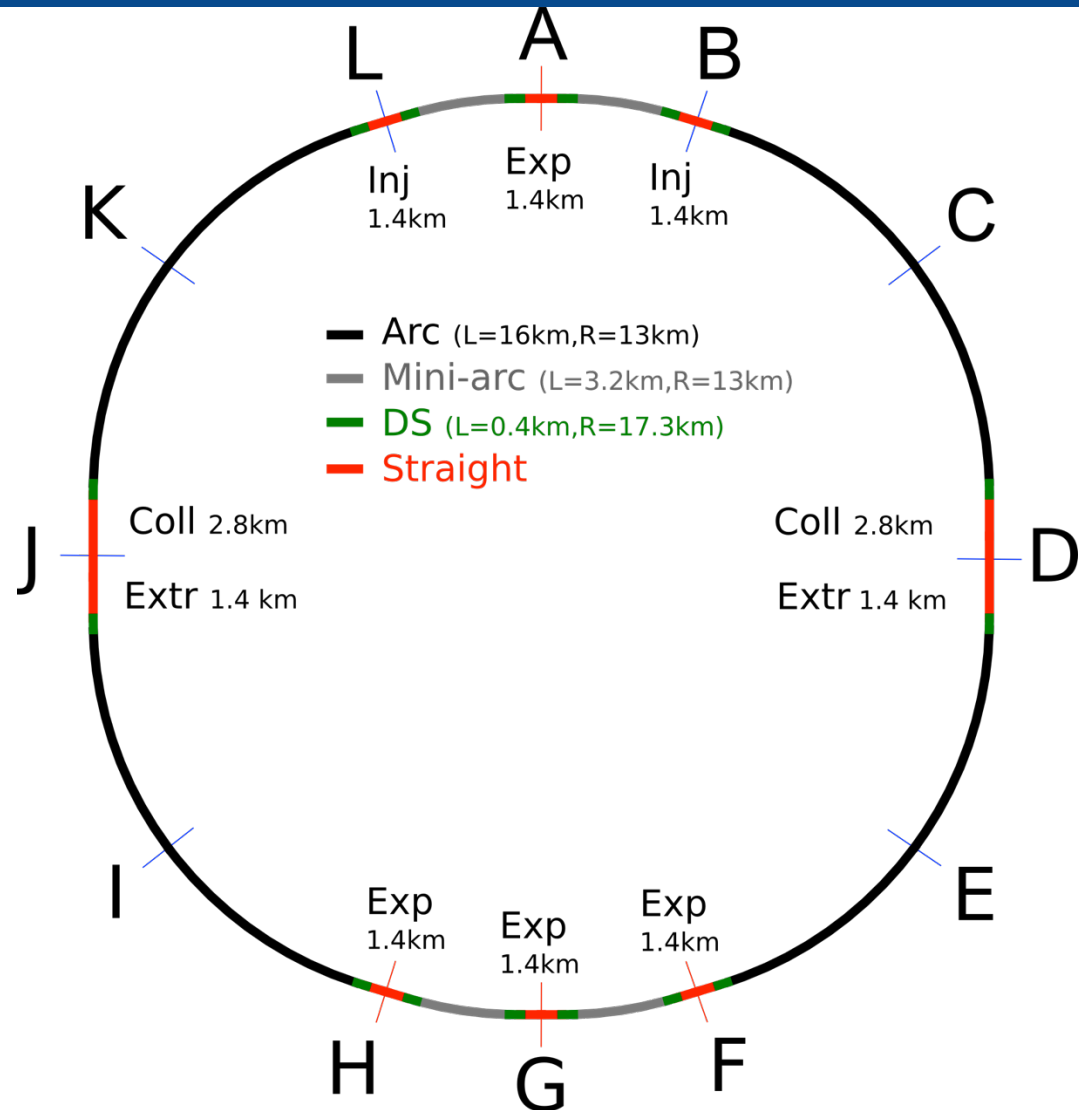
100 km layout for FCC-hh
(different sizes under investigation)

⇒ **Two high-luminosity experiments (A and G)**

⇒ **Two other experiments (F and H) grouped with main experiment in G**

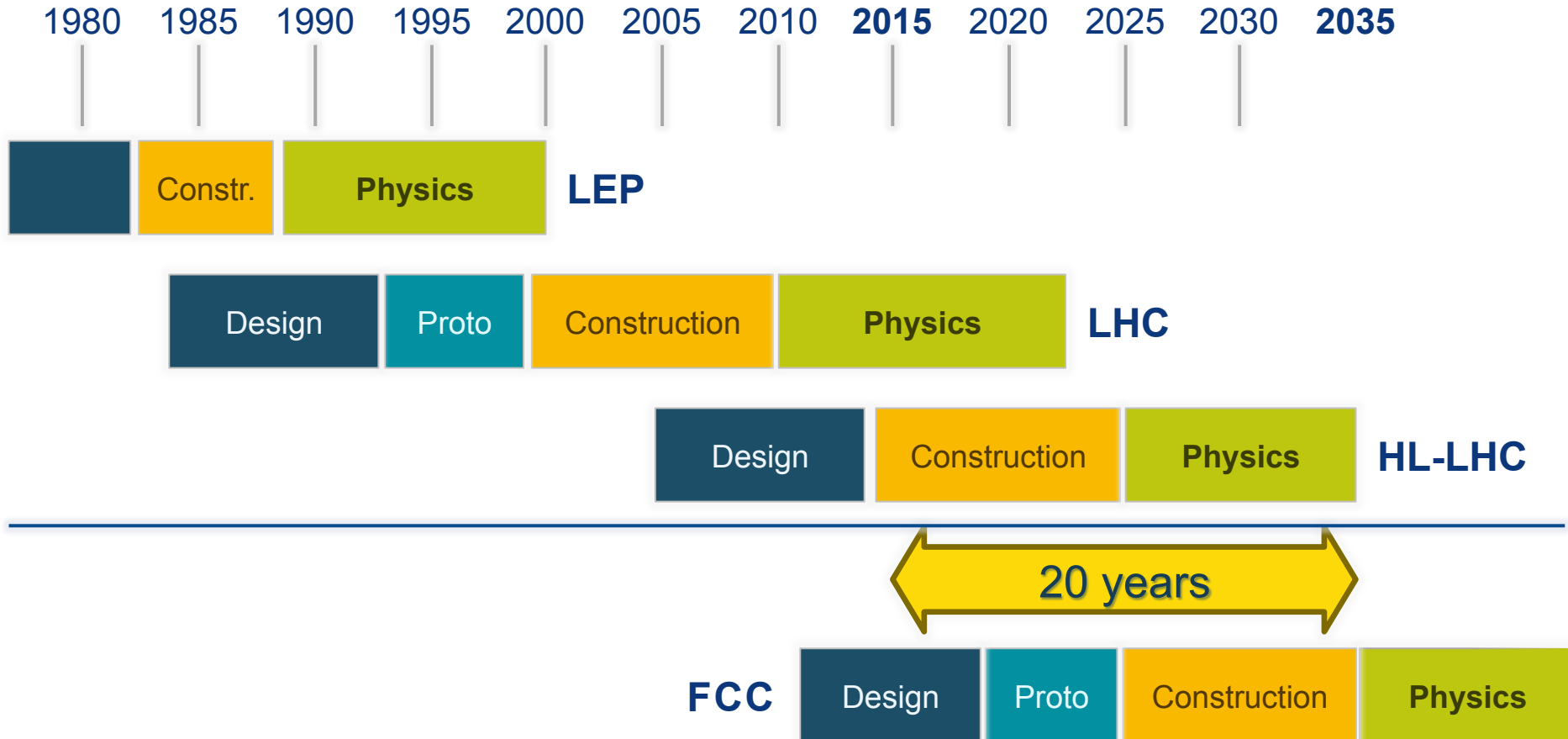
⇒ **Two collimation lines**

⇒ **Two injection and two extraction lines**



FCC-hh Machine Parameters

Parameter	FCC-hh		SPPC	LHC	HL LHC
collision energy cms [TeV]	100		71.2	14	
dipole field [T]	16		20	8.3	
# IP	2 main & 2		2	2 main & 2	
bunch intensity [10^{11}]	1	1 (0.2)	2	1.1	2.2
bunch spacing [ns]	25	25 (5)	25	25	25
luminosity/lp [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$]	5	~25	12	1	5
events/bunch crossing	170	~850 (170)	400	27	135
stored energy/beam [GJ]	8.4		6.6	0.36	0.7
synchrotron radiation [W/m/aperture]	30		58	0.2	0.35



FCC Conceptual Design Report by end 2018 for the European strategy update

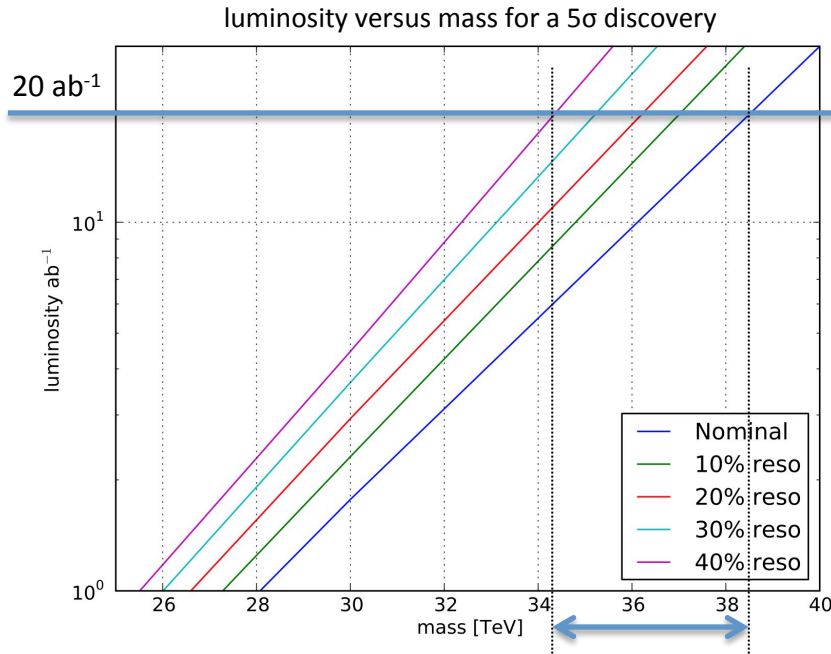
Summary of Requirements from Physics for the FCC Detectors (very preliminary)

**...please see yesterday's lectures for many more details on
possible measurements**

Physics at the $L\sigma$ Limit

Exploration potential through higher energy, increased statistics, increased precision

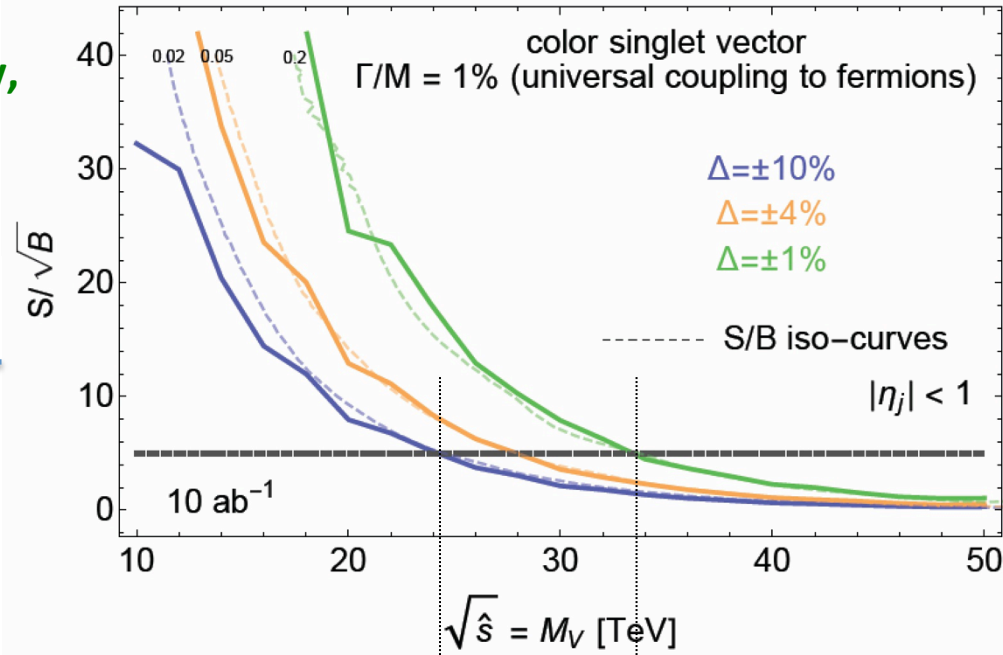
Example: Z' $_{SSM}$ discovery



$$\frac{\Delta p}{p} \propto \frac{p}{BL^2}$$

Muon momentum resolution:

- $O(15\%)$ at 10TeV.
- Compare to 10% at 1TeV spec. at LHC



$$\frac{\Delta E}{E} \propto \frac{1}{\sqrt{E}} \oplus k$$

Di-jet resonances: Extend discovery potential by 10TeV between mass resolutions of $\Delta = \pm 10\%$ to $\Delta = \pm 1\%$

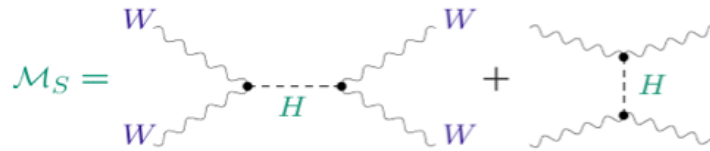
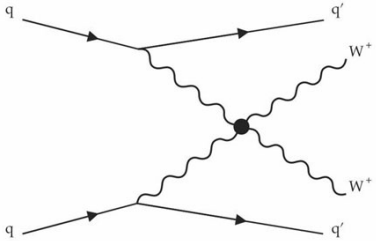
2% jet resolution a reasonable choice ($\Delta = \pm 4\%$)

- Constant term dominates, $\approx 2\%$ goal
- \rightarrow full shower containment is mandatory !
- \rightarrow HCAL depth of $12 \lambda_{\text{int}}$!

WW Scattering by VBF Mechanism

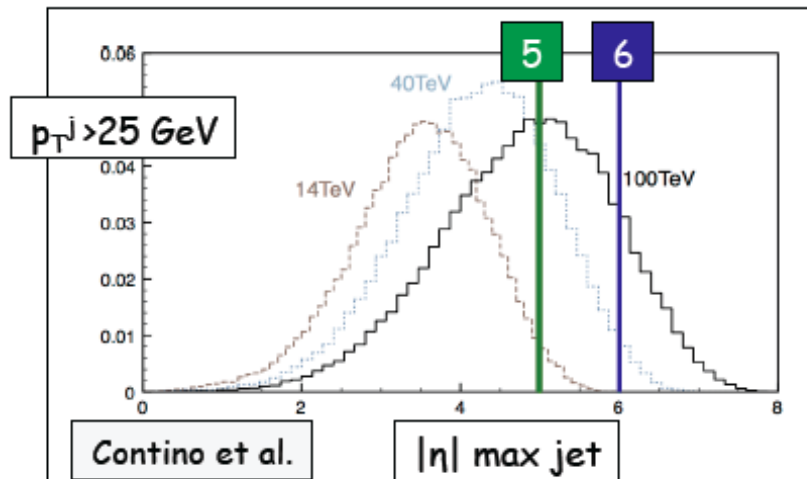
WW→WW scattering violates unitarity at high energies

- A scalar, such as the Higgs boson, fixes this (partially)
- Probing characteristics of VV scattering is an important test of the nature of electroweak symmetry breaking
- New Physics would modify interferences between diagrams → modified $V p_T$ and di-boson mass. Also: Are there high mass resonances WW, ZZ, HH, ...



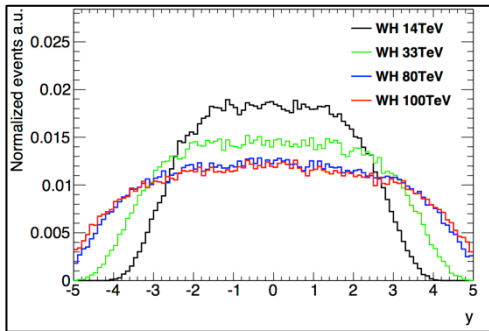
VBF jets also important for tagging of Higgs produced through VBF, like $H \rightarrow b\bar{b}$, $H \rightarrow \tau\tau$ etc.

VBF jets between $\eta \sim 2$ and $\eta \sim 6$ need to be well measured and separated from pile-up



Higgs Measurements

H → 4l acceptance vs η coverage (p_T cuts applied)



	14 TeV		100 TeV	
	2.5	4	2.5	4
ggF	0.74	0.99	0.56	0.88
WH	0.66	0.97	0.45	0.77
ZH	0.69	0.98	0.48	0.80
ttH	0.84	1	0.56	0.90
VBF	0.75	0.98	0.55	0.87

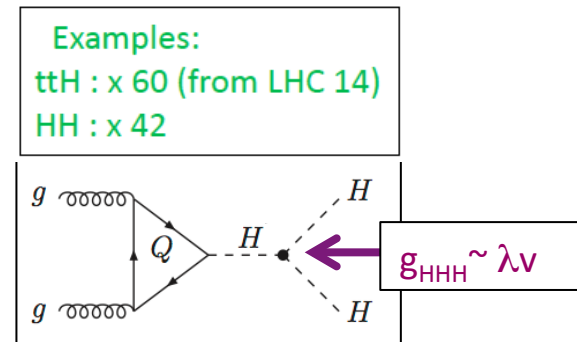
		η < 2.5	η < 4	η < 5
YY	100 TeV	0.74	0.95	0.99
	14 TeV	0.90	1	1

- 30-50% acceptance loss for H → 4l at 100 TeV wrt 14 TeV if **tracking and precision EM calorimetry** limited to |η| < 2.5 (as ATLAS and CMS)
- can be recovered by extending to |η| ~ 4

“Heavy” final states require high v_s, e.g.:

HH production (including measurements of self-couplings λ)

ttH (note: ttH → ttμμ, ttZZ “rare” and particularly clean)



FCC

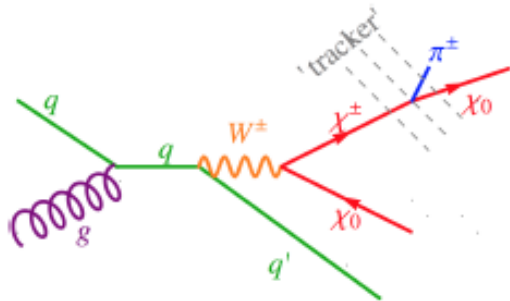
	HL-LHC	ILC500	ILC500-up	ILC1000	ILC1000-up	CLIC1400	CLIC3000	HE-LHC	VLHC
\sqrt{s} (GeV)	14000	500	500	500/1000	500/1000	1400	3000	33,000	100,000
$\int \mathcal{L} dt$ (fb ⁻¹)	3000	500	1600 [‡]	500/1000	1600/2500 [‡]	1500	+2000	3000	3000
λ		83%	46%	21%	13%	21%	10%	20%	8%

More Exotic

Disappearing Tracks - Introduction

$$M_{\chi^\pm} - M_{\chi_0} = 165 \text{ MeV} > m_\pi \Rightarrow \text{lifetime } \tau \simeq 6 \text{ cm} \simeq 0.2 \text{ ns}$$

Almost all χ^\pm s decay to χ_0 + soft pions before reaching detectors



Feng Strassler 1994

Feng Moroi Randall Strassler Su 1999

...

Low Wang 1404.0682

Filippo Sala

→ Missing E_T Measurement important! E_T^{miss} distributions with smallest tails possible to keep sensitive to very rare processes → high eta coverage!

Physics at a 100 TeV Hadron Collider

Exploration + Higgs as a tool for discovery

Numerous physics opportunities with a large number of possible measurements.

How to specify detectors for such a machine ?

ATLAS and CMS are general purpose detectors that were benchmarked with the 'hypothetical' Higgs in different mass regions with precision tracking and calorimetry up to $\eta=2.5$.

The Higgs is also key benchmark for the FCC detectors, with highly forward boosted features ($E_{\text{cm}}=100\text{TeV}$, Higgs mass = 125GeV)

FCC detectors must be 'general general' purpose detectors with very large η acceptance and extreme granularity.

Approximate Overall Needs

Tracking: Momentum resolution $\approx 15\%$ at $p_T=10\text{TeV}$

Precision tracking (momentum spectroscopy) and ECAL up to $\eta=4$

ECAL fine granularity for track-cluster matching (or particle flow) to mitigate pile-up and recover Bremsstrahlungs losses

**Tracking and calorimetry for jets up to $\eta=6$.
12 λ_{int} calorimetry $\approx 2\%$ constant term.**

HCAL granularity of 0.05×0.05 or 0.025×0.025 to mitigate pileup and measure jet substructure and boosted objects.

B-tagging, timing for pileup rejection etc. ...

What Do Inelastic Collisions at 100TeV Look Like

Minimum Bias events scaling 14TeV \rightarrow 100TeV:

Inelastic cross-section changes from 80 \rightarrow 108mb.

Multiplicity changes from 5.4 \rightarrow 8 charged particles per rapidity unit.

Average p_T of charged particles changes from 0.6 \rightarrow 0.8 GeV/c.

Hard scatter events (events of interest) with p_T up to 7 times higher (100/14).

\rightarrow Transverse energy sum increases by about a factor of 2.

\rightarrow The Min. Bias events at FCC are quite similar to the Min. Bias events at LHC.

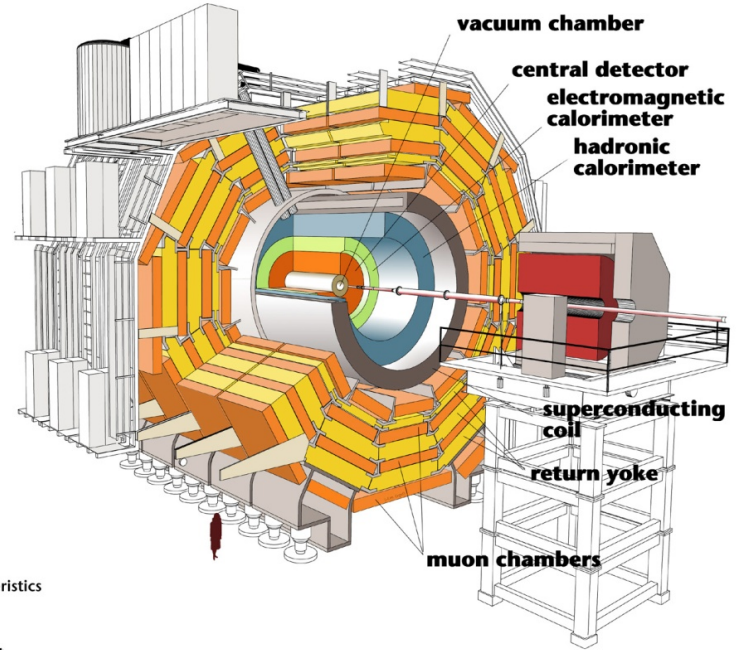
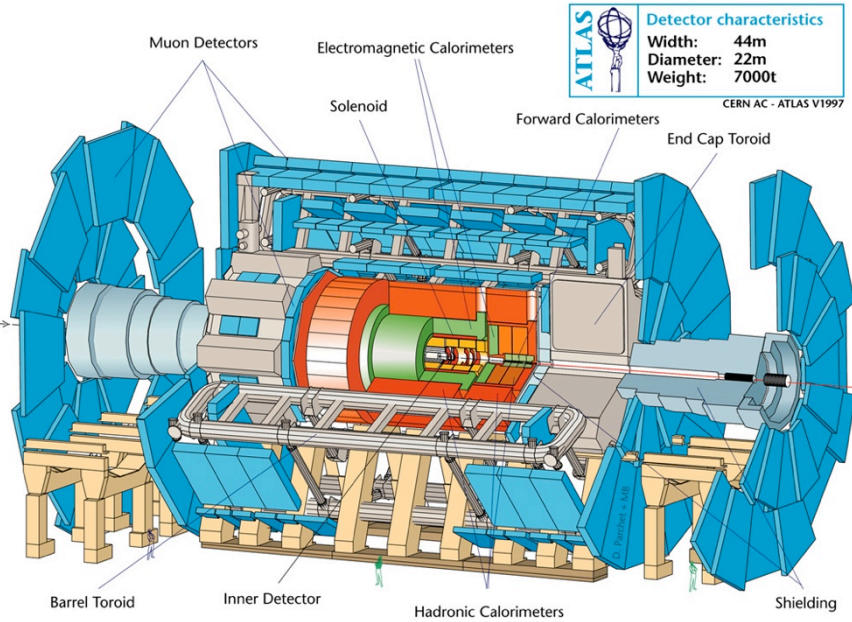
Key Point and Strategy

If the FCC hadron machine with 16T magnets, 5MW synchrotron radiation and a 100km tunnel can be realized, there is no doubt that a detector, that makes full use of the physics potential, can be built.

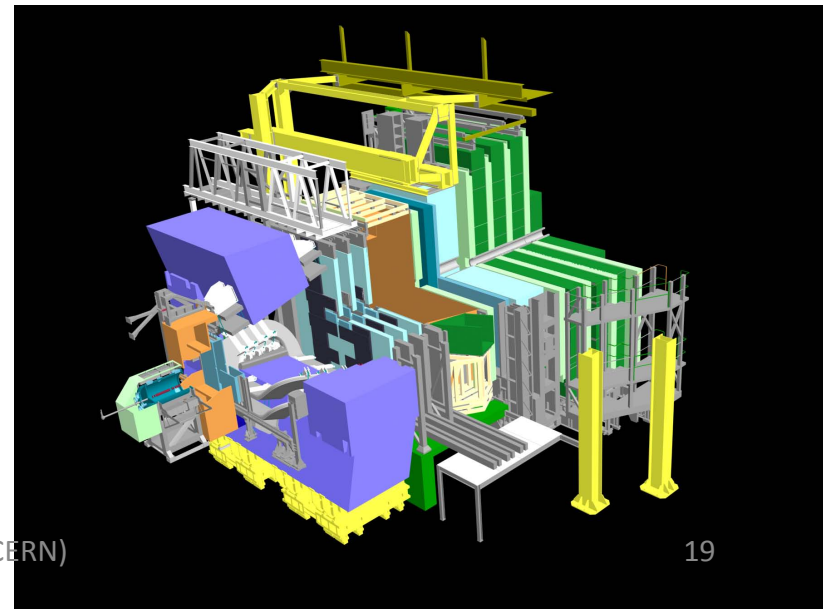
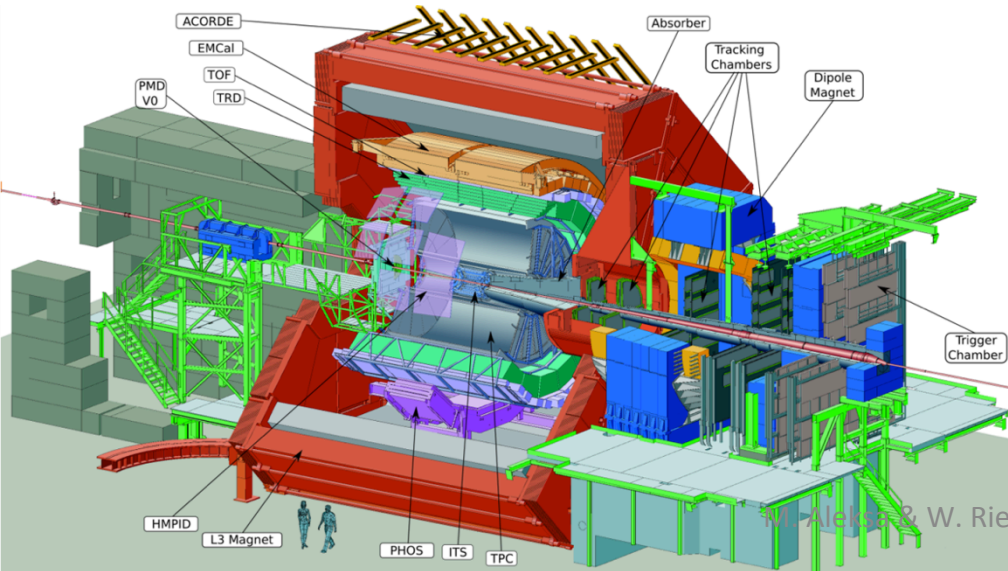
Much of detector technology is driven by silicon technology and computing power i.e. we can count on significant improvements.

Since the maximum energy and delivered luminosity are the key goals for the FCC-hh machine, **the detector efforts should put minimal constraints at the machine efforts.**

Guidance and Scaling from ATLAS, CMS, ALICE, LHCb

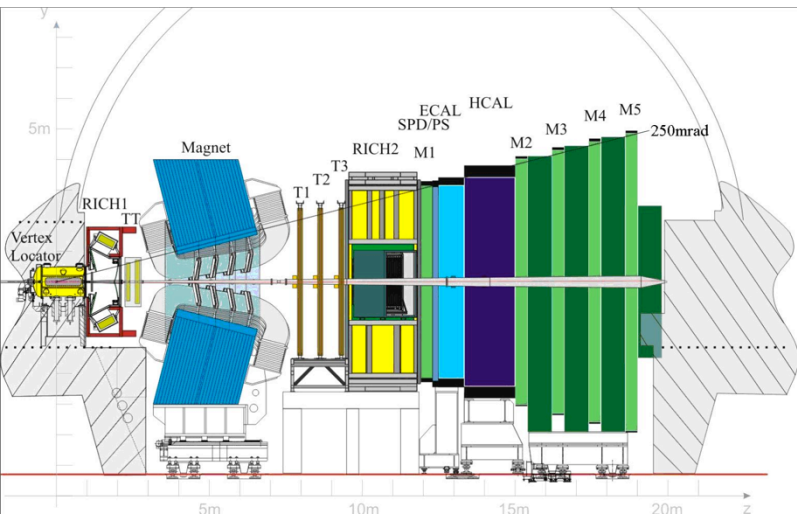


Detector characteristics
 Width: 22m
 Diameter: 15m
 Weight: 14500t

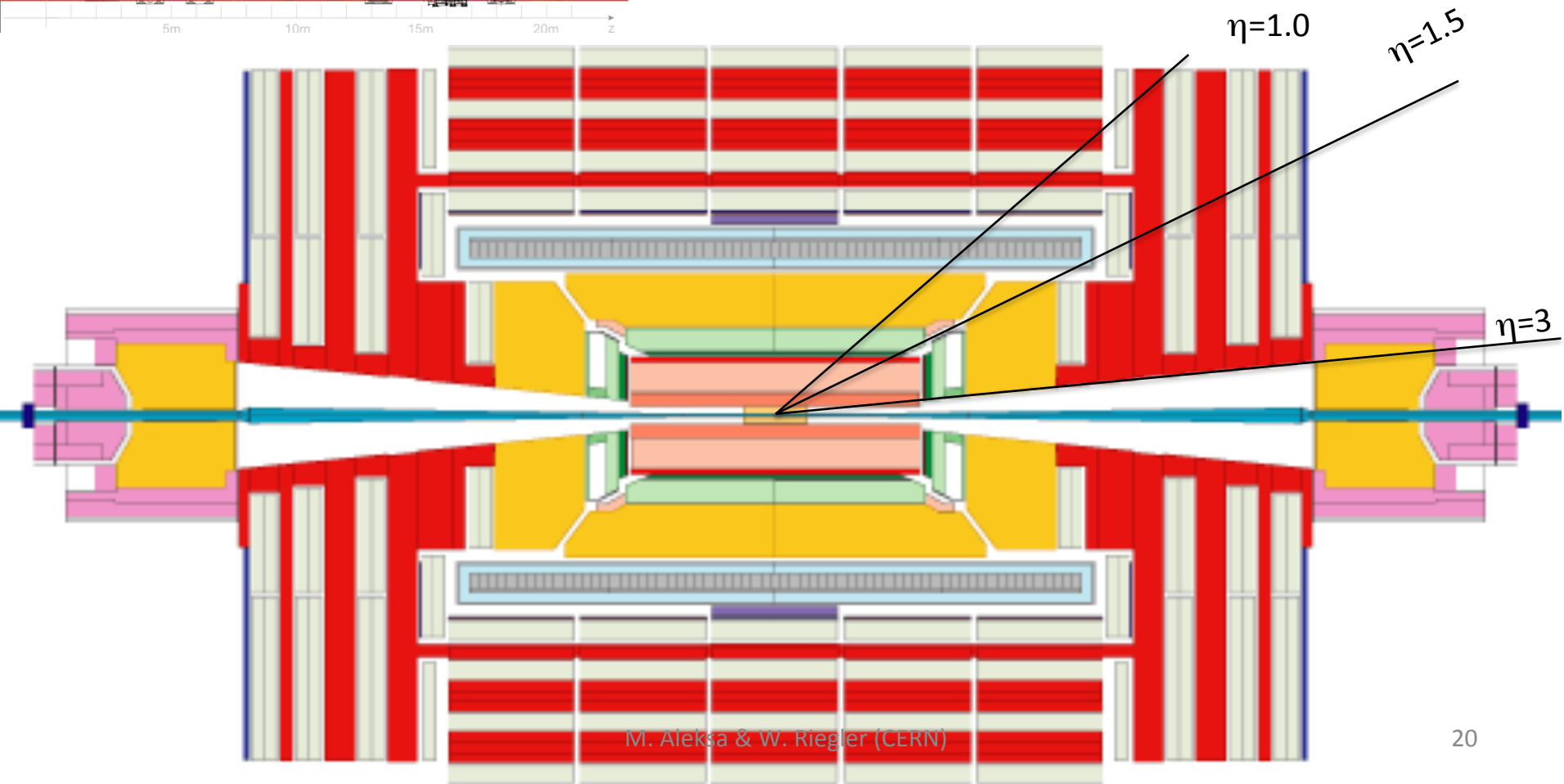


LHCb: Tracking, Calo $\eta = 2 - 5$

... all with impressive performance ...



ATLAS, CMS: tracking, calo $\eta -2.5, 2.5$



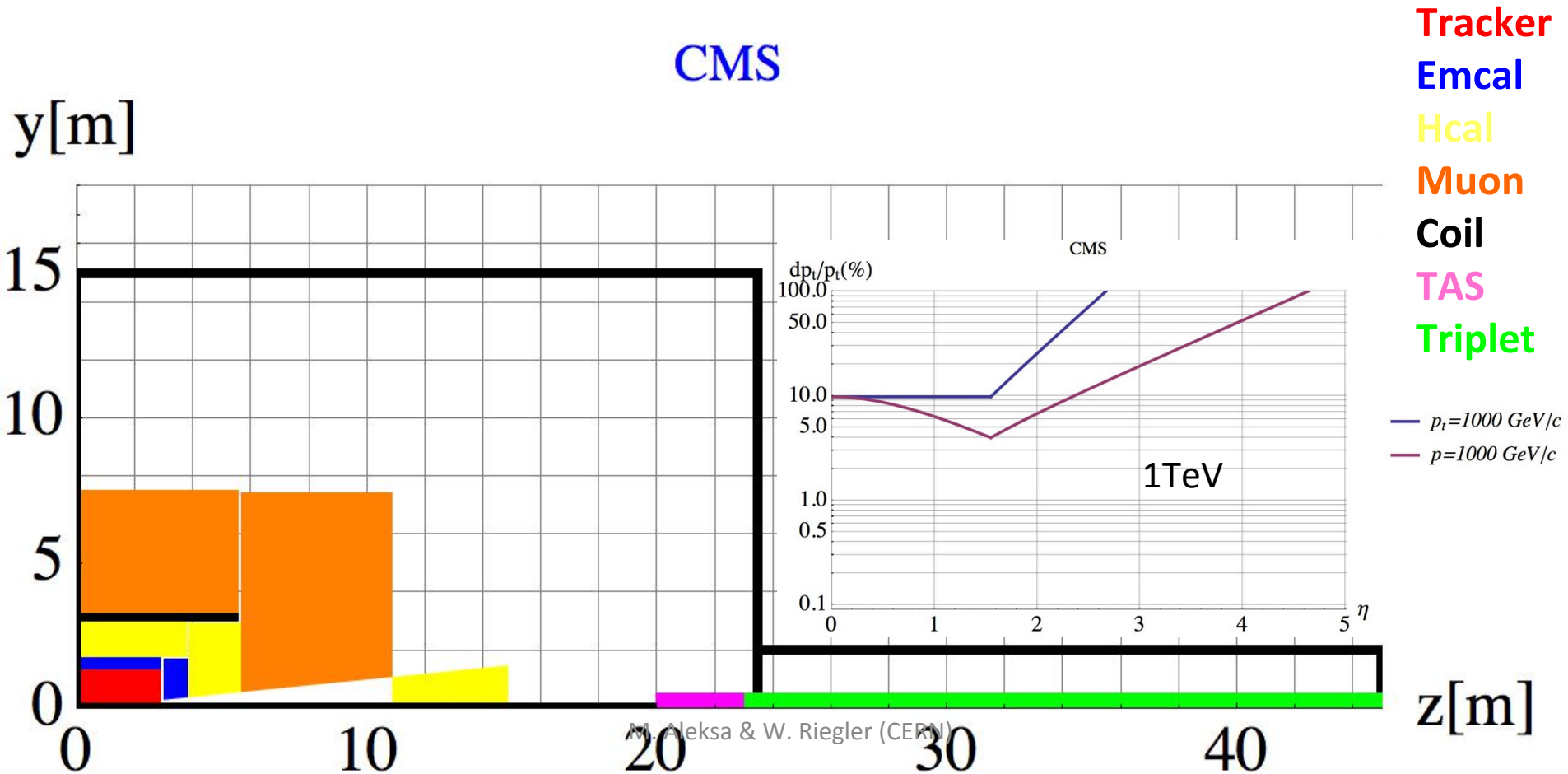
ATLAS

- Tracker $r=1\text{m}$, $B=2\text{T}$ **thin solenoid coil** in front of the calorimeters
- LArg ECAL, HCAL and $7.4 \lambda_{\text{int}}$ that returns the flux
- **Large air core toroid**, $B=0.5\text{T}$ 'standalone muon system'



CMS

- Tracker $r=1.2\text{m}$
- **Compact Crystal ECAL, 'short' HCAL** of and $5.82 \lambda_{\text{int}}$, cut at $\eta = 3$ to move FCAL away.
- **R=3m solenoid coil** with 3.8T field.
- **Iron Yoke to return Flux**, instrumented with muon chambers.
- CMS muons are relying on a properly working tracker.



How to Scale LHC Experiments to FCC ?

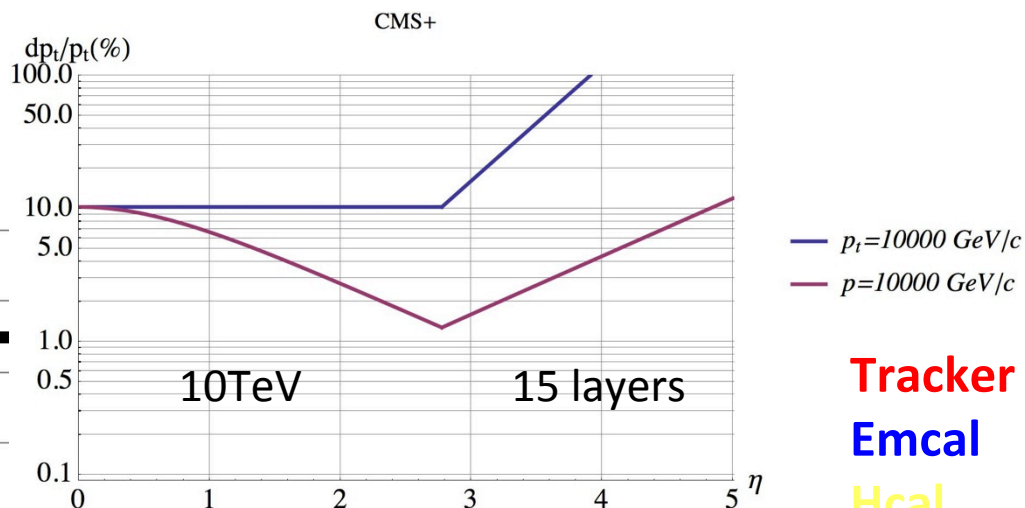
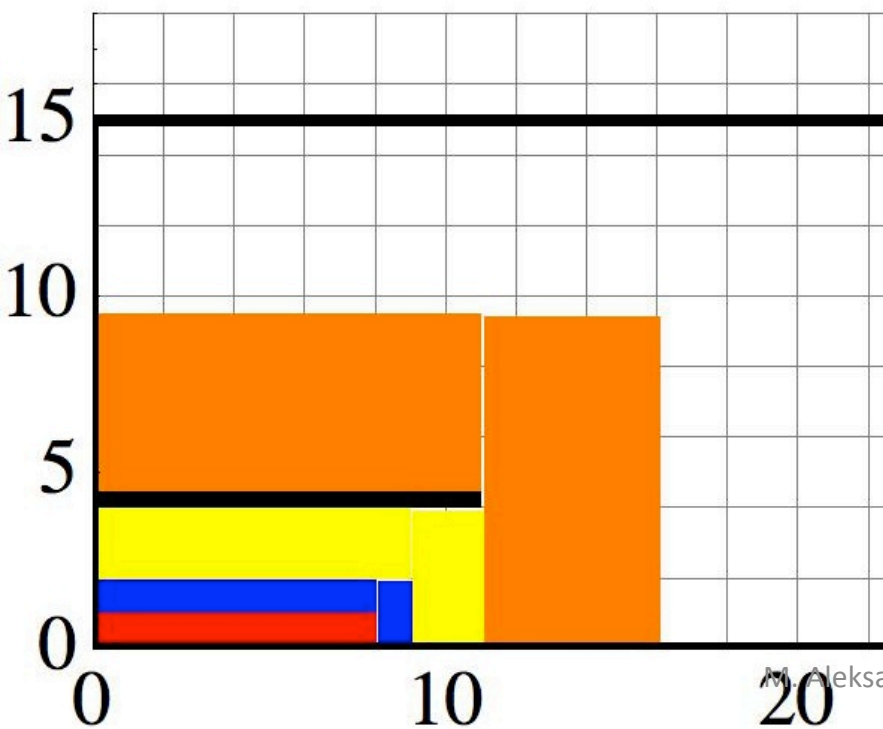
Let's assume a tracking resolution of **10-15%** for **10TeV** particles and a calo **constant term of $\approx 2\%$** which requires full shower containment and therefore **$12 \lambda_{\text{int}}$** of calo i.e. $\geq 3\text{m}$

- Coil with high B-field and low material budget in front of ECAL/HCAL seems very difficult, so scaling the **ATLAS approach is questionable**.
- Leaving the **tracker radius** similar to LHC values of **$r=1\text{m}$** , which is **extremely challenging**, with **$12\lambda_{\text{int}}$** calo a coil radius of at least 4m is needed (**$\rightarrow \text{CMS+}$**).
 - \rightarrow An iron yoke to return the flux for such a coil might still be affordable.
- With a more realistic approach for calorimetry and tracking we end up with coil radii of 6m, which requires an iron yoke that is probably unaffordable.
 - \rightarrow In this case one uses either **active shielding (twin solenoid)** or a yoke that only returns part of the flux (**partial shielding**) - stringent requirements on the equipment in the environment.

CMS Scaled Detector with Very Long Extreme Resol. Tracker

- Maximum coil producing 6T with affordable iron yoke ($r=4\text{m}$)
 - Tracker radius 1m, 6T \rightarrow resolution has to be improved by factor 6 with respect to CMS $\rightarrow 5\mu\text{m}$ layer resolution and less material (multiple scattering)
 - 8m long tracker gives large η acceptance.
 - 2.8m available for EMCAL+HCAL e.g. very compact W/Si particle flow calorimeters
 - Very high granularity forward calorimeters needed
 - Muon system a'la CMS
- \rightarrow **'extreme' technology challenge.**

y[m]



Tracker

Emcal

Hcal

Muon

Coil

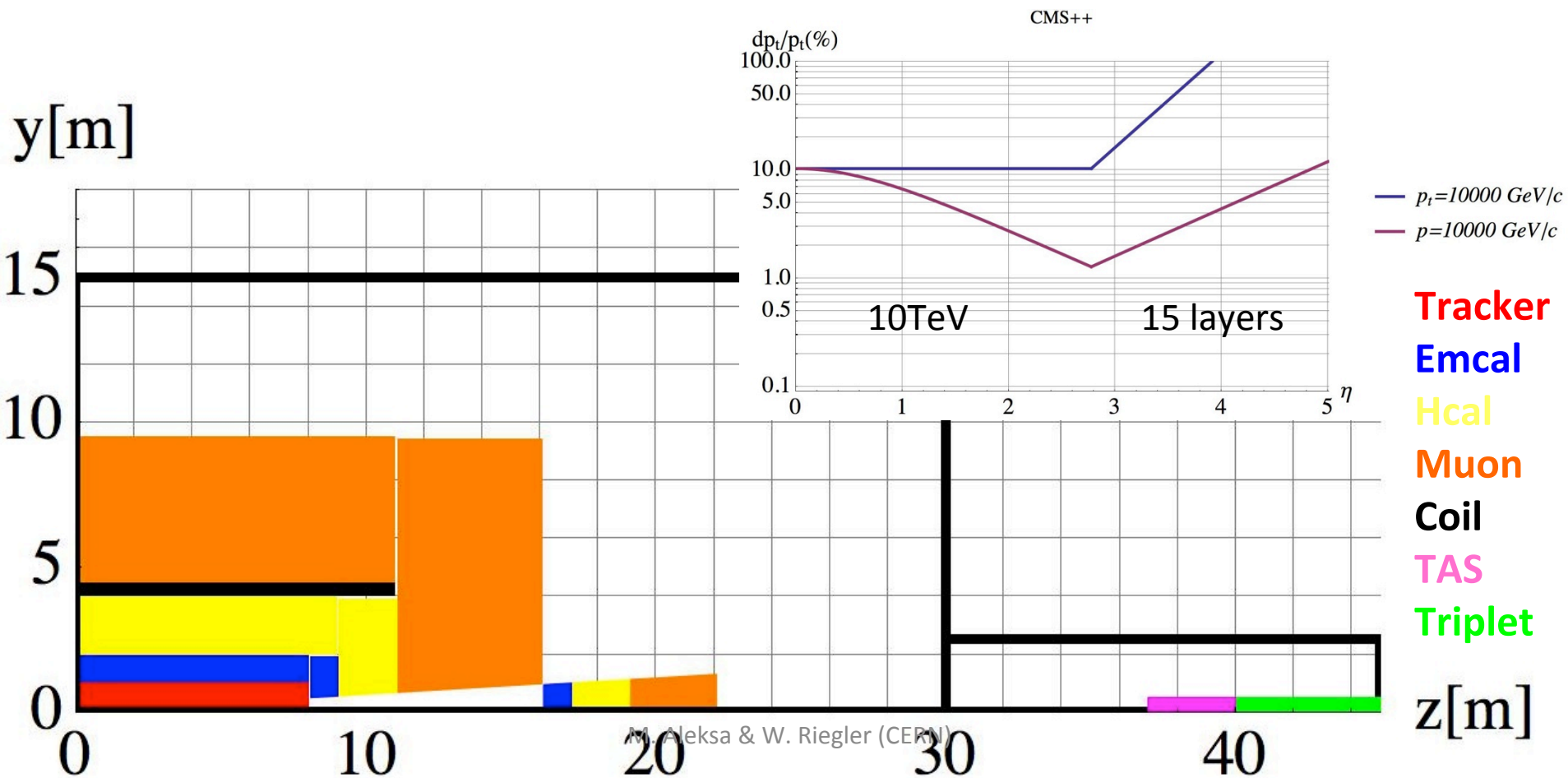
TAS

Triplet

z[m]

CMS Scaled Detector, Forward Calorimetry Moved Out

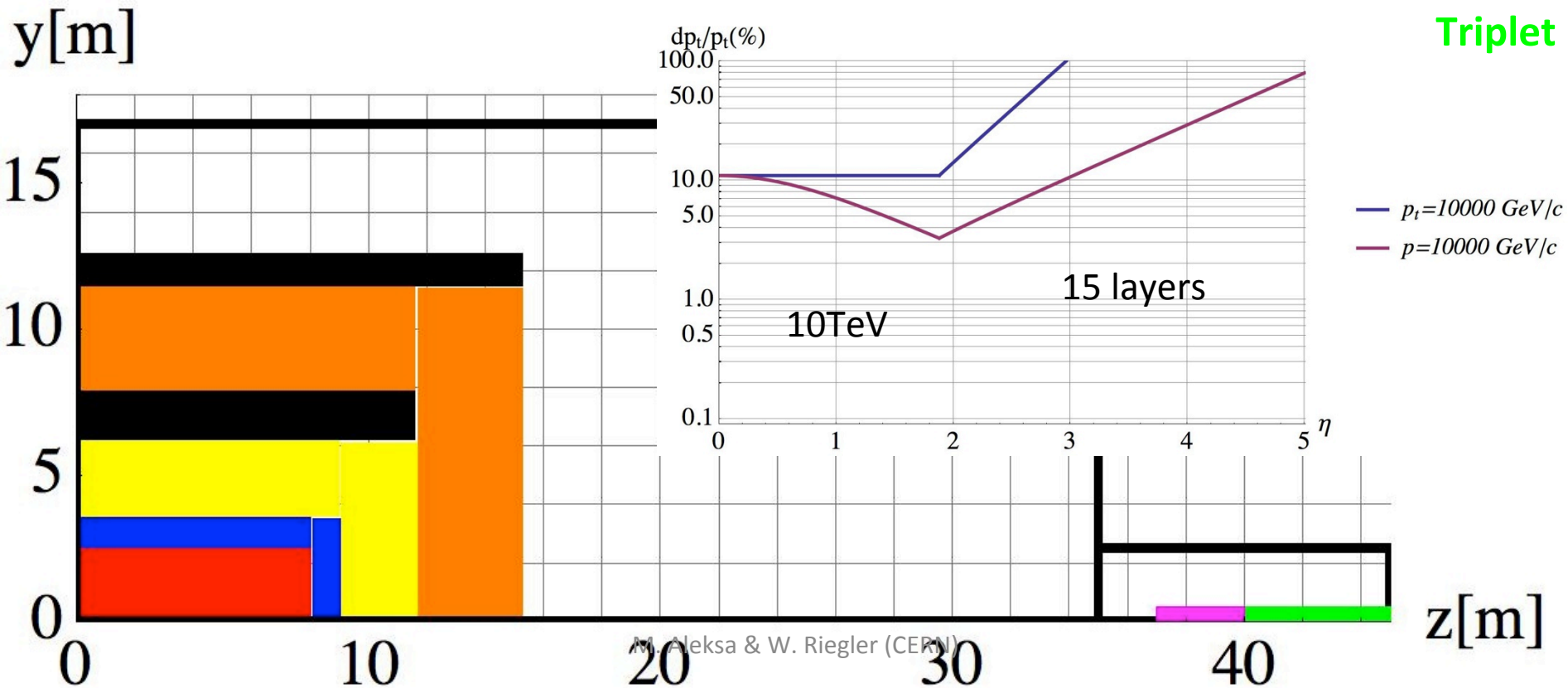
- Forward calorimetry moved to large distance from $\eta = 3.5$ for reduced occupancy and radiation load



Twin Solenoid BL² Scaling

- How to achieve 10% for a 10TeV charged particle **assuming tracker with nowadays layer resolution** (~20 μ m)?
- Solenoid and shielding solenoid with B=6T in Tracker and B=2.5T in Muon System
- Tracker r=2.5m, L=16m, tracking layer resolution similar to CMS detector
- ECAL+HCAL = 3.4m = 12 λ_{int}
- Momentum resolution gets marginal at $\eta > 3$.

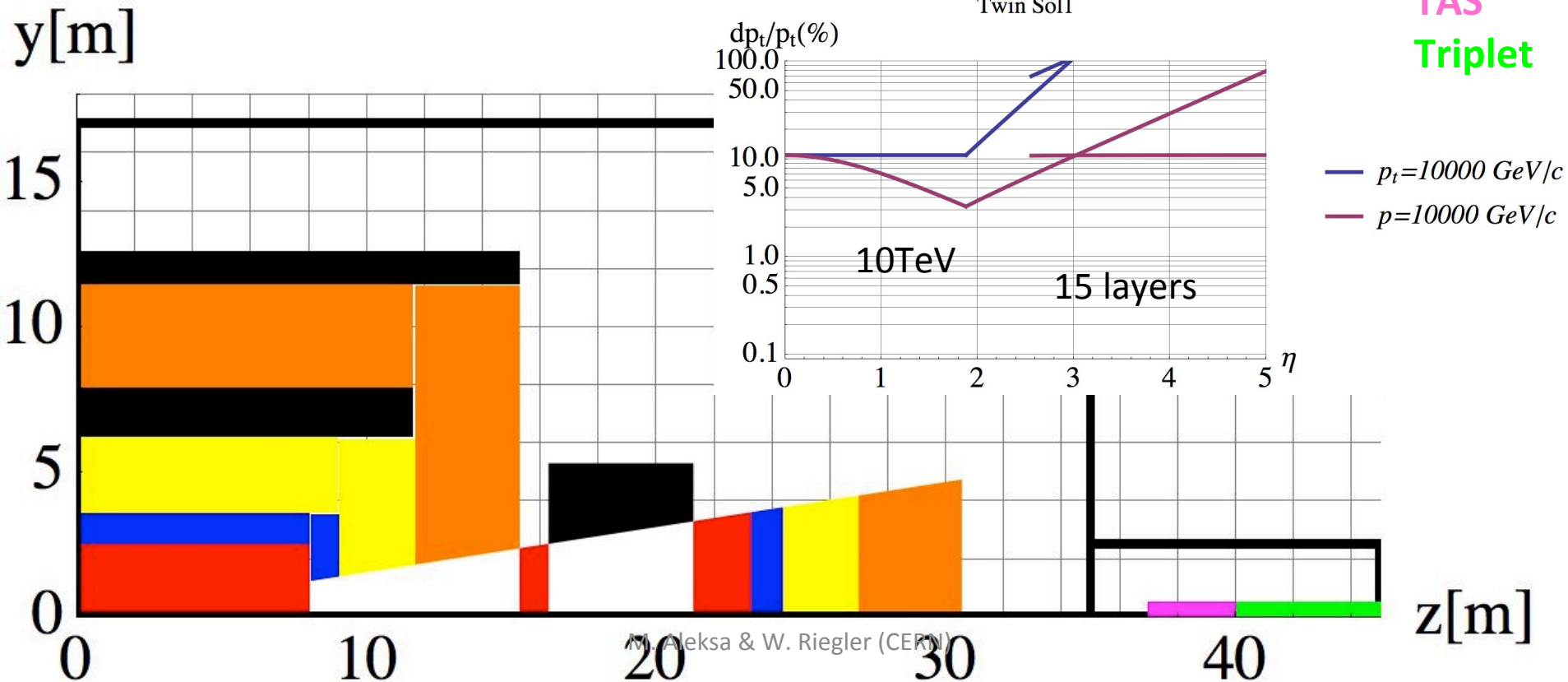
Tracker
Emcal
Hcal
Muon
Coil
TAS
Triplet

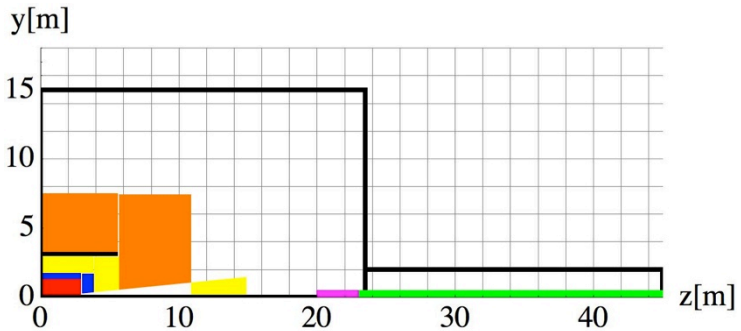


Twin Solenoid BL² Scaling + Forward Dipole

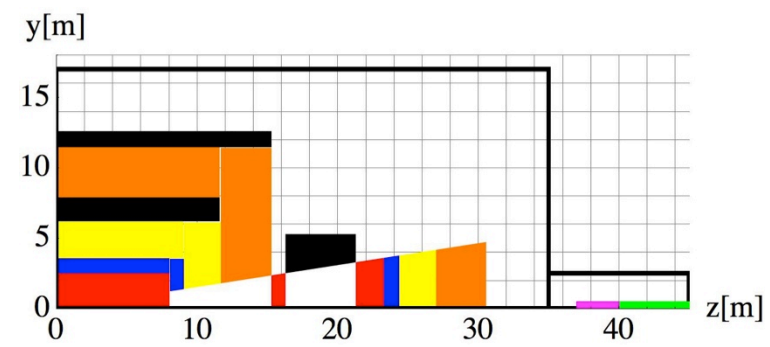
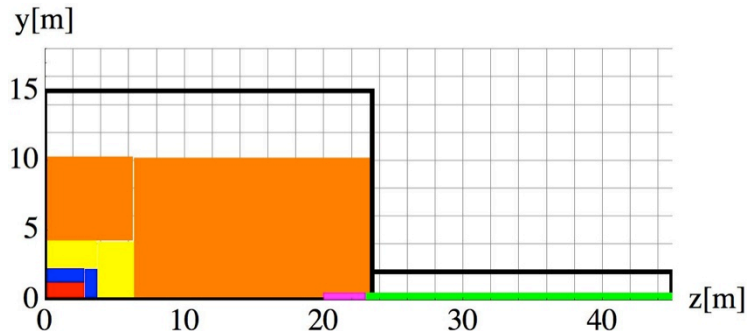
- Opening at $\eta = 2.5$
- Adding a **forward Dipole for momentum spectroscopy**.
- Moving forward calorimeters to larger distance decreasing the particle densities and overlaps.
- Allows separate instrumentation and upgrade of forward detectors
- Integration and maintenance is a challenge

Tracker
Emcal
Hcal
Muon
Coil
TAS
Triplet



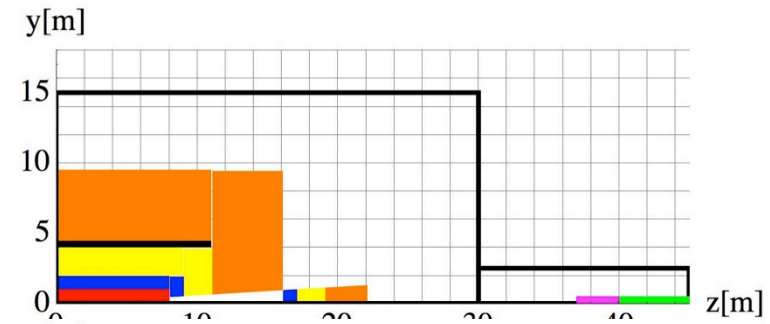


CMS & ATLAS

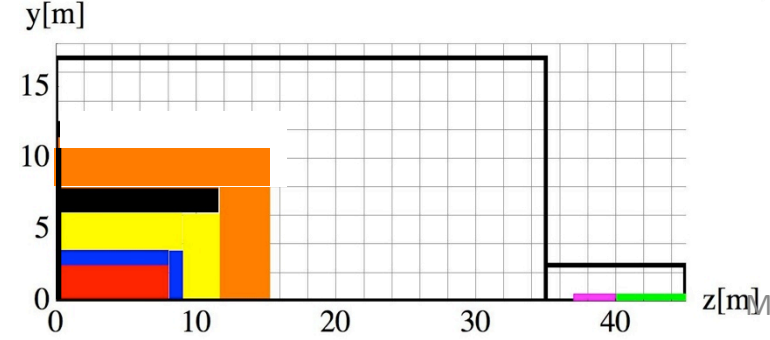


Twin Solenoid + Dipole

Popular at Present



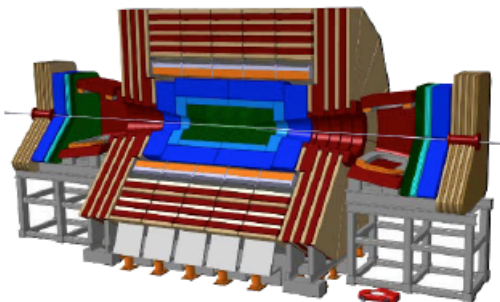
CMS+



Partially shielded large solenoid

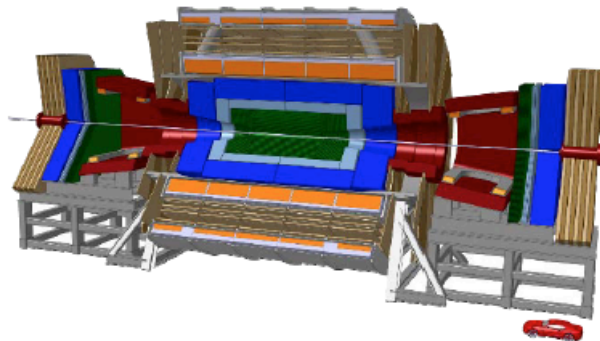
FCC Magnet System Concepts

Inclusion of Dipoles in the Forward region for momentum measurement over a large eta range.



Large coil ($r=6\text{m}$, 6T) with Yoke that returns all the flux:

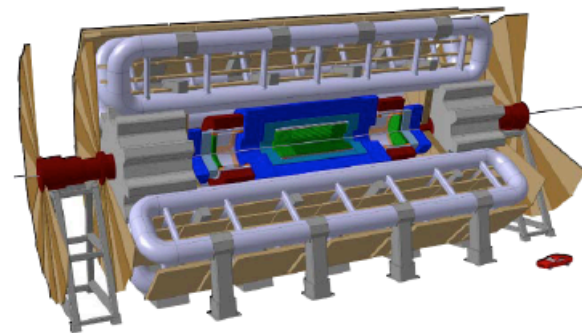
Huge mass,
Iron very expensive



Large coil ($r=6\text{m}$, 6T) with active shielding



This concept is at present studied in quite some detail to have a baseline reference design.
But we have to stay very open for alternative designs



Scaling the ATLAS approach.

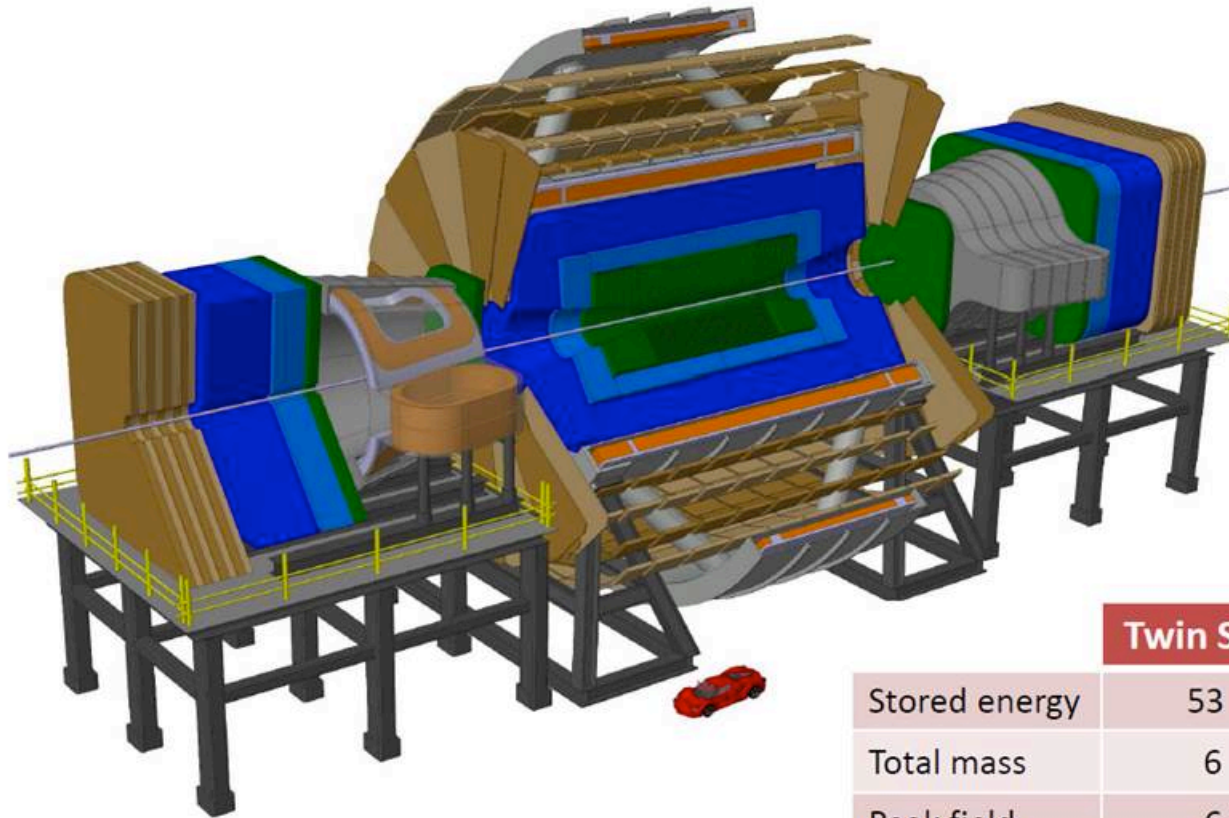
The ATLAS 'standalone' Muon Toroid was motivated by things like:

- worries that trackers might not work at LHC rate
- Space for excellent HCAL, good jet calorimetry
- Independent magnet system

These points are not very strong as of today.

Twin Solenoid + Dipole Magnet System

Matthias Mentink, Alexey Dudarev, Helder Filipe Pais Da Silva, Christophe Paul Berriaud, Gabriella Rolando, Rosalinde Pots, Benoit Cure, Andrea Gaddi, Vyacheslav Klyukhin, Hubert Gerwig, Udo Wagner, and Herman ten Kate



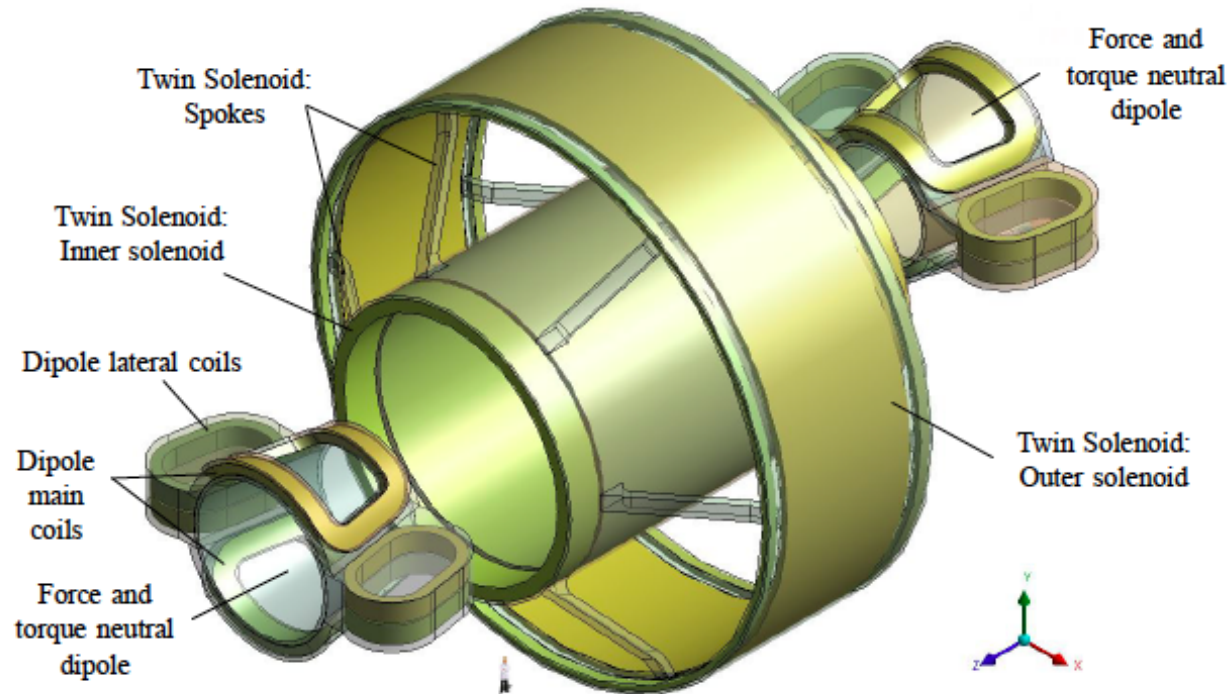
FCC Air core Twin solenoid and Dipoles

State of the art high stress / low mass design.

	Twin Solenoid	Dipole
Stored energy	53 GJ	2 x 1.5 GJ
Total mass	6 kt	0.5 kt
Peak field	6.5 T	6.0 T
Current	80 kA	20 kA
Conductor	102 km	2 x 37 km
Bore x Length	12 m x 20 m	6 m x 6 m

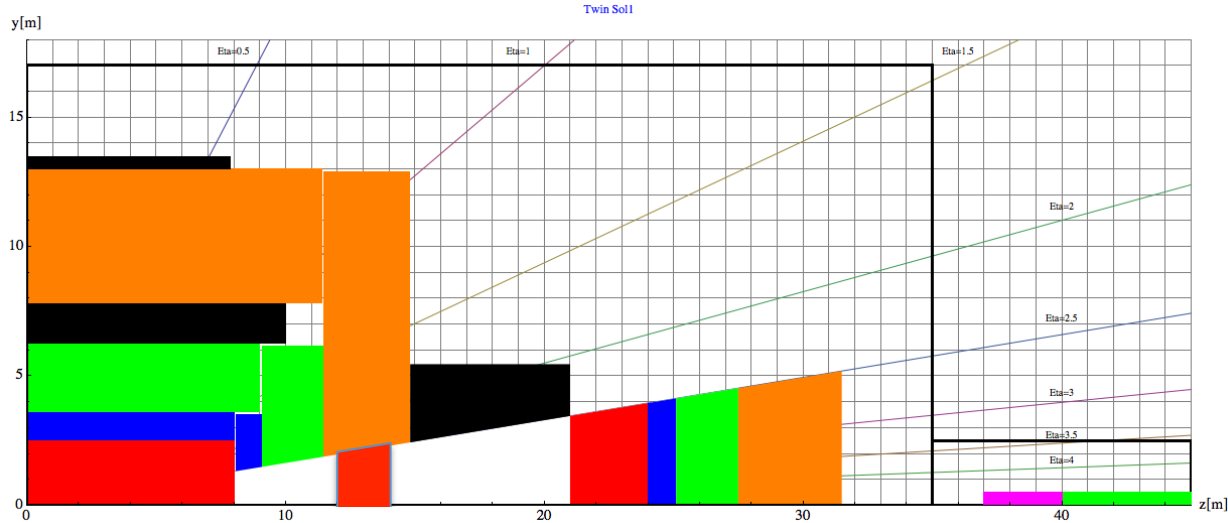
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Dipoles are also actively shielded with SC coils → No Iron Yoke → Decoupling of mechanical forces between solenoid and dipole.

Baseline Geometry, Twin Solenoid



Barrel:

Tracker available space:
R=2.1cm to R=2.5m, L=8m

EMCAL available space:
R=2.5m to R= 3.6m → dR= 1.1m

HCAL available space:
R= 3.6m to R=6.0m → dR=2.4m

Coil+Cryostat:
R= 6m to R= 7.825 → dR = 1.575m, L=10.1m

Muon available space:
R= 7.825m to R= 13m → dR = 5.175m
Revision of outer radius is ongoing.

Coil2:
R=13m to R=13.47m → dR=0.475m, L=7.6m

Endcap:

EMCAL available space:
z=8m to z= 9.1m → dz= 1.1m

HCAL available space:
z= 9.1m to z=11.5m → dz=2.4m

Muon available space:
z= 11.5m to z= 14.8m → dz = 3.3m

Forward:

Dipole:
z= 14.8m to z= 21m → dz=6.2m

FTracker available space:
z=21m to R=24m, L=3m

FEMCAL available space:
Z=24m to z= 25.1m → dz= 1.1m

FHCAL available space:
z= 25.1m to z=27.5m → dz=2.4m

FMuon available space:
z= 27.5m to z=31.5m → dz=4m

Tracking

Radiation Estimate for Inner Tracker Layers

Scaling radiation load of first Pixel layer at $r=3.7\text{cm}$ from ATLAS PHASE II tracker numbers to find the orders of magnitude:

HL-LHC 3ab^{-1}

1MeVneq Fluence (NIEL) = $1.5 \times 10^{16} \text{ cm}^{-2}$

Dose = **5MGy**

FCC 3ab^{-1}

1MeVneq Fluence = $3 \times 10^{16} \text{ cm}^{-2}$

Dose = **10MGy**

FCC 30ab^{-1}

1MeVneq Fluence = $3 \times 10^{17} \text{ cm}^{-2}$

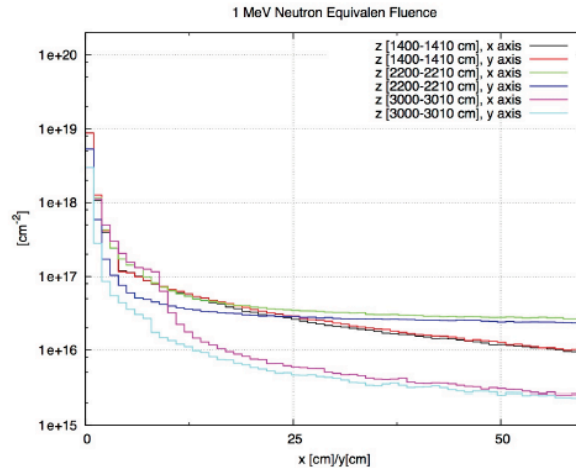
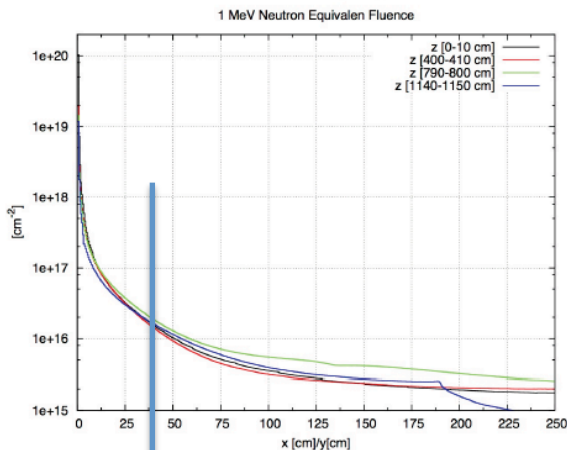
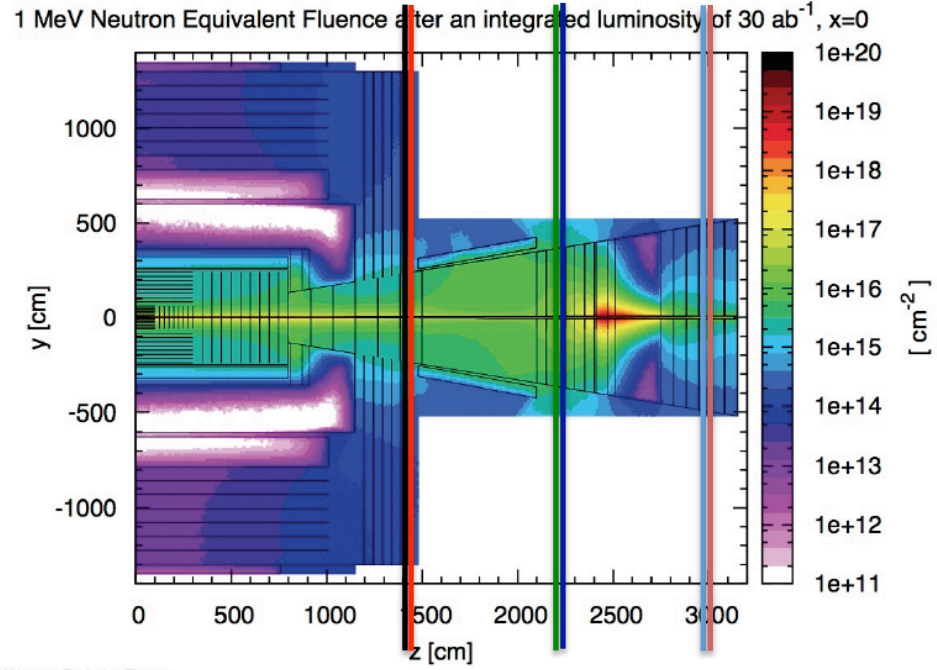
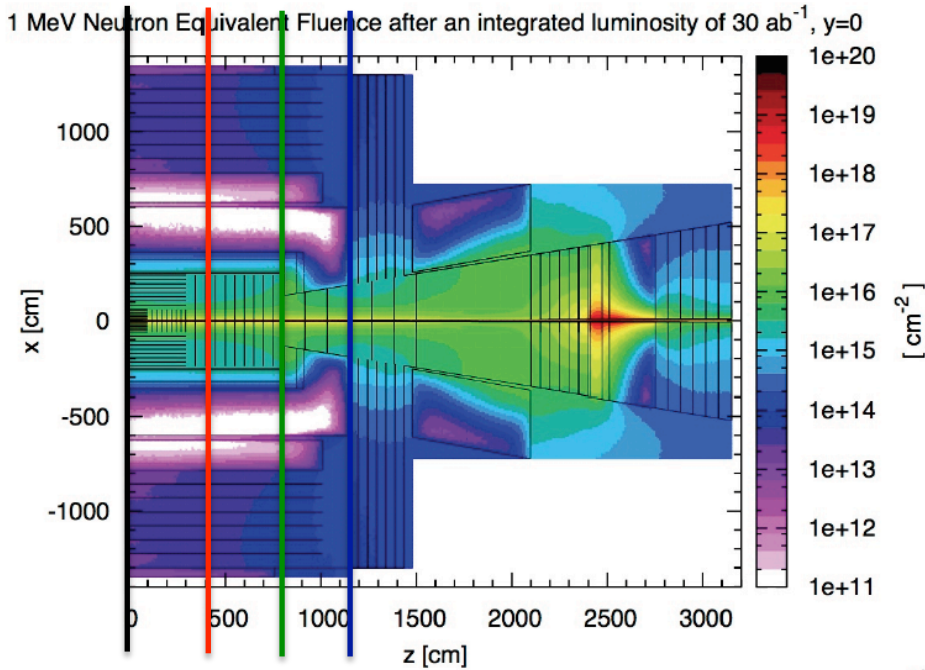
Dose = **100MGy**

FCC 30ab^{-1} $r_{\text{pixel}}=2.1\text{cm}$:

1MeVneq Fluence = 10^{18} cm^{-2}

Dose = **350MGy**

1 MeV Neutron Equivalent Fluence (FLUKA simulation, M.I.Besana)



ATLAS Phase II Tracker innermost pixel layer

Detailed radiation simulations with FLUKA for the baseline detector exist.

Radiation load in the trackers shows primarily radial dependence from the beamline, weak dependence on z (as expected).

For radii $< 50 \text{ cm}$ we exceed the HL-LHC numbers (10^{16} cm^{-2}) by up to 2 orders of magnitude

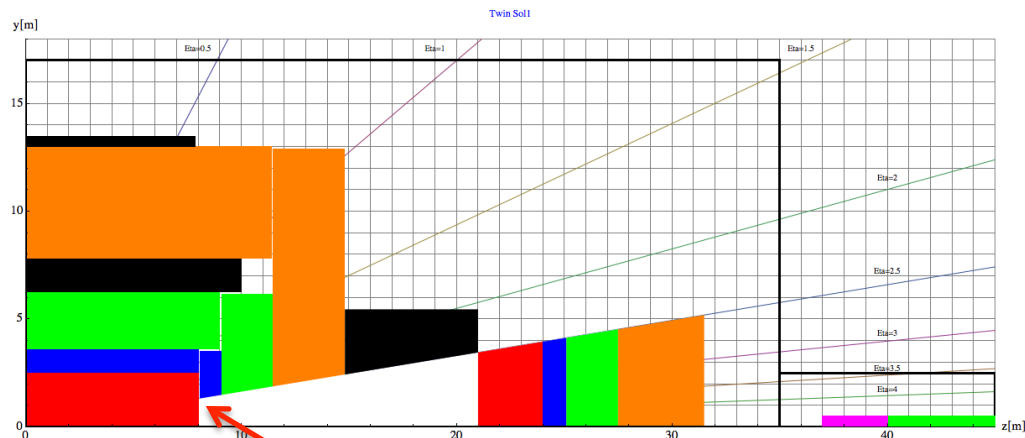
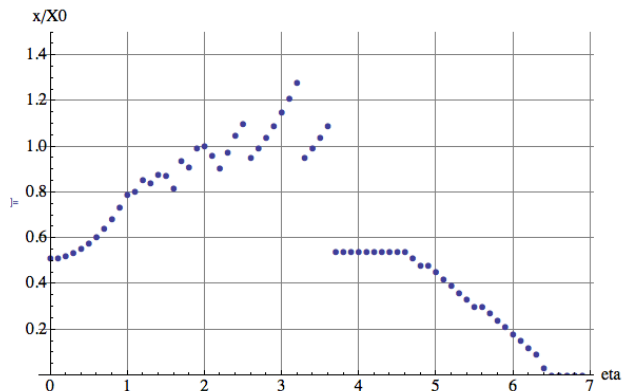
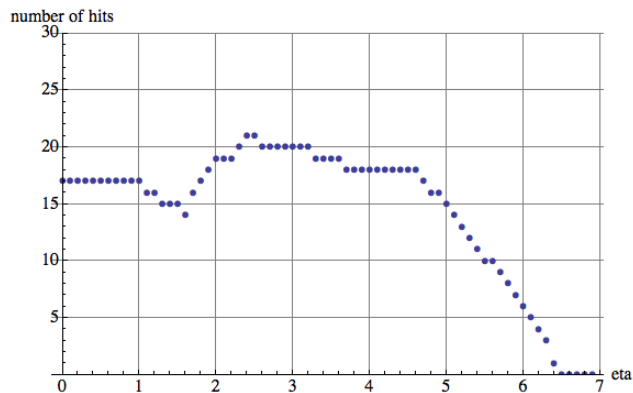
→ Technology challenge !

Simplified Tracker Assumptions

Neglecting radiation for a moment: is 10% resolution achievable (for 10TeV)?

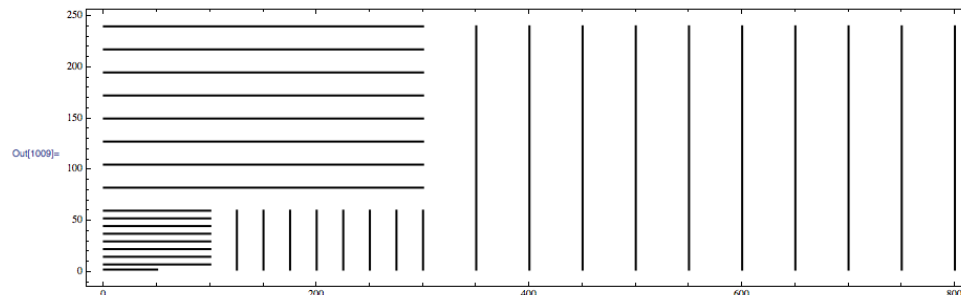
Material composition in Volume (%):
 Si 20%, C 42%, Cu 2%, Al 6%, Plastic 30%
 X_0 of this mix: 14.37cm

We assume 3% of radiation length per layer,
 i.e. each layer has a thickness of 0.43cm.

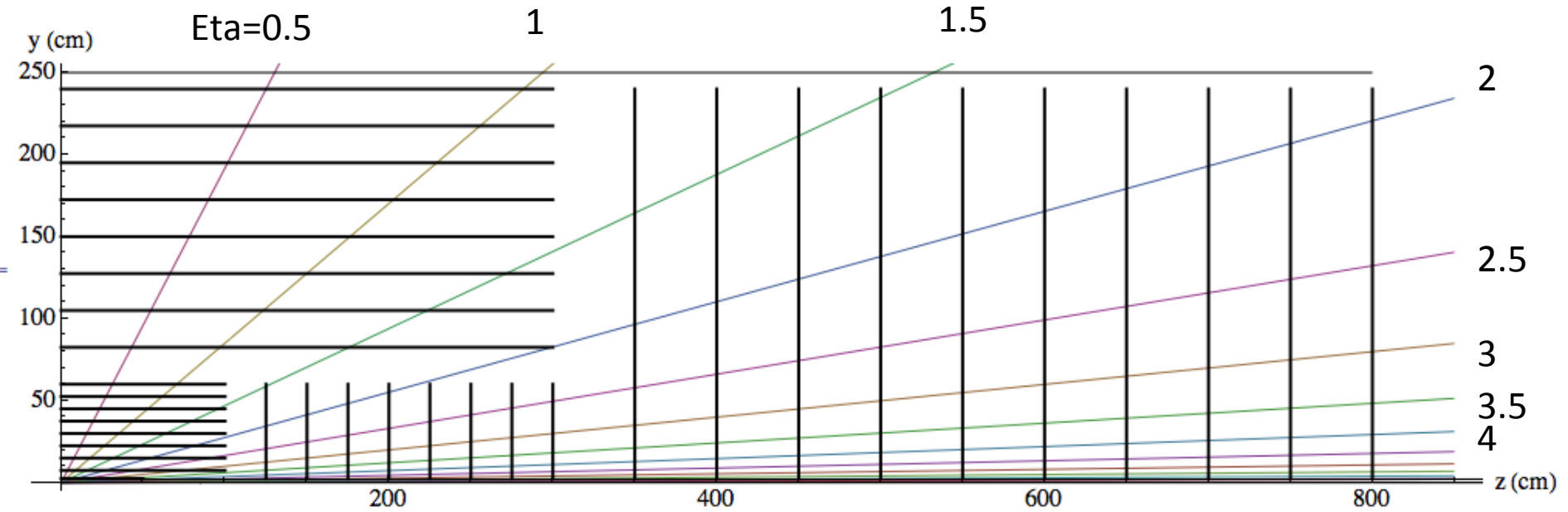


$R_{out}=2.4m$
 Half the lever arm at $\eta=2.6 \rightarrow L=8m$

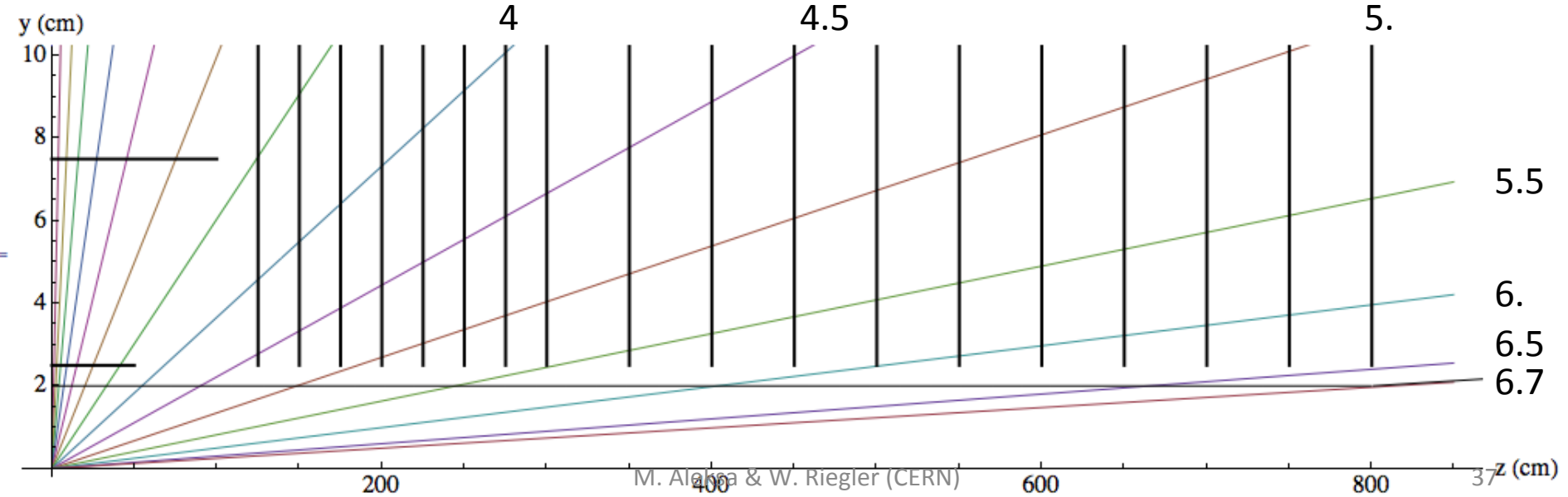
```
In[1009]:= gr1 = Graphics[1, AspectRatio -> 2.5 / 8, Frame -> True]
```



Tracker



```
Show[p11, gr1, p10, plbeampipe, PlotRange -> {{0, 850}, {0, 10}}, AspectRatio -> 250 / 850]
```



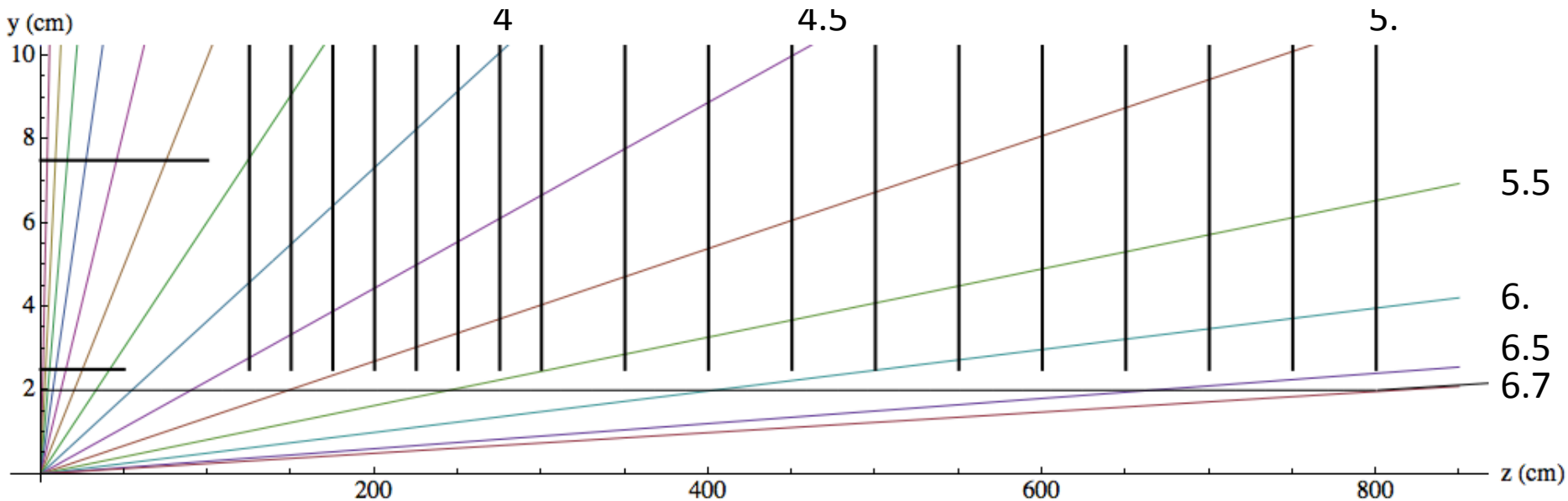
Tracker

Side remark: A track at $\eta=5$ hits the first detector layer only at 200cm distance from the IP. We cannot dream of B-tagging a'la LHCb.

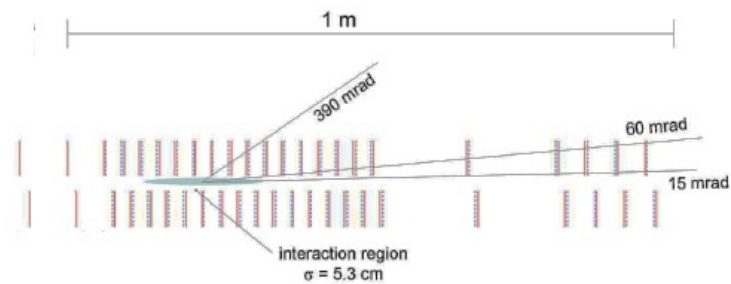
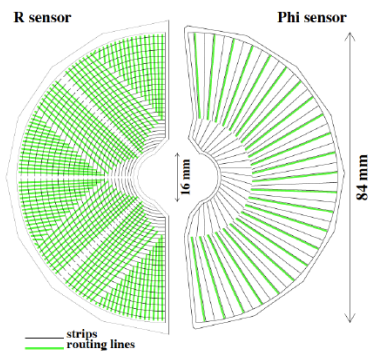
LHCb has the VELO with discs only a few mm from the beam in a secondary vacuum.

This arrangement has significant infrastructure around the IP which is not compatible with a co-existent central detector.

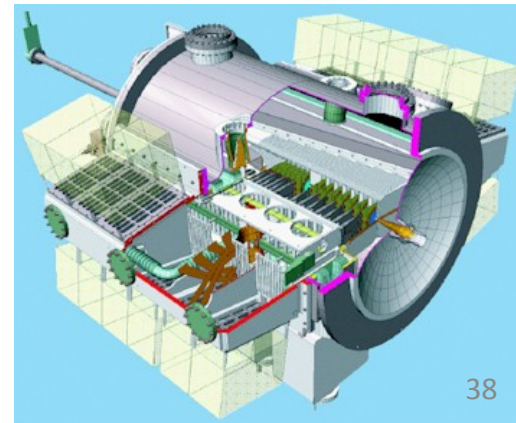
→ Clever ideas needed !!



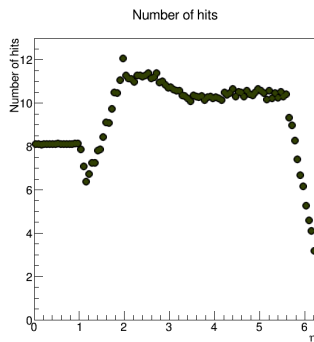
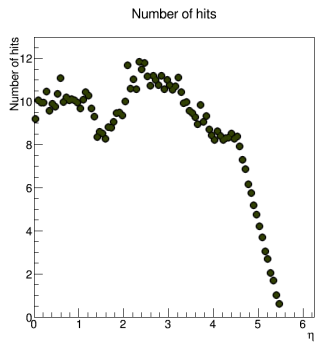
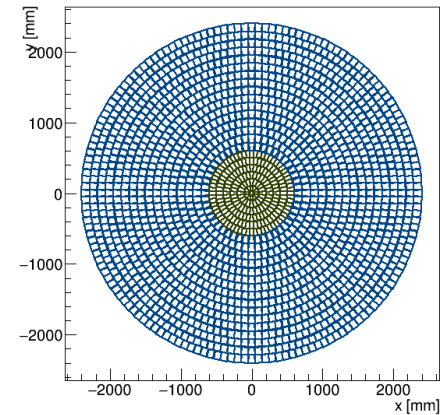
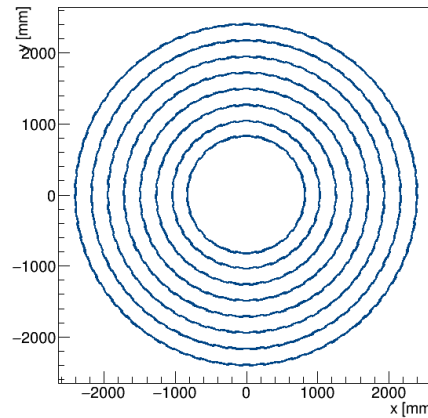
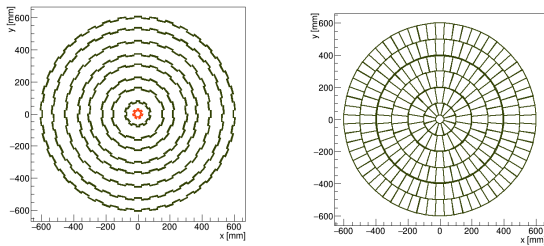
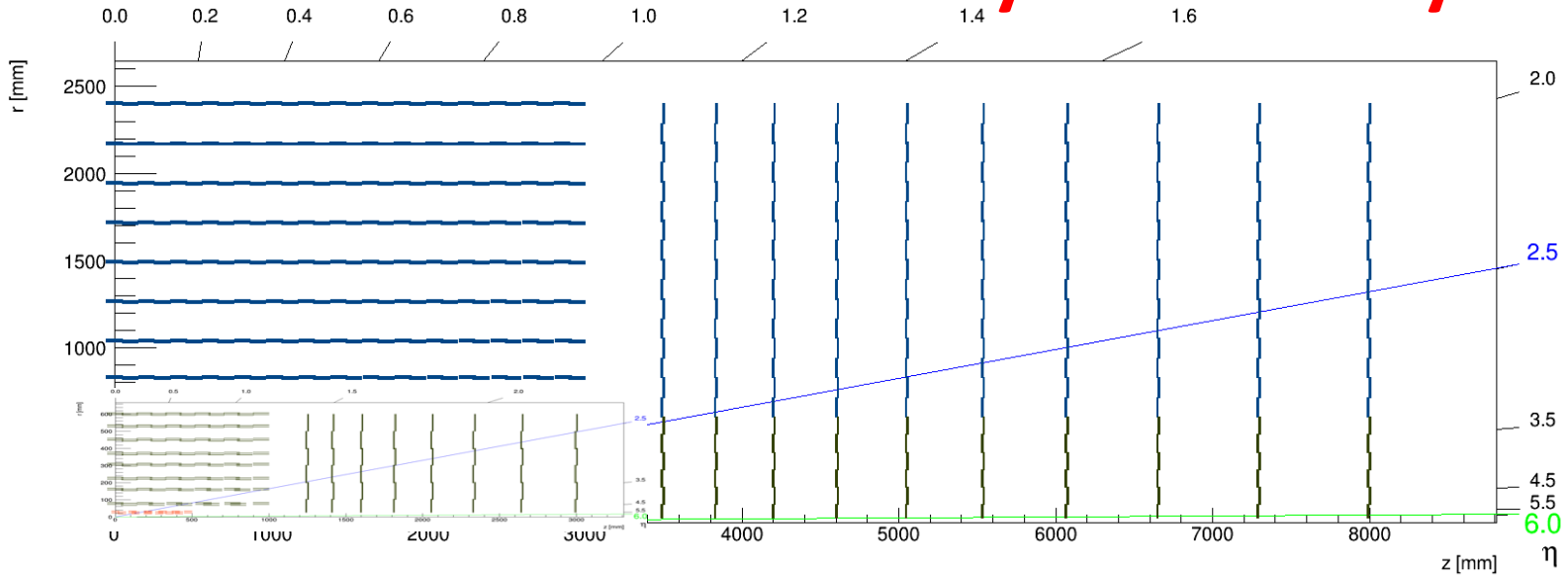
LHCb:



M. Aleksa & W. Riegler (CERN)



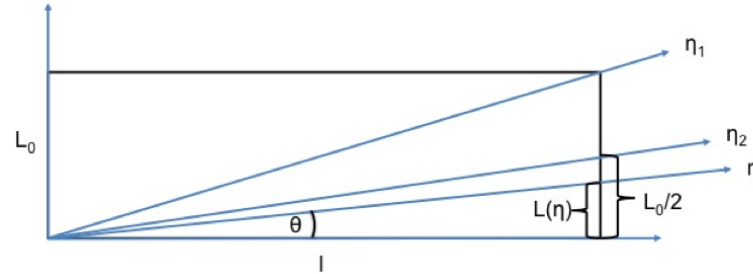
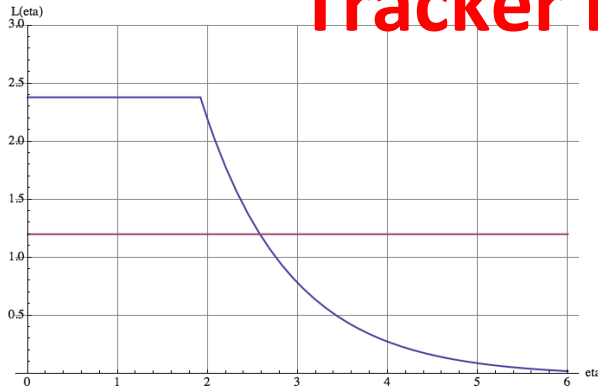
Tracker – Realistic FCC Layout in TKLayout



Realistic Layout with correct modules using TKLayout (CMS PhaseII upgrade tool)

http://fcc-tklayout.web.cern.ch/fcc-tklayout/FCChh_Option2/errorsTRK.html

Tracker Resolution First Principles



$$\eta_1 = -\ln \tan \left(\frac{1}{2} \arctan \frac{L_0}{l} \right) \quad \eta_2 = -\ln \tan \left(\frac{1}{2} \arctan \frac{L_0}{2l} \right)$$

For a geometry with $L_0 = 2.4m$ and $l = 8m$ we have $\eta_1 = 1.9$ and $\eta_2 = 2.6$

$$L(\eta) = L_0 \quad \eta < \eta_1 \quad L(\eta) = L_0 \frac{\sinh \eta_1}{\sinh \eta} \quad \eta > \eta_1$$

$$\left. \frac{\Delta p_T}{p_T} \right|_{reso.} = \frac{\sigma p_T}{0.3BL(\eta)^2} \sqrt{\frac{720}{N(\eta) + 4}} \quad \left. \frac{\Delta p_T}{p_T} \right|_{m.s.} = \frac{0.0136}{0.3BL(\eta)} \sqrt{\frac{x}{X_0}}(\eta)$$

$$\frac{\Delta p_T}{p_T} = \sqrt{\left(\left. \frac{\Delta p_T}{p_T} \right|_{reso.} \right)^2 + \left(\left. \frac{\Delta p_T}{p_T} \right|_{m.s.} \right)^2}$$

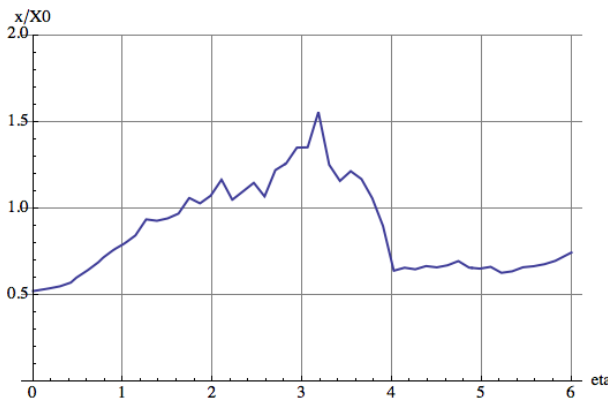
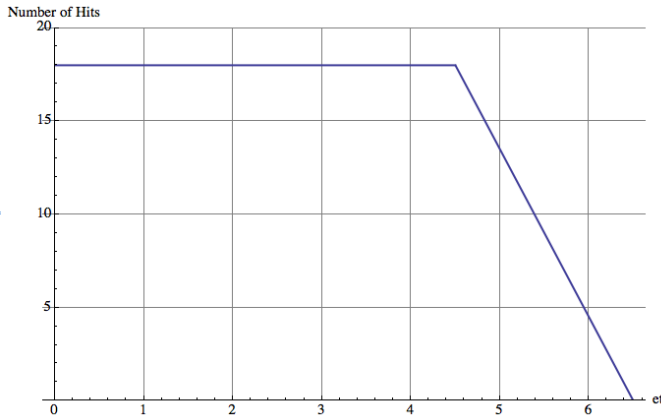
ln[60]:= **L0 = 2.4;**

ln[61]:= **l = 8;**

ln[62]:= **B = 6;**

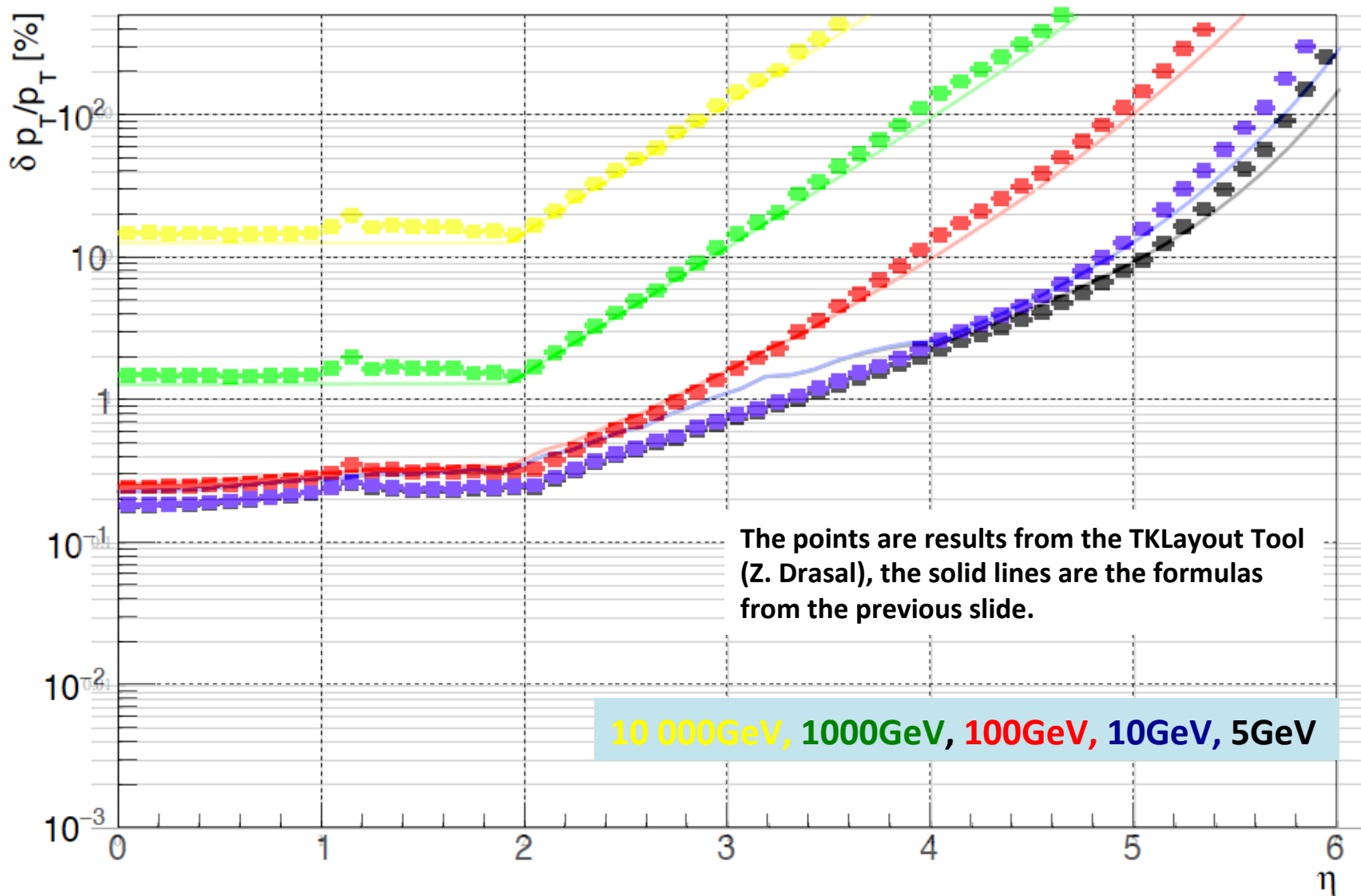
ln[63]:= **sig = 23 * 10 ^ (-6);**

40



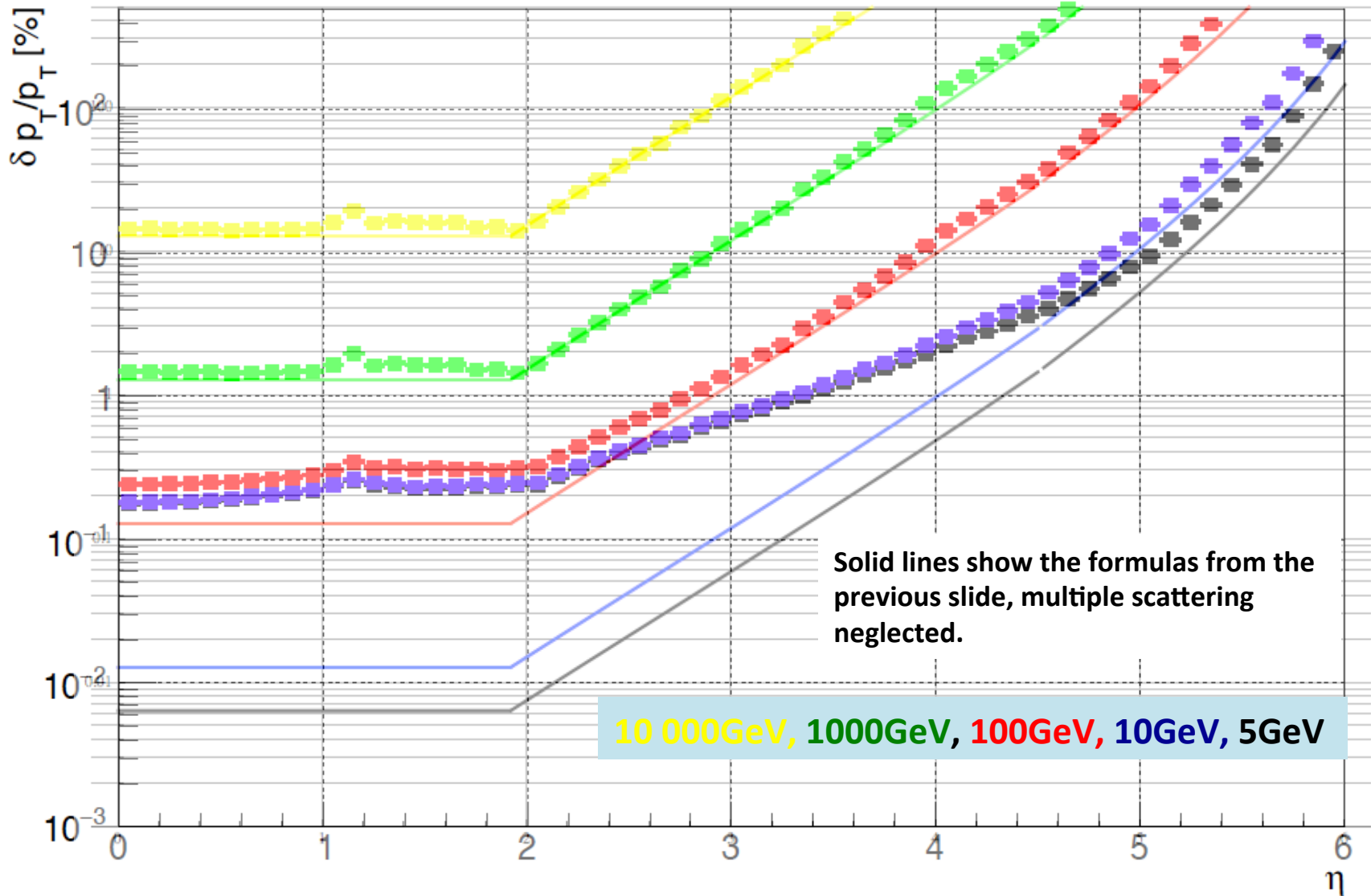
Tracker

p_T resolution versus η - const P_T across η



Tracker

p_T resolution versus η - const P_T across η



Tracker

$$\frac{\Delta p_T}{p_T} \Big|_{reso.} = \frac{\sigma p_T}{0.3BL(\eta)^2} \sqrt{\frac{720}{N(\eta) + 4}} \quad \frac{\Delta p_T}{p_T} \Big|_{m.s.} = \frac{0.0136}{0.3BL(\eta)} \sqrt{\frac{x}{X_0}(\eta)}$$

Large BL^2 needed for high momenta, but large BL also key to minimize multiple scattering contribution \rightarrow 10% for a 10TeV charged particle within reach for tracker radius of 2.5m ($|\eta| < 2$)!

Could we also reach that with a CMS like design (smaller tracker radius): How to scale the system and keep the performance constant ?

At constant B and 1/2 the tracker radius (2.5m \rightarrow 1.2m) \rightarrow free bore of solenoid from 12m to 10m we need:

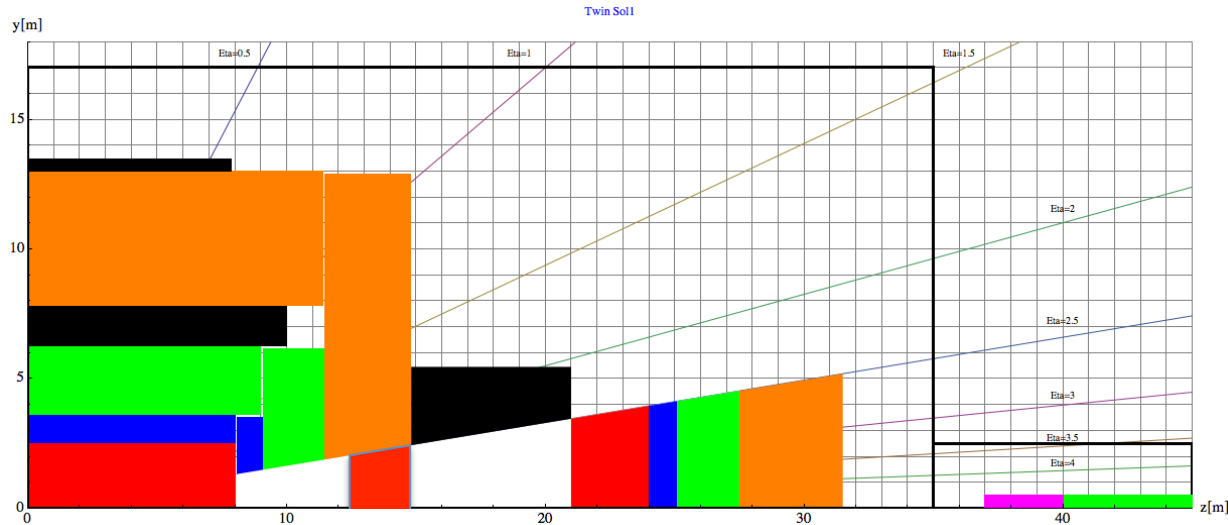
- 4 times the tracker resolution ($20\mu\text{m} \rightarrow 5\mu\text{m}$) and
- 4 times less material budget ($x/X_0=50\%$ at $\eta=0$ to $x/X_0=12.5\%$ at $\eta=0$ i.e. 3% per Layer to 0.75% per layer)

These values are challenging but not out of reach.

\rightarrow A final choice is part of an optimization that depends on future technologies

Forward Tracking

Forward Tracking Resolution, Position Resolution

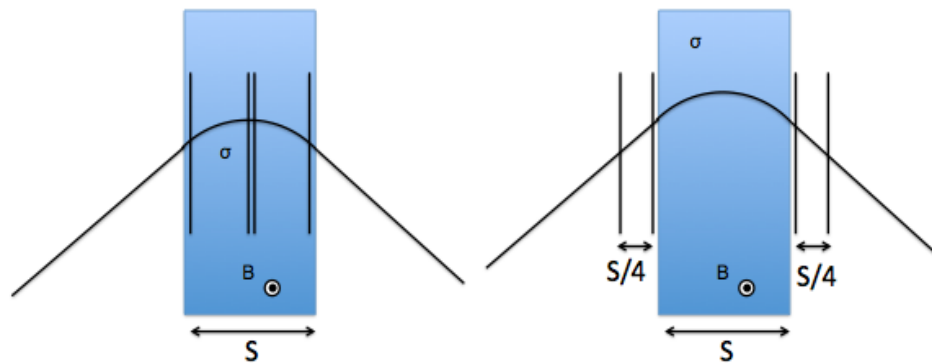


Using 4 tracking stations for a dipole with constant magnetic field and length S , the optimum spectrometer resolution is achieved by placing 2 stations in the center and one on each end to measure the sagitta.

The same performance is achieved by placing the chambers outside the dipole at separation of $S/4$.

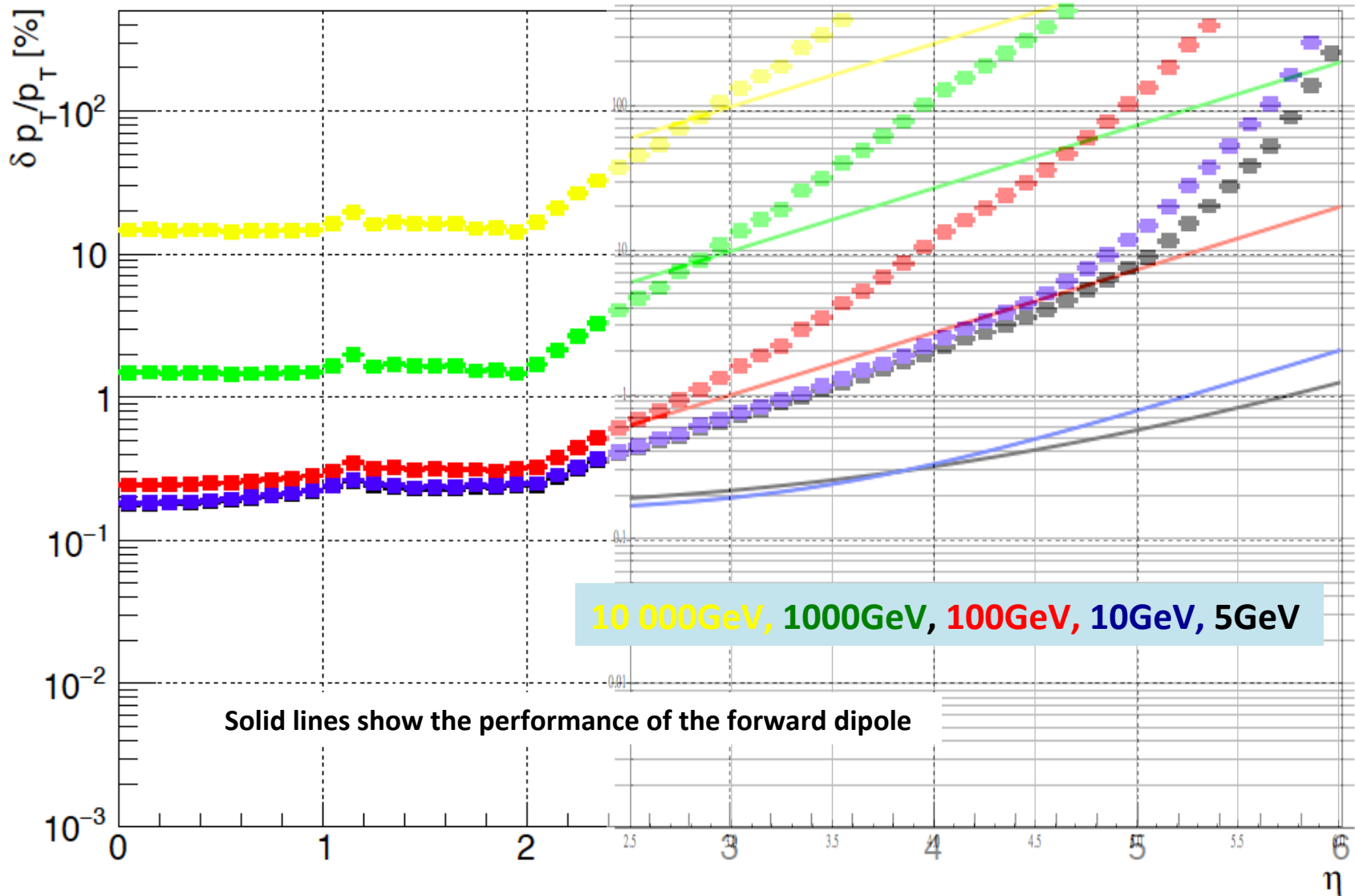
This is what LHCb uses, because if space is available it is more easy to implement the detectors outside, and also avoid occupancy from loopers in the field (details on catching Ks etc. are of course to be considered ...)

We use this idea for now (is also easier to calculate ! It is just the $\int B dl$ that counts)



Forward Tracker Resolution

p_T resolution versus η - const P_T across η

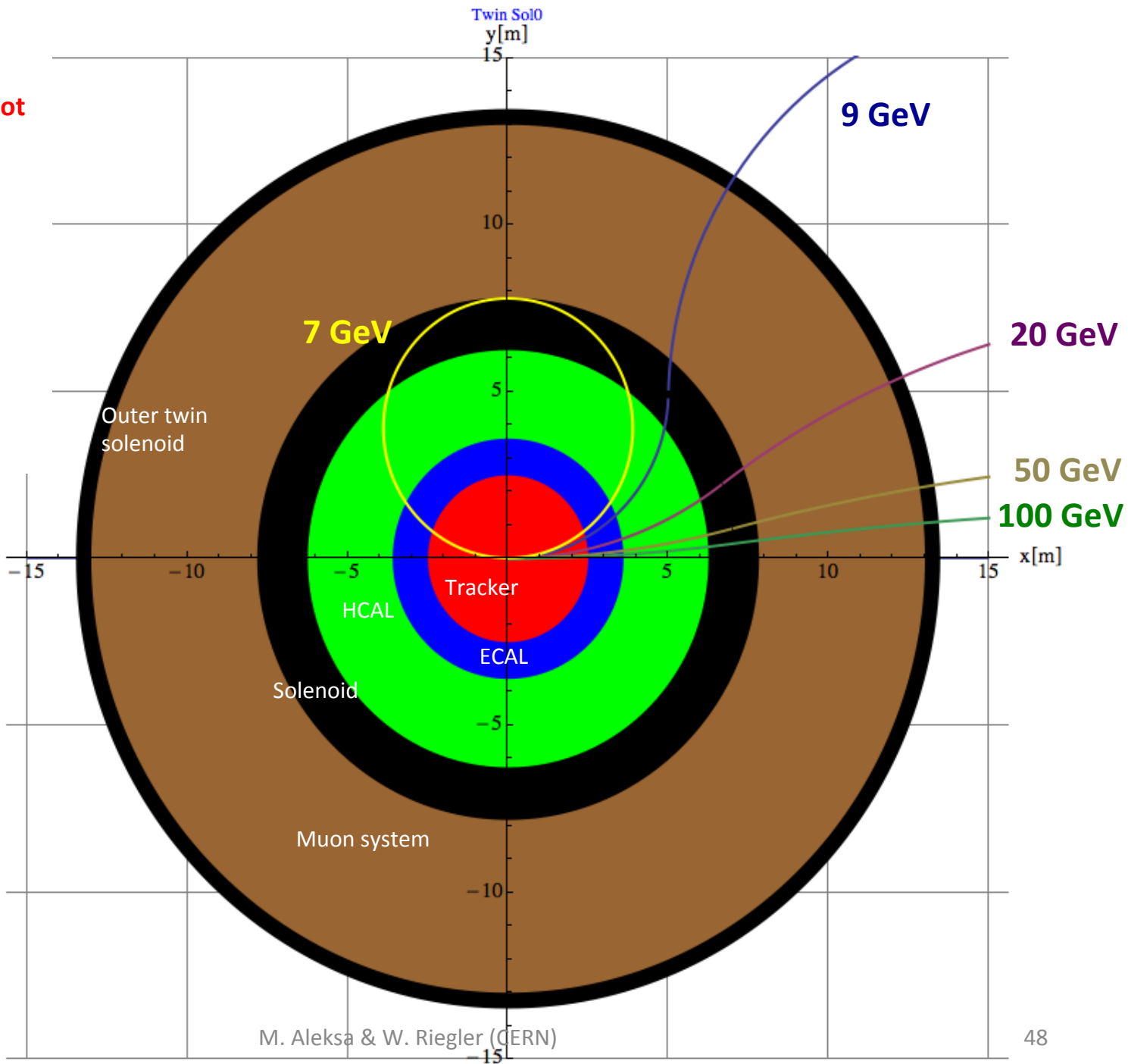


Muon System

At $B_0=6T$ and $R_0=6m$,
Muons below 7GeV do not
enter the muon system.

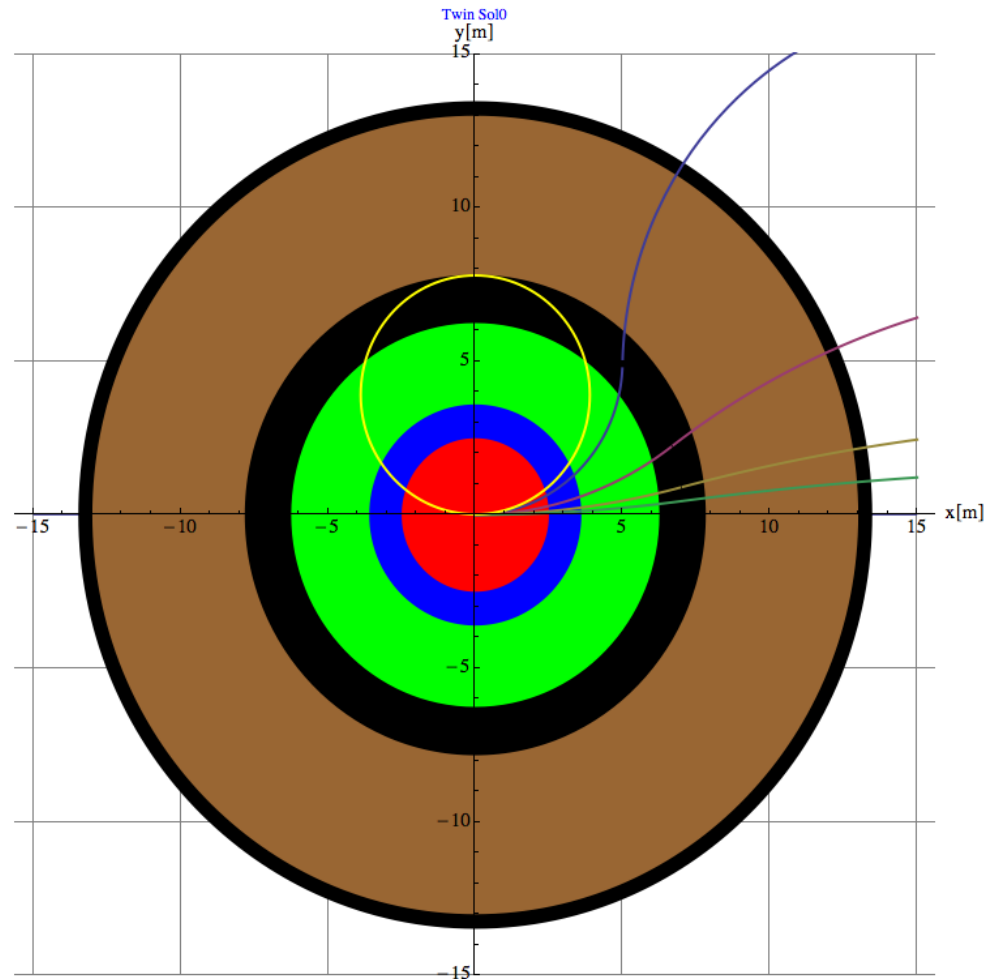
No Muon Trigger below
7GeV.

Possibly muon ID with
HCAL.



Muon Momentum Can Be Measured by...

- 1) The inner tracker
→ resolution plots from before
- 2) A 'standalone' sagitta measurement in the muon system (no iron → precise !)
- 3) The track angle at the entrance of the muon system → Trigger
- 4) The combined fit of inner tracker and outer layers of the muon system.



CMS Muon Performance

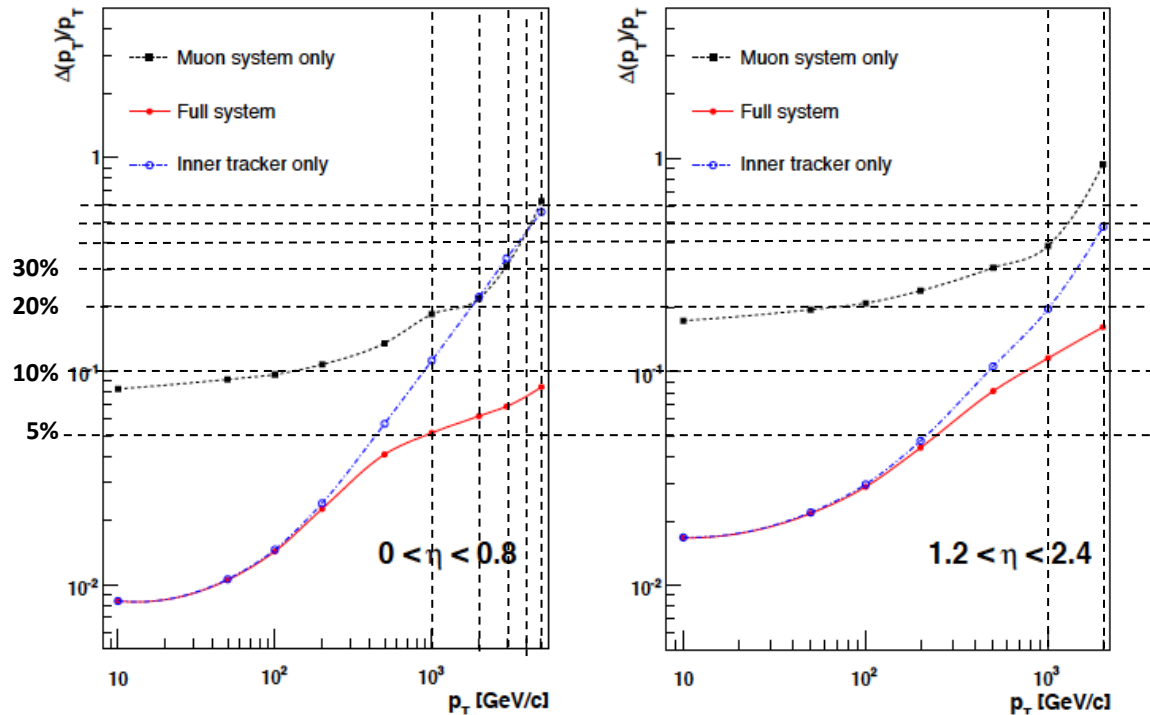


Figure 1.2: The muon transverse-momentum resolution as a function of the transverse-momentum (p_T) using the muon system only, the inner tracking only, and both. Left panel: $|\eta| < 0.8$, right panel: $1.2 < |\eta| < 2.4$.

$P_T=1\text{TeV}/c$, $0 < \eta < 0.8$:

20% muon standalone (angle)
10% inner tracker only
5% combined

$P_T=1\text{TeV}/c$, $\eta < \eta < 2.4$:

40% muon standalone (angle)
20% inner tracker only
10% combined

Sagitta Measurement in the Muon System

The return field is 2.45T

Measuring over the 5m lever arm with stations of $\sigma=50\mu\text{m}$ resolution we have

$$\frac{dp_T}{p_T} = \frac{\sigma \cdot p_T}{(0.3 \cdot B \cdot L^2)} \cdot 8$$

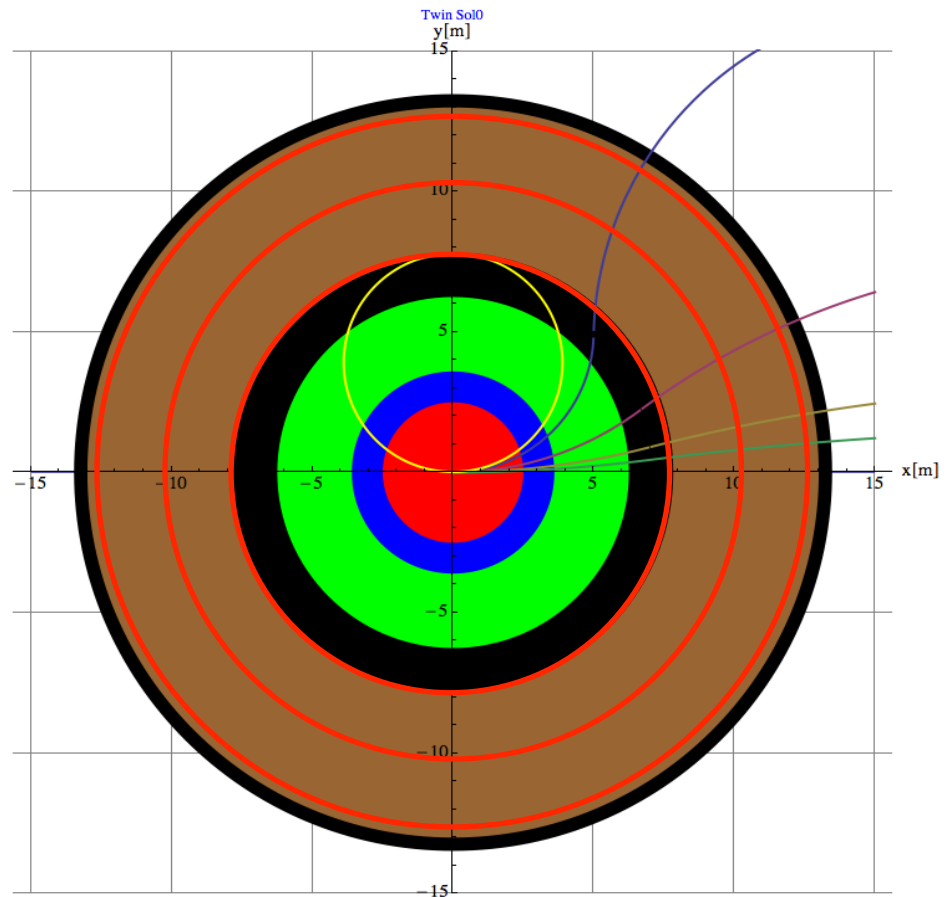
= 20% @ 10TeV

with possibly excellent performance at low p_T due to the absence of iron (vs. CMS) .

but very hard to beat the angular measurement and the inner tracker (10% at 10TeV)

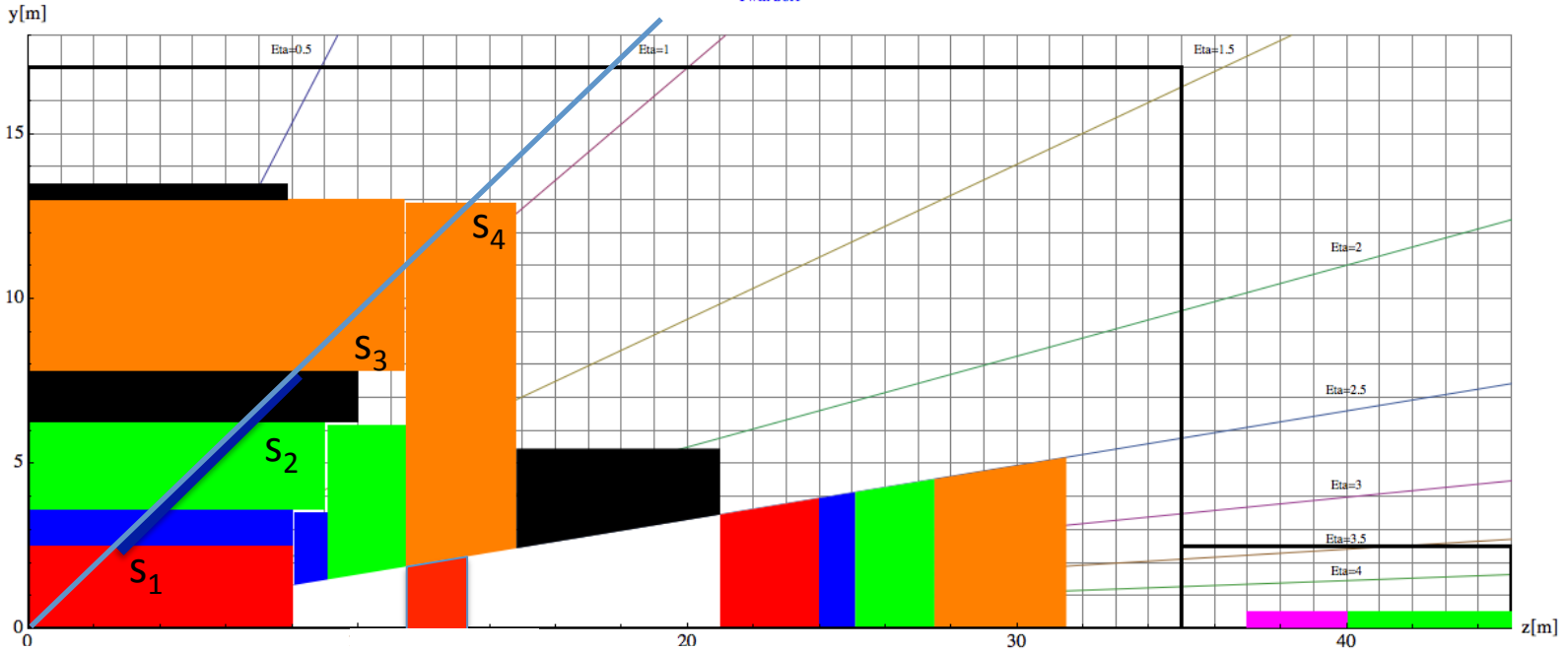
Surface > 5000 m²

CMS sagitta measurement in the muon system is limited to $dp_T/p_T = 20\%$ due to multiple scattering alone.

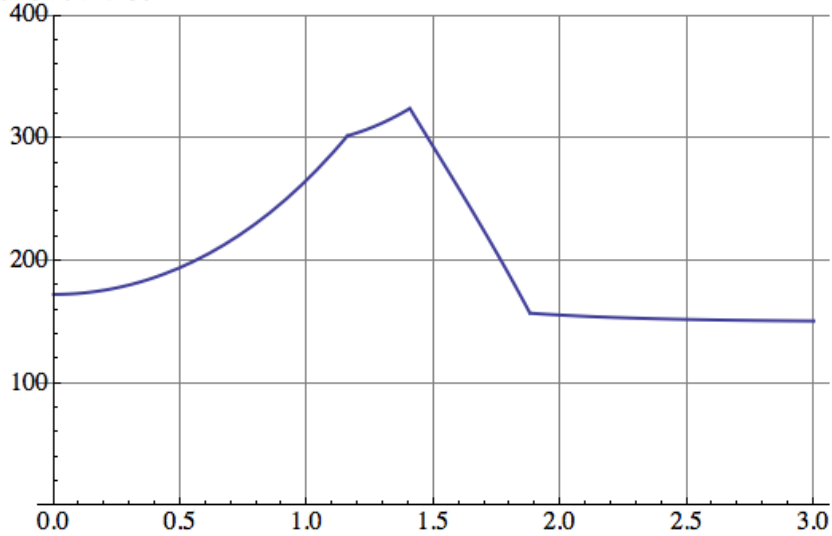


Radiation Length and Angular Deflection (Mult. Scattering)

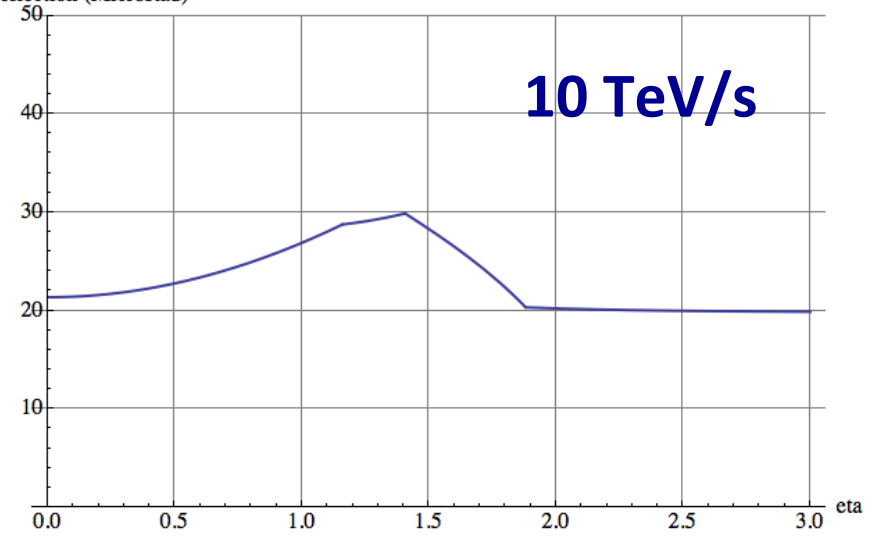
Twin Sol1



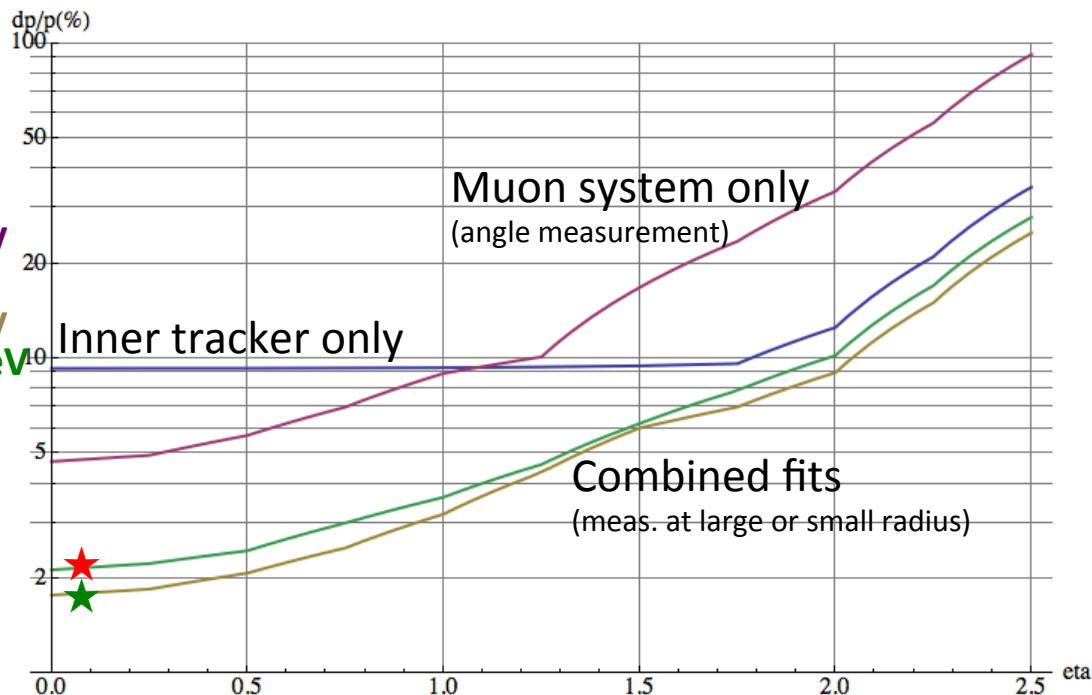
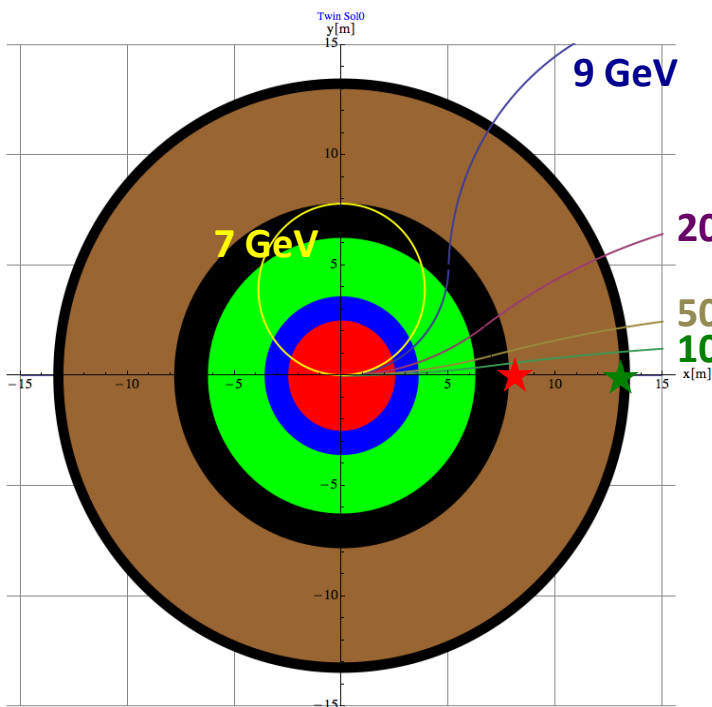
x/X0 of Calo and Coil



Angular Deflection (MicroRad)



Momentum Resolution for a 10 TeV/s Muon



Twin Solenoid assuming inner tracker with baseline resolution curves and multiple scattering limit in the muons system.

$P_T=10\text{TeV}/c$ $\eta = 0$:

5% muon standalone (angle)
10% inner tracker only
2% combined

$P_T=10\text{TeV}/c$ $\eta=2$:

35% muon standalone (angle)
12.5% inner tracker only
8% combined

Compare to the CMS numbers:

$P_T=1\text{TeV}/c$, $0<\eta < 0.8$:

20% muon standalone (angle)
10% inner tracker only
5% combined

$P_T=1\text{TeV}/c$, $\eta 1.2<\eta < 2.4$:

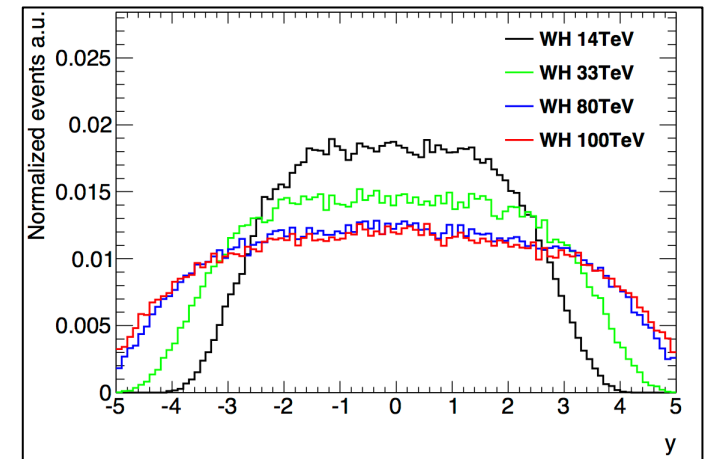
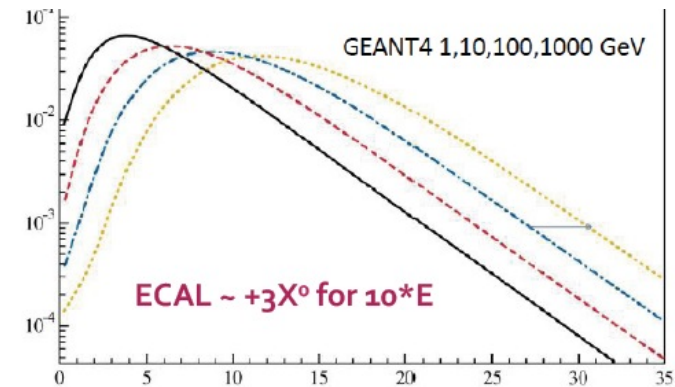
40% muon standalone (angle)
20% inner tracker only
10% combined

Calorimetry

Requirements ECAL

ECAL:

- Depth only moderately sensitive to ν_s : $30X_0$ enough for fully contained e/γ (ATLAS $\sim 22X_0$)
- Large acceptance up to $|\eta|=6$
- High granularity
 - highly collimated final states (**high boost**)
 - **Pile-up mitigation** (up to 1000 events per BC)
 - **Track-cluster matching**, position resolution
 - Pointing resolution
 - Tau reconstruction
- Excellent timing resolution could help for pile-up mitigation.
- High radiation tolerance and stability
- L1 triggering (low p_T thresholds for W and Z will be challenging!)



Requirements ECAL

Some general thoughts:

- High magnetic field and large radius: **Bremsstrahlungs** photons will end up far away from electron (i.e. will mostly not be contained in the same cluster)
 - e.g. distance of e^- and brem γ is up to $\sim 30\text{cm}$ for $20\text{GeV } e^-$, similar problem for photon conversions
- **High pile-up**: pile-up rejection (e.g. for isolation requirement for EM objects) will also need to rely on tracker information

→ EM energy measurement will not be able to rely on the ECAL only → EM energy measurement in FCC will consist in an intelligent combination between tracker measurement and ECAL measurement (of course the jet and E_T^{miss} measurement even more so)

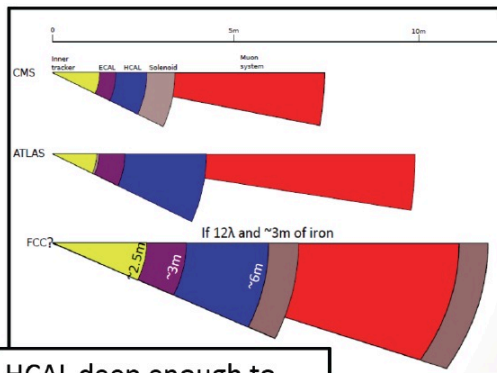
- Track-cluster matching is essential to achieve the above → **fine (lateral) granularity** and good position resolution should be achieved



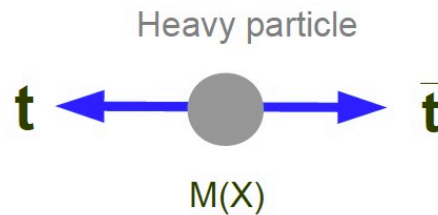
Requirements HCAL

HCAL:

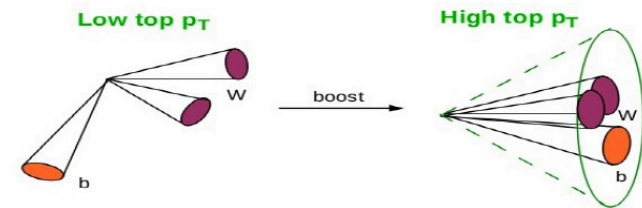
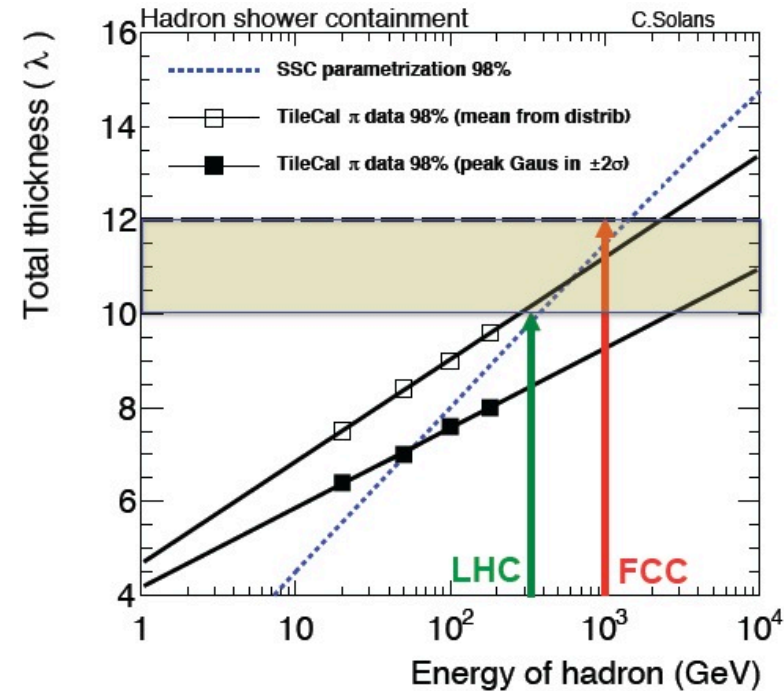
- Jet containment: 8% of single hadron constituents of 30TeV jets have $E > 1\text{TeV}$. 98% containment requires 12λ
- Large acceptance up to $|\eta|=6$
- Highly collimated (boosted) final states
 - Minimal distance between two partons proportional to m/p_T (e.g. top)
- → high granularity also in the HCAL
 - Sub-structure identification will become difficult as the jet cone tends to be very narrow when particles enter the calorimeter → object overlap
 - Tau reconstruction



HCAL deep enough to prevent punch-through



M. Aleksa & W. Riegler (CERN)



HCAL Energy Resolution

Performance of calorimeters improves with energy

$$\frac{\sigma(E)}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

a – stochastic/sampling term,
b - electronic noise term
c - constant term

Single hadrons:

ATLAS: $\sigma_E/E \sim 50\%/\sqrt{E} + 3.0\%$

(small noise term for both)

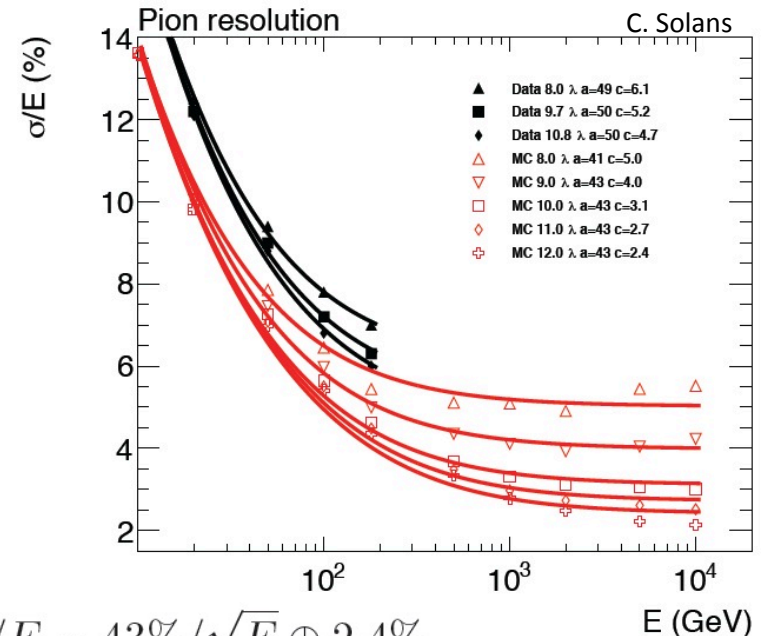
CMS: $\sigma_E/E \sim 100\%/\sqrt{E} + 4.5\%$

Jet $p_T > 5\text{TeV}$: constant term dominates

Reduction of the constant term- need to work on:

- $e/h \neq 1$
- dead material,
- longitudinal and lateral energy leakage,
- non-uniformity calibration,
- transition region, etc.

Achievable resolution at 12λ (ATLAS like HCAL): $\sigma_E/E \sim 43\%/\sqrt{E} \oplus 2.4\%$



Conclusions

Studies of accelerators and detectors for the post-LHC energy frontier are on-going.

A conceptual design report is planned for 2018.

Basic concepts for detectors at these future colliders are being worked on and have been shown.

Detector technology choices will depend on the requirements from physics – further refinement under way.

Concentrating on few example designs while staying open for innovative concepts.

New ideas and person-power are highly welcome!

Join in!

Backup

LUMINOSITY GOALS FOR A 100-TeV PP COLLIDER

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Institut de Physique Théorique Philippe Meyer, École Normale Supérieure
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April 24, 2015

Abstract

We consider diverse examples of science goals that provide a framework to assess luminosity goals for a future 100-TeV proton-proton collider.

An integrated luminosity goal of 20ab^{-1}
matches very well the 100TeV c.m. Energy

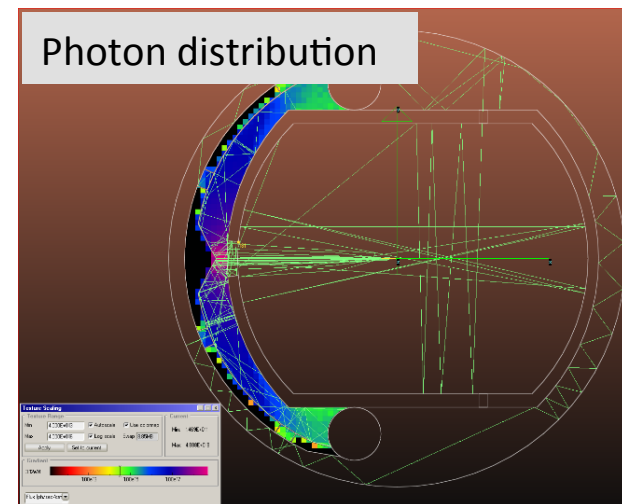
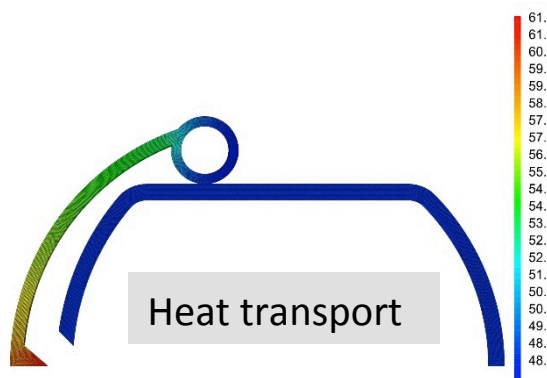
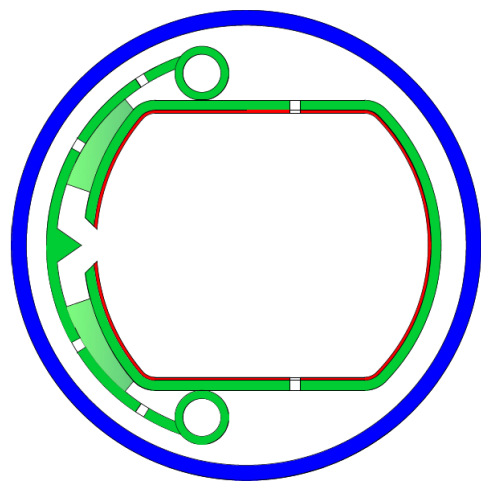
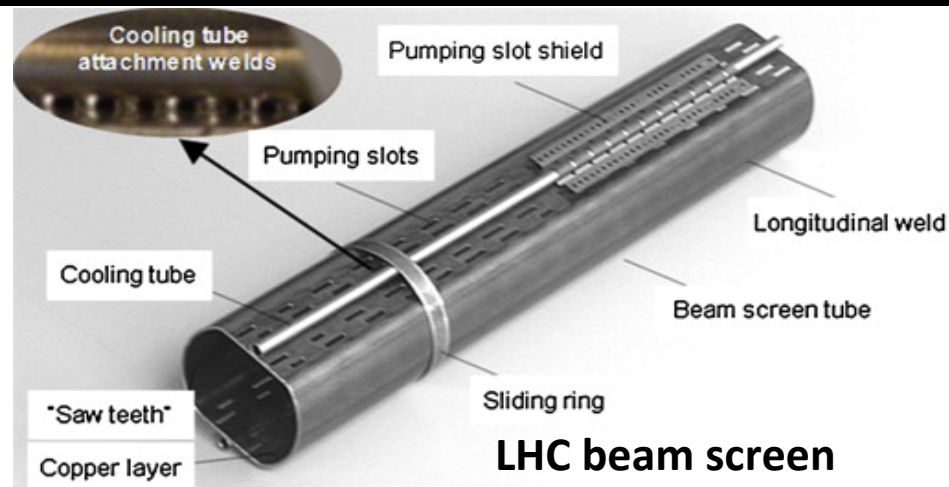
High synchrotron radiation load (SR) of protons @ 50 TeV:

~30 W/m/beam @16 T (LHC <0.2W/m)

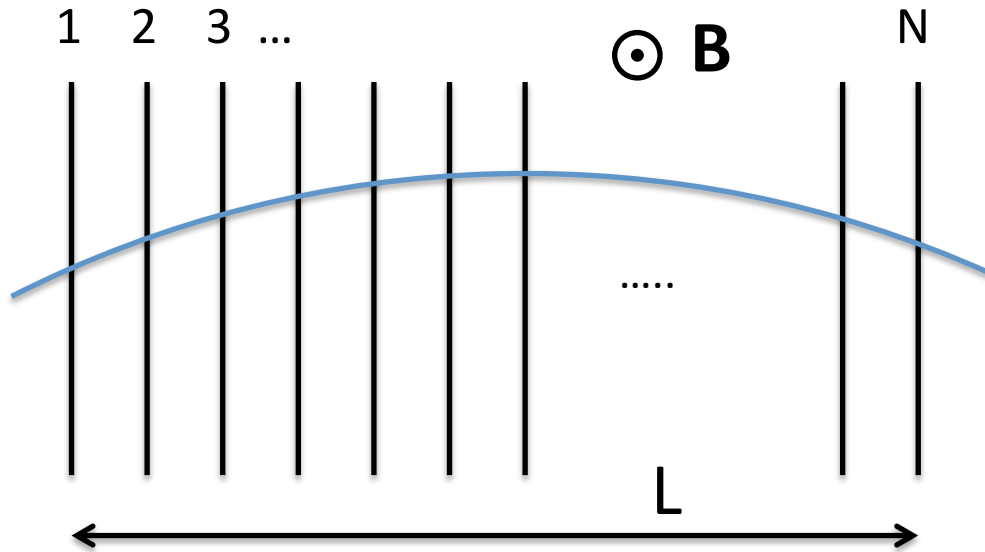
→ **5 MW total in arcs**

New type of chamber

- absorption of synchrotron radiation
- avoids photo-electrons, helps vacuum



Point Resolution and Multiple Scattering



σ ... point resolution/plane

X_{tot}/X_0 ... total material budget

N equidistant layers

Position Resolution

$$\frac{\Delta p_t}{p_t} = \frac{\sigma [m] p [\text{GeV}/c]}{0.3 B [T] L^2 [m^2]} \sqrt{\frac{720 (N - 1)^3}{(N - 2) N (N + 1) (N + 2)}}$$

$$\approx \frac{\sigma [m] p [\text{GeV}/c]}{0.3 B [T] L^2 [m^2]} \sqrt{\frac{720}{N + 4}}$$

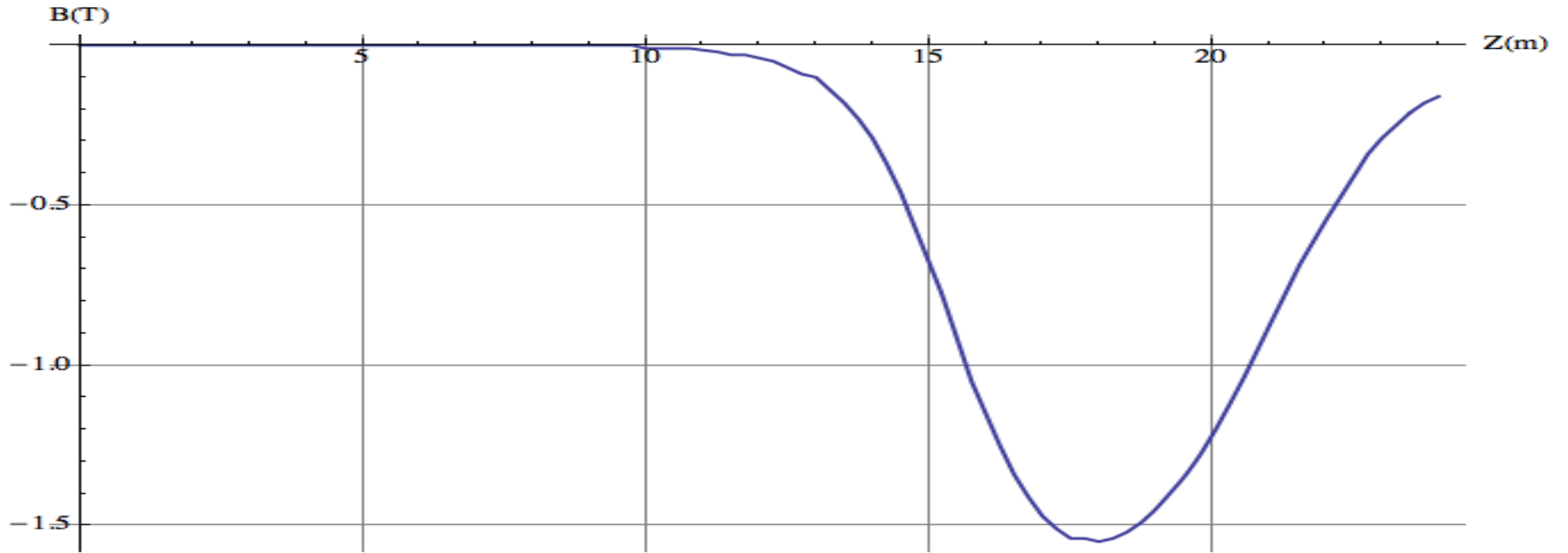
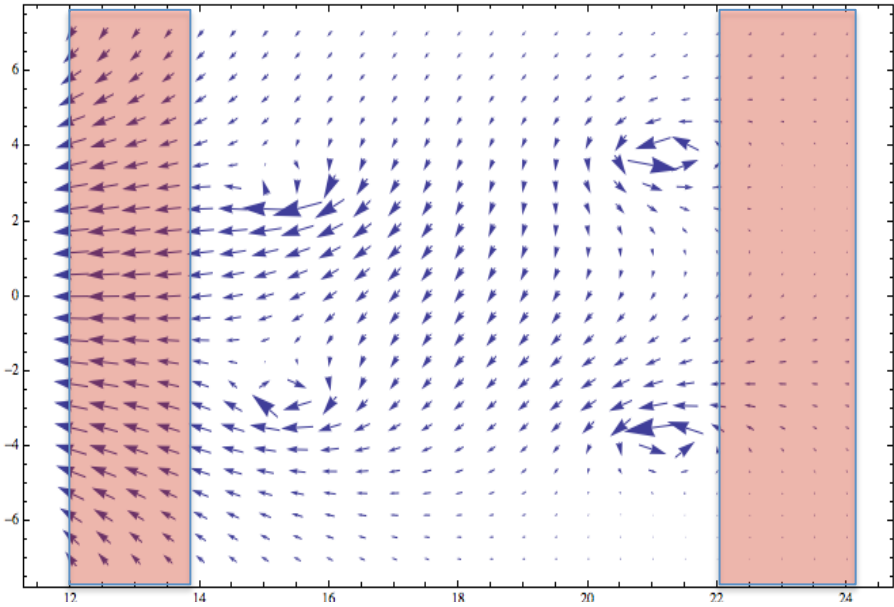
**Multiple Scattering,
Equal weighting**

$$\frac{\Delta p_t}{p_t} = \frac{0.0136}{0.3 \beta B [T] L [m]} \sqrt{\frac{X_{\text{tot}}}{X_0}} \sqrt{\frac{10}{7} \frac{12 + (N - 1) N^2 (N + 1)}{(N - 2) N (N + 1) (N + 2)}}$$

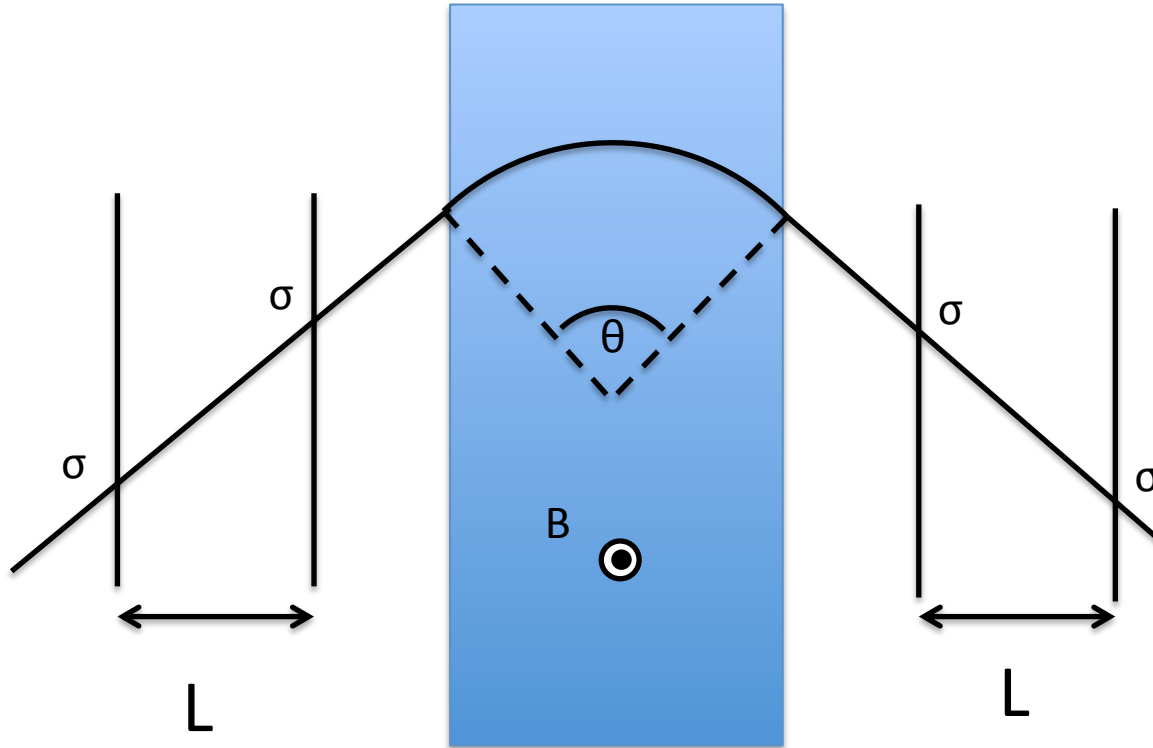
$$\approx \frac{0.0136}{0.3 \beta B [T] L [m]} \sqrt{\frac{X_{\text{tot}}}{X_0}} \sqrt{\frac{10}{7}} \quad \left. \vphantom{\frac{\Delta p_t}{p_t}} \right\} \approx 1.2$$

$$\approx \frac{0.0542}{\beta B [T] L [m]} \sqrt{\frac{X_{\text{tot}}}{X_0}}$$

Taking into account the correlation (Kalman filter etc.) this number can be reduced to 1.0⁶³

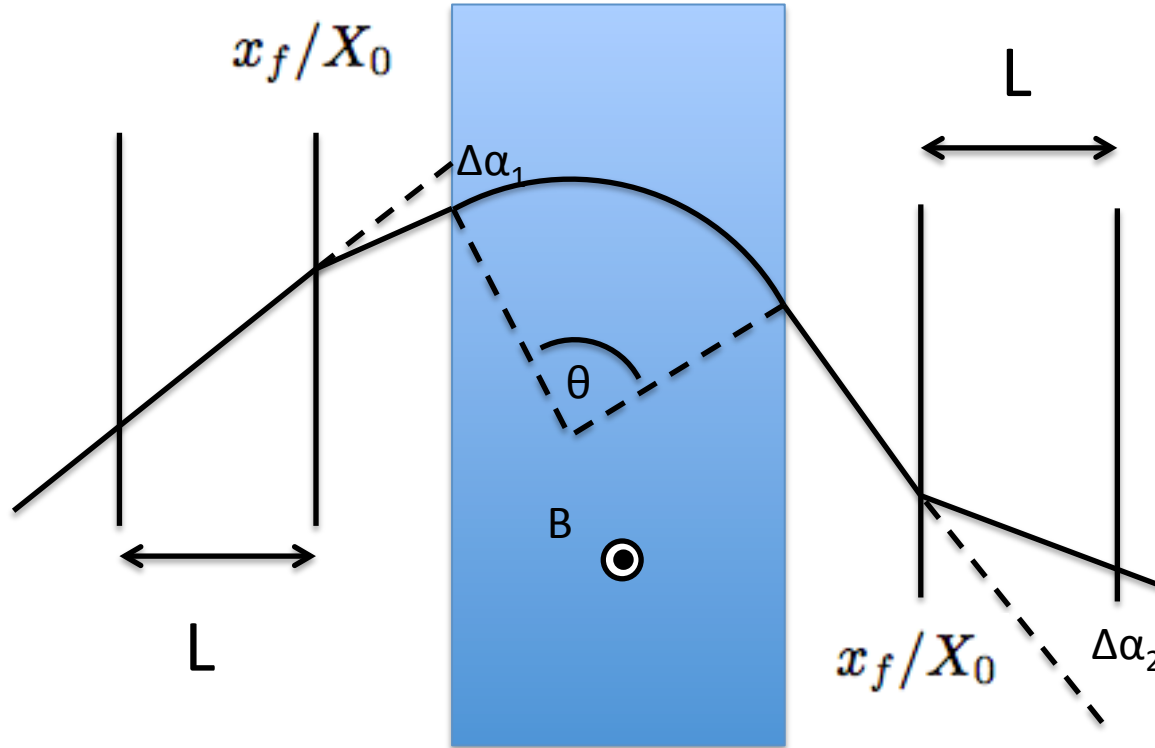


Forward Tracker Resolution, Position Resolution



$$\left(\frac{\Delta p_L}{p_L}\right)^2 = \left(\frac{p_L}{0.3 \int B_T dl}\right)^2 (\Delta\alpha_{res}^2 + \Delta\alpha_{ms}^2) \quad \Delta\alpha_{res} = 2\sigma/L.$$

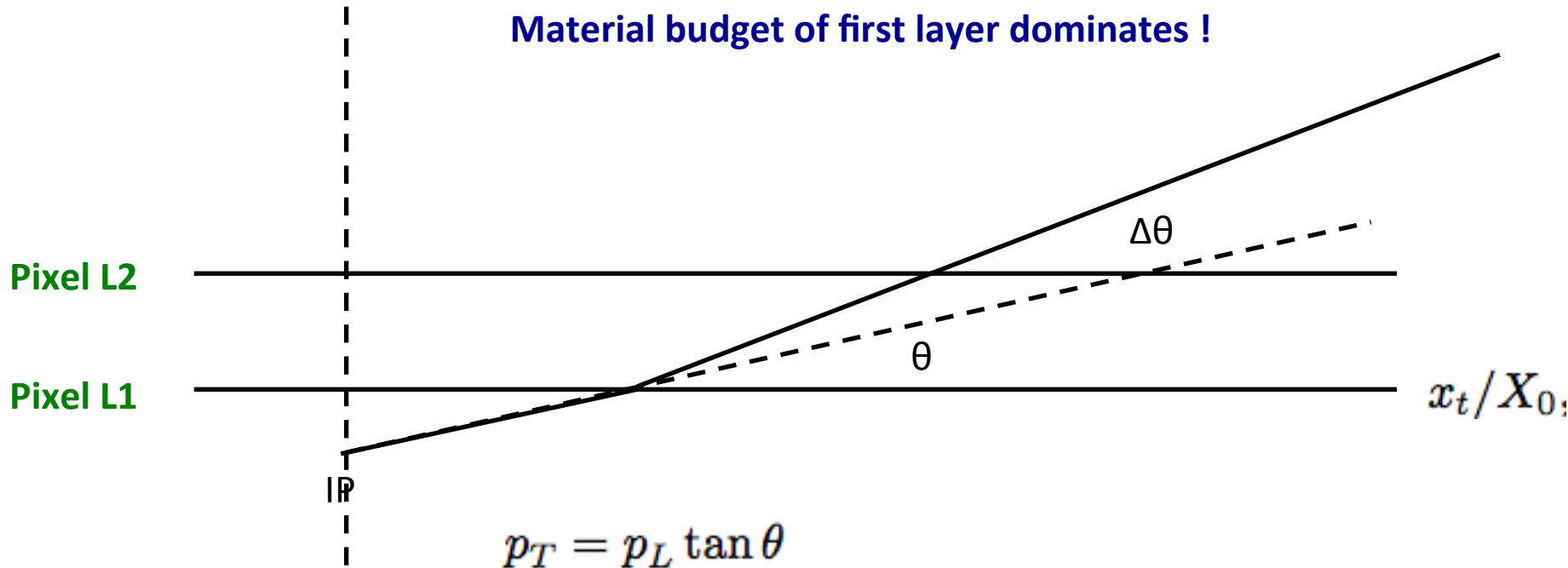
Forward Tracker Resolution, Multiple Scattering



$$\left(\frac{\Delta p_L}{p_L}\right)^2 = \left(\frac{p_L}{0.3 \int B_T dl}\right)^2 (\Delta\alpha_{res}^2 + \Delta\alpha_{ms}^2) \quad \Delta\alpha_{ms} = 0.0136/p_L \sqrt{2x_f/X_0}$$

Forward Tracker Resolution, Measurement of angle

Material budget of first layer dominates !



$$\left(\frac{\Delta p_T}{p_T}\right)^2 = \left(\frac{\Delta p_L}{p_L}\right)^2 + \left(\frac{1}{\sin \theta \cos \theta}\right)^2 \Delta\theta^2$$

$$\Delta\theta = \frac{0.0136}{p} \sqrt{\frac{x_t}{X_0} \frac{1}{\sin \theta}} = \frac{0.0136 \sin \theta}{p_T} \sqrt{\frac{x_t}{X_0} \frac{1}{\sin \theta}}$$

Forward Tracker Resolution

$$\left(\frac{\Delta p_T}{p_T}\right)^2 = \left(\frac{2\sigma p_T}{\tan\theta \cdot 0.3L \int B_T dl}\right)^2 + \left(\frac{0.0136}{0.3 \int B_T dl} \sqrt{2\frac{x_f}{X_0}}\right)^2 + \left(\frac{0.0136}{p_T \cos\theta} \sqrt{\frac{x_t}{X_0} \frac{1}{\sin\theta}}\right)^2$$

$$\left(\frac{\Delta p_T}{p_T}\right)^2 = \left(\frac{2\sigma p_T \sinh\eta}{0.3L \int B_T dl}\right)^2 + \left(\frac{0.0136}{0.3 \int B_T dl} \sqrt{2\frac{x_f}{X_0}}\right)^2 + \left(\frac{0.0136 \coth\eta}{p_T} \sqrt{\frac{x_t}{X_0} \cosh\eta}\right)^2$$

$$\sigma = 30\mu\text{m}$$

$$\text{Int Bdl} = 10 \text{ Tm}$$

$$L = 2\text{m}$$

$$x_f/X_0 = 0.06$$

$$\text{Int Bdl} = 10 \text{ Tm}$$

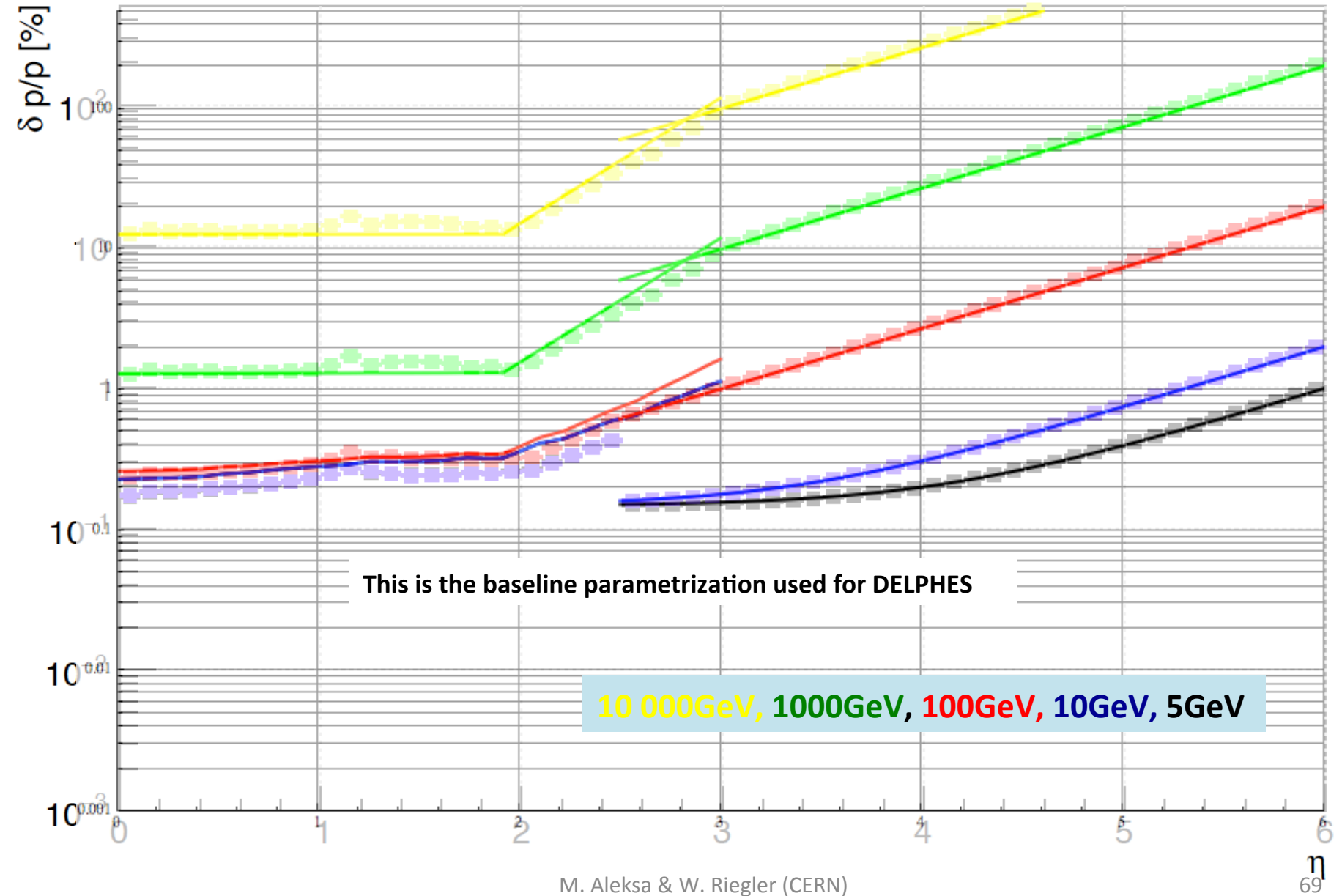
$$x_t/X_0 = 0.03$$

Using $L = 2 \text{ m}$, $\sigma = 30 \mu\text{m}$, $x_f/X_0 = 0.06$, $x_t/X_0 = 0.03$

$$\frac{\Delta p_T}{p_T} = 10^{-3} \sqrt{1.5^2 + (10^{-2} p_T \sinh\eta)^2 + \left(2.4 \frac{\coth\eta}{p_T} \sqrt{\cosh\eta}\right)^2}$$

Forward Tracker Resolution

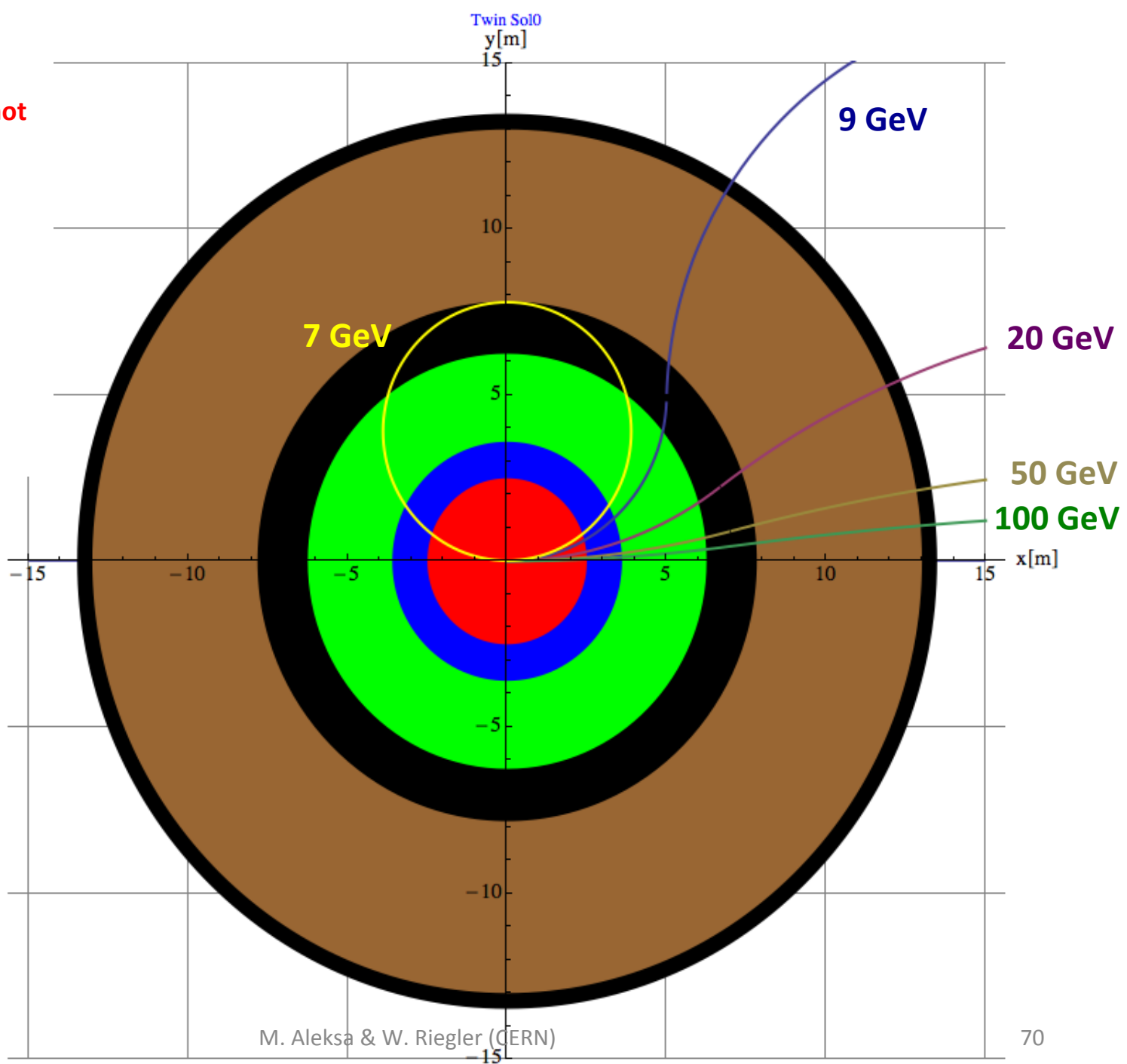
Total (Dipole+Central) p_T resolution versus η - const P_T across η



At $B_0=6T$ and $R_0=6m$,
Muons below 7GeV do not
enter the muon system.

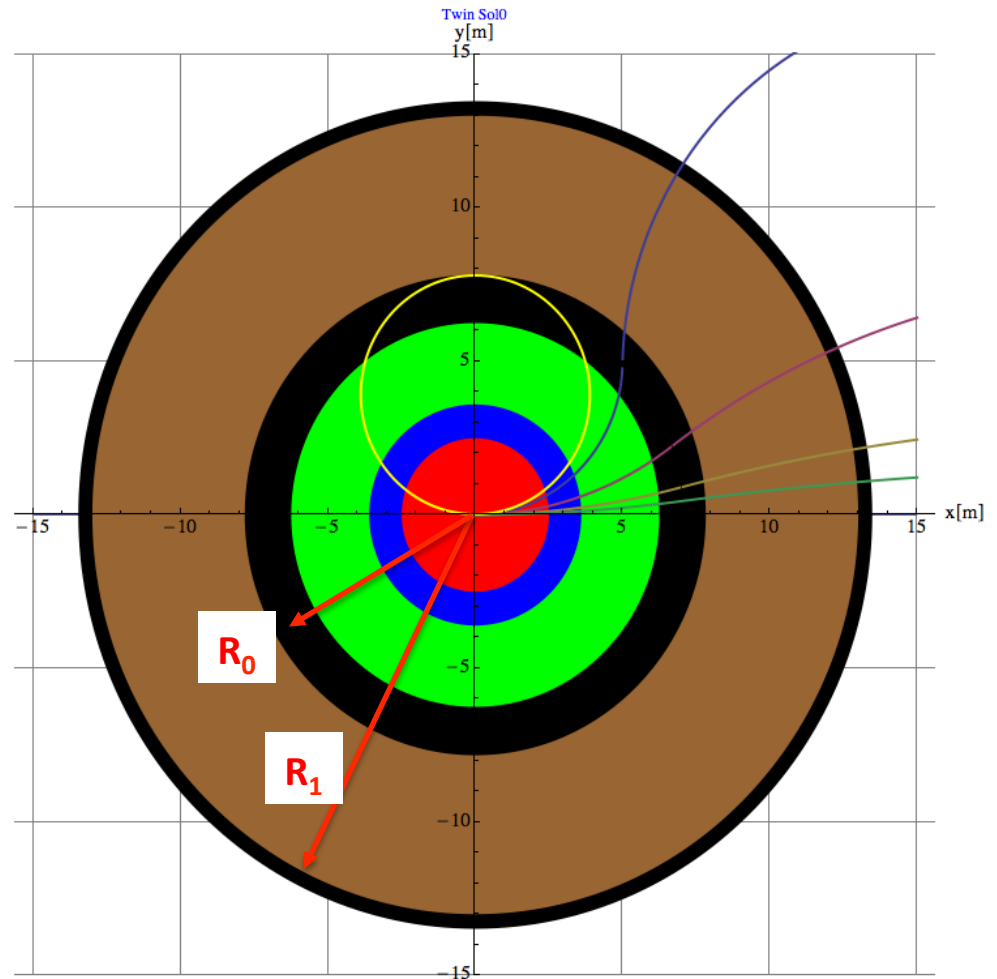
No Muon Trigger below
7GeV.

Possibly muon ID with
HCAL.



Muon Momentum can be measured by

- 1) The inner tracker
→ resolution plots from before
- 2) The track angle at the entrance of the muon system → Trigger
- 3) A sagitta measurement in the muon system (no iron → precise !)
- 4) The combined fit of inner tracker and outer layers of the muon system.



2) Track angle at the entrance of the muon system

$$\frac{\Delta p_T}{p_T} = \Delta\theta \sqrt{\left(\frac{2p_T}{0.3B_0R_0}\right)^2 - 1} \approx \frac{2p_T}{0.3B_0R_0} \Delta\theta \quad \text{for a large } p_T$$

10% at 10TeV, $B_0=6\text{T}$, $R_0=6\text{m}$
 $\Delta\theta=50\mu\text{Rad}$

→ 2 stations at 1.5m distance with 50um position resolution

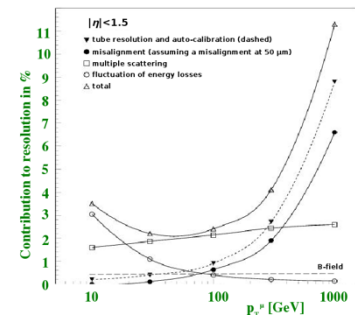
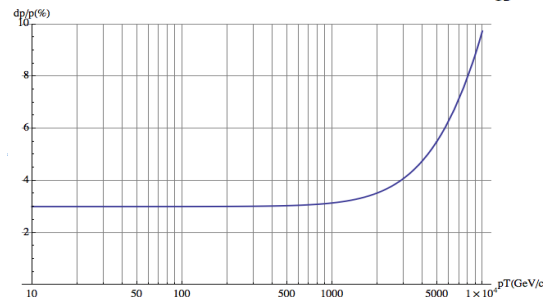
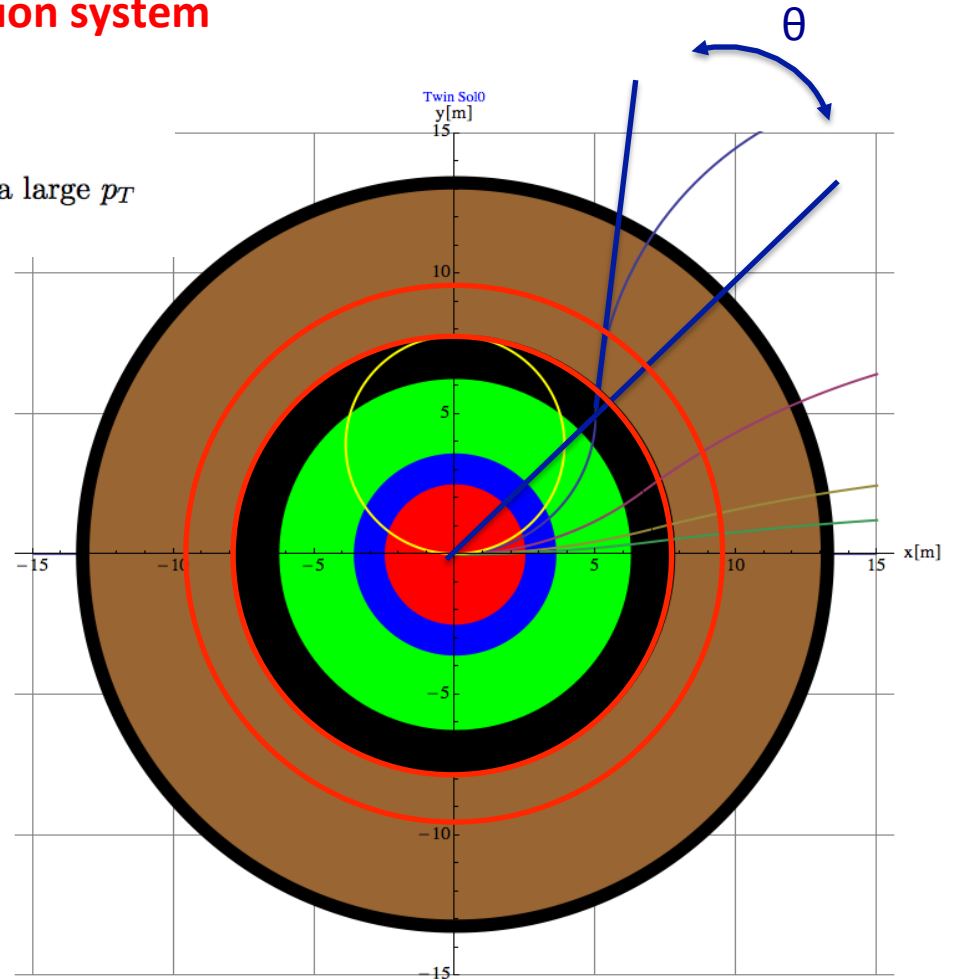
For low momentum, limit due to multiple scattering in the calorimeters and coil:

Calorimeter+Cryostat: $35X_0$
 HCAL: $110X_0$
 Coil: $5X_0$
 → $x_{\text{tot}}/X_0 \approx 150$

$$\frac{\Delta p_T}{p_T} = \frac{2 \times 0.0136}{0.3B_0R_0} \sqrt{\frac{x_{\text{tot}}}{X_0}}$$

$B_0=6\text{T}$, $R_0=6\text{m}$
 → $dp/p=3\% !!!$

(CMS 9% because $B_0R_0=1/3$)



M. Aleksa & W. Riegler (CERN)

At eta=0 ATLAS type standalone Muon Performance up to 10TeV !!!

3) Sagitta measurement in the muon system

The return field is 2.45T

Measuring over the 5m lever arm with stations of $\sigma=50\mu\text{m}$ resolution we have

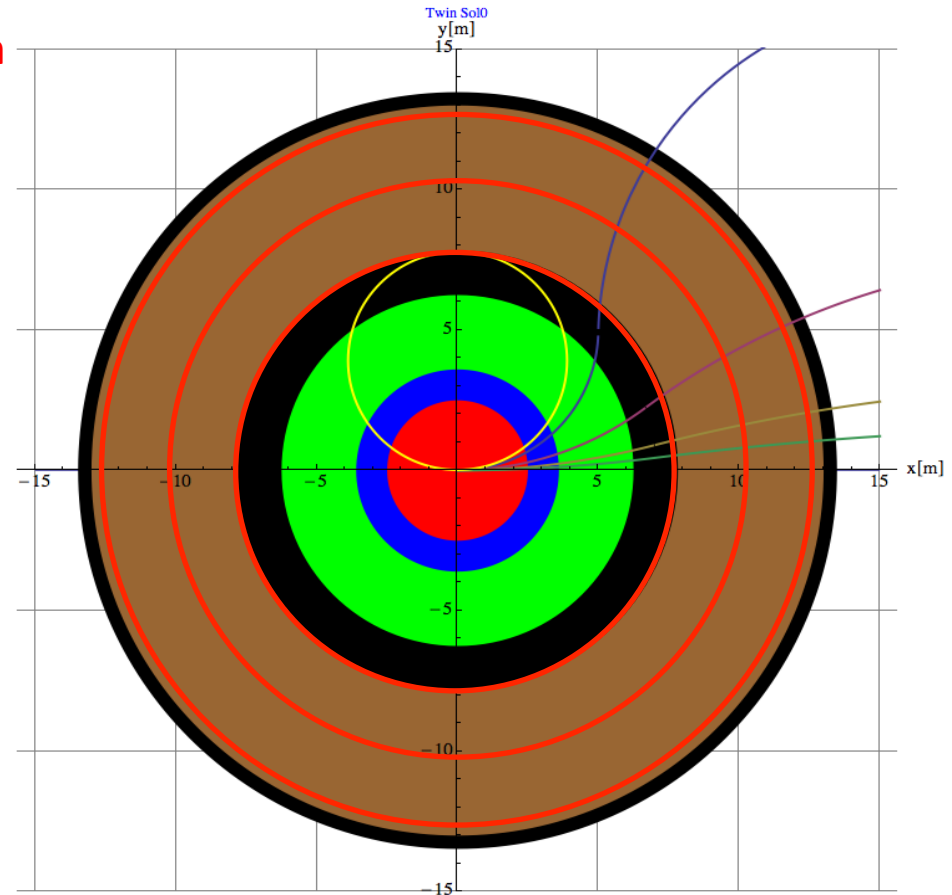
$$\frac{dp_T}{p_T} = \frac{\sigma * p_T}{(0.3 * B * L^2) * 8} \\ = 20\% @ 10\text{TeV}$$

with possibly excellent performance at low p_T due to the absence of iron (vs. CMS) .

but very hard to beat the angular measurement at high p_T and the inner tracker at low p_T .

Surface > 5000 m²

CMS sagitta measurement in the muon system is limited to $dp_T/p_T = 20\%$ due to multiple scattering alone.



Combined Measurement

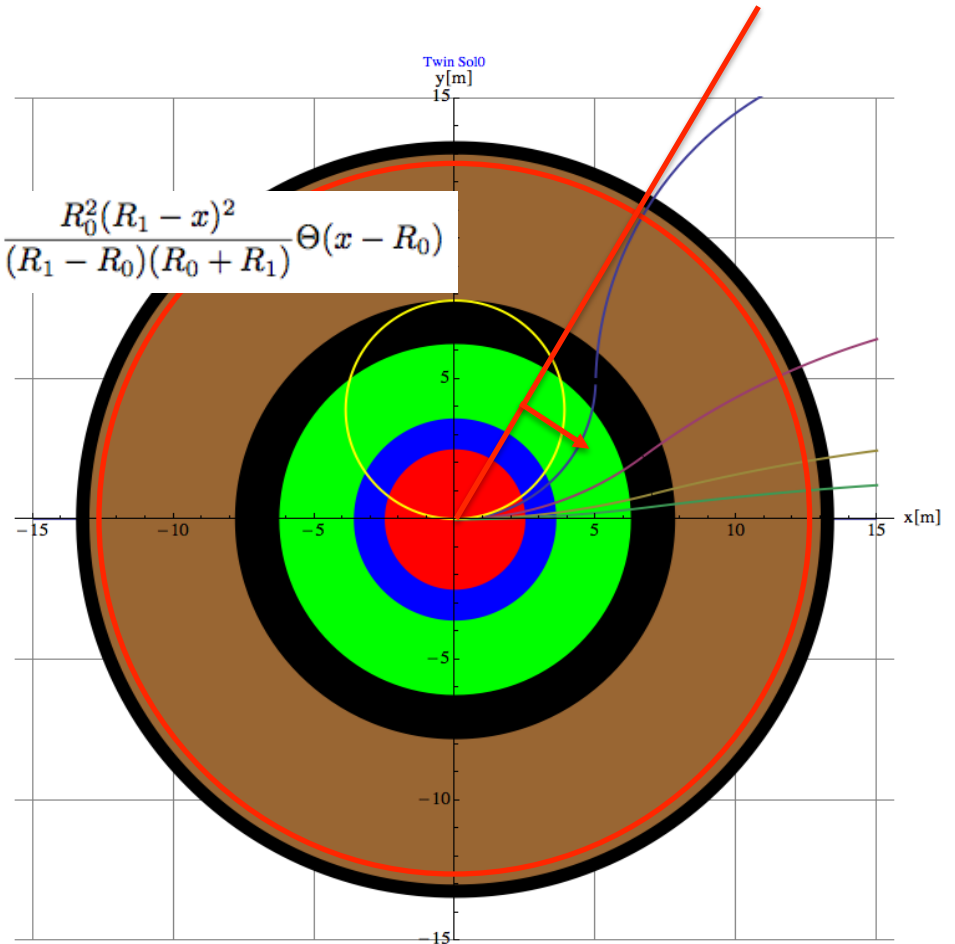
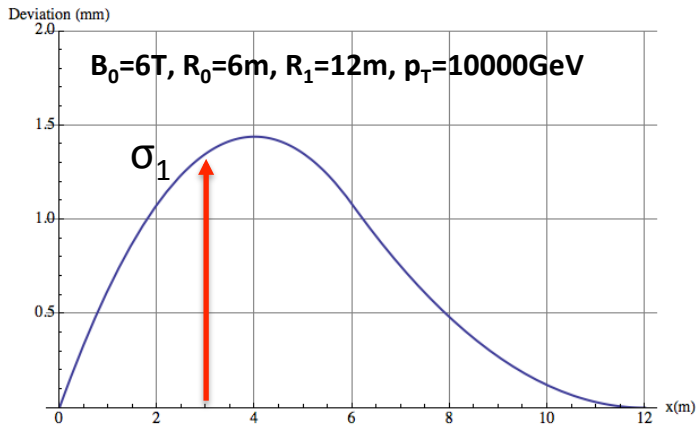
If the full flux is returned through the muon system, the muon trajectory at the exit of the system points exactly to the IP !

$$y_t(x) = \frac{0.3B_0}{2p_T} \left(x^2 - \frac{2R_0R_1}{R_0 + R_1}x \right) \Theta(R_0 - x) - \frac{0.3B_0}{2p_T} \frac{R_0^2(R_1 - x)^2}{(R_1 - R_0)(R_0 + R_1)} \Theta(x - R_0)$$

The maximum excursion $y_t(x_0)$ is always at the same radial distance of x_0

$$x_0 = \frac{R_0R_1}{R_0 + R_1} \quad y_t(x_0) = -\frac{0.3B_0}{2p_T} x_0^2 = -\frac{0.3B_0}{2p_T} \left(\frac{R_0R_1}{R_0 + R_1} \right)^2$$

For values below: $x_0=4\text{m}$, $y_t(x_0)=1.44\text{mm}$
 Ideal measurement point is at the peak,
 but $y_t(2.4\text{m})= 1.24\text{mm}$ still good !

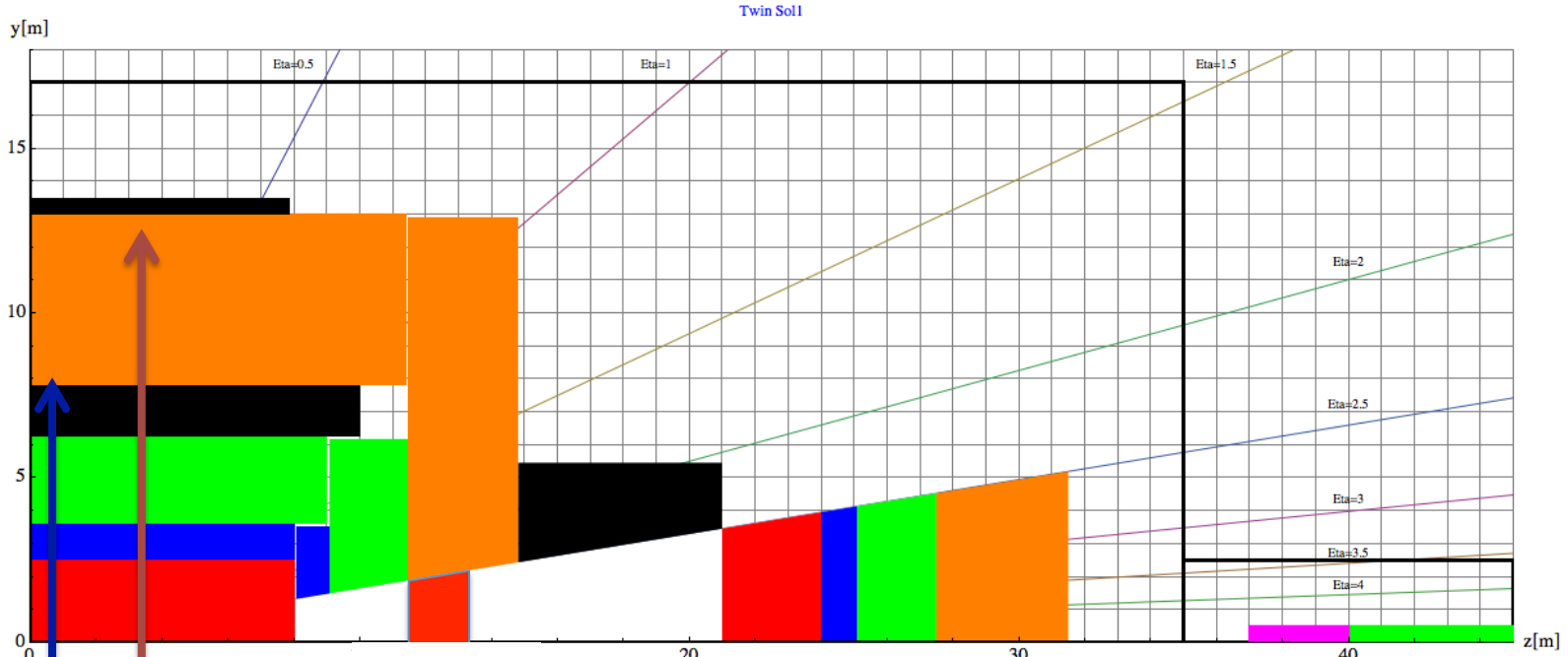


$$\sigma^2 = \sigma_1^2 + (x/R_1\sigma_2)^2$$

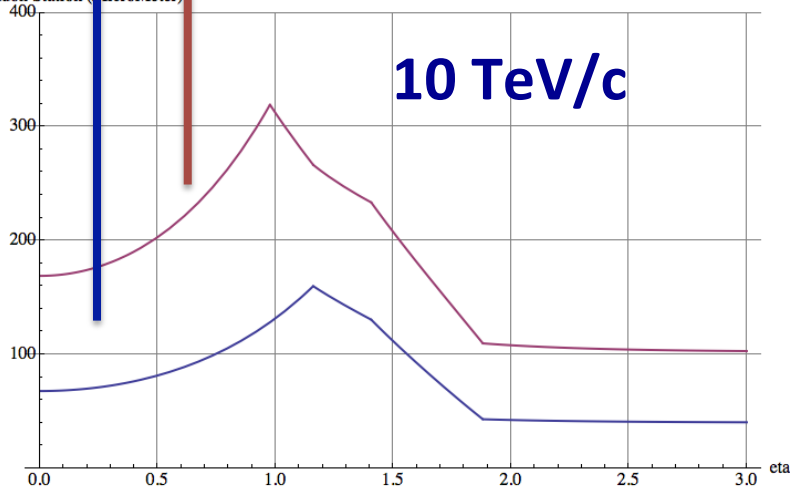
$x=2.4\text{m}, R_1=12\text{m}, \sigma_1=50\mu\text{m}, \sigma_2=250\mu\text{m},$
 $\sigma=64\mu\text{m}, dp_T/p_T=5\%$ at 10TeV !

Measuring just in the last tracker layer and in the outermost muon station already beats the full inner tracker performance (14 layers, 23um).

Measurement error in the muon system due to multiple scattering



MS Error at Muon Station (MicroMeter)



The r.m.s. (in plane) angular deflection θ_0 and spatial displacement σ_y of a particle passing through material of thickness x are given by

$$\theta_0 = \frac{0.0136}{\beta p [\text{GeV}/c]} \sqrt{\frac{x}{X_0}} \left(1 + 0.038 \log \frac{x}{X_0} \right) \quad (184)$$

$$\sigma_y = \frac{1}{\sqrt{3}} x \theta_0 \quad (185)$$

For the proper correlated M.C. generation of these two quantities, one uses two independent Gaussian random variables z_1, z_2 with mean of zero and sigma of one, and calculates

$$y = z_1 x \frac{\theta_0}{\sqrt{12}} + z_2 x \frac{\theta_0}{2} \quad (186)$$

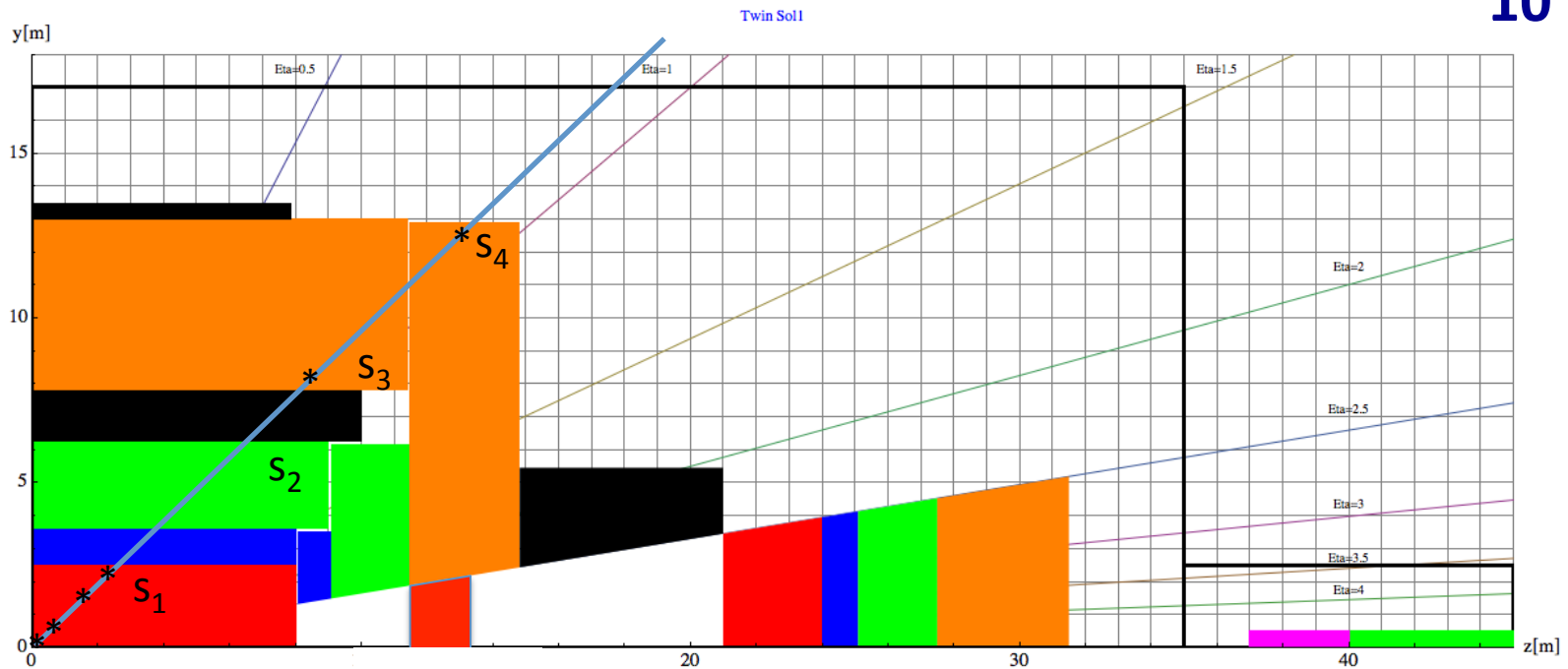
$$\theta = z_2 \theta_0 \quad (187)$$

In case a muon passes the calorimeter of thickness L_1 and then propagates through the (material less) muons system of length L_2 the spatial deflection is given by

$$s = z_1 L_1 \frac{\theta_0}{\sqrt{12}} + z_2 L_1 \frac{\theta_0}{2} + z_2 \theta_0 L_2 \quad (188)$$

and therefore a position error of

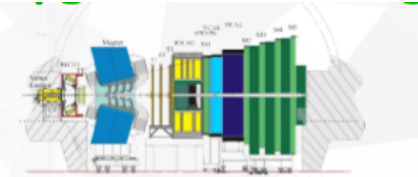
$$\sigma_s^2 = \theta_0^2 \left[\frac{L_1^2}{12} + \left(\frac{L_1}{2} + L_2 \right)^2 \right] \quad (189)$$



Fit of the template function to 14 measurement points of 23um in the inner tracker and 1 layer in the muon system with a resolution that is much better than the multiple scattering limit.

i.e. assuming the resolution at s3 to be better than 50um and at s4 to be better than 100um and the resolution of the angular measurement better than 20uRad.

LHCb & ALICE in 2018



40 MHz

DAQ

40 MHz

LLT: $\mu, e/\gamma$,
hadrons

5-40 MHz

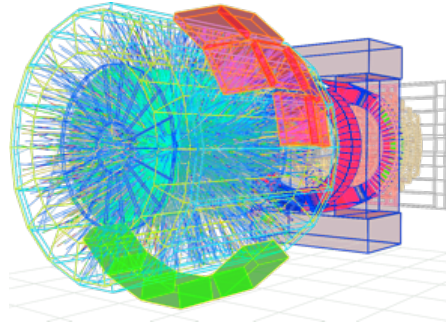
HLT: event
recon.

20 kHz (0.1 MB/event)

Storage

2 GB/s

4 TByte/s into PC farm for HLT selection.



50 kHz

Reconstruction
+
Compression



50 kHz (1.5 MB/event)

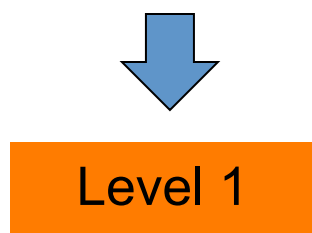
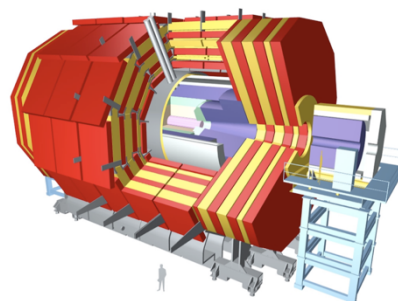
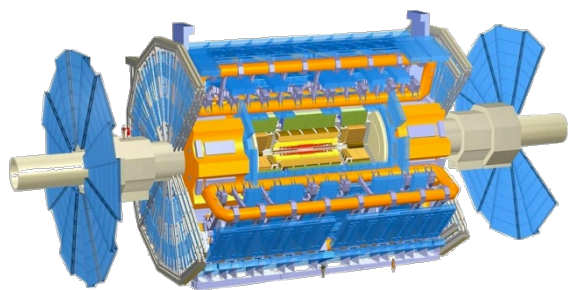
Storage

1 TByte/s into PC farm for data compression. All events to disc.

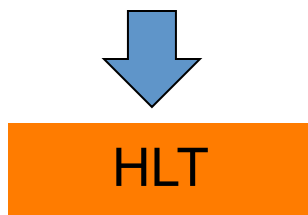
← PEAK OUTPUT →

75 GB/s

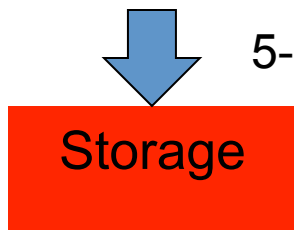
ATLAS & CMS in 2018



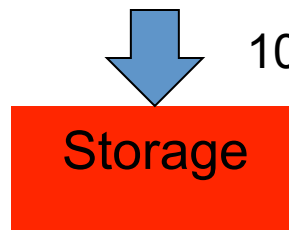
40 MHz



0.5-1 MHz



5-10 kHz (2MB/event)



10 kHz (4MB/event)

5 TByte/s into PC farm for HLT selection.

Would be 200TByte/s without Level1

10-20 GB/s

← PEAK OUTPUT →

40 GB/s

Hardware Trigger ?

CMS HL-LHC results in 200TByte/s into the online system for a triggerless readout.

For 2022 this is considered too difficult.

Assuming that the total track rate for 100TeV pp collisions (Phase I) is only a factor 2 larger, there is very little doubt that by 2035 and FCC-hh detector can be read out in a triggerless fashion.

In 2035 maybe no hardware trigger necessary ! All data to the online system, synchronous or asynchronous, where a sophisticated selection and compression can be done.

N.b. the techniques to get the data out of the detector with a small amount of material is a key question to be solved.

Even if one would afford to read all data to HLT for Phase II, the amount of copper lines to get all the signals out of the silicon detector would destroy the tracker performance.

Moore's Law

<http://www.livescience.com/23074-future-computers.html>

“If the doubling of computing power every two years continues to hold, then by 2030 whatever technology we're using will be sufficiently small that we can fit all the computing power that's in a human brain into a physical volume the size of a brain”,

explained Peter Denning, distinguished professor of computer science at the Naval Postgraduate School and an expert on innovation in computing.

"Futurists believe that's what you need for artificial intelligence. At that point, the computer starts thinking for itself."

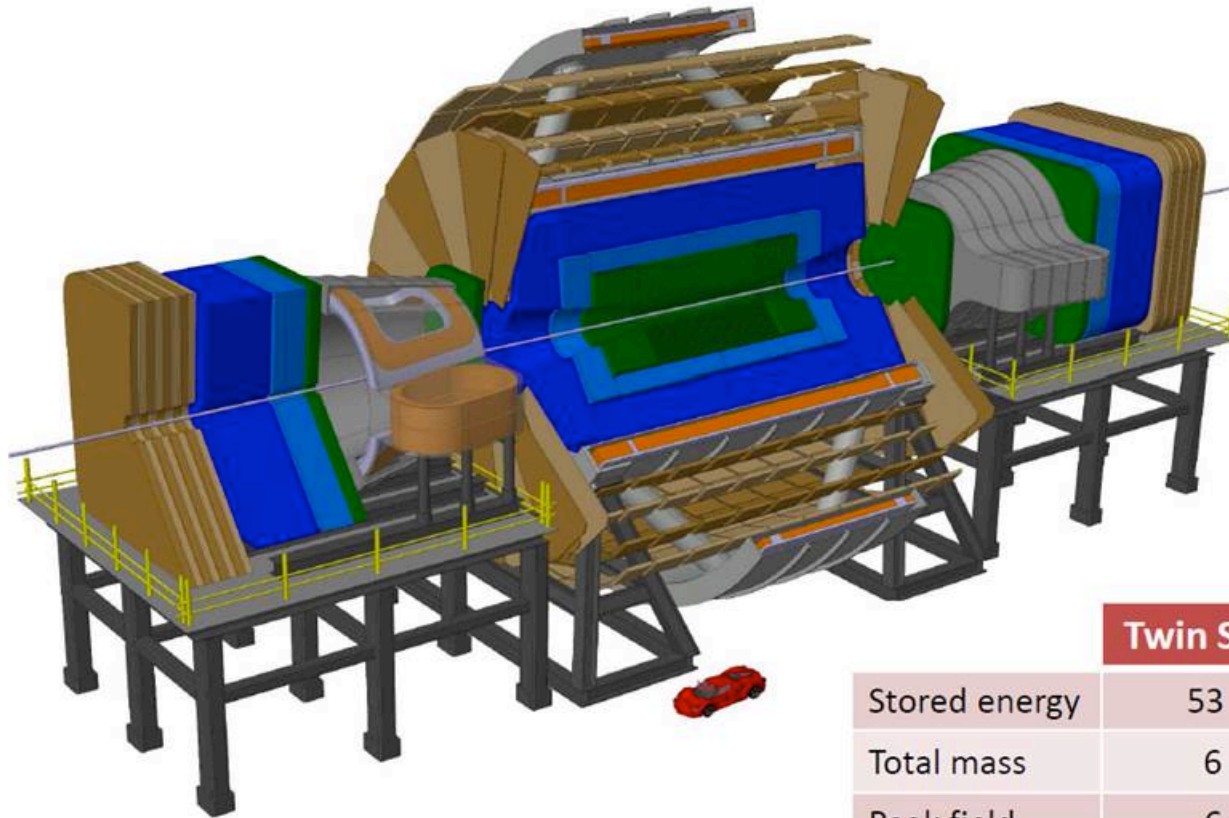
→ Computers will anyway by themselves figure out what to do with the data by 2035.

Magnet systems and shielding will be rather conventional and can be worked out to some detail now.

For detector technology and computing power we are allowed to dream a bit.

Twin Solenoid + Dipole Magnet System

Matthias Mentink, Alexey Dudarev, Helder Filipe Pais Da Silva, Christophe Paul Berriaud, Gabriella Rolando, Rosalinde Pots, Benoit Cure, Andrea Gaddi, Vyacheslav Klyukhin, Hubert Gerwig, Udo Wagner, and Herman ten Kate



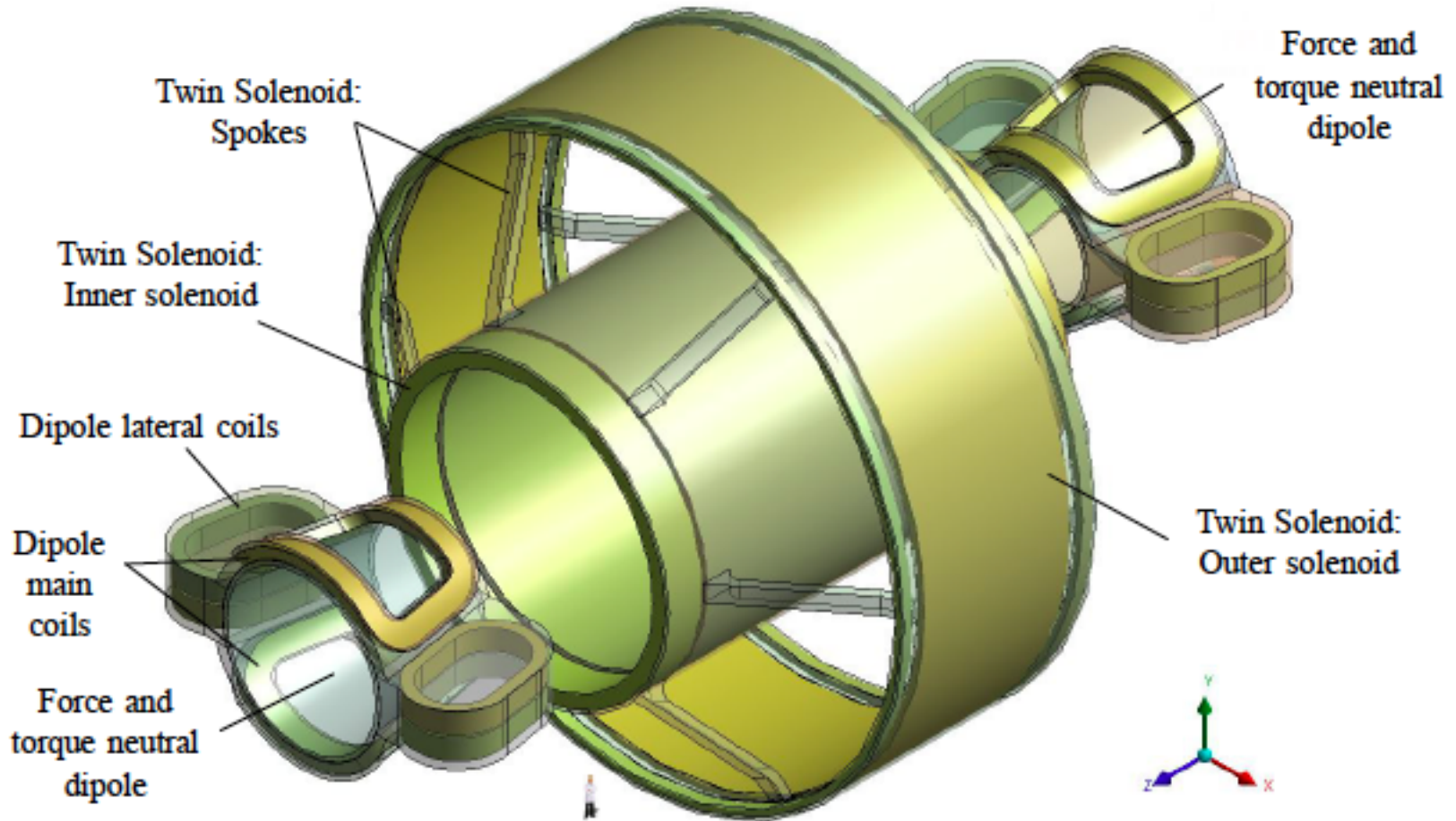
FCC Air core Twin solenoid and Dipoles

State of the art high stress / low mass design.

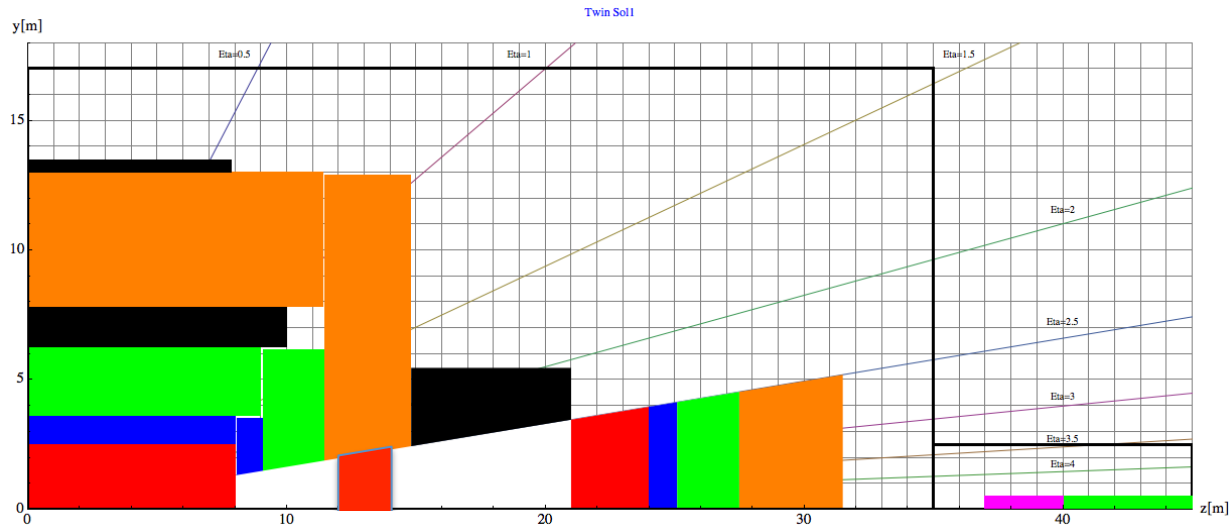
	Twin Solenoid	Dipole
Stored energy	53 GJ	2 x 1.5 GJ
Total mass	6 kt	0.5 kt
Peak field	6.5 T	6.0 T
Current	80 kA	20 kA
Conductor	102 km	2 x 37 km
Bore x Length	12 m x 20 m	6 m x 6 m

Twin Solenoid + Dipole Magnet System

Matthias Mentink, Alexey Dudarev, Helder Filipe Pais Da Silva, Christophe Paul Berriaud, Gabriella Rolando, Rosalinde Pots, Benoit Cure, Andrea Gaddi, Vyacheslav Klyukhin, Hubert Gerwig, Udo Wagner, and Herman ten Kate



Baseline Geometry, Twin Solenoid



Barrel:

Tracker available space:
 $R=2.1\text{cm to } R=2.5\text{m}, L=8\text{m}$

EMCAL available space:
 $R=2.5\text{m to } R=3.6\text{m} \rightarrow dR=1.1\text{m}$

HCAL available space:
 $R=3.6\text{m to } R=6.0\text{m} \rightarrow dR=2.4\text{m}$

Coil+Cryostat:
 $R=6\text{m to } R=7.825 \rightarrow dR=1.575\text{m}, L=10.1\text{m}$

Muon available space:
 $R=7.825\text{m to } R=13\text{m} \rightarrow dR=5.175\text{m}$

Coil2:
 $R=13\text{m to } R=13.47\text{m} \rightarrow dR=0.475\text{m}, L=7.6\text{m}$

Endcap:

EMCAL available space:
 $z=8\text{m to } z=9.1\text{m} \rightarrow dz=1.1\text{m}$

HCAL available space:
 $z=9.1\text{m to } z=11.5\text{m} \rightarrow dz=2.4\text{m}$

Muon available space:
 $z=11.5\text{m to } z=14.8\text{m} \rightarrow dz=3.3\text{m}$

Forward:

Dipole:
 $z=14.8\text{m to } z=21\text{m} \rightarrow dz=6.2\text{m}$

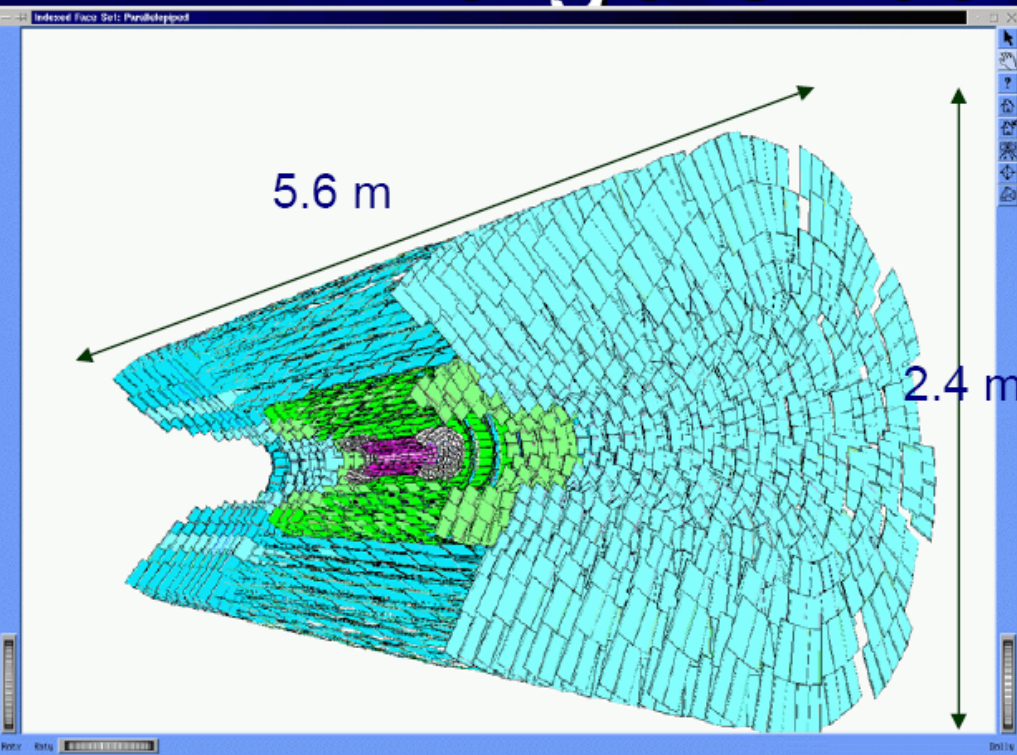
FTracker available space:
 $z=21\text{m to } R=24\text{m}, L=3\text{m}$

FEMCAL available space:
 $Z=24\text{m to } z=25.1\text{m} \rightarrow dz=1.1\text{m}$

FHCAL available space:
 $z=25.1\text{m to } z=27.5\text{m} \rightarrow dz=2.4\text{m}$

FMuon available space:
 $z=27.5\text{m to } z=31.5\text{m} \rightarrow dz=4\text{m}$

Large Silicon Systems



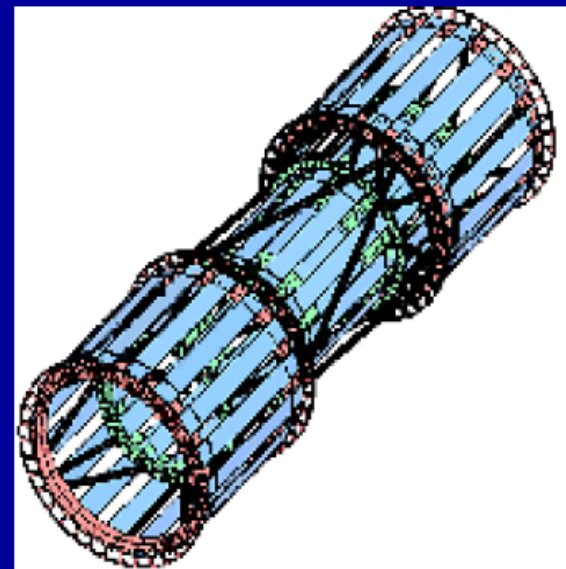
CMS tracker (~2007)

12000 modules

~ 445 m² silicon area

~ 24,328 silicon wafers

~ 60 M readout channels



CDF SVX IIa (2001-)

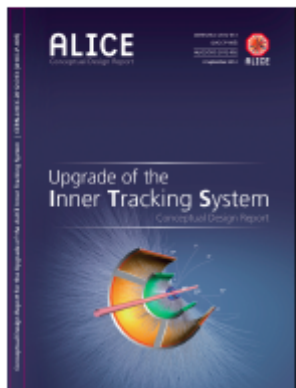
~ 11m² silicon area

~ 750 000 readout channels

ALICE 2018 upgrade, 20x20um monolithic pixels

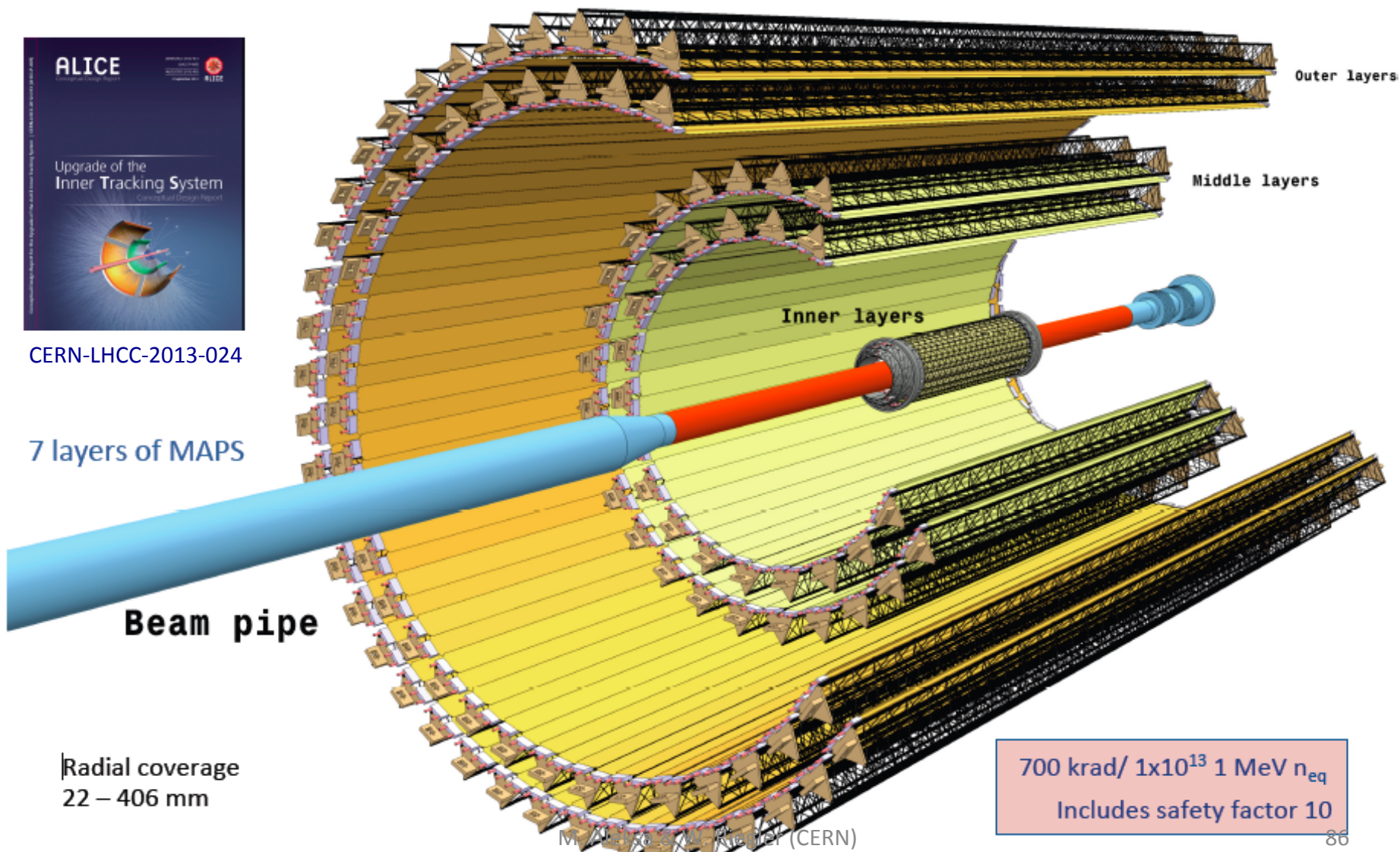
New ITS Layout

25 G-pixel camera
(10.3 m²)



CERN-LHCC-2013-024

7 layers of MAPS



Radial coverage
22 – 406 mm

700 krad/ 1×10^{13} 1 MeV n_{eq}
Includes safety factor 10

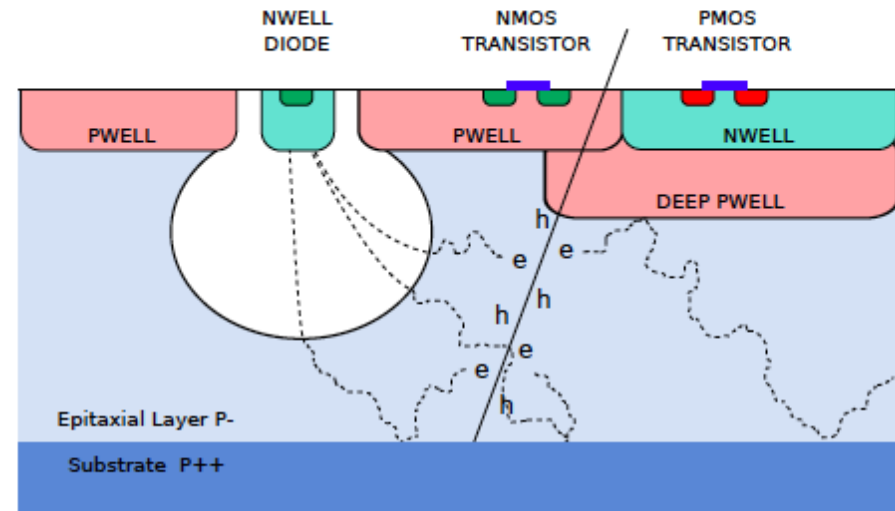
PIXEL Chip - technology

Monolithic PIXEL chip using Tower/Jazz 0.18 μm technology

- feature size 180 nm
- gate oxide < 4nm
- metal layers 6
- high resistivity epi-layer
 - thickness 18-40 μm
 - resistivity 1-6 k $\Omega \times \text{cm}$
- “special” deep p-well layer to shield PMOS transistors (allows in-pixel truly CMOS circuitry)
- Possibility to build single-die circuit larger than reticle size

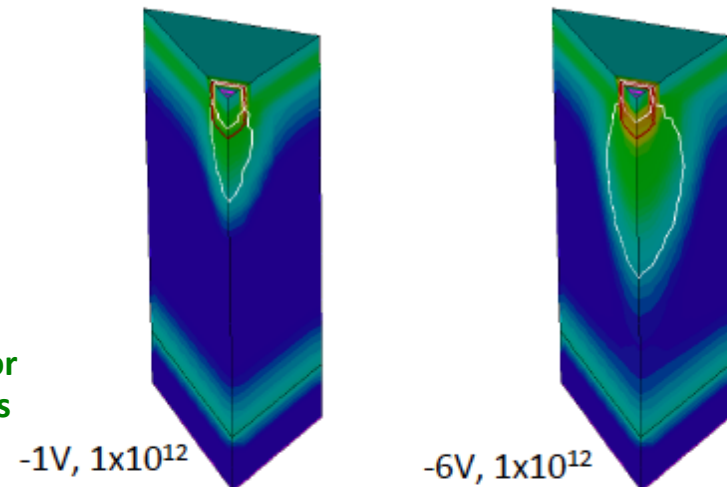
Standard processing, no bump bonding (>>50% of Pixel detector cost). Allows implementation of complex processing electronics inside the entire pixel area.

→ Revolution !



Schematic cross-section of CMOS pixel sensor (ALICE ITS Upgrade TDR)

TCAD simulation of total diode reverse bias (ALICE ITS Upgrade TDR)



Pixel Revolution Hybrid → Monolythic

Table 2.2: Chip design options.

Architecture (discriminator, read-out)	Pitch ($r\phi \times z$) (μm^2)	Integration time (ps)	Power consumption (mW cm^{-2})
MISTRAL (end-of-column, rolling-shutter)	22×33.3	30	200
ASTRAL (in-pixel, rolling-shutter)	24×31	20	85
CHERWELL (in-striixel ^a , rolling-shutter)	36×31	60	60
ALPIDE (in-pixel, in-matrix sparsification)	20×20	30	90
ALPIDE (in-pixel, in-matrix sparsification)	28×28	4	< 50

^a A striixel is a 128-pixel column over which the electronics are distributed.

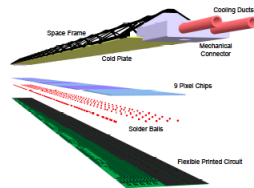
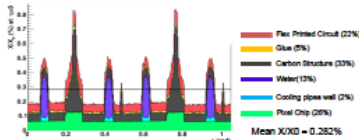
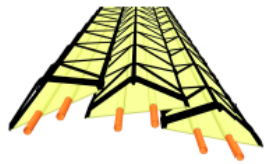


Figure 4.1: Schematic view of the Inner Barrel Slave.



Dramatic decrease in cost.

Very low power consumption, possibly $<100\text{mW/cm}^2$ i.e. simple water cooling

Ultra low material budget $<0.5\%$ for inner layers, $<1\%$ for outer layers.

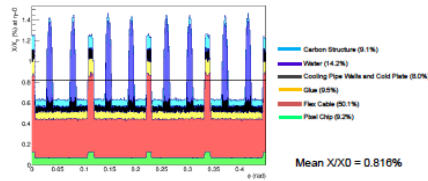
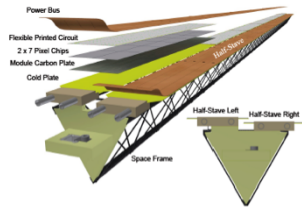
Question of speed and radiation hardness:

At present, integration time of $4\mu\text{s}$ (noise, electron diffusion) radiation resistance up to few 10^{13} neq.

Development (next 20 years) towards larger (full) depletion will improve speed and radiation hardness significantly.

Also – in case one has a full pixel tracker one can use 1 or 2 layers with ,fast' pixels to do the BCID (25ns or even 5ns) and then match the other hits.

With a full pixel tracker of $20 \times 20 \mu\text{m}$ pixels one can pile up a fair amount of events before occupancy gets to large !!!



→ Technical design report for the upgrade of the ALICE inner tracking system CERN-LHCC-2013-024

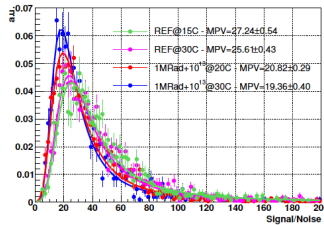


Figure 2.22: SNR of seed pixel measured with MIMOSA-32ter at the CERN-SPS, at two operating temperatures, before and after irradiation with the combined load of 1 Mrad and 10^{13} $1\text{ MeV } n_{\text{eq}}/\text{cm}^2$.