Semi-holography for heavy-ion collisions

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Thermalization in strongly coupled CFT

Gauge/gravity duality (aka holography)

Maldacena (1997): Duality between
• large-N_c maximally supersymmetric Yang-Mills theory at infinite 't Hooft coupling
and
• type-IIB supergravity on $\mathrm{AdS}_5 \times S^5$

used extensively to study
isotropization/thermalization/hydrodynamization in (infinitely) strongly coupled QFT

seminal paper: P. Chesler & L. Yaffe, "Horizon formation and far-from-equilibrium isotropization in supersymmetric Yang-Mills plasma", Phys.Rev.Lett. 102 (2009) 211601 [arXiv:0812.2053]

→ **talk by D. Mateos later today**

Semi-holographic models

Semi-holography:

dynamical boundary theory (weakly self-coupled in the UV)
coupled to a strongly coupled conformal sector with gravity dual
oxymoron coined by Faulkner & Polchinski, JHEP 1106 (2011) 012 [arXiv:1001.5049]
in study of holographic non-Fermi-liquid models

- retains only the universal low energy properties, which are most likely to be relevant to the realistic systems
- allows more flexible model-building

further developed for NFLs in:
A. Mukhopadhyay, G. Policastro, PRL 111 (2013) 221602 [arXiv:1306.3941]

Semi-holographic model for heavy-ion collisions

Aim: hybrid strong/weak coupling model of quark-gluon plasma formation
(QCD: strongly coupled in IR, weakly coupled in UV)

(different) successful example:
J. Casalderrey-Solana et al., JHEP 1410 (2014) 19 and JHEP 1603 (2016) 053

Idea of semi-holographic model by
E. Iancu, A. Mukhopadhyay, JHEP 1506 (2015) 003 [arXiv:1410.6448]:

combine pQCD (Color-Glass-Condensate) description of initial stage of HIC through overoccupied gluons with AdS/CFT description of thermalization

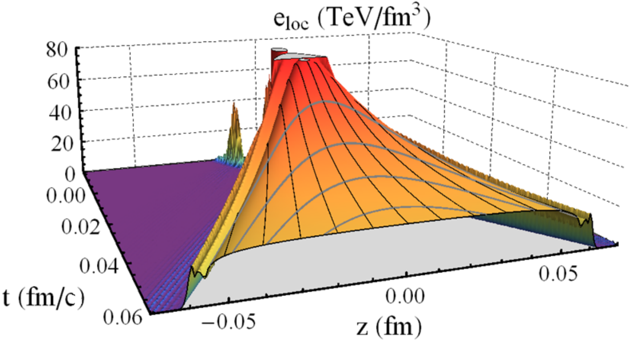
modified and extended recently in
A. Mukhopadhyay, F. Preis, A.R., S. Stricker, JHEP 1605 (2016) 141 [arXiv:1512.06445]

- s.t. $\exists$ conserved local energy-momentum tensor for combined system
- verified in (too) simple test case

Gravity dual of heavy-ion collisions

pioneered and developed in particular by P. Chesler & L. Yaffe [JHEP 1407 (2014) 086]

most recent attempt towards quantitative analysis along these lines:
Wilke van der Schee, Björn Schenke, PRC92 (2015) 064907 [1507.08195]



had to scale down energy density by a factor of 20 (6) for the top LHC (RHIC) energies

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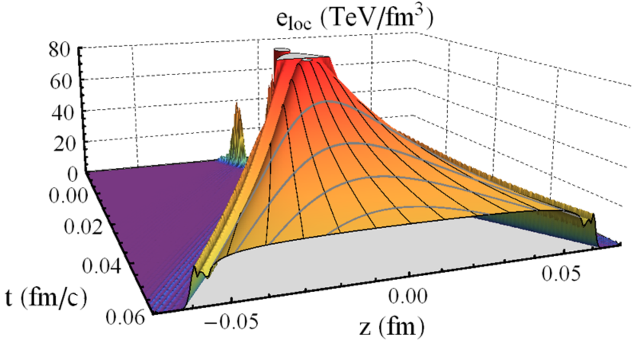


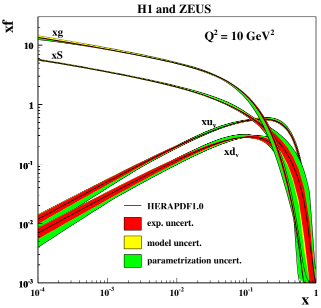
had to scale down energy density by a factor of 20 (6) for the top LHC (RHIC) energies

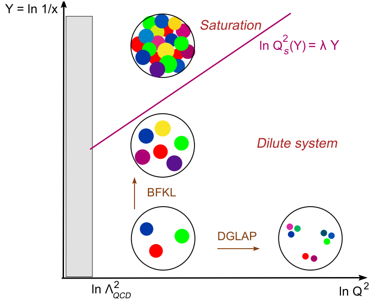
perhaps improved by involving pQCD for (semi-)hard degrees of freedom?

pQCD and Color-Glass-Condensate framework

recap: [e.g. F. Gelis et al., arXiv:1002.033]

gluon distribution $xG(x, Q^2)$ in a proton rises very fast with decreasing longitudinal momentum fraction x at large, fixed Q^2





HIC: high gluon density $\sim \alpha_s^{-1}$ at "semi-hard" scale Q_s ($\sim$ few GeV)

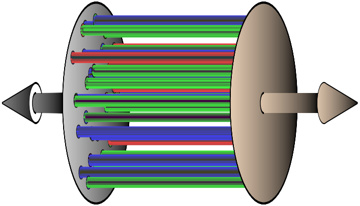
weak coupling $\alpha_s(Q_s) \ll 1$ but highly nonlinear because of large occupation numbers

description in terms of classical YM fields as long as gluon density nonperturbatively high

Color-Glass-Condensate evolution of HIC at LO

effective degrees of freedom in this framework:

1. color sources ρ at large x (frozen on the natural time scales of the strong interactions and distributed randomly from event to event)
2. gauge fields A^μ at small x
   (saturated gluons with large occupation numbers $\sim 1/\alpha_s$, with typical momenta peaked about $k_\perp Q_s$)

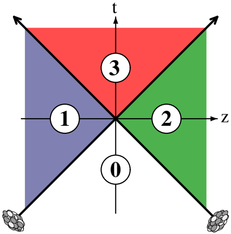


glasma: non-equilibrium matter, with high occupation numbers $\sim 1/\alpha_s$

initially longitudinal chromo-electric and chromo-magnetic fields that are screened at distances $1/Q_s$ in the transverse plane of the collision

Color-Glass-Condensate evolution of HIC at LO

- colliding nuclei as shock waves with frozen color distribution



classical YM field equations

$$D_\mu F^{\mu\nu}(x) = \delta^{\nu+}\rho_{(1)}(x^-, \mathbf{x}_\perp) + \delta^{\nu-}\rho_{(2)}(x^+, \mathbf{x}_\perp)$$

in Schwinger gauge $A^\tau = (x^+A^- + x^-A^+)/\tau = 0$
with ρ from random distribution (varying event-by-event)

outside the forward light-cone (3):
(causally disconnected from the collision)
pure-gauge configurations

$$\begin{aligned}
&A^+ = A^- = 0\\
&A^i(x) = \theta(-x^+)\theta(x^-)A^i_{(1)}(\mathbf{x}_\perp) + \theta(-x^-)\theta(x^+)A^i_{(2)}(\mathbf{x}_\perp)\\
&A^i_{(1,2)}(\mathbf{x}_\perp) = \tfrac{\mathrm{i}}{g}\, U_{(1,2)}(\mathbf{x}_\perp)\partial_i U^\dagger_{(1,2)}(\mathbf{x}_\perp)\\
&U_{(1,2)}(\mathbf{x}_\perp) = \mathrm{P}\ \exp\left(-\mathrm{i}g\int \mathrm{d}x^\mp \tfrac{1}{\nabla^2_\perp}\rho_{(1,2)}(x^\mp, \mathbf{x}_\perp)\right)
\end{aligned}$$

inside forward light-cone:
numerical solution with initial conditions at $\tau = 0$:

$$A^i = A^i_{(1)} + A^1_{(2)}, \qquad A^\eta = \tfrac{\mathrm{i}g}{2}\left[A^i_{(1)}, A^i_{(2)}\right], \qquad \partial_\tau A^i = \partial_\tau A^\eta = 0$$

Color-Glass-Condensate evolution of HIC at LO

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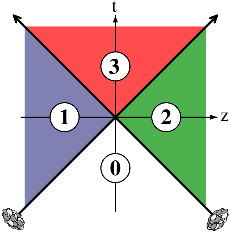

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- **Aim of semi-holographic model:** include bottom-up thermalization from relatively soft gluons with higher α_s and their backreaction when they build up thermal bath

Semi-holographic glasma evolution

[E. Iancu, A. Mukhopadhyay, JHEP 1506 (2015) 003]

[A. Mukhopadhyay, F. Preis, AR, S. Stricker, JHEP 1605 (2016) 141]

UV-theory=classical Yang-Mills theory for overoccupied gluon modes with $k \sim Q_s$
IR-CFT=effective theory of strongly coupled soft gluon modes $k \ll Q_s$,
modelled by N=4 SYM gravity dual marginally deformed by:
① boundary metric $g^{(\mathrm{b})}_{\mu\nu}$, ② dilaton $\phi^{(\mathrm{b})}$, and ③ axion $\chi^{(\mathrm{b})}$ which are functions of A_μ

$$S = S_{\mathrm{YM}}[A] + W_{\mathrm{CFT}}\left[g^{(\mathrm{b})}_{\mu\nu}[A], \phi^{(\mathrm{b})}[A], \chi^{(\mathrm{b})}[A]\right]$$

W_{CFT}: generating functional of the IR-CFT (on-shell action of its gravity dual)
minimalistic coupling through gauge-invariant dimension-4 operators

① IR-CFT energy-momentum tensor $\frac{1}{2\sqrt{-g^{(\mathrm{b})}}}\frac{\delta W_{\mathrm{CFT}}}{\delta g^{(\mathrm{b})}_{\mu\nu}} = \mathcal{T}^{\mu\nu}$ coupled to energy-momentum tensor $t_{\mu\nu}$ of YM (glasma) fields through

$$g^{(\mathrm{b})}_{\mu\nu} = \eta_{\mu\nu} + \frac{\gamma}{Q_s^4} t_{\mu\nu}, \quad t_{\mu\nu}(x) = \frac{1}{N_c}\mathrm{Tr}\Big(F_{\mu\alpha}F_\nu{}^\alpha - \frac{1}{4}\eta_{\mu\nu}F_{\alpha\beta}F^{\alpha\beta}\Big);$$

② $\phi^{(\mathrm{b})} = \frac{\beta}{Q_s^4}h, \quad h(x) = \frac{1}{4N_c}\mathrm{Tr}(F_{\alpha\beta}F^{\alpha\beta})$; ③ $\chi^{(\mathrm{b})} = \frac{\alpha}{Q_s^4}a, \quad a(x) = \frac{1}{4N_c}\mathrm{Tr}\left(F_{\mu\nu}\tilde{F}^{\mu\nu}\right)$

α, β, γ dimensionless and $O(1/N_c^2)$

Semi-holographic glasma evolution

IR-CFT: marginally deformed AdS/CFT

in Fefferman-Graham coordinates:

$$
\begin{aligned}
\chi(z,x) &= \frac{\alpha}{Q_s^4} a(x) + \cdots + z^4 \frac{4\pi G_5}{l^3} \mathcal{A}(x) + \mathcal{O}(z^6), \\
\phi(z,x) &= \frac{\beta}{Q_s^4} h(x) + \cdots + z^4 \frac{4\pi G_5}{l^3} \mathcal{H}(x) + \mathcal{O}(z^6), \\
G_{rr}(z,x) &= \frac{l^2}{z^2}, \\
G_{r\mu}(z,x) &= 0, \\
G_{\mu\nu}(z,x) &= \frac{l^2}{z^2} \Big(\underbrace{\eta_{\mu\nu} + \frac{\gamma}{Q_s^4} t_{\mu\nu}(x)}_{g^{(\mathrm{b})}_{\mu\nu} = g_{(0)\mu\nu}} + \cdots + z^4 \Big(\underbrace{\frac{4\pi G_5}{l^3}}_{2\pi^2/N_c^2} \mathcal{T}_{\mu\nu}(x) + P_{\mu\nu}(x) \Big) \\
&\qquad + \mathcal{O}(z^4 \ln z) \Big),
\end{aligned}
$$

with $P_{\mu\nu} = \frac{1}{8} g_{(0)\mu\nu} \left((\mathrm{Tr}\, g_{(2)})^2 - \mathrm{Tr}\, g_{(2)}^2 \right) + \frac{1}{2} (g_{(2)}^2)_{\mu\nu} - \frac{1}{4} g_{(2)\mu\nu} \mathrm{Tr}\, g_{(2)}$

[de Haro, Solodukhin, Skenderis, CMP 217 (2001) 595]

Semi-holographic glasma evolution

Modified YM (glasma) field equations

$$\frac{\delta S}{\delta A_\mu(x)} = \frac{\delta S_{\rm YM}}{\delta A_\mu(x)} + \int \mathrm{d}^4 y \left(\frac{\delta W_{\rm CFT}}{\delta g^{(\rm b)}_{\alpha\beta}(y)} \frac{\delta g^{(\rm b)}_{\alpha\beta}(y)}{\delta A_\mu(x)} + \frac{\delta W_{\rm CFT}}{\delta \phi^{(\rm b)}(y)} \frac{\delta \phi^{(\rm b)}(y)}{\delta A_\mu(x)} + \frac{\delta W_{\rm CFT}}{\delta \chi^{(\rm b)}(y)} \frac{\delta \chi^{(\rm b)}(y)}{\delta A_\mu(x)} \right)$$

gives

$$D_\mu F^{\mu\nu} = \frac{\gamma}{Q_s^4} D_\mu \left(\hat{\mathcal{T}}^{\mu\alpha} F_\alpha{}^\nu - \hat{\mathcal{T}}^{\nu\alpha} F_\alpha{}^\mu - \frac{1}{2} \hat{\mathcal{T}}^\alpha_\alpha F^{\mu\nu} \right) + \frac{\beta}{Q_s^4} D_\mu \left(\hat{\mathcal{H}} F^{\mu\nu} \right) + \frac{\alpha}{Q_s^4} \left(\partial_\mu \hat{\mathcal{A}} \right) \tilde{F}^{\mu\nu}$$

with $\hat{\mathcal{T}}^{\alpha\beta} = \frac{\delta W_{\rm CFT}}{\delta g^{(\rm b)}_{\alpha\beta}} = \sqrt{-g^{(b)}}\mathcal{T}^{\alpha\beta}, \quad \hat{\mathcal{H}} = \frac{\delta W_{\rm CFT}}{\delta \phi^{(\rm b)}} = \sqrt{-g^{(b)}}\mathcal{H}, \quad \hat{\mathcal{A}} = \frac{\delta W_{\rm CFT}}{\delta \chi^{(\rm b)}} = \sqrt{-g^{(b)}}\mathcal{A}$

Total energy-momentum tensor of combined system

IR-CFT, like glasma EFT, interpreted as living in Minkowski space

covariant conservation equation for energy-momentum tensor

$$\nabla_\mu \mathcal{T}^{\mu\nu}(x) = -\frac{\beta}{Q_s^4}\,\mathcal{H}(x)\nabla^\nu h(x),$$

with metric $g^{(\mathrm{b})}_{\mu\nu}(x) = \eta_{\mu\nu} + \frac{\gamma}{Q_s^4} t_{\mu\nu}(x)$

$\rightarrow$ nonconservation in Minkowski space, with driving forces derived from UV $t_{\mu\nu}[A]$

$$\partial_\mu \mathcal{T}^{\mu\nu} = -\frac{\beta}{Q_s^4}\,\mathcal{H}\, g^{\mu\nu}_{(\mathrm{b})}[t]\,\partial_\mu h - \mathcal{T}^{\alpha\nu}\Gamma^\gamma_{\alpha\gamma}[t] - \mathcal{T}^{\alpha\beta}\Gamma^\nu_{\alpha\beta}[t]$$

with $\Gamma^\mu_{\nu\rho}[t] = \frac{\gamma}{2Q_s^4}\left(\partial_\nu t^\mu{}_\rho + \partial_\rho t^\mu{}_\nu - \partial^\mu t_{\nu\rho}\right) + \mathcal{O}(t^2)$

Total conserved energy-momentum tensor in Minkowski space ($\partial_\mu T^{\mu\nu} = 0$):

$$T^{\mu\nu} = t^{\mu\nu} + \mathcal{T}^{\mu\nu} + \text{hard-soft interaction terms}$$

(but $\mathcal{T}^{\mu\nu}$ not purely soft, contains also some hard-soft pieces through $g^{\mu\nu}_{(\mathrm{b})}[t]$)

Total energy-momentum tensor of combined system

Temporarily replacing Minkowski metric $\eta_{\mu\nu}$ by $g^{\rm YM}_{\mu\nu}$:

$$
\begin{aligned}
T^{\mu\nu} \;=\;& \frac{2}{\sqrt{-g^{\rm YM}}}\Bigg[\frac{\delta S_{\rm YM}}{\delta g^{\rm YM}_{\mu\nu}(x)} \\
&+\int {\rm d}^4y \left(\frac{\delta W_{\rm CFT}}{\delta g^{(b)}_{\alpha\beta}(y)}\frac{\delta g^{(b)}_{\alpha\beta}(y)}{\delta g^{\rm YM}_{\mu\nu}(x)} + \frac{\delta W_{\rm CFT}}{\delta \phi^{(b)}(y)}\frac{\delta \phi^{(b)}(y)}{\delta g^{\rm YM}_{\mu\nu}(x)} + \frac{\delta W_{\rm CFT}}{\delta \chi^{(b)}(y)}\frac{\delta \chi^{(b)}(y)}{\delta g^{\rm YM}_{\mu\nu}(x)}\right)\Bigg]
\end{aligned}
$$

At $g^{\rm YM}_{\mu\nu} = \eta_{\mu\nu}$, this gives

$$
\begin{aligned}
T^{\mu\nu} \;=\;& t^{\mu\nu} + \hat{\mathcal{T}}^{\mu\nu} \\
&-\frac{\gamma}{Q_s^4 N_c}\hat{\mathcal{T}}^{\alpha\beta}\left[{\rm Tr}(F_\alpha{}^\mu F_\beta{}^\nu) - \frac{1}{2}\eta_{\alpha\beta}{\rm Tr}(F^{\mu\rho}F^\nu{}_\rho) + \frac{1}{4}\delta^\mu_{(\alpha}\delta^\nu_{\beta)}{\rm Tr}(F^2)\right] \\
&-\frac{\beta}{Q_s^4 N_c}\hat{\mathcal{H}}\,{\rm Tr}(F^{\mu\alpha}F^\nu{}_\alpha) - \frac{\alpha}{Q_s^4}\eta^{\mu\nu}\hat{\mathcal{A}}\,a
\end{aligned}
$$

Can indeed prove $\boxed{\partial_\mu T^{\mu\nu} = 0}$

Iterative solution

Practical implementation will have to be done presumably in iterative procedure

1. Solve LO glasma evolution with $\gamma = \beta = \alpha = 0$
2. solve gravity problem with boundary condition provided by glasma $t^{\mu\nu}(\tau)$,...
   to obtain $\mathcal{T}^{\mu\nu}(\tau), \ldots$

   As in previous thermalization studies:

   need thermal AdS/CFT with small nonzero temperature to start from (seed black hole)
3. solve glasma evolution with $\gamma, \beta, \alpha \neq 0$ and given $\mathcal{T}^{\mu\nu}(\tau), \ldots$
4. goto 2) until convergence reached

Simple test case

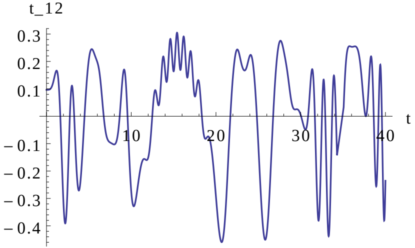
First test with dimensionally reduced (spatially homogeneous) YM fields $A^a_\mu(t)$ which already have nontrivial (in general chaotic) dynamics

SU(2) gauge fields, $a = 1, 2, 3$, using temporal gauge $A^a_0 = 0$, $g = 1$

$$D^\mu F_{\mu j} = 0 \Rightarrow \ddot{A}^a_j - A^a_i A^b_i A^b_j + A^a_j A^b_i A^b_i = 0,$$

Gauss law
$D^\mu F^d_{\mu 0} = 0 \Rightarrow \epsilon^{dea} A^{ei} \dot{A}^a_i = 0$, satisfied by initial conditions $A(t_0) = 0$ or $\dot{A}(t_0) = 0$

General case: 9 degrees of freedom with chaotic dynamics
conserved, traceless energy-momentum tensor with $\varepsilon = const.$, $t^{0i} = 0$,
but otherwise wildly fluctuating:



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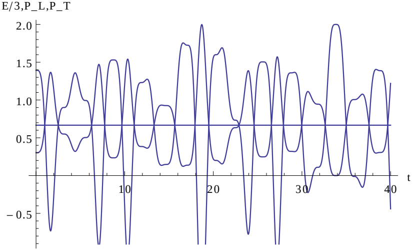
Color-spin-locked: $A^a_i \propto \delta^a_i$, 3 degrees of freedom
$\rightarrow$ diagonal but anisotropic traceless energy-momentum tensor
e.g., only one direction singled out:

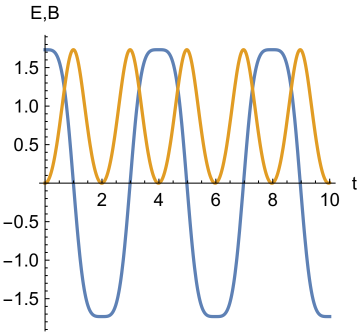


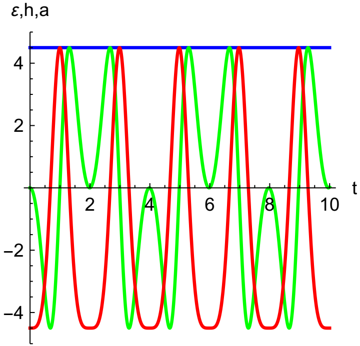
Simplest test case

$\exists$ a nontrivial solution with homogeneous isotropic energy-momentum tensor ($p = \varepsilon/3$) by homogeneous and isotropic color-spin locked oscillations $A_0^a = 0, \quad A_i^a = \delta_i^a f(t)$

$$f(t) = C\,\mathrm{sn}(C(t-t_0)|-1) \qquad \text{(Jacobi elliptic function sn)}$$

$$E_i^a = \delta_i^a f', \quad B_i^a = \delta_i^a f^2 \qquad \varepsilon = const., \; h = -\tfrac{1}{2}(\mathbf{E}^a \cdot \mathbf{E}^a - \mathbf{B}^a \cdot \mathbf{B}^a), \; a = -\mathbf{E}^a \cdot \mathbf{B}^a$$





Simplest test case – gravitational solution

Switching off $\alpha = 0 = \beta$ (otherwise also nontrivial sources for dilaton and axion!)

IR-CFT $\hat{\mathcal{T}}^{\mu\nu} = \mathrm{diag}(\hat{\mathcal{E}}, \hat{\mathcal{P}}, \hat{\mathcal{P}}, \hat{\mathcal{P}})$ to be determined by gravitational problem with

$$g^{(\mathrm{b})}_{\mu\nu} = \eta_{\mu\nu} + \frac{\gamma}{Q_s^4} t_{\mu\nu} = \mathrm{diag}\left(-1 + \frac{3\gamma}{Q_s^4} p(t), 1 + \frac{\gamma}{Q_s^4} p(t), 1 + \frac{\gamma}{Q_s^4} p(t), 1 + \frac{\gamma}{Q_s^4} p(t)\right)$$

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Result:

$$\begin{aligned}
\hat{\mathcal{E}} \quad &:= \quad \hat{\mathcal{T}}^{tt} = \frac{N_c^2}{2\pi^2}\left(\frac{3c}{4r_{(0)}(t)^2 v'_{(0)}(t)} + \frac{3r'_{(0)}(t)^4}{16 r_{(0)}(t)^6 v'_{(0)}(t)^5}\right), \\
\hat{\mathcal{P}} \quad &:= \quad \hat{\mathcal{T}}^{xx} = \hat{\mathcal{T}}^{yy} = \hat{\mathcal{T}}^{zz} = \\
&= \quad \frac{N_c^2}{2\pi^2}\left\{\frac{c v'_{(0)}(t)}{4r_{(0)}(t)^2} + \frac{r'_{(0)}(t)^2\left[4r_{(0)}(t) r'_{(0)}(t) v''_{(0)}(t) + r_{(0)}(t)\left(5r'_{(0)}(t)^2 - 4r_{(0)}(t) r''_{(0)}(t)\right)\right]}{16 r_{(0)}(t)^6 v'_{(0)}(t)^4}\right\},
\end{aligned}$$

with

$$r_{(0)}(t) = \sqrt{1 + (\gamma/Q_s^4) p(t)}\,, \qquad v'_{(0)}(t) = \sqrt{\frac{1 - (\gamma/Q_s^4) 3p(t)}{1 + (\gamma/Q_s^4) p(t)}}$$

because of isotropy and homogeneity, gravity solution locally diffeomorphic to AdS-Schwarzschild with integration constant c corresponding to mass of black hole

Convergence of iterations

Coupled glasma equation of test case is 4th order nonlinear ODE
— no reasonable solutions found directly —

Convergence of iterations

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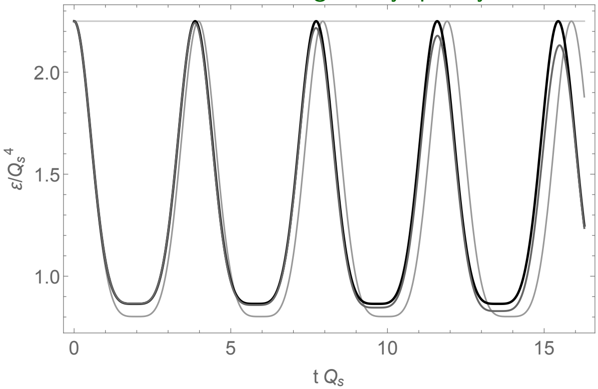
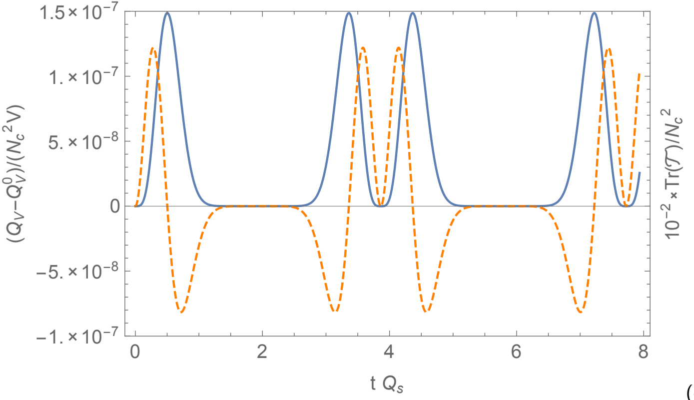
— no reasonable solutions found directly —
but iterative solution converges very quickly:



UV not able to give off energy to IR permanently because of isotropy and homogeneity:
gravity dual does not have propagating degrees of freedom!

Energy exchanges but no thermalization

No test of thermalization yet:
entropy (area of the black hole) is conserved,
canonical charge of the black hole changes only according to trace anomaly:



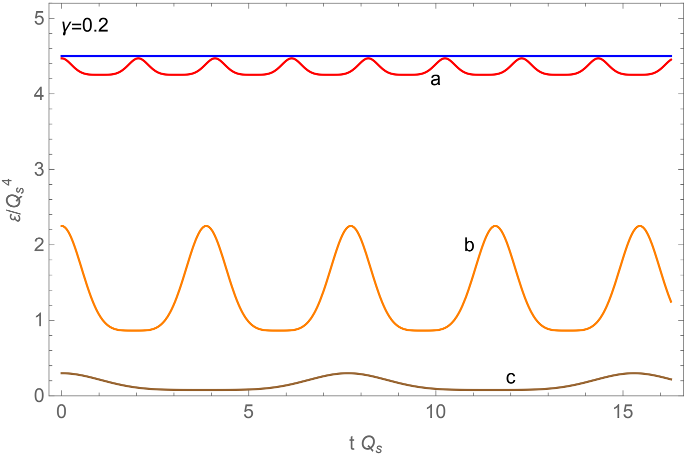
(case b)

blue: canonical charge (thermal energy) returns to same value at stationary points (at different extrema of $\varepsilon^{\mathrm{YM}}$)
orange: $\mathrm{Tr}\,\mathcal{T}$

Solutions with different IR entropy

initial conditions with
little (a) - medium (b) - large (c)
thermal (IR) contribution to total (conserved) energy

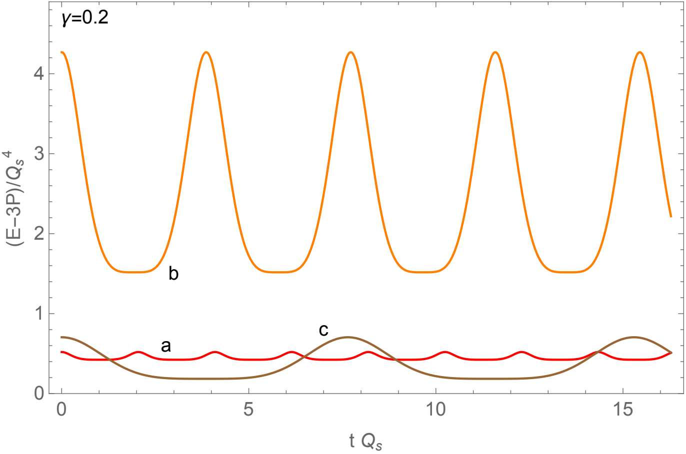
(ε=YM energy density)



Solutions with different IR entropy

initial conditions with
little (a) - medium (b) - large (c)
thermal (IR) contribution to total energy

trace of the full energy-momentum tensor (“interaction measure”):



Outlook

Expect thermalization (actual growth of black hole entropy) as soon as $\alpha, \beta \neq 0$ (propagating scalar d.o.f. in bulk)
or when no longer homogeneous and isotropic (propagating tensor d.o.f. in bulk)

Numerical studies of homogeneous, isotropic case with $\beta \neq 0$ currently ongoing
(with C. Ecker, A. Mukhopadhyay, F. Preis, S. Stricker)

Related previous works (w/o backreaction on source):

▷ R. Auzzi, S. Elitzur, S.B. Gudnason, E. Rabinovici,
"Periodically Driven AdS/CFT", JHEP 1311 (2013) 016

▷ M. Rangamani, M. Rozali, A. Wong,
"Driven Holographic CFTs", JHEP 1504 (2015) 093

studied entropy production when thermal system is driven by given external scalar source

Conclusions

Pure gauge-gravity thermalization treats infinite coupling limit
Long-term goal: hybrid description with less strongly coupled UV sector

- Semi-holographic framework of Iancu and Mukhopadhyay proposes to combine LO glasma evolution with thermalization of soft degrees of freedom in AdS/CFT
- New scheme has conserved total energy-momentum tensor (in Minkowski space)
  formal proof + numerical verification in simple test case
- First tests suggest that proposed iterative scheme can be numerically stable and convergent
- Next step: test case with propagating modes in bulk
- Also ongoing: Semi-holographic setup for coupling Yang-Mills fluctuations to late-time hydro evolution (with Y. Hidaka, A. Mukhopadhyay, F. Preis, A. Soloviev, D.-L. Yang)

Details of gravitational side of test case

Homogeneous isotropic ansatz in Eddington-Finkelstein coordinates

$$ds^2 = -A(r,v)dv^2 + 2dr\,dv + \Sigma(r,v)d\vec{x}^2$$

equations of motion take the compact form

$$\begin{aligned} 0 &= \Sigma(\dot{\Sigma})' + 2\Sigma'\dot{\Sigma} - 2\Sigma^2 \,, \\ 0 &= A'' - 12\Sigma'\dot{\Sigma}/\Sigma^2 + 4 \\ 0 &= 2\ddot{\Sigma} - A'\dot{\Sigma} \\ 0 &= \Sigma'' \end{aligned}$$

Power series ansatz in r involves only finite number of terms, one integration constant c

$$\begin{aligned} A(r,v) &= r^2\left(1 - \frac{c}{r^4\Sigma_0(v)}\right) - 2r\frac{\partial_v\Sigma_0(v)}{\Sigma_0(v)} \\ \Sigma(r,v) &= r\,\Sigma_0(v) \end{aligned}$$

Details of gravitational side of test case

Transformation to Fefferman-Graham coordinates with

$$g_{(0)\mu\nu} = \mathrm{diag}\left(-1 + \frac{\gamma}{Q_s^4}3p, 1 + \frac{\gamma}{Q_s^4}p, 1 + \frac{\gamma}{Q_s^4}p, 1 + \underbrace{\frac{\gamma}{Q_s^4}p}_{\tilde{p}}\right)$$

$$T^\mu_\nu = \frac{N_c^2}{2\pi^2}\mathrm{diag}(-\mathcal{E}, P, P, P),$$

$$r^{-1} = u = u_1(t)z + O(z^2), \quad v = t + O(z)$$

$$\begin{aligned}
\mathcal{E} &= \frac{3}{4}\frac{c}{r^4\Sigma_0(v)}u_1^4 - \frac{2(u_1\Sigma_0^{(1)} - \Sigma_0 u_1^{(1)})^4}{16\Sigma_0^4}\,, \\
P &= \frac{1}{4}\frac{c}{r^4\Sigma_0(v)}u_1^4 \\
&\quad + \frac{1}{16\Sigma_0^4}\Big[\left(u_1\Sigma_0^{(1)} - \Sigma_0 u_1^{(1)}\right)^2\Big(-3\Sigma_0^2(u_1^{(1)})^2 + u_1^2\left((\Sigma_0^{(1)})^2 - 4\Sigma_0\Sigma_0^{(2)}\right) \\
&\qquad + 2\Sigma_0 u_1\left(\Sigma_0^{(1)}u_1^{(1)} + 2\Sigma_0 u_1^{(2)}\right)\Big)\Big]
\end{aligned}$$

with $u_1 = \frac{1}{\sqrt{1-3\tilde{p}}}$, $\qquad \Sigma_0 = \sqrt{\frac{1+\tilde{p}}{1-3\tilde{p}}}$

Details of gravitational side of test case

Solution in Eddington-Finkelstein coordinates

$$\begin{aligned} A(r,v) &= r^2\left(1-\frac{c}{r^4\Sigma_0(v)}\right)-2r\frac{\partial_v\Sigma_0(v)}{\Sigma_0(v)} \\ \Sigma(r,v) &= r\,\Sigma_0(v) \end{aligned}$$

is locally diffeomorphic to Schwarzschild solution

$$ds^2=-R^2\left(1-\frac{c}{R^4}\right)dV^2+2dR\,dV+R^2d\vec{x}^2$$

through coordinate transformation $R=r\Sigma_0(v)$ and $V=\int\frac{dv}{\Sigma_0(v)}$

Time dependence of boundary metric has however
nontrivial effect on Brown-York stress tensor $\mathcal{T}^{\mu\nu}$