Holographic Heavy Ion Collisions

David Mateos ICREA & University of Barcelona

with Maximilian Attems, Jorge Casalderrey, Michal Heller, Daniel Santos-Olivan, Carlos Sopuerta, Miquel Triana, Wilke van der Schee and Miguel Zilhao

• Holography applied to QCD — limitations.

- Holography applied to heavy ion collisions overview.
	- ‣ For newest results (non-conformal theories) see talk by Jorge Casalderrey on Thursday.
- Holography applied to cosmology not for today.

Animation by Jeffery Mitchell (Brookhaven National Laboratory). Simulation by the UrQMD Collaboration

Heavy ion collisions

- How long is t_{hydro} ? Data indicates $t_{\text{hydro}}T_{\text{hydro}} \leq 1$.
- What determines when hydro becomes applicable?
 i *P/E* ⇠ 1 (2) *The The Theodores and The Theodores and The Theodores* and The Theodores and
- *b*(*t*) = *B*(*t*) = *B*(*t*) = *R*) = *B*(*t*) = *B*(*t* • What is the nature of the hydro expansion?
	- What are the initial conditions for hydro? *i cⁱ bi*(*r*)*ei*!*i^t* (3) **P**² (3) **P**²
	- *t* trong & weak coupling physics. • Mixture of strong & weak coupling physics.
	- All explained by QCD, but QCD is hard. **i i** $\frac{1}{2}$ (5) \frac

From viewpoint of a *theorist*

• Duality is a remarkable development:

$$
Quantum gravity = \boxed{\text{Ordinary QFT}}
$$

In terms of *applications* to QCD

At present the duality has its own limitations

 $(N_c \rightarrow \infty)$

esses quanti Suppresses quantum corrections. Suppresses string corrections.

 $\lambda = g_{\scriptscriptstyle\text{{YM}}}^2 N_{\scriptscriptstyle\text{{C}}} \to \infty$

g corre **g**
esses string correction $\frac{2}{\pi}$

 $(N_c \rightarrow \infty)$

 $\lambda = g_{\scriptscriptstyle\text{{YM}}}^2 N_{\scriptscriptstyle\text{{C}}} \to \infty$

esses quanti Suppresses quantum corrections. Suppresses string corrections.

g corre **g**
esses string correction $\frac{2}{\pi}$

$$
\boxed{N_{\rm c}\to\infty}
$$

 $\lambda = g_{\rm YM}^2 N_{\rm c} \to \infty$

*R*4

esses quanti Suppresses quantum corrections. Makes the string tiny.

 $\frac{1}{2}$ g tiny. $\frac{1}{2}$ s the $\frac{1}{2}$ ring

Solving large- N_c would be great progress!

$$
\begin{array}{c}\hline\\[-1.2mm]\hline\\[-1.
$$

esses quanti Suppresses quantum corrections. Makes the string tiny.

- Asymptotically free.
- Dynamically generated scale.
- Confinement.
- *N*e *Decont* • Deconfinement phase transition.

 $\lambda = g_{\rm YM}^2 N_{\rm c} \to \infty$

$\frac{1}{2}$ g tiny. $\frac{1}{2}$ s the $\frac{1}{2}$ ring

• …

$$
\boxed{N_{\rm c}\to\infty}
$$

 = *a d*⌅ *N*^c Suppresses quantum corrections. Makes the string tiny.

$$
\lambda = g_{\rm YM}^2 N_{\rm c} \to \infty
$$

 $\frac{1}{2}$ g tiny. $\frac{1}{2}$ s the $\frac{1}{2}$

Strong coupling means no asymptotic freedom!

Therefore

• At present gauge/string duality is not a tool for *precision* QCD physics.

• However, it may still provide useful insights.

• In particular, if strong coupling + far from equilibrium then holography is the *only* first-principle tool.

What we would like to do

Heavy ion collisions in QCD

What we can do

Formation and evolution of the QGP

Holographic heavy ion collisions in CFT

Chesler & Yaffe '10

Toy model for collisions of infinite nuclei with no baryon charge:

Holographic heavy ion collisions in CFT

Chesler & Yaffe '10

• No transverse dynamics.

1 ua

• CFT implies EOS obeyed in and out of equilibrium: 1. st order viscous hydrogeness.

$$
T^{\mu}_{\mu} = 0 \qquad \rightarrow \qquad \bar{P} = P_{eq}(\mathcal{E}) = \frac{1}{3}\mathcal{E}
$$

$$
\bar{P} = \frac{1}{3}(P_L + 2P_T)
$$

- *av* \mathbf{e} $\overline{}$ *t*hydro <u>'</u> *f P P P P A statement about average pressure.* size: I P *POS* • I emphasize: EOS is a statement about average pressure.
- *P*_{*L*} (5) *P_L* (5) *P_L* (5) *P_L* (5) *P_L* (5) *P_L* (5) *P_L* (5) *P* • Therefore P_L and P_T can deviate a lot from P_{eq} ! $P_{\rm eq}$!

Holographic heavy ion collisions in CFT *P*eq (3) *PT* 1st order viscous hydro

Chesler & Yaffe '10 $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

Holographic heavy ion collisions in CFT *P*eq (3) *PT* 1st order viscous hydro Γ ⇢!CY ' 0*.*64 (10)

Chesler & Yaffe '10 $\frac{1}{2}$ $\frac{1}{2}$ $\frac{1}{2}$

Pressures at mid rapidity $\frac{1}{2}$ **Pressures at mighting the Pressures at mighting of the Pressures at mighting** \mathbf{P}

- Hydro applies at $t_{\text{hydro}}T_{\text{hydro}} \simeq 0.65$. $\qquad 0.5$
- Hydrodynamization without isotropization:

$$
\left. \frac{P_T}{P_L} \right|_{t_{\rm hydro}} \simeq
$$

• Hydro works when gradients are still very large: $P_{\perp}^{\text{hyd}} = P$

adients are still very large:

\n
$$
P_L^{\text{hyd}} = P_{\text{eq}} + P_{\eta} + \cancel{\mathbb{X}}
$$
\n
$$
P_T^{\text{hyd}} = P_{\text{eq}} - \frac{1}{2}P_{\eta} + \cancel{\mathbb{X}}
$$
\n
$$
\int_{\text{shear}}^{\text{hyd}} \frac{1}{\text{shear}} \, \text{bulk}
$$

tiscosity viscosity

A dynamical cross-over Casalderrey, Heller, D.M. & van der Schee '13

Qualitatively different dynamics depending on the collision energy:

Low energy collision (thick shocks)

—

- Realizes Landau model approximately: Energy gets compressed, stops and explodes hydrodynamically.
- No clear separation between plasma and receding fragments.
- The receding maxima move at $v \sim 0.88$.

Transparency scenario

- Shocks pass through one another and plasma gets created in between.
- The receding maxima move at v ~ 1 despite infinite coupling.
- · Clear separation between receding fragments and plasma.

- Motivation: p+A collisions have asymmetric longitudinal extent/structure.
- Motivation: In fact, A+A collisions also have longitudinal structure (albeit symmetric).
- Question: Does any of this leave an imprint on the resulting plasma?
- Compare the following collisions (at fixed total energy):

• Answer: Longitudinal structure leaves no imprint if $\ell_{\rm char} \lesssim 0.26/T_{\rm hyd}$ (coherence).

Coherent regime

• Answer: Longitudinal structure leaves no imprint if $\ell_{\rm char} \lesssim 0.26/T_{\rm hyd}$ (coherence).

Coherent regime

• Answer: Longitudinal structure leaves no imprint if $\ell_{\rm char} \lesssim 0.26/T_{\rm hyd}$ (coherence).

Coherent regime

• Answer: Longitudinal structure leaves no imprint if $\ell_{\rm char} \lesssim 0.26/T_{\rm hyd}$ (coherence).

Coherent regime Incoherent regime

• Answer: Longitudinal structure leaves no imprint if $\ell_{\rm char} \lesssim 0.26/T_{\rm hyd}$ (coherence).

Coherent regime Incoherent regime

- Answer: Longitudinal structure leaves no imprint if $\ell_{\rm char} \lesssim 0.26/T_{\rm hyd}$ (coherence).
- Implication: In coherent regime c.o.m. of QGP equals c.o.m. of all participating nucleons.

Coherent regime Incoherent regime

Collisions with baryon charge Casalderrey, D.M., van der Schee & Triana '16

Toy model for collisions of infinite nuclei with baryon charge:

Collisions with baryon charge Casalderrey, D.M., van der Schee & Triana '16

- We find significant stopping of baryon number.
- Hence good model for low- and moderate-energy collisions but not for high-energy.

• At high energies, rapidity shifts of valence quarks involve large momentum transfers and are suppressed by asymptotic freedom.

• Suggests using a hybrid model. Casalderrey, Gulhan, Milhano, Pablos & Rajagopal '14

- *z z* 15 Iancu & Mukhopadhyay '15
	- Mukhopadhyay, Preis, Rebhan & Stricker '16

Beyond conformal symmetry

Attems, Casalderrey, D.M., Santos-Olivan, Sopuerta, Triana & Zilhao '16

For details see talk by Jorge Casalderrey on Thursday.

Beyond conformal symmetry \blacksquare

Attems, Casalderrey, D.M., Santos-Olivan, Sopuerta, Triana & Zilhao '16 time [21]. We choose *t* = 0 as the time at which the two

- Main conclusions:
	- ‣ EOS does NOT hold out of equilibrium.

Beyond conformal symmetry \blacksquare

Attems, Casalderrey, D.M., Santos-Olivan, Sopuerta, Triana & Zilhao '16 time [21]. We choose *t* = 0 as the time at which the two

- Main conclusions:
	- ‣ EOS does NOT hold out of equilibrium.
	- ‣ Hydrodynamization without equilibration.

Beyond conformal symmetry \blacksquare

Attems, Casalderrey, D.M., Santos-Olivan, Sopuerta, Triana & Zilhao '16 time [21]. We choose *t* = 0 as the time at which the two

- 1st order viscous hydro • Main conclusions:
	- \triangleright EOS does NOT hold out of equilibrium.
	- **P** Hydrodynamization without equilibration.

 $P_L^{\text{hyd}} = P_{\text{eq}} + P_{\eta} + P_{\zeta}$ $P_{T}^{\rm hyd} \;\; = \;\; P_{\rm eq} - \frac{1}{2}$ 2 **The properties of the contract of the** *Responsible* for anisotropy 1st order viscous hydro

Beyond conformal symmetry λ

Attems, Casalderrey, D.M., Santos-Olivan, Sopuerta, Triana & Zilhao '16 time [21]. We choose *t* = 0 as the time at which the two

- Main conclusions: 1st order viscous hydro
	- \triangleright EOS does NOT hold out of equilibrium.
	- **P** Hydrodynamization without equilibration.

 $P_L^{\text{hyd}} = P_{\text{eq}} + P_{\eta} + P_{\zeta}$ $P_{T}^{\rm hyd} \;\; = \;\; P_{\rm eq} - \frac{1}{2}$ 2 **The properties of the contract of the** *Responsible* for anisotropy 1st order viscous hydro

- $\frac{1}{2}$ *R* vis \overline{a} *^T* ⁼ *^P*eq ¹ Required bulk viscosity about 1/10 of QCD at Tc.
- ▶ Hydro time 2.5 longer than in CFT.

Off-centre collisions of finite nuclei

Chesler & Yaffe '15

Off-centre collisions of finite nuclei Chesler & Yaffe '15

But essentially no elliptic flow. (perhaps due to transverse Gaussians).

See development of transverse flow.

p+A collisions and the smallest drops of QGP Chesler '15

• Produce droplets of size $R \sim 1/T_{\rm hyd}$ that are well described by hydro.

Thank you.