Viscous Cosmology — Is the dark sector perfect?

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- Cosmic expansion and dissipation
- A unified dark sector?
- -Viscous dark matter?
- Observables (focus on bulk viscosity)

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Basic observations

- I expansion history (SNIa, BAO, ...)
 - → accelerated expansion of the Universe
- 2- geometry (CMB)
 - → spatial flatness
- 3- growth of structures (CMB, LSS surveys)
 → non-relativistic (cold/warm) dark matter

Composition of the Universe



dark sector: cosmological constant & cold dark matter

Why does vanilla cosmology model the cosmic substratum as perfect fluids?

hydrodynamic limit: $1/\tau_{\rm relax} \gg H \equiv \frac{d \ln V}{dt}$

cosmological constant o.k. but for dark matter ?

Physical Cosmology

I- equations of motion

 → general relativity & energy/matter content
 (or modification of GR)

 2- initial/boundary conditions

 → ??? (inflation?, QG???)

instead

extra symmetry principles to avoid initial/boundary data

cosmological principle(s) !!!

Cosmological principle(s)

Exact: The Universe is spatially isotropic and homogeneous. ruled out by the fact that we see cosmic structures, reasonable 1st approximation

Statistical: The distribution of mass and light in the Universe is statistically isotropic and homogeneous.

now:

consider comoving cells of volume V, large enough that statistical averages are useful, but much smaller than scales of interest

Isotropic & Homogeneous Universe

 $T^{\mu}_{\ \nu} = (\epsilon + P)u^{\mu}u_{\nu} + \delta^{\mu}_{\nu}P, \qquad \epsilon = \epsilon(t), P = P(t), u^{\mu} = u^{\mu}(t)$

 $V \mathrm{d} \epsilon = -(\epsilon + P) \mathrm{d} V + \delta' Q$ Ist law of TD for each voxel

δ'Q heat flow: must vanish if no direction is preferred ⇒ isotropic & homogeneous fluid

But does that mean that there is no dissipation?

dS > 0 does allow for bulk viscosity:

Eckart 1940; $P = p + \Pi = p - 3H\zeta, \quad \zeta = \zeta(t) \ge 0$ Landau & Lifshitz 1958

st or 2nd order?

Müller 1967; Israel & Stewart 1976

Interest in cosmic bulk viscosity arose in the in the context of cosmological inflation (acceleration of the expansion!)

see e.g. review by Maartens 1995

$$\Pi + \tau \dot{\Pi} = -3H\zeta - \frac{1}{2}\tau \Pi \left(3H + \frac{\dot{\tau}}{\tau} - \frac{\dot{\zeta}}{\zeta} - \frac{\dot{T}}{T} \right)$$

$$t \gg \tau$$
: $\Pi = \frac{-3H\zeta}{1 + \frac{3}{2}H\tau} \approx -3H\zeta_{\text{eff}}$

If relaxation time is comparable or larger than Hubble time, 2nd order required, but then fluid assumption is very questionable



Disconzi, Kephart & Scherrer 2015 based on Lichnerowitz 1944, 1957

Why should we expand in gradients of u^{μ} ?

claim:

gradient expansion in u^{μ} is reason for violation of causality, instead expand in

$$Fu^{\mu} \equiv \frac{\epsilon + p}{\mu} u^{\mu}, \quad \nabla_{\nu}(\mu u^{\nu}) \equiv 0$$

for dissipationless fluids $\mu = \epsilon + p$

now:
$$\Pi = -\dot{F}\zeta - F3H\zeta$$

causal at 1st order, promising idea!

Standard model or unified dark matter

SM assumes no dissipation: $p = -\epsilon$ no fluctuations1 - cosmological constant $p = -\epsilon$ no fluctuations2 - cold dark matterp = 0thermal origin (?), free fall3 - atoms (visible matter)p = 0gas,
viscosity in astrophysics

Could the dark sector be unified?

Generalized Chaplygin gas: $p = -A\epsilon_0(\epsilon_0/\epsilon)^{\alpha}$

Kamenshchik, Moschella & Pasquier 2001

Viscous dark fluid: $p = -3H\zeta$, $\zeta = \zeta_0(\epsilon_0/\epsilon)^{\nu}$, $H = H_0(\epsilon_0/\epsilon)^{1/2}$

Zimdahl, Schwarz, Balakin & Pavon 2001; Fabris, Goncalves & de Sa Ribeiro 2006

Unified dark matter

At background level, both models are equivalent and include the standard model ($\alpha = 0$; $\nu = -1/2$)



Dissipative fluctuations

But dissipative terms in VDF model are different from (dissipationless) GCG perturbations.

$$\psi'' + 3\mathcal{H}\psi' = \qquad \qquad \text{Velten \& Schwarz 2011}$$

$$w_{v} \left[\left[-\frac{1}{2} + \frac{k^{2}}{(1+w_{v})9\mathcal{H}^{2}} \right] 3\mathcal{H}\psi' + \left[\frac{3\mathcal{H}^{2}}{2} + \frac{k^{2}}{1+w_{v}} \right] \psi + \frac{3\mathcal{H}^{2}}{2} \frac{\delta\zeta}{\zeta} \right]$$

equation for metric potential at $z \ll 1000$

rhs gives rise to observable modifications wrt Λ CDM and allows us to distinguish VDF ($\delta \zeta \neq 0$) from GCG ($\delta \zeta = 0$)

Observable effects |

Both models give rise to a modification of the late time integrated Sachs-Wolfe effect Li & Barrow 2009, Velten & Schwarz 2011



$$\left(\frac{\delta T}{T}\right)_{\rm ISW} = 2 \int_{\rm dec}^{\rm today} {\rm d}\eta \,\partial_\eta \psi$$



Planck 2015

Observable effects ||

The VDF gives also rise to an extra damping of small scale structure at late times Velten & Schwarz 2011



Sub-galactic scales are damped exponentially might be useful to smooth galactic cores and to reduce number of satellites and dwarfs

Viscous dark matter

Hot and cold dark matter behave differently:

relativistic neutrinos:

Weinberg 1971, Straumann 1976

$$c_{\rm s}^2 = \frac{1}{3}, \quad \eta \approx \frac{4}{15}\epsilon\tau, \quad \zeta \approx 0, \quad \chi \approx \frac{4}{3}\frac{\epsilon}{T}\tau, \quad \tau \approx \tau_{\rm coll}$$

WIMPs in radiation background:

Hofmann, Schwarz & Stoecker 2001; Green, Hofmann & Schwarz 2004,2005; Bringmann & Hofmann, 2006

$$c_{\rm s}^2 \approx \frac{5}{3} \frac{T}{m}, \quad \eta \approx nT\tau, \quad \zeta \approx \frac{3}{5}nT\tau, \quad \chi \approx 0; \quad \tau \approx \sqrt{\frac{2}{3}} \frac{m}{T}\tau_{\rm coll}$$

 $\Rightarrow L \approx \frac{3}{2} \frac{T}{m} \tau \qquad \text{damping length}$

both examples relevant in early universe

Viscous dark matter

study a AvCDM model

viscosity in todays universe?



Velten & Schwarz 2012; Velten, Schwarz, Fabris & Zimdahl 2013

Viscous dark matter

subhorizon vCDM equation: scale dependent terms !

$$\begin{aligned} a^{2} \frac{d^{2} \Delta_{v}}{da^{2}} + \left[\frac{a}{H} \frac{dH}{da} + 3 + A(a) + B(a)k^{2}\right] a \frac{d\Delta_{v}}{da} + \left[+C(a) + D(a)k^{2} - \frac{3}{2}\right] \Delta_{v} &= P(a), \\ A(a) &= -6w_{v} + \frac{a}{1+w_{v}} \frac{dw_{v}}{da} - \frac{2a}{1+2w_{v}} \frac{dw_{v}}{da} + \frac{3w_{v}}{2(1+w_{v})} \\ B(a) &= -\frac{w_{v}}{3a^{2}H^{2}(1+w_{v})} \\ C(a) &= \frac{3w_{v}}{2(1+w_{v})} - 3w_{v} - 9w_{v}^{2} - \frac{3w_{v}^{2}}{1+w_{v}} \left(1 + \frac{a}{H} \frac{dH}{da}\right) - 3a \left(\frac{1+2w_{v}}{1+w_{v}}\right) \frac{dw_{v}}{da} + \frac{6aw_{v}}{1+2w_{v}} \frac{dw_{v}}{da} \\ D(a) &= \frac{w_{v}^{2}}{a^{2}H^{2}(1+w_{v})} \\ P(a) &= -3\nu w_{v}a \frac{d\Delta_{v}}{da} + 3\nu w_{v}\Delta_{v} \left[-\frac{1}{2} + \frac{9w_{v}}{2} + \frac{-1-4w_{v}+2w_{v}^{2}}{w_{v}(1+w_{v})(1+2w_{v})}a \frac{dw_{v}}{da} - \frac{k^{2}(1-w_{v})}{3H^{2}a^{2}(1+w_{v})}\right], \end{aligned}$$

CDM: Meszaros equation for A = B = C = D = P = 0

Velten & Schwarz 2012

Constrain on dark matter bulk viscosity



existance of dwarf galaxies excluded unless

 $\zeta_0 < 10^{-14} H_0/G$

consistency check: $\zeta_0 \sim \varepsilon_v \tau \ll H_0/G$ thus $\tau H_0 \ll \varepsilon_0/\varepsilon_v \sim I$

Velten & Schwarz 2012

Open issues

study the more general cases $\zeta = \zeta(\varepsilon, S, ...)$

- could huge ISW contribution be avoided?
- put quantitative limits on unified dark matter from final Planck release and LSS surveys
 update limits on ζ and η of dark matter from final Planck release and LSS surveys
 calculate ζ and η for realistic dark matter candidates
 (so far done for generic WIMP and SDM)



Unified dark matter/viscous inflation require unrealistic assumptions on bulk viscosity/relaxation times

Alternative solution to causality problem?

Viscous dark matter is a real option, but bulk viscosity must be small and is in accord with hydrodynamic assumption

Constraints on viscous cosmology from ISW (CMB - LSS xcorrelation) and dwarf galaxies

Is there a good candidate for VDM? Not WIMPs!