



Quasielastic Scattering at MiniBooNE Energies

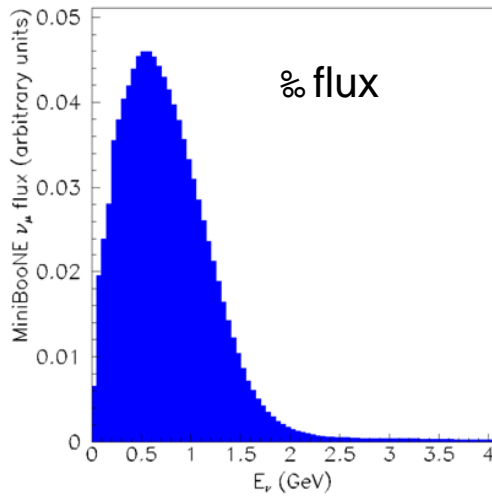
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1. U. Murcia

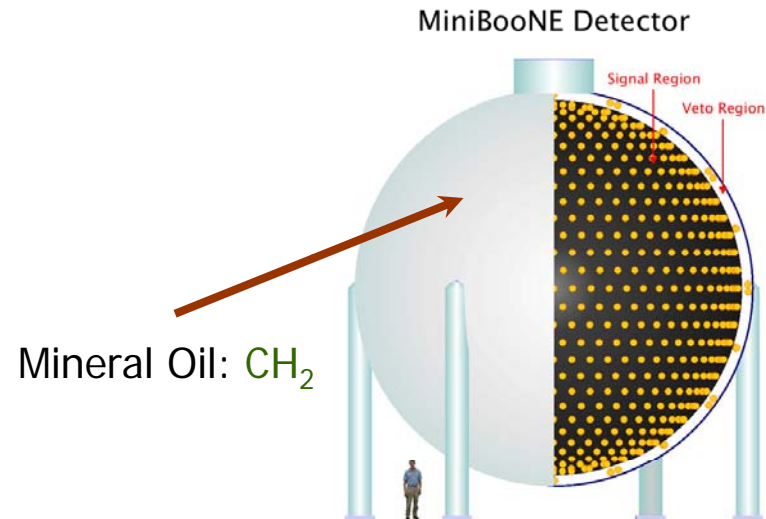
2. U. Giessen

Introduction

- MiniBooNE has collected the largest sample of low energy ν_{μ} CCQE events to date. Aguilar-Arevalo et. al., PRL 100 (2008) 032301



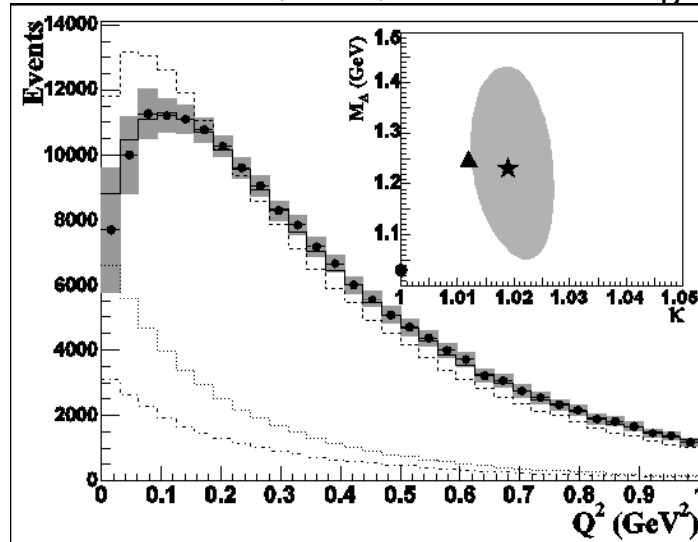
$$\langle E_{\nu} \rangle \sim 750 \text{ MeV}$$



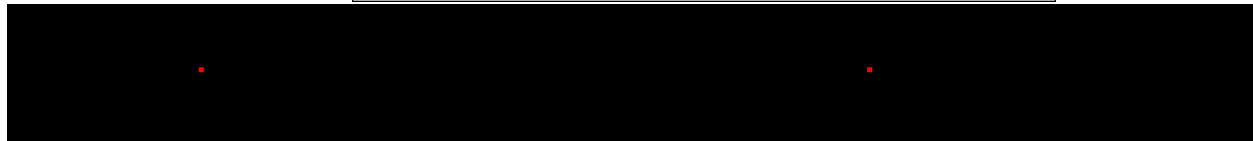
- CCQE is relevant for **oscillation** experiments
- CCQE is interesting by itself:
 - **Axial form factor** of the nucleon (M_A)
 - Nuclear **correlations**

Introduction

- The shape of $k d^2/d\cos\theta_{lab} dE_{lab}$ is accurately described by a **Global Fermi Gas Model** Smith, Moniz, NPB 43 (1972) 605 with: $E_B = 34$ MeV, $p_F = 220$ MeV



But



- $M_A = 1.23 \ll 0.20$ GeV
consistent with $M_A = 1.2 \ll 0.12$ GeV (K2K)
higher than $M_A = 1.01 \ll 0.01$ GeV (0d ; 0p ; $(e, e^{01}/4)$; BBBA07)
 $M_A = 1.05 \ll 0.08$ GeV (mainly ^{12}C , 3-100 GeV, NOMAD)

The model

- **Our aim:** realistic description of **CCQE** in nuclei
compare to **MiniBooNE** data (modified Smith-Moniz ansatz)
- Relevant hadronic degrees of freedom: $\pi, N, f(1232)$
- Ingredients:
 - Elementary process $\nu, n \rightarrow l, p$
 - Fermi motion
 - Pauli blocking
 - Nuclear binding
 - Nucleon spectral functions
 - Medium polarization (RPA)

and also

- Non **CCQE** background (cannot be separated from **CCQE** in a model independent way)

The model

Elementary amplitude for $\bar{p}_1(k) n(p) \rightarrow \bar{p}_1(k') p(p')$

$$i \mathcal{M}^{\otimes} J^{\otimes}$$

where $i \mathcal{M}^{\otimes}$ and

$$J^{\otimes} = F_1^V \quad F_2^V \quad F_A \quad F_P$$

$$F_{1;2}^V \leftarrow \text{BBBA07} \quad Q^2 = i(k_i - k^0)^2$$

$$F_A \quad F_P \quad F_A \leftarrow \text{PCAC}$$

$$g_A = 1.26 \leftarrow \text{neutron } \beta \text{ decay} \quad M_A = 1 \text{ GeV}$$

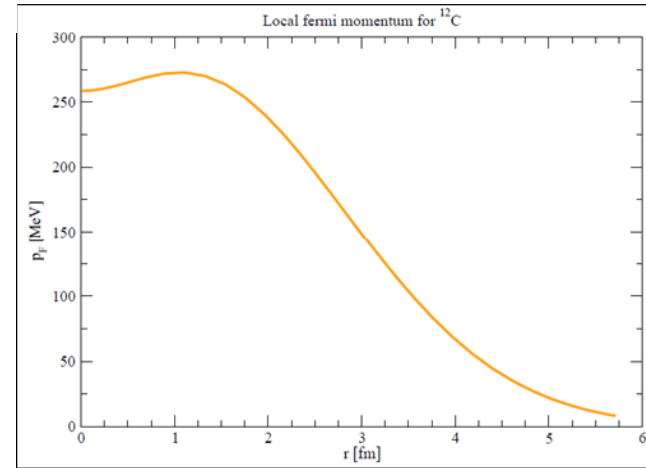
Using PCAC and ρ -pole dominance: $\dots \leftarrow \text{Goldberger-Treiman}$

The model

■ Local Fermi Gas

- Fermi Motion of initial nucleons:

- Pauli blocking of final nucleons:



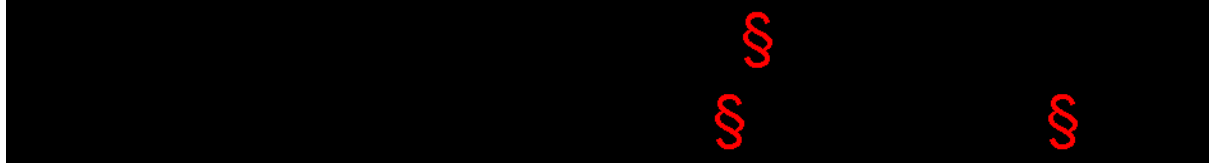
■ Mean field potential

- Density and momentum dependent
- Parameters fixed in p -Nucleus scattering [Teis et. al., Z. Phys. A 346 \(1997\) 421](#)
- Nucleons acquire **effective masses**

The model

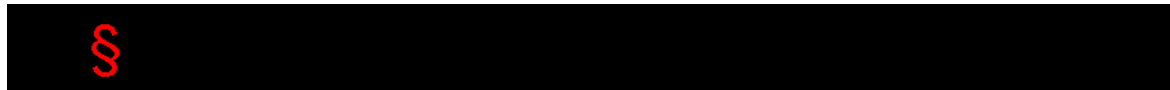
■ Spectral functions

- The momentum distribution of a nucleon with 4-momentum p is



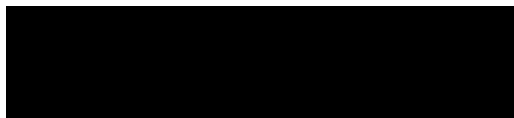
§ ← nucleon selfenergy

- For final nucleons (above the Fermi sea)



← collisional broadening
GiBUU

- § is obtained from § with a dispersion relation assuming that at the pole position:



- For initial nucleons (holes) we take §



The model

- RPA nuclear correlations

"In nuclei, the strength of electroweak couplings may change from their free nucleon values due to the presence of strongly interacting nucleons"

Singh, Oset, NPA 542 (1992) 587

- For the **axial coupling** g_A :



χ_3 dipole susceptibility

g' Lorentz-Lorenz factor $\sim 1/3$

Ericson, Weise, Pions in Nuclei

- The **quenching of g_A** in Gamow-Teller \dagger **decay** is well established

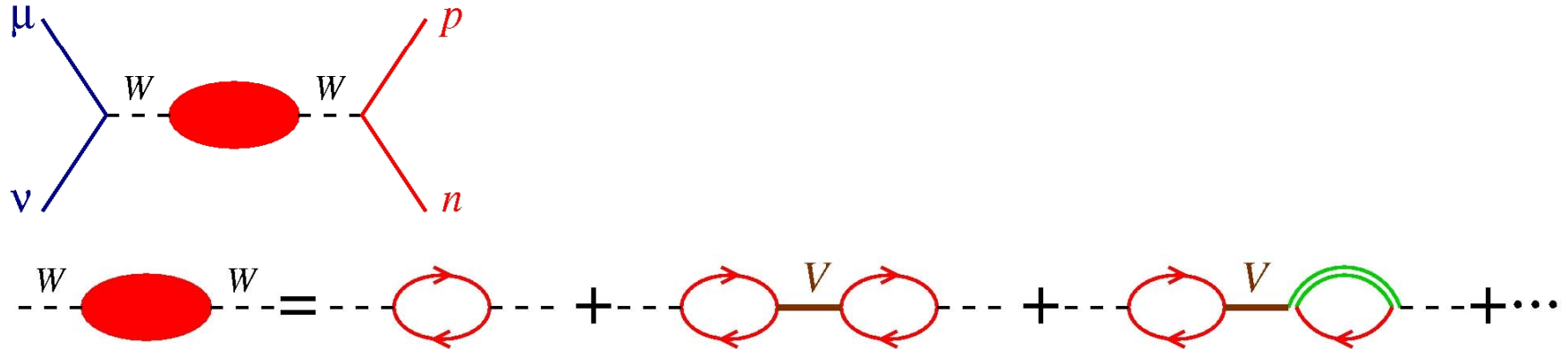


Wilkinson, NPA 209 (1973) 470

The model

- RPA nuclear correlations

- Following Nieves et. al. PRC 70 (2004) 055503 :



V_{NN}

- In particular



- spectral function **changes** in the **nuclear medium** \rightarrow so does J^A



The model

■ RPA nuclear correlations

- **RPA approach** built up with single-particle states in a Fermi sea
- Simplified vs. some theoretical models (e.g. continuum RPA)
- Applies to inclusive processes; not suitable for transitions to discrete states

But

- Incorporates explicitly \bar{p} and \bar{n} exchange and f -hole states
- Can be inserted in a unified framework to study **QE**, $1\bar{n}$, **N-knockout**, etc
- Has been successfully applied to \bar{p} , \bar{n} and **electro**-nuclear reactions
- Describes correctly ... capture on ^{12}C and **LSND** CCQE

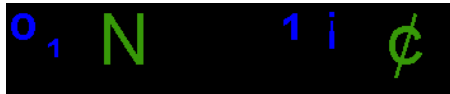
Nieves et. al. PRC 70 (2004) 055503

- **Important at low Q^2** at **MiniBooNE** energies

The model

- Non CCQE background (GiBUU transport model)

- Most relevant processes:



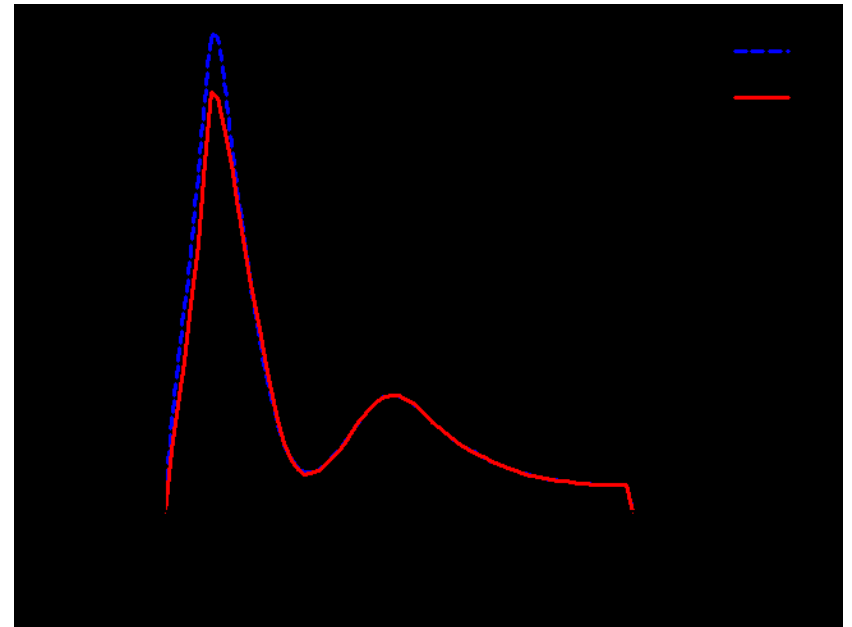
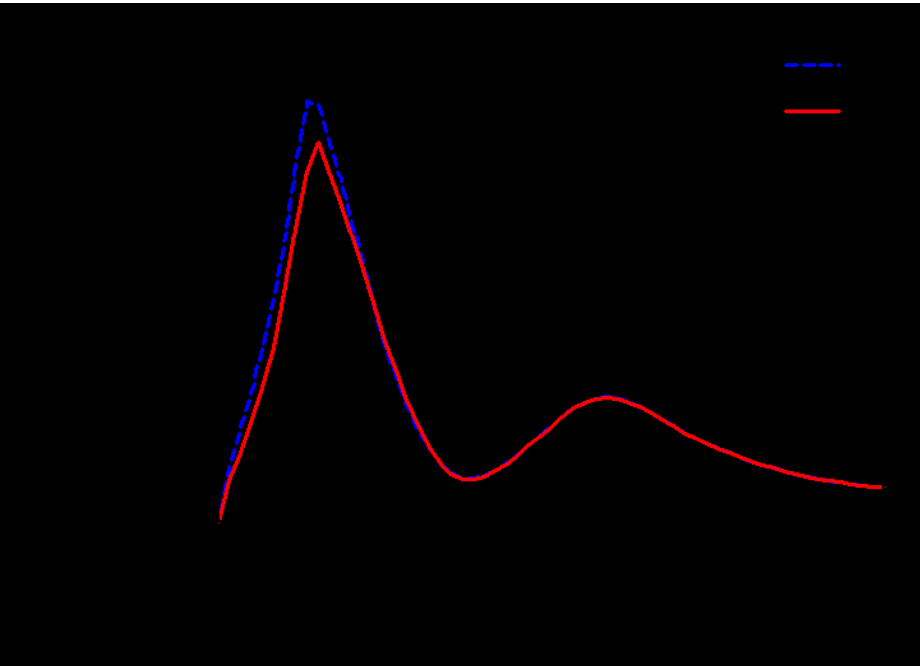
followed by $\phi N \rightarrow N N$ (\sim less decay mode)

or by $\phi N \rightarrow N \frac{1}{4}$ and then $\frac{1}{4} N N \rightarrow N N$ (\sim absorption)

- Details in the talk by Tina Leitner

Results

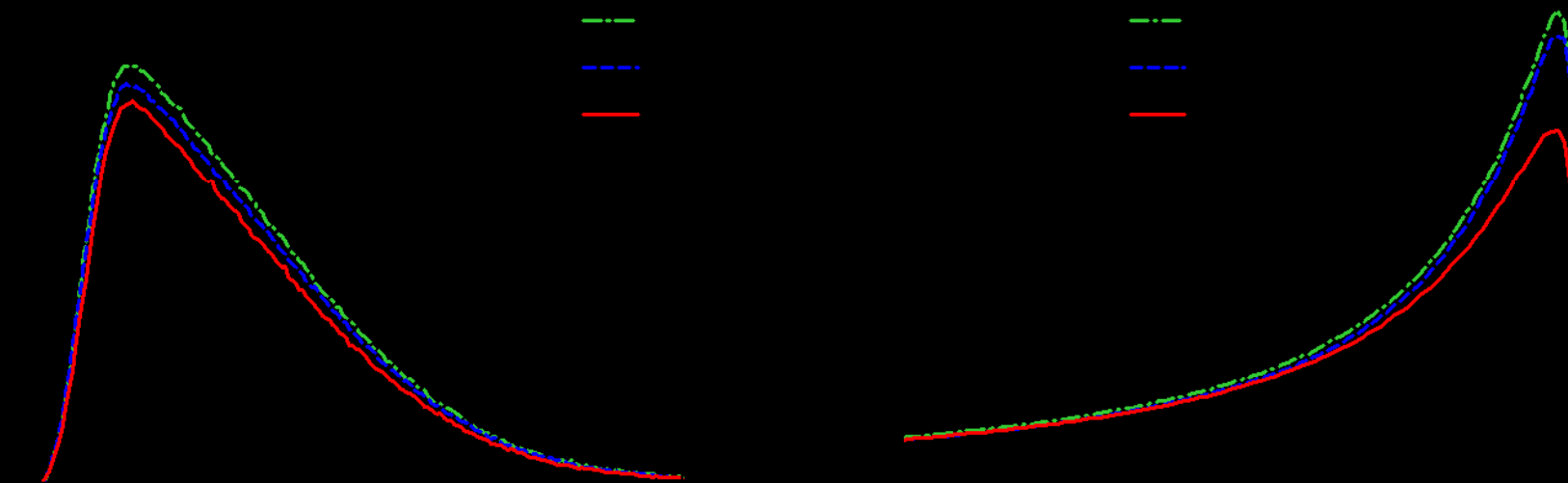
- Comparison to inclusive electron scattering data



- Missing strength in the **deep region**:
more complete many-body dynamics required [Gil et. al., NPA 627 \(1997\) 543](#)
- RPA less relevant than in the weak case

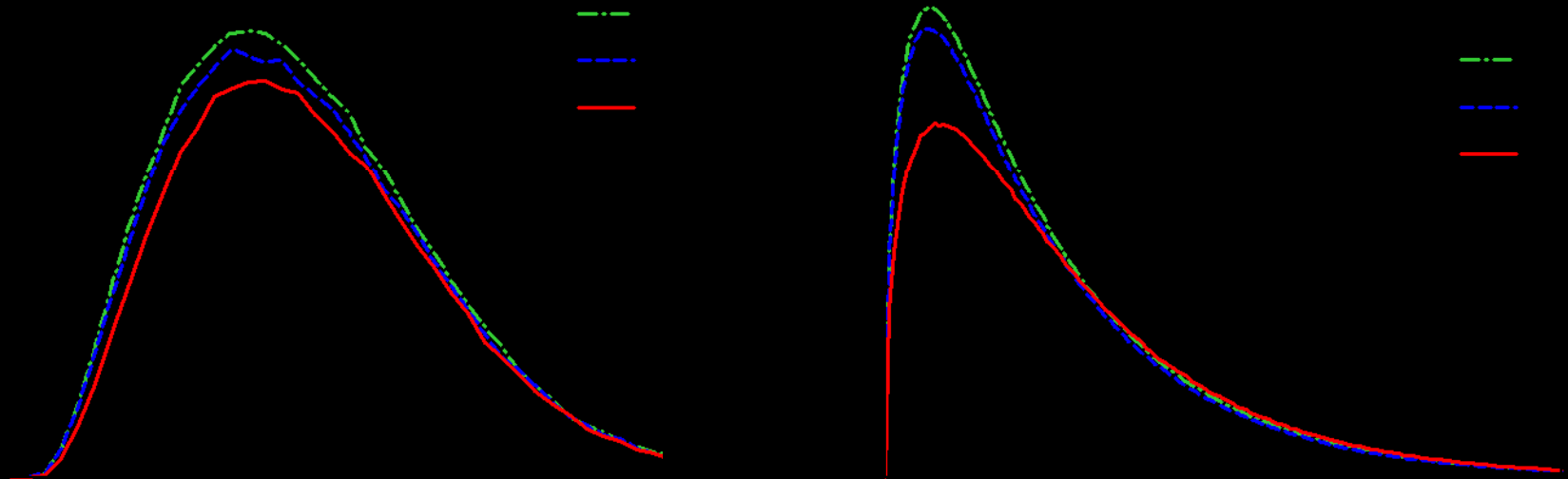
Results

- Differential cross sections averaged over the MiniBooNE flux



Results

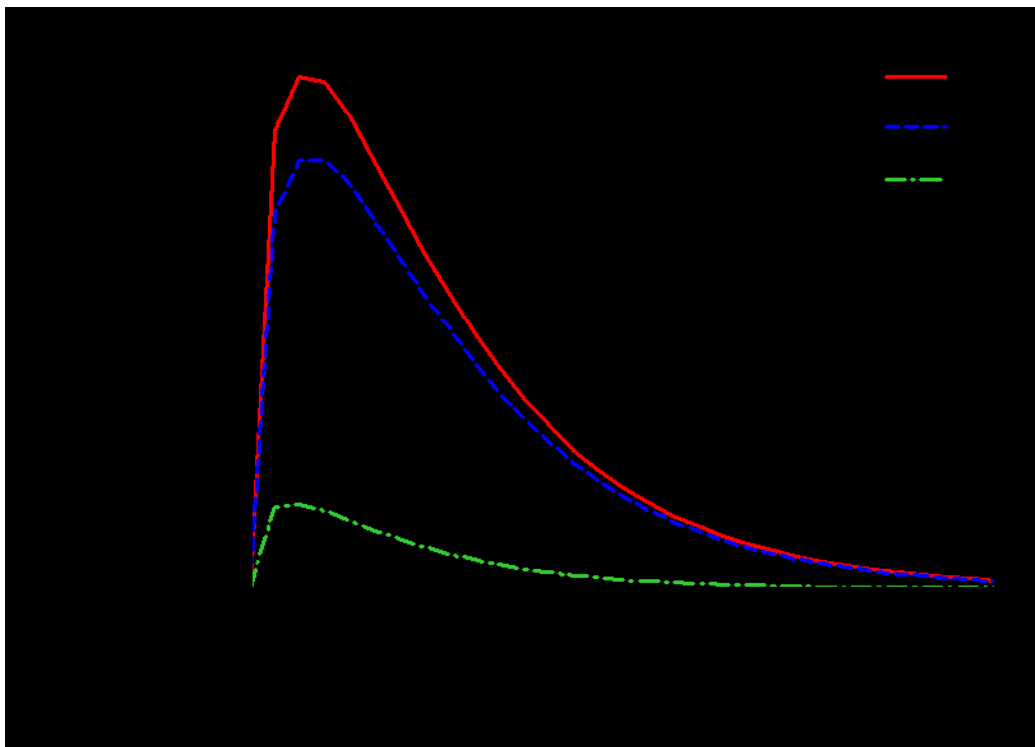
- Differential cross sections averaged over the MiniBooNE flux



- RPA correlations cause a **considerable reduction** of the c.s. at **low Q^2** and **forward angles**

Results

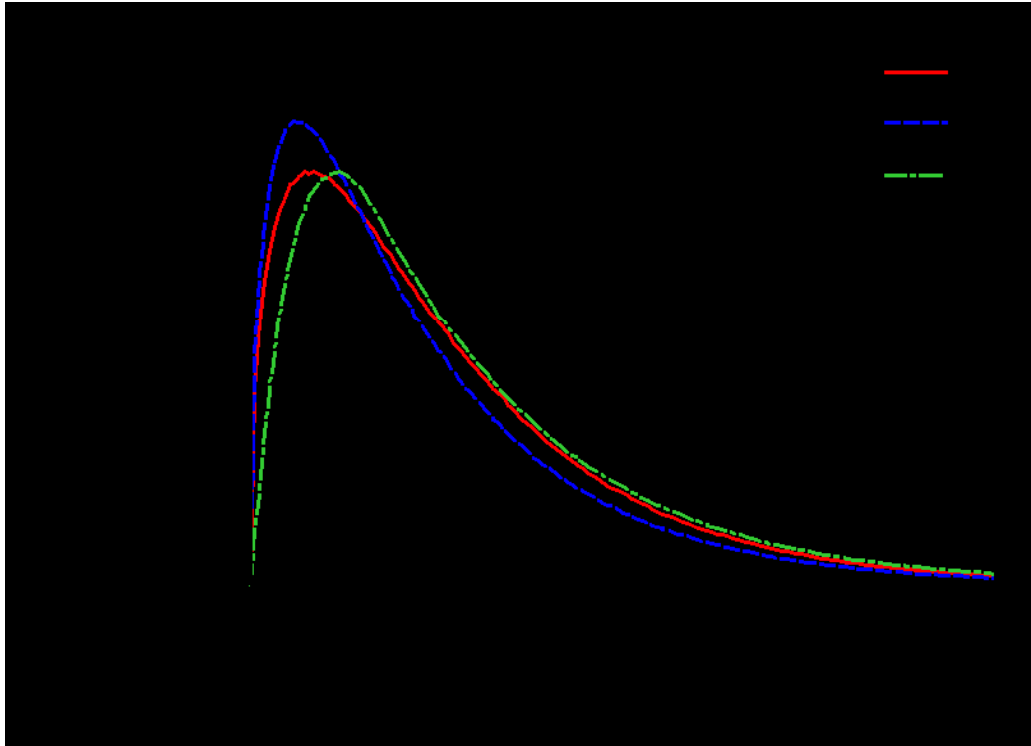
- CCQE + non-CCQE background (measured quantity)



non $\bar{\nu}_e$ CCQE
total CCQE $\bar{\nu}_e$ like

Results

- Comparison to the modified Smith-Moniz ansatz (shape)

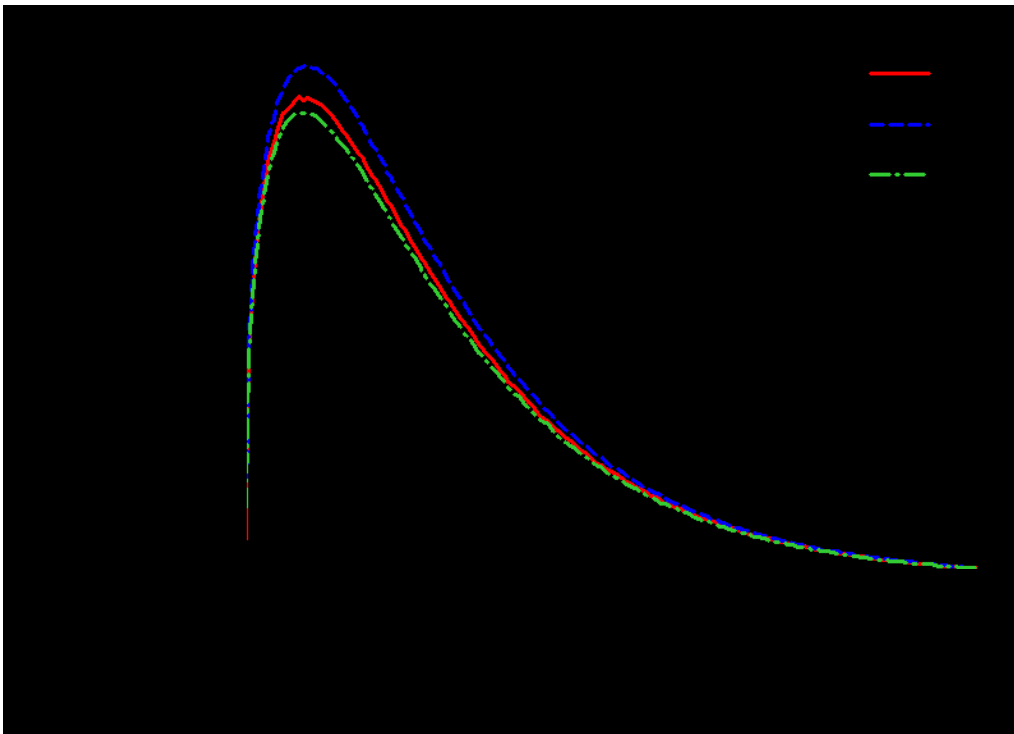


All curves are normalized
to the **same area**

- The effect of **RPA** brings the **shape** of the Q^2 distribution closer to experiment keeping $M_A = 1 \text{ GeV}$

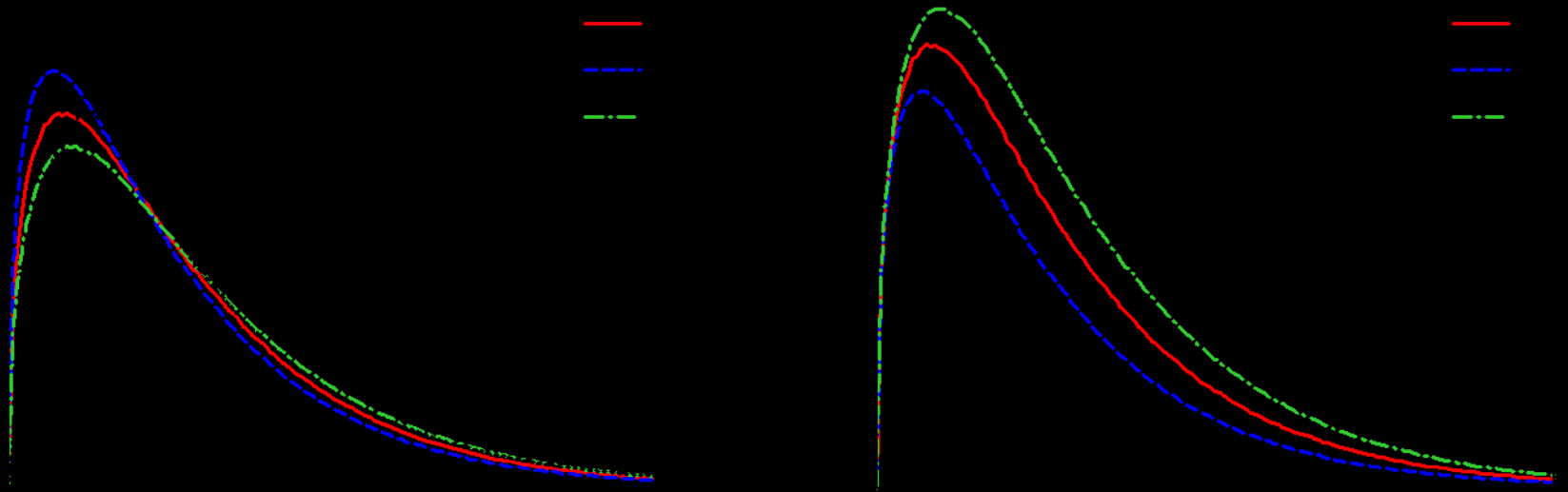
Results

- Changing $g' = 0.5-0.7$
 - Leaves the **shape unaltered**
 - **Small changes** in the integrated **cross section** $\hat{\kappa} \approx 3.1-3.4 \times 10^{-38} \text{ cm}^2$



Results

■ Changing M_A



- The experimental shape is better described with $M_A=1$ GeV

- There are large differences in the integrated cross section

M_A [GeV]	0.8	1.0	1.2
$k^{\perp 2}$ [μm^{-2}]	2.5	3.2	3.7

Conclusions

- We have developed a model for **CCQE** starting from a Local Fermi Gas but incorporating important many body corrections: **nucleon spectral functions**, medium polarization (**RPA**).
- **RPA**: strong reduction at low q^2 (quenching)
- Good description of the **shape** of the **MiniBooNE** q^2 distribution with **$M_A = 1$ GeV**
- A larger integrated CCQE cross section would require **$M_A > 1$ GeV**

Conclusions

- There is (nuclear) physics beyond the naive Fermi Gas Model.

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