

Do we theoretically
understand/control CC and NC
coherent pion production at low
and intermediate neutrino
energies?

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It depends on what we mean by **control/understand**. We have a deep understanding of the process and sophisticated calculations are available, though there still exist some uncertainties...

NC Rein–Sehgal model [NPB 223 (1983) 29]

- PCAC

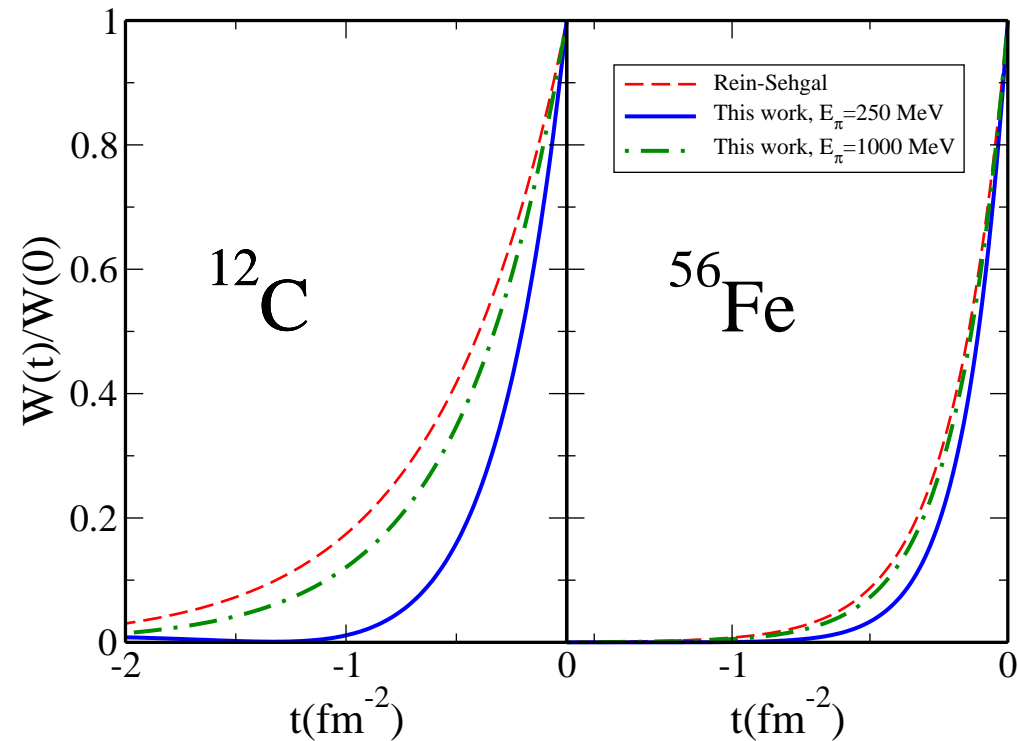
$$\left\langle \mathcal{N}_{gs} \pi^0(k_\pi) \left| q_\mu A_{NC}^\mu \right| \mathcal{N}_{gs} \right\rangle_{\mathbf{q}^2=0} = 2f_\pi \left\langle \mathcal{N}_{gs} \pi^0(k_\pi) \left| -iT \right| \pi^0(q) \mathcal{N}_{gs} \right\rangle_{\mathbf{q}^2=0}$$

elastic $\pi^0 \mathcal{N}_{gs}$ Xsect and pion distortion effects...?

- $t = (q - p_\pi)^2 = 0$; $q^2 = (k_\nu - k'_\nu)^2 = 0$ approximations

$$\frac{d\sigma}{dq^2 dq^0 dt} \propto \frac{1}{\left(1 - \frac{q^2}{1 \text{ GeV}^2}\right)^2} \left(\left| \mathbf{F}_{\mathcal{A}}(\mathbf{t}) \right|^2 \mathbf{F}_{\text{abs}} \frac{d\sigma(\pi^0 N \rightarrow \pi^0 N)}{dt} \Big|_{\mathbf{E}_\pi = \mathbf{q}^0, \mathbf{t}=0} \right)$$

$$F_{\mathcal{A}}(t) = \int d^3 \vec{r} e^{i(\vec{q} - \vec{k}_\pi) \cdot \vec{r}} \{ \rho_p(\vec{r}) + \rho_n(\vec{r}) \}$$



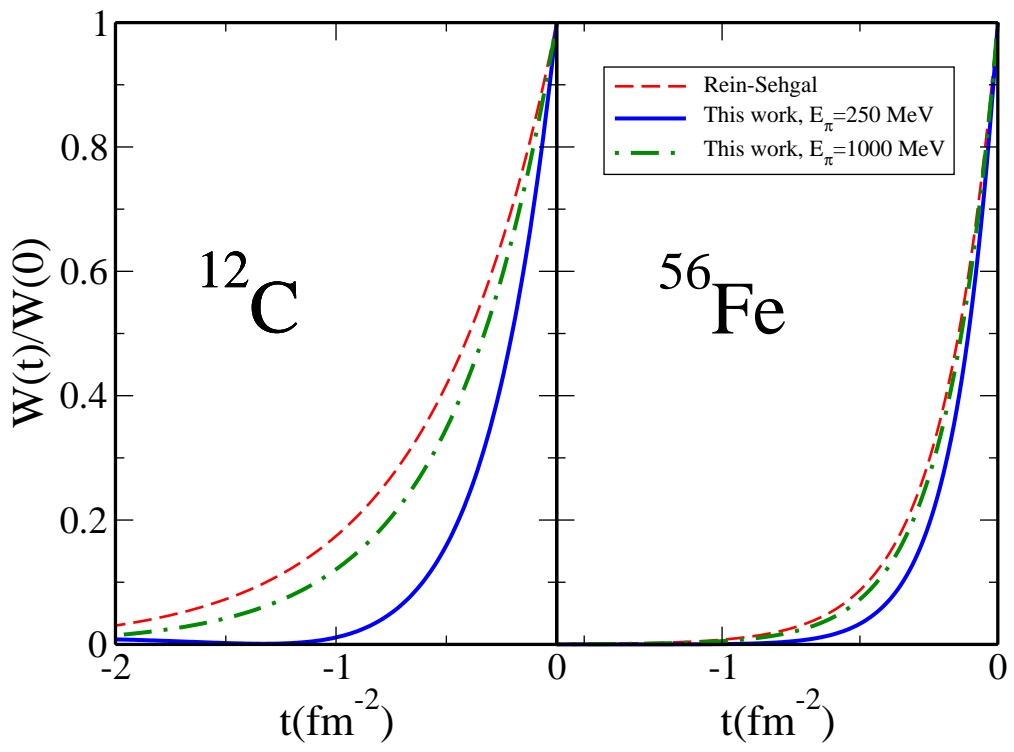
(Hernández, Nieves and Vicente-Vacas, arXiv:0903.5285)

$$W(t) = |F_{\mathcal{A}}(t)|^2 \frac{d\sigma_{nsf}(\pi^0 N \rightarrow \pi^0 N)}{dt},$$

Large (and negative) q^2 values are suppressed by the elastic pion–nucleus differential cross section that strongly favours $t = 0$.

(Hernández, Nieves and Vicente–Vacas, arXiv:0903.5285)

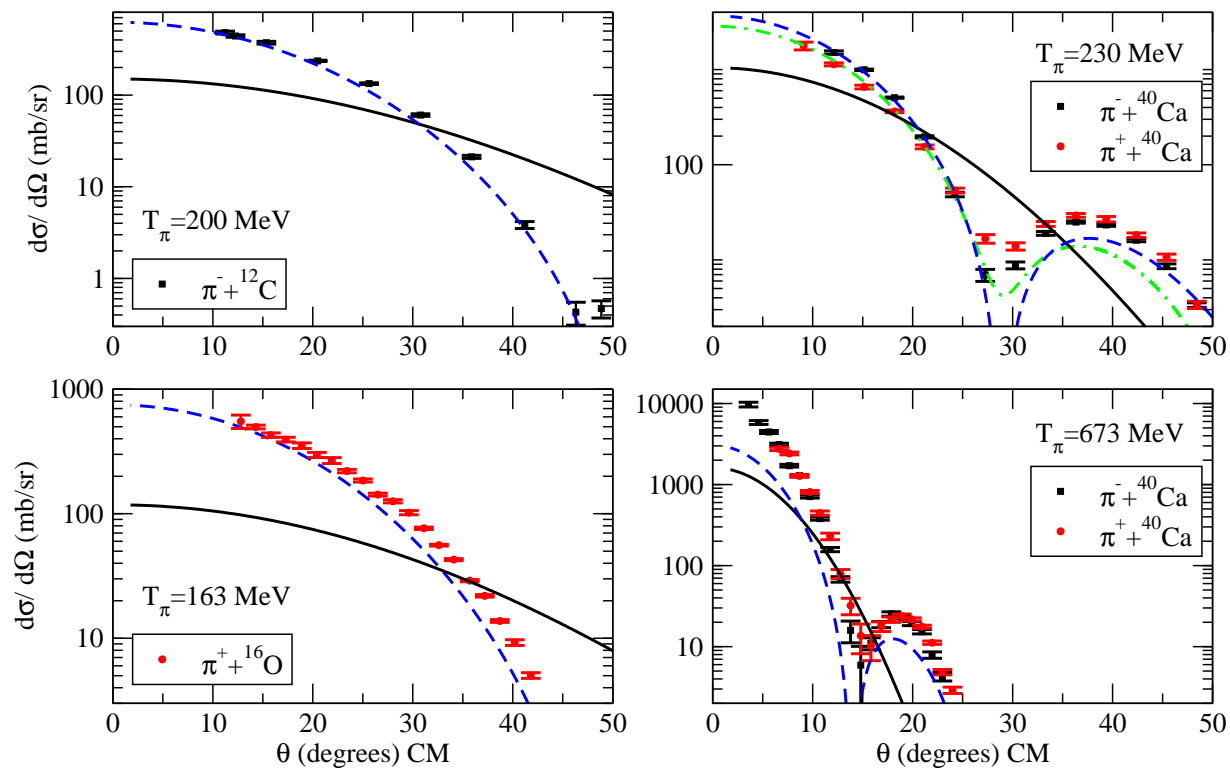
$$W^{\text{RS}}(t) = |F_{\mathcal{A}}(t)|^2 \frac{d\sigma_{nsf}(\pi^0 N \rightarrow \pi^0 N)}{dt} \Big|_{t=0}$$



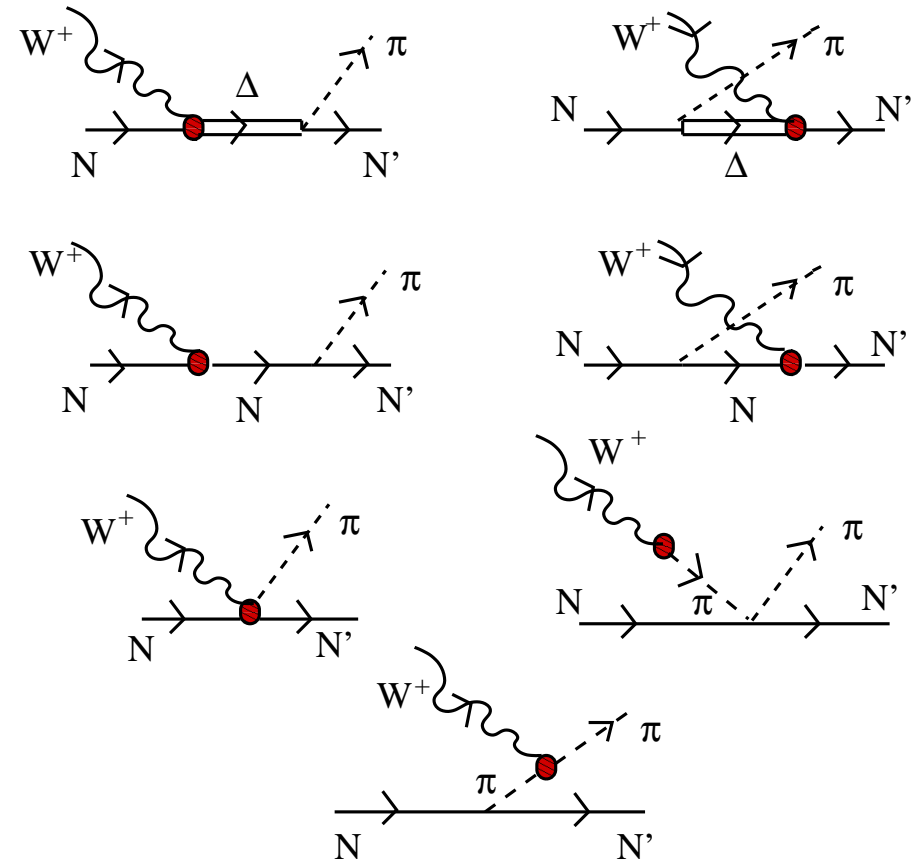
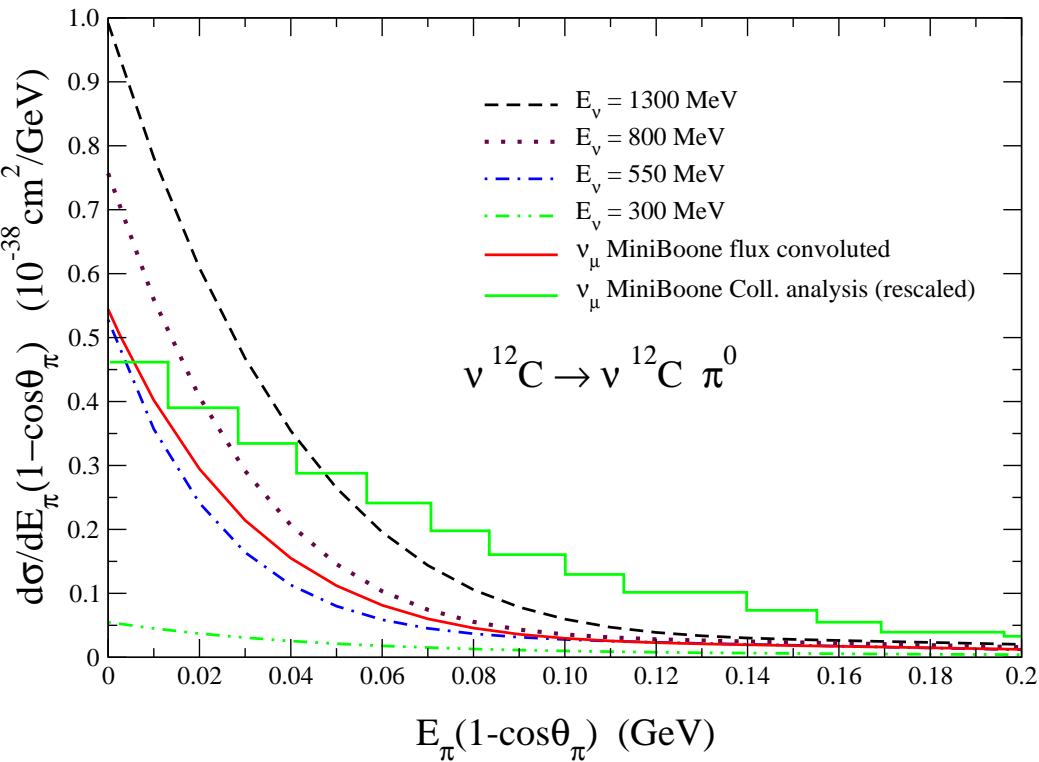
Original RS work: medium size nucleus, aluminum, and neutrino energies above 2 GeV ⇒ OK

However, for neutrino energies below 1 GeV and lighter nuclei, like carbon or oxygen, the nuclear form factor is **not enough forward peaked to render the finite t -dependence of the pion–nucleon cross section negligible**, and even in the forward direction the t value is not close enough to zero. **$q^2 \neq 0$ contributions beyond PCAC and further $t \neq 0$ dependences neglected within the RS model turn out to be now relevant.**

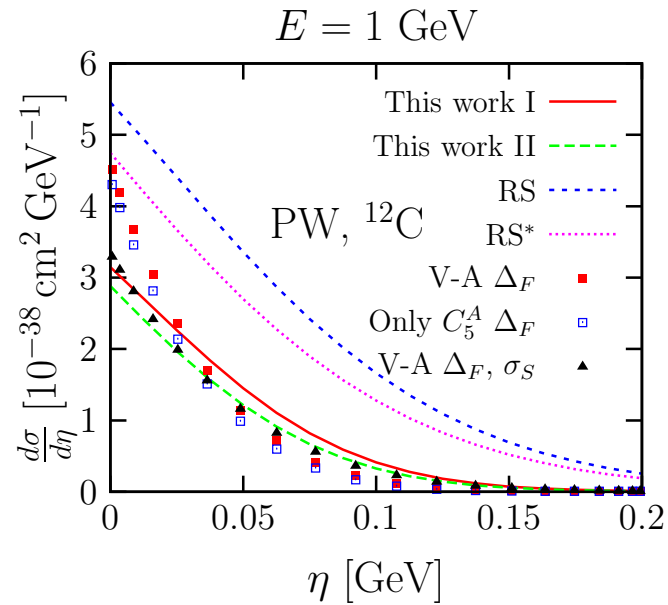
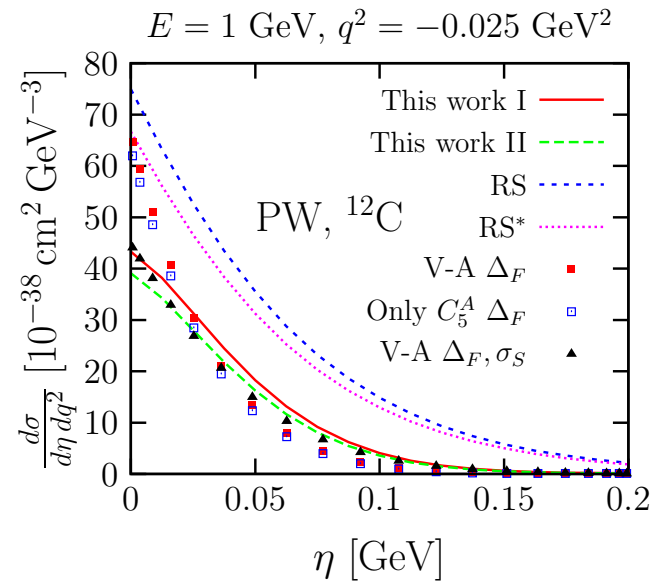
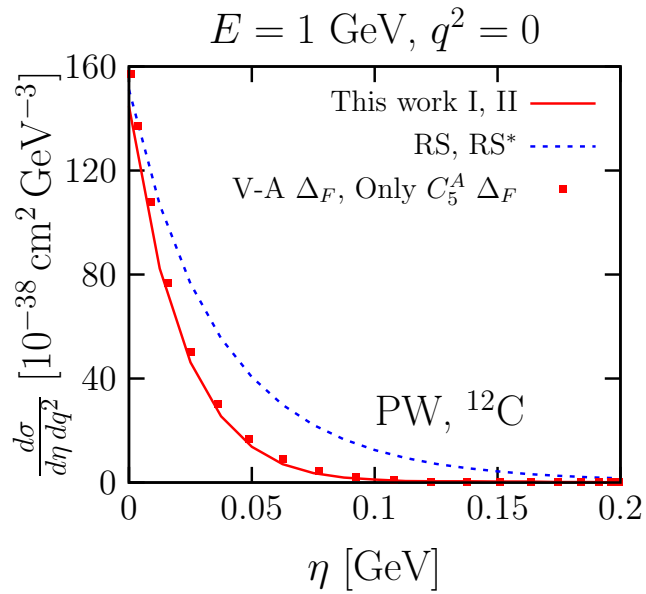
Thus at low neutrino energies, for instance those relevant in MiniBooNE, RS model would predict **elastic $\pi\mathcal{N}_{gs}$ Xsects** strongly disfavored by data ...



(Amaro, Hernández, Nieves and Valverde, PRD 79: 013002)



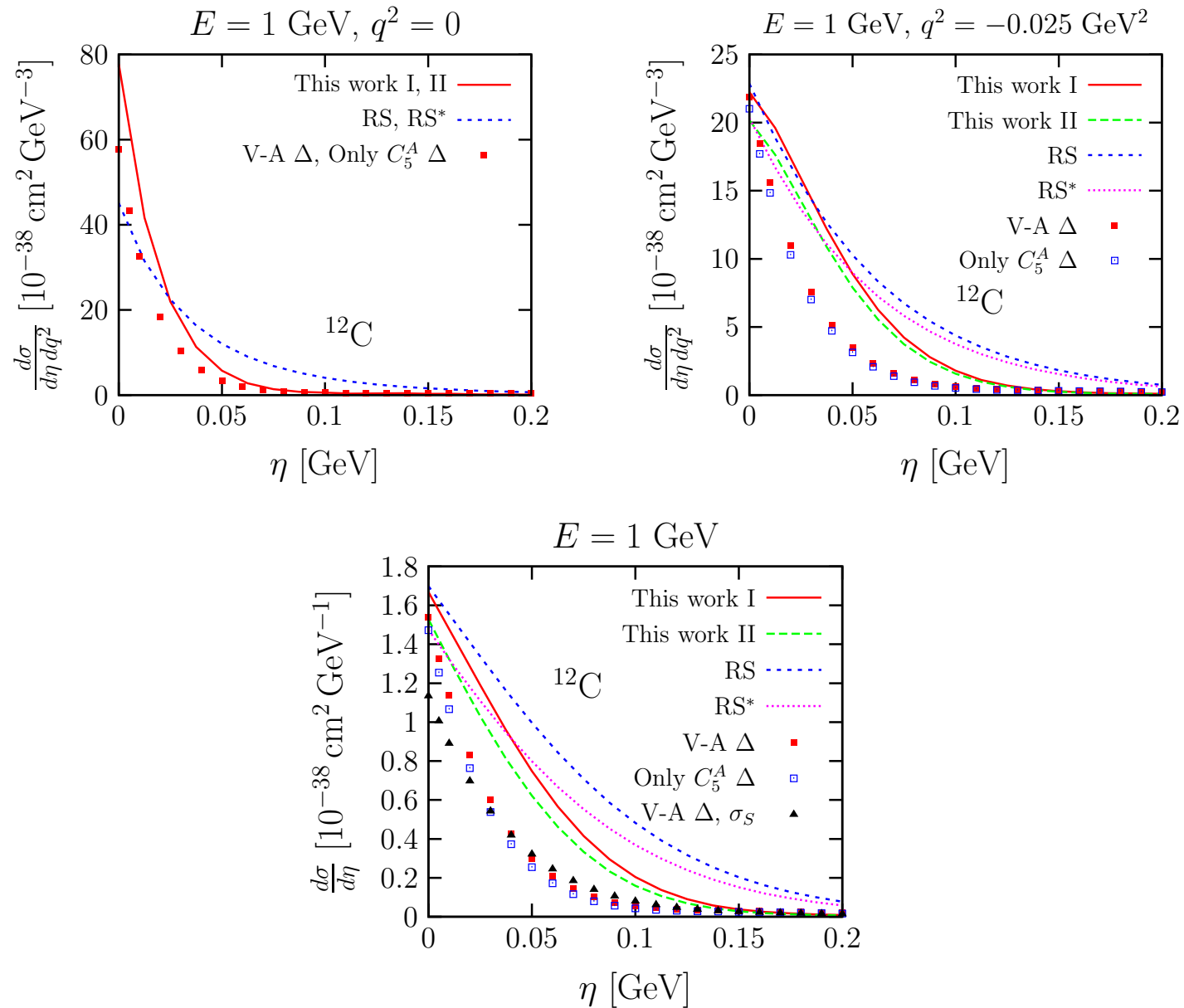
CC vs NC ?



Remarks:

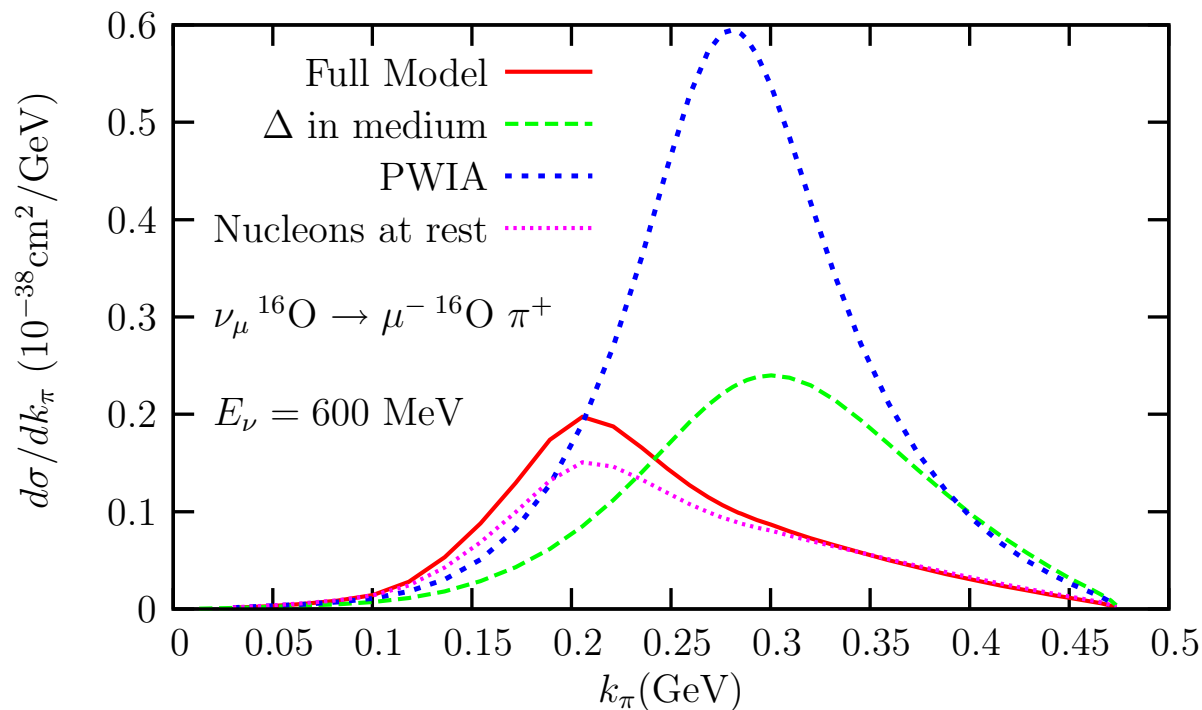
- No distortion for outgoing pion
- RS*: RS model including some kinematical corrections that vanish at $q^2 = 0$ and recently proposed by Berger and Sehgal (PRD 79: 053003)
- $L_{\mu\nu}H^{\mu\nu} \neq q_\mu q_\nu A^\mu A^\nu$ except for $q^2 = 0$
 - PCAC based models only work when the $q^2 = 0$ contribution is totally dominant.
 - When $q^2 \neq 0$ contributions cannot be neglected, PCAC models cannot provide the incoming neutrino–outgoing pion angular distribution.

Once the pion distortion is taken into account...



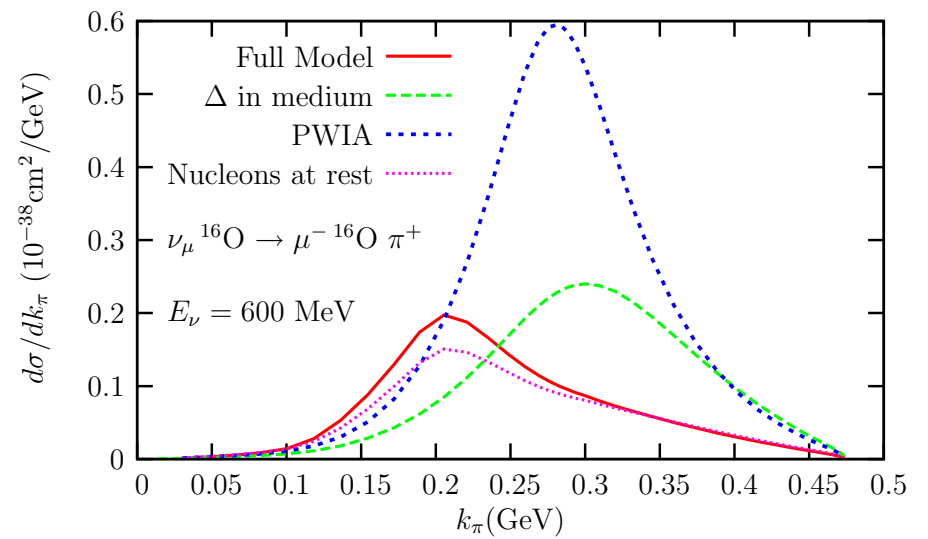
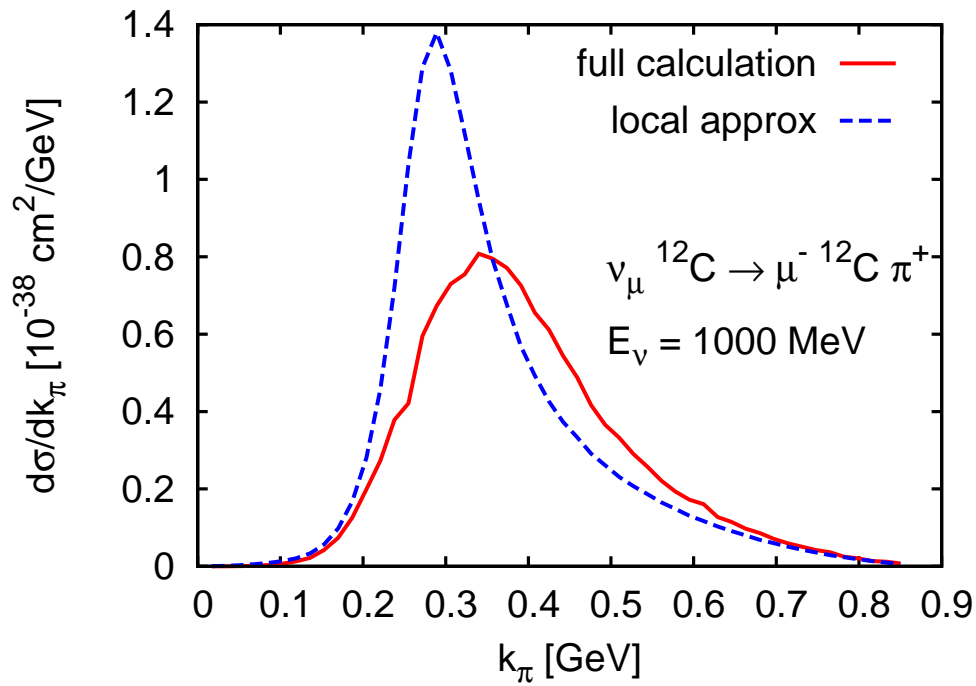
Microscopical models

- Singh, Athar and Ahmad, PRL 96: 241801
- Alvarez-Ruso, Geng and Vicente-Vacas, PRC76:068501
- Amaro, Hernández, Nieves and Valverde, PRD 79: 013002

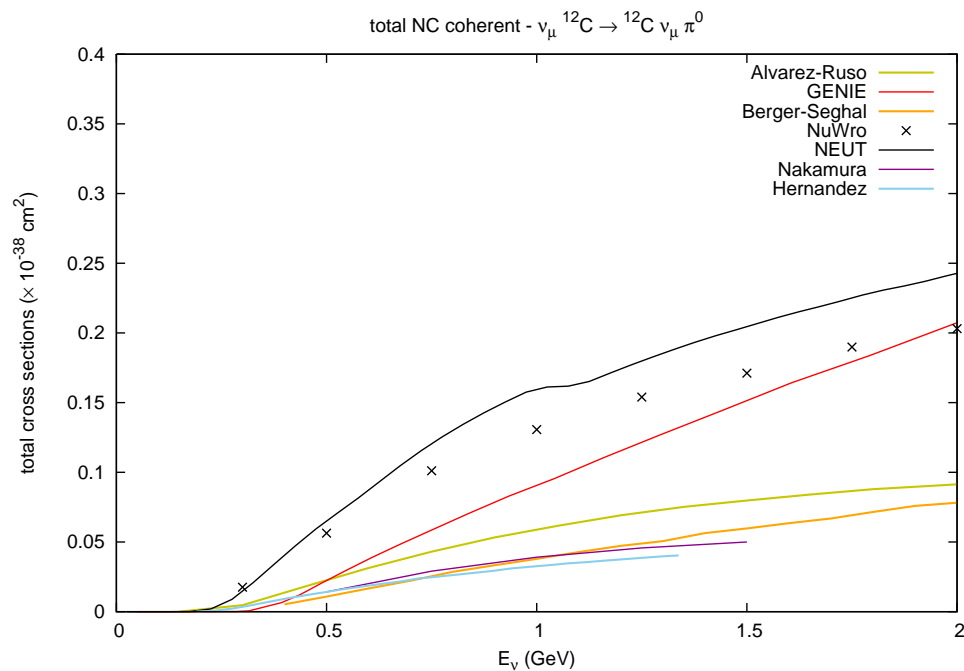


- Δ excitation dominant mechanism for $E_\nu \leq 1 \text{ GeV}$
- Many nuclear effects,
 - Δ -selfenergy in the medium
 - outgoing pion distortion
 - non-localities: momentum of the pion inside of the nucleus, momentum of the nucleons...
 - $C_5^A(q^2) = ?$ (axial ΔN transition vertex)

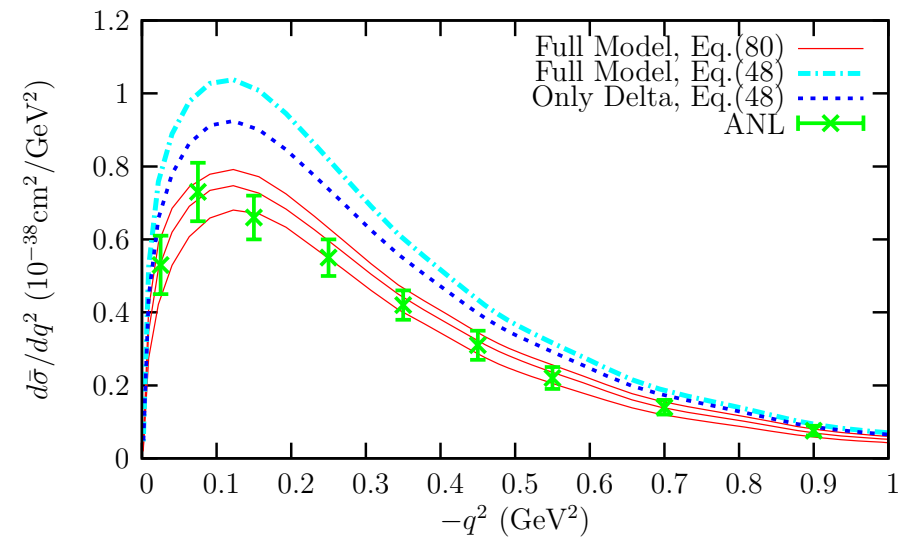
(Leitner, Mosel and Winkelmann, arXiv:0901.2837)



NUINT09



- $C_5^A(q^2)$ is crucial!! Uncertainties in Xsects of about a factor of 2.



- Beyond, let us say $E_\nu > 1.3 \text{ GeV}$ for NC (phase space peaks around $W_{\pi N} \sim 1.5 \text{ GeV}$), other resonances and mechanisms might play a significant role.
- Above $E_\nu > 2 \text{ GeV}$: PCAC based models should work.