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# **Calibration of the Pythia 8.2 Top Quark Mass Using 2-Jettiness in $e^+e^-$**

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# Outline

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- Introduction
- Monte Carlo generators and the top quark mass
- Calibration of the Monte Carlo top mass parameter
- Preliminary detailed results of first serious systematic analysis
- Summary, future plans

In collaboration with:

M. Butenschön

B. Dehnadi,

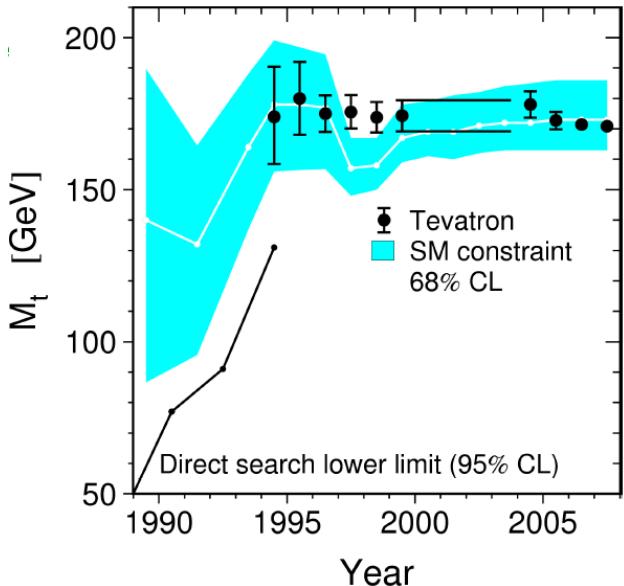
V. Mateu,

M. Preisser

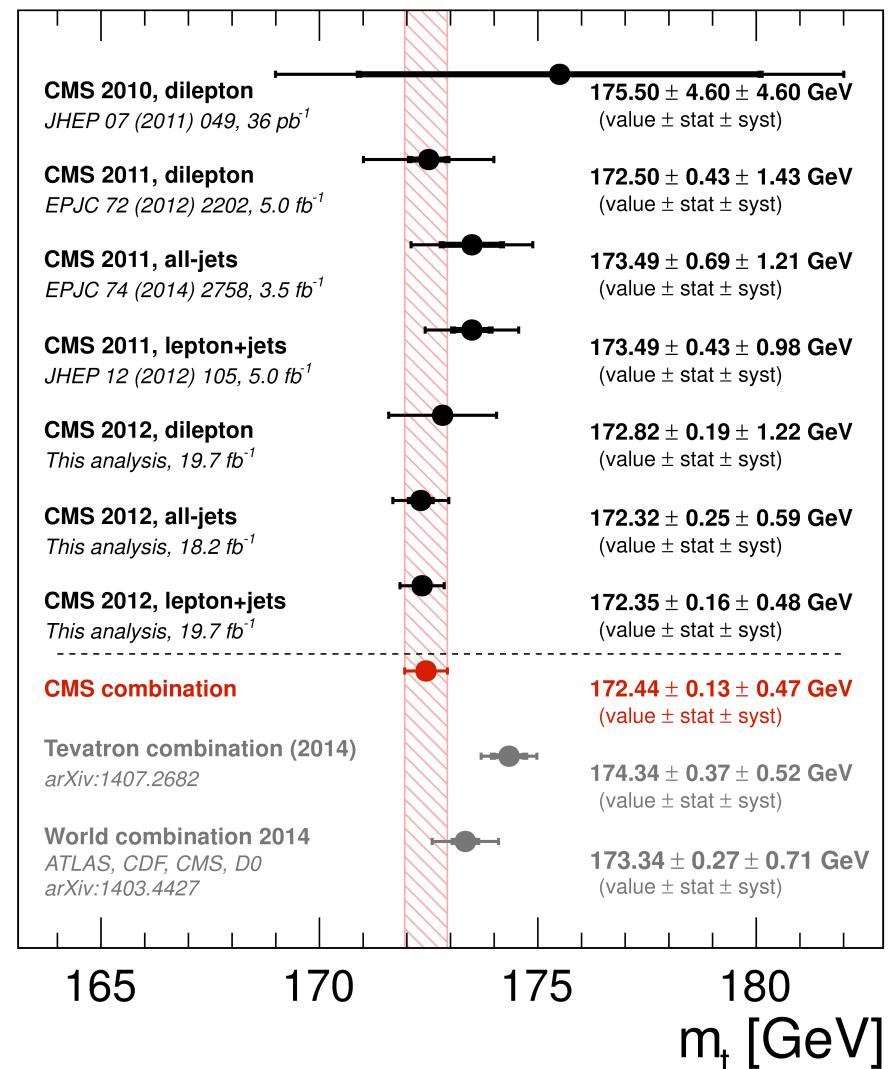
I. Stewart



# A small history on top mass reconstruction



- Many individual measurements with uncertainty below 1 GeV.
- Some discrepancies between LHC and Tevatron
- Reached <500MeV range.



# Main Top Mass Measurements Methods

[LHC+Tevatron](#)

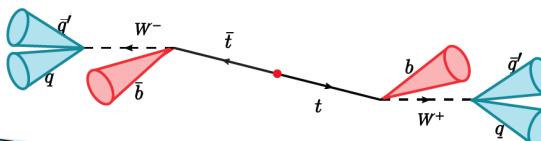
## Direct Reconstruction:

### Kinematic Fit

- Selected objects:
  - 4 untagged jets
  - 2 b-tagged jets



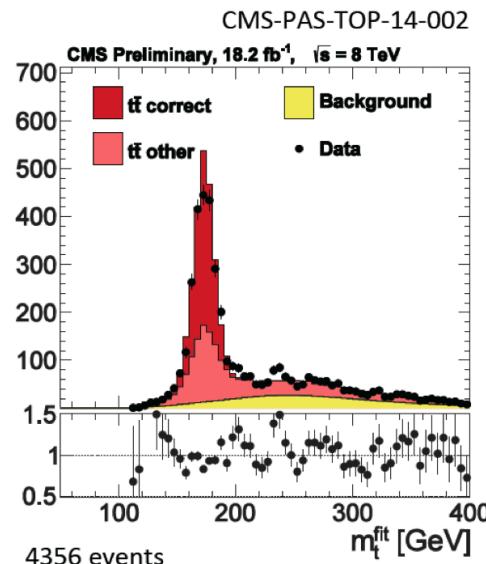
- Constraints:
  - $2x m_{jj} = m_W$
  - $m_{top} = m_{jjb,1} = m_{jjb,2} = m_{antitop}$



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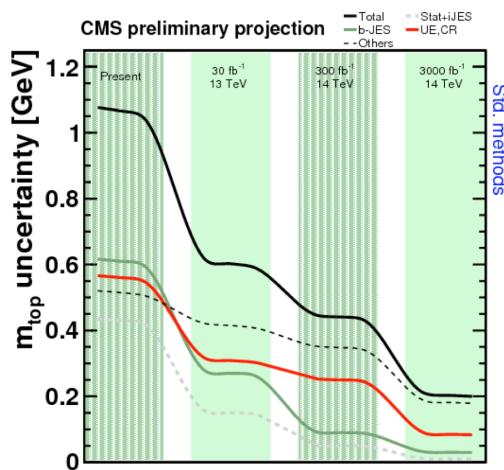
Eike Schlieckau - Universität Hamburg

September 30th 2014



Determination of the best-fit value of the Monte-Carlo top quark mass parameter

Kinematic mass determination



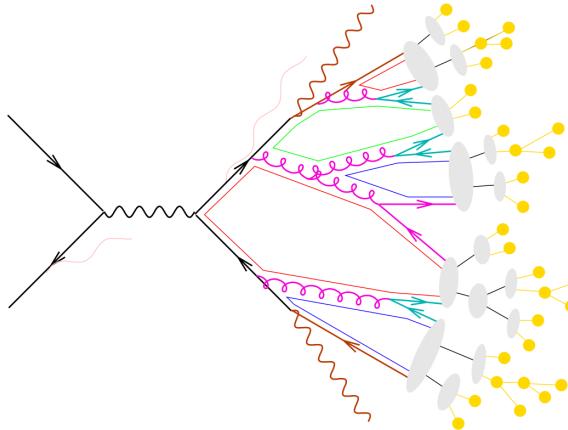
⊕ High top mass sensitivity

⊖ Precision of MC ?

⊖ Meaning of  $m_t^{\text{MC}}$  ?

$\Delta m_t \sim 0.5 \text{ GeV}$

# Monte-Carlo Event Generators



- hard scattering
- (QED) initial/final state radiation
- partonic decays, e.g.  $t \rightarrow bW$
- parton shower evolution
- nonperturbative gluon splitting
- colour singlets
- colourless clusters
- cluster fission
- cluster  $\rightarrow$  hadrons
- hadronic decays

- Full simulation of all processes (all experimental aspects accessible)
- QCD-inspired: partly first principles QCD  $\Leftrightarrow$  partly model (observable-dependent)
- Description power of data better than intrinsic theory accuracy.
- Top quark: treated like a real particle ( $m_t^{\text{MC}} \approx m_t^{\text{pole}} + ?$ ).

But pole mass ambiguous by  $O(1 \text{ GeV})$  due to confinement.

Better mass definition needed.

Uncertainty (a): But how precise is modelling?  $\rightarrow$  Part of exp. Analyses

Unvertainty (b): What is the meaning of MC QCD parameters?  $\rightarrow$

Depends strictly speaking on the observable, because of model character of MCs !

Must be adressed for each type of observable (until we have better MCs).

# MC Top Quark Mass

$$m_t^{\text{MC}} = m_t^{\text{MSR}}(R = 1 \text{ GeV}) + \Delta_{t,\text{MC}}(R = 1 \text{ GeV})$$

$$\Delta_{t,\text{MC}}(1 \text{ GeV}) \sim \mathcal{O}(1 \text{ GeV})$$

AHH, Stewart 2008

AHH, 2014

- small size of  $\Delta_{t,\text{MC}}$
- Renormalon-free
- little parametric dependence on other parameters

## MSR Mass Definition

MS Scheme:  $(\mu > \overline{m}(\overline{m}))$

$$\overline{m}(\overline{m}) - m^{\text{pole}} = -\overline{m}(\overline{m}) [0.42441 \alpha_s(\overline{m}) + 0.8345 \alpha_s^2(\overline{m}) + 2.368 \alpha_s^3(\overline{m}) + \dots]$$

MSR Scheme:  $(R < \overline{m}(\overline{m}))$



$$m_{\text{MSR}}(R) - m^{\text{pole}} = -R [0.42441 \alpha_s(R) + 0.8345 \alpha_s^2(R) + 2.368 \alpha_s^3(R) + \dots]$$

$$m_{\text{MSR}}(m_{\text{MSR}}) = \overline{m}(\overline{m})$$

→  $m_{\text{MSR}}(R)$  Short-distance mass that smoothly interpolates all R scales

# Calibration of the MC Top Mass

## Method:

- ✓ 1) Strongly mass-sensitive hadron level observable (as closely as possible related to reconstructed invariant mass distribution !)
- ✓ 2) Accurate analytic hadron level QCD predictions at  $\geq$  NLL/NLO with full control over the quark mass scheme dependence.
- ✓ 3) QCD masses as function of  $m_t^{\text{MC}}$  from fits of observable.
- 4) Cross check observable independence

$$m_t^{\text{MC}} = m_t^{\text{MSR}}(R = 1 \text{ GeV}) + \Delta_{t,\text{MC}}(R = 1 \text{ GeV})$$

$$\Delta_{t,\text{MC}}(1 \text{ GeV}) = \bar{\Delta} + \delta\Delta_{\text{MC}} + \delta\Delta_{\text{pQCD}} + \delta\Delta_{\text{param}}$$

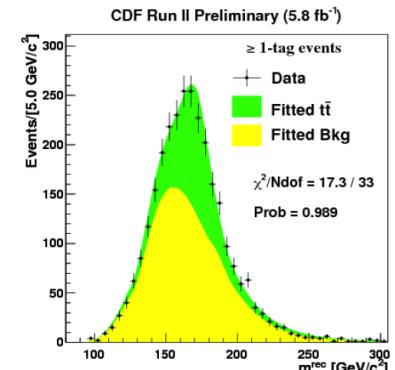
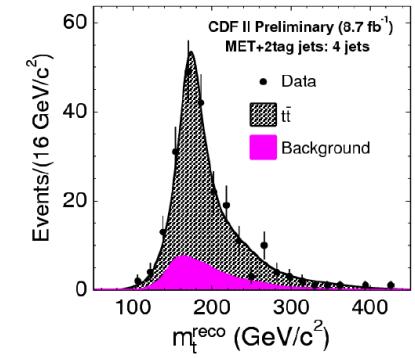
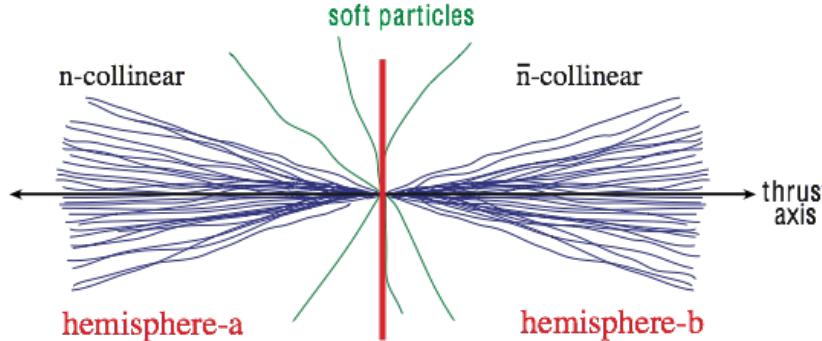
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- The diagram illustrates the decomposition of the total uncertainty  $\Delta_{t,\text{MC}}(1 \text{ GeV})$  into three components: Monte Carlo errors, QCD errors, and Parametric errors. The total uncertainty is represented by a horizontal line segment. Three red arrows point to this segment from below, each pointing to a list of errors:
- Monte Carlo errors:**
    - different tunings
    - parton showers
    - color reconnection
    - Intrinsic error, ...
  - QCD errors:**
    - perturbative error
    - scale uncertainties
    - electroweak effects
  - Parametric errors:**
    - strong coupling  $\alpha_s$
    - Non-perturbative parameters

# Thrust Distribution

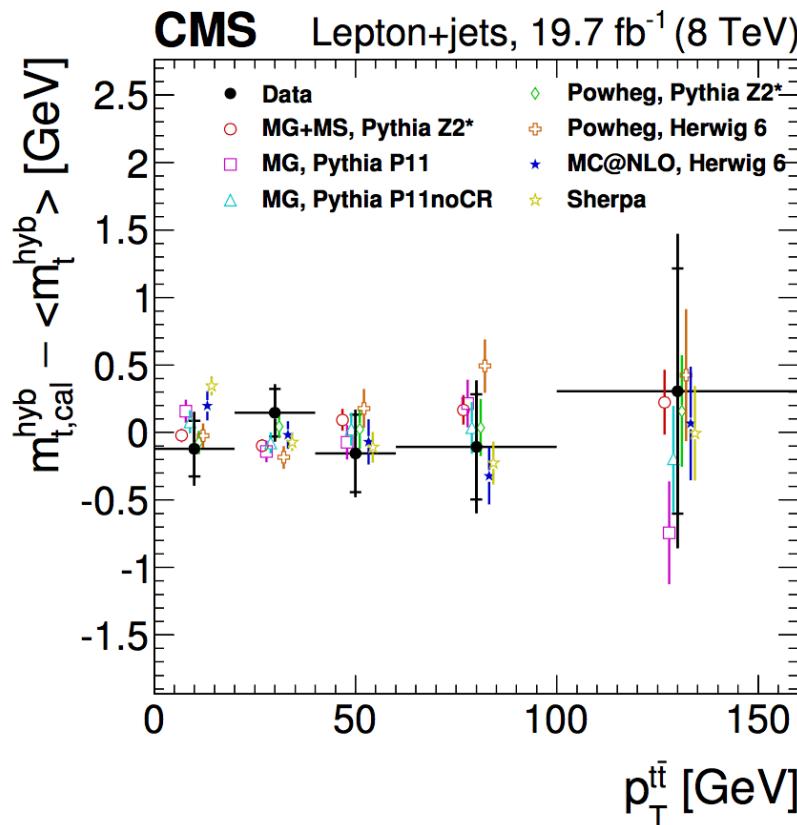
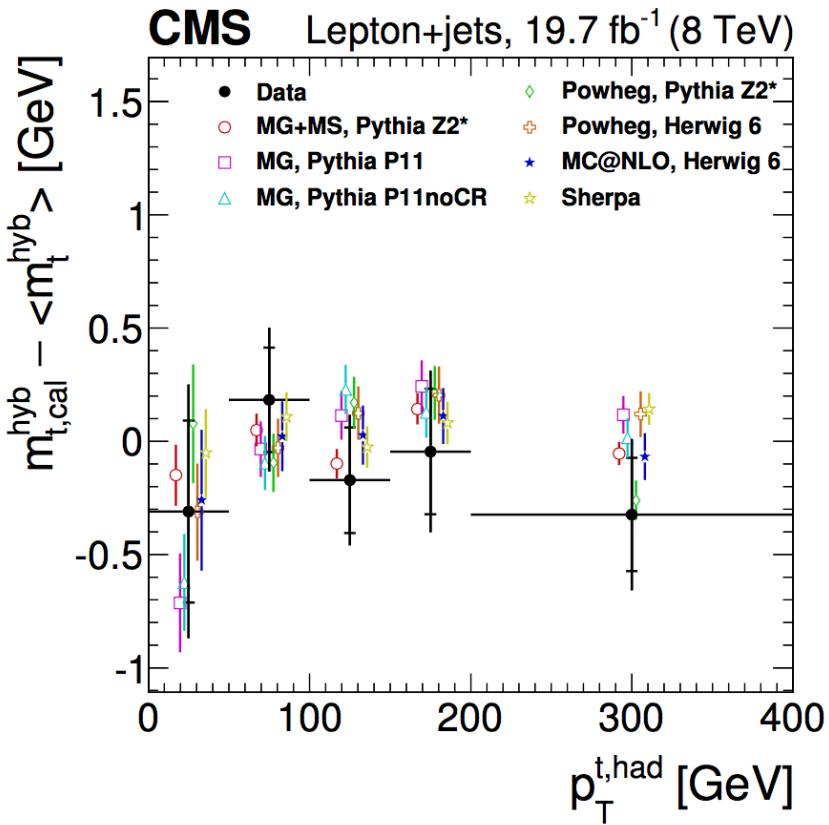
Observable: 2-jettiness in  $e^+e^-$  for  $Q \sim p_T \gg m_t$  (boosted tops)

$$\begin{aligned}\tau &= 1 - \max_{\vec{n}} \frac{\sum_i |\vec{n} \cdot \vec{p}_i|}{Q} \\ \tau \rightarrow 0 &\approx \frac{M_1^2 + M_2^2}{Q^2}\end{aligned}$$

Invariant mass distribution in the resonance region  
of wide hemisphere jets !

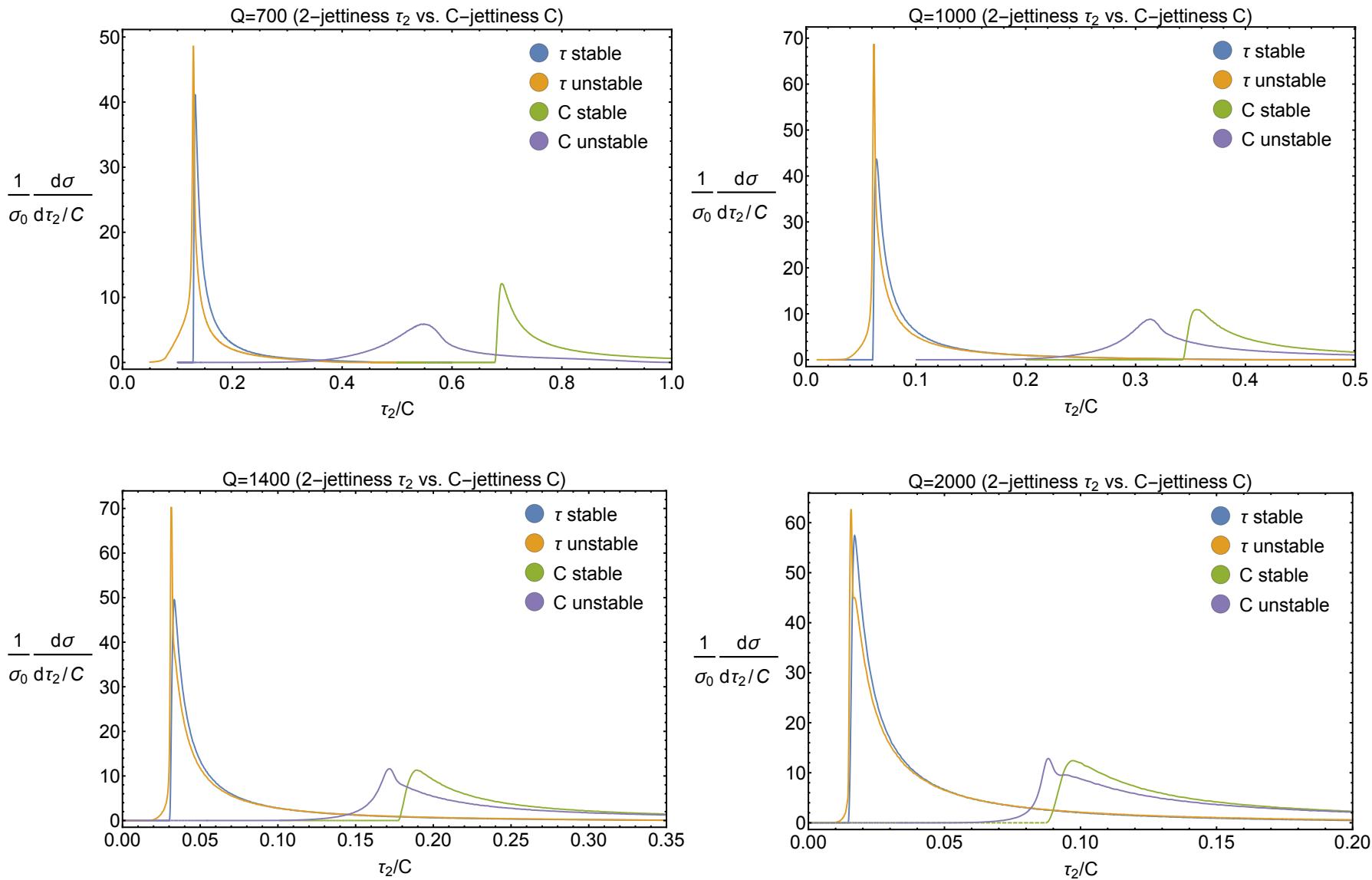


# Boosted Top Mass Measurements at CMS



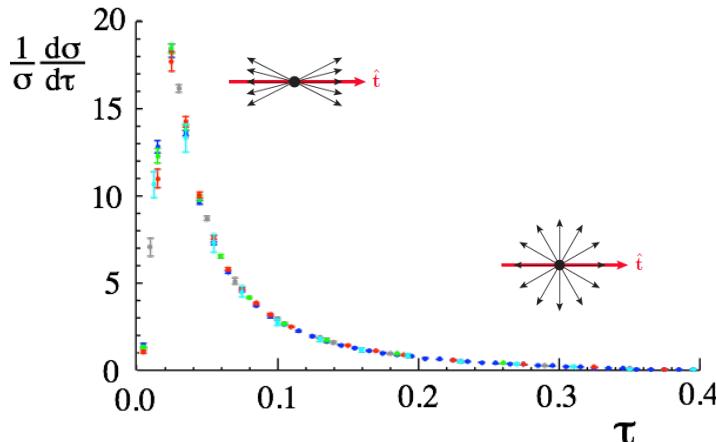
- Top mass from reconstruction of boosted tops consistent with low  $p_T$  results.
- More precise studies possible with more statistics from Run2.

# Event Shape Distributions (Pythia 8.2)



# Factorization for Event Shapes

$$\frac{d\sigma}{d\tau} = Q^2 \sigma_0 H_0(Q, \mu) \int d\ell J_0(Q\ell, \mu) S_0(Q\tau - \ell, \mu)$$



Extension to massive quarks:

- VFNS for final state jets (with massive quarks): log summation incl. mass
- Boosted fat top jets

Fleming, AHH, Mantry, Stewart 2007

Gritschacher, AHH, Jemos, Mateu Pietrulewicz 2013-2014

Butenschön, Dehnadi, AHH, Mateu 2016 (to appear soon)

→ NNLL + NLO + non-singular + hadronization + renormalon-subtraction

Massless quarks:

Korshenski, Sterman 1995-2000

Bauer, Fleming, Lee, Sterman  
(2008)

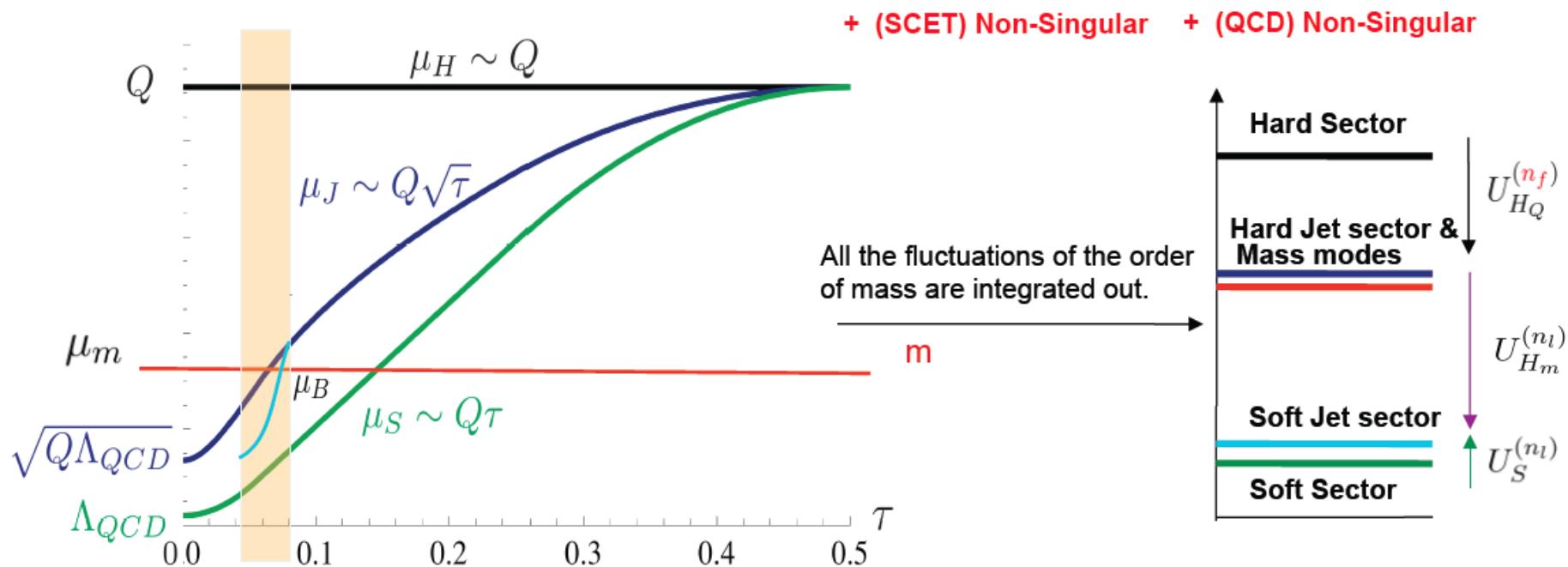
Becher, Schwartz (2008)

Abbate, AHH, Fickinger, Mateu,  
Stewart 2010

# b(oothed)HQET Factorization

$$\left| \frac{1}{\sigma_0} \frac{d\hat{\sigma}(\tau)}{d\tau} \right|^{\text{bHQET}} = Q H_Q^{(n_f)}(Q, \mu_Q) U_{H_Q}^{(n_f)}(Q, \mu_Q, \mu_m) H_m^{(n_f)}(\bar{m}^{(n_f)}, \mu_m) U_{H_m}^{(n_l)}\left(\frac{Q}{\bar{m}^{(n_l)}}, \mu_m, \mu_B\right)$$

$$+ \int ds \int dk B^{(n_l)}\left(\frac{s}{m_J^{(n_l)}}, \mu_B, m_J^{(n_l)}\right) U_S^{(n_l)}(k, \mu_B, \mu_S) S_{\text{part}}^{(n_l)}(Q\tau - Q\tau_{\text{MIN}} - \frac{s}{Q} - k, \mu_S)$$



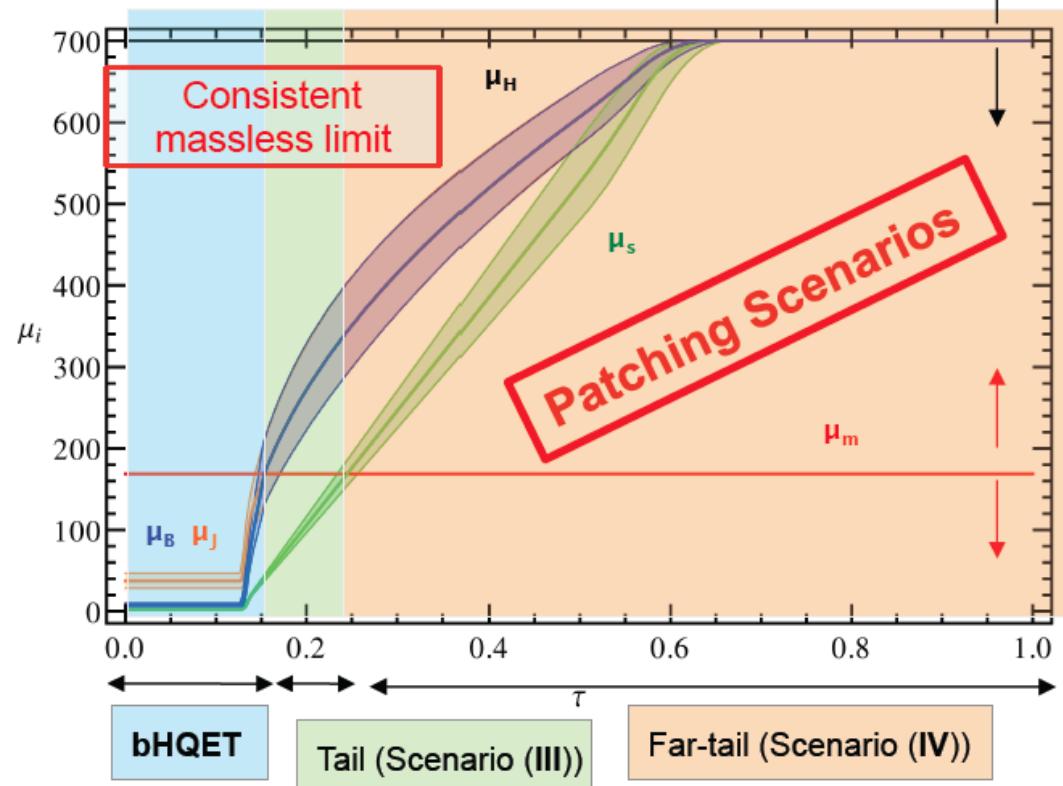
- Matching coefficient of SCET and bHQET have a large log from secondary corrections.

# Profile Functions

Profile functions should sum up large logarithms and achieve smooth transition between the peak, tail and far-tail.

$$\log\left(\frac{Q}{\mu_H}\right) \quad \log\left(\frac{m_J}{\mu_m}\right) \quad \log\left(\frac{\mu_J^2}{Q\mu_s}\right) \quad \log\left(\frac{m_J\mu_B}{Q\mu_s}\right) \quad \log\left(\frac{Q(\tau - \tau_{\min}) + 2\Lambda_{\text{QCD}}}{\mu_s}\right)$$

$Q = 700 \text{ GeV}$

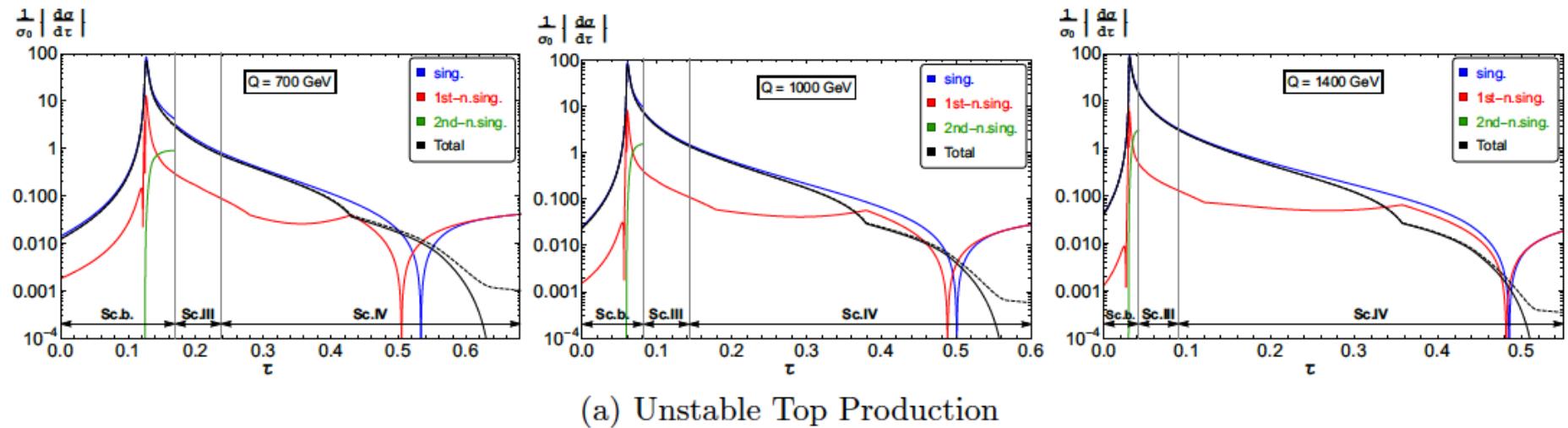


## Scales Variation

- ✓ Generalized to arbitrary mass values
- ✓ Compatible with massless profiles

Proper scale variations are essential in reliable estimation of missing higher order terms.

# Large Log Resummation

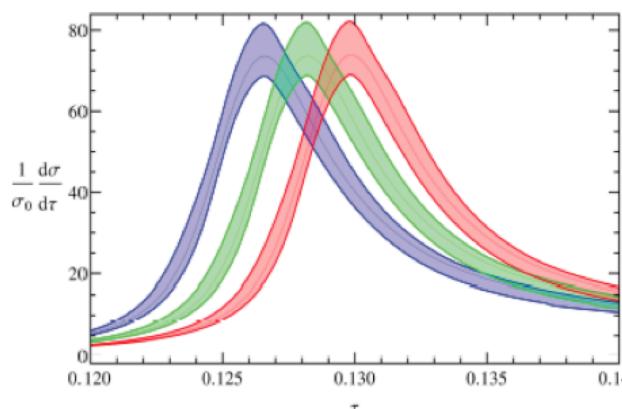


- Contributions contained in singular factorized cross section more than 1 order of magnitude larger than the non-singular contributions
- Confirmation that for  $Q=700 \text{ GeV}$  the top quark are already boosted that that the correct treatment of large collinear logs is important.

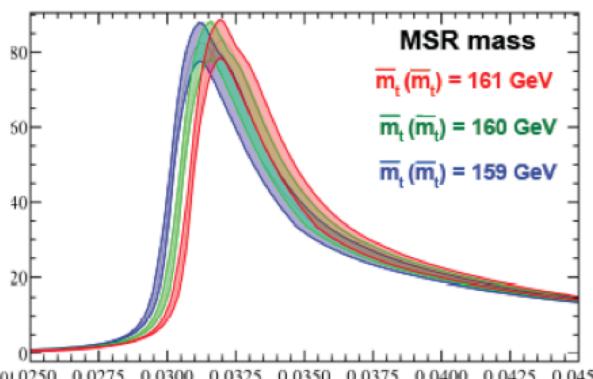
# 2-Jettiness for Top Production (QCD)

$$\frac{d\sigma}{d\tau_2} = f(\underbrace{m_t^{\text{MSR}}(R), \alpha_s(M_Z), \Omega_1, \Omega_2, \dots}_{\text{any scheme possible}}, \underbrace{\mu_h, \mu_j, \mu_s, \mu_m}_{\text{Non-perturbative}}, \underbrace{R, \Gamma_t}_{\text{renorm. scales finite lifetime}})$$

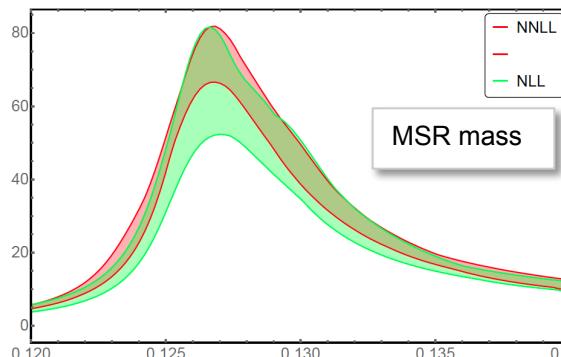
$Q=700 \text{ GeV}$



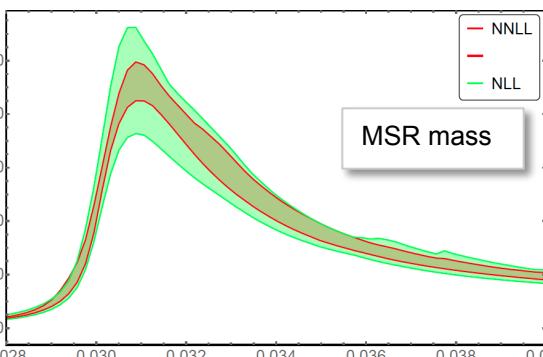
$Q=1400 \text{ GeV}$



$Q=700 \text{ GeV}$



$Q=1400 \text{ GeV}$



- Higher mass sensitivity for lower  $Q$  ( $p_T$ )
- Finite lifetime effects included
- Dependence on non-perturbative parameters
- Convergence:  $\Omega_{1,2,\dots}$
- Good convergence
- Reduction of scale uncertainty (NLL to NNLL)
- Control over whole distribution

# Fit Procedure Details

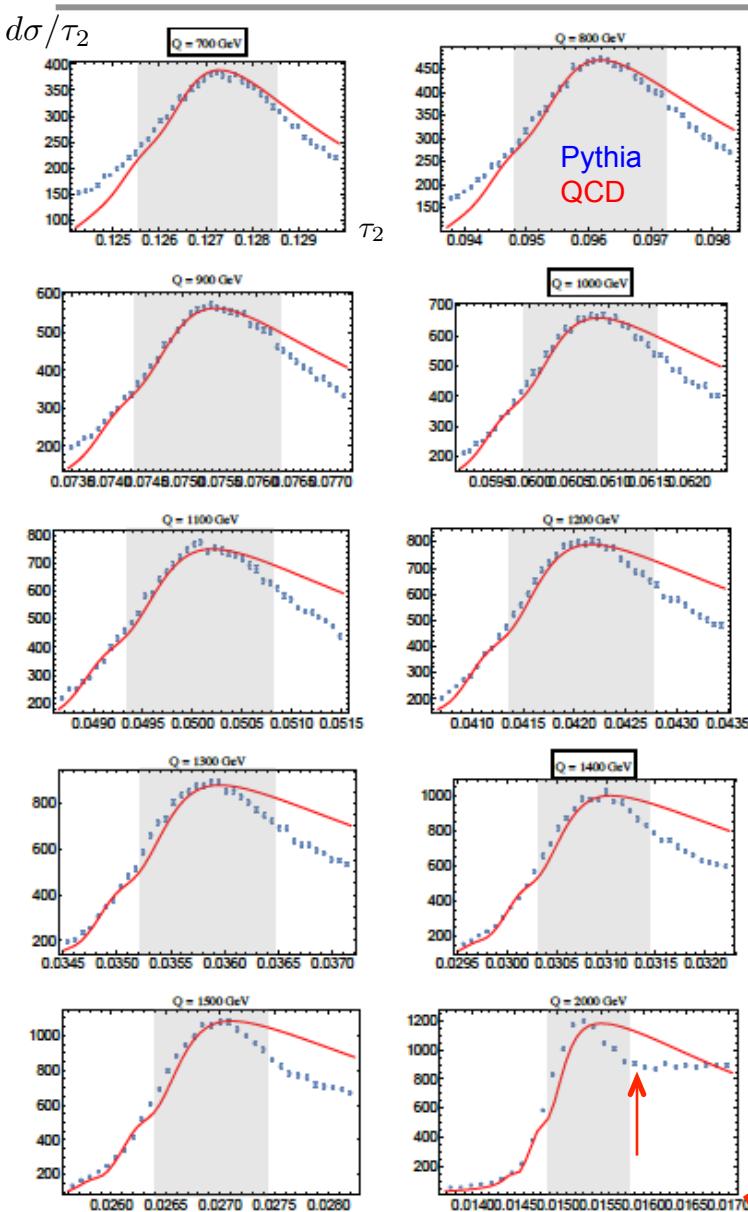
$$\frac{d\sigma}{d\tau_2} = f(\underbrace{m_t^{\text{MSR}}(R), \alpha_s(M_Z), \Omega_1, \Omega_2, \dots}_{\text{any scheme possible}}, \underbrace{\mu_h, \mu_j, \mu_s, \mu_m}_{\text{Non-perturbative}}, \underbrace{R, \Gamma_t}_{\text{renorm. scales finite lifetime}})$$

QCD parameters measured from Pythia

- Fit parameters:  $m_t^{\text{MSR}}(R), \alpha_s(M_Z), \Omega_1, \Omega_2, \dots$ ,
- Perturbative error: fits for 500 randomly picked sets of renor. scales
- Tunings: 1 (“very old”), 3 (“LEP”), 7 (“Monash”)
- Top quark width:  $\Gamma_t$  = dynamical (default), 0.7, 1.4, 2.0 GeV
- External smearing (Detector effects):  $\Omega_{1,\text{smear}}$  = 0, 0.5, …, 3.0, 3.5, GeV (just for cross checks)
- Pythia masses:  $m_t^{\text{Pythia}} = 170, \dots, 175$  GeV
- Strong coupling:  $\alpha_s(M_Z) = 0.114, 0.116, 0.118, 0.120, 0.122$
- Fit possible for any order / mass scheme (so far NLL+NNLL / MSR)

Number of fits entering the first analysis:  $2.8 \cdot 10^6$

# Peak Fits



Default renormalization scales;  $\Gamma_t = 1.4$  GeV, tune 3,  $\Omega_{1,\text{smear}} = 0$  GeV,  $m_t^{\text{Pythia}} = 170$  GeV,  $Q = \{700, 1000, 1400\}$  GeV, peak fit (60/80)%, normalized to fit range

- Good agreement of Pythia 8.2 with NNLL+NLO QCD description
- Pythia statistics:  $10^6$  events
- Discrepancies in distribution tail and for higher energies (Pythia is less reliable where fixed-order results valid, well reliable in soft-collinear limit)
- Pythia kink issue ?
- Excellent sensitivity to the top quark mass.
- Tree-Level:

$$\tau_2^{\text{peak}} = 1 - \sqrt{1 - \frac{4m_t^2}{Q^2}}$$

tune = 3

$m_{\text{MC}} = 170$ .

$\Gamma_t = -1$ . GeV

$\alpha = 0.118$

$m^{\text{SR}}(5 \text{ GeV}) = 169.138 \pm 0.099$

$$\frac{\chi^2}{\text{dof}} = 35.36$$

$\Omega_1 = 0.434 \pm 0.060$  GeV

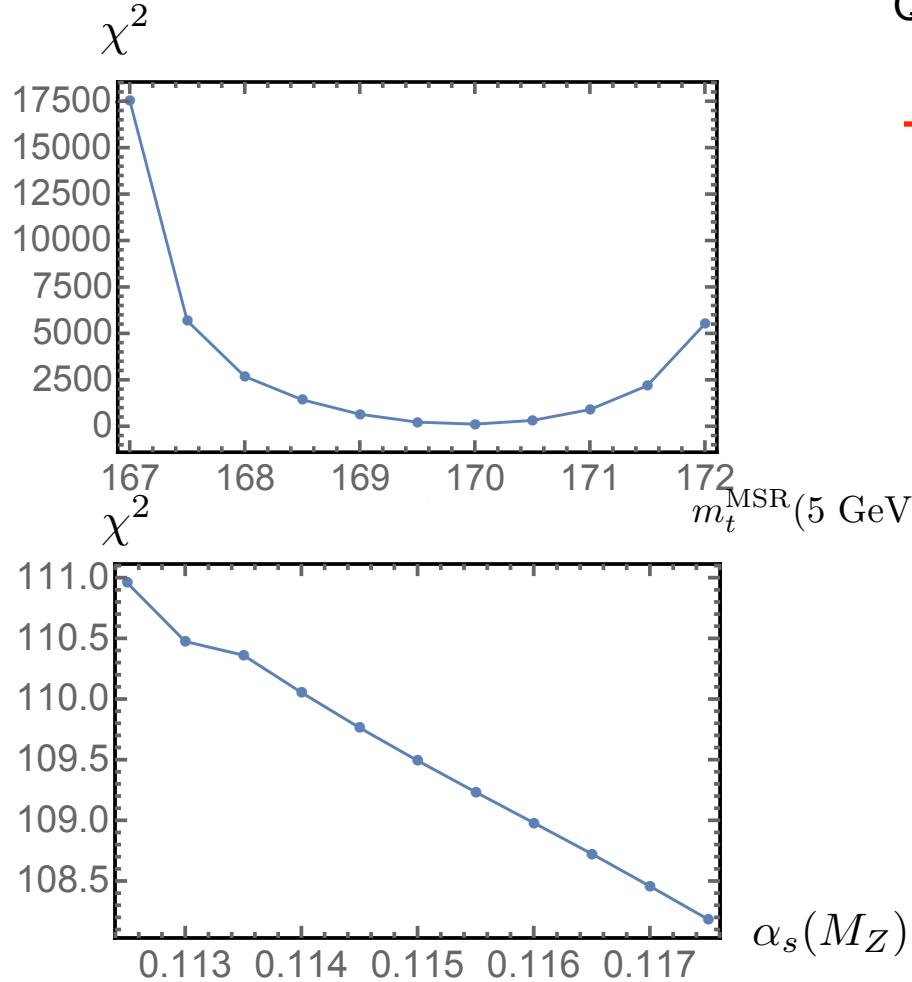
$\Omega_2 = 0.473 \pm 0.060$  GeV

$\Omega_3 = -0.158 \pm 0.300$  GeV

$\Omega_4 = -2.226 \pm 1.000$  GeV

Preliminary

# Peak Fits



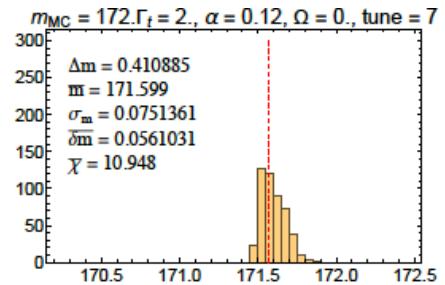
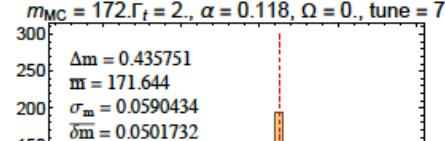
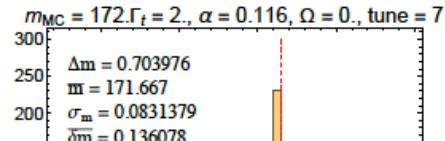
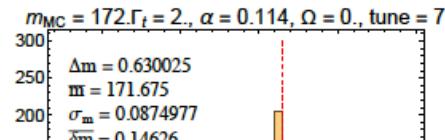
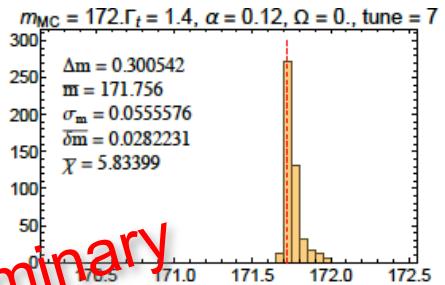
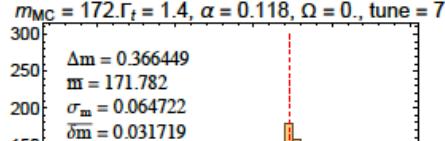
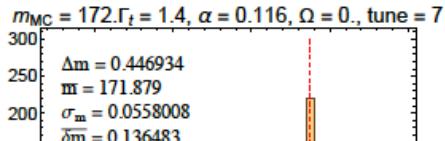
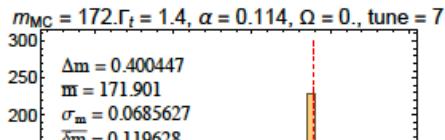
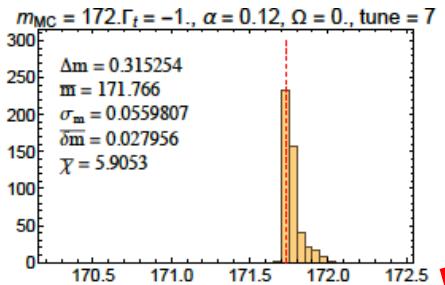
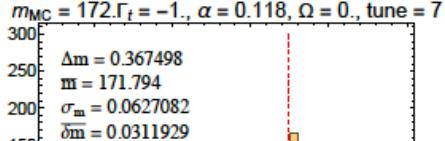
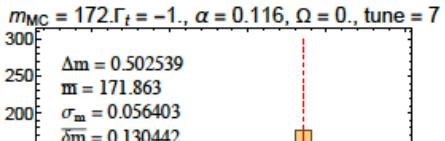
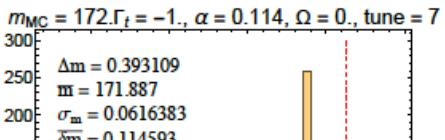
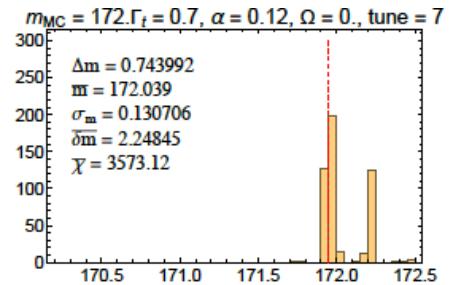
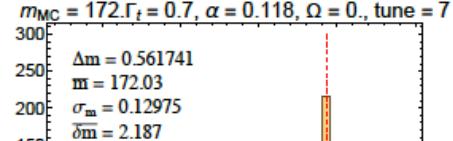
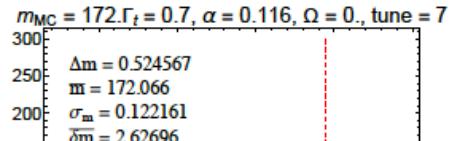
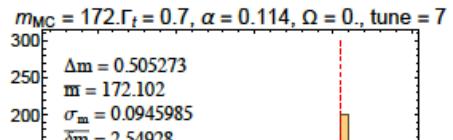
Default renormalization scales;  $\Gamma_t=1.4 \text{ GeV}$ , tune 7,  $\Omega_{1,\text{smear}}=2.5 \text{ GeV}$ ,  $m_t^{\text{Pythia}}=171 \text{ GeV}$ ,  $Q=\{700, 1000, 1400\} \text{ GeV}$ , peak fit (60/80)%

→  $\chi^2_{\min} \sim O(100)$

- Very strong sensitivity to  $m_t$
- Low sensitivity to strong coupling
- Take strong coupling as input
- $\chi^2_{\min}$  and  $\delta m_t^{\text{stat}}$  do not have any physical meaning
- We use rescaled  $\chi^2/\text{dof}$  (PDG prescription) to define “**intrinsic MC compatibility uncertainty**”

Preliminary

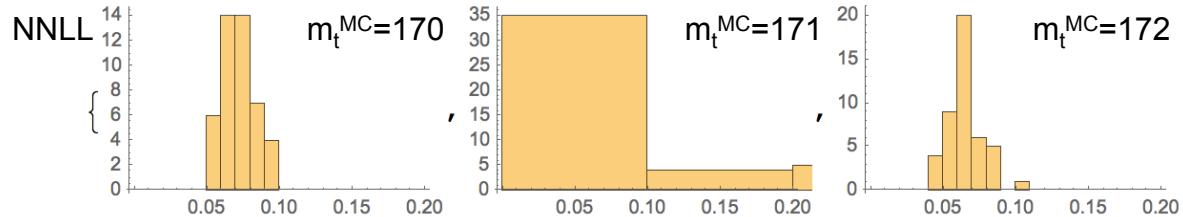
# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$



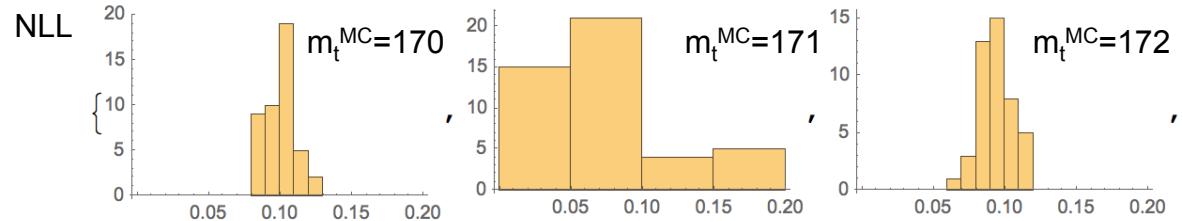
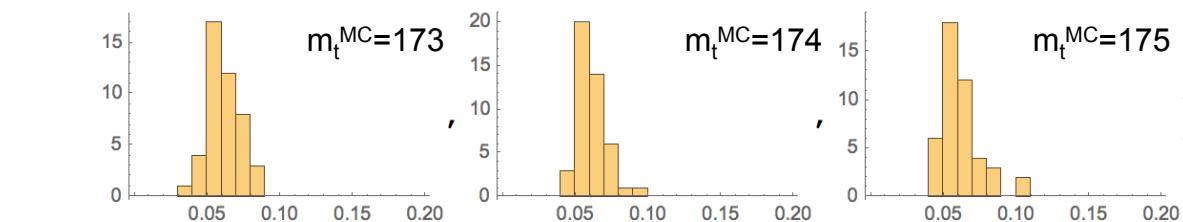
Preliminary

# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$

Distribution of standard deviations: each from scan over 500 profile functions



- Measure for scale error
- NNLL: 60-75 MeV
- NLL: 90-110 MeV
- Probably to be multiplied by factor 2 for scale uncertainty



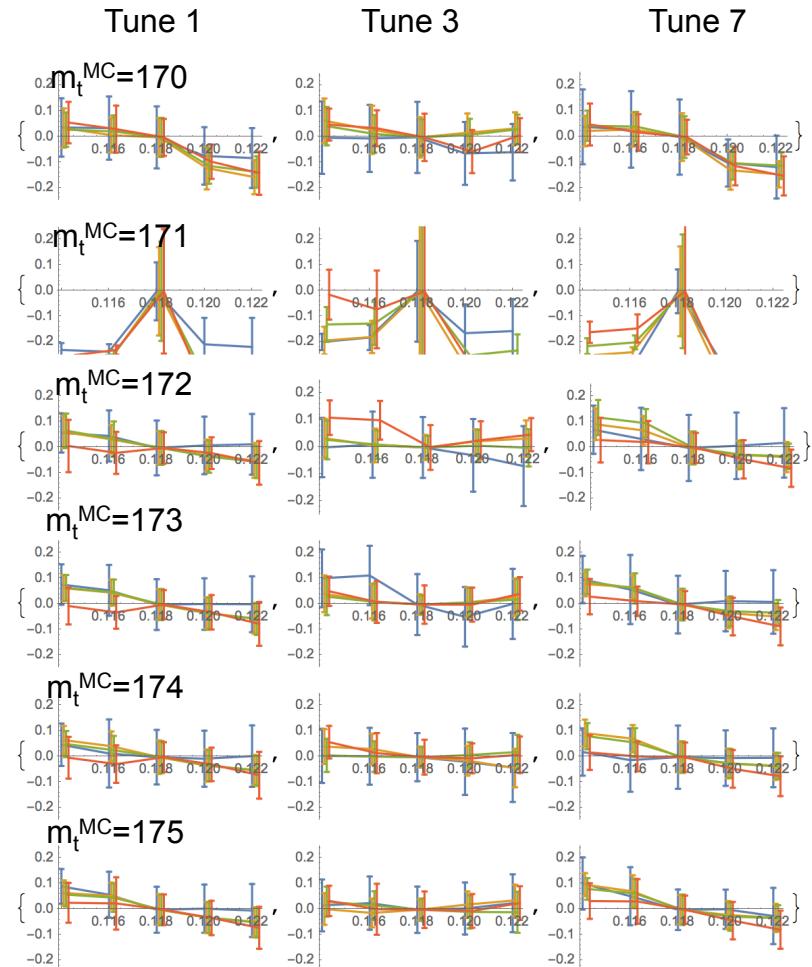
Preliminary

# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$

Parametric dependence on strong coupling

$$m_t^{\text{MSR}}[\alpha_s(M_Z)] - m_t^{\text{MSR}}[0.118]$$

- Small sensitivity of  $m_t^{\text{MSR}}(1\text{GeV})$  on  $\alpha_s(M_Z)$ . [~50 MeV error] ✓



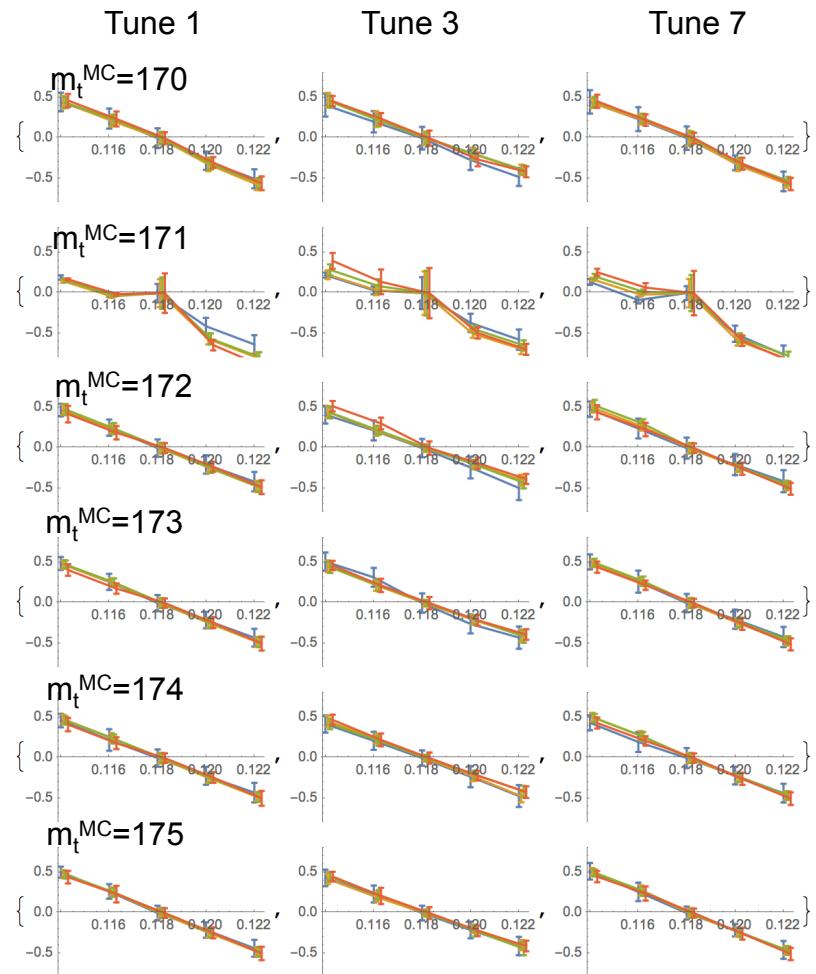
Preliminary

# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$

Parametric dependence on strong coupling

$$m_t(m_t)[\alpha_s(M_Z)] - m_t(m_t)[0.118]$$

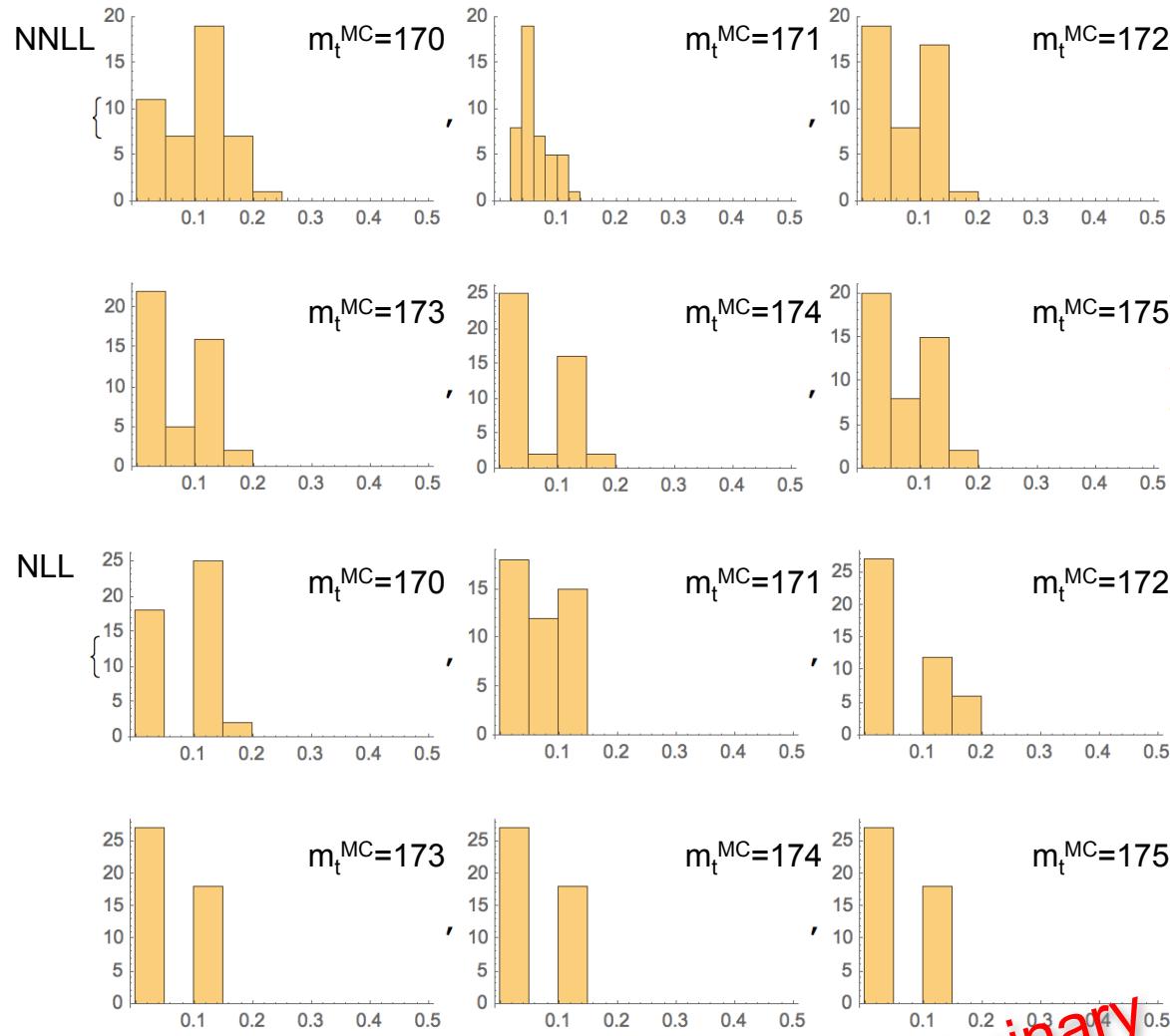
- Large sensitivity of MSbar mass on  $\alpha_s(M_Z)$ . [not an error, but calculated from MSR mass] ✓
- The MC top mass IS FAR AWAY from the Msbar mass.



Preliminary

# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$

## Intrinsic MC Compatibility Error (distribution of mean values)



- Measure for intrinsic MC uncertainty
- NNLL: ~150 MeV
- NLL: ~ 150 MeV
- Probably never before accounted in reconstruction analyses
- **Measure for ultimate precision (MC dependent !)**

Preliminary

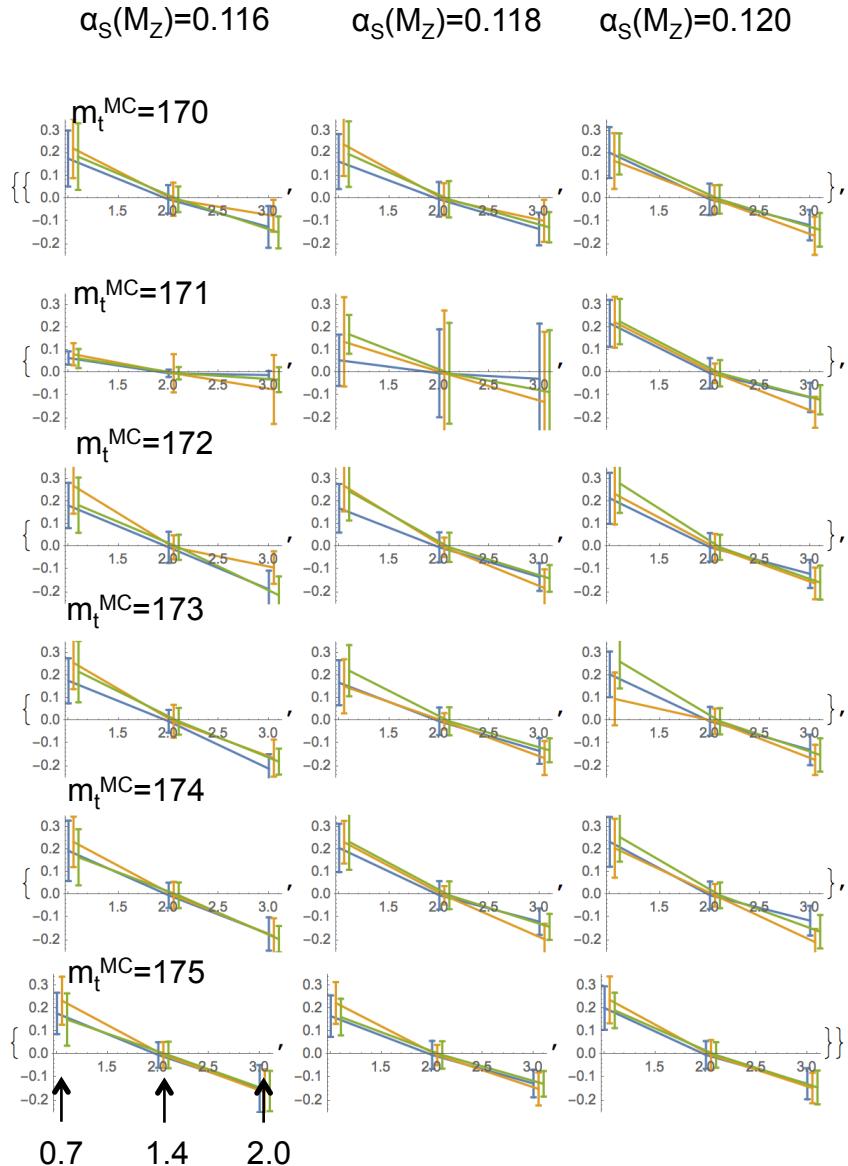
# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$

Top width dependence

$$m_t^{\text{MSR}}[\Gamma_t] - m_t^{\text{MSR}}[\Gamma_t=1.4]$$

- Clear sensitivity to top width value.
- Can be interpreted as observable dependence or MC modelling dependence.
- Should be accounted for as an additional uncertainty [ $\sim 150 \text{ MeV}$ ].

Preliminary



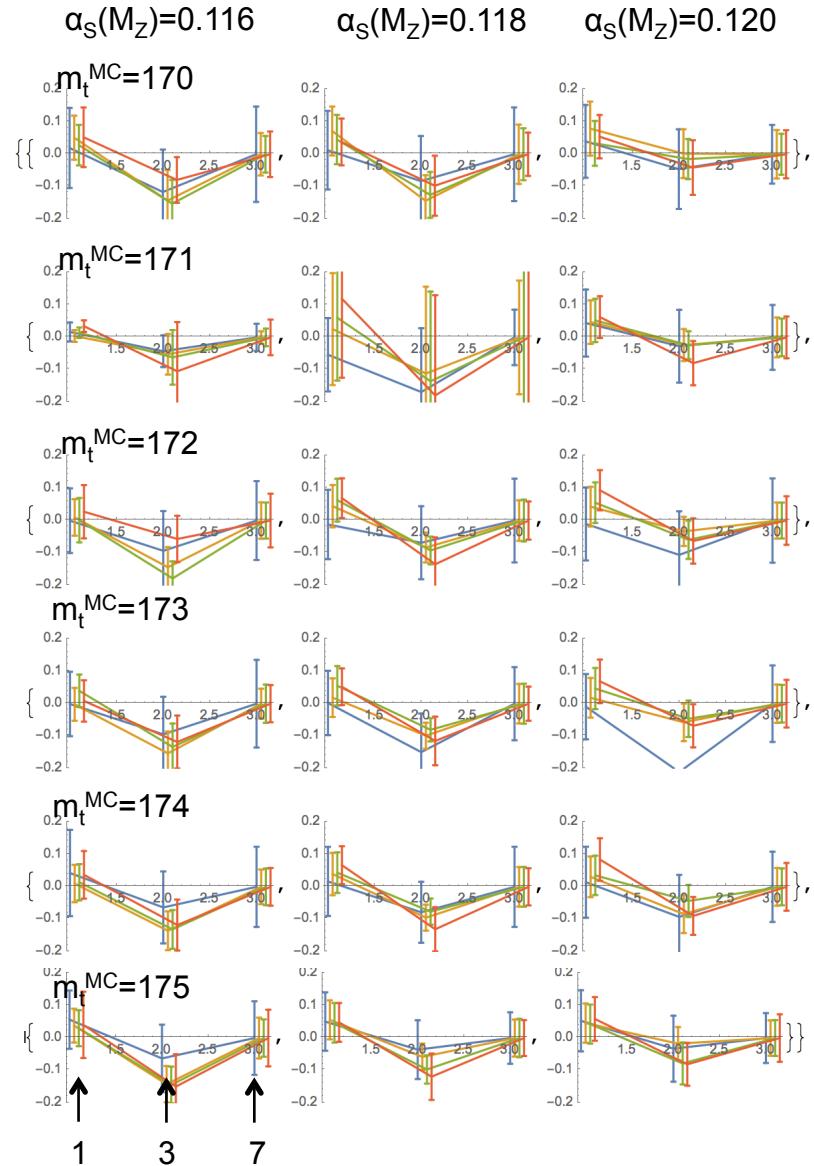
# Peak Fits: $m_t^{\text{MSR}}(1 \text{ GeV})$

## Tune dependence

$$m_t^{\text{MSR}}[\text{tune}] - m_t^{\text{MSR}}[\text{tune 7}]$$

- Clear sensitivity to tune.
- MC top mass is tune-dependent !
- Tune-dependence is not an error !
- Opposite dependence should be visible in MC top mass determinations from experimental data.  
(highly nontrivial validation)

Preliminary



# Summary

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- First serious precise MC top quark mass calibration based on  $e^+e^-$  2-jettiness (large  $p_T$ ): preliminary results.
- NNLL+NLO QCD calculations based on an extension of the SCET approach concerning massive quark effects (all large logs incl.  $\ln(m)$ 's summed systematically).
- The Monte Carlo top mass calibration in terms of  $m_t^{\text{MSR}}(1\text{GeV})$ :
  - Scale dependence (NNLL): ~ 150 MeV
  - $\alpha_s$  dependence ( $\delta\alpha_s = 0.002$ ): ~ 50 MeV
  - Intrinsic MC error: ~ 150 MeV
  - Observable dependence: ~ 150 MeV
- MC top mass is tune-dependent and MC dependent !  
*Using MC top mass calibration might eliminate these error sources from the experimental analyses.*  
*Confirmation of the dependence predicted by calibration provides highly non-trivial cross check concerning the universality of the calibration.*

Preliminary !!!

# Outlook & Plans

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- Full verified error analysis @ NNLL+NLO on the way
  - Different sets of  $Q(p_T)$  values
  - Different fit ranges
  - Bug fixes
- Calibration Package for public use
  - Calibration  $m_t^{\text{MC}} \rightarrow m_t^{\text{MSR}}(1\text{GeV})$
  - Code  $m_t^{\text{MSR}}(1\text{GeV}) \rightarrow$  any other scheme
- Heavy jet mass, C-parameter (NNLL), pp-2-jettiness analysis (NLL) w.i.p.
- NNNLL+NNLO (2-jettiness for  $e^+e^-$ ) w.i.p
- Mass (+ Yukawa coupling) conversions w. QCD + electroweak

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# Backup Slides

# Pole Mass from MSR Mass

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$$\begin{aligned}\alpha_s(M_Z) &= 0.118 \\ n_f &= 5\end{aligned}$$

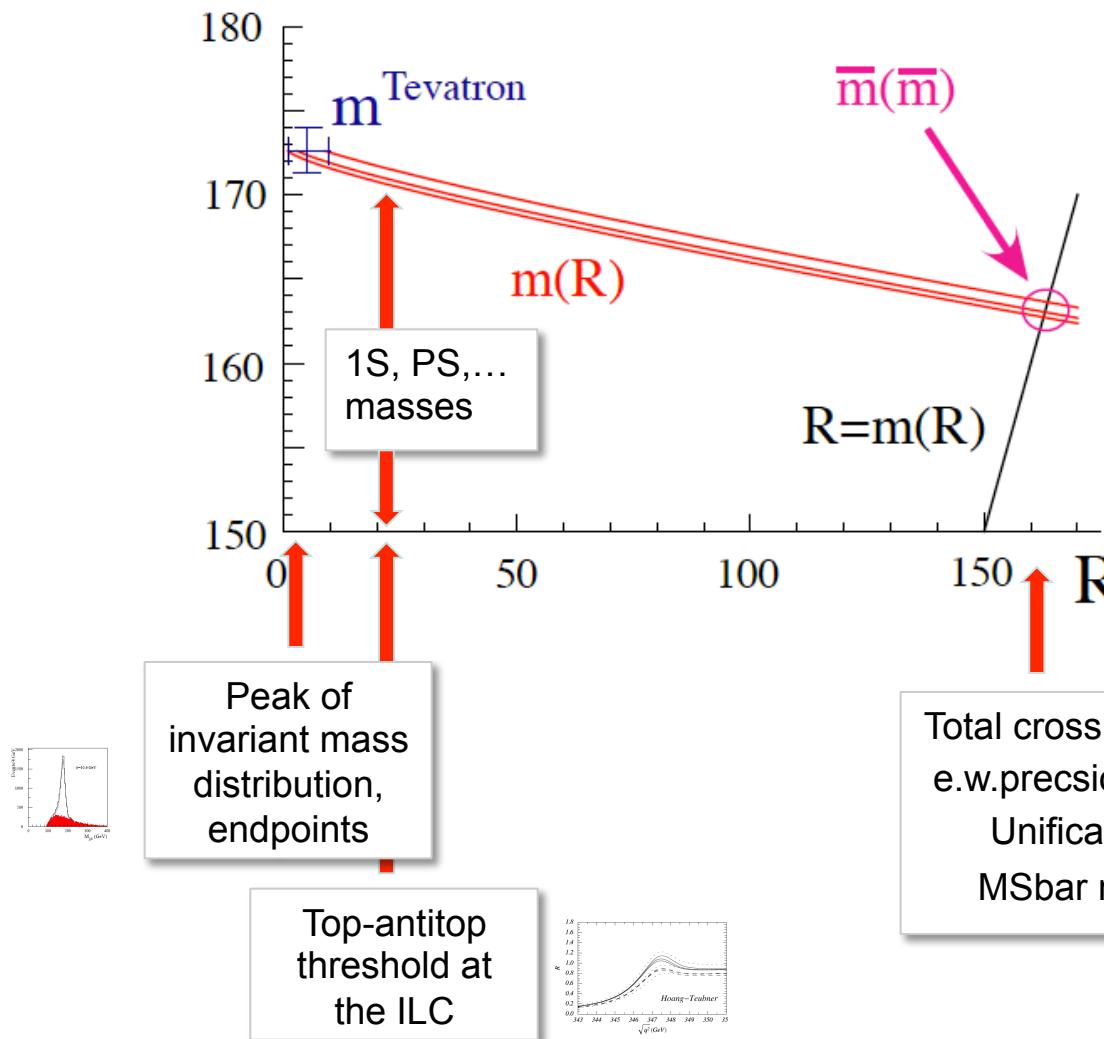
$$\begin{aligned}m_t^{\text{pole}} - m_t^{\text{MSR}}(1 \text{ GeV}) &= 0.173 + 0.138 + 0.159 + 0.23 \text{ GeV} \xleftarrow{\text{calculated}} \\ &\quad + 0.53 + 1.43 + 4.54 + 16.6 \text{ GeV} \xleftarrow{\text{extrapolated}} \\ &\quad + 68.6 + 317.7 + 1629 + 9158 \text{ GeV}\end{aligned}$$

- Size of terms consistent with scale error estimate of calibration.
- No stable determination of pole mass.

# MSR Mass Definition

$$m_t^{\text{MC}} = m_t^{\text{MSR}}(3^{+6}_{-2} \text{ GeV}) = m_t^{\text{MSR}}(3 \text{ GeV})^{+0.6}_{-0.3}$$

AH, Stewart: arXive:0808.0222



Good choice for R:

Of order of the typical scale  
of the observable used to  
measure the top mass.

# Masses Loop-Theorists Like to use

## Total cross section (LHC/Tev):

$$m_t^{\text{MSR}}(R = m_t) = \overline{m}_t(\overline{m}_t)$$

- more inclusive
- sensitive to top production mechanism (pdf, hard scale)
- indirect top mass sensitivity
- large scale radiative corrections

$$M_t = M_t^{(O)} + M_t(0)\alpha_s + \dots$$

## Threshold cross section (ILC):

$$m_t^{\text{MSR}}(R \sim 20 \text{ GeV}), \ m_t^{\text{1S}}, \ m_t^{\text{PS}}(R)$$

$$M_t = M_t^{(O)} + \langle p_{\text{Bohr}} \rangle \alpha_s + \dots$$

$$\langle p_{\text{Bohr}} \rangle = 20 \text{ GeV}$$

## Inv. mass reconstruction (ILC/LHC):

$$m_t^{\text{MSR}}(R \sim \Gamma_t), \ m_t^{\text{jet}}(R)$$

$$M_t = M_t^{(O)} + \Gamma_t \alpha_s + \dots$$

$$\Gamma_t = 1.3 \text{ GeV}$$

- more exclusive
- sensitive to top final state interactions (low scale)
- direct top mass sensitivity
- small scale radiative corrections

