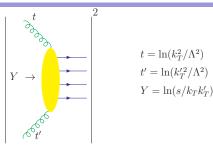
Fitting the Discrete BFKL Pomeron to Low-x HERA Data

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$$\mathcal{A}(Y,t,t') = \int_{\mathcal{C}} d\omega e^{\omega Y} f_{\omega}(t) f_{\omega}^{*}(t'), \qquad \int dt' \, \mathcal{K}(\alpha_{s},t,t') f_{\omega}(t') = \omega \, f_{\omega}(t)$$

For fixed coupling

$$f_{\omega}(t) \sim e^{i V_{\omega} t}$$

with ν_{ω} fixed for fixed ω .

For running coupling v_{ω} becomes t dependent, decreasing as t increases to t_c .

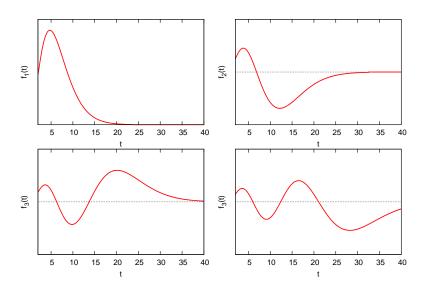
For $t > t_c$, v_{ω} is imaginary \longrightarrow evanescence $(f_{\omega}(t) \to 0 \text{ as } t \to \infty)$.

Assume that the IR (non-perturbative) properties of QCD determine the phase, η , of oscillations at some small $t=t_0$) (Lipatov 1986) This leads to a discrete set of eigenfunctions

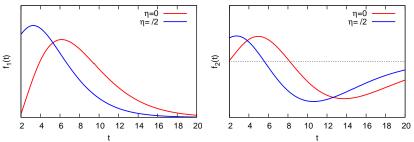
$$\mathcal{A}(Y,t,t') = \sum_{n} e^{\omega_{n} Y} f_{\omega}(t) f_{\omega}^{*}(t')$$

[Regge Poles]

First 4 Eigenfunctions



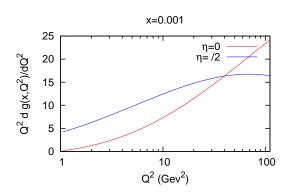
Sensitivity to Infrared phase, η



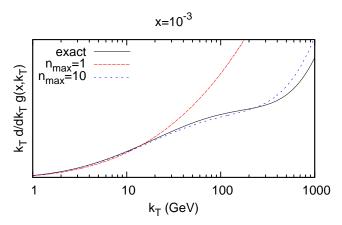
 η cannot be determined in perturbative QCD.

Phases must be treated as free parameters in a fit to data.

Infrared Phase Sensitivity of Unintegrated Gluon Density

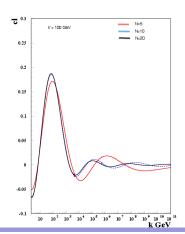


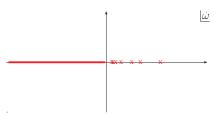
Very good convergence for the un-integrated gluon density after 10 eigenfunctions.



Completeness of Eigenfunctions

$$\sum_{n} f_{\omega_{n}}(t) f_{\omega_{n}}^{*}(t') = \delta(t - t')$$



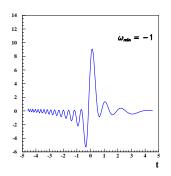


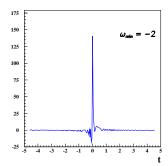
As well as a set of discrete poles for positive ω there is a cut along negative real axis.

The contribution from this cut is needed in order to reproduce the required completeness relation.

This cut contribution generates a small but non-negligible contribution to unintegrated gluon density for low-x (despite $x^{-|\omega|}$ suppression).

$$\sum_{n} f_{\omega_n}(t) f_{\omega_n}^*(t') + \int_{\omega_{min}}^0 d\omega f_{-|\omega|}(t) f_{-|\omega|}^*(t') = \delta(t - t')$$





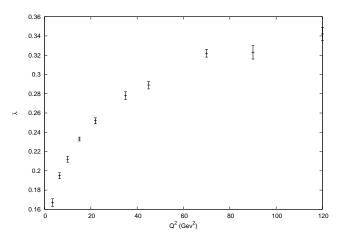
A popular parameterization of structure functions:

$$F_2(x,Q^2) = A(Q^2)x^{-\lambda(Q^2)}$$

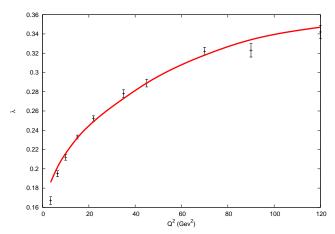
Not motivated by either a BFKL or DGLAP analysis, but nevertheless seems to work very well (Caldwell, 2015) To match the discrete BFKL pomeron we need to find a fit such that

$$A(Q^2)x^{-\lambda(Q^2)} \approx \sum_n C_n f_{\omega_n}(Q^2)x^{-\omega_n}$$

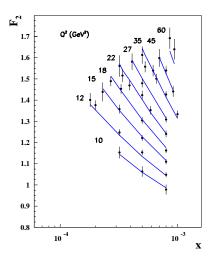
Since $f_{\omega}(t)$ decreases for sufficiently large t, we expect that for sufficiently large Q^2 , $\lambda(Q^2)$ is a decreasing function of Q^2 .



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A global fit is only possible for data with $x \le 10^{-3}$ (37 data points from HERA) Infrared phases selected using a 3-parameter ansatz. $\chi^2/\text{DOF} = 0.67$

It may only be possible to get a good fit for $x \le \sim 10^{-3}$. BFKL is an expansion in $1/\ln|x|$. Corrections to NLO BFKL expected to be $\sim 1/|\ln(x)|^2$. $\sim 2\%$ for $x = 10^{-3}$ (comparable with experimental accuracy) $\sim 5\%$ for $x = 10^{-2}$

There are very large corrections to the photon impact factor at NLO (Chirilli & Balitsky; Bartels & Chachamis) - these need some sort of collinear resummation analogous to the NLO characteristic function (Salam, 1999)

Summary

- ➤ Eigenfunctions of discrete BFKL Pomeron are very sensitive to infrared phases.
- ➤ Rapid convergence of sum of eigenfunctions to generate unintegrated gluon density. 10 eigenfunctions sufficient for good accuracy
- Continuum from cut along negative real axis in Mellin transform variable, ω, is needed for completeness - makes a small but significant contribution to unintegrated gluon density.
- ▶ For $F_2 \sim x^{-\lambda(Q^2)}$ we expect $\lambda(Q^2)$ decrease at large Q^2 reflecting dacay of leading eigenfunctions. Data so far only gives a hint of this perhaps this could be confirmed at LHeC!!
- ► Fit with $\chi^2/DOF = 0.67$ has been found for HERA DIS data with $x < \sim 10^{-3}$.