



***Photon - induced interactions  
in the run II at the LHC***

*Victor P. Goncalves*

*Theory High Energy Physics - Lund University - Sweden*

*and*

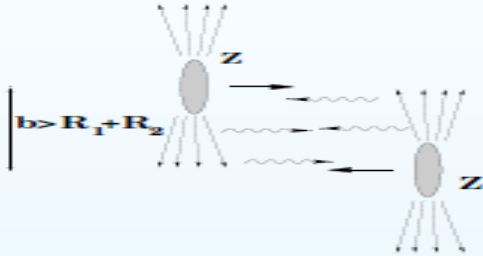
*High and Medium Energy Group - UFPel - Brazil*

*Based on arXiv:1605.05840, arXiv:1605.08186 and paper in preparation*

*In collaboration with B. Moreira, F. Navarra, M. Machado, G. S. dos Santos and D. Spiering.*

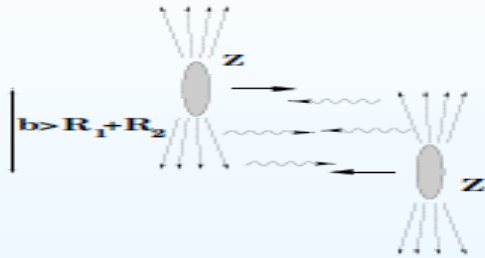
***Gyöngyös  
07 June  
2016***

# Motivation



1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

# Motivation

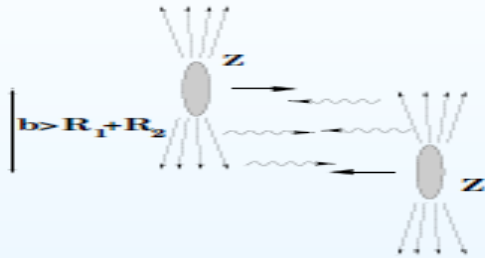


Center of mass energies

1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

LHC	$pp$	$W_{\gamma p} \lesssim 8390 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 4504 \text{ GeV}$
LHC	$pPb(Ar)$	$W_{\gamma A} \lesssim 1500 (2130) \text{ GeV}$	$W_{\gamma\gamma} \lesssim 260 (480) \text{ GeV}$
LHC	$PbPb$	$W_{\gamma A} \lesssim 950 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 160 \text{ GeV}$
HERA	$ep$	$W_{\gamma p} \lesssim 200 \text{ GeV}$	—

# Motivation



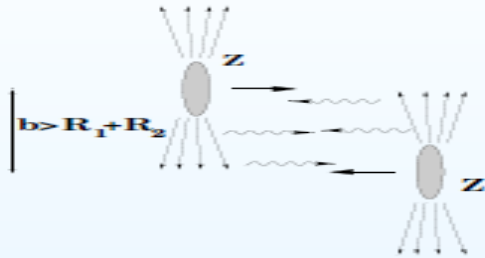
Center of mass energies

1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

LHC	$pp$	$W_{\gamma p} \lesssim 8390 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 4504 \text{ GeV}$
LHC	$pPb(Ar)$	$W_{\gamma A} \lesssim 1500 (2130) \text{ GeV}$	$W_{\gamma\gamma} \lesssim 260 (480) \text{ GeV}$
LHC	$PbPb$	$W_{\gamma A} \lesssim 950 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 160 \text{ GeV}$
HERA	$ep$	$W_{\gamma p} \lesssim 200 \text{ GeV}$	-

- Photoproduction in  $pp$  collisions at LHC probes energies one order of magnitude larger than HERA.

# Motivation



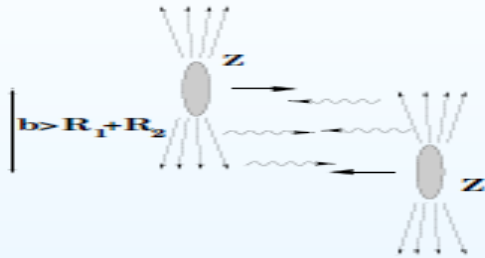
Center of mass energies

1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

LHC	<i>pp</i>	$W_{\gamma p} \lesssim 8390 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 4504 \text{ GeV}$
LHC	<u><i>pPb(Ar)</i></u>	$W_{\gamma A} \lesssim 1500 (2130) \text{ GeV}$	$W_{\gamma\gamma} \lesssim 260 (480) \text{ GeV}$
LHC	<u><i>PbPb</i></u>	$W_{\gamma A} \lesssim 950 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 160 \text{ GeV}$
HERA	<i>ep</i>	$W_{\gamma p} \lesssim 200 \text{ GeV}$	—

- Photoproduction in  $pA$  and  $AA$  collisions probes an unexplored regime of center of mass energies.

# Motivation



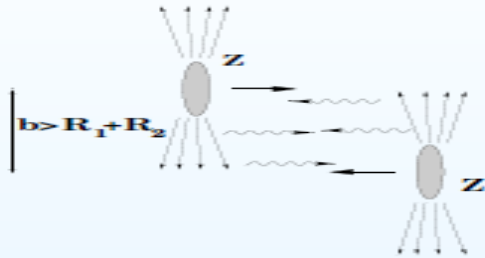
Center of mass energies

1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

LHC	$pp$	$W_{\gamma p} \lesssim 8390 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 4504 \text{ GeV}$
LHC	$pPb(Ar)$	$W_{\gamma A} \lesssim 1500 (2130) \text{ GeV}$	$W_{\gamma\gamma} \lesssim 260 (480) \text{ GeV}$
LHC	$PbPb$	$W_{\gamma A} \lesssim 950 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 160 \text{ GeV}$
HERA	$ep$	$W_{\gamma p} \lesssim 200 \text{ GeV}$	—

- ➔ The LHC is the world's most powerful collider not only for protons and lead ions but also for  $\gamma\gamma$  and  $\gamma h$  collisions.

# Motivation



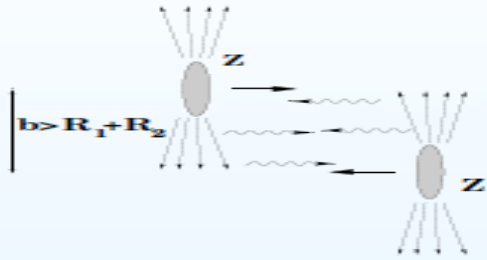
Center of mass energies

1.  $\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$
2.  $\gamma\gamma$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_1(\omega) \otimes n_2(\omega) \otimes \sigma^{\gamma\gamma \rightarrow X}(W_{\gamma\gamma})$

LHC	$pp$	$W_{\gamma p} \lesssim 8390 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 4504 \text{ GeV}$
LHC	$pPb(Ar)$	$W_{\gamma A} \lesssim 1500 (2130) \text{ GeV}$	$W_{\gamma\gamma} \lesssim 260 (480) \text{ GeV}$
LHC	$PbPb$	$W_{\gamma A} \lesssim 950 \text{ GeV}$	$W_{\gamma\gamma} \lesssim 160 \text{ GeV}$
HERA	$ep$	$W_{\gamma p} \lesssim 200 \text{ GeV}$	—

- Photon - induced interactions at LHC allows to study Quantum Chromodynamics in an unexplored regime of center of mass energies.

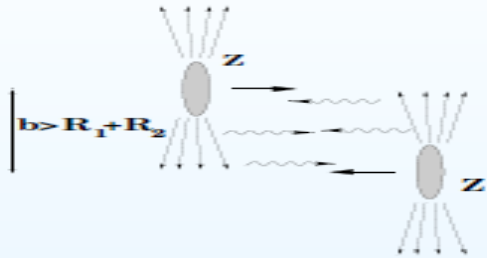
# Photon - Hadron Interactions: Inclusive and exclusive processes



$\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$



# Photon - Hadron Interactions: Inclusive and exclusive processes



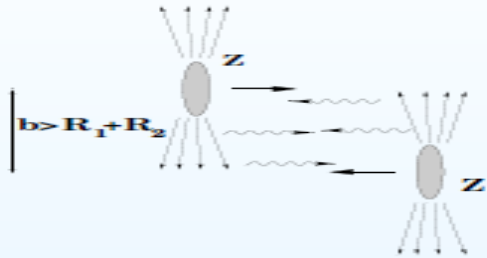
$$\gamma h \text{ Processes: } \sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$$

- Inclusive processes:  $\gamma p \rightarrow XY$

⇒ Heavy quark photoproduction ( $X = c\bar{c}, b\bar{b}$ )

The final state is characterized by **one rapidity gap** due to the dissociation of the hadron target ( $pp \rightarrow p \otimes XY$ ).

# Photon - Hadron Interactions: Inclusive and exclusive processes

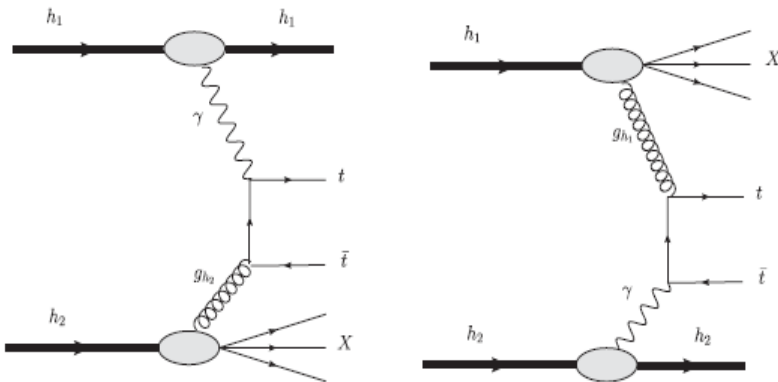


$$\gamma h \text{ Processes: } \sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$$

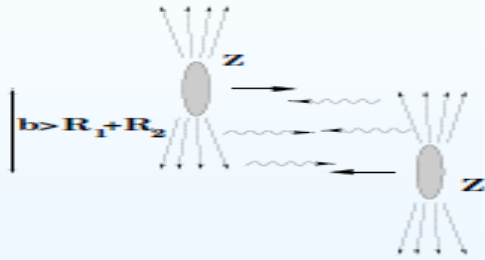
- Inclusive processes:  $\gamma p \rightarrow XY$

$\Rightarrow$  Heavy quark photoproduction ( $X = c\bar{c}, b\bar{b}$ )

The final state is characterized by **one rapidity gap** due to the dissociation of the hadron target ( $pp \rightarrow p \otimes XY$ ).



# Photon - Hadron Interactions: Inclusive and exclusive processes

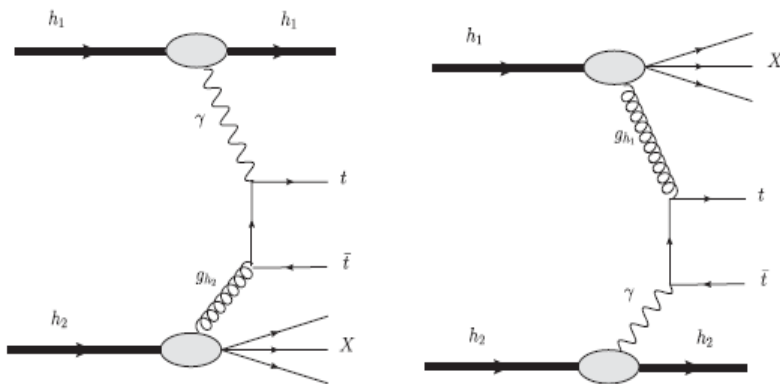


$$\gamma h \text{ Processes: } \sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$$

- Inclusive processes:  $\gamma p \rightarrow XY$

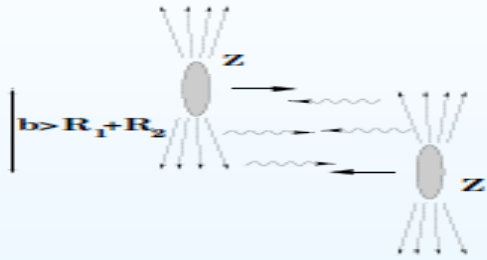
$\Rightarrow$  Heavy quark photoproduction ( $X = c\bar{c}, b\bar{b}$ )

The final state is characterized by **one rapidity gap** due to the dissociation of the hadron target ( $pp \rightarrow p \otimes XY$ ).



$X = t\bar{t}$  <sup>a</sup>VPG, PRD88, 054025 (2013)

# Photon - Hadron Interactions: Inclusive and exclusive processes



$$\gamma h \text{ Processes: } \sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$$

- Inclusive processes:  $\gamma p \rightarrow XY$

$\Rightarrow$  Heavy quark photoproduction ( $X = c\bar{c}, b\bar{b}$ )

The final state is characterized by **one rapidity gap** due to the dissociation of the hadron target ( $pp \rightarrow p \otimes XY$ ).

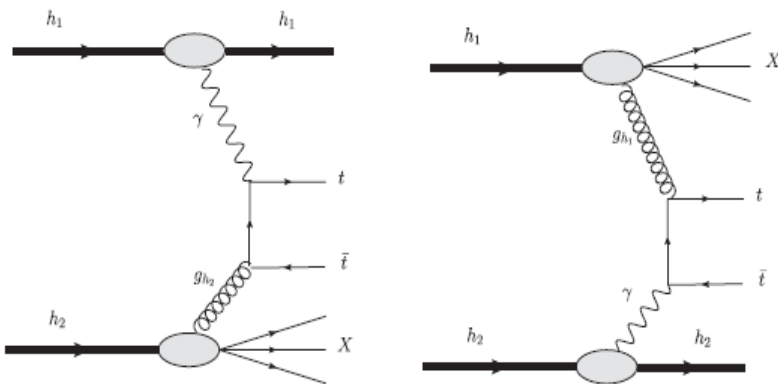
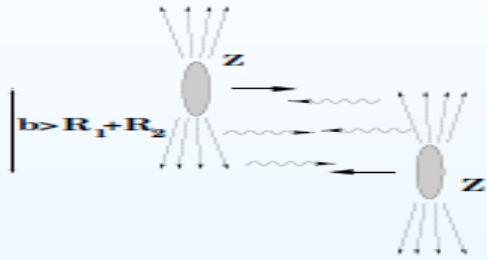


TABLE I. The integrated cross section (events rate) for the photoproduction of top quarks in  $pp$ ,  $pPb$ , and  $PbPb$  collisions at LHC energies.

$pp$	MRST	CT10
$\sqrt{s} = 8 \text{ TeV}$	0.739 pb (73900)	0.764 pb (76400)
$\sqrt{s} = 14 \text{ TeV}$	2.50 pb (250000)	2.53 pb (253000)
$pPb$	MRST	MRST + EPS09
$\sqrt{s} = 5.5 \text{ TeV}$	0.036 nb (5.4/3600)	0.038 nb (5.7/3800)
$\sqrt{s} = 8.8 \text{ TeV}$	0.159 nb (23.85/15900)	0.165 nb (24.75/16500)
$PbPb$	MRST	MRST $\pm$ EPS09
$\sqrt{s} = 5.5 \text{ TeV}$	0.42 nb (0.18)	0.40 nb (0.17)

$X = t\bar{t}$  <sup>a</sup>VPG, PRD88, 054025 (2013)

# Photon - Hadron Interactions: Inclusive and exclusive processes



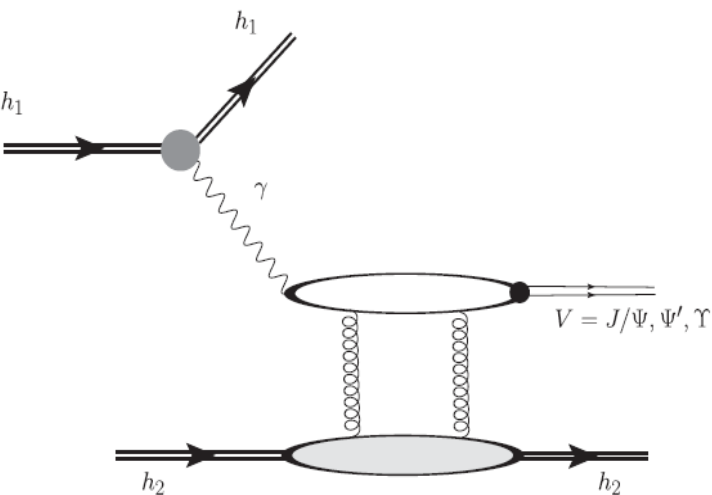
$\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$

- Exclusive processes:  $\gamma p \rightarrow X p$

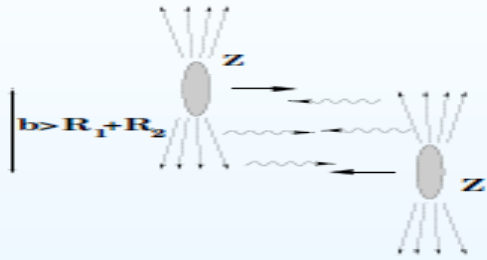
$\Rightarrow$  Heavy vector meson photoproduction ( $X = J/\Psi, \Upsilon$ )

The final state is characterized by two rapidity gaps

( $pp \rightarrow p \otimes X \otimes p$ ).



# Photon - Hadron Interactions: Inclusive and exclusive processes



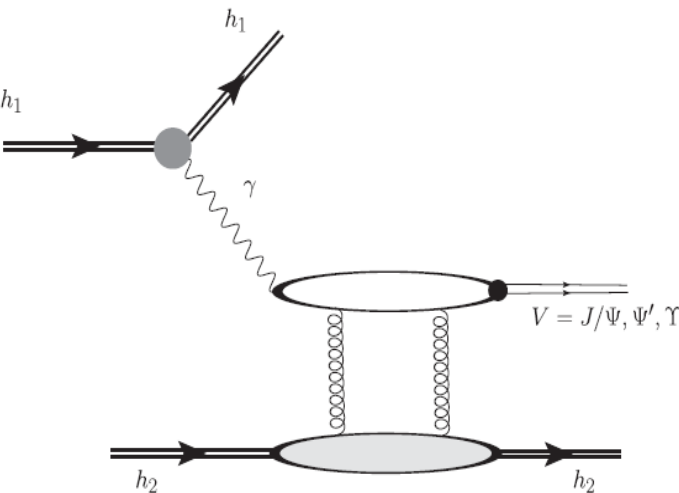
$$\gamma h \text{ Processes: } \sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$$

- Exclusive processes:  $\gamma p \rightarrow X p$

⇒ Heavy vector meson photoproduction ( $X = J/\Psi, \Upsilon$ )

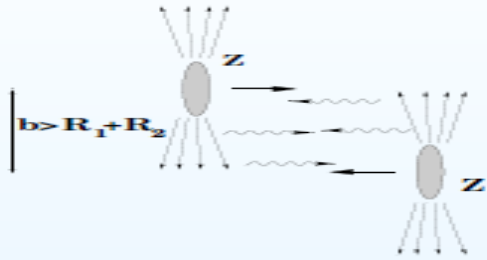
The final state is characterized by two rapidity gaps

( $pp \rightarrow p \otimes X \otimes p$ ).



- Cross section is proportional to the **square** of the proton/nuclear gluon distribution.

# Photon - Hadron Interactions: Inclusive and exclusive processes



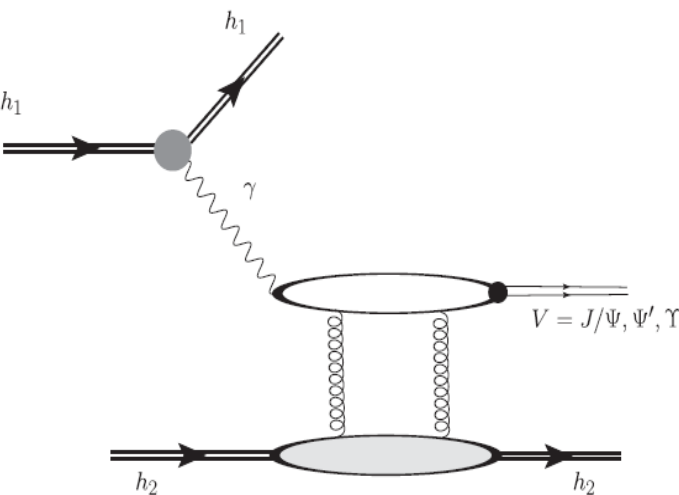
$\gamma h$  Processes:  $\sigma(h_1 h_2 \rightarrow X) = n_h(\omega) \otimes \sigma^{\gamma h \rightarrow X}(W_{\gamma h})$

- Exclusive processes:  $\gamma p \rightarrow X p$

$\Rightarrow$  Heavy vector meson photoproduction ( $X = J/\Psi, \Upsilon$ )

The final state is characterized by two rapidity gaps

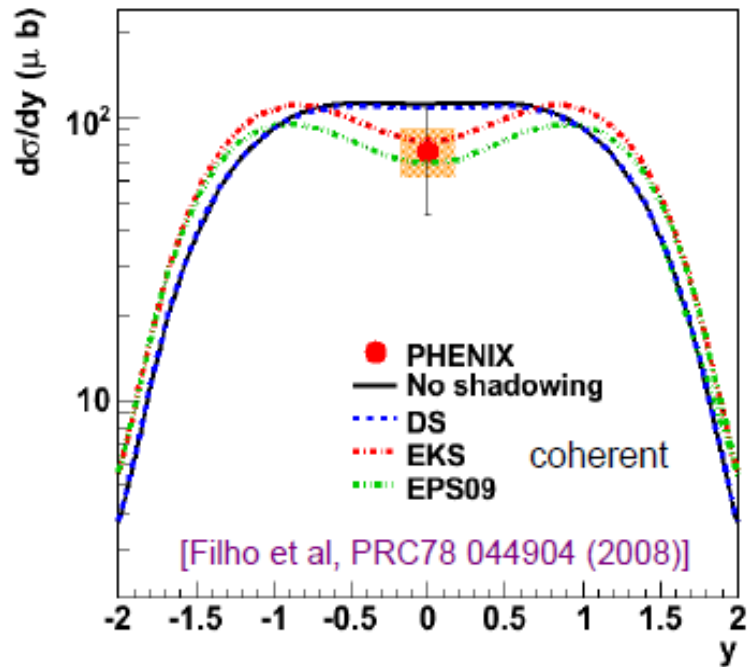
( $pp \rightarrow p \otimes X \otimes p$ ).



- Cross section is proportional to the **square** of the proton/nuclear gluon distribution.
- Diffractive vector meson photoproduction in UPHIC is a **probe** of the gluon distribution <sup>a</sup>

# Diffraction vector meson photoproduction in UPHIC

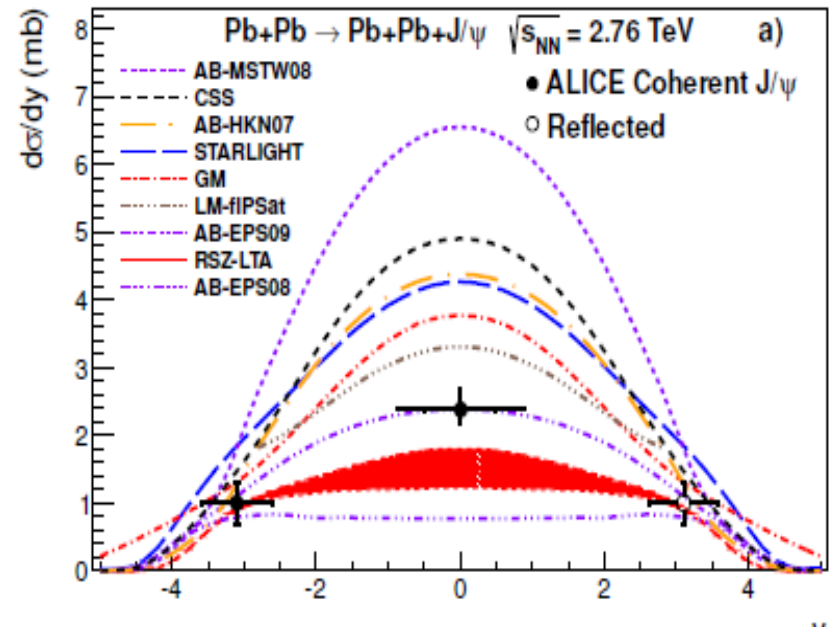
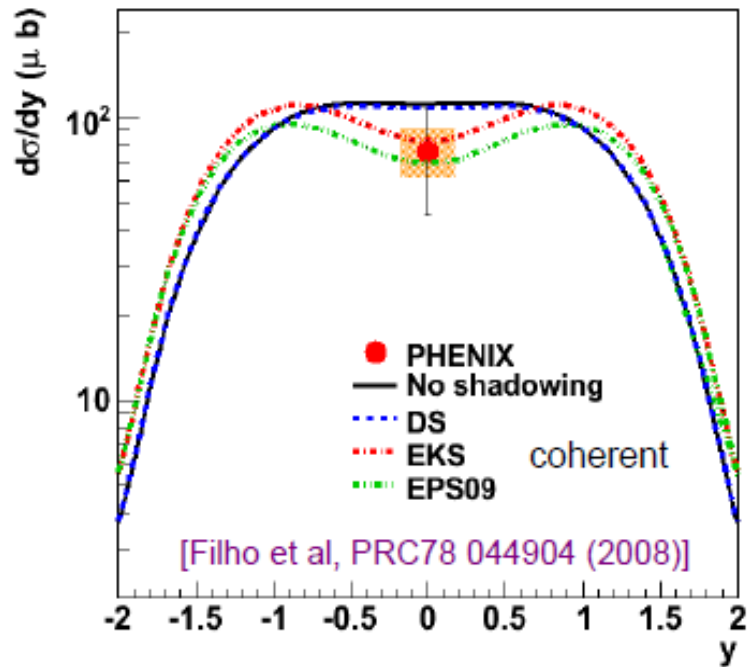
## Probing the nuclear gluon distribution





# Diffraction vector meson photoproduction in UPHIC

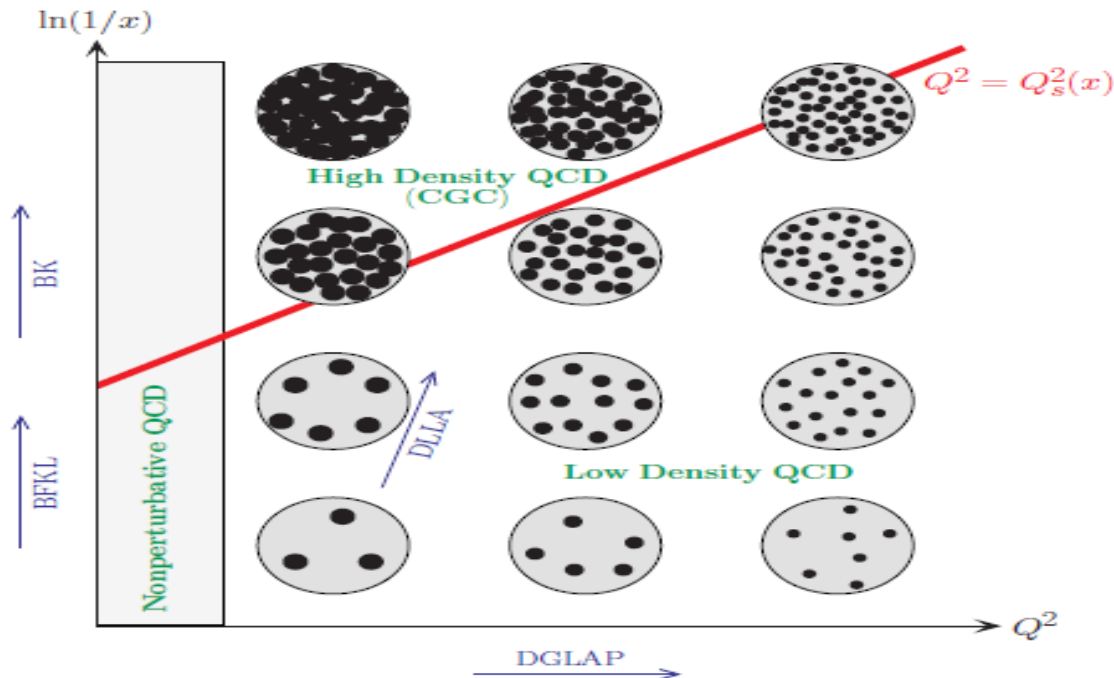
## Probing the nuclear gluon distribution



- Since  $x = M_{J/\psi} / \sqrt{s} \exp(-y)$  we have:  
 $y = -3 \Rightarrow x = 0.02$   
 $y = 0 \Rightarrow x = 0.001$  in  $xg_A(x, Q^2)$ .

# Diffraction vector meson photoproduction in UPHIC

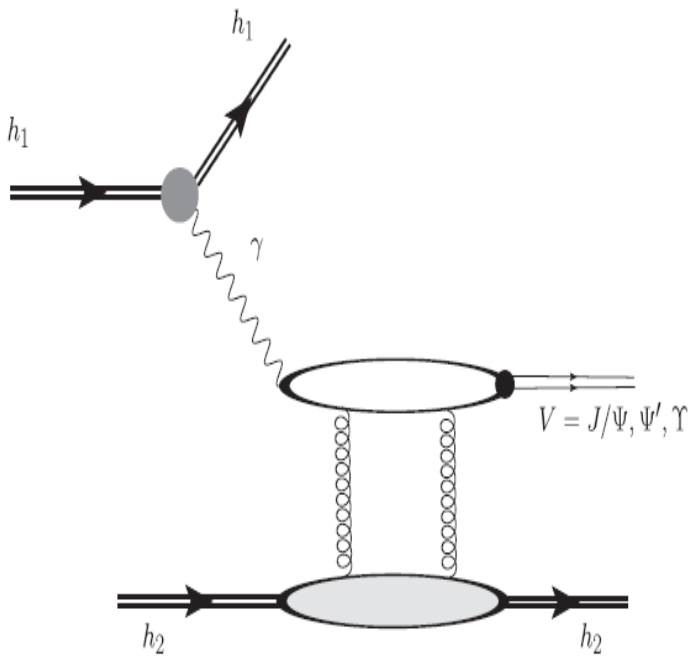
## Probing the QCD dynamics at high energies



- Linear QCD evolution equations predict a power growth of gluon distribution as  $x \rightarrow 0$  (violates unitarity).
- Number of gluons in the nucleon becomes so large that gluon recombine  $\Rightarrow$  Nonlinear effects
- Saturation scale  $Q_s$  (energy and atomic number dependent) defines the onset of nonlinear QCD dynamics.

# Diffractive vector meson photoproduction in UPHIC

## Probing the QCD dynamics at high energies



- Diffractive vector meson photoproduction in photon-induced interactions is a **probe** of the nonlinear effects in the QCD dynamics at high energies and the vector meson wave function <sup>a</sup>.

# Diffraction vector meson photoproduction in UPHIC

## Color Dipole Formalism

$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{d^2b dy} = [\omega N_{h_1}(\omega, b) \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega)]_{\omega_L} + [\omega N_{h_2}(\omega, b) \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega)]_{\omega_R}$$

$$\sigma(\gamma h \rightarrow V h) = \int_{-\infty}^0 \frac{d\sigma}{dt} dt = \frac{1}{16\pi} \int_{-\infty}^0 |\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta)|^2 dt$$

$$\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta) = i \int dz d^2\mathbf{r} d^2\mathbf{b}_h e^{-i[\mathbf{b}_h - (1-z)\mathbf{r}] \cdot \Delta} \underbrace{(\Psi^{V*} \Psi)_T}_{\text{Overlap functions for Vector Mesons}} 2\mathcal{N}_h(x, \mathbf{r}, \mathbf{b}_h)$$

Overlap functions for **Vector Mesons**:

$$(\Psi_V^* \Psi)_T = \frac{\hat{e}_f e}{4\pi} \frac{N_c}{\pi z(1-z)} \{m_f^2 K_0(\epsilon r) \phi_T(r, z) - [z^2 + (1-z)^2] \epsilon K_1(\epsilon r) \partial_r \phi_T(r, z)\}$$

# Diffraction vector meson photoproduction in UPHIC

## Color Dipole Formalism

$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{d^2b dy} = [\omega N_{h_1}(\omega, b) \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega)]_{\omega_L} + [\omega N_{h_2}(\omega, b) \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega)]_{\omega_R}$$

$$\sigma(\gamma h \rightarrow V h) = \int_{-\infty}^0 \frac{d\sigma}{dt} dt = \frac{1}{16\pi} \int_{-\infty}^0 |\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta)|^2 dt$$

$$\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta) = i \int dz d^2\mathbf{r} d^2\mathbf{b}_h e^{-i[\mathbf{b}_h - (1-z)\mathbf{r}] \cdot \Delta} (\Psi^{V*} \Psi)_T 2\mathcal{N}_h(x, \mathbf{r}, \mathbf{b}_h)$$



Forward dipole - hadron scattering amplitude: **Determined by the QCD dynamics**

# Diffraction vector meson photoproduction in UPHIC

## Color Dipole Formalism

$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{d^2b dy} = [\omega N_{h_1}(\omega, b) \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega)]_{\omega_L} + [\omega N_{h_2}(\omega, b) \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega)]_{\omega_R}$$

$$\sigma(\gamma h \rightarrow V h) = \int_{-\infty}^0 \frac{d\sigma}{dt} dt = \frac{1}{16\pi} \int_{-\infty}^0 |\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta)|^2 dt$$

$$\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta) = i \int dz d^2\mathbf{r} d^2\mathbf{b}_h e^{-i[\mathbf{b}_h - (1-z)\mathbf{r}] \cdot \Delta} (\Psi^{V*} \Psi)_T 2\mathcal{N}_h(x, \mathbf{r}, \mathbf{b}_h)$$



Forward dipole - hadron scattering amplitude: **Determined by the QCD dynamics**

- Proton: **Constrained by Hera data for inclusive and exclusive processes**

# Diffraction vector meson photoproduction in UPHIC

## Color Dipole Formalism

$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{d^2b dy} = [\omega N_{h_1}(\omega, b) \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega)]_{\omega_L} + [\omega N_{h_2}(\omega, b) \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega)]_{\omega_R}$$

$$\sigma(\gamma h \rightarrow V h) = \int_{-\infty}^0 \frac{d\sigma}{dt} dt = \frac{1}{16\pi} \int_{-\infty}^0 |\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta)|^2 dt$$

$$\mathcal{A}_T^{\gamma h \rightarrow V h}(x, \Delta) = i \int dz d^2\mathbf{r} d^2\mathbf{b}_h e^{-i[\mathbf{b}_h - (1-z)\mathbf{r}] \cdot \Delta} (\Psi^{V*} \Psi)_T \mathcal{N}_h(x, \mathbf{r}, \mathbf{b}_h)$$



Forward dipole - hadron scattering amplitude: **Determined by the QCD dynamics**

- Proton: **Constrained by Hera data for inclusive and exclusive processes**

- Nucleus:  $\mathcal{N}_A(x, \mathbf{r}, \mathbf{b}_A) = 1 - \exp \left[ -\frac{1}{2} \sigma_{dp}(x, \mathbf{r}^2) T_A(\mathbf{b}_A) \right]$

$$\sigma_{dp}(x, \mathbf{r}^2) = 2 \int d^2\mathbf{b}_p \mathcal{N}_p(x, \mathbf{r}, \mathbf{b}_p)$$

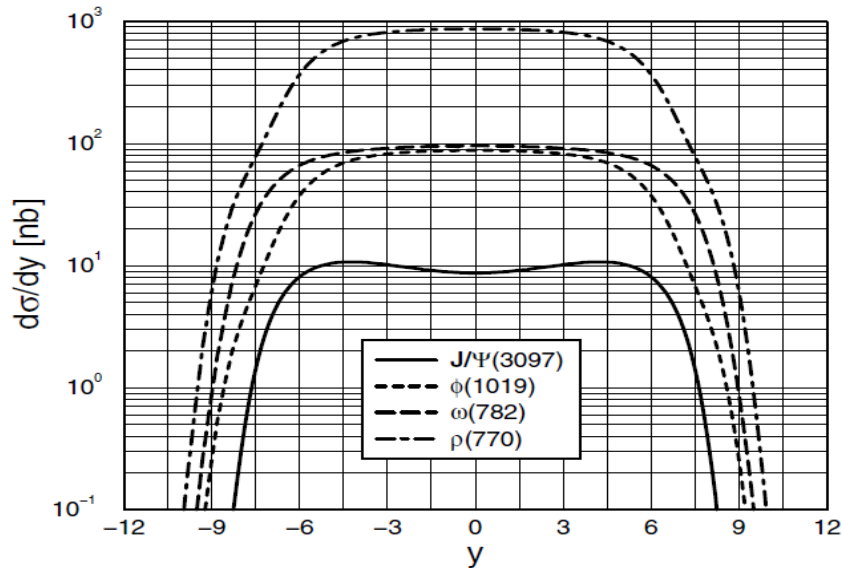


*Sums all multiple elastic rescatterings of the dipole.*

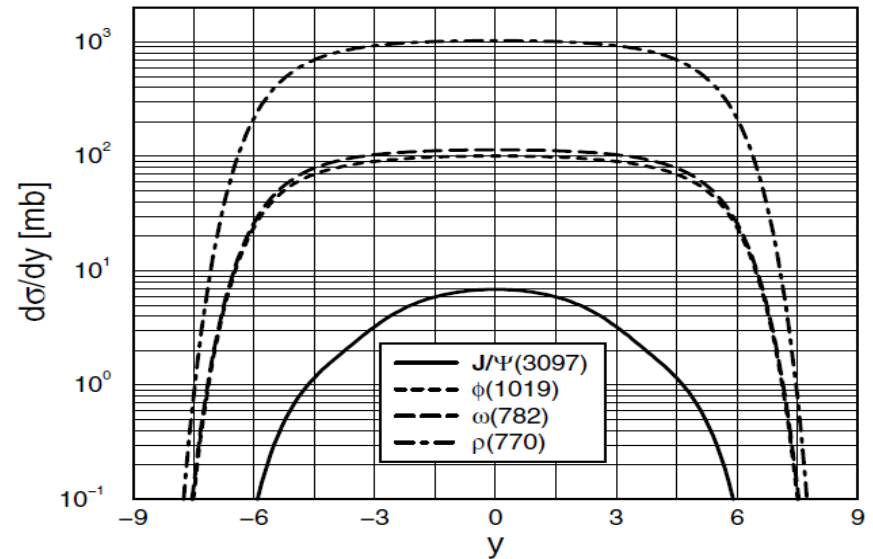
# Diffraction vector meson photoproduction in UPHIC

## First predictions using the Color Dipole Formalism

LHC:  $pp \rightarrow pp + (V_M = \rho, \omega, \phi, J/\Psi)$



LHC:  $PbPb \rightarrow PbPb + (V_M = \rho, \omega, \phi, J/\Psi)$



- Golec-Biernat - Wusthoff (GBW) :

$$\mathcal{N}^P(\hat{x}, \mathbf{r}, \mathbf{b}) = \mathcal{N}^P(\hat{x}, \mathbf{r}) S(\mathbf{b})$$

$$\mathcal{N}^P(x, \mathbf{r}) = 1 - \exp\left[-\frac{Q_s^2 r^2}{4}\right]$$

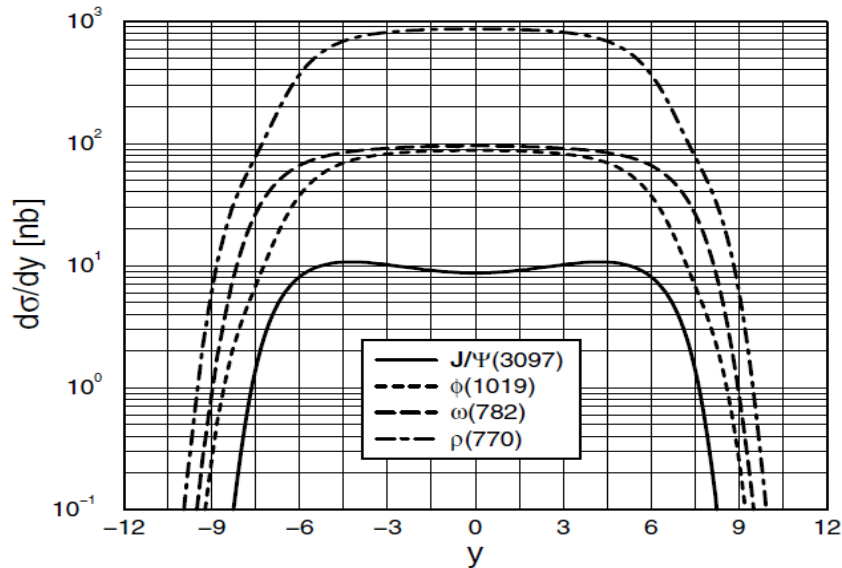
- DGKP model for the vector meson WF



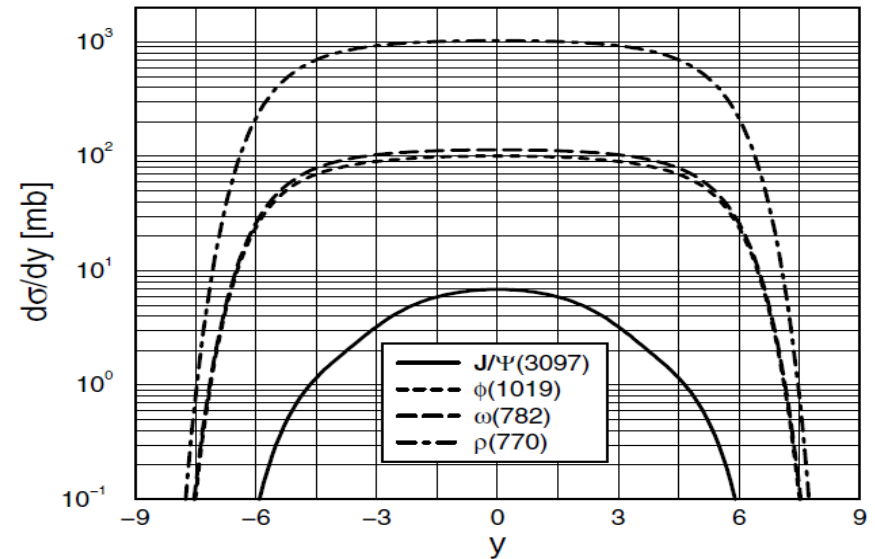
# Diffraction vector meson photoproduction in UPHIC

## First predictions using the Color Dipole Formalism

LHC:  $pp \rightarrow pp + (V_M = \rho, \omega, \phi, J/\Psi)$



LHC:  $PbPb \rightarrow PbPb + (V_M = \rho, \omega, \phi, J/\Psi)$



During the last 11 years, several authors have updated these predictions taking into account the improvements in our understanding of the QCD dynamics and for the description of exclusive processes in electron - proton scattering (e.g. impact parameter dependence, skewness, vector meson W. F., ...)

# Diffraction vector meson photoproduction in UPHIC

# Diffractive vector meson photoproduction in UPHIC

## Overlap functions for Vector Mesons:

$$(\Psi_V^* \Psi)_T = \frac{\hat{e}_f e}{4\pi} \frac{N_c}{\pi z(1-z)} \{m_f^2 K_0(\epsilon r) \phi_T(r, z) - [z^2 + (1-z)^2] \epsilon K_1(\epsilon r) \partial_r \phi_T(r, z)\}$$

- Gaus - LC :

$$\phi_T(r, z) = N_T [z(1-z)]^2 \exp(-r^2/2R_T^2)$$

- Boosted Gaussian :

$$\phi_{T,L}(r, z) = N_{T,L} z(1-z) \exp\left[-\frac{m_f^2 R^2}{8z(1-z)} - \frac{2z(1-z)r^2}{R^2} + \frac{m_f^2 R^2}{2}\right]$$

# Diffractive vector meson photoproduction in UPHIC

## Overlap functions for Vector Mesons:

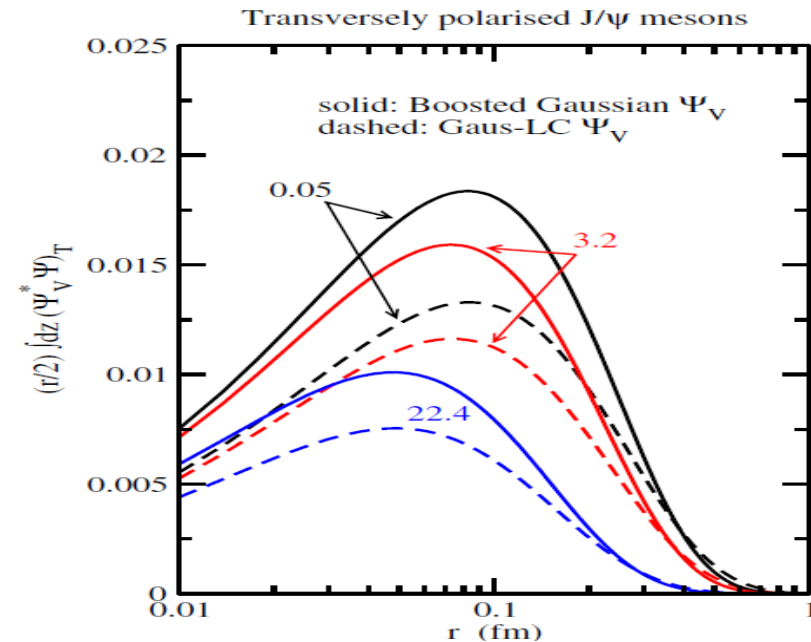
$$(\Psi_V^* \Psi)_T = \frac{\hat{e}_f e}{4\pi} \frac{N_c}{\pi z(1-z)} \{ m_f^2 K_0(\epsilon r) \phi_T(r, z) - [z^2 + (1-z)^2] \epsilon K_1(\epsilon r) \partial_r \phi_T(r, z) \}$$

- Gaus - LC :

$$\phi_T(r, z) = N_T [z(1-z)]^2 \exp(-r^2/2R_T^2)$$

- Boosted Gaussian :

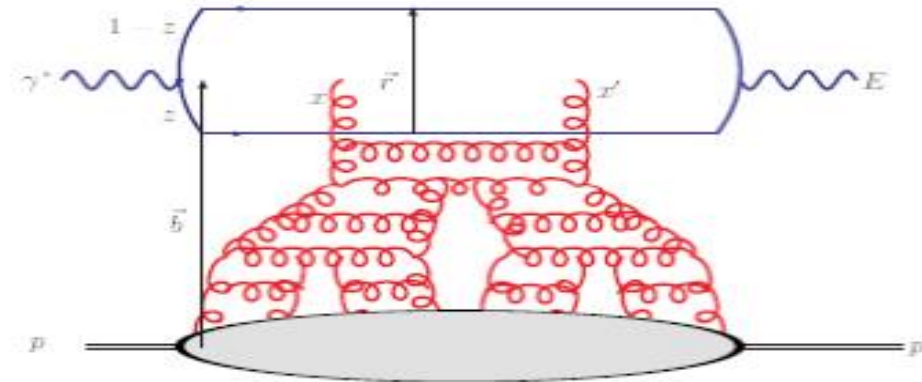
$$\phi_{T,L}(r, z) = N_{T,L} z(1-z) \exp \left[ -\frac{m_f^2 R^2}{8z(1-z)} - \frac{2z(1-z)r^2}{R^2} + \frac{m_f^2 R^2}{2} \right]$$



Kowalski, Motyka and Watt (06)

# Diffractive vector meson photoproduction in UPHIC

Dipole - proton scattering amplitude:

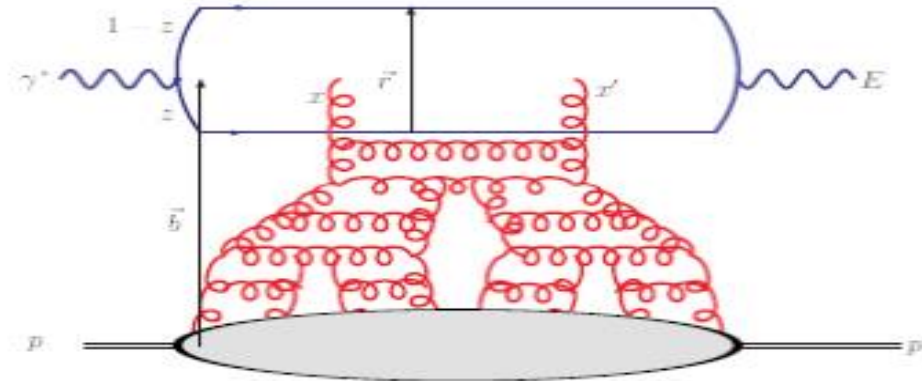


- bCGC : 
$$\mathcal{N}^P(\hat{x}, r, b) = \begin{cases} \mathcal{N}_0 \left( \frac{rQ_s(b)}{2} \right)^{2\left(\gamma_s + \frac{\ln(2/rQ_s(b))}{\kappa\lambda Y}\right)} & rQ_s(b) \leq 2 \\ 1 - e^{-A \ln^2(BrQ_s(b))} & rQ_s(b) > 2 \end{cases}$$

$$Q_s(b) \equiv Q_s(\hat{x}, b) = \left( \frac{x_0}{\hat{x}} \right)^{\frac{\lambda}{2}} \left[ \exp \left( -\frac{b^2}{2B_{\text{CGC}}} \right) \right]^{\frac{1}{2\gamma_s}}.$$

# Diffractive vector meson photoproduction in UPHIC

Dipole - proton scattering amplitude:

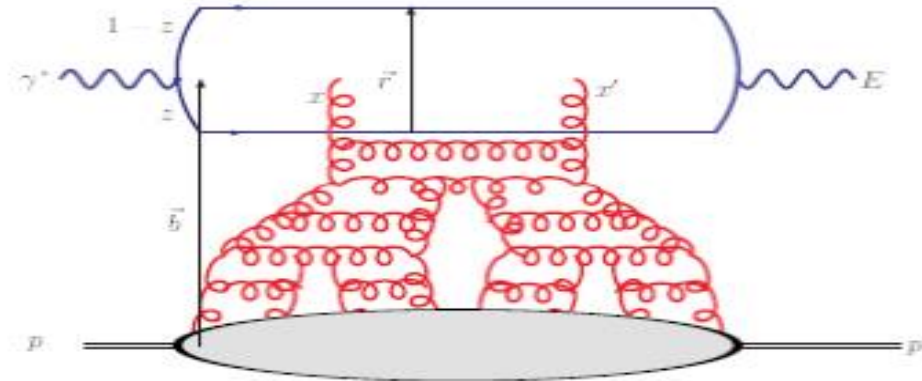


• bCGC: 
$$\mathcal{N}^P(\hat{x}, r, \mathbf{b}) = \begin{cases} \mathcal{N}_0 \left( \frac{rQ_s(b)}{2} \right)^{2(\gamma_s + \frac{\ln(2/rQ_s(b))}{\kappa\lambda Y})} & rQ_s(b) \leq 2 \\ 1 - e^{-A \ln^2(BrQ_s(b))} & rQ_s(b) > 2 \end{cases}$$

- Proposed originally by Kowalski, Motyka and Watt (06)
- Parameters of the model updated considering the high precision combined HERA data (Rezaeian, Schmidt, 13)

# Diffractive vector meson photoproduction in UPHIC

Dipole - proton scattering amplitude:



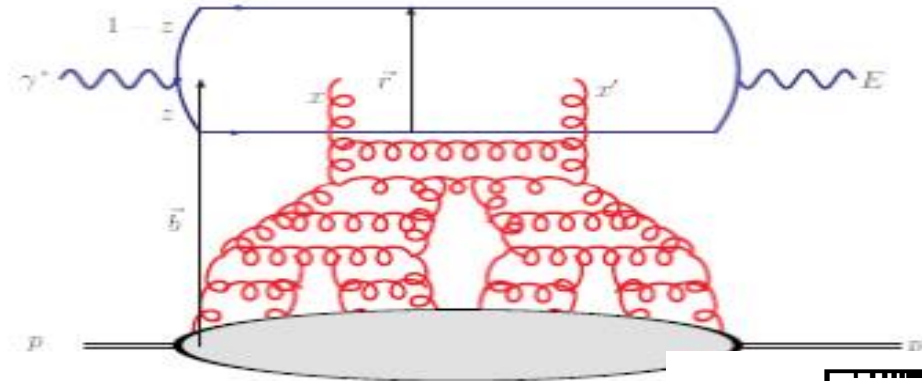
- Iancu - Itakura - Munier (IIM/CGC):

$$\mathcal{N}^p(\hat{x}, \mathbf{r}, \mathbf{b}) = \mathcal{N}^p(\hat{x}, \mathbf{r}) S(\mathbf{b})$$

$$\mathcal{N}^p(x, \mathbf{r}) = \begin{cases} \mathcal{N}_0 \left( \frac{r Q_s}{2} \right)^{2 \left( \gamma_s + \frac{\ln(2/r Q_s)}{\kappa \lambda Y} \right)}, & \text{for } r Q_s(x) \leq 2, \\ 1 - e^{-a \ln^2(b r Q_s)}, & \text{for } r Q_s(x) > 2, \end{cases}$$

# Diffractive vector meson photoproduction in UPHIC

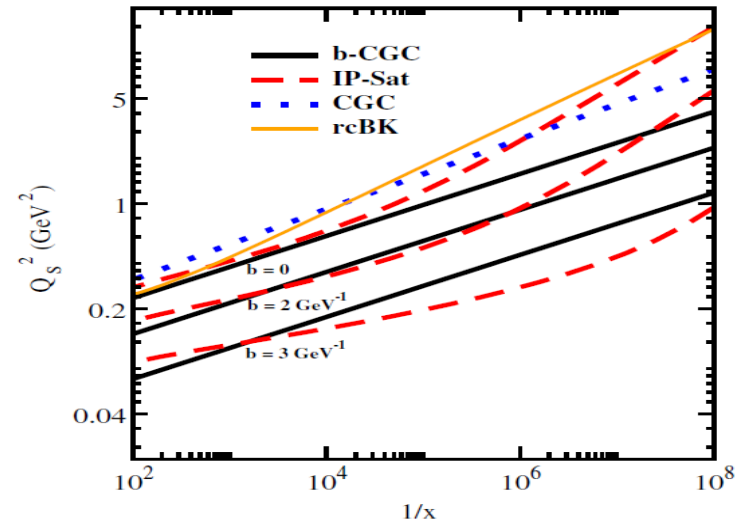
Dipole - proton scattering amplitude:



- Iancu - Itakura - Munier (IIM/CGC):

$$\mathcal{N}^P(\hat{x}, \mathbf{r}, \mathbf{b}) = \mathcal{N}^P(\hat{x}, \mathbf{r}) S(\mathbf{b})$$

$$\mathcal{N}^P(x, \mathbf{r}) = \begin{cases} \mathcal{N}_0 \left( \frac{r Q_s}{2} \right)^{2 \left( \gamma_s + \frac{\ln(2/r Q_s)}{\kappa \lambda Y} \right)}, & \text{for } r Q_s(x) \leq 2, \\ 1 - e^{-a \ln^2(br Q_s)}, & \text{for } r Q_s(x) > 2, \end{cases}$$



Rezaeian, Schmidt (13)

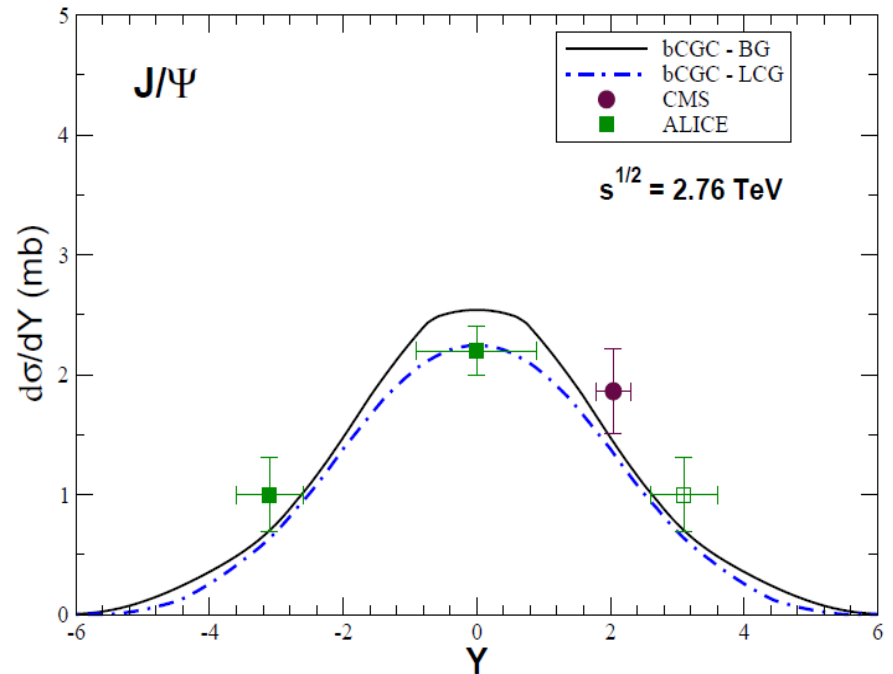
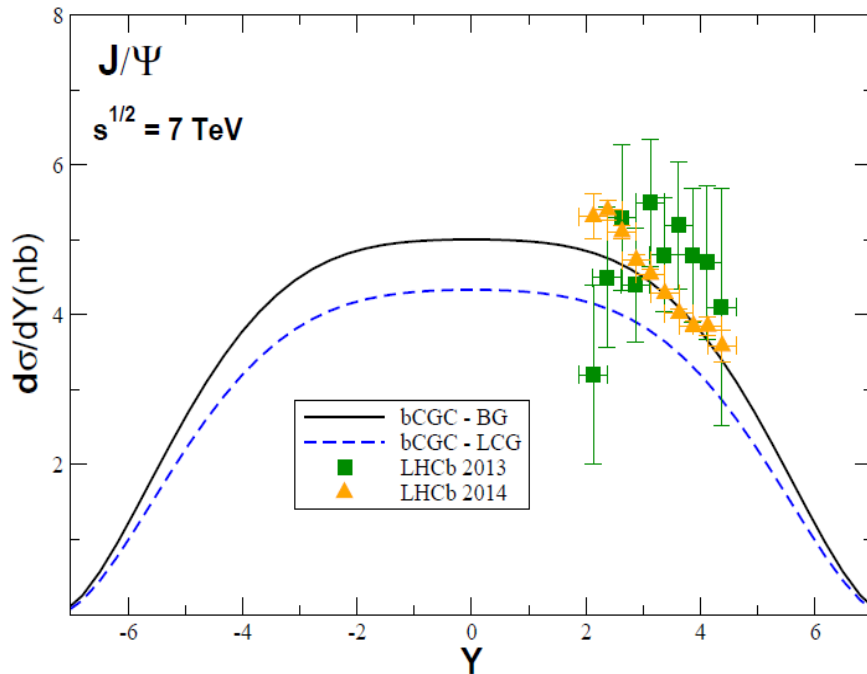


# Comparison with the Run I data

- Diffractive  $J/\Psi$  photoproduction in hadronic collisions <sup>a</sup>

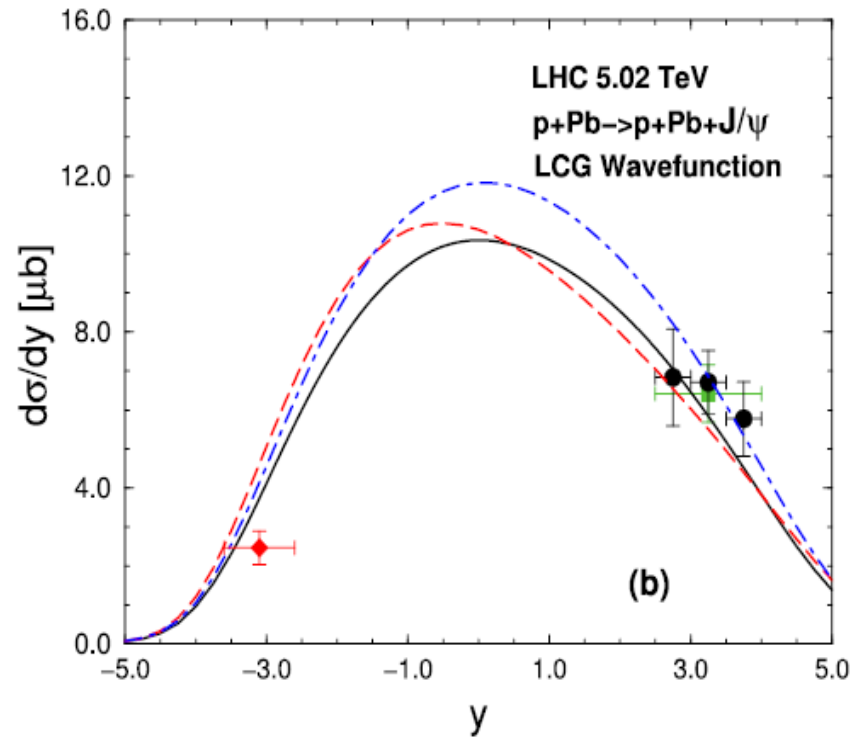
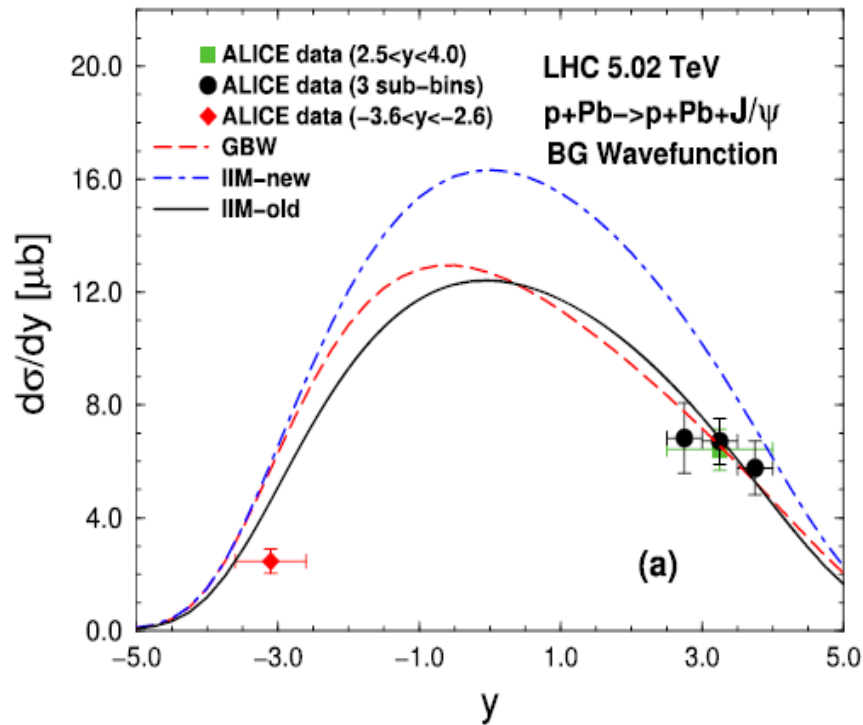
# Comparison with the Run I data

● Diffractive  $J/\Psi$  photoproduction in hadronic collisions <sup>a</sup>



# Comparison with the Run I data

● Diffractive  $J/\Psi$  photoproduction in hadronic collisions <sup>a</sup>

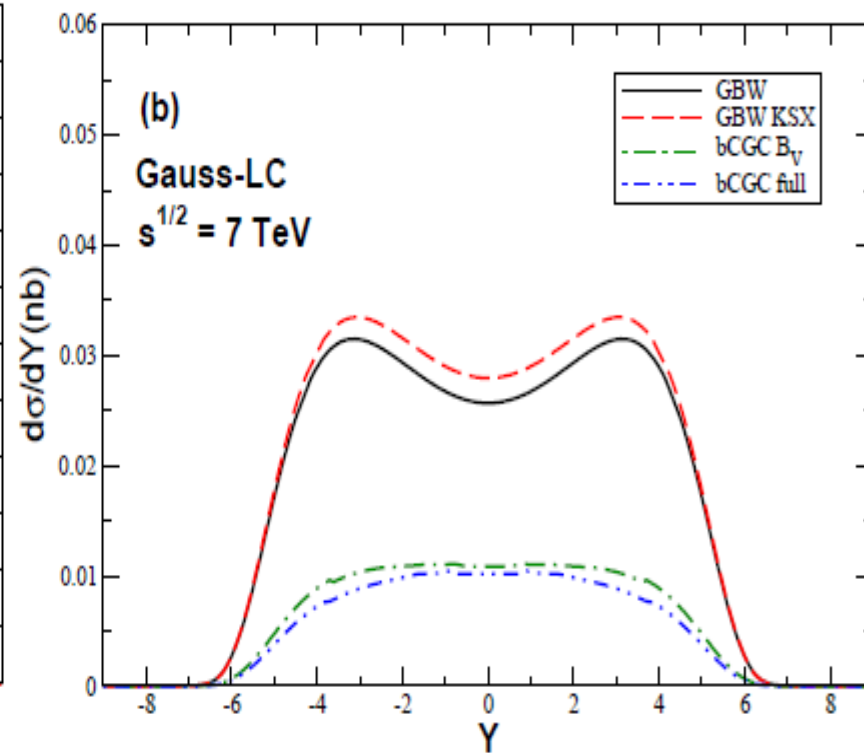
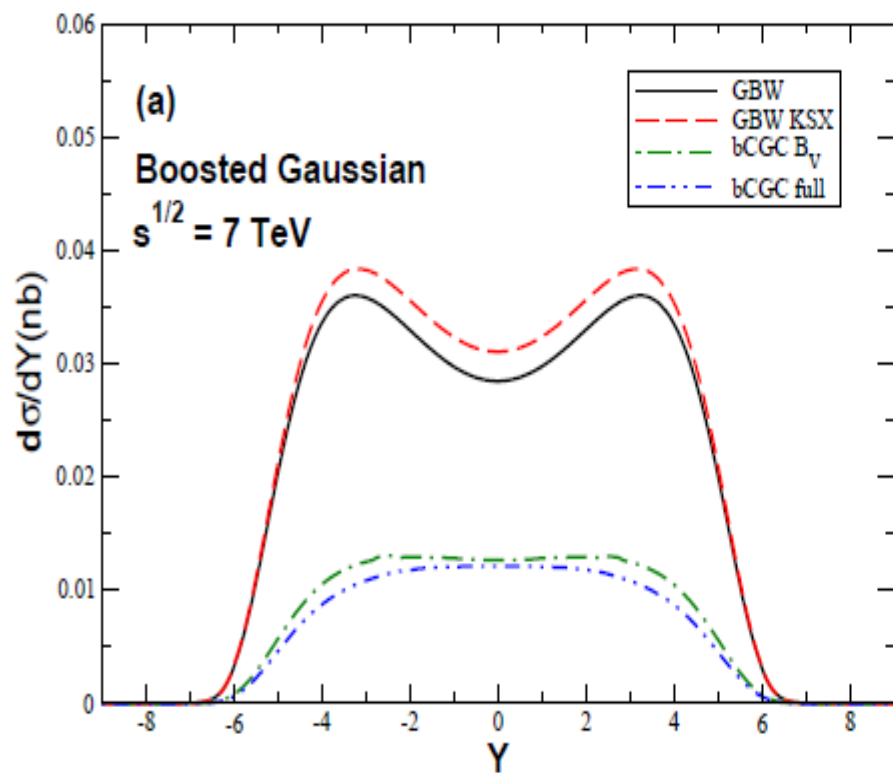


# Comparison with the Run I data

- Diffractive  $\Upsilon$  photoproduction in hadronic collisions <sup>b</sup>

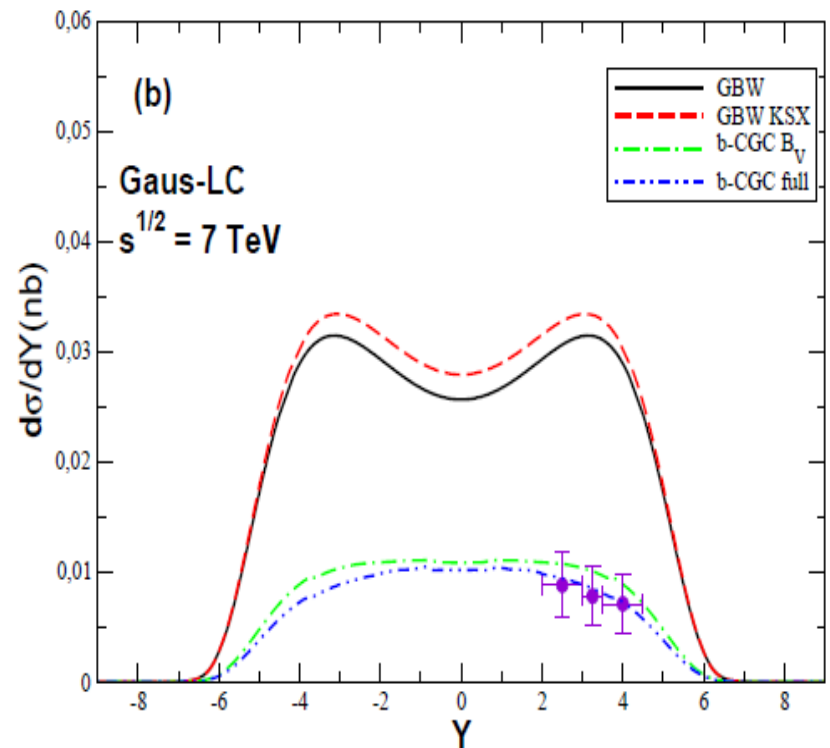
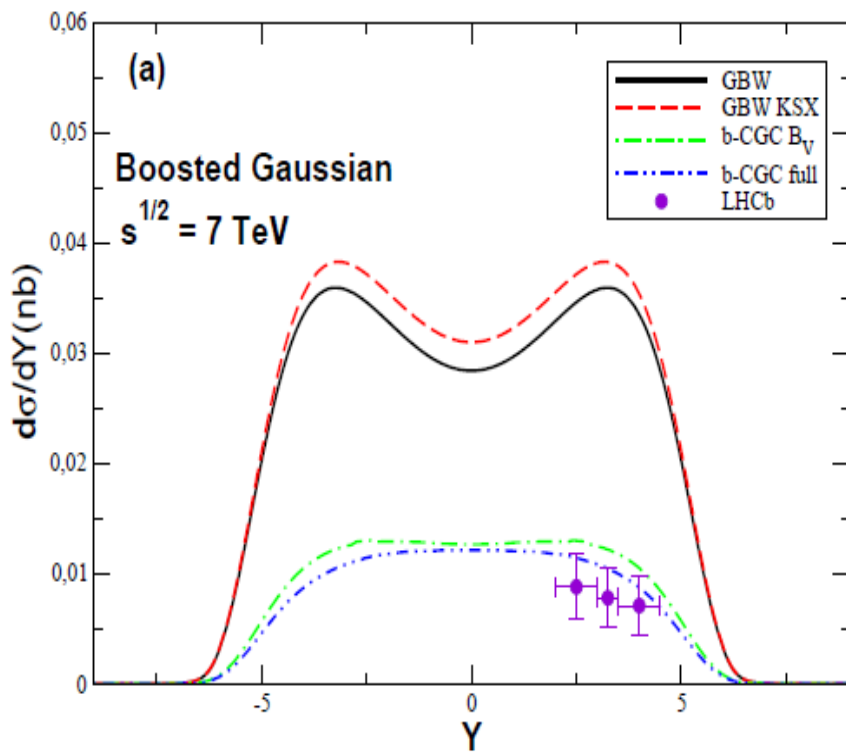
# Comparison with the Run I data

● Diffractive  $\Upsilon$  photoproduction in hadronic collisions <sup>b</sup>



# Comparison with the Run I data

● Diffractive  $\Upsilon$  photoproduction in hadronic collisions <sup>b</sup>



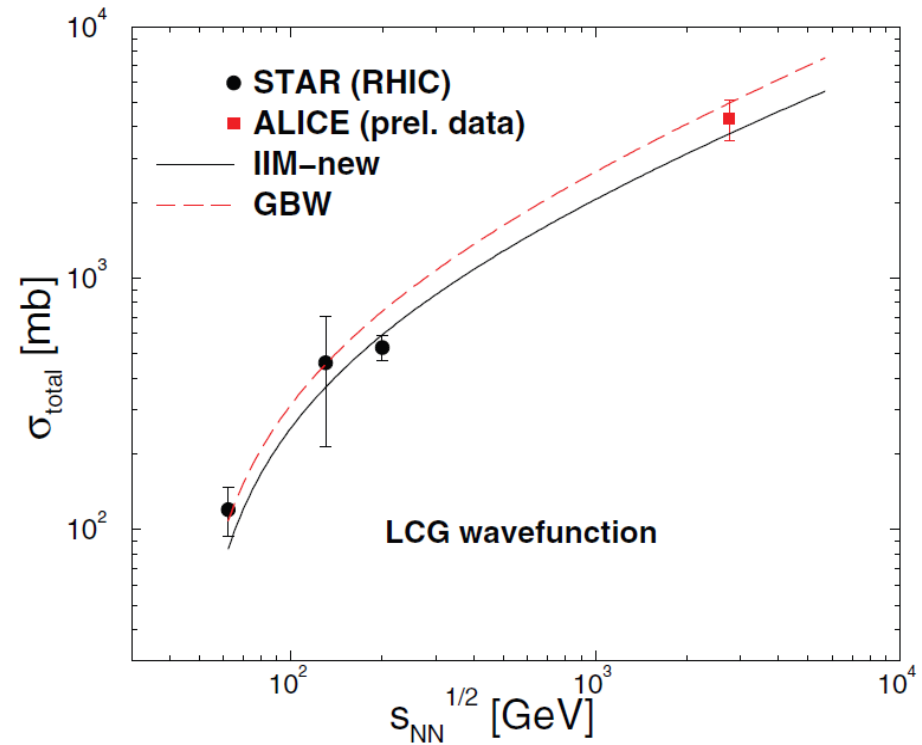
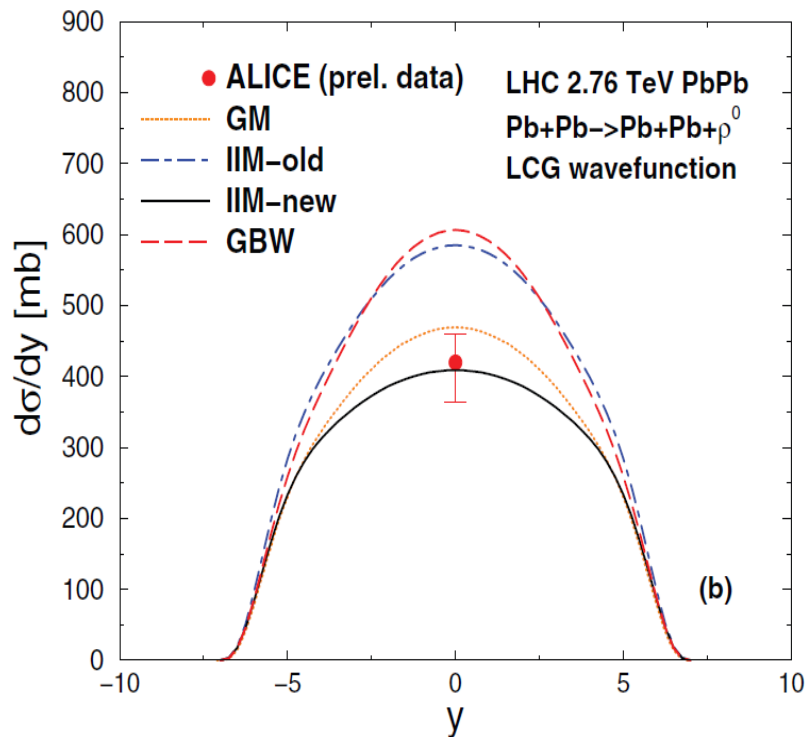
# Comparison with the Run I data

- Diffractive  $\rho$  photoproduction in hadronic collisions <sup>c</sup>

(<sup>c</sup>) VPG, Machado, EPJC 40, 519 (2005); PRC80, 054901 (2009); PRC84, 011902 (2011)  
Machado, dos Santos, PRC91, 025203 (2015)

# Comparison with the Run I data

● Diffractive  $\rho$  photoproduction in hadronic collisions <sup>c</sup>



(<sup>c</sup>) VPG, Machado, EPJC 40, 519 (2005); PRC80, 054901 (2009); PRC84, 011902 (2011)  
 Machado, dos Santos, PRC91, 025203 (2015)



# Predictions for the Run II

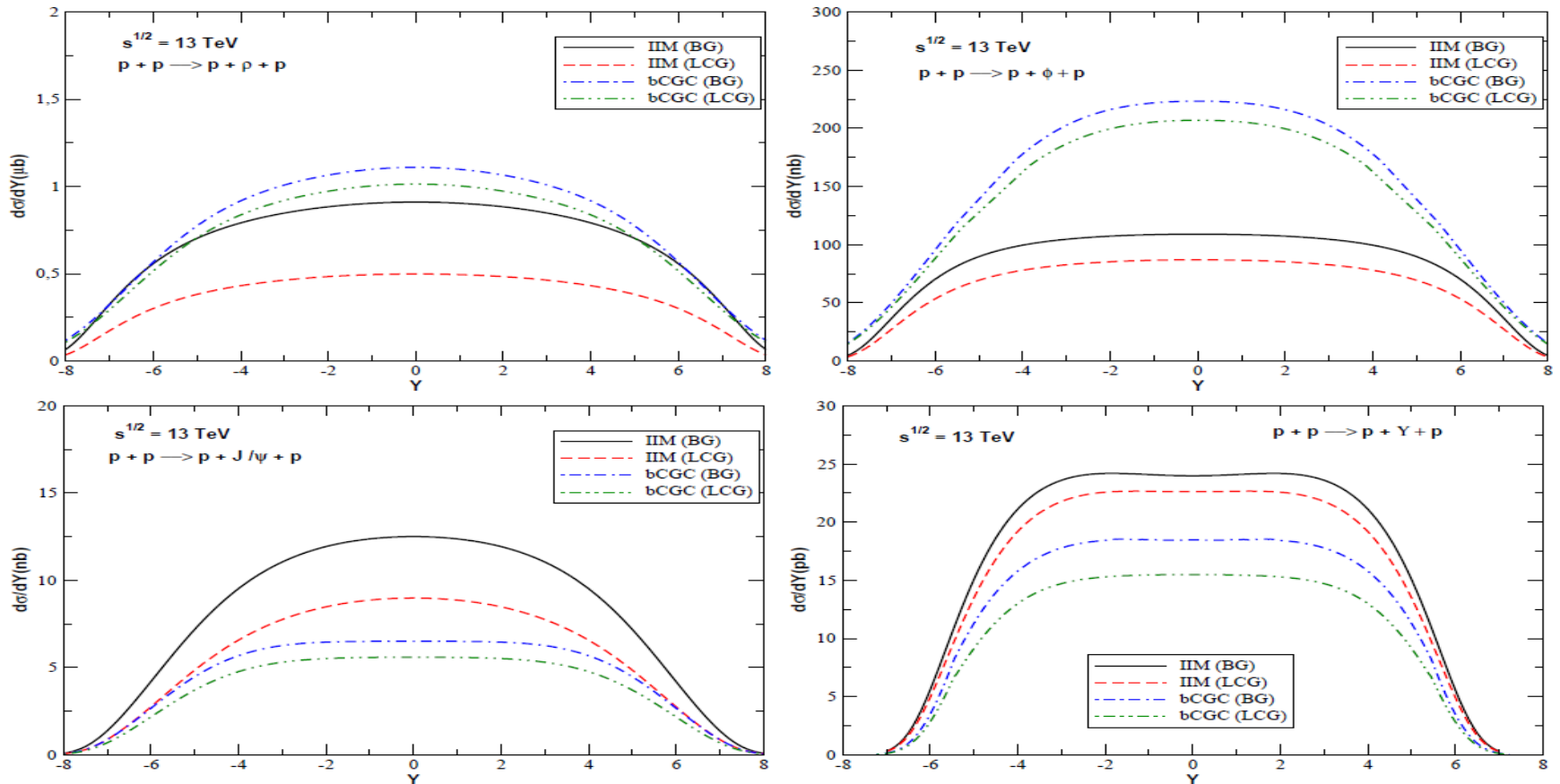
(\* ) VPG, Machado, Moreira, Navarra, dos Santos (paper in preparation)

# Predictions for the Run II

- Diffractive vector meson photoproduction in proton - proton collisions

# Predictions for the Run II

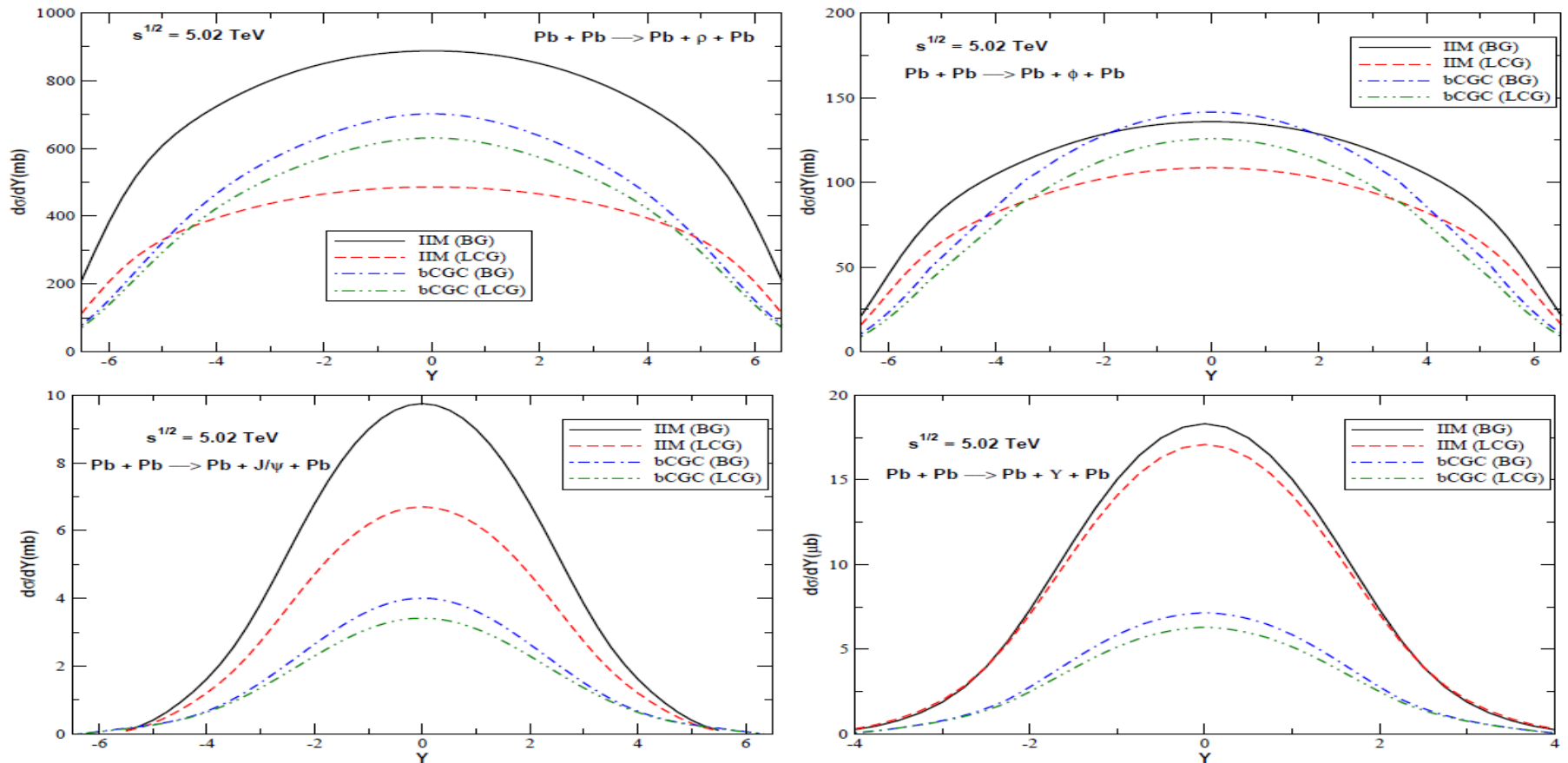
➤ Diffractive vector meson photoproduction in proton - proton collisions



(\* ) VPG, Machado, Moreira, Navarra, dos Santos (paper in preparation)

# Predictions for the Run II

## ➤ Diffractive vector meson photoproduction in heavy ion collisions



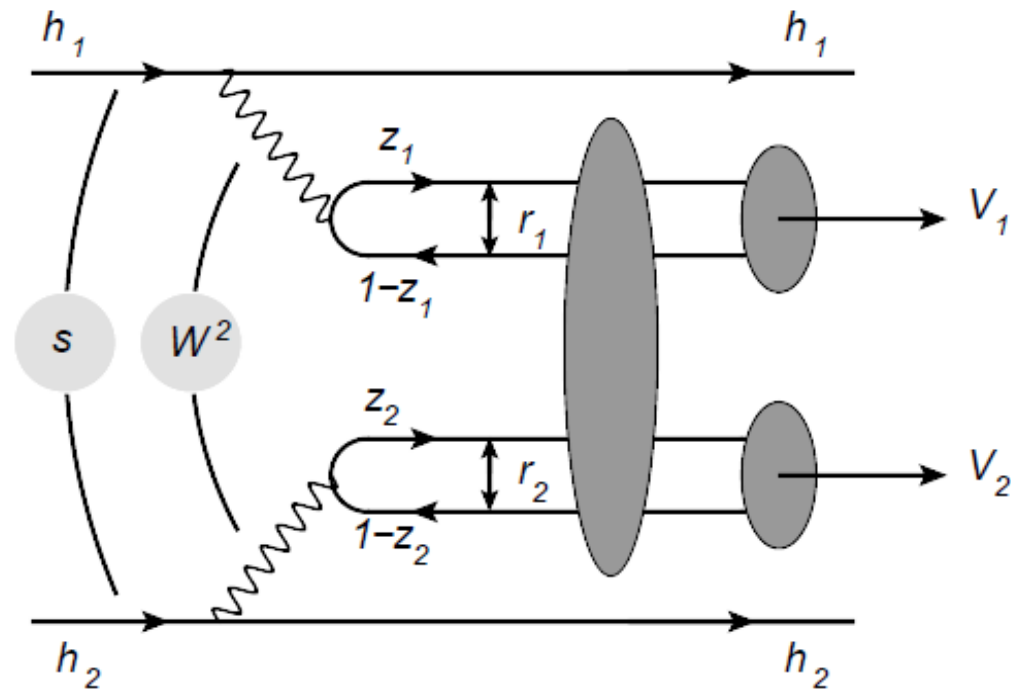
(\* ) VPG, Machado, Moreira, Navarra, dos Santos (paper in preparation)

# Predictions for the Run II

- Double Vector Meson production

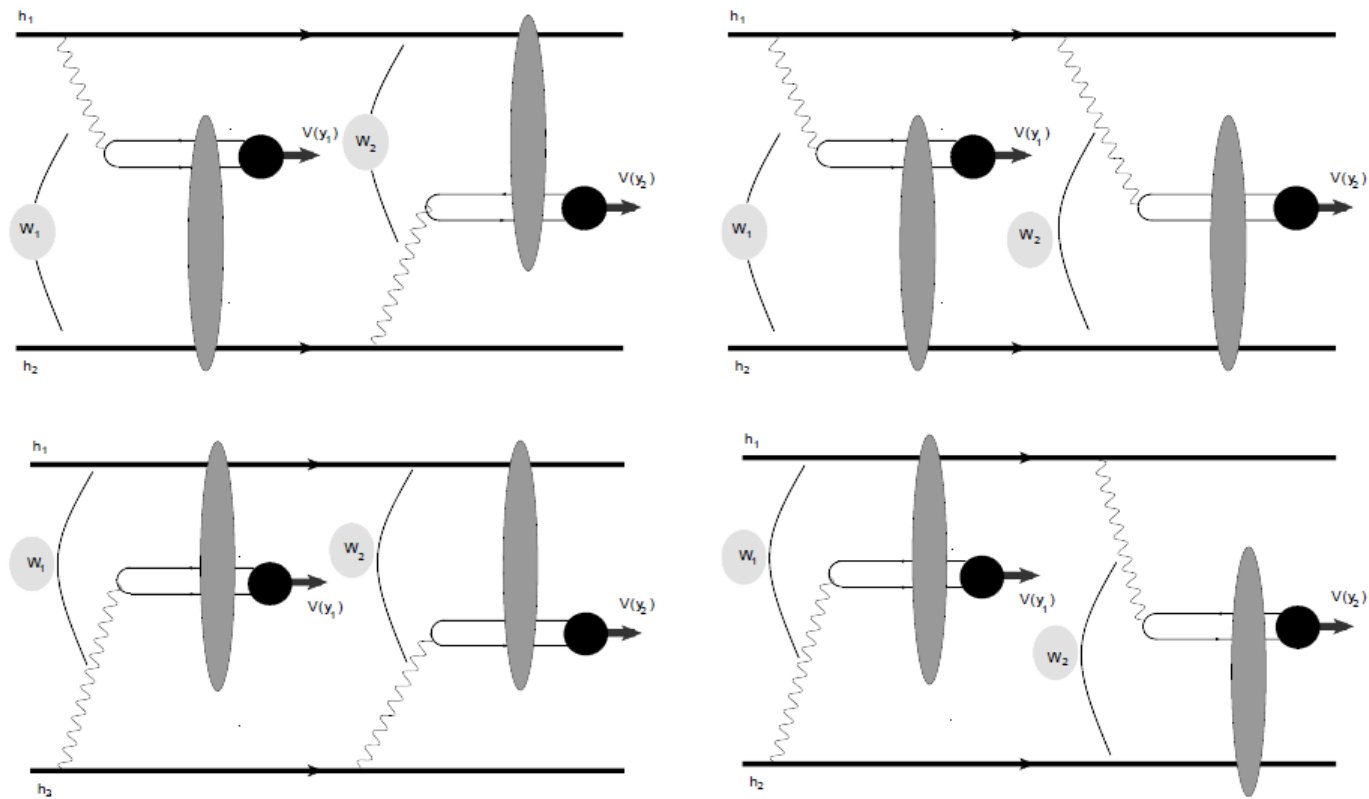
# Predictions for the Run II

## ➤ Double Vector Meson production



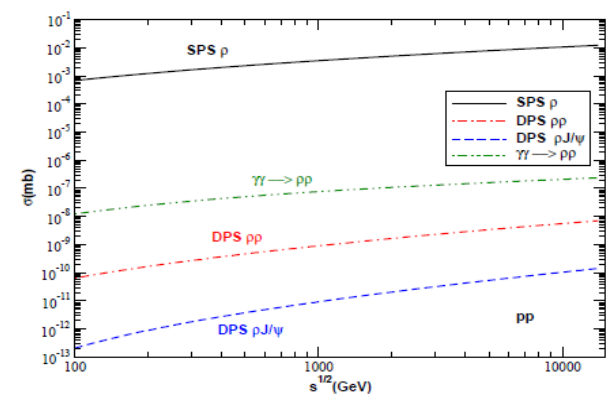
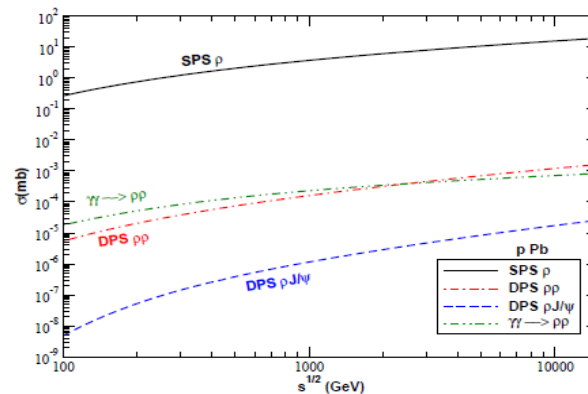
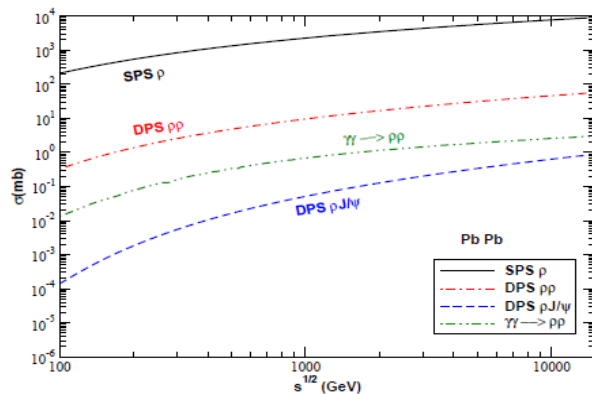
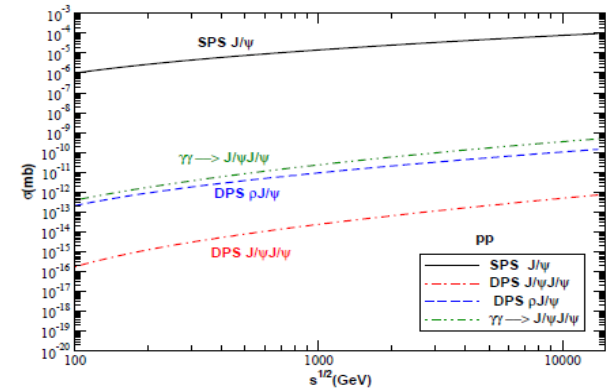
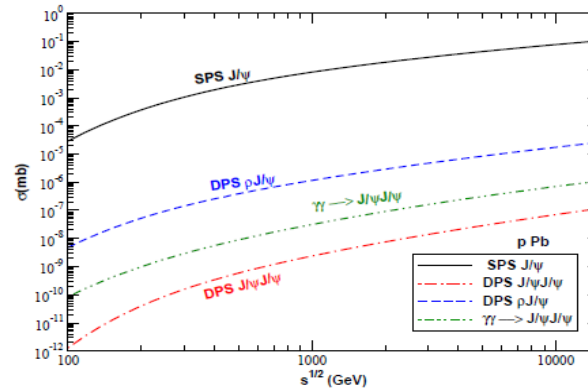
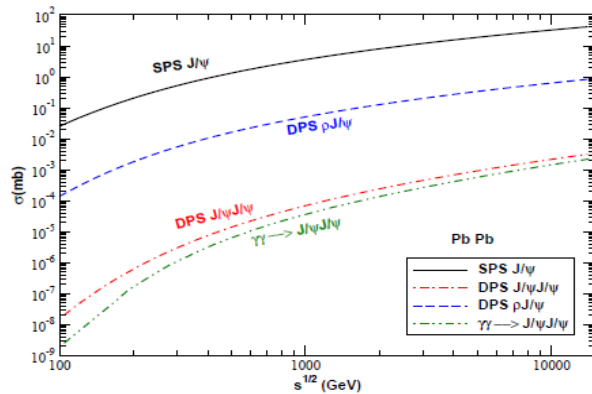
# Predictions for the Run II

## ➤ Double Vector Meson production



# Predictions for the Run II

## ➤ Double Vector Meson production



(\* ) VPG, Moreira, Navarra, arXiv:1605.05840[hep-ph]



# Double Vector Meson Production

## Predictions for the total cross sections

Final state	Mechanism	$PbPb$ $\sqrt{s} = 2.76$ TeV	$PbPb$ $\sqrt{s} = 5.5$ TeV	$pPb$ $\sqrt{s} = 5$ TeV	$pp$ $\sqrt{s} = 7$ TeV	$pp$ $\sqrt{s} = 14$ TeV
$J/\Psi J/\Psi$	DPS	402.301 nb	1054.951 nb	28.473 pb	$3.223 \times 10^{-4}$ pb	$7.256 \times 10^{-4}$ pb
	$\gamma\gamma$	235.565 nb	658.589 nb	310.194 pb	0.2412 pb	0.4793 pb
$\rho\rho$	DPS	21.150 mb	29.421 mb	702.595 nb	4.354 pb	7.083 pb
	$\gamma\gamma$	1.389 mb	1.973 mb	536.432 nb	182.442 pb	237.006 pb
$\rho J/\Psi$	DPS	0.18 mb	0.35 mb	8.929 nb	$7.469 \times 10^{-2}$ pb	$14.288 \times 10^{-2}$ pb

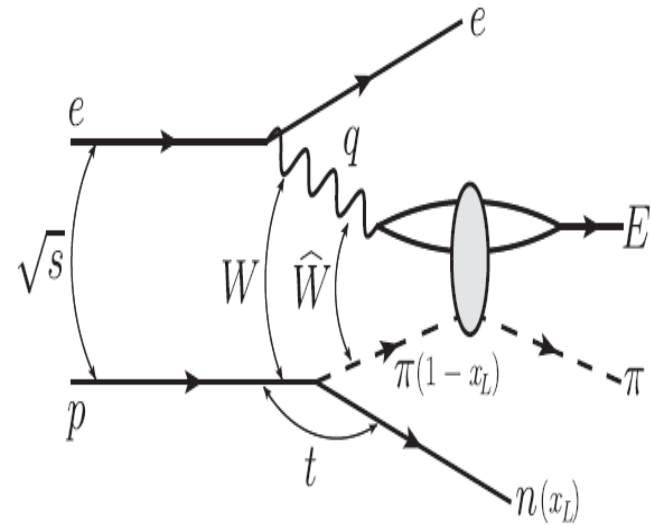
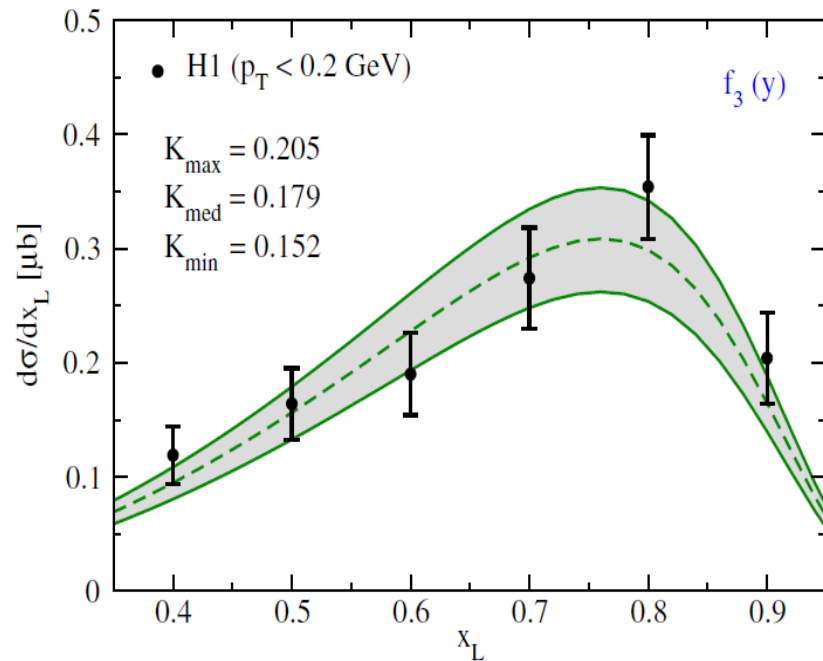
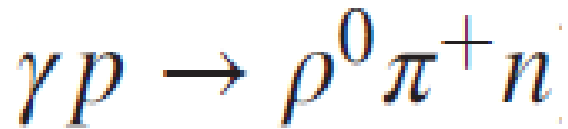
Final state		LHCb $2 < y_{1,2} < 4.5$	ATLAS/CMS $-2 < y_{1,2} < 2$	ALICE1 $-1 < y_{1,2} < 1$	ALICE2 $-1 < y_1 < 1$ and $-3.6 < y_2 < -2.6$
$J/\Psi J/\Psi$	$PbPb$ ( $\sqrt{s} = 2.76$ TeV)	5.51 nb	234.94 nb	69.91 nb	6.94 nb
	$PbPb$ ( $\sqrt{s} = 5.5$ TeV)	30.85 nb	446.11 nb	118.03 nb	25.45 nb
	$pPb$ ( $\sqrt{s} = 5$ TeV)	3.25 pb	8.87 pb	2.16 pb	0.37 pb
$\rho\rho$	$PbPb$ ( $\sqrt{s} = 2.76$ TeV)	0.93 mb	6.08 mb	1.58 mb	0.54 mb
	$PbPb$ ( $\sqrt{s} = 5.5$ TeV)	1.50 mb	7.06 mb	1.79 mb	0.73 mb
	$pPb$ ( $\sqrt{s} = 5$ TeV)	84.09 nb	122.03 nb	30.11 nb	8.53 nb
$\rho J/\Psi$	$PbPb$ ( $\sqrt{s} = 2.76$ TeV)	4.48 $\mu$ b	75.17 $\mu$ b	20.94 $\mu$ b	2.06 (7.25) $\mu$ b
	$PbPb$ ( $\sqrt{s} = 5.5$ TeV)	13.42 $\mu$ b	112.00 $\mu$ b	29.06 $\mu$ b	6.21 (11.86) $\mu$ b
	$pPb$ ( $\sqrt{s} = 5$ TeV)	1.02 nb	2.08 nb	0.51 nb	87.31 (144.56) pb

# Predictions for the Run II

- Vector Meson photoproduction with a leading neutron

# Predictions for the Run II

- Vector Meson photoproduction with a leading neutron at HERA



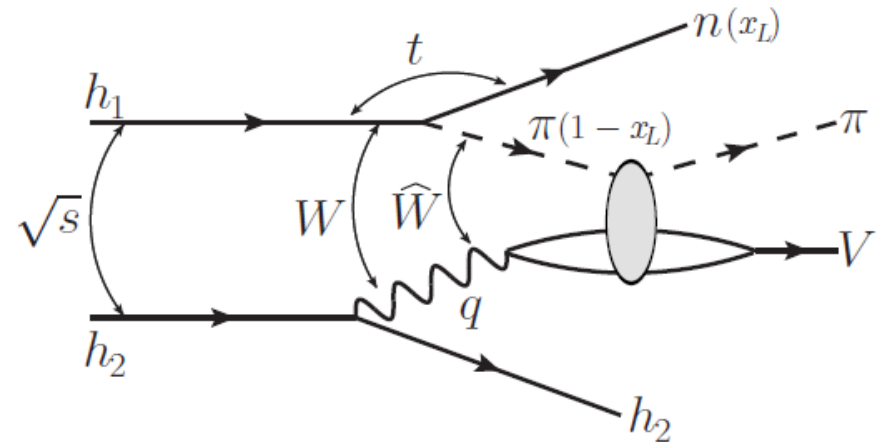
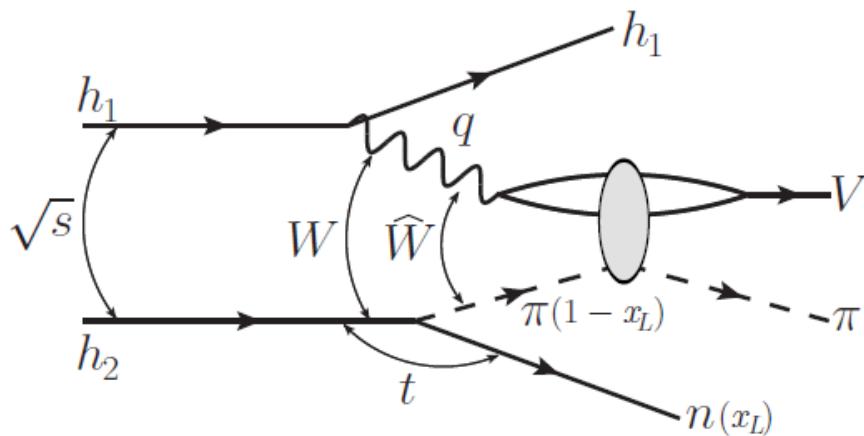
$$\sigma_{\gamma p \rightarrow V \otimes \pi + n}(W^2) = \mathcal{K} \cdot \int dx_L dt f_{\pi/p}(x_L, t) \cdot \sigma_{\gamma \pi \rightarrow V \otimes \pi}(\hat{W}^2)$$

VPG, Navarra, Spiering, PRD 93, 054025 (2016)

See talk by F. Navarra

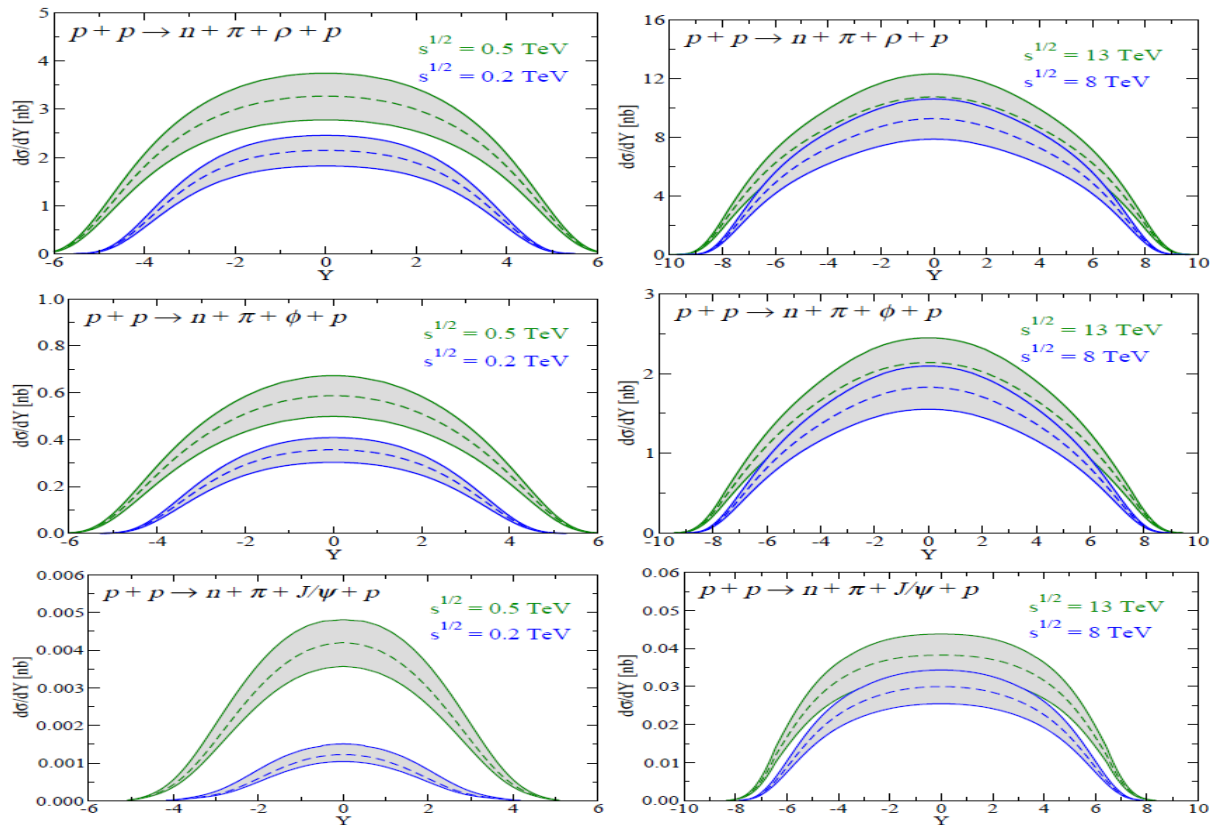
# Predictions for the Run II

- Vector Meson photoproduction with a leading neutron in UPHIC



# Predictions for the Run II

➤ Vector Meson photoproduction with a leading neutron in UPHIC



(\*\*) VPG, Moreira, Navarra, Spiering arXiv:1605.08186[hep-ph]

# Predictions for the Run II

➤ Vector Meson photoproduction with a leading neutron in UPHIC

$\sigma(V)$ [nb]		$\sqrt{s} = 0.2$ TeV	$\sqrt{s} = 0.5$ TeV	$\sqrt{s} = 8.0$ TeV	$\sqrt{s} = 13.0$ TeV
$\rho$	$K_{min}$	12.17	22.06	90.12	110.51
	$K_{med}$	14.34	25.98	106.12	130.14
	$K_{max}$	16.42	29.75	121.54	149.04
$\phi$	$K_{min}$	1.83	3.58	16.67	20.73
	$K_{med}$	2.15	4.21	19.63	24.42
	$K_{max}$	2.46	4.83	22.48	27.96
$J/\psi$	$K_{min}$	0.0042	0.019	0.25	0.35
	$K_{med}$	0.0049	0.022	0.30	0.42
	$K_{max}$	0.0064	0.026	0.34	0.48

(\*\*) VPG, Moreira, Navarra, Spiering arXiv:1605.08186[hep-ph]

# Summary

- ✓ The diffractive vector meson photoproduction in photon - induced interactions at the LHC is an important probe of the QCD dynamics at high energies.

# Summary

- ✓ The diffractive vector meson photoproduction in photon - induced interactions at the LHC is an important probe of the QCD dynamics at high energies.
- ✓ The Run I data can be successfully described by the color dipole formalism taking into account the nonlinear effects in the QCD dynamics.



# Summary

- ✓ The diffractive vector meson photoproduction in photon - induced interactions at the LHC is an important probe of the QCD dynamics at high energies.
- ✓ The Run I data can be successfully described by the color dipole formalism taking into account the nonlinear effects in the QCD dynamics.
- ✓ The Run II data can be used to constrain the description of the dipole - hadron scattering amplitude and the vector meson wave function.

# Summary

- ✓ The diffractive vector meson photoproduction in photon - induced interactions at the LHC is an important probe of the QCD dynamics at high energies.
- ✓ The Run I data can be successfully described by the color dipole formalism taking into account the nonlinear effects in the QCD dynamics.
- ✓ The Run II data can be used to constrain the description of the dipole - hadron scattering amplitude and the vector meson wave function
- ✓ Complementary studies can be performed by the analysis of the double vector meson production and the vector meson production associated to a leading neutron.

# Summary

- ✓ The diffractive vector meson photoproduction in photon - induced interactions at the LHC is an important probe of the QCD dynamics at high energies.
- ✓ The Run I data can be successfully described by the color dipole formalism taking into account the nonlinear effects in the QCD dynamics.
- ✓ The Run II data can be used to constrain the description of the dipole - hadron scattering amplitude and the vector meson wave function
- ✓ Complementary studies can be performed by the analysis of the double vector meson production and the vector meson production associated to a leading neutron.

**Thank you for your attention !**

# Extras

# Double Vector Meson Production Formalism

$$\frac{d^2\sigma_{h_1 h_2 \rightarrow h_1 V_1 V_2 h_2}}{dy_1 dy_2} = c \int_{b_{min}} \frac{d\sigma [h_1 + h_2 \rightarrow h_1 V_1 h_2]}{d^2 b dy_1} \times \frac{d\sigma [h_1 + h_2 \rightarrow h_1 V_2 h_2]}{d^2 b dy_2} d^2 b$$

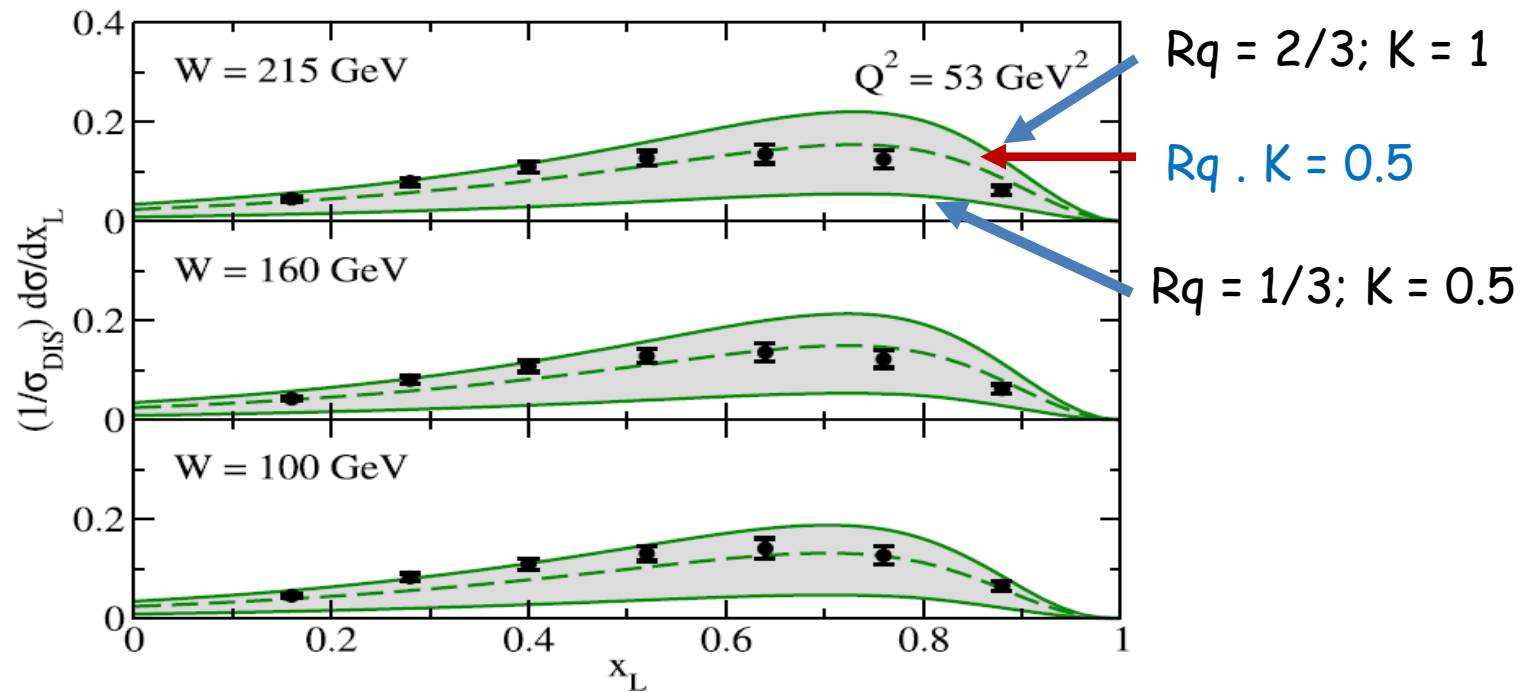
$$\frac{d\sigma [h_1 + h_2 \rightarrow h_1 \otimes V \otimes h_2]}{d^2 b dy} = [\omega N_{h_1}(\omega, b) \sigma_{\gamma h_2 \rightarrow V \otimes h_2}(\omega)]_{\omega_L} + [\omega N_{h_2}(\omega, b) \sigma_{\gamma h_1 \rightarrow V \otimes h_1}(\omega)]_{\omega_R}$$

$$N(\omega, b) = \frac{Z^2 \alpha_{em}}{\pi^2} \frac{1}{b^2 \omega} \cdot \left[ \int u^2 J_1(u) F \left( \sqrt{\frac{\left(\frac{b\omega}{\gamma_L}\right)^2 + u^2}{b^2}} \right) \frac{1}{\left(\frac{b\omega}{\gamma_L}\right)^2 + u^2} du \right]^2$$

# Leading neutron production in ep collisions

## Results for Inclusive Processes

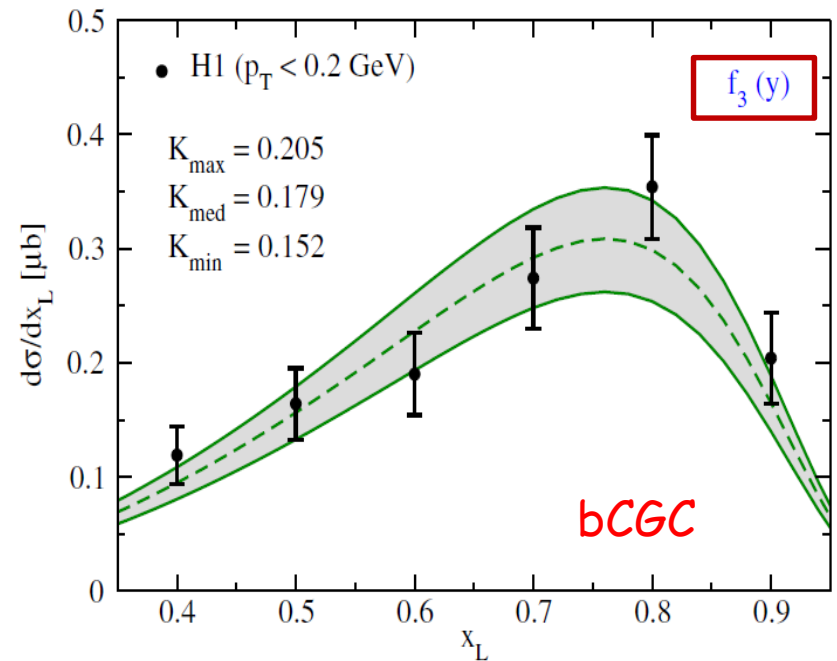
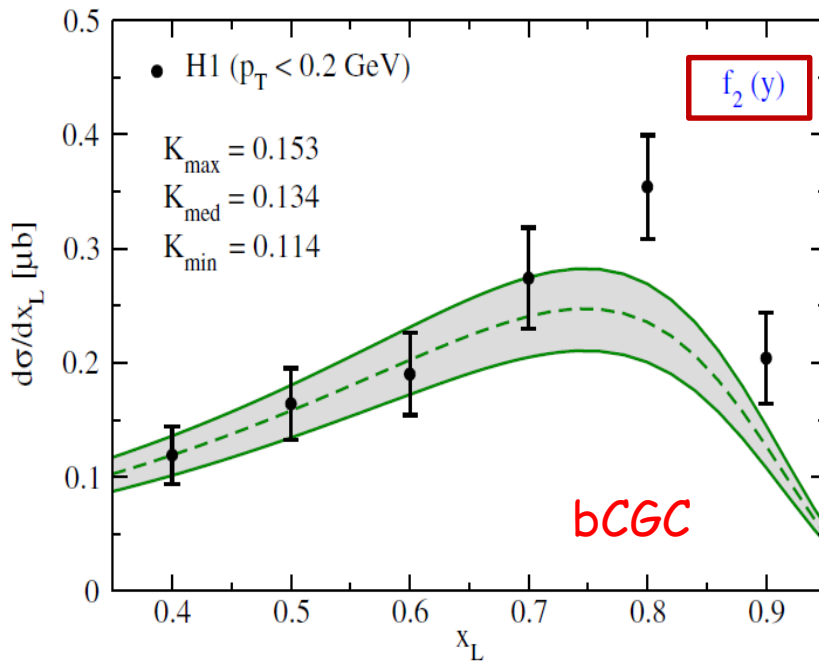
$$\frac{d^2\sigma(W, Q^2, x_L, t)}{dx_L dt} = f_{\pi/p}(x_L, t) \sigma_{\gamma^* \pi}(\hat{W}^2, Q^2)$$



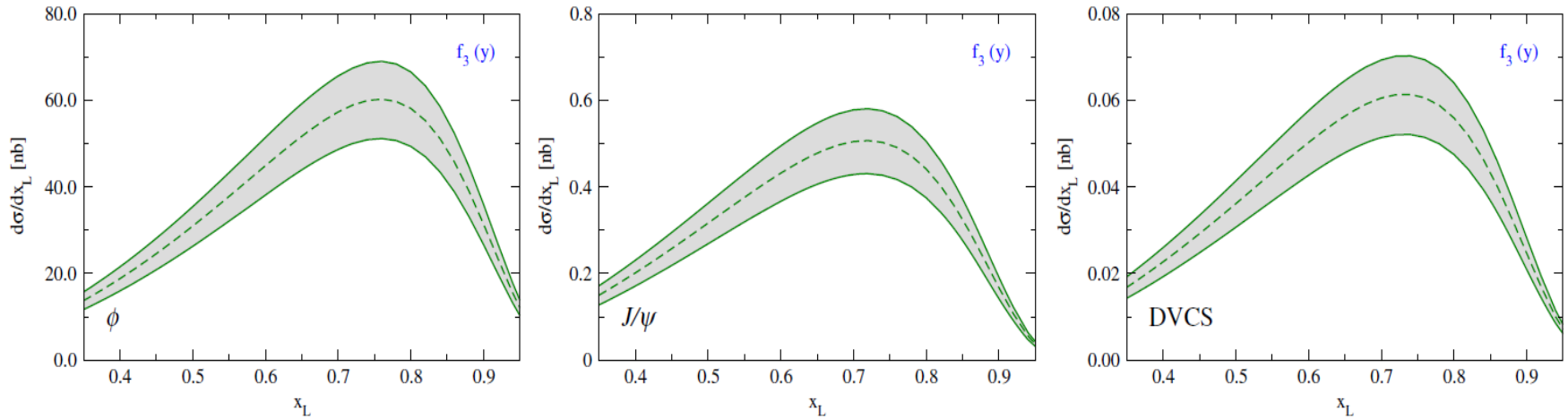
# Leading neutron production in ep collisions

## Results for Exclusive Processes

$$\gamma p \rightarrow \rho^0 \pi^+ n$$



# Predictions for Exclusive Processes with a leading neutron at HERA



$$Q^2 = 0.04 \text{ GeV}^2$$

$$\sigma(\gamma p \rightarrow \phi \pi n) = 25.47 \pm 3.70 \text{ nb}$$

$$\sigma(\gamma p \rightarrow J/\Psi \pi n) = 0.22 \pm 0.03 \text{ nb}$$

$$Q^2 = 10 \text{ GeV}^2$$

$$\sigma(\gamma^* p \rightarrow \gamma \pi n) = 0.008 \pm 0.001 \text{ nb}$$