

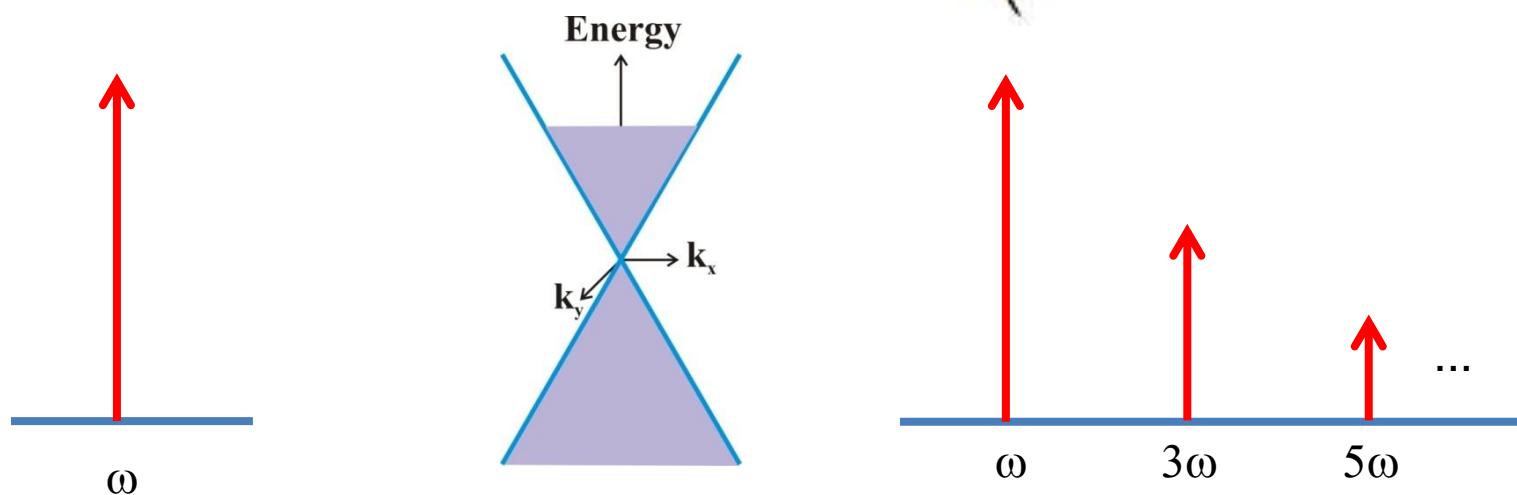
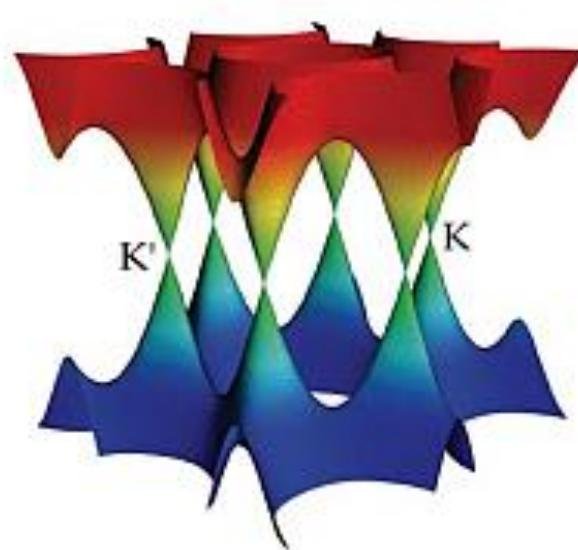
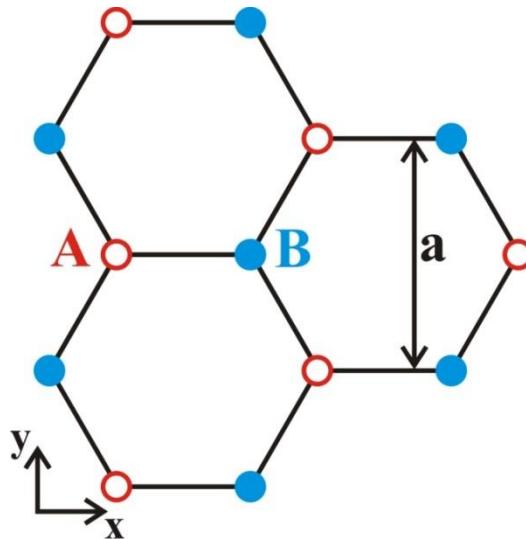
Terahertz Response of Monolayer Graphene: Velocity Gauge Vs Length Gauge

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Basics of Graphene



S. A. Mikhailov, Europhysics Lett. 79, 27002 (2007)

Motivation & Outline

- Find the best method to model the nonlinear response of Graphene

- Velocity Gauge:
 - F.H.M. Faisal, arXiv, 1, 0810.07881 (2008)
 - A.R. Wright, Appl. Phys. 44, 083001 (2011)

- Length Gauge:
 - C. Aversa and J. E. Sipe, Phys. Rev. B 52, 14636 (1995)
 - K. S. Virk and J. E. Sipe, Phys. Rev. B 76, 035213 (2007)

- Response of Graphene to THz field by employing these two gauges

Velocity & Length Gauges

- Hamiltonian in the present of electric field:

$$H = H_0 + H'$$

- Hamiltonian of the system in the absence of electric field:

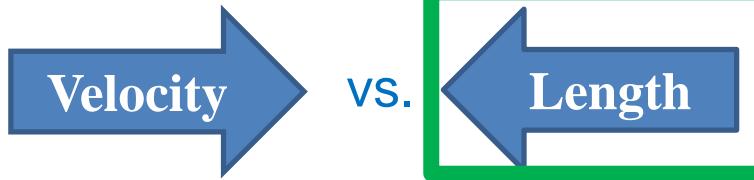
$$H_0 = \frac{p^2}{2m} + V(r)$$

- H' is the interaction Hamiltonian.

$$H' = -\frac{e}{m} p \cdot A(t) \quad \text{Velocity Gauge}$$

$$H' = -e \mathbf{r} \cdot \mathbf{E}(\mathbf{r}, t) \quad \text{Length gauge}$$

Gauge Choice



- Quantum mechanics is gauge invariant and both gauges should give, in principle, identical results
- Convergence of dynamic equations using the velocity gauge requires a large number of bands
- Divergences at zero frequency arise using the velocity gauge in the nonlinear response
- The divergence can be removed by developing sum rules, which become extremely complicated at high-order in the field

C. Aversa and J. E. Sipe, Phys. Rev. B 52, 14636 (1995)

K. S. Virk and J. E. Sipe, Phys. Rev. B 76, 035213 (2007)

Populations Dynamic Equations

Length Gauge:

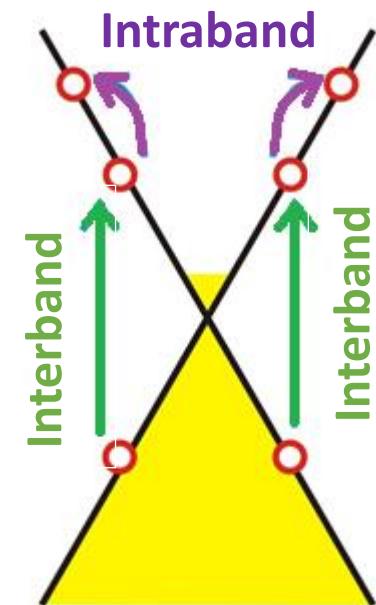
$$\frac{d\rho_{nn}(\mathbf{k})}{dt} = -\frac{i\sigma_n e \cdot \mathbf{E}(t)}{\hbar} [\xi_{cv}(\mathbf{k}) \rho_{vc}(\mathbf{k}) - \xi_{vc}(\mathbf{k}) \rho_{cv}(\mathbf{k})] - \frac{e(\mathbf{E}(t) \cdot \nabla_{\mathbf{k}}) \rho_{nn}(\mathbf{k})}{\hbar} - \frac{[\rho_{nn}(\mathbf{k}) - f_n(\mathbf{k}, t)]}{\tau_n}$$

Velocity Gauge:

$$\frac{d\rho_{nn}(\mathbf{k})}{dt} = \frac{e}{\hbar} \sigma_n \omega_{cv}(\mathbf{k}) (\xi_{cv}(\mathbf{k}) \rho_{vc}(\mathbf{k}) + \rho_{cv}(\mathbf{k}) \xi_{vc}(\mathbf{k})) \cdot \mathbf{A}(t) - \frac{\rho_{nn}(\mathbf{k}) - f_n(\mathbf{k}, t)}{\tau_n}$$

$$\xi_{nm}(\mathbf{k}) = [2\delta_{n,m} - 1] \frac{\hat{\theta}}{2k}$$

$$\mathbf{A} = -\frac{i\mathbf{E}}{\omega}$$



I. Al-Naib et al., Phys. Rev. B 90, 245423, 2014

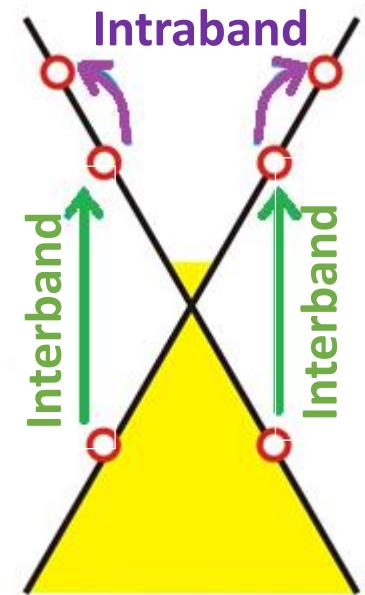
Coherences Dynamic Equations

Length Gauge:

$$\frac{d\rho_{cv}(\mathbf{k})}{dt} = \frac{ie\mathbf{E}(t)\xi_{cv}(\mathbf{k})}{\hbar} [\rho_{vv}(\mathbf{k}) - \rho_{cc}(\mathbf{k})] - i\omega_{cv}(\mathbf{k})\rho_{cv}(\mathbf{k}) - \frac{e\mathbf{E}(t) \cdot \nabla_{\mathbf{k}}\rho_{cv}(\mathbf{k})}{\hbar} - \frac{\rho_{cv}(\mathbf{k})}{\tau}$$

Velocity Gauge:

$$\frac{d\rho_{cv}(\mathbf{k})}{dt} = -\frac{e\omega_{cv}(\mathbf{k})\mathbf{A}(t) \cdot \xi_{cv}(\mathbf{k})}{\hbar} (\rho_{vv}(\mathbf{k}) - \rho_{cv}(\mathbf{k})) - i\omega_{cv}(\mathbf{k})\rho_{cv}(\mathbf{k}) + \frac{ie\mathbf{A}(t)}{\hbar} \cdot \nabla_{\mathbf{k}}\omega_{cv}(\mathbf{k})\rho_{cv}(\mathbf{k}) - \frac{\rho_{cv}(\mathbf{k})}{\tau}$$



I. Al-Naib et al., Phys. Rev. B 90, 245423, 2014

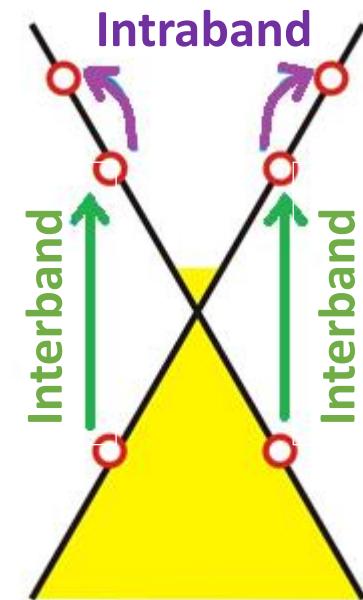
Current Density in Length Gauge

Interband Current:

$$\mathbf{J}_e(t) = \frac{8|e|}{A} \operatorname{Re} \left\{ \sum_{\mathbf{k}} \frac{\hat{\boldsymbol{\theta}}}{2k} \frac{d\rho_{cv}(\mathbf{k}, t)}{dt} \right\}$$

Intraband Current:

$$\begin{aligned} \mathbf{J}_i(t) &= \frac{-4|e|v_F}{A} \sum_{\mathbf{k}} \{ \rho_{cc}(\mathbf{k}, t) - \rho_{vv}(\mathbf{k}, t) \} \hat{\mathbf{k}} \\ &+ \frac{8|e|^2}{A\hbar} \sum_{\mathbf{k}} \operatorname{Re} \{ \rho_{cv}(\mathbf{k}, t) \nabla_{\mathbf{k}} [\mathbf{E}(t) \cdot \boldsymbol{\xi}_{vc}(\mathbf{k})] \} \end{aligned}$$



Current Density in Velocity Gauge

Interband Current:

$$\begin{aligned} J_e(t) = & \frac{4ie}{A} \sum_{\mathbf{k}} (\xi_{cv}(\mathbf{k})\rho_{vc}(\mathbf{k}) - \xi_{vc}(\mathbf{k})\rho_{cv}(\mathbf{k}))\omega_{cv}(\mathbf{k}) \\ & - \frac{4e^2}{\hbar} \sum_{\mathbf{k}} \omega_{cv}(\mathbf{k}) (\xi_{cv}^i(\mathbf{k})\xi_{vc}^j(\mathbf{k}) + \xi_{cv}^j(\mathbf{k})\xi_{vc}^i(\mathbf{k})) \mathbf{A}(t) (\rho_{vv}(\mathbf{k}, t) - \rho_{cc}(\mathbf{k}, t)) \\ & + \frac{4e^2}{\hbar} \sum_{\mathbf{k}} (\rho_{vc}(\mathbf{k}, t)\xi_{cv}(\mathbf{k}) - \rho_{cv}(\mathbf{k}, t)\xi_{vc}(\mathbf{k})) \nabla_{\mathbf{k}} \omega_{cv}(\mathbf{k}) \cdot \mathbf{A}(t) \end{aligned}$$

Intraband Current:

$$\begin{aligned} J_i(t) = & -\frac{4ev_f}{A} \sum_{\mathbf{k}} \{\rho_{cc}(\mathbf{k}, t) - \rho_{vv}(\mathbf{k}, t)\} \hat{\mathbf{k}} \\ & + \frac{8e^2}{\hbar A} \sum_{\mathbf{k}} \text{Im} \{ \rho_{cv}(\mathbf{k}, t) \nabla_{\mathbf{k}} [\omega_{cv}(\mathbf{k}) \xi_{vc}(\mathbf{k})] \cdot \mathbf{A}(t) \} \\ & - \frac{e^2}{A} \sum_{\mathbf{k}} \left(\frac{1}{m} \{\rho_{vv}(\mathbf{k}, t) + \rho_{cc}(\mathbf{k}, t)\} \mathbf{A}(t) \right) \end{aligned}$$

$$\mathbf{A} = -\frac{iE}{\omega}$$

Current Density

Effective Mass Sum Rule:

$$\frac{1}{m} = \frac{1}{\hbar^2} \frac{\partial^2 E_n(\mathbf{k})}{\partial k^x \partial k^j} - \sum_{m \neq n} \frac{\omega_{nm}^2(\mathbf{k})(\xi_{nm}^x(\mathbf{k})\xi_{mn}^j(\mathbf{k}) + \xi_{nm}^j(\mathbf{k})\xi_{mn}^x(\mathbf{k}))}{m^2(E_m(\mathbf{k}) - E_n(\mathbf{k}))}$$

Total Current Density:

$$\begin{aligned} J(t) &= \frac{4ie}{A} \sum_{\mathbf{k}} (\xi_{cv}(\mathbf{k})\rho_{vc}(\mathbf{k}) - \xi_{vc}(\mathbf{k})\rho_{cv}(\mathbf{k}))\omega_{cv}(\mathbf{k}) \\ &+ \frac{4e^2}{A\hbar} \sum_{\mathbf{k}} (\rho_{vc}(\mathbf{k}, t)\xi_{cv}(\mathbf{k}) - \rho_{cv}(\mathbf{k}, t)\xi_{vc}(\mathbf{k}))\nabla_{\mathbf{k}}\omega_{cv}(\mathbf{k}) \cdot \mathbf{A}(t) \\ &- \frac{4ev_f}{A} \sum_{\mathbf{k}} \{\rho_{cc}(\mathbf{k}, t) - \rho_{vv}(\mathbf{k}, t)\}\hat{\mathbf{k}} \\ &+ \frac{8e^2}{A\hbar} \sum_{\mathbf{k}} \text{Im} \{ \rho_{cv}(\mathbf{k}, t) \nabla_{\mathbf{k}} [\omega_{cv}(\mathbf{k})\xi_{vc}(\mathbf{k})] \cdot \mathbf{A}(t) \} \\ &- \frac{4e^2}{A\hbar} \sum_{\mathbf{k}} \left(\frac{\partial^2 \omega_v(\mathbf{k})}{\partial \mathbf{k}^2} \rho_{vv}(\mathbf{k}, t) + \frac{\partial^2 \omega_c(\mathbf{k})}{\partial \mathbf{k}^2} \rho_{cc}(\mathbf{k}, t) \right) \mathbf{A}(t) \end{aligned}$$

J. Callaway, Academic Press, Second edition, (1974)

First order Intraband Conductivity

➤ Length Gauge:

$$\sigma_i^{(1)}(\omega) = \frac{2ie^2k_B T}{\pi\hbar^2(\omega + \frac{i}{\tau})} \ln(2\cosh(\frac{\beta\mu}{2}))$$

➤ Velocity Gauge:

$$\sigma_i^{(1)}(\omega) = \frac{2ie^2k_B T}{\pi\hbar^2\omega} \ln(2\cosh(\frac{\beta\mu}{2}))$$

First Order Interband Conductivity

➤ Length Gauge:

$$\sigma_e^{(1)} = -\frac{ie^2}{4\pi\hbar} \int \frac{\omega}{\omega_{cv}} \frac{d\omega_{cv}}{(\omega_{cv} - \omega - \frac{i}{\tau}) \cosh \beta \mu + \cosh \beta (\frac{\hbar \omega_{cv}}{2})}$$

➤ Velocity Gauge:

$$\sigma_e^{(1)} = -\frac{ie^2}{4\pi\hbar} \int \frac{\omega_{cv}}{\omega} \frac{d\omega_{cv}}{(\omega_{cv} - \omega - \frac{i}{\tau}) \cosh \beta \mu + \cosh \beta (\frac{\hbar \omega_{cv}}{2})}$$

➤ $\tau \rightarrow \infty$

$$\sigma_e^{(1)} = \frac{e^2}{4\hbar}$$

Summary

- If one uses the mass sum rule and neglects scattering, two gauges yield identical linear conductivities.
- If we consider scattering term, even the linear interband conductivity is quite different in two gauges.
- One should do calculations in the interaction picture when employing velocity gauge.
- Nonlinear response can be quite different for the two approaches, due to the diverges arise at zero frequency in the velocity gauge when one uses a basis with a finite number of bands.
- One should use the length gauge for graphene when calculating the nonlinear THz response.

Thank You!

Current

- $\mathbf{J}(t) = e \int \frac{d\vec{k}}{(2\pi)^2} V_{nm}(k) \rho_{nm}(k, t)$
- $V_{nm}(k) = \mathbf{v}_{nm}(k) - e \frac{1}{m} \delta_{nm} \mathbf{A}(t)$
- $\langle n\mathbf{k} | \mathbf{r} | m\mathbf{k}' \rangle = \delta(\mathbf{k} - \mathbf{k}') \xi_{nm}(\mathbf{k}) + i \delta_{nm} \nabla_{\mathbf{k}} \delta(\mathbf{k} - \mathbf{k}')$

Last term in the intraband current:

$$\frac{4ie^2}{A} \sum_{\mathbf{k}} \frac{1}{m\omega} \{ \rho_{cc}(\mathbf{k}, t) + \rho_{vv}(\mathbf{k}, t) \} \mathbf{E}(t)$$

$$\mathbf{E} = \mathbf{e}_x E_0 e^{-i\omega t} + \mathbf{e}_y E_0 e^{-i\omega t}$$

Mass Sum Rule

- $\frac{1}{m} = \frac{1}{\hbar} \frac{\partial^2 E_n(k)}{\partial k^x \partial k^j}$
+ $\sum_{m \neq n} \frac{p_{nm}^x(k) p_{mn}^j(k) + p_{nm}^j(k) p_{mn}^x(k)}{m^2(E_m(k) - E_n(k))}$

$$\langle n, k | p_e | m, k' \rangle = i m \omega_{nm}(k) \langle n, k | r_e | m, k' \rangle$$

$$\langle n, k | r_e | m, k' \rangle = (1 - \delta_{nm}) \delta(k - k') \xi_{nm}(k, k')$$

$$\langle n\mathbf{k} | \mathbf{r} | m\mathbf{k}' \rangle = \delta(\mathbf{k} - \mathbf{k}') \boldsymbol{\xi}_{nm}(\mathbf{k}) + i \delta_{nm} \nabla_{\mathbf{k}} \delta(\mathbf{k} - \mathbf{k}')$$

Linear Response

Velocity Gauge

$$E_c(k) = \hbar v_f k \quad \& \quad E_v = -\hbar v_f k$$
$$J_i^{(1)} = -\frac{4e^2 v_f}{\hbar A} \sum_k \frac{1}{k} \left(1 - \frac{k_x^2}{k^2}\right) \left(1 - \rho_{hh}^{(0)}(\mathbf{k}) + \rho_{cc}^{(0)}(\mathbf{k})\right) A(t)$$
$$J_e^{(1)} = -\frac{8e}{A} \sum_k \omega_{cv}(\mathbf{k}, t) \text{Im}\{\rho_{cv}^{(1)}(\mathbf{k}, t) \xi_{vc}(\mathbf{k}, t)\}$$

Length Gauge

$$J_i^{(1)} = 4e v_f \sum_k \{\rho_{cc}^{(1)}(\mathbf{k}) + \rho_{hh}^{(1)}(\mathbf{k})\} \hat{k}$$
$$J_e^{(1)} = 8e Re \sum_k \{\xi_{cv}(\mathbf{k}) \dot{\rho}_{cc}^{(1)}(\mathbf{k})\}$$

First order Intraband conductivity (ω)

➤ Velocity Gauge:

$$\sigma_i^{(1)}(\omega) = \frac{2ie^2k_B T}{\pi\hbar^2\omega} \ln\left(2\cosh\left(\frac{\beta\mu}{2}\right)\right)$$

➤ Length Gauge

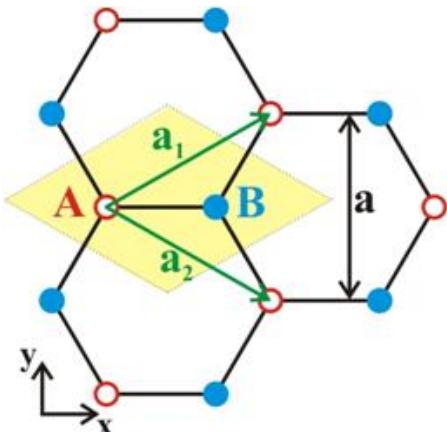
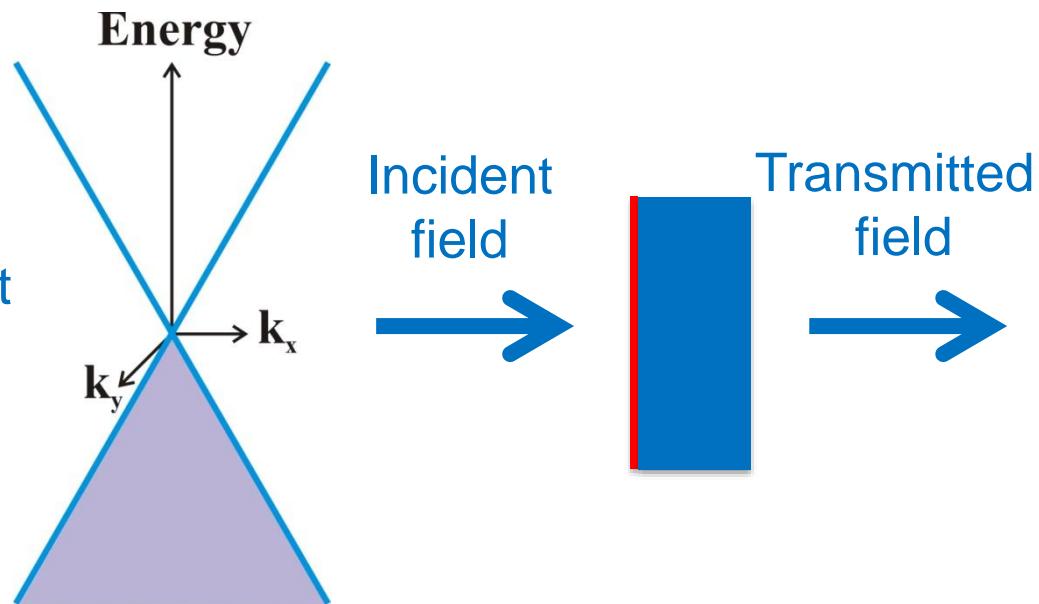
$$\sigma_i^{(1)}(\omega) = \frac{2ie^2k_B T}{\pi\hbar^2\left(\omega + \frac{i}{\tau}\right)} \ln\left(2\cosh\left(\frac{\beta\mu}{2}\right)\right)$$

Outline

- Motivation
- Velocity gauge and Length Gauge
- Comparing response of Graphene in Velocity gauge and Length gauge
- Conclusion

System Under Study

- Graphene on a substrate
- Fermi level at the Dirac point
- Temperature of 30 K
- Scattering time of 50 fs



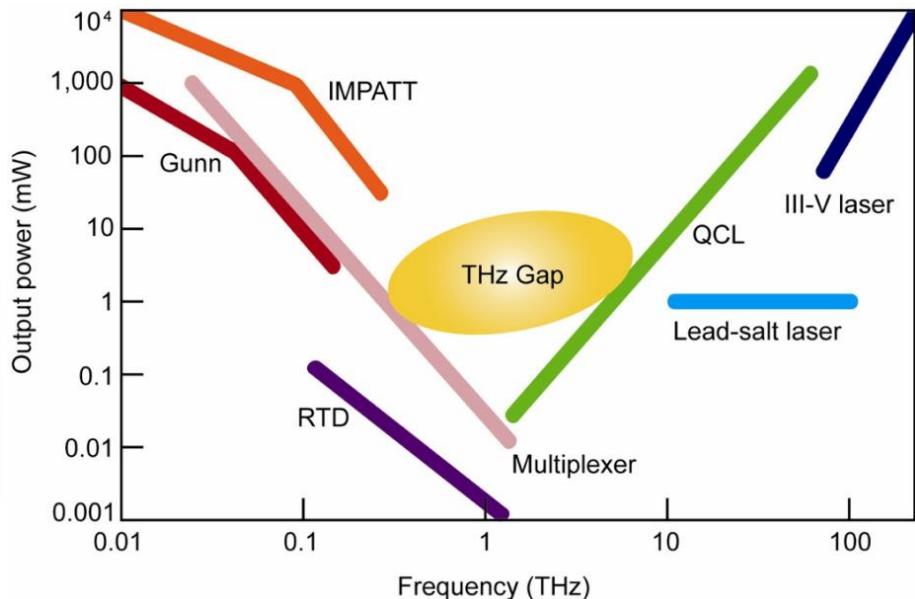
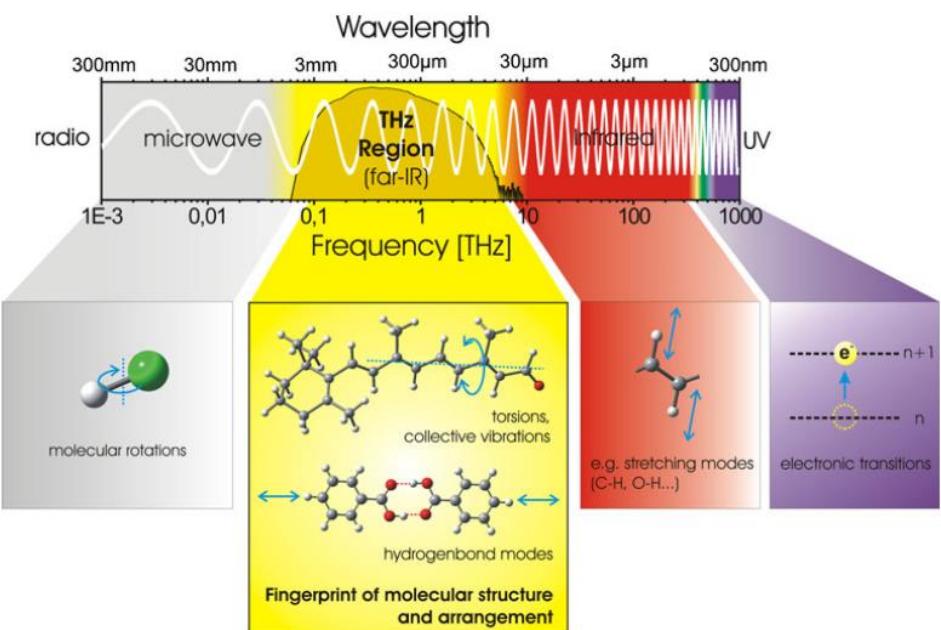
$$\phi_{n\mathbf{k}}(\mathbf{r}) = N \sum_{\mathbf{R}} [\varphi_{pz}(\delta\mathbf{r}_A) + \sigma_n e^{-i\chi(\mathbf{k})} \varphi_{pz}(\delta\mathbf{r}_B)] e^{i\mathbf{k}\cdot\mathbf{R}}$$

$$\sigma_v = 1 \text{ and } \sigma_c = -1$$

$$\chi(\mathbf{k}) = \arg [f(\mathbf{k})]$$

$$f(\mathbf{k}) \equiv 1 + e^{-i\mathbf{k}\cdot\mathbf{a}_1} + e^{-i\mathbf{k}\cdot\mathbf{a}_2}$$

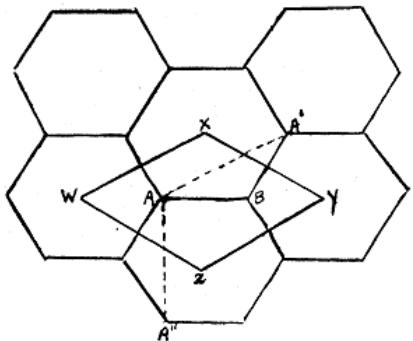
Introduction: THz Band



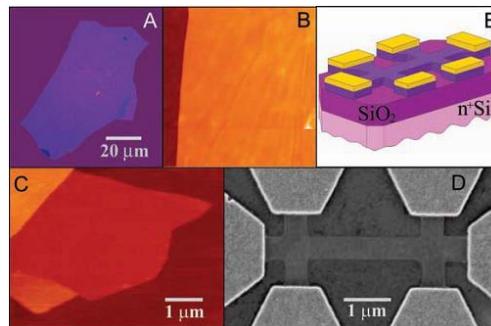
Moloney *et al.*, DOI: 10.1117/2.1201102.003523

- Numerous organic molecules exhibit strong absorption and dispersion due to rotational and vibrational transitions

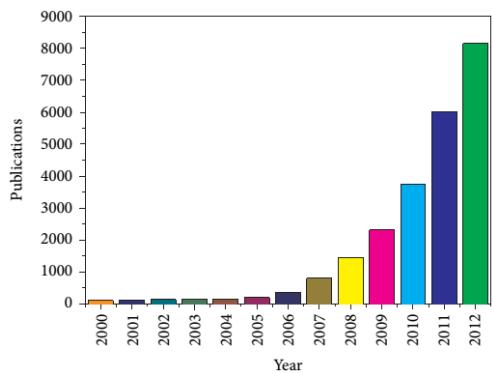
Introduction: Graphene



P. R. Wallace, Physical Review 71, 622 (1947)

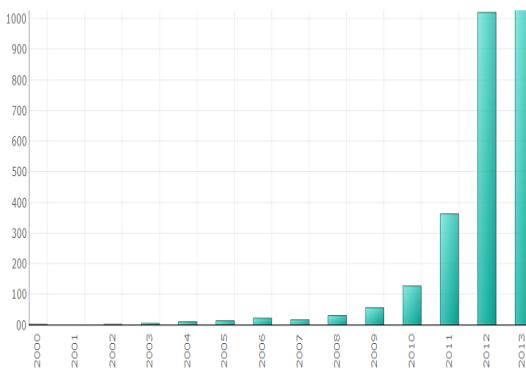


K. S. Novoselov *et al.*, Science 306, 666 (2004)



Publication trend

V. Dhand *et al.*, J. of Nanomaterials 763953 (2013)



Patents trend

Technology Insight Report: Graphene (2013)

First order Intraband Conductivity ($T \rightarrow 0$)

Velocity Gauge in the Schrodinger Picture:

$$\sigma_i^{(1)} = \frac{ie^2 E_f}{\pi \hbar^2 \omega}$$

Velocity Gauge in the Interaction Picture:

$$\sigma_i^{(1)} = \frac{ie^2 E_f}{\pi \hbar^2 (\omega + \frac{i}{\tau_c})}$$

Length gauge in the Schrodinger picture:

$$\sigma_i^{(1)} = \frac{ie^2 E_f}{\pi \hbar^2 (\omega + \frac{i}{\tau_c})}$$