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Estimation of initial temperature of high multiplicity pp event at 7 TeV in the light of Color String Percolation Model (CSPM)

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Outline

- Introduction
- Physics Motivation: Temperature estimation using Color String Percolation Model (CSPM)
- Analysis details
- Results
 - >> Determination of color suppression factor $F(\xi)$
- Summary and outlook

Introduction

- The main goal of relativistic heavy ion collision is to study the de-confined matter, called Quark Gluon Plasma (QGP)
- Percolation theory can be used to describe the transition from hadronic to QGP state[1].
- Several objects can form a cluster of communication.
- At a certain density of the object a spanning cluster appears, which marks percolation phase transition.
- This is defined by percolation density parameter.

[1] T. Celik, F. Karsch, H. Satz, Phys. Lett. B (1980) 128

Color string percolation model (CSPM)

- The CSPM describes the initial collision of two nucleon in terms of color string stretched between the projectile and the target.
- In the transverse plane, the color string may be viewed as a small disk filled with the color field created by the colliding partons (Fig.1).

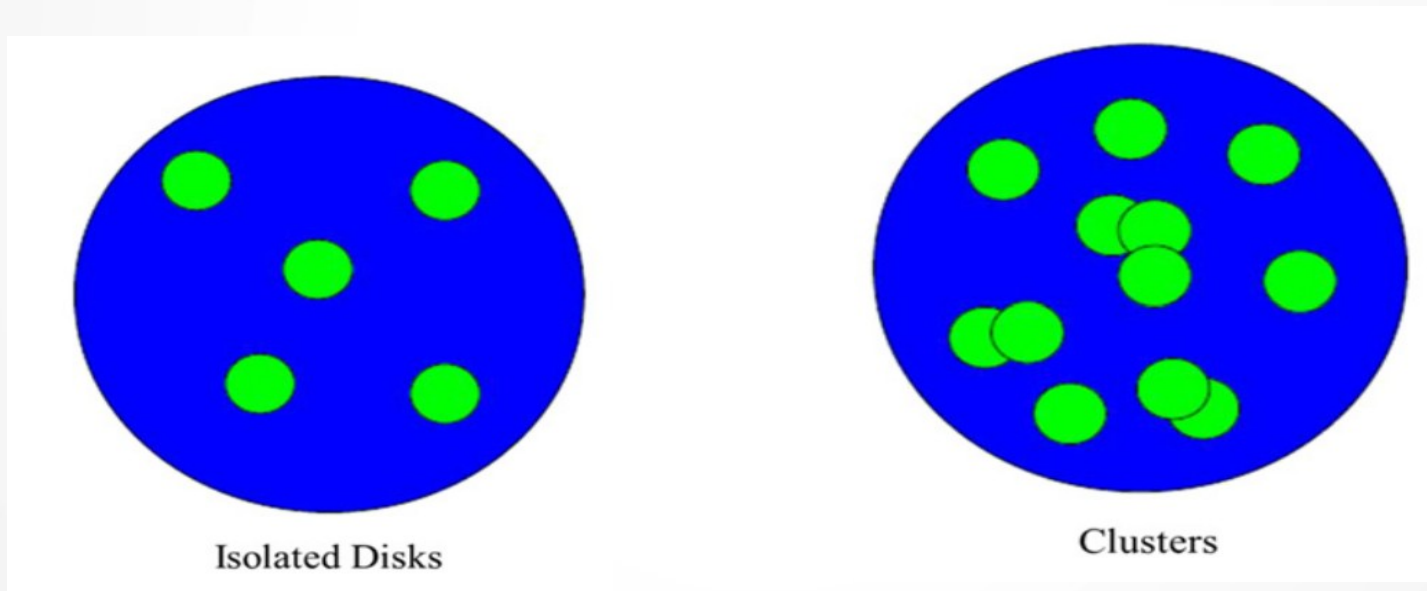
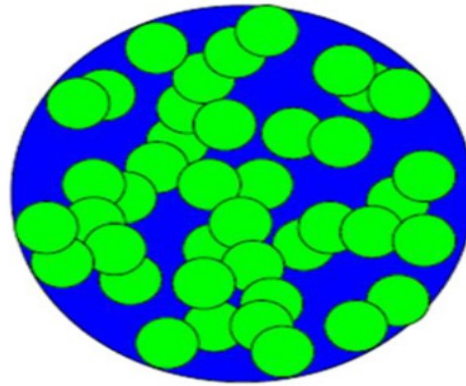


Fig.1

Fig.2

- With the growing energy and size, the number of string grows and start to overlap to form clusters (Fig.2).

- At a certain critical density, a macroscopic cluster appears that marks the percolation phase transition (Fig.3).



Percolation

Fig.3

- The onset of percolation is given by the percolation density (ξ) parameter,

$$\xi = \frac{NS_1}{S}$$

Where N = number of strings formed in the collision,
 S_1 = transverse area of a single string,
 S = transverse nuclear overlap area.

➤ Cluster of n strings behaves as a single string with energy momentum that corresponds to the sum of the energy momentum of the overlapping strings.

➤ The color charge (Q_n) of such strings, which cover the same area S_n as the area of single string S_1 is

$$Q_n^2 = nQ_1^2 \Rightarrow Q_n = \sqrt{n}Q_1$$

which reduces the color charge by factor a of \sqrt{n}

➤ In general, $S_n > S_1$ hence the reduction factor becomes

$$\sqrt{\frac{nS_n}{S_1}}$$

➤ In the framework of SM for color field, the multiplicity (μ_n) is proportional to color charge (Q_n) hence,

$$\mu_n = \sqrt{\frac{nS_n}{S_1}} \mu_1$$

➤ The average transverse momentum squared ($\langle p_T^2 \rangle$) is proportional to the string tension (σ_n). Again,

$$Q_n \sim S_n \sigma_n$$

➤ Thus the transverse momentum squared ($\langle p_T^2 \rangle$) produced by cluster of n strings having single string transverse momentum squared $\langle p_T^2 \rangle_1$ is,

$$\langle p_T^2 \rangle = \sqrt{\frac{nS_1}{S_n}} \langle p_T^2 \rangle_1$$

► At high string density,

$$\left\langle \frac{nS_1}{S_n} \right\rangle = \frac{1}{F(\xi)^2}$$

Where $F(\xi)$, is the color suppression factor (due to overlapping of the strings) and is given by,

$$F(\xi) = \sqrt{\frac{1 - e^{-\xi}}{\xi}}$$

► The net effect of $F(\xi)$ is the reduction of multiplicity (μ_n) and increase in the transverse momentum squared ($\langle p_T^2 \rangle$).

$$\mu_n = nF(\xi)\mu_1 \quad \text{and} \quad \langle p_T^2 \rangle = \frac{\langle p_T^2 \rangle_1}{F(\xi)}$$

Determination of temperature

- ▶ The Schwinger mechanism for massless particles can be expressed in terms of p_T^2

$$\frac{dN}{dp_T^2} \sim e^{-\pi p_T^2 / x^2}$$

Where $\langle x^2 \rangle$ is the average string tension.

- ▶ Gaussian fluctuation of the string tension for the cluster gives rise to the thermal distribution.

$$\frac{dN}{dp_T^2} \sim e^{-p_T \sqrt{\frac{2\pi}{\langle x^2 \rangle}}}$$

with $\langle x^2 \rangle = \pi \langle p_T^2 \rangle / F(\xi)$

► The temperature can be expressed as

$$T(\xi) = \sqrt{\frac{\langle p_T^2 \rangle_1}{2F(\xi)}}$$

where $\langle p_T^2 \rangle_1$ is calculated using above equation at $\xi=1.2$ with universal freeze out temperature $T_f=167.7 \pm 2.6$ MeV which gives

$$\sqrt{\langle p_T^2 \rangle_1} = 207.2 \pm 3.3 \text{ MeV}$$

Analysis Details :

- $0.25 < p_T < 1.3$ (GeV/c)
- $|\eta| < 0.5$
- Detector used : TPC
- Z Vertex cut : 10. cm
- Range used to fit $\langle p_T^2 \rangle$ distribution = $0.145 - 1.5$ (GeV/c)²
- ALICE data: p-p at $\sqrt{S_{NN}} = 7$ TeV

LHC10d, pass2, ESD data

Run Number : 125085, 126007, 126082, 126097

- MC data : p-p at $\sqrt{S_{NN}} = 7$ TeV

LHC10f6a, ESD data

Run Number : 126097, 126007, 126082, 126088

- Multiplicity cut : V0A multiplicity cut

ALICE detector

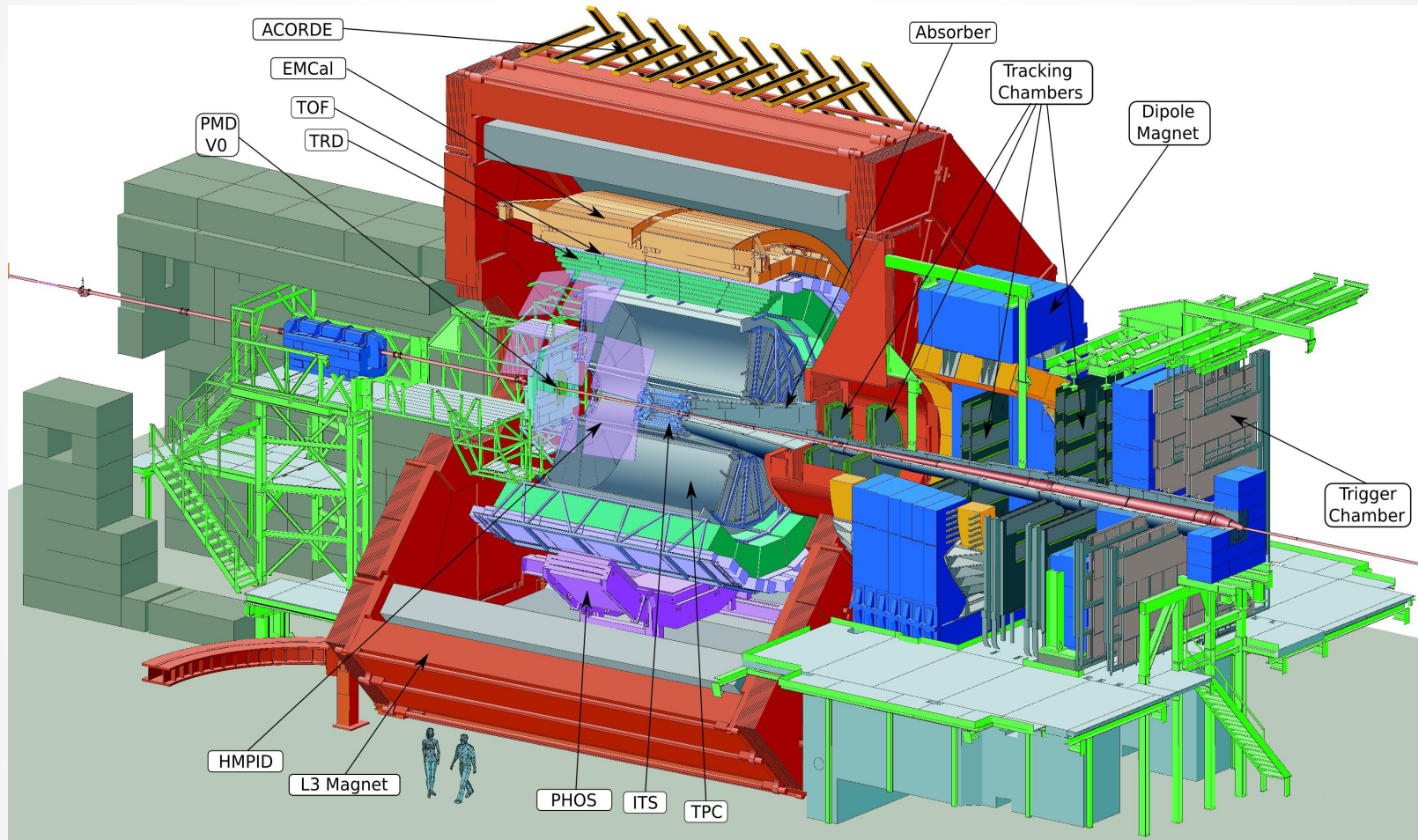


Fig. 1. ALICE detector system

Particles Identification

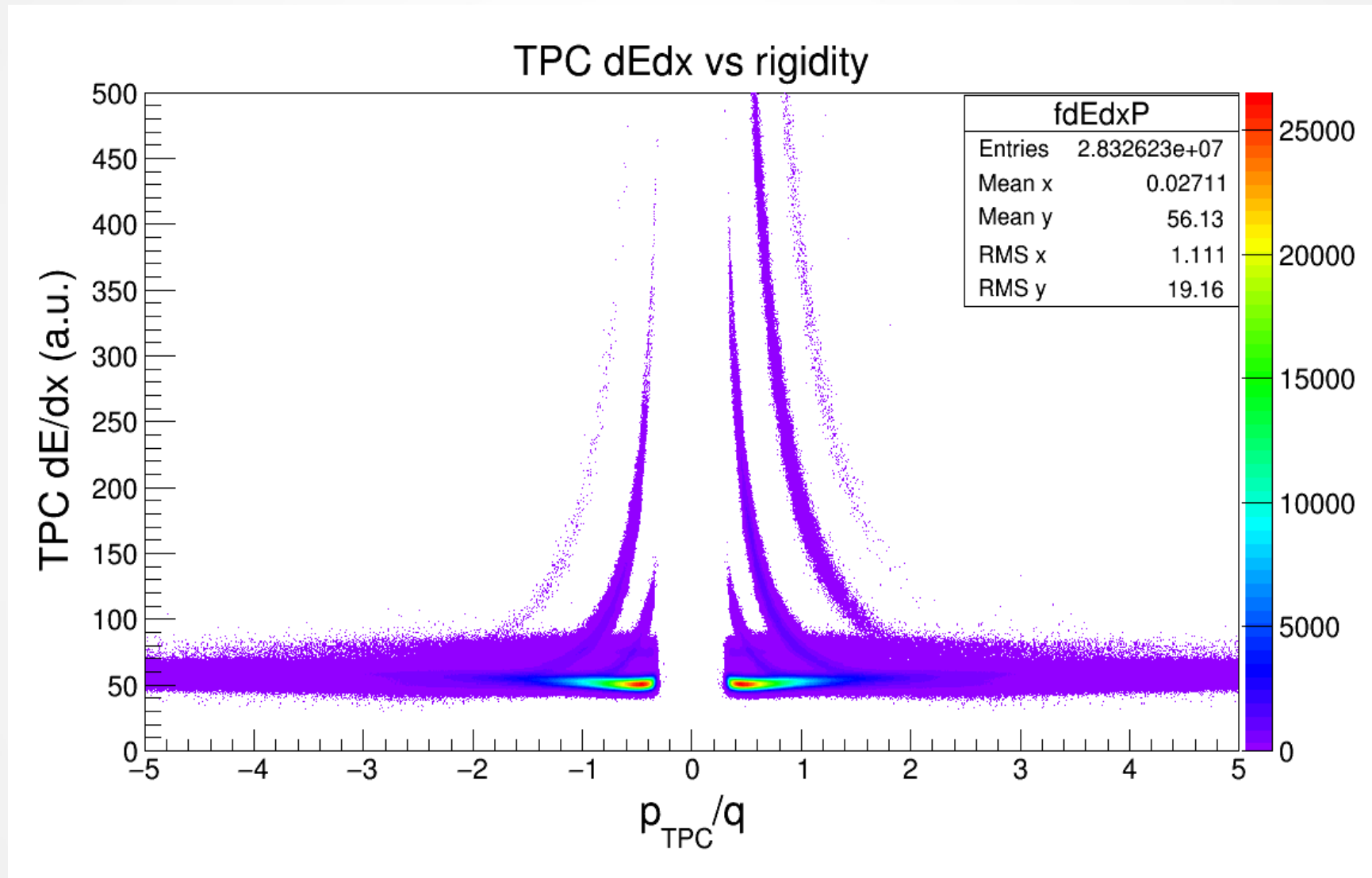


Fig. 2. Specific energy loss dE/dx in the TPC vs rigidity

Combined PID for TPC and TOF done with,

$$n\sigma_{combined} = \sqrt{(n\sigma_{TPC}^2 + n\sigma_{TOF}^2)}$$

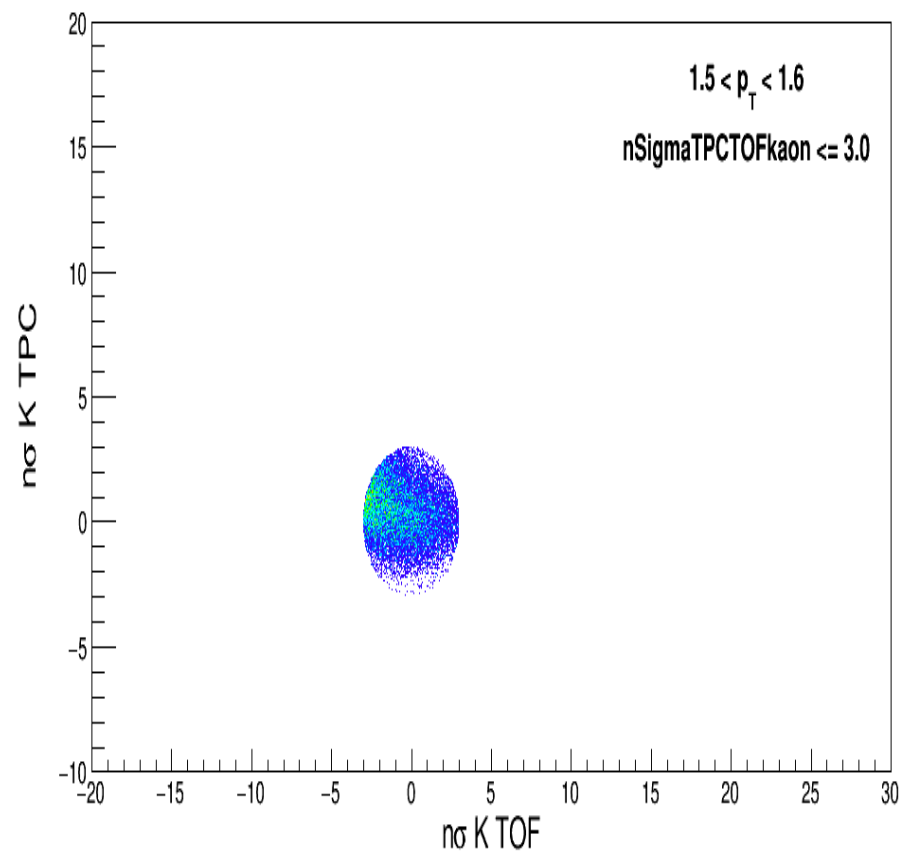
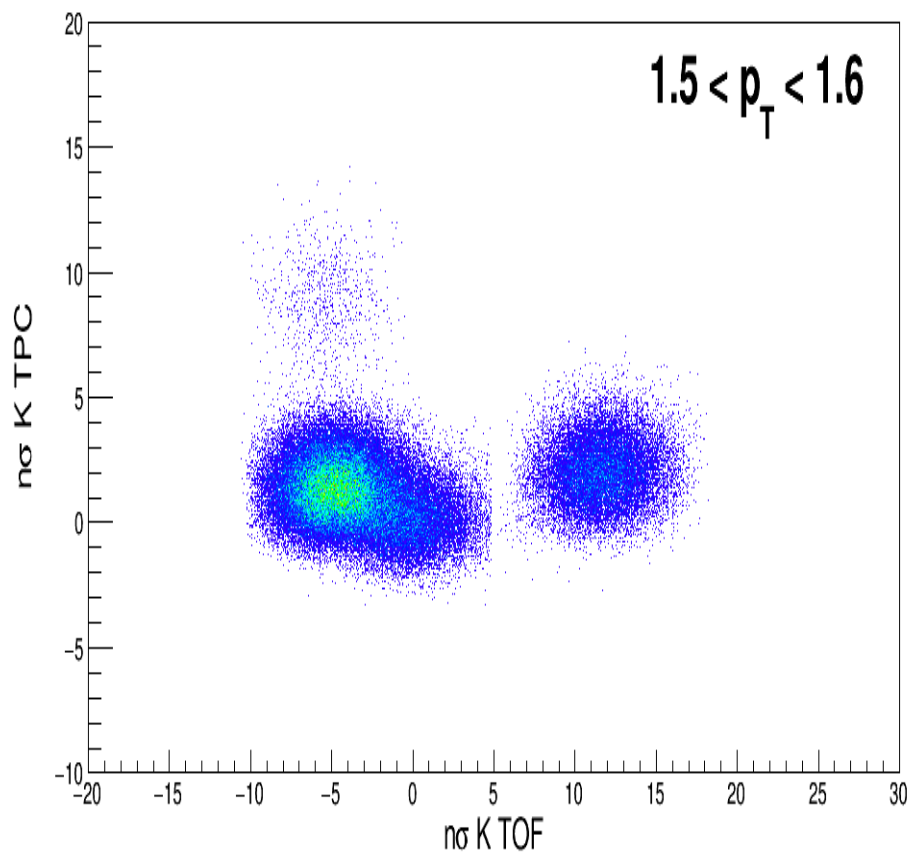


Fig. 3. PID with combined TPC and TOF detector

p_T dependent efficiency for all charged particles

Efficiency = Reconstructed / Generated

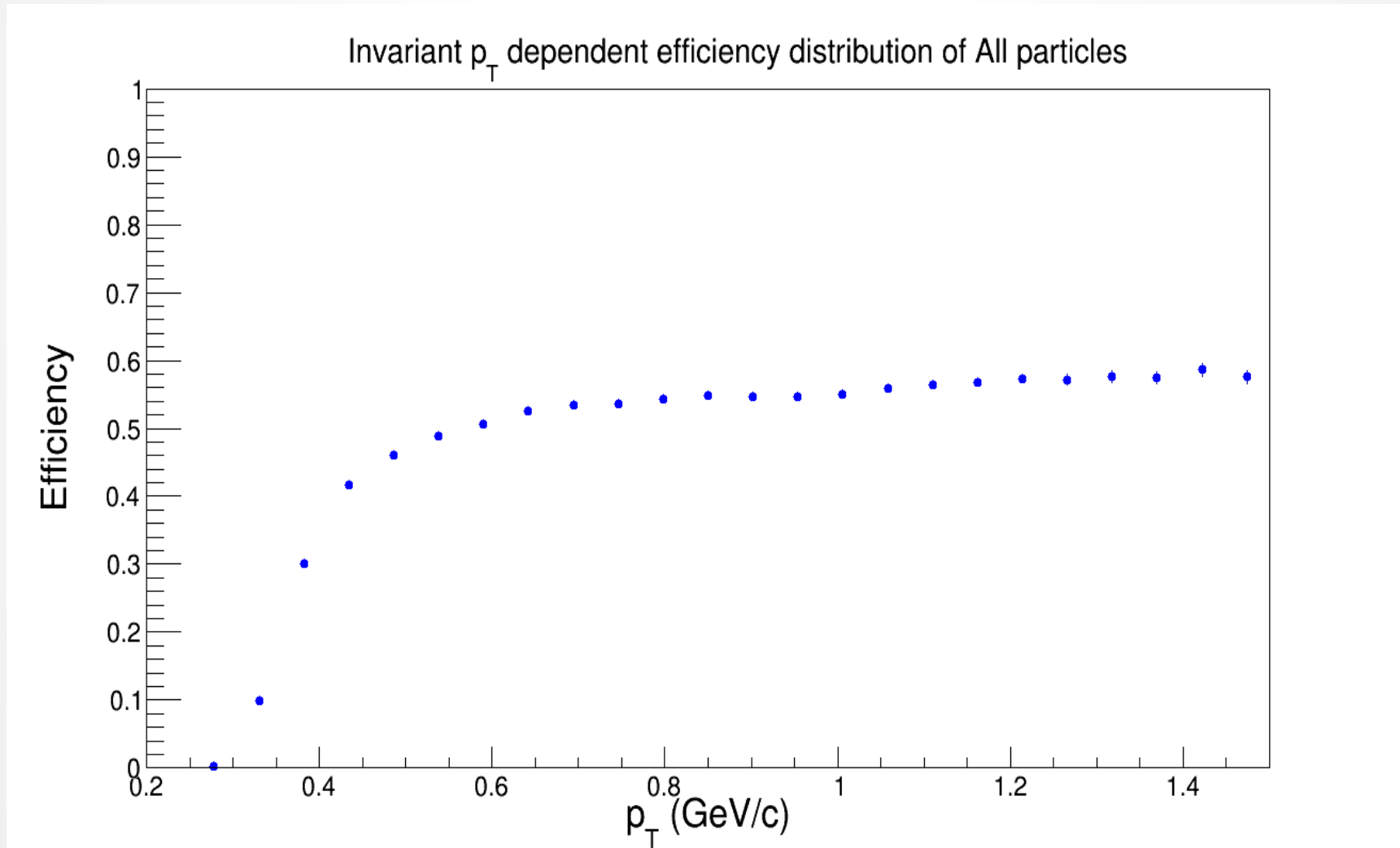


Fig. 4. p_T dependent efficiency distribution for all charged particles

p_T distribution of all charged particles

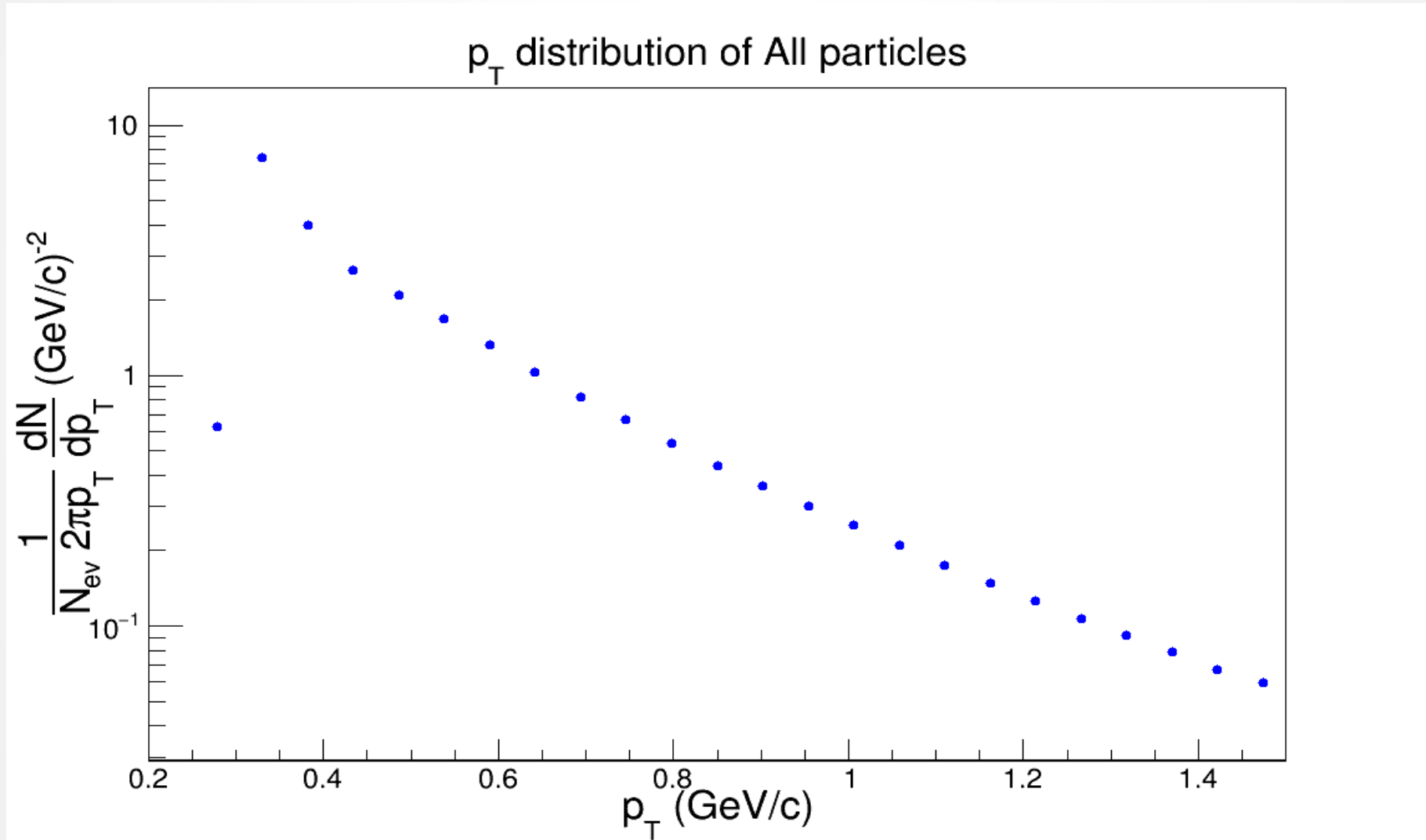


Fig. 5. Invariant p_T distribution for all charged particles

Comparison of p_T distribution with ALICE preliminary preliminary results for K^-

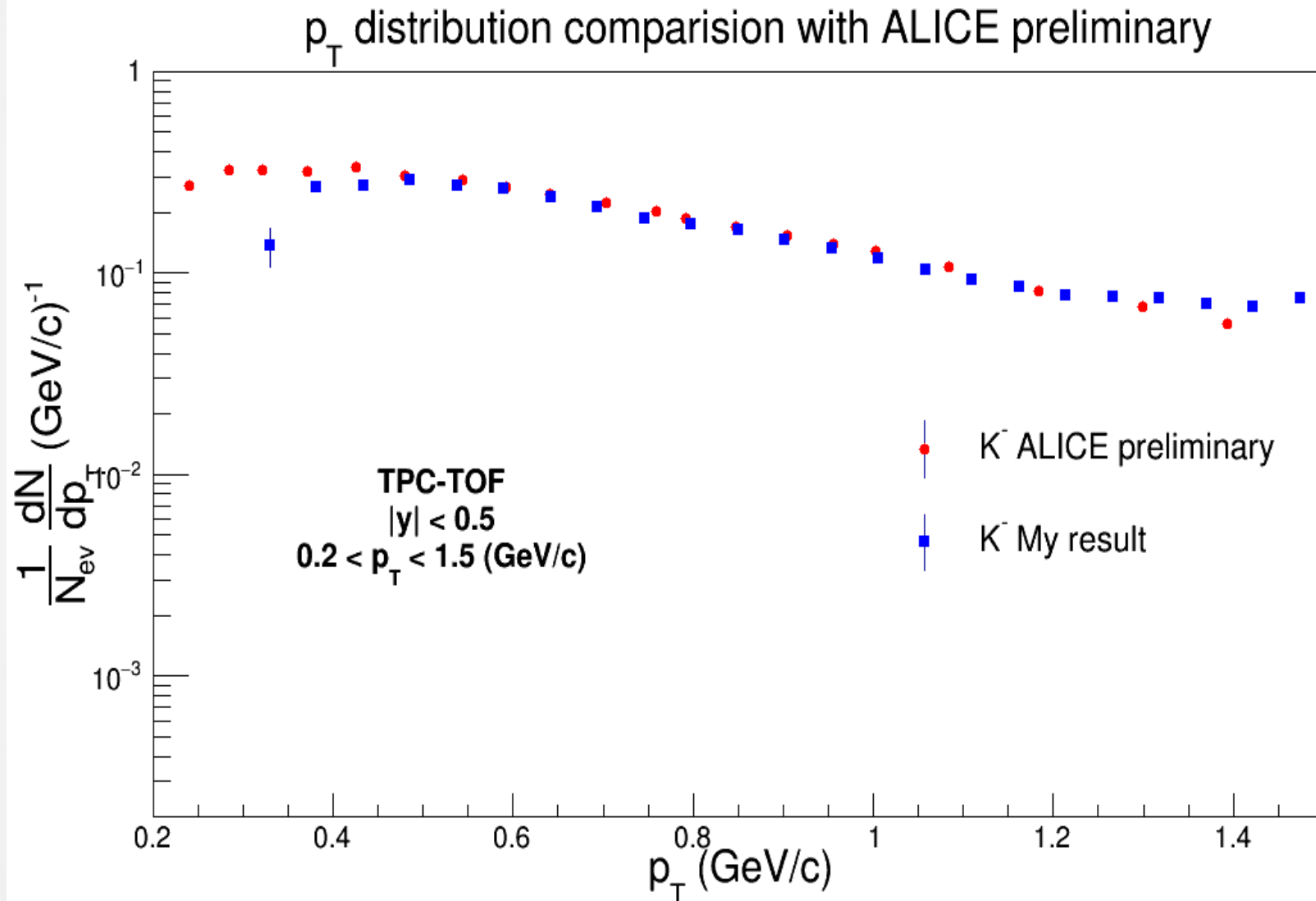


Fig. 6. Comparison of p_T distribution of K^-

Determination of color suppression factor $F(\xi)$

- ▶ To evaluate the initial value of ξ from data, a parameterization of pp collision at $\sqrt{s} = 200$ GeV is used to compute the p_T distribution,

$$\frac{dN}{dp_T^2} = \frac{a}{(p_0 + p_T)^n}$$

Where a is normalization constant, $p_0 = 1.98$ and $n = 12.88$ are parameter used to fit the data.

► The above parameter can be used for the high multiplicity pp event to take account of the interaction of the strings.

$$p_0 \rightarrow p_0 \left(\frac{\langle \frac{nS_1}{S_n} \rangle_{HM}}{\langle \frac{nS_1}{S_n} \rangle_{pp}} \right)^{\frac{1}{4}}$$

► Thus we get,

$$\frac{dN}{dp_T^2} = \frac{b}{\left(p_0 \sqrt{\frac{F(\xi)_{pp}}{F(\xi)_{HM}}} + p_T \right)^n}$$

Here we take $F(\xi)_{pp} = 1$

V0A Multiplicity

V0A multiplicity distribtuion

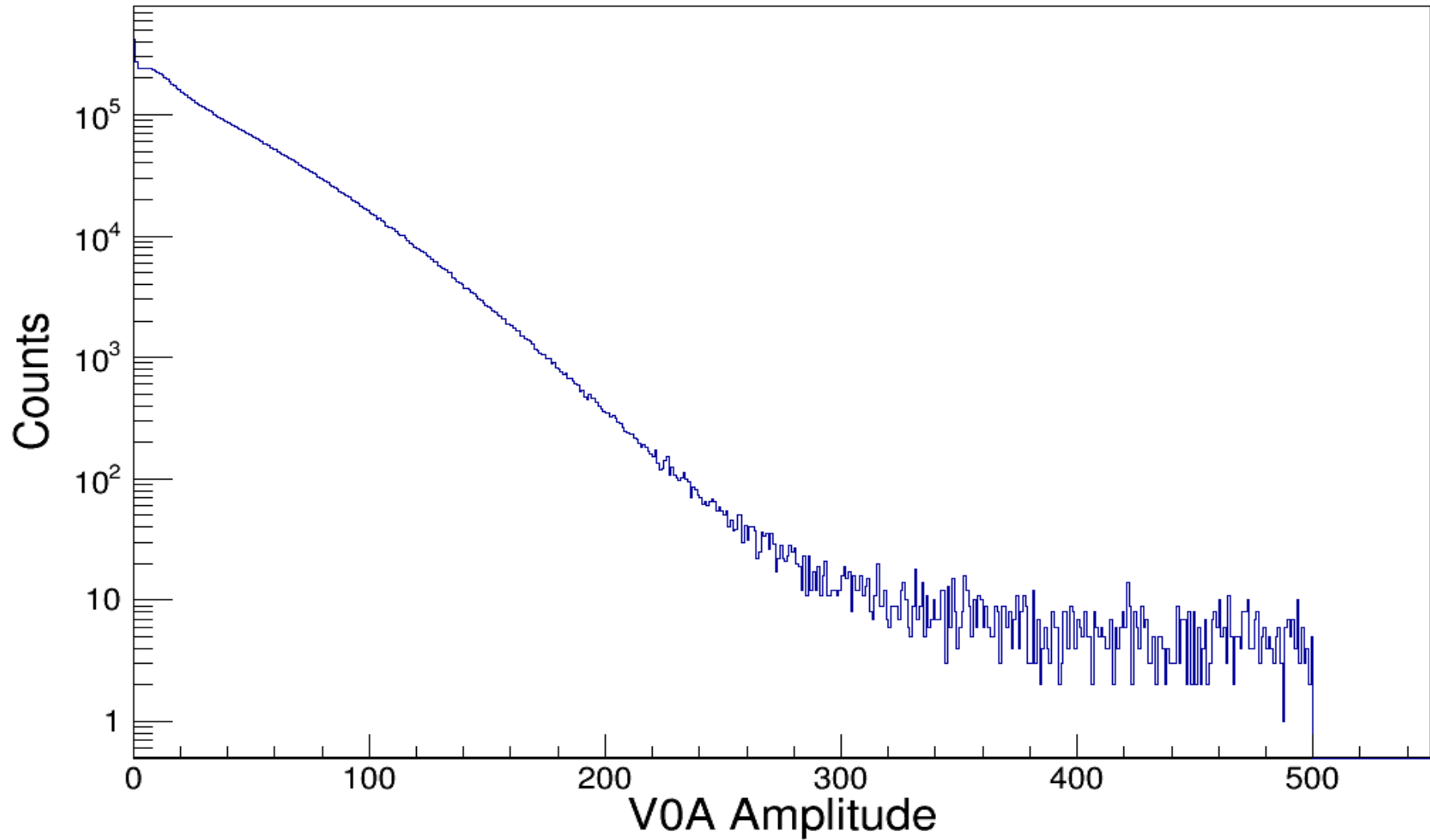


Fig. 7. V0A Multiplicity

p_T^2 dependent efficiency for all charged particles

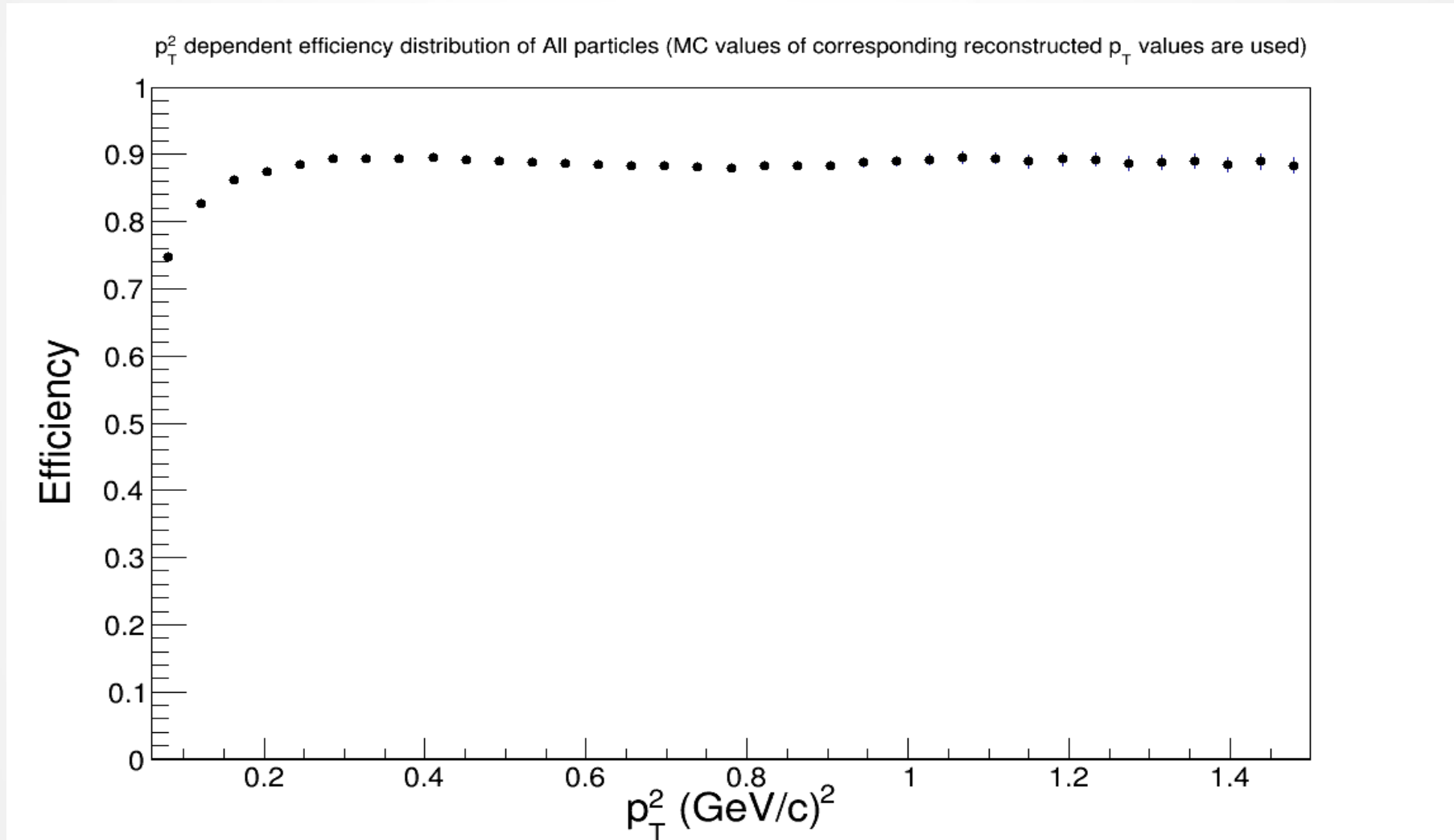


Fig. 8. p_T^2 dependent efficiency distribution for all charged particles

p_T^2 distribution of all charged particles for different multiplicity cut

The fitting equation is,
$$\frac{dN}{dp_T^2} = \frac{b}{\left(p_0 \sqrt{\frac{F(\xi)pp}{F(\xi)HM}} + p_T \right)^n}$$

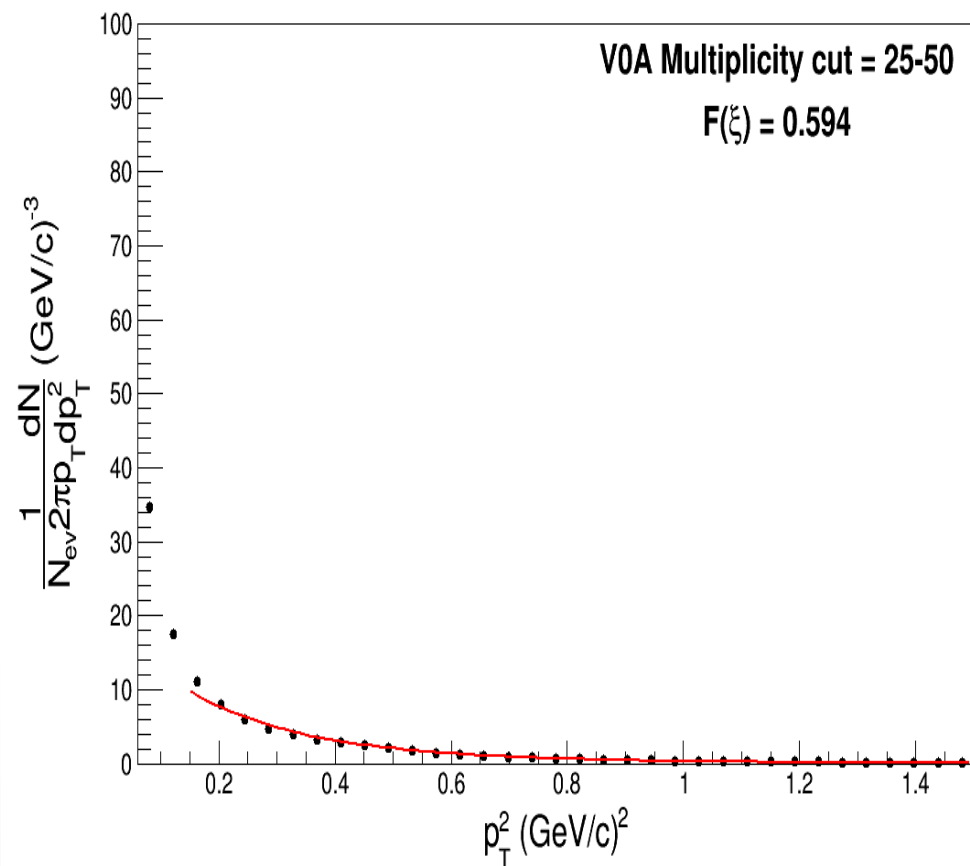
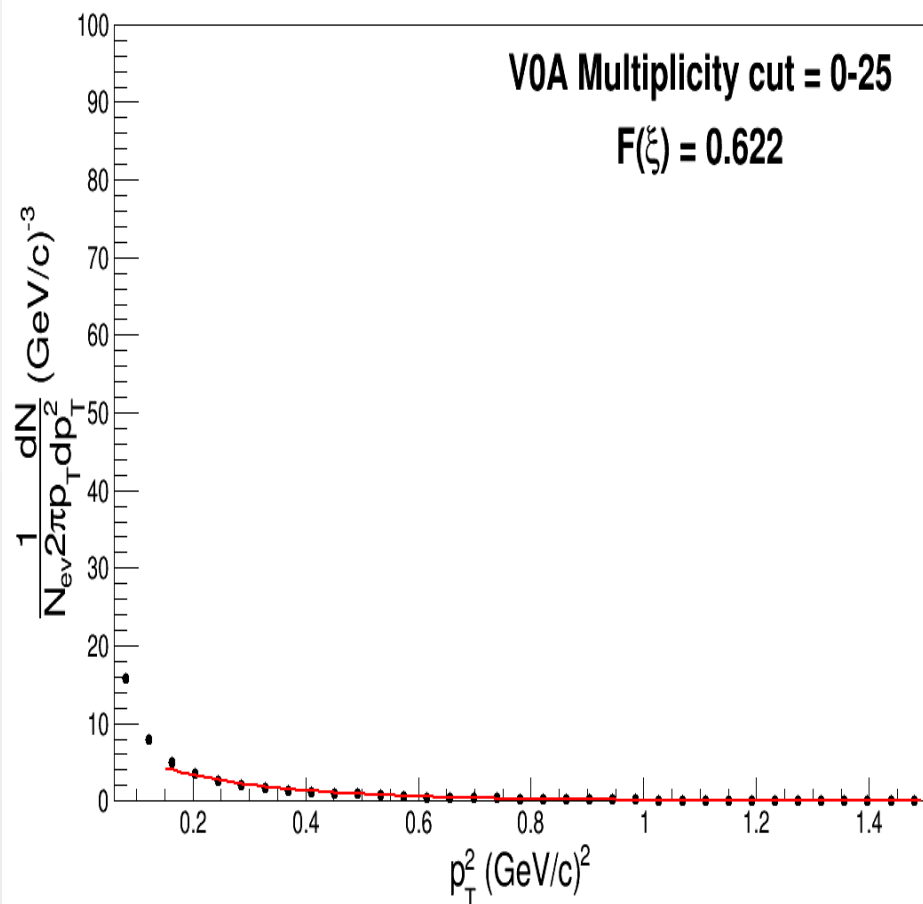


Fig. 9. Invariant p_T^2 distribution for all charged particles at different multiplicity

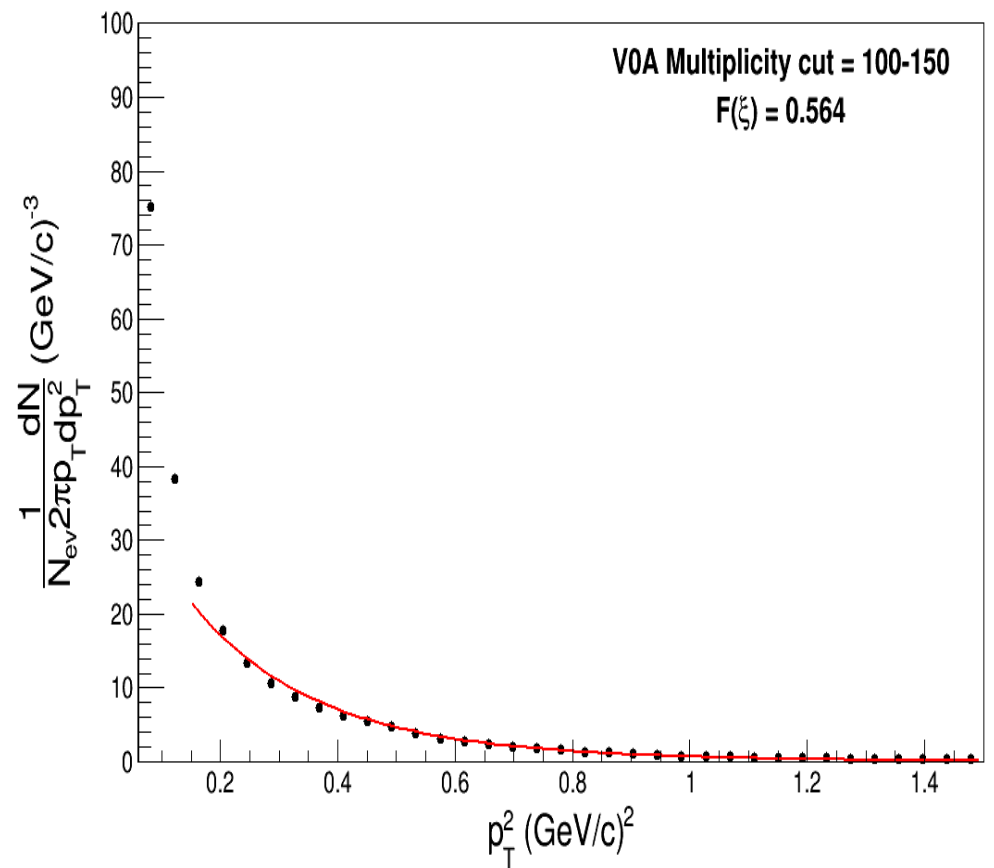
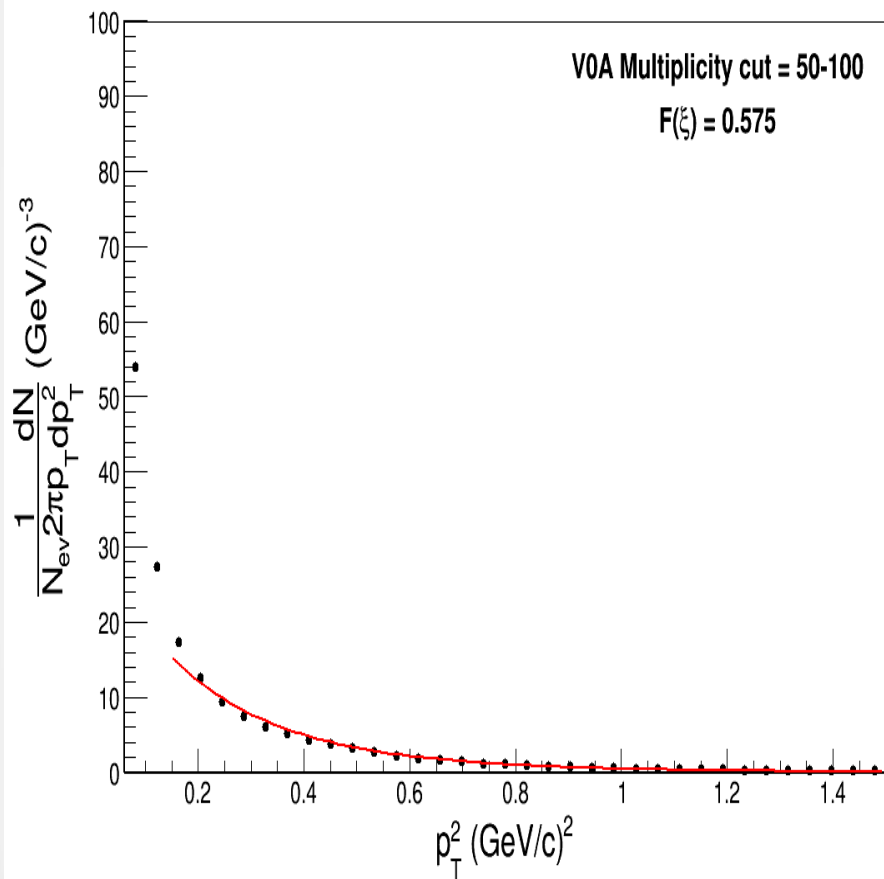


Fig. 10. Invariant p_T^2 distribution for all charged particles at different multiplicity

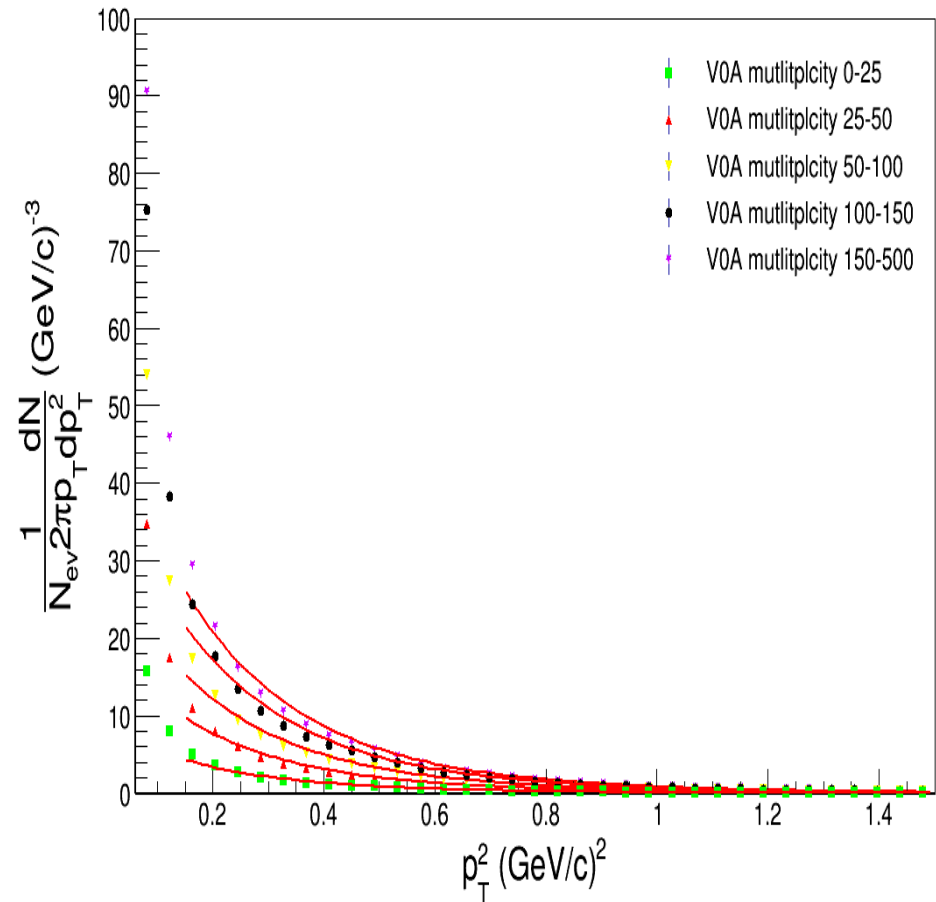
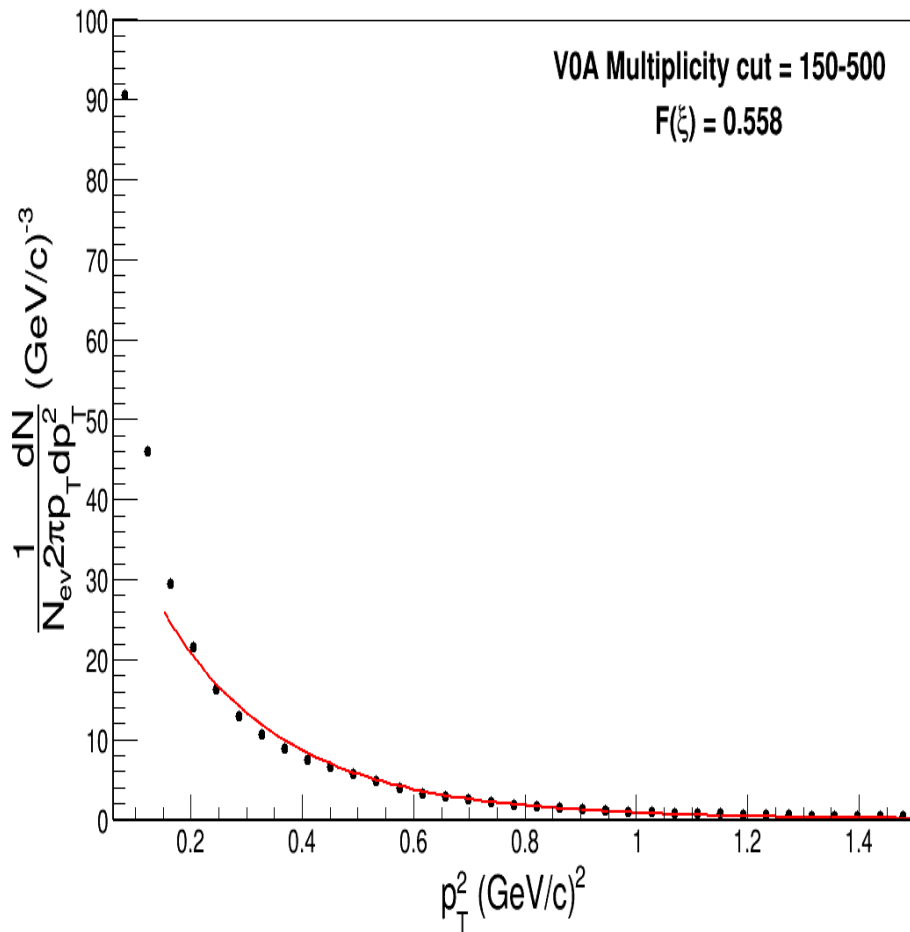


Fig. 11. Invariant p_T^2 distribution for all charged particles at multiplicity 150-500 and at all multiplicity

Table 1: V0A Multiplicity cut, $F(\xi)$, ξ and initial temperature

V0A Multiplicity cut	$F(\xi)$	ξ	Temperature (MeV)
0-25	0.622	2.34	185.8 ± 3.35
25-50	0.594	2.62	190.1 ± 3.44
50-100	0.575	2.85	193.2 ± 3.49
100-150	0.564	2.99	195.1 ± 3.56
150-500	0.558	3.06	196.1 ± 3.68

Table 2: Initial temperature of other system (published results)

System	Temperature (MeV)
Au+Au (0-10%) @ 200 GeV	193.6 ± 3.0
Pb+Pb (0-5%) @ 2.76 TeV (estimated)	262.2 ± 13.0

Ref: J. D. de Deus, A. S. Hirsch, C. Pajares, R. P. Scharenberg, B. K. Srivastava, Eur.Phys. J. C 72 (2012) 2123.

Summary

- Successfully extract the color suppression factor and percolation density.
- Temperature is also determined successfully.
- The temperature obtained is much above the freeze out temperature.
- The temperature obtained from the analysis are in agreement with the published results.

Outlook

- Analyze the all available energy. (2.76 TeV and 13 TeV)
- To check the variation of energy density with temperature.
- Variation of shear viscosity to entropy density ratio with temperature.

References

- R. P. Scharenberg *et al.*, Eur. Phys. J. C 71 (2011) 1510.
- J. D. de Deus *et al.*, Eur. Phys. J. C 72 (2012) 2123.
- B. K. Srivastava, Nucl. Phys. A 926 (2014) 142.
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- B. K. Srivastava *et al.*, *Int.J. Mod. Phys. E Vol. 24, No. 12 (2015) 1550101*
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